

# The **L**orentz **I**ntegral **T**ransform and Resonances

- Introduction
- LIT method: controlled resolution
- $0^+$  resonance of  $^4\text{He}$
- Case study for two-nucleon case

# main point of the LIT :

Schrödinger-like equation with a source

$$(H - E_0 - \omega_0 + i\Gamma) \tilde{\Psi} = S$$

The  $\tilde{\Psi}$  solution is unique and has **bound state like** asymptotic behavior



one can apply **bound state methods**

# Reformulation of the LIT

$$\text{LIT}(\sigma_R, \sigma_I) = -\frac{1}{\sigma_I} \text{Im} \left\{ \langle \Psi_0 | \Theta^\dagger (\sigma_R + E_0 - H + i \sigma_I)^{-1} \Theta | \Psi_0 \rangle \right\}$$

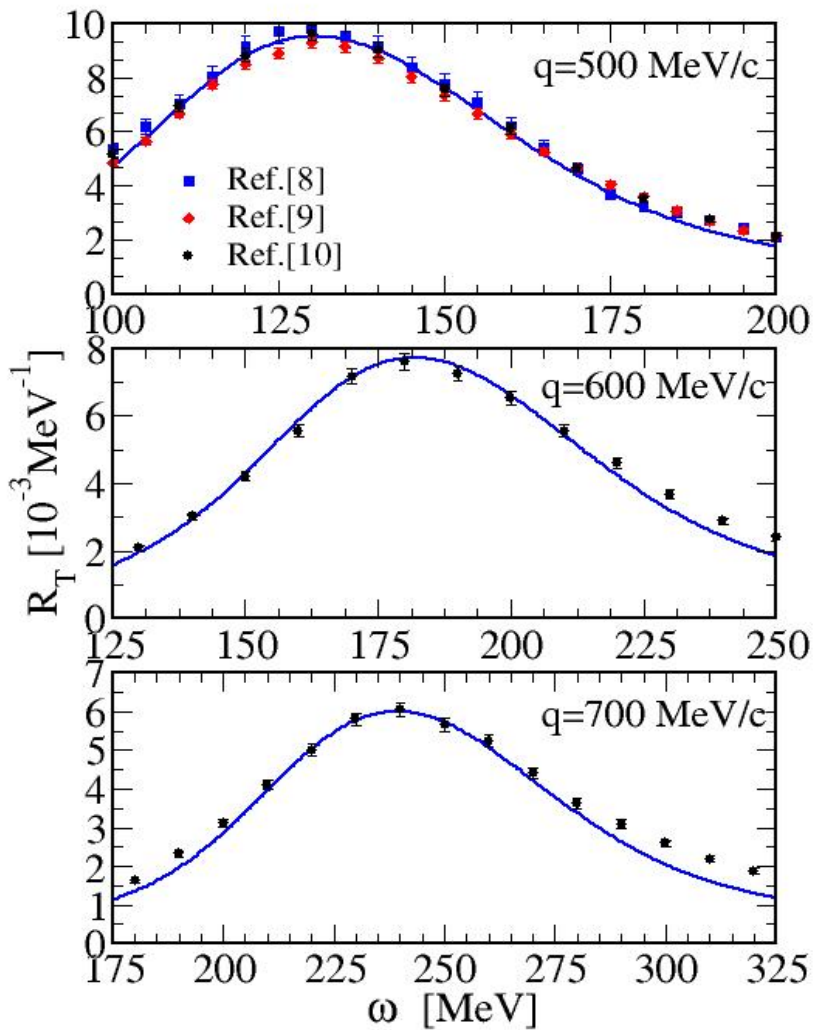
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$$R(E = \sigma_R) = -\frac{1}{\pi} \text{Im} \left\{ \lim_{\sigma_I \rightarrow 0} \langle \Psi_0 | \Theta^\dagger (\sigma_R + E_0 - H + i \sigma_I)^{-1} \Theta | \Psi_0 \rangle \right\}$$

LIT method allows calculation  
up into the far continuum!

## $R_T(q, \omega)$ of ${}^3\text{He}(e, e')$



NN potential AV18  
Three-nucleon force UIX

L. Yuan et al., PLB 706, 90 (2011)

Experimental data:  
Bates, Saclay,  
world data (J. Carlson et al.)

# LIT - Inversion

## Standard LIT inversion method

Take the following ansatz for the response function  $R(\omega)$  (or  $F_{fi}(E,E')$ )

$$R(\omega') = \sum_{m=1}^{M_{\max}} c_m \chi_m(\omega', \alpha_i)$$

with  $\omega' = \omega - \omega_{th}$ , given set of functions  $\chi_m$ , and unknown coefficients  $c_m$

Define: 
$$\tilde{\chi}_m(\sigma_R, \sigma_I, \alpha_i) = \int_0^{\infty} d\omega' \frac{\chi_m(\omega', \alpha_i)}{(\omega' - \sigma_R)^2 + \sigma_I^2}$$

Calculate LIT  $L(\sigma_R, \sigma_I) = \langle \tilde{\Psi} | \tilde{\Psi} \rangle$  for many  $\sigma_R$  and fixed  $\sigma_I$

and expand in set  $\tilde{\chi}_m$ : 
$$L(\sigma_R, \sigma_I) = \sum_{m=1}^{M_{\max}} c_m \tilde{\chi}_m(\omega', \alpha_i)$$

Determine  $c_m$  via best fit

Increase  $M_{\max}$  up to the point that stable result is obtained for  $R(\omega)$ . Even further increase of  $M_{\max}$  might lead to new structures in  $R(\omega)$



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As basis set  $\chi_m$  we normally use

$$\chi_m(\omega', \alpha_i) = (\omega')^\alpha \exp(-\alpha_2 \omega'/m) \quad \text{with } m = 1, 2, \dots, M_{\max}$$

# **LIT method: controlled resolution**

# LIT - Example

deuteron photodisintegration in unretarded dipole approximation

unretarded dipole approximation  $\Rightarrow \Theta = \sum_{i=1}^A z_i \frac{1+\tau_{i,z}}{2}$ ,

$z_i, \tau_{i,z}$ : 3<sup>rd</sup> components of position and isospin coordinates of i-th nucleon

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$$\stackrel{\Theta}{\Rightarrow} \sigma_{\gamma}(\omega) = 4\pi^2 \alpha \omega R(\omega) \quad \text{with} \quad R(\omega) = \sum_f |\langle f | \Theta | 0 \rangle|^2 \delta(\omega - E_f - E_0)$$

with  $|0\rangle$  and  $E_0$  bound-state wave function and energy

$|f\rangle$  and  $E_f$  final-state wave function and energy

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$|\mathbf{f}\rangle$  and  $E_{\mathbf{f}}$  final-state wave function and energy

In unretarded dipole approximation  $|\mathbf{f}\rangle$  contains only  ${}^3P_0, {}^3P_1, {}^3P_2 - {}^3F_2$  NN states

Solution of LIT equation via expansion on HH basis

⇒ Discretization of the continuum

Example: deuteron photodisintegration in unretarded dipole approximation. Expansion of radial part **Laguerre polynomials up to order N** times **an exponential fall-off  $\exp(-\rho/b)$**  with  $\rho \equiv r$

Solution of LIT equation via expansion on HH basis

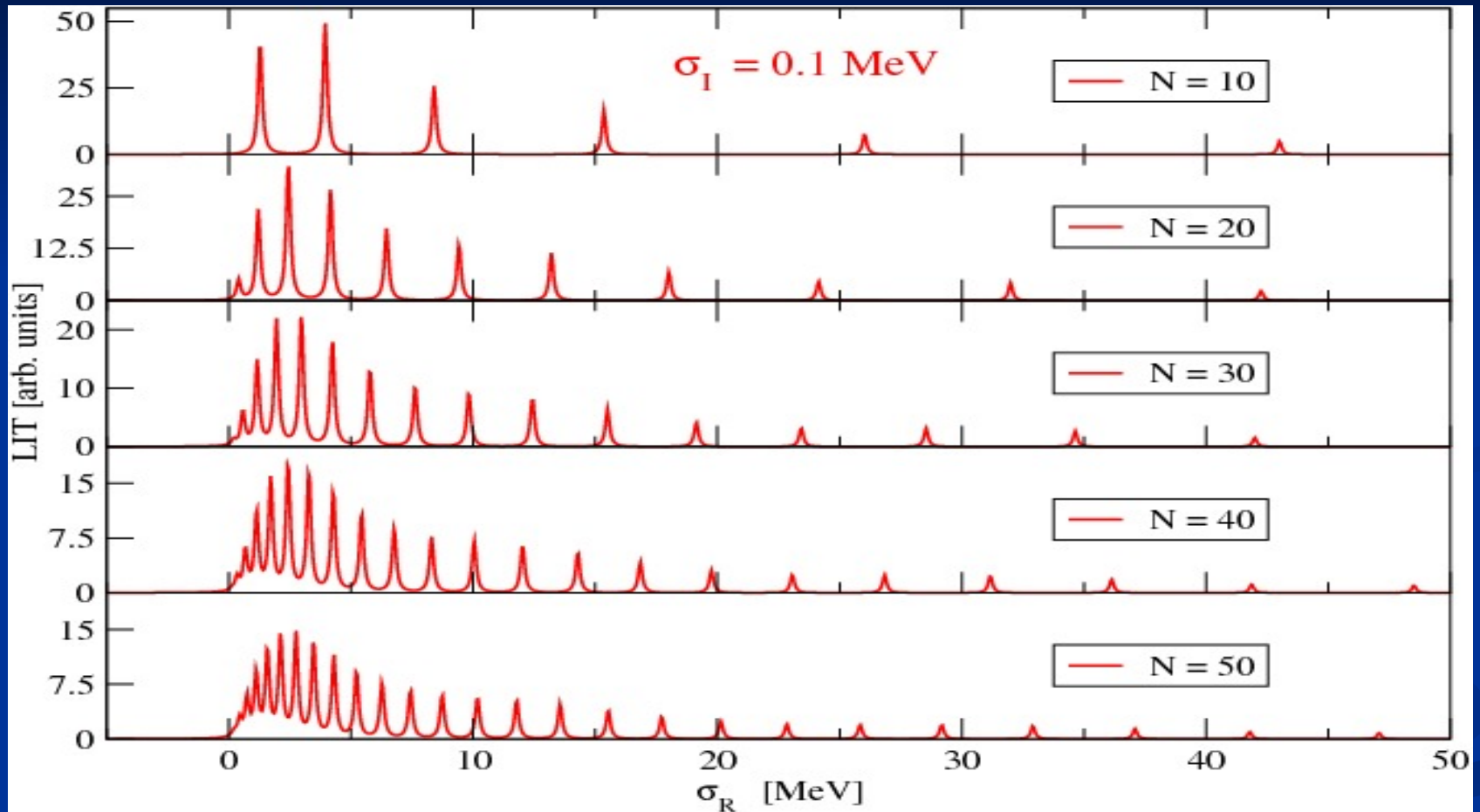
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LIT results with various resolutions  $\sigma_l$

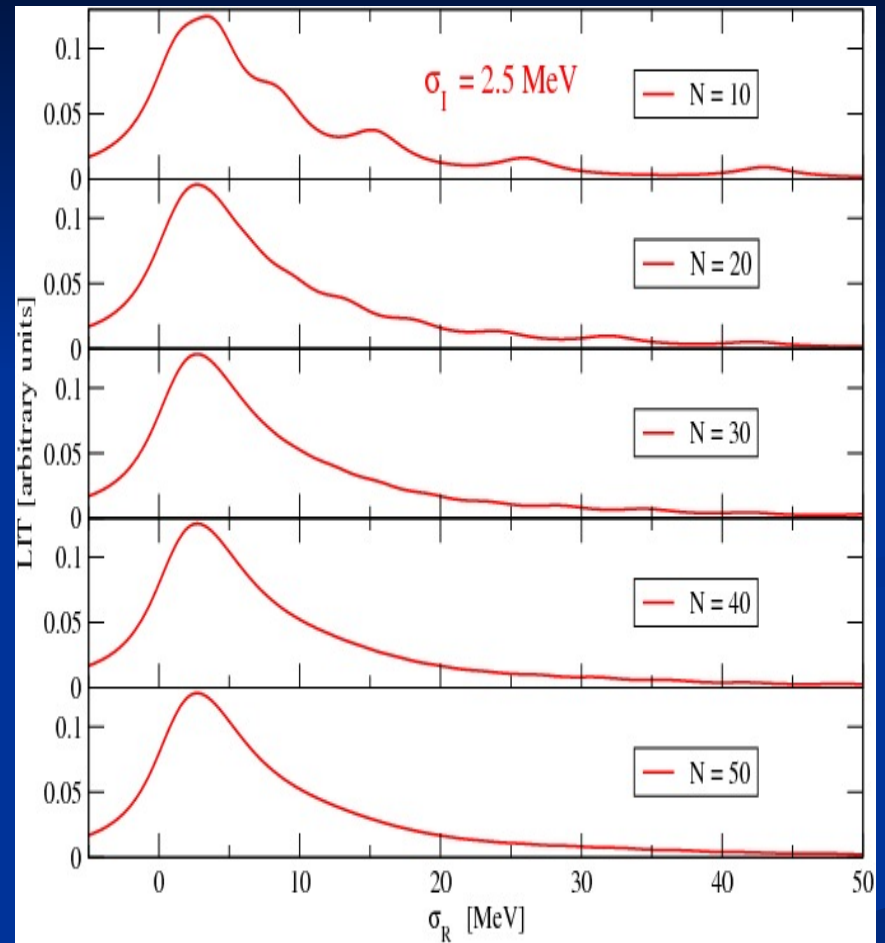
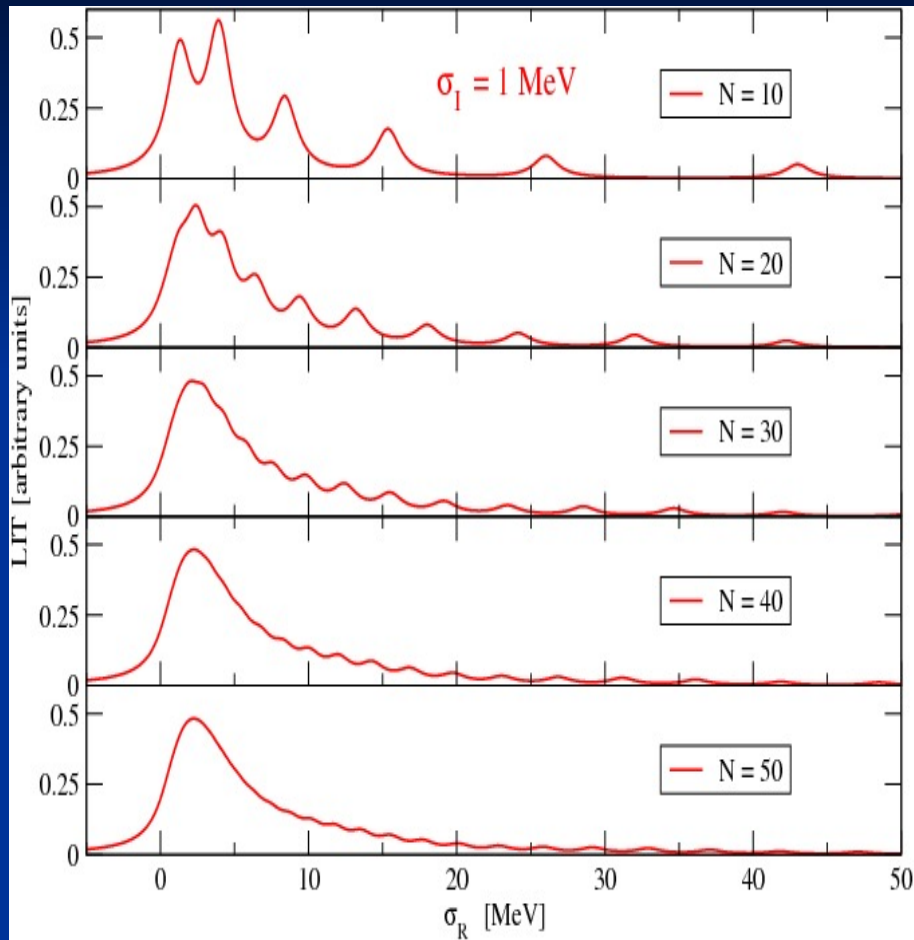
(only transitions to  ${}^3P_1$ )

This leads to the following LITs with Laguerre polynomials up to order  $N$  with exponential fall-off ( $b=0.5$  fm):

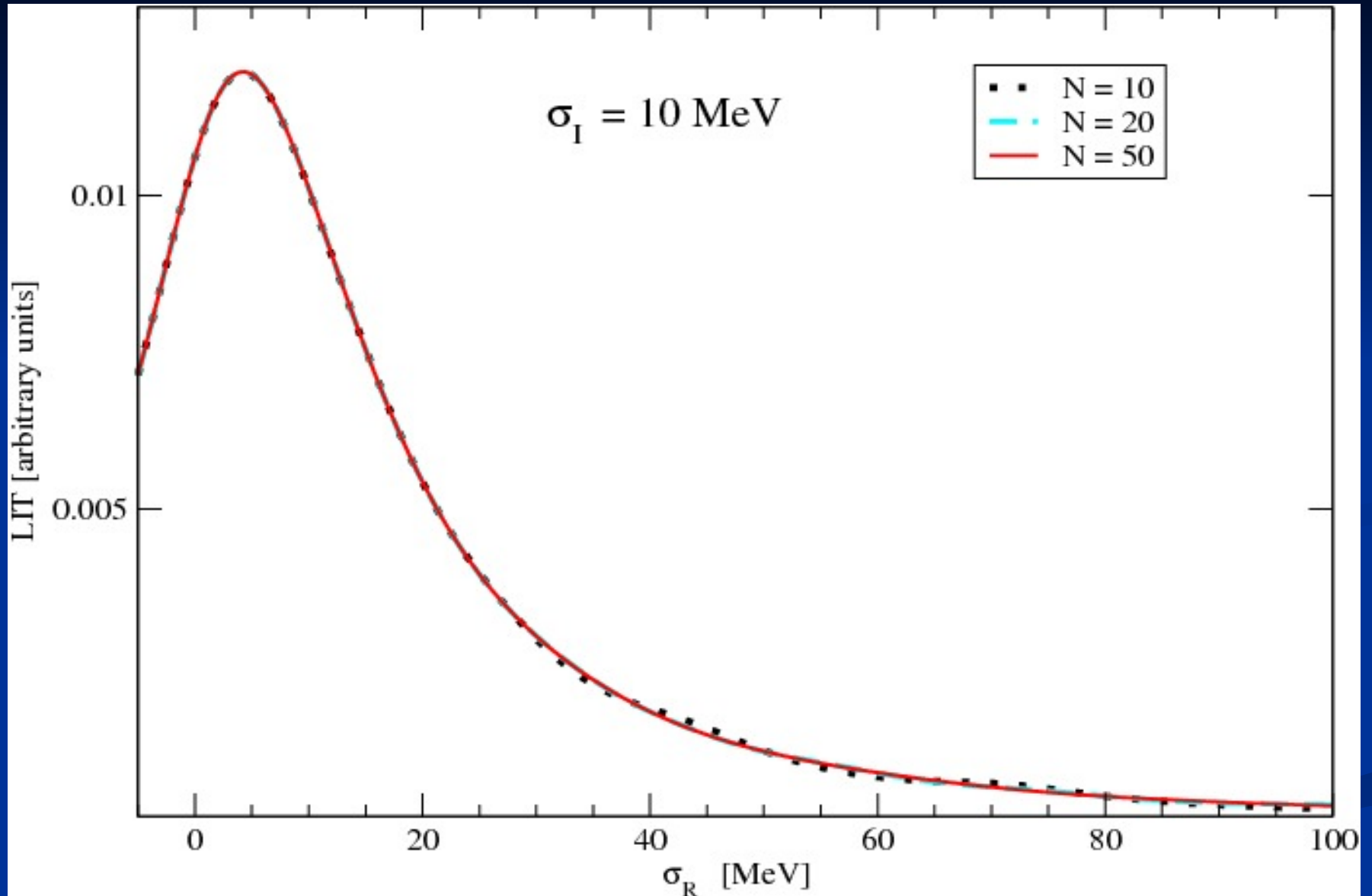




## Laguerre polynomials up to order N (exponential fall-off)



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LIT method is an approach with a controlled resolution!

Strength for a given discrete state of energy  $E$  is **not** the actual strength for this energy, but can only be interpreted correctly within an integral transform approach.

The **correct distribution** of strength is obtained via the **inversion** of the integral transform.

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**Correct procedure:** Ensure convergence of expansion and search for the smallest possible resolution

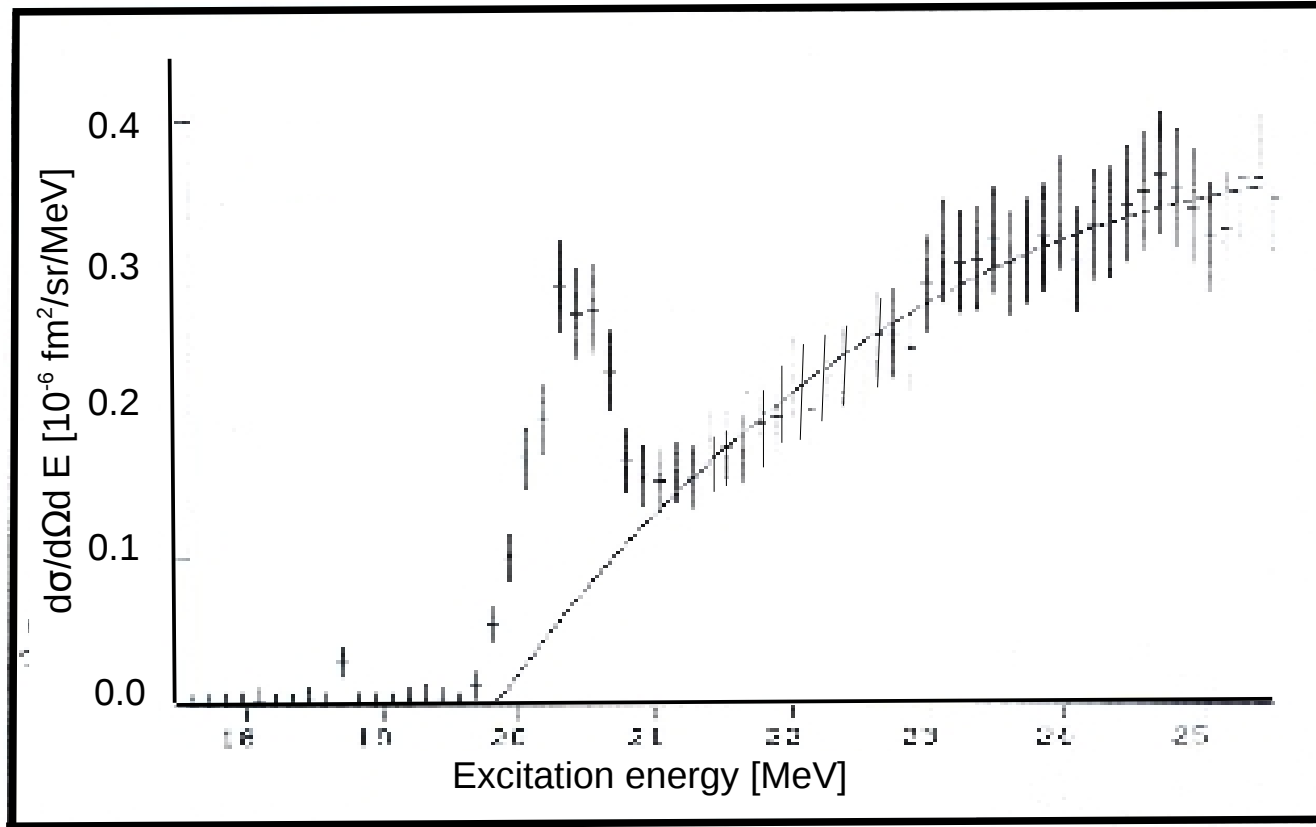
# LIT method and resonances

# $O^+$ resonance in longitudinal response function $R_L$ in ${}^4\text{He}(e,e')$

S. Bacca, N. Barnea, WL, G. Orlandini, PRL 110, 042503 (2013)

# $0^+$ Resonance in the $^4\text{He}$ compound system

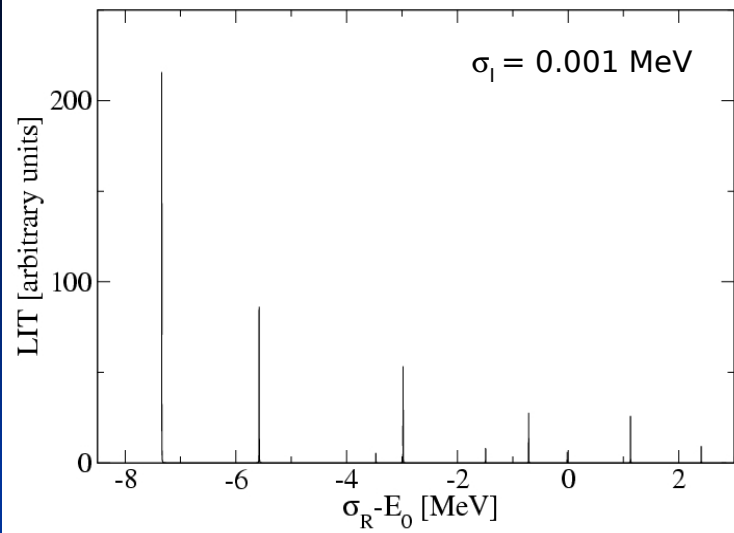
G. Köbschall et al./ Quasi bound state in  $^4\text{He}$  - Nucl. Phys. A405, 648 (1983)

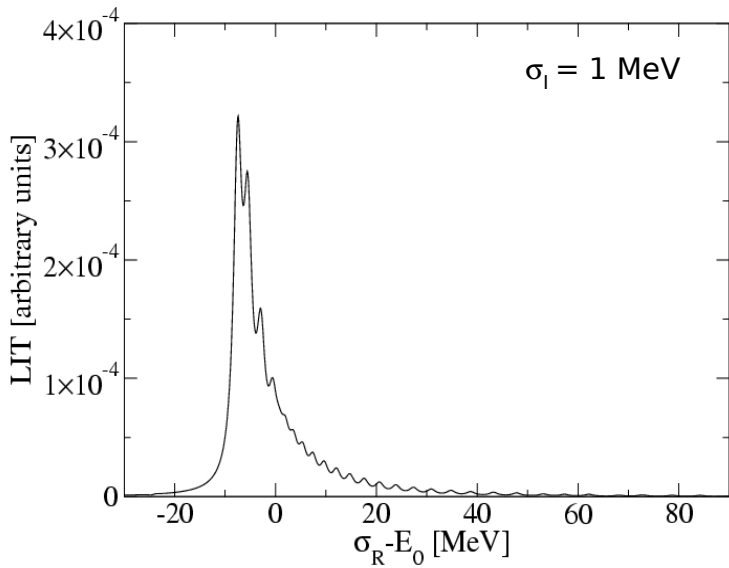
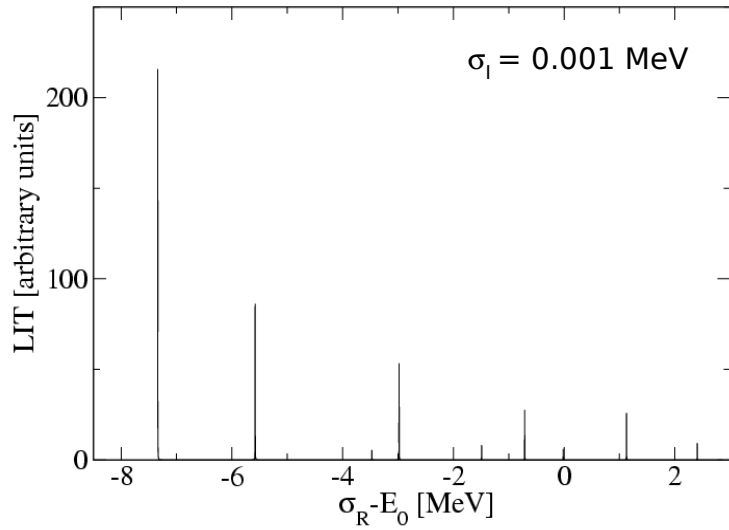


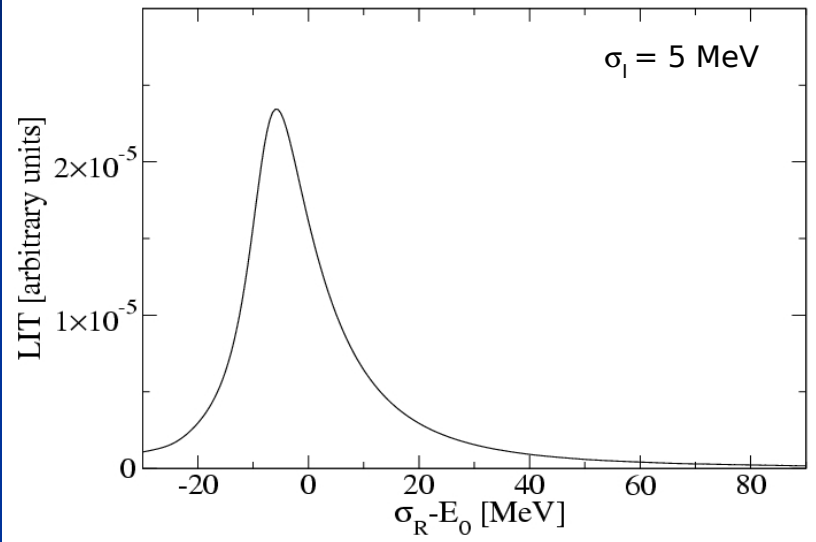
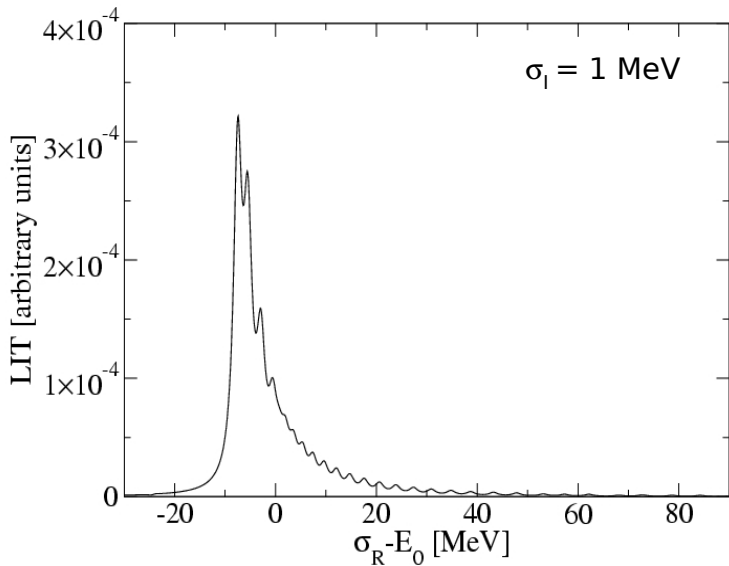
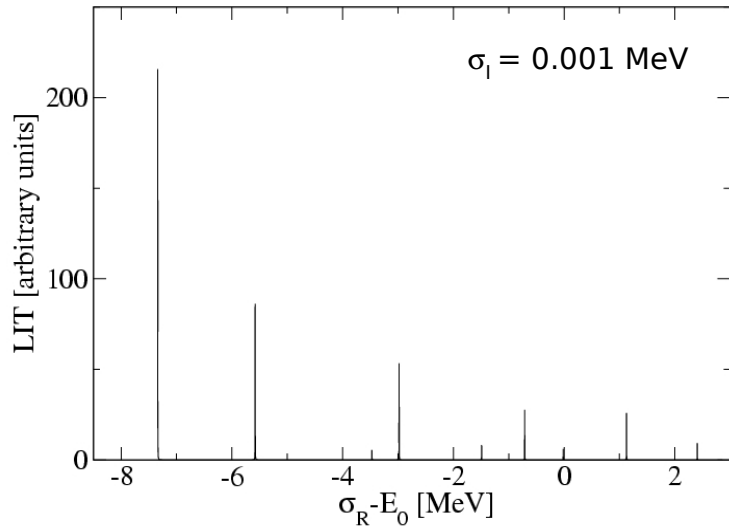
Resonance at  $E_R = -8.2$  MeV, i.e. above the  $^3\text{H}$ -p threshold. **Strong evidence** in electron scattering off  $^4\text{He}$ ,  $\Gamma = 270 \pm 70$  keV

# Results of our LIT calculation









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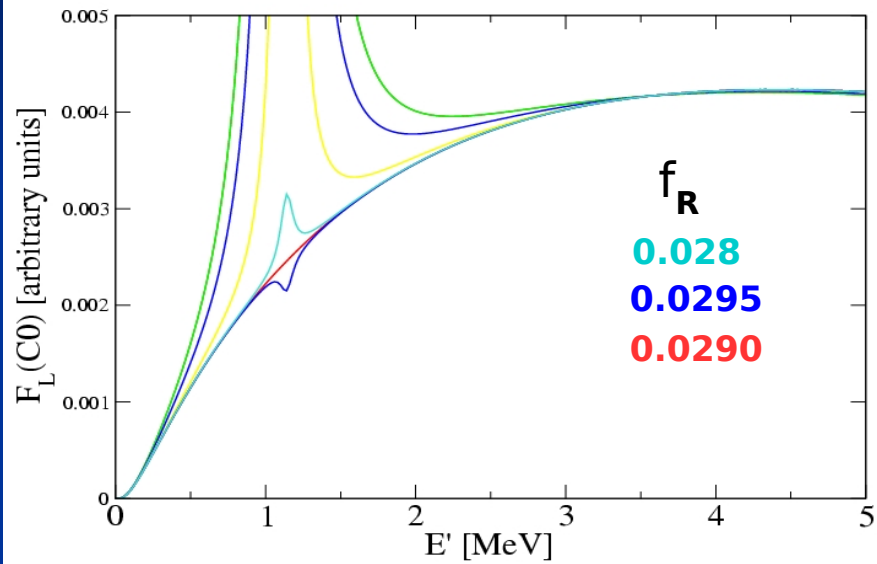
However, the strength of the resonance can be determined!

Of course not by taking the strength to the discretized state, but by rearranging the inversion in a suitable way:

Reduce strength of the first state up to the point that the inversion does not show any resonant structure at the resonance energy  $E_R$ :

$$\text{LIT}(\sigma_R, \sigma_I) \rightarrow \text{LIT}(\sigma_R, \sigma_I) - f_R / [(E_R - \sigma_R)^2 + \sigma_I^2] \equiv \text{LIT}(\sigma_R, \sigma_I, f_R)$$

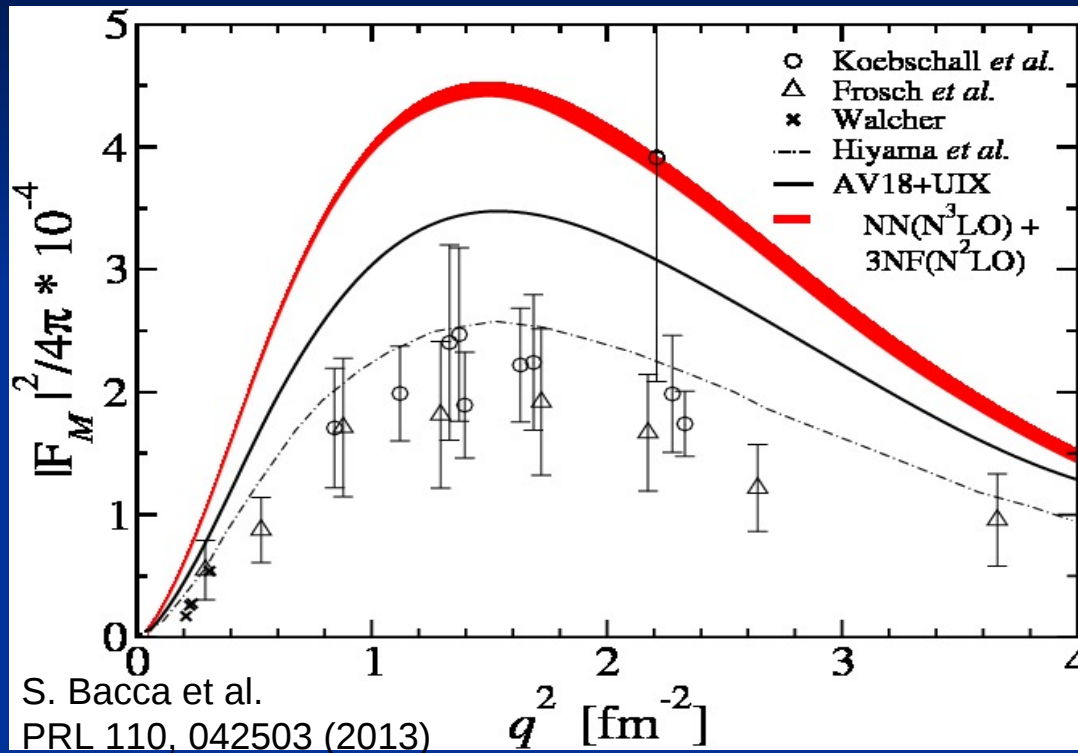
with resonance strength  $f_R$



Inversion results with  
different  $f_R$  values  
AV18+UIX,  $q=300$  MeV/c



# Comparison to experimental results



LIT/EIHH Calculation for AV18+UIX and Idaho-N3LO+N2LO

Dotted: AV8' + central 3NF (Hiyama et al.)

# **Study of Resonance for the deuteron case**

## NN potential with fictitious resonance in ${}^3P_1$ partial wave

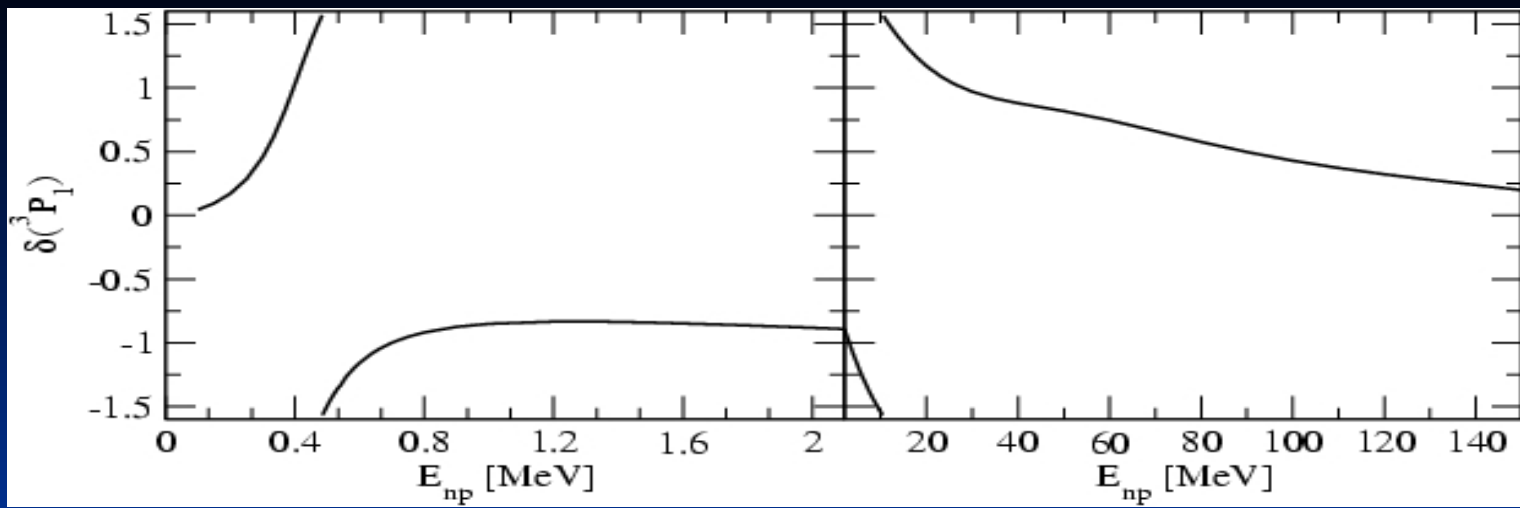
$$V({}^3P_1) \longrightarrow V({}^3P_1) + V_{\text{add}}$$

$$\text{With } V_{\text{add}} = -\frac{57.6 \text{ MeV}}{r} (1 - \exp(-2r^2))(1 + \exp(\frac{r-5}{0.2}))^{-1}$$

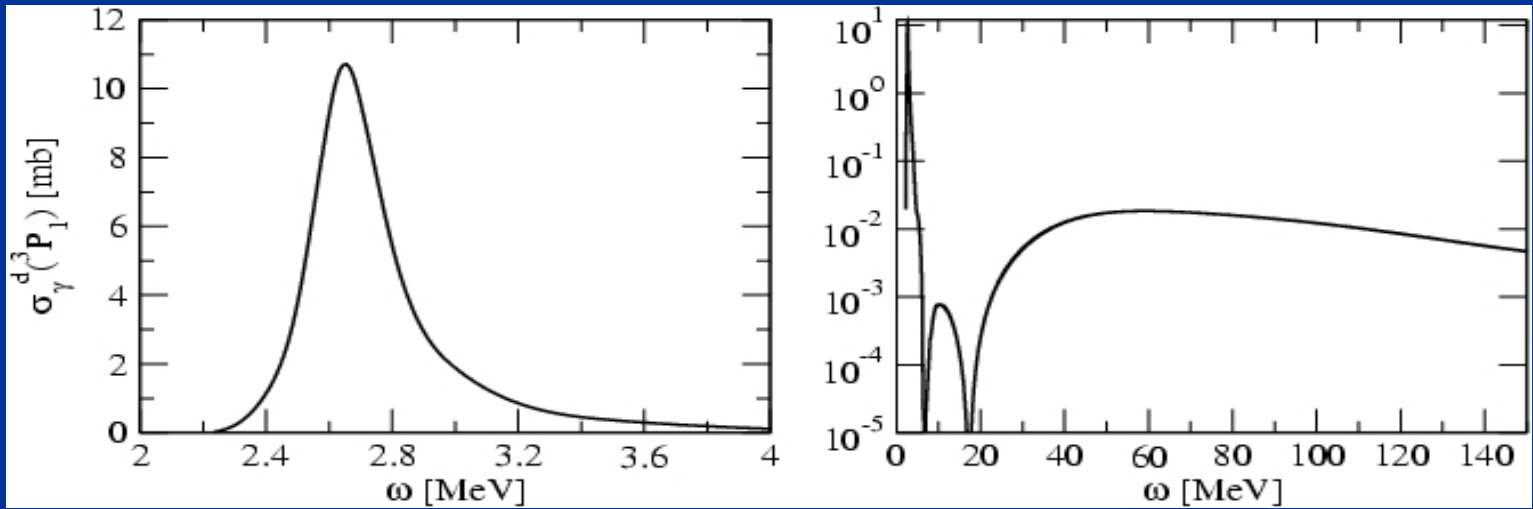
and relative coordinate  $r$  in units of fm

### Why such a potential?

To understand this better let us have a look on corresponding phaseshift  ${}^3P_1$  and deuteron photoabsorption cross section in  ${}^3P_1$  partial wave



Phase shifts shows two resonances one at  $E_{np} = 0.48, 10.5$  MeV

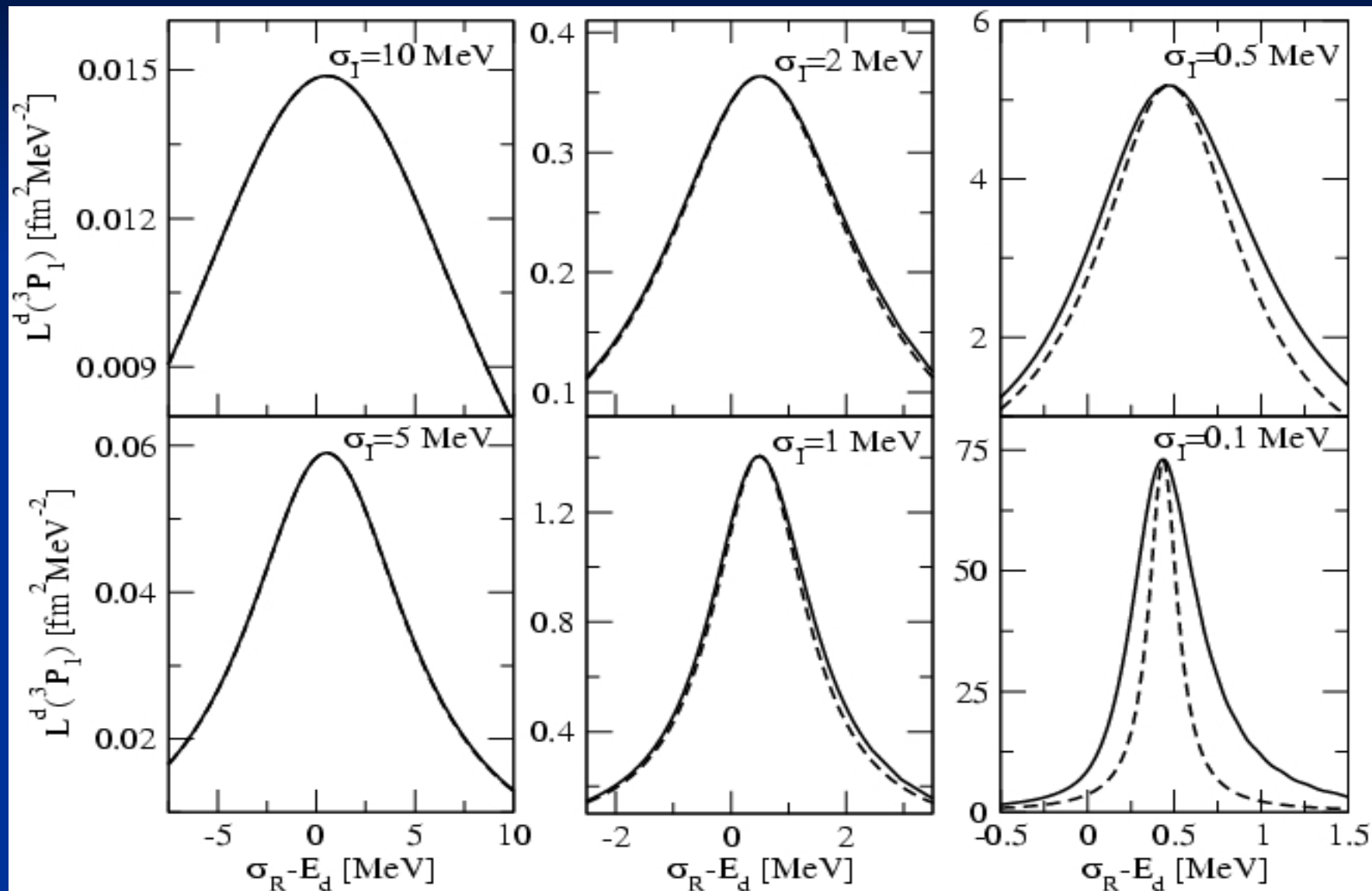


$\sigma_{\gamma}^d(^3P_1)$  shows two corresponding resonances: low-energy resonance very pronounced with small width  $\Gamma=270$  KeV, the other one is much weaker and has a larger width

# Results with modified ${}^3P_1$ potential

First LIT in the region of the low-energy resonance

LITs in the resonance region with various  $\sigma_I$  (full curves);  
 comparison with single Lorentzians of corresponding  $\sigma_I$  (dashed curves)

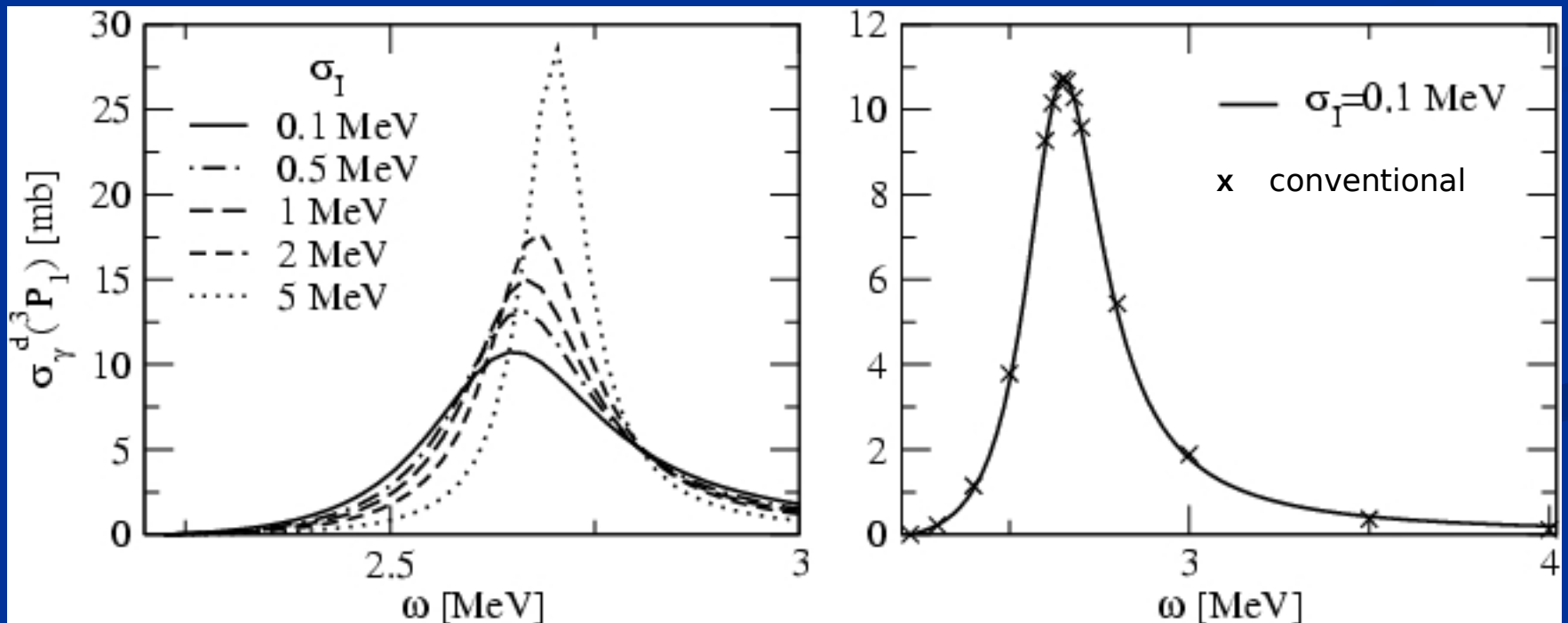


# Incomplete Inversion

Instead of using set  $\chi_m$  defined previously we take  $M_{\max}=1$  and take

$$\chi_1^{\text{res}} = \frac{1}{(E_{\text{np}} - E_{\text{res}})^2 + (\Gamma/2)^2} \left( \frac{1}{1 + \exp(-1)} - \frac{1}{1 + \exp((E_{\text{np}} - \alpha_3)/\alpha_3)} \right)$$

$E_{\text{res}}$ ,  $\Gamma$ , and  $\alpha_3$  are fit parameters



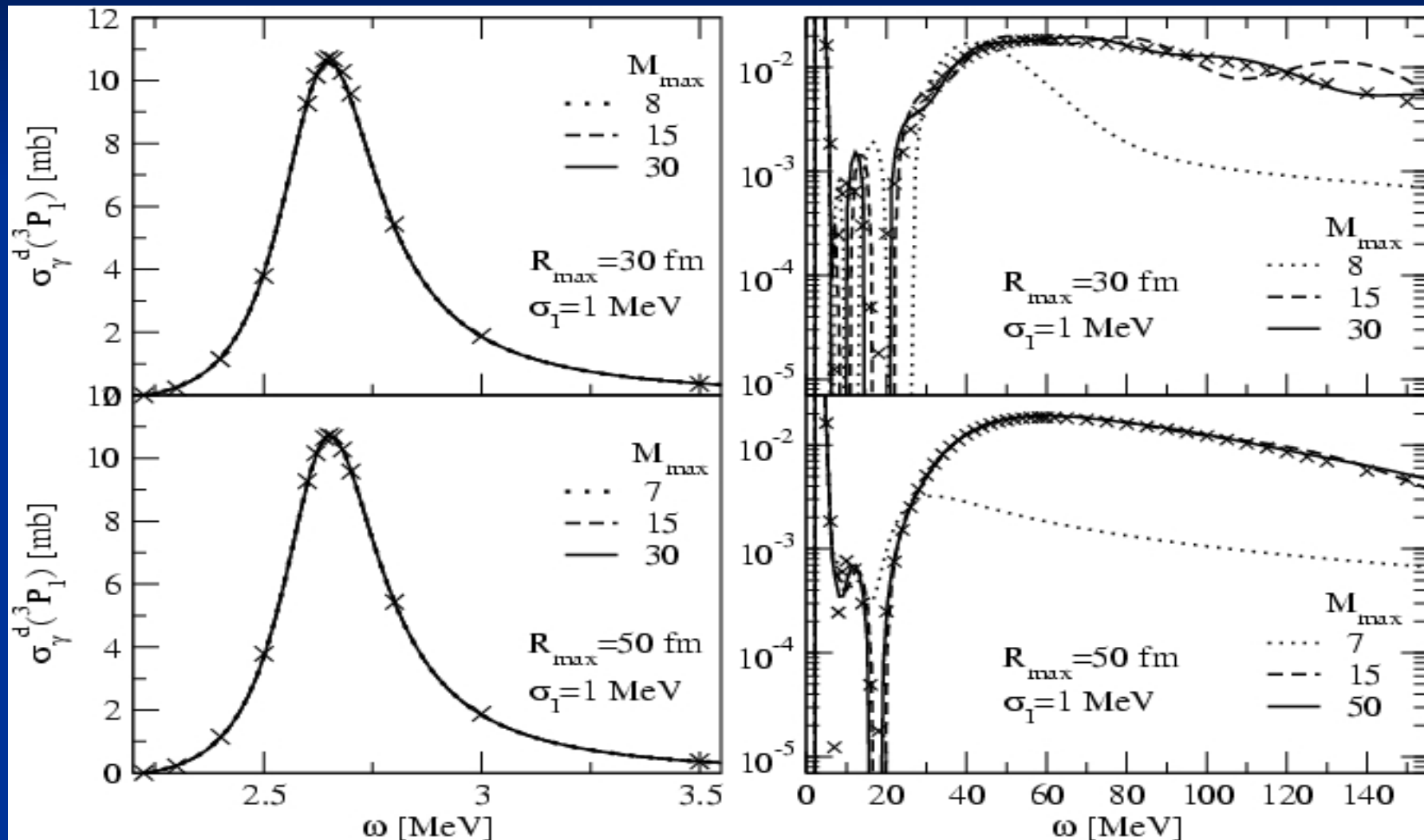
# Results with modified $^3P_1$ potential

Now to the LIT results beyond low-energy resonance



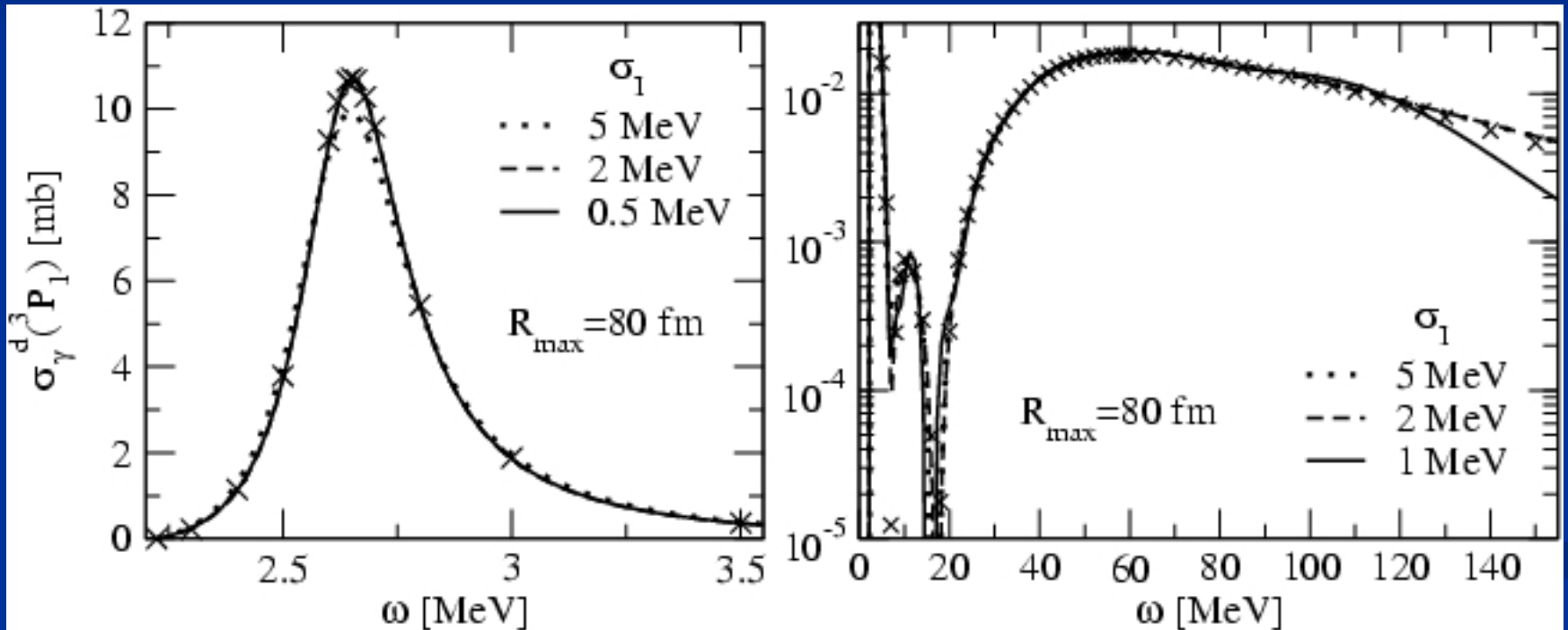
Complete inversion with set  $\chi_m$  defined previously using in addition as new first basis function  $\chi_1^{\text{res}}$

$\sigma_1 = 1 \text{ MeV}$ ,  $R_{\text{max}} = 30 \text{ and } 50 \text{ fm}$ , various  $M_{\text{max}}$



Complete inversion with set  $\chi_m$  defined previously using in addition as new first basis function  $\chi_1^{\text{res}}$

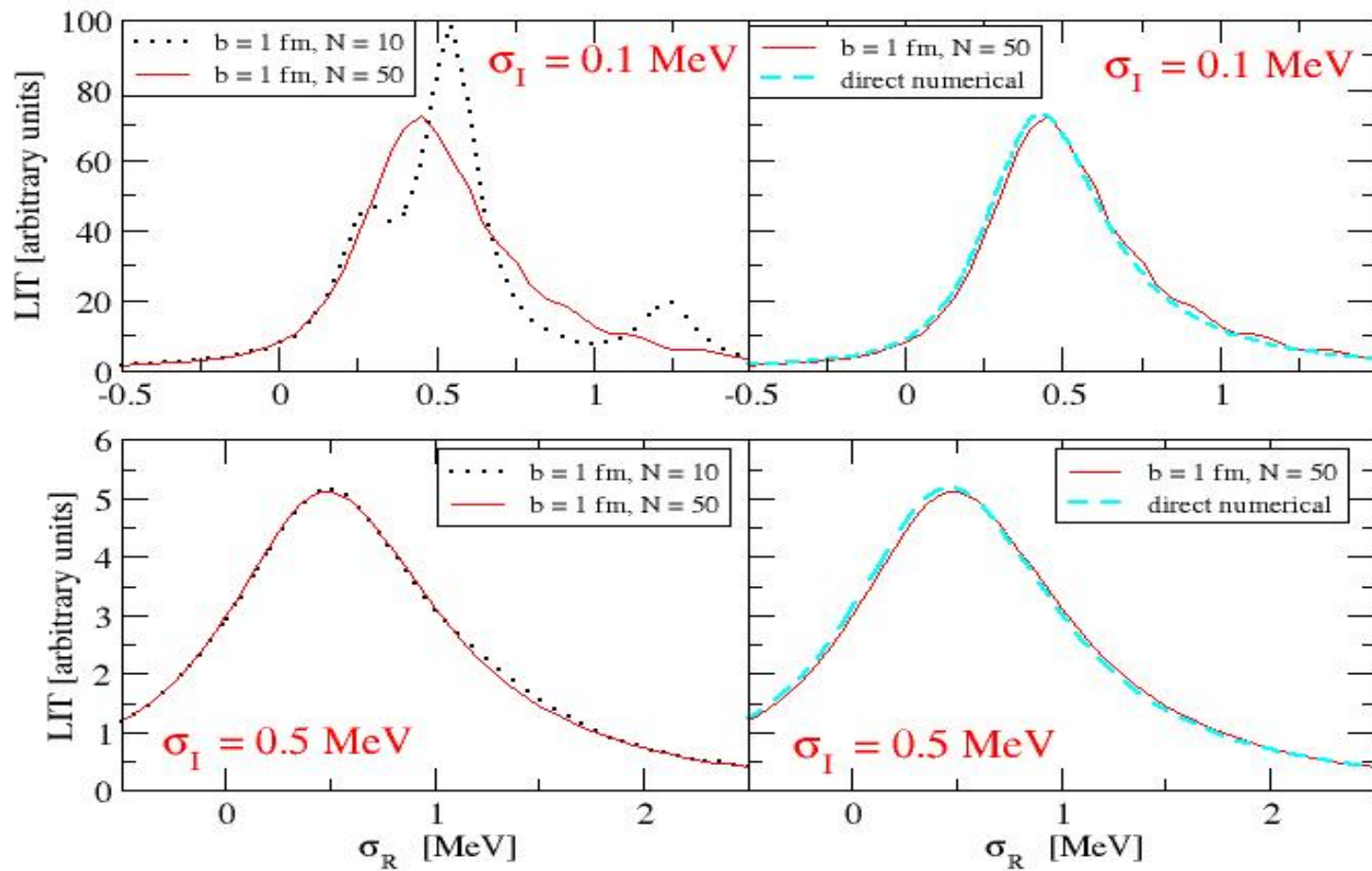
various  $\sigma_I$ ,  $R_{\text{max}} = 80$  fm,  $M_{\text{max}} = 30$



Up to now **direct numerical solutions** of Schrödinger equation for bound state and LIT equation for  $\tilde{\Psi}$

For  $A > 2$  it is more convenient to use **expansions** in complete sets using expansions in **HH** or **HO** functions

Expansion of radial part of  $\tilde{\Psi}$  in **Laguerre polynomials up to order N** times **an exponential fall-off  $\exp(-\rho/b)$**  with  $\rho \equiv r$



# Conclusions

- the **LIT** method opens up the possibility to carry out ab-initio calculations of reactions into the **A-body continuum for  $A > 2$**
- only **bound states** techniques are needed
- the LIT is a method with controlled resolution
- Exact determination of resonant shapes might be difficult, but is in principle possible