

Selfconsistent mean field calculations of the nuclear response using a realistic nucleon- nucleon interaction with a density dependent corrective term

P. Veselý and collaborators

INFN (Istituto Nazionale di Fisica Nucleare), Naples, Italy

Workshop “From Few-Nucleon Forces to Many-Nucleon Structure“

Trento, June 2013

List of collaborators: F. Knapp, N. Lo Iudice, D. Bianco, F. Andreozzi, A. Porino

Basic ingredients: HFB, TDA, RPA methods, realistic NN interaction, density-dependent (DD) phenomenological interaction as simulation of the 3-body terms, spin-orbit (s.o.) splitting, TDA phonons, Equation of Motion, Hilbert space divided to 0,1,2,3,...,n – phonon subspaces, exact treating of Pauli principle (no any quasiboson approximation)

Motivation of study: using the self-consistent methods like HFB, TDA, RPA in the context of the realistic NN potentials, comparison of TDA & RPA, dependence of the s.p. and TDA spectrum on the phenomenological DD and s.o. terms, TDA phonons as a basis for calculations beyond one-phonon approaches, fragmentation of giant resonance, pygmy resonances and possibly other types of transitions with strong anharmonic effects

Formalism & DD interaction

$$H = T_{int} + h_{s.o.} + V + V^{DD}[\rho]$$

$$T_{int} = \left(1 - \frac{1}{A}\right) \sum_i \frac{p_i^2}{2m} - \frac{1}{mA} \sum_{i < j} p_i \cdot p_j$$

intrinsic kinetic term

$$h_{s.o.} = Cl.s$$

phenom. spin-orbit term

2-body “realistic” NN interaction

2-body density-dependent int.

$$V^{DD}[\rho] = \frac{C_{3N}}{6} (1 + P_\sigma) \rho \left(\frac{\mathbf{r}_1 + \mathbf{r}_2}{2} \right) \delta^{(3)}(\mathbf{r}_1 - \mathbf{r}_2)$$

2-body density-dependent int.
contact in S=1, T=0 channel

$$V^{3N} = C_{3N} \delta^{(3)}(\mathbf{r}_1 - \mathbf{r}_2) \delta^{(3)}(\mathbf{r}_2 - \mathbf{r}_3)$$

the same contribution to E_{HF}
energy as pure contact 3-body
interaction



Formalism & DD interaction

$$H = T_{int} + h_{s.o.} + V + V^{DD}[\rho]$$

$$T_{int} = \left(1 - \frac{1}{A}\right) \sum_i \frac{p_i^2}{2m} - \frac{1}{mA} \sum_{i < j} p_i \cdot p_j$$

intrinsic kinetic term

$$h_{s.o.} = Cl.s$$

phenom. spin-orbit term

2-body “realistic” NN interaction

2-body density-dependent int.

$$V^{DD}[\rho] = \frac{C_0}{6} (1 + P_\sigma) \rho \left(\frac{\mathbf{r}_1 + \mathbf{r}_2}{2} \right) \delta^{(3)}(\mathbf{r}_1 - \mathbf{r}_2) + \frac{C_1}{6} (1 - P_\sigma) \rho \left(\frac{\mathbf{r}_1 + \mathbf{r}_2}{2} \right) \delta^{(3)}(\mathbf{r}_1 - \mathbf{r}_2)$$

(S=1, T=0) term

(S=0, T=1) term

The DD interaction can be introduced phenomenologically without any relation to some 3-body term.



Formalism & DD interaction



The interaction matrix elements for DD force – like-wise particles & proton-neutron

$$\langle l_i j_i, J | \hat{V}^{DD} | l_k j_k, J \rangle_{pp} = \frac{V_1}{4\pi} \frac{1}{12(2J+1)} \frac{1}{\sqrt{(1+\delta_{ij})(1+\delta_{kl})}} \sqrt{(2j_i+1)(2j_j+1)(2J_k+1)(2j_l+1)} I_{rad} \\ (-1)^{l_j+l_i+j_j+j_l} (1+(-1)^{l_i+l_j+J}) (1+(-1)^{l_k+l_l+J}) C_{j_i 1/2 j_j -1/2}^{J0} C_{j_k 1/2 j_l -1/2}^{J0}$$

$$\langle l_i j_i, J | \hat{V}^{DD} | l_k j_k, J \rangle_{pn} = \frac{V_0}{4\pi} \frac{1}{24(2J+1)} \sqrt{(2j_i+1)(2j_j+1)(2J_k+1)(2j_l+1)} I_{rad} [2(-1)^{j_i+j_j+j_k+j_l} (1+(-1)^{l_i+l_j+l_k+l_l}) \\ C_{j_i 1/2 j_j 1/2}^{J1} C_{j_k 1/2 j_l 1/2}^{J1} + (-1)^{l_i+l_k+j_i+j_k+1} (1+(-1)^{l_i+l_j+J+1}) (1+(-1)^{l_k+l_l+J+1}) C_{j_i 1/2 j_j -1/2}^{J0} C_{j_k 1/2 j_l -1/2}^{J0}] \\ + \frac{V_1}{4\pi} \frac{1}{24(2J+1)} \sqrt{(2j_i+1)(2j_j+1)(2J_k+1)(2j_l+1)} I_{rad} \\ (-1)^{l_j+l_i+j_j+j_l} (1+(-1)^{l_i+l_j+J}) (1+(-1)^{l_k+l_l+J}) C_{j_i 1/2 j_j -1/2}^{J0} C_{j_k 1/2 j_l -1/2}^{J0}$$

radial integral

$$I_{rad} = \sum_{n m l m j_m} V_{n m l m j_m}^2 \frac{2j_m+1}{4\pi} \int dr r^2 R_{n_i l_i}(r) R_{n_j l_j}(r) R_{n_k l_k}(r) R_{n_l l_l}(r) R_{n_m l_m}(r) R_{n_m l_m}(r)$$

Formalism & DD interaction

relation DD 2-body & 3-body interaction

$$V_{ijkl}^{DD,pp} = \sum_{m < E_F} V_{ijmklm}^{ppp} + \sum_{m < E_F} V_{ijmklm}^{ppn}$$

$$V_{ijkl}^{DD,nn} = \sum_{m < E_F} V_{ijmklm}^{nnn} + \sum_{m < E_F} V_{mijmkl}^{pnn}$$

$$V_{ijkl}^{DD,pn} = \sum_{m < E_F} V_{mijmkl}^{ppn} + \sum_{m < E_F} V_{ijmklm}^{pnn}$$



Due to Wick theorem:

- contribution to HF energy

$$E_{HF} = \frac{1}{3} \left(\frac{1}{2} \sum_{ij < E_F} V_{ijij}^{DD,pp} + \frac{1}{2} \sum_{ij < E_F} V_{ijij}^{DD,nn} + \sum_{ij < E_F} V_{ijij}^{DD,pn} \right)$$

- contribution to s.p. energies:

$$e_i^p = \frac{1}{2} \sum_{k < E_F} \left(V_{ikik}^{DD,pp} + V_{ikik}^{DD,pn} \right)$$

- contribution to 2-body part of Hamiltonian:

$$\begin{aligned} \hat{H}^{(2)} = & \frac{1}{4} \sum_{ijkl} \left(V_{ijkl}^{pp} + V_{ijkl}^{DD,pp} \right) a_i^\dagger a_j^\dagger a_k a_l \\ & + \frac{1}{4} \sum_{ijkl} \left(V_{ijkl}^{nn} + V_{ijkl}^{DD,nn} \right) a_i^\dagger a_j^\dagger a_k a_l \\ & + \sum_{ijkl} \left(V_{ijkl}^{pn} + V_{ijkl}^{DD,pn} \right) a_i^\dagger a_j^\dagger a_k a_l \end{aligned}$$

Formalism & DD interaction

relation DD 2-body & 3-body interaction

$$V_{ijkl}^{DD,pp} = \sum_{m < E_F} V_{ijmklm}^{ppp} + \sum_{m < E_F} V_{ijmklm}^{ppn}$$

$$V_{ijkl}^{DD,nn} = \sum_{m < E_F} V_{ijmklm}^{nnn} + \sum_{m < E_F} V_{mijmkl}^{pnn}$$

$$V_{ijkl}^{DD,pn} = \sum_{m < E_F} V_{mijmkl}^{ppn} + \sum_{m < E_F} V_{ijmklm}^{pnn}$$



Due to Wick theorem:

- contribution to HF energy

$$E_{HF} = \frac{1}{3} \left(\frac{1}{2} \sum_{ij < E_F} V_{ijij}^{DD,pp} + \frac{1}{2} \sum_{ij < E_F} V_{ijij}^{DD,nn} + \sum_{ij < E_F} V_{ijij}^{DD,pn} \right)$$

- contribution to s.p. energies:

$$e_i^p = \frac{1}{2} \sum_{k < E_F} \left(V_{ikik}^{DD,pp} + V_{ikik}^{DD,pn} \right)$$

- contribution to 2-body part of Hamiltonian:

$$\begin{aligned} \hat{H}^{(2)} = & \frac{1}{4} \sum_{ijkl} \left(V_{ijkl}^{pp} + V_{ijkl}^{DD,pp} \right) a_i^\dagger a_j^\dagger a_k a_l \\ & + \frac{1}{4} \sum_{ijkl} \left(V_{ijkl}^{nn} + V_{ijkl}^{DD,nn} \right) a_i^\dagger a_j^\dagger a_k a_l \\ & + \sum_{ijkl} \left(V_{ijkl}^{pn} + V_{ijkl}^{DD,pn} \right) a_i^\dagger a_j^\dagger a_k a_l \end{aligned}$$

Formalism & EMPM model

Basic building blocks - TDA phonons

$$O_{\lambda}^{\dagger} = \sum_{ph} c_{ph}^{\lambda} a_p^{\dagger} a_{\bar{h}}$$

Construction of n-phon states
as 1phon excitation on (n-1)-phon

$$|n; \beta\rangle = \sum_{\lambda\alpha} C_{\lambda\alpha}^{(\beta)} O_{\lambda}^{\dagger} |n-1; \alpha\rangle$$

EoM: generate n-phon energies

$$\langle n, \beta | [H, O_{\lambda}^{\dagger}] | n-1, \alpha \rangle = (E_{\beta} - E_{\alpha}) X_{\lambda\alpha}^{(\beta)}$$

Transitions from one correlated state to another correlated state:

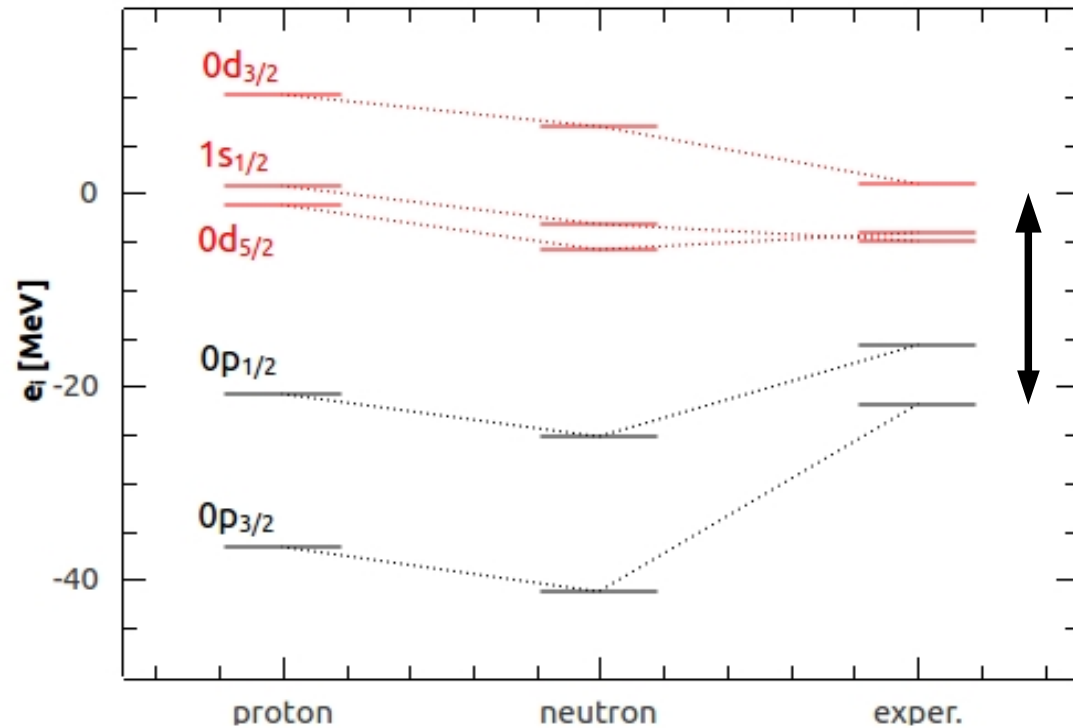
$$\begin{aligned} \mathcal{M}_{\lambda}(i \rightarrow f) &= \langle \Psi_{v_f} | \mathcal{M}_{\lambda} | \Psi_{v_i} \rangle \\ &= \sum_{\beta_i \beta_f} C_{\beta_i}^{(v_i)} C_{\beta_f}^{(v_f)} \langle n_f, \beta_f | \mathcal{M}_{\lambda} | n_i, \alpha_i \rangle \end{aligned}$$



$E^{(0)}$	$\langle 0ph H 1ph \rangle$	$\langle 0ph H 2ph \rangle$	0
$\langle 1ph H 0ph \rangle$	$E^{(1)}$ $E_1^{(1)}$ 0 $E_2^{(1)}$	$\langle 1ph H 2ph \rangle$	$\langle 1ph H 3ph \rangle$
	0	$E_k^{(1)}$	
$\langle 2ph H 0ph \rangle$	$\langle 2ph H 1ph \rangle$	$E^{(2)}$ $E_1^{(2)}$ 0 $E_2^{(2)}$	$\langle 2ph H 3ph \rangle$
		0	$E_k^{(2)}$
0	$\langle 3ph H 1ph \rangle$	$\langle 3ph H 2ph \rangle$	$E^{(3)}$ $E_1^{(3)}$ 0 $E_2^{(3)}$
			0
			$E_k^{(3)}$

Spectra without DD interaction

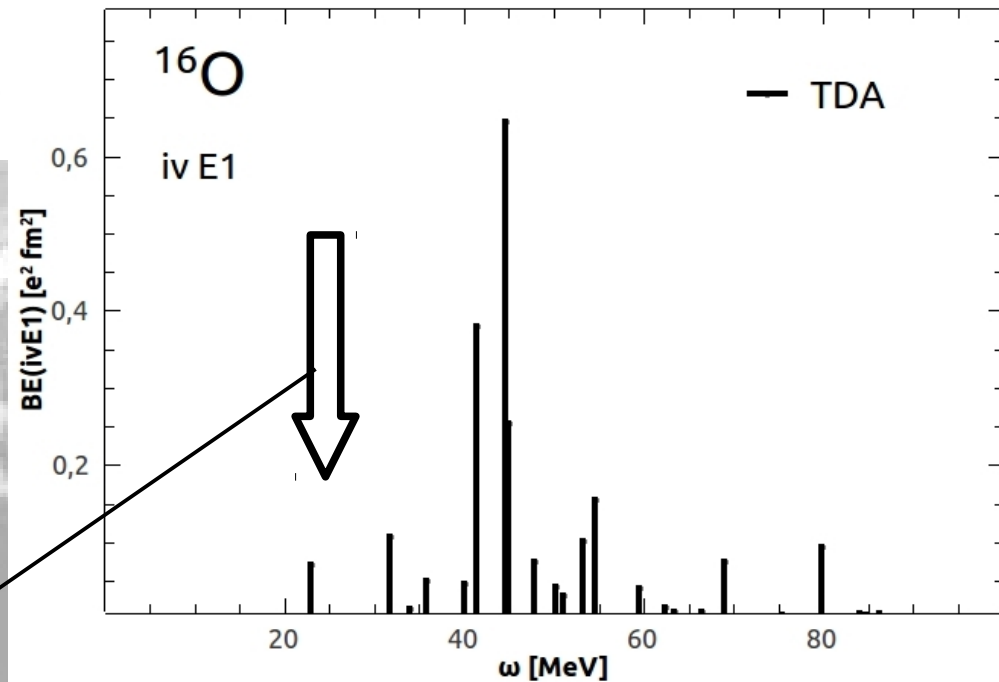
s.p. energies



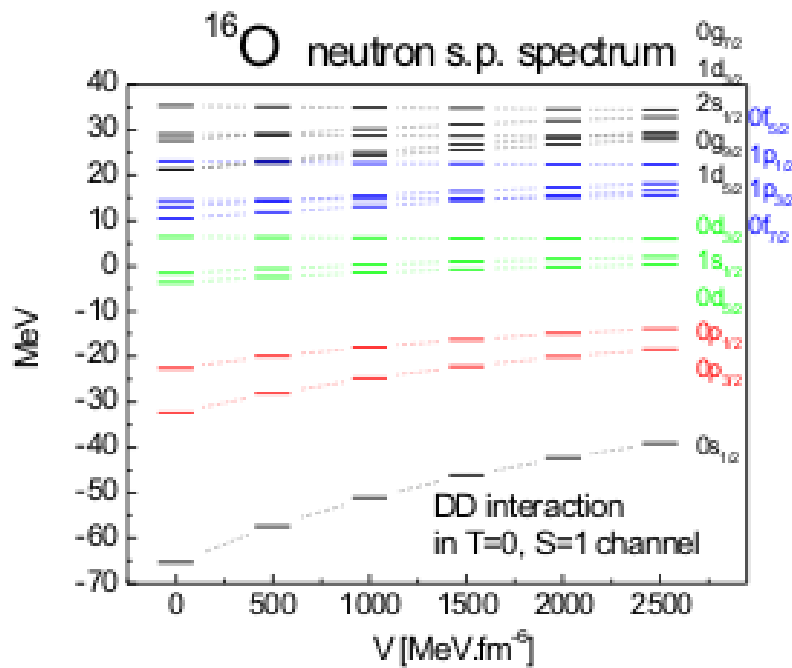
NN interaction: CD-Bonn + V_{lowk}
with $k_{cut-off} = 1.9 \text{ fm}^{-1}$

“experimental gap“ between p & (sd) shells much smaller than in HF calculation with realistic NN force

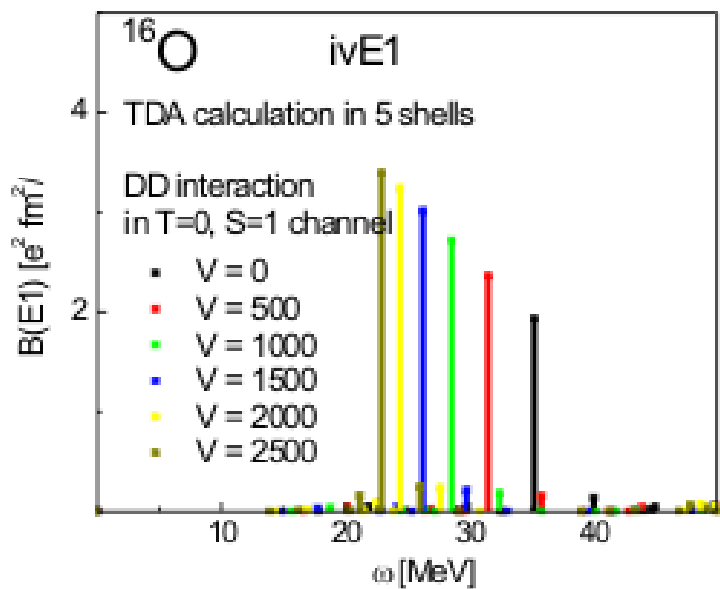
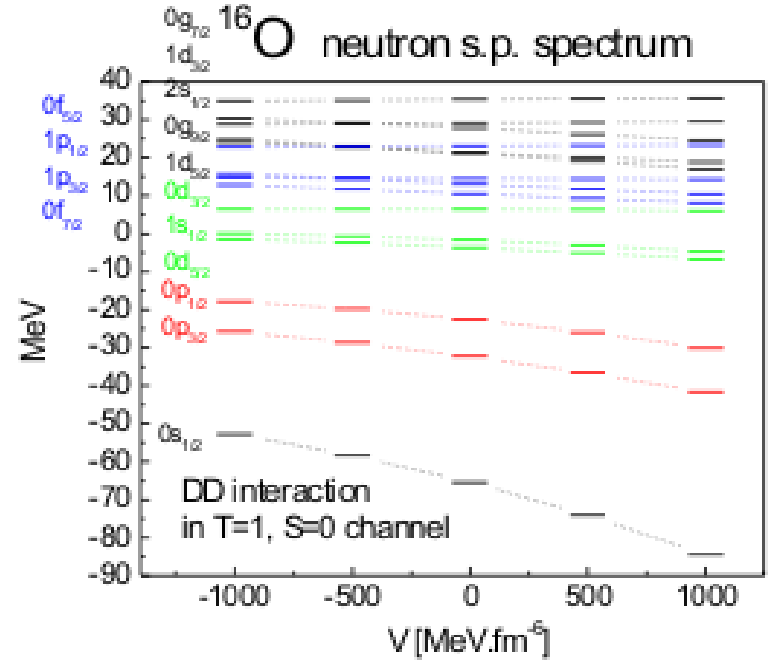
In experiment the E1 resonance occurs at $\sim 25 \text{ MeV}$



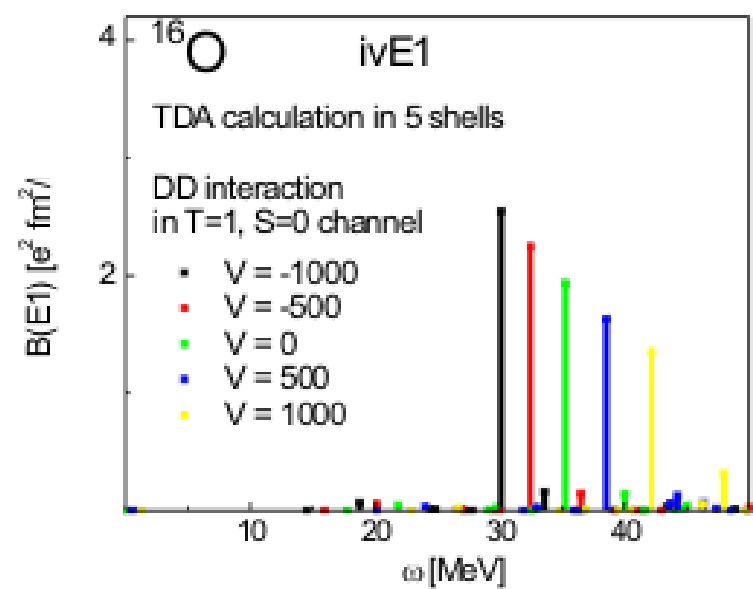
Test of the DD term



The spectrum becomes more realistic but



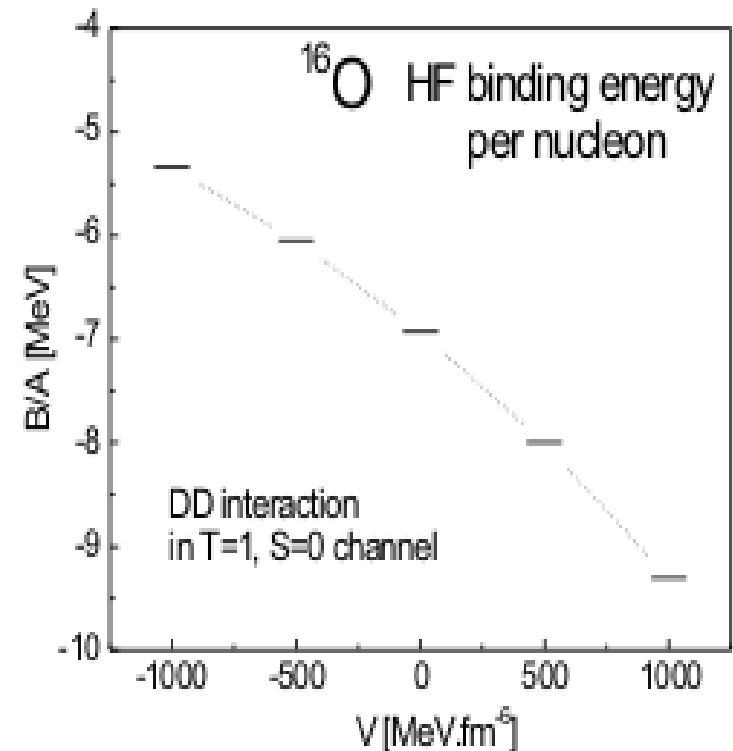
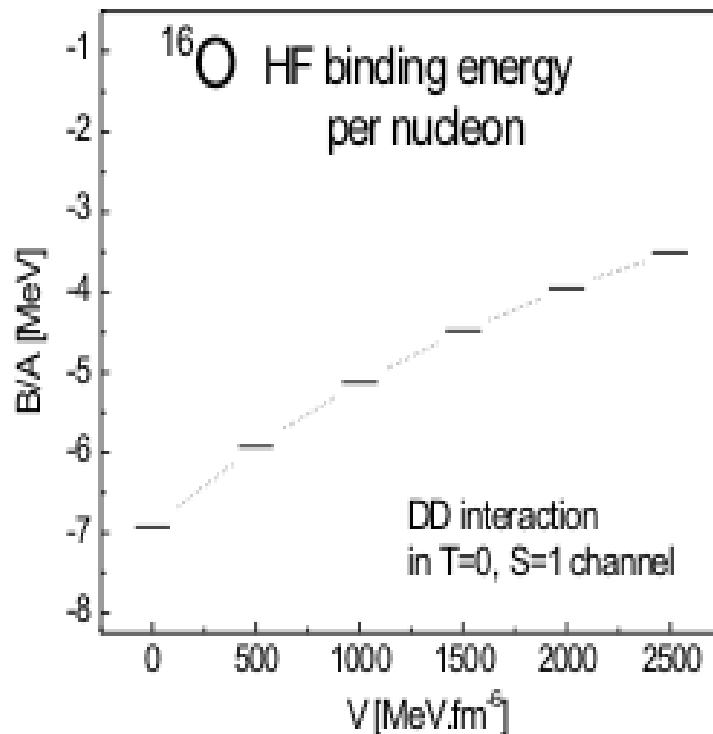
we pay a price for it (see next page)...



Test of the DD term

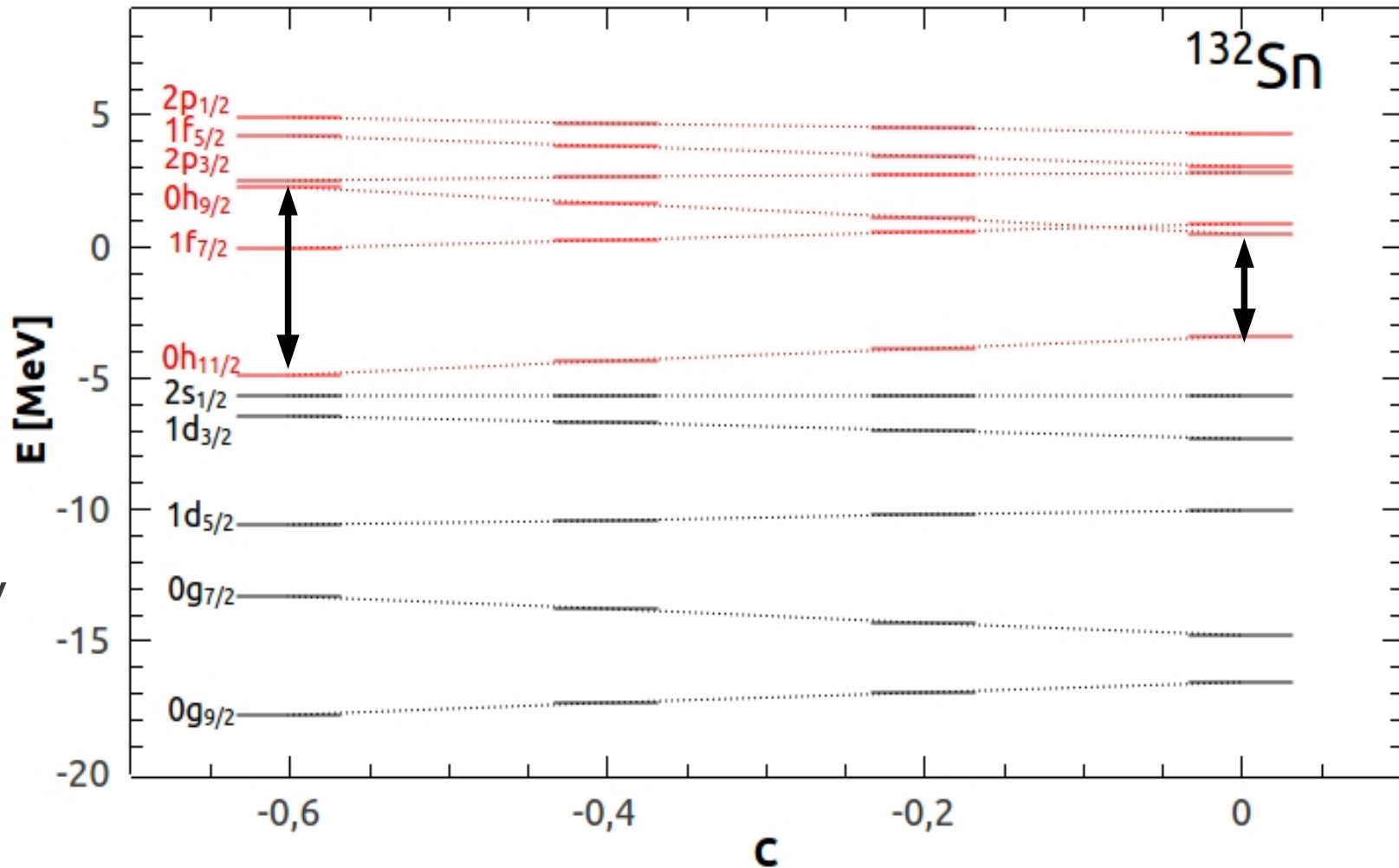
the better description of spectra the worse (i.e. less) of the binding energy (HF ground state energy per nucleon)

For ^{16}O $(B/A)_{\text{exp.}} = 7.98 \text{ MeV}$



Test of the spin-orbit term

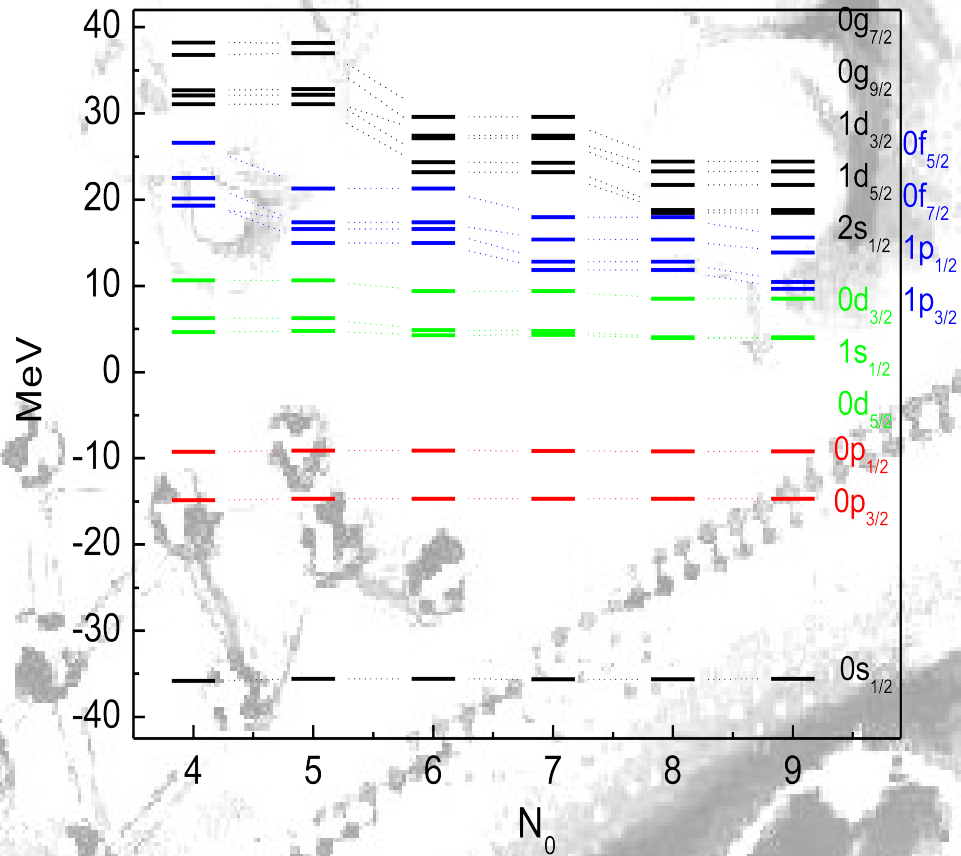
neutron s.p. energies



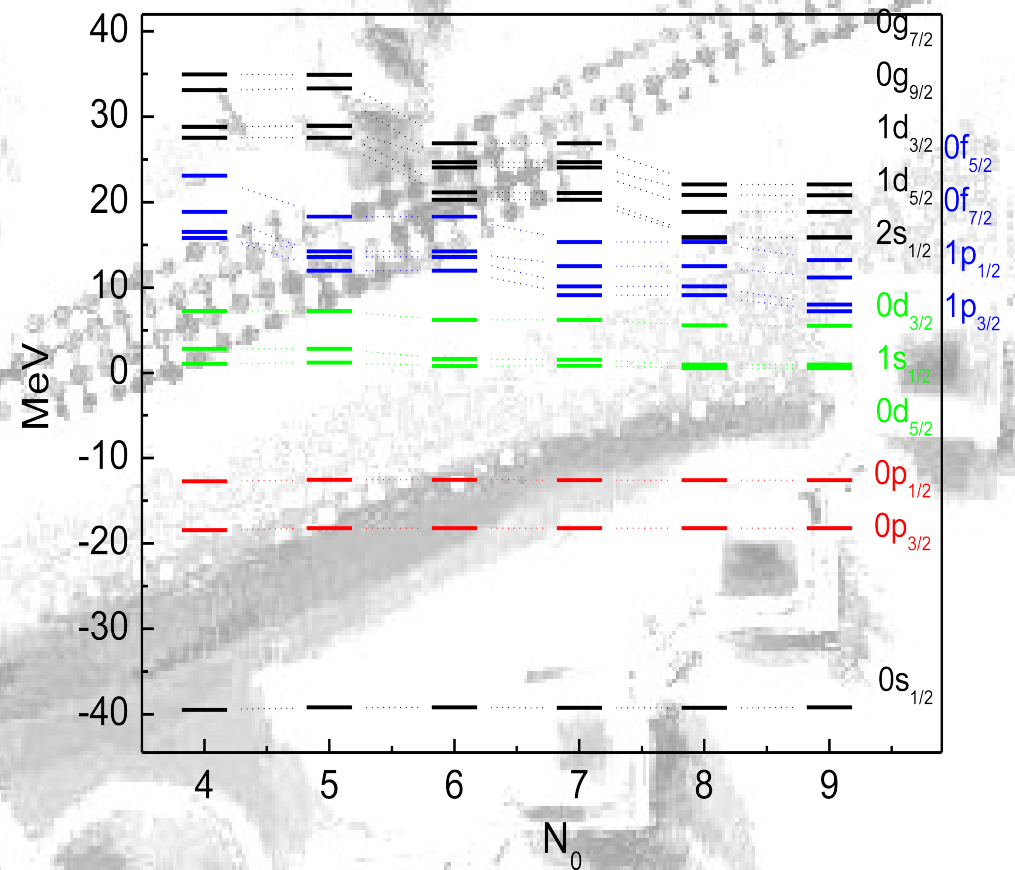
Example of spin-orbit splitting:
neutron $0h$ in ^{132}Sn
exper. = 6.53 MeV

Effect of the configuration space

^{16}O proton s.p. spectrum



^{16}O neutron s.p. spectrum



The s.p. spectrum is comprised with enlarging the configuration space.
(dissappearance of the shell structure above the Fermi level)

DD int. :

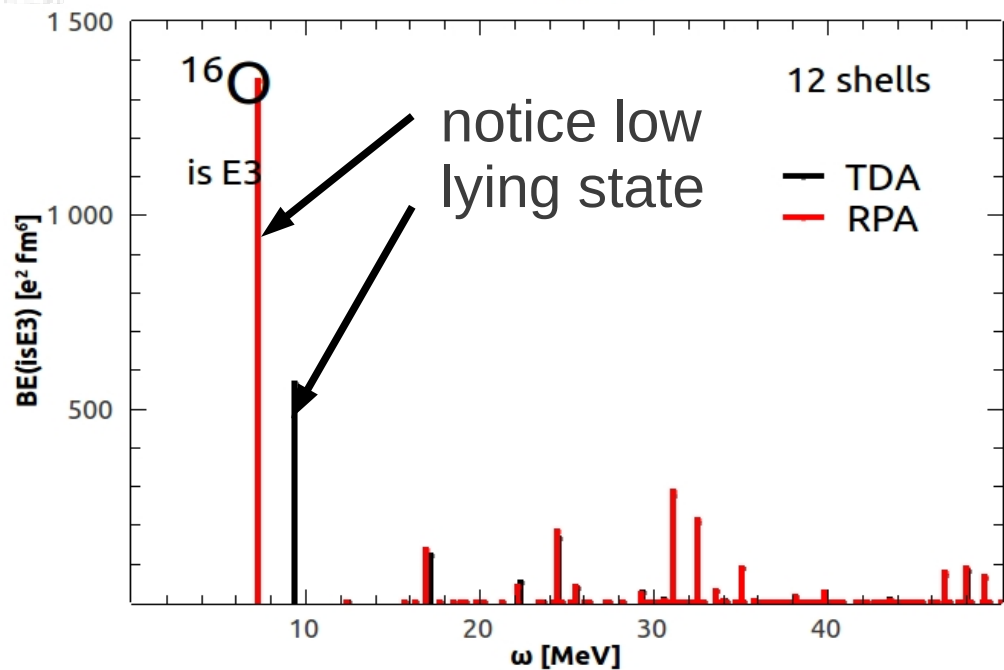
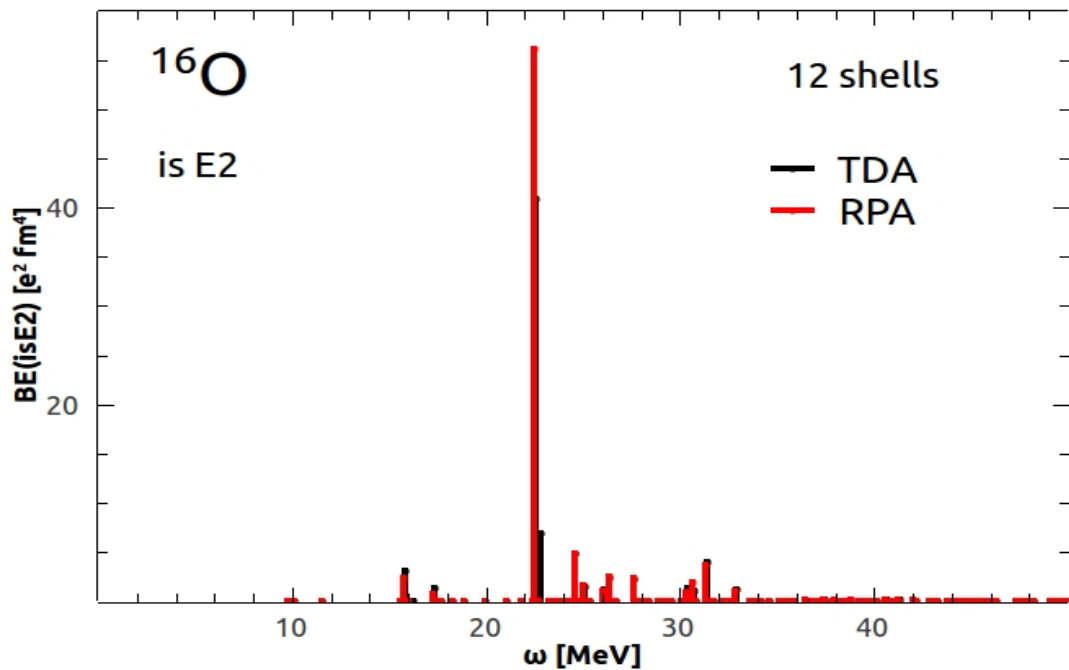
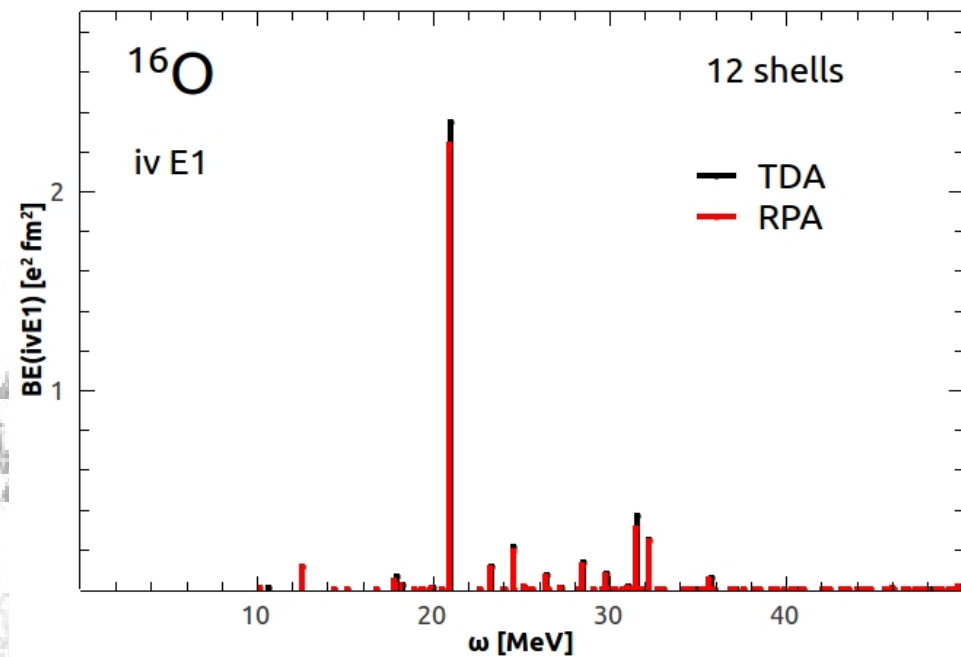
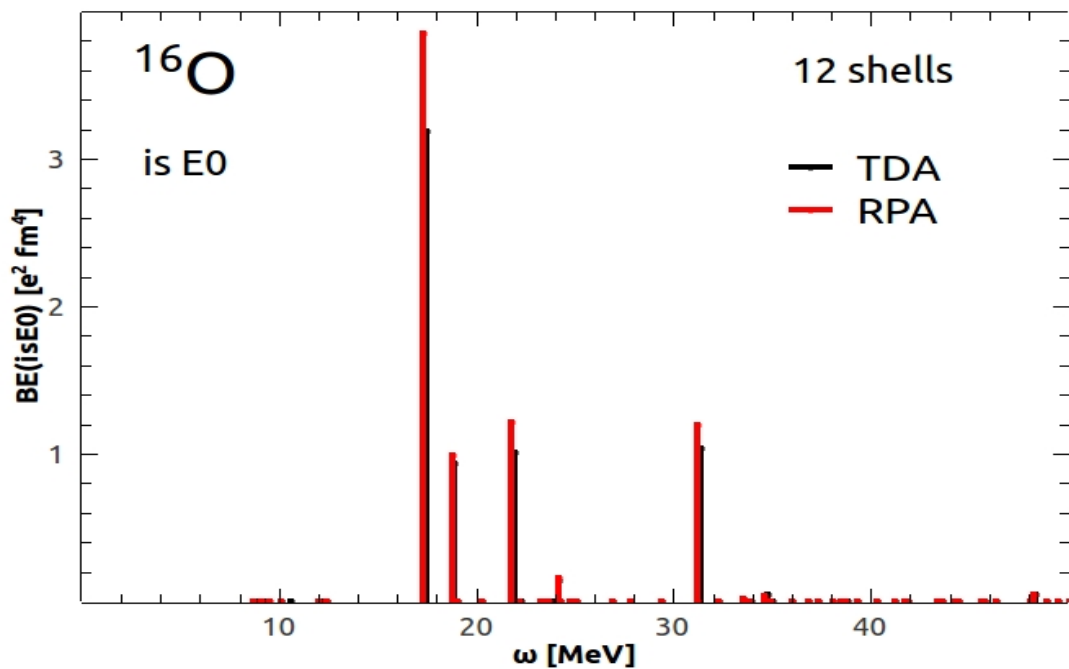
$$V(S=1, T=0) = 1300 \text{ MeV}\cdot\text{fm}^6$$

$$V(S=0, T=1) = -1300 \text{ MeV}\cdot\text{fm}^6$$

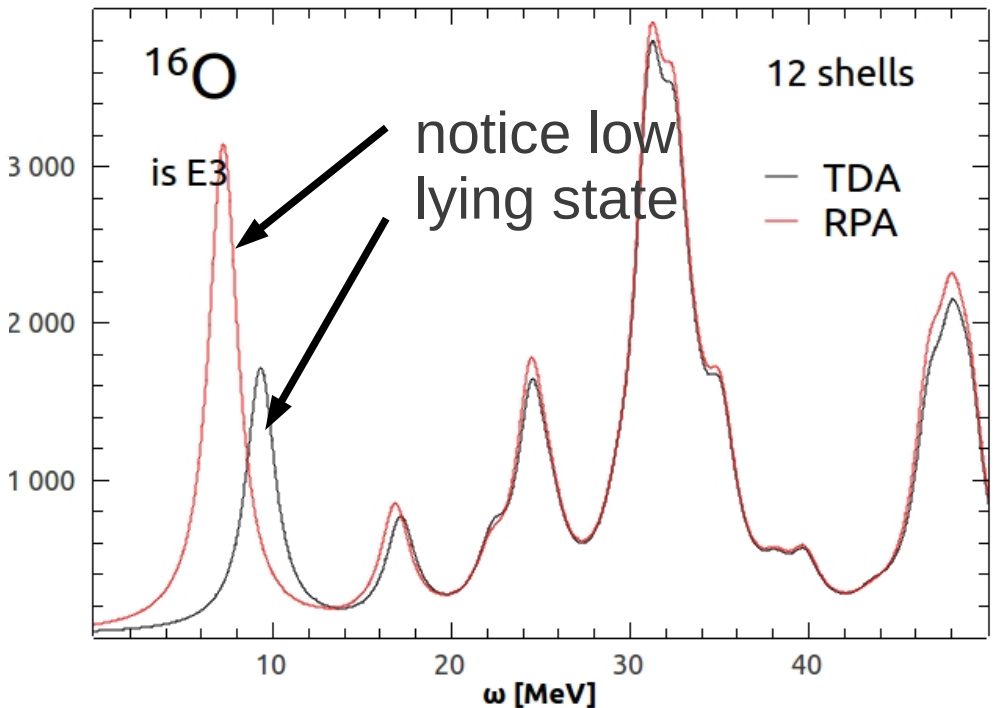
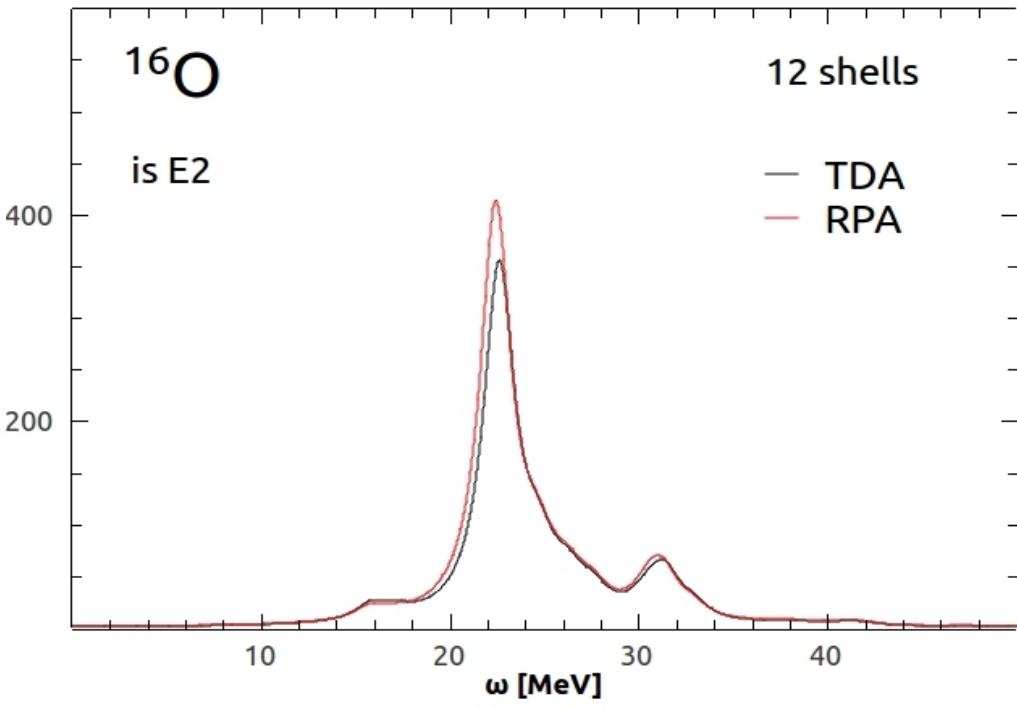
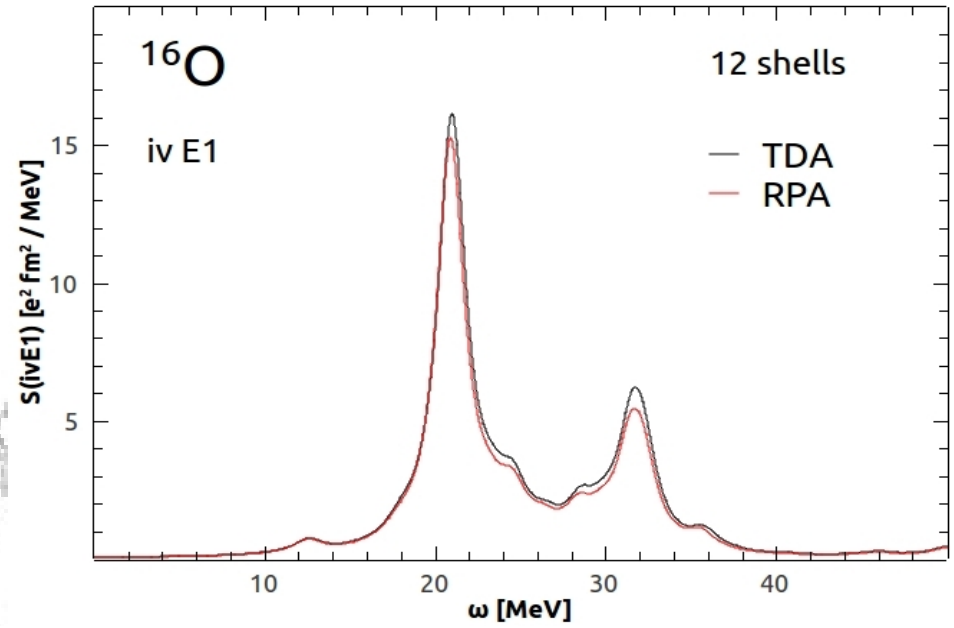
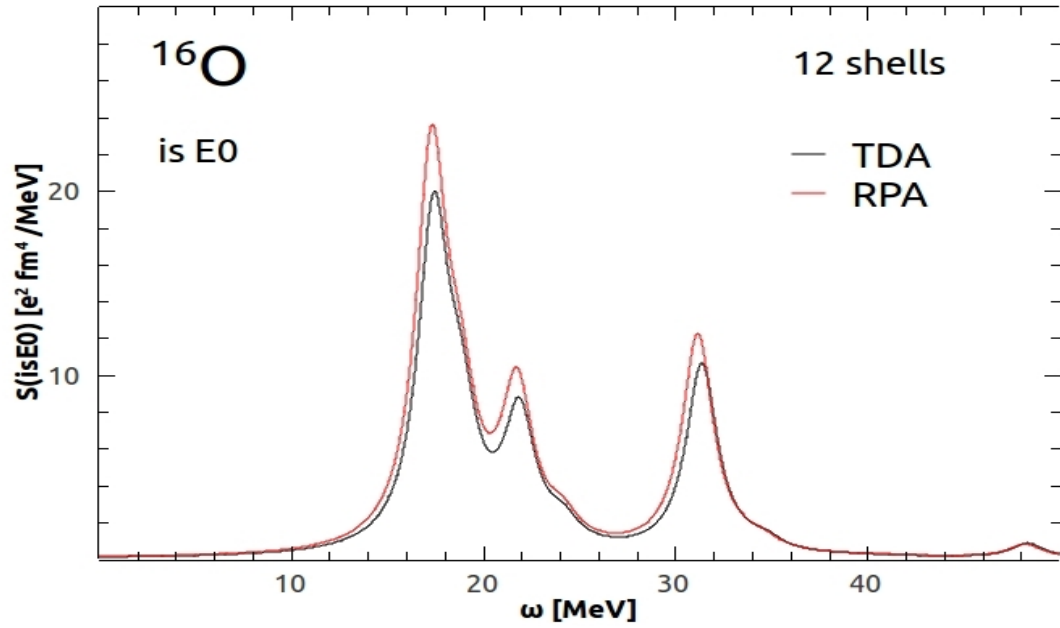
TDA & RPA methods

s.o. term :

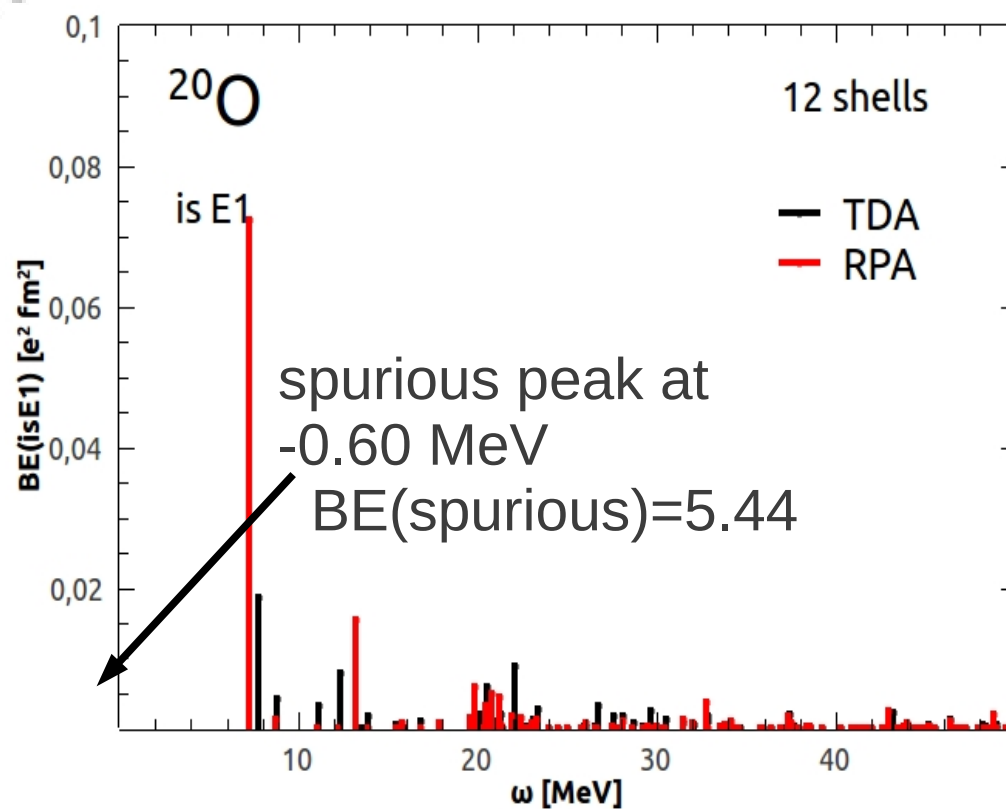
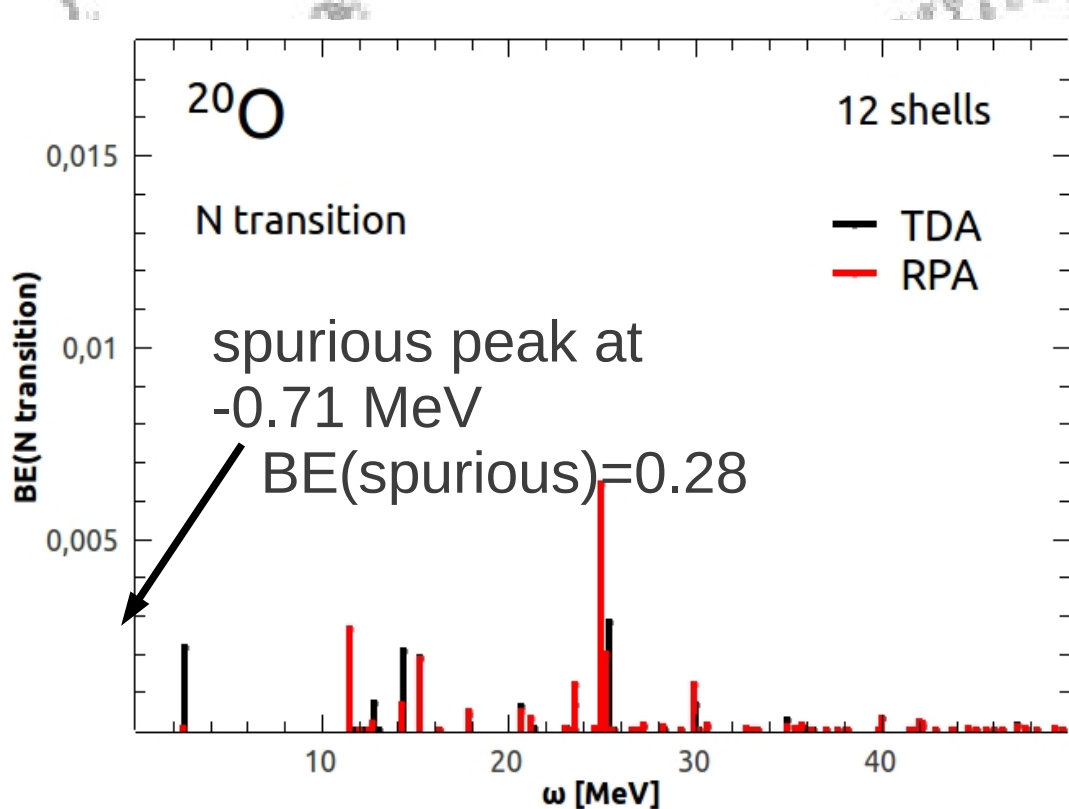
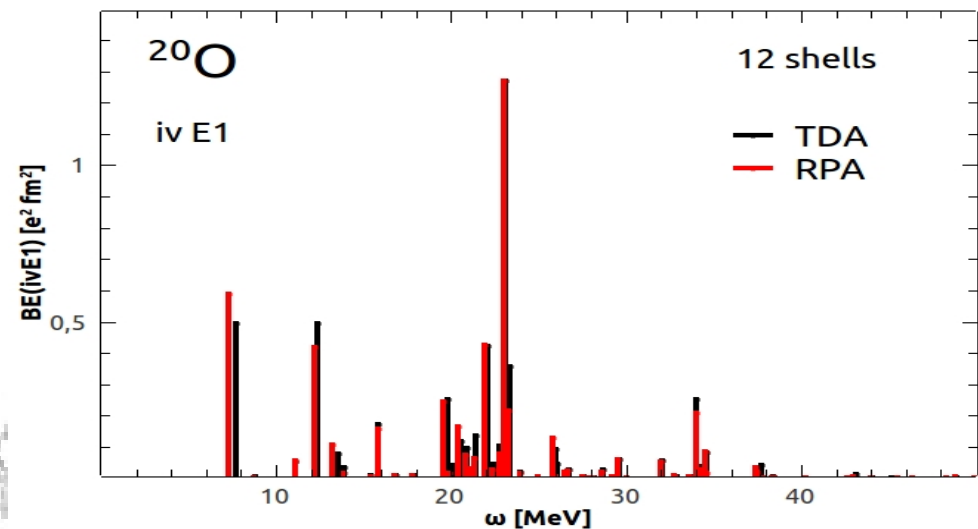
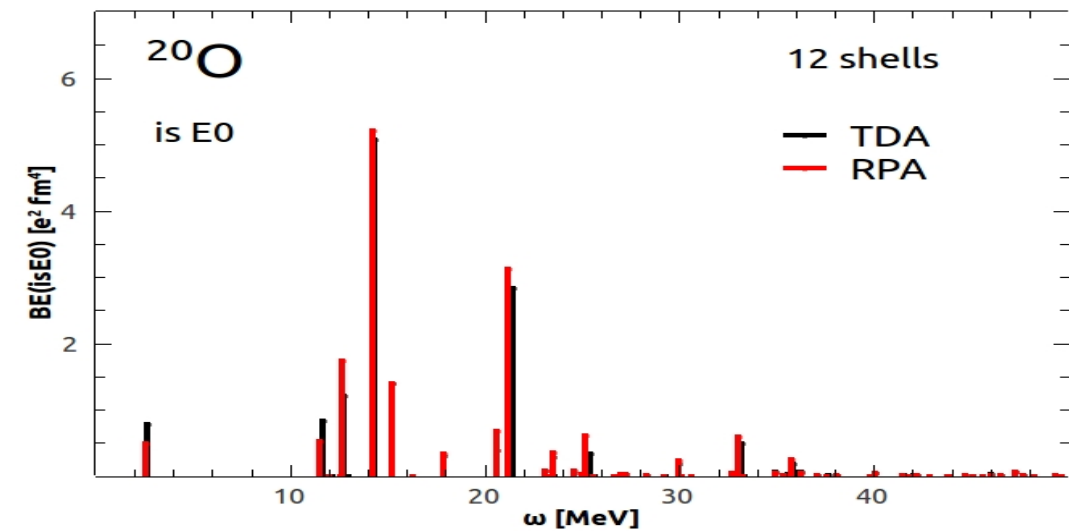
$$C = -0.9 \text{ MeV}$$



TDA & RPA methods (the same with Lorentz averaging)



TDA & RPA methods



DD int. :

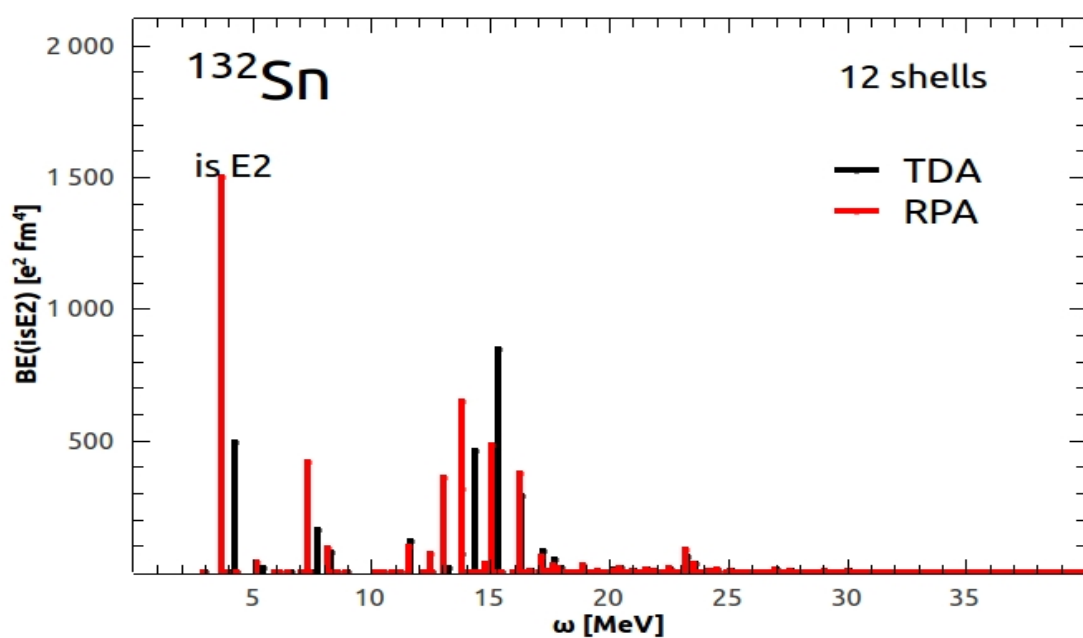
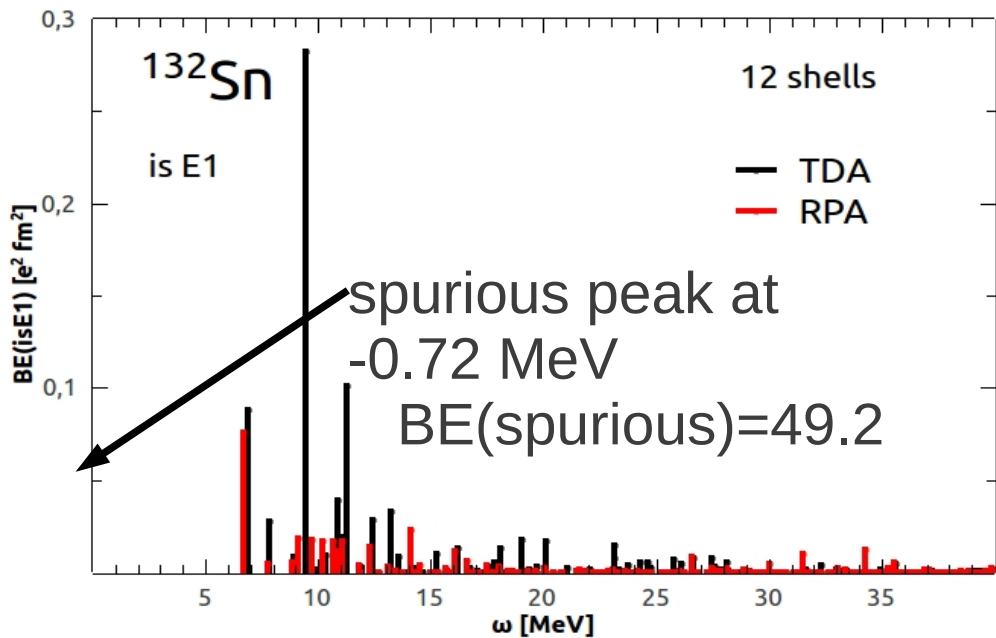
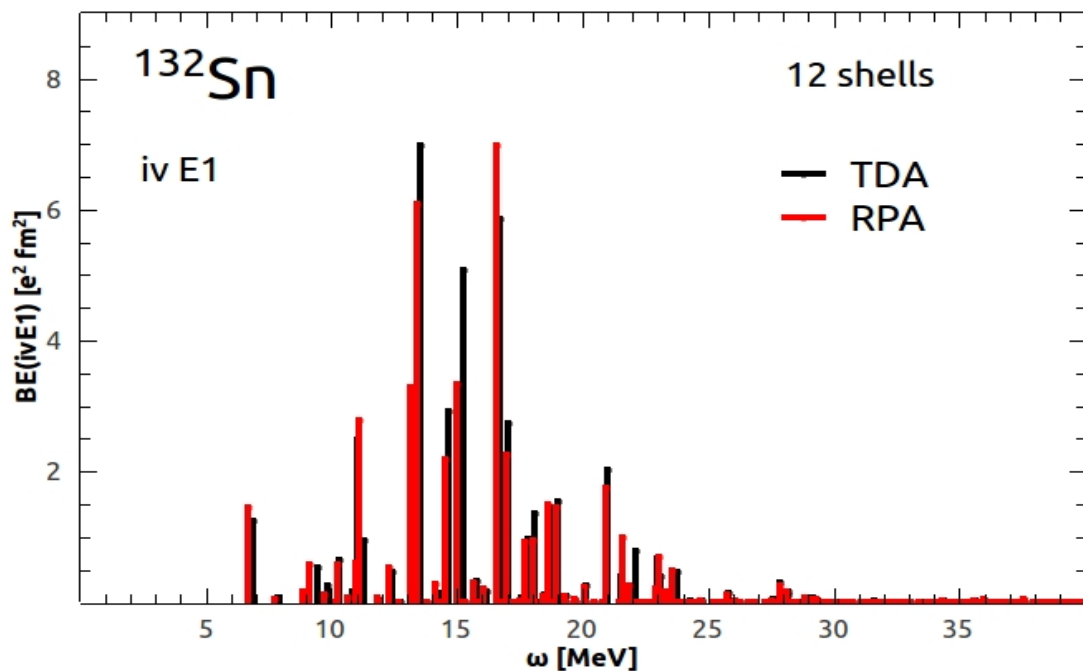
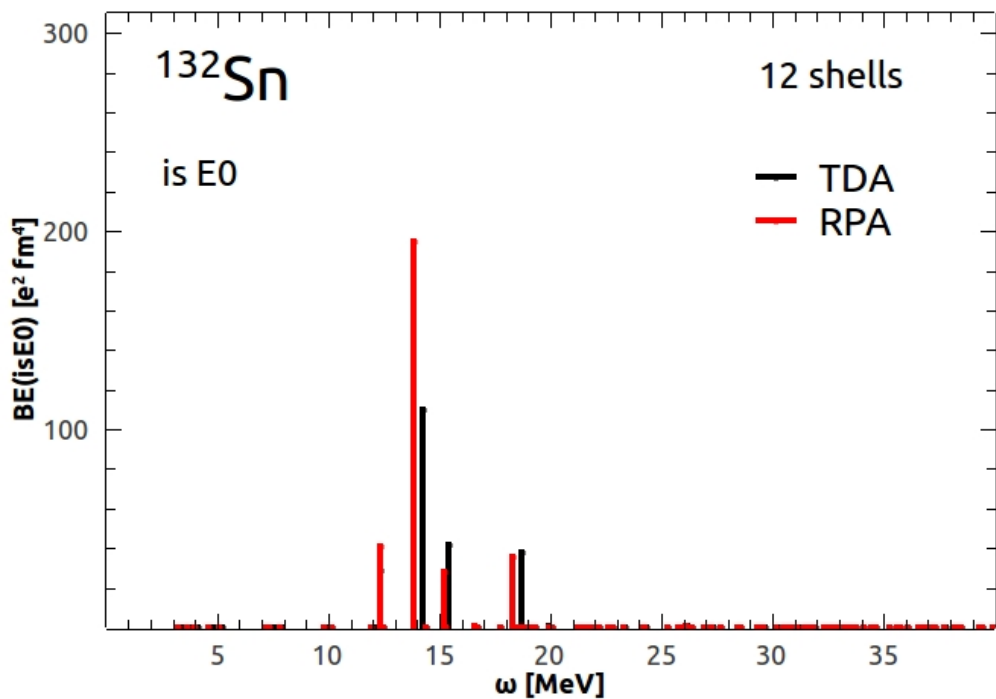
$$V(S=1, T=0) = 2400 \text{ MeV}\cdot\text{fm}^6$$

$$V(S=0, T=1) = -1800 \text{ MeV}\cdot\text{fm}^6$$

TDA & RPA methods

s.o. term :

$$C = -0.6 \text{ MeV}$$

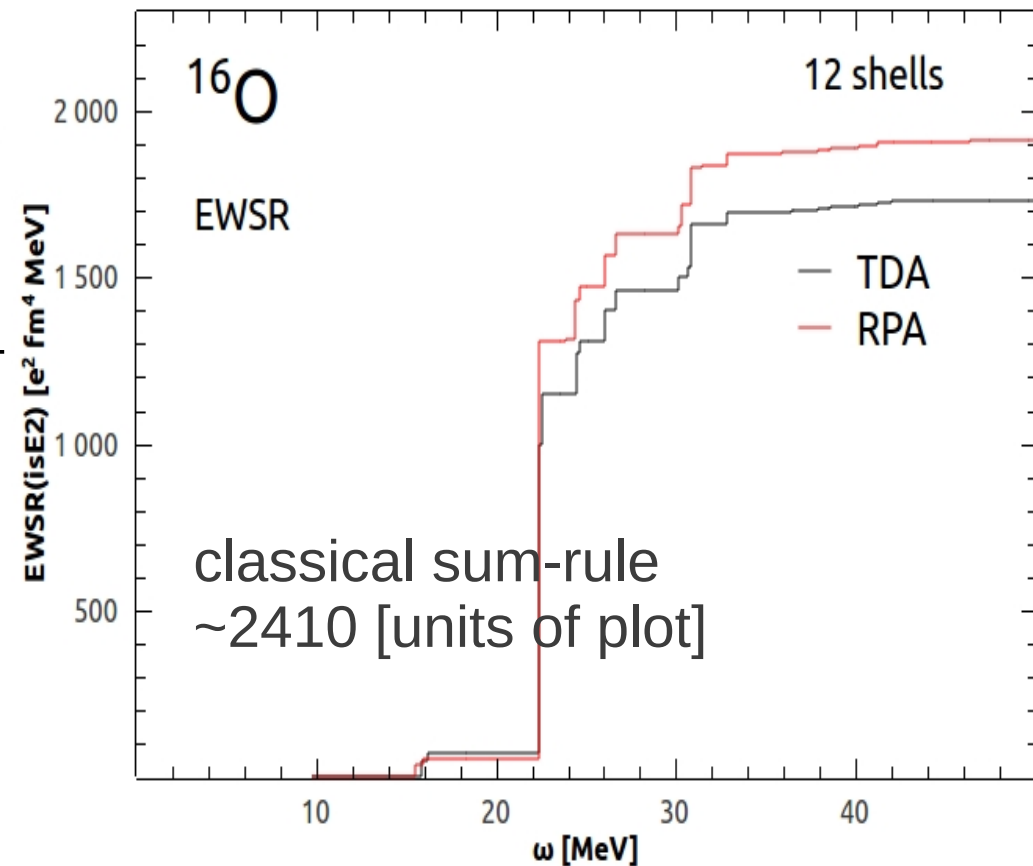
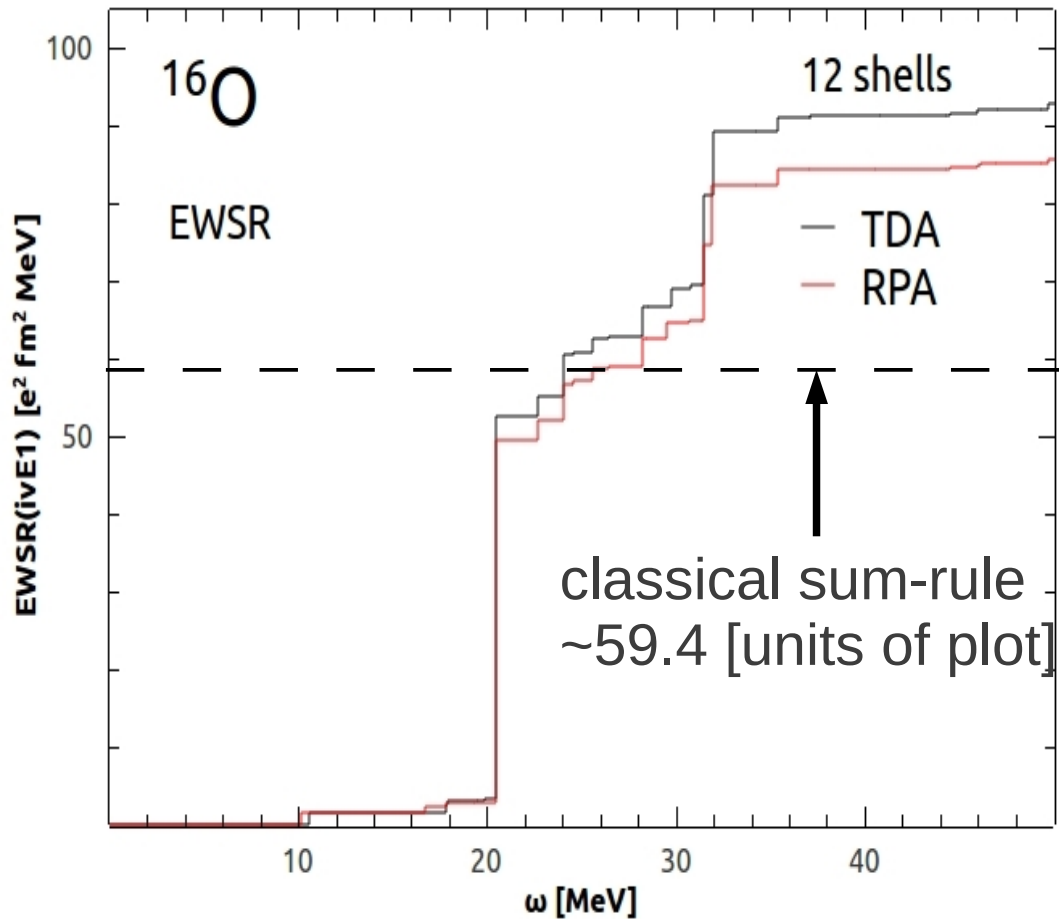


TDA & RPA methods

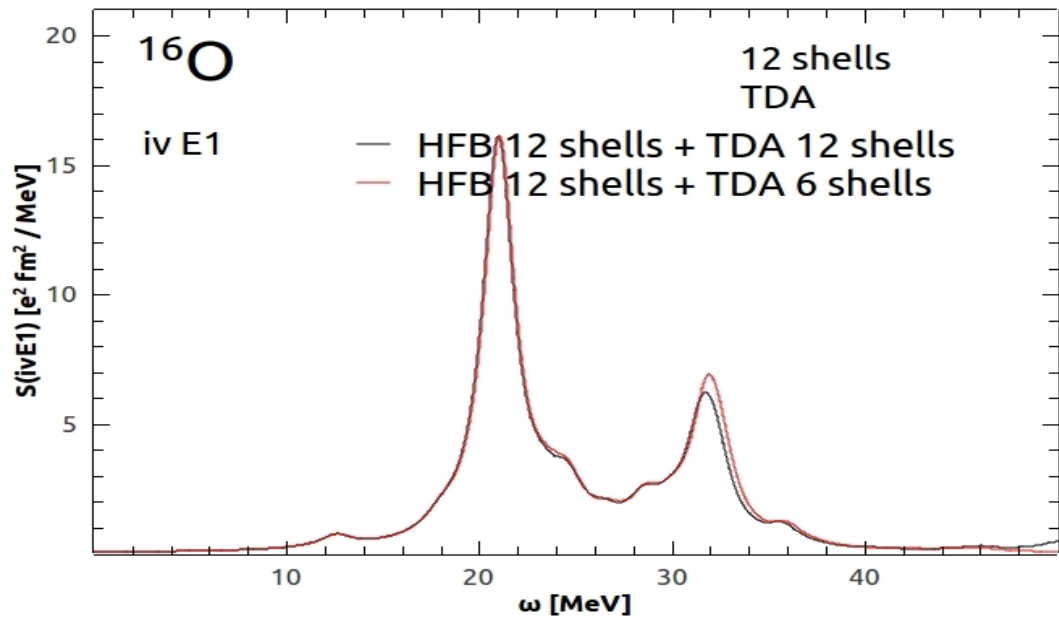
running energy-weighted sum rule (EWSR) for ivE1 and isE2

ivE1 ... overestimate
isE2 ... underestimate

Should EWSR be precisely reached?
EWSR derived for unperturbed g.s.
EWSR in shell model?



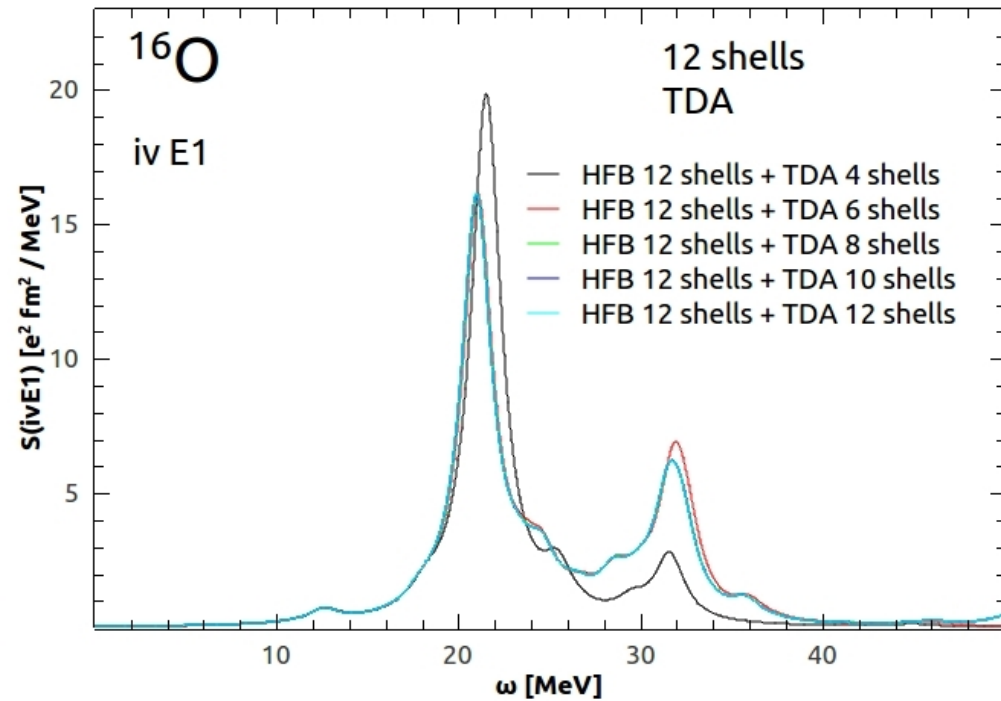
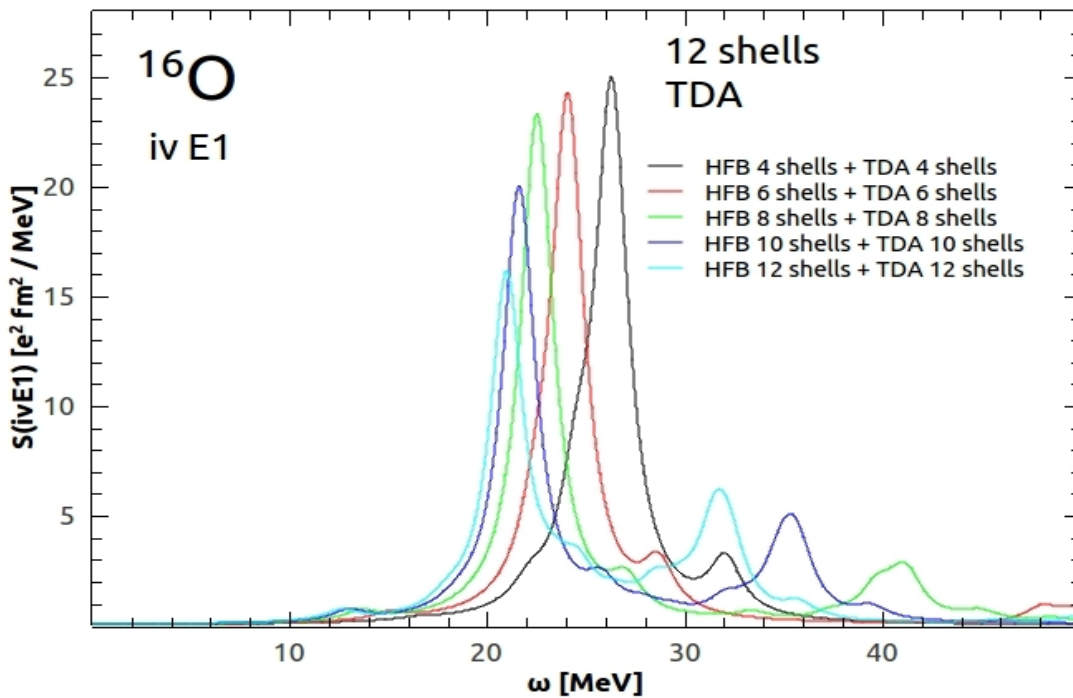
TDA – effect of configuration space



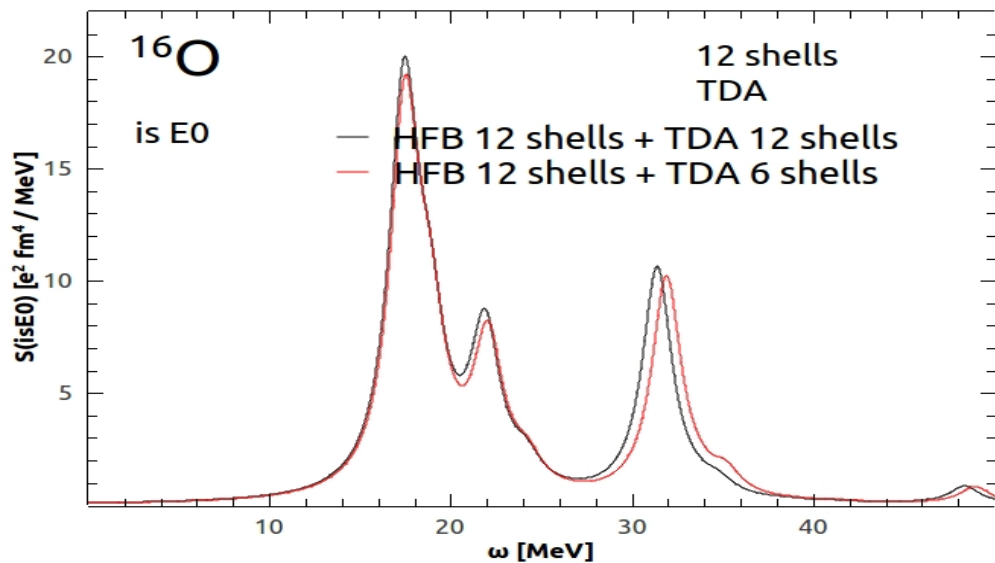
“natural cut“ of TDA configuration space

HFB calculated in maximal space cut-off only at the level of TDA

much quicker convergence due to cut-off only for TDA

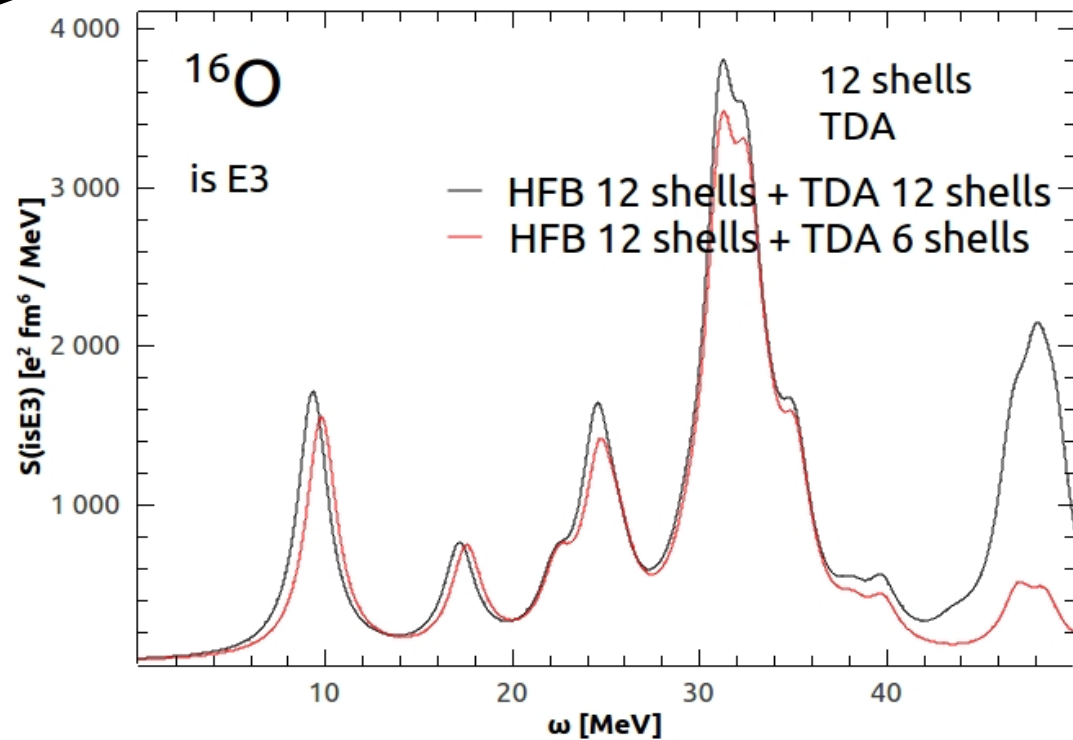
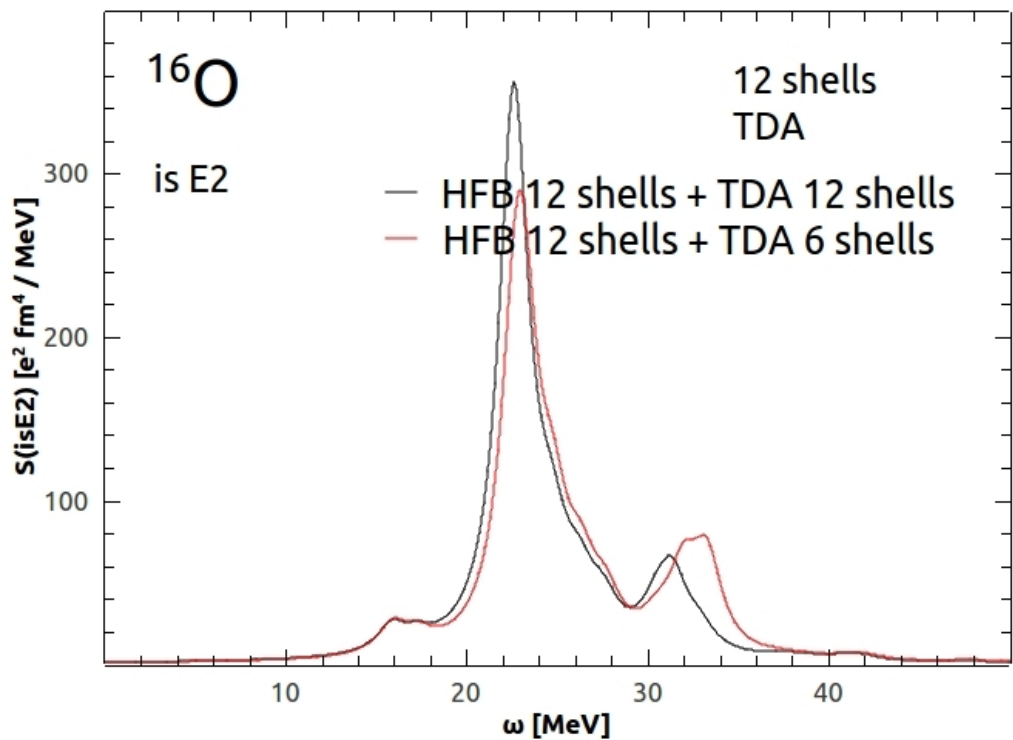


TDA – effect of configuration space

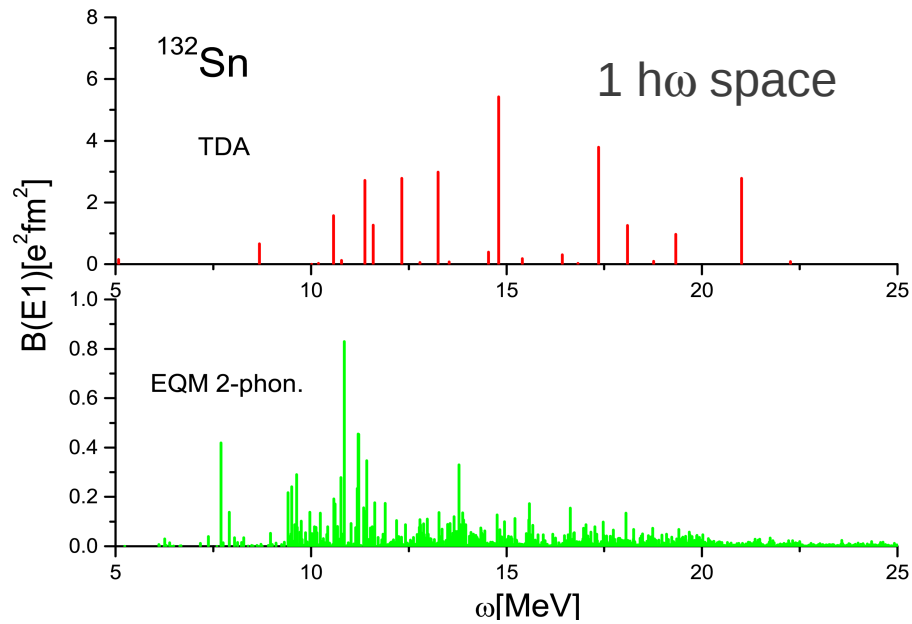


the effect on other multipolarities

E0, E2, E3



EMPM calculations in HF basis



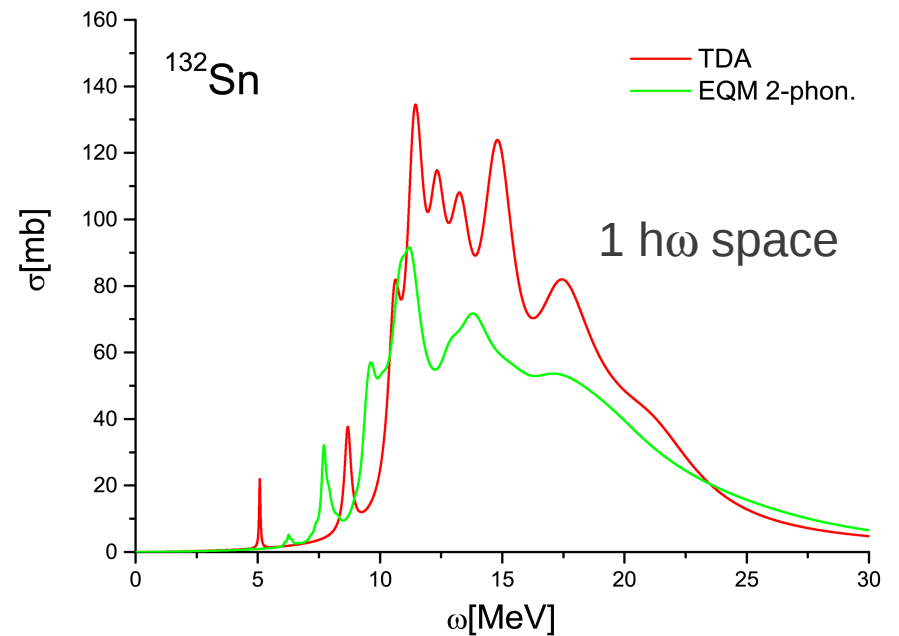
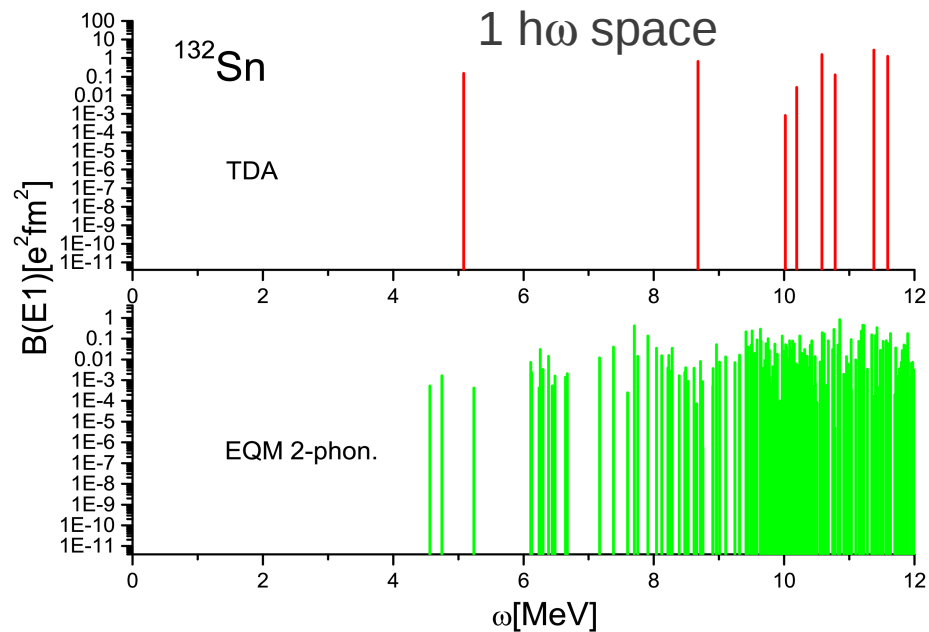
E1 calculation up to 2-phonons in ¹³²Sn...

Excitation energies measured from HF vacuum energy.

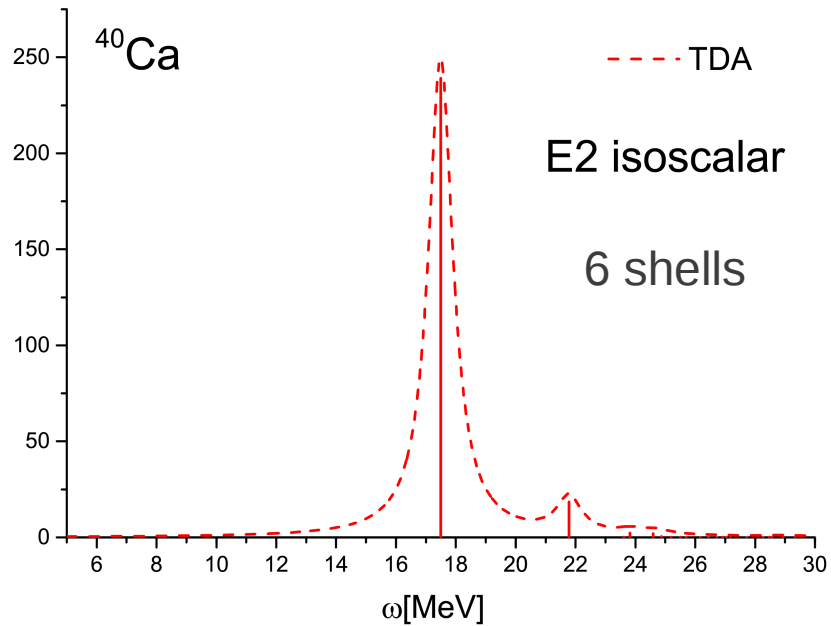
$$E_{\text{corr.}} = -5.8 \text{ MeV}$$

Systematical downshift of spectra expected at the 3-phonon and 4-phonon level

Visible multifragmentation of the spectrum...



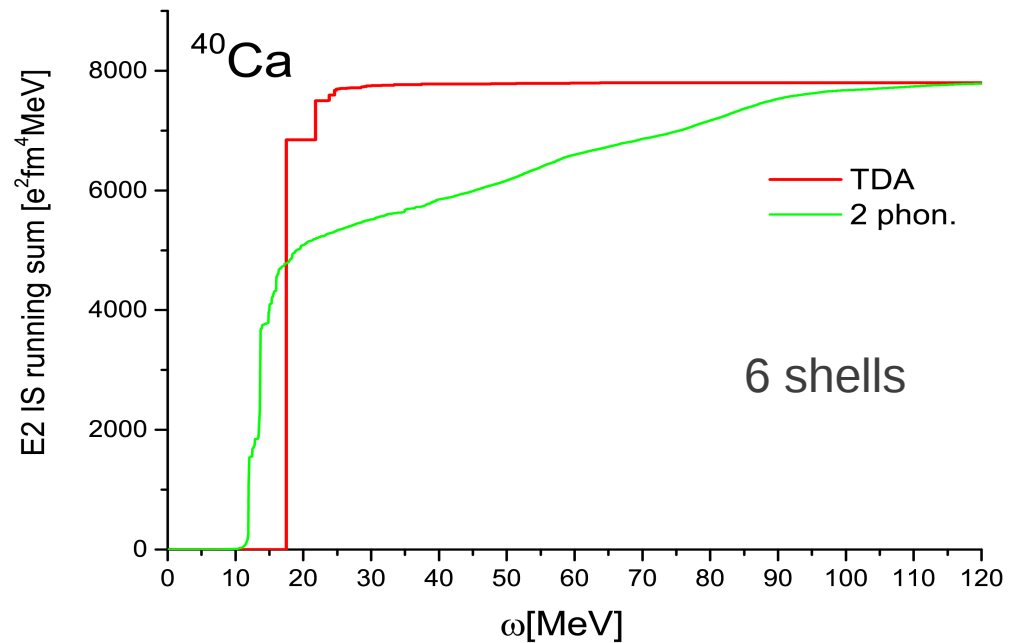
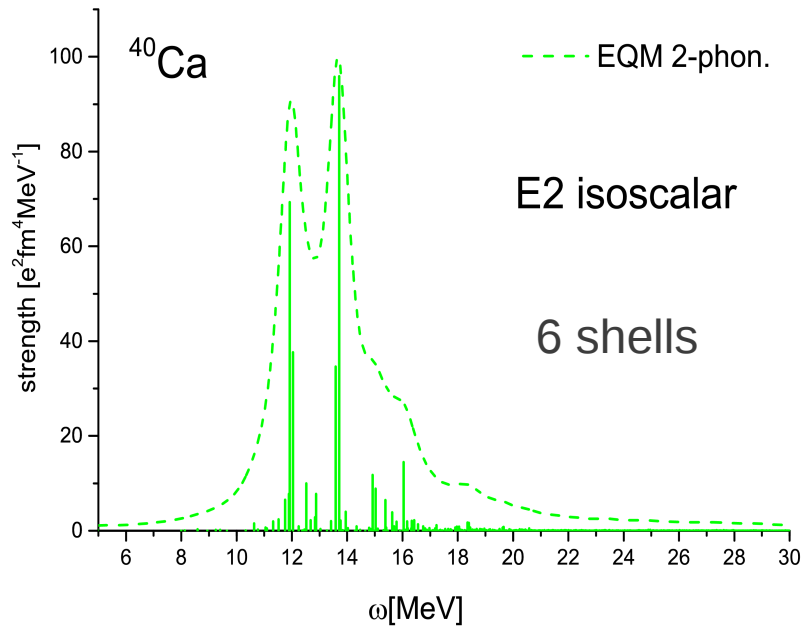
EMPM calculations in HF basis



E2 calculation up to 2-phonons in ⁴⁰Ca...
Excitation energies measured from HF
vacuum energy.

$$E_{\text{corr.}} = -43.9 \text{ MeV}$$

Multifragmentation effect – visible also
on the running sum



Plans for the near future:

EMPM calculations in self-consistent basis for isotopes of calcium, tin, lead on a way

Quasiparticle formulation of the EMPM (application of the model to semi-magic and neutron rich nuclei)

Study of pygmy resonances, M1 transitions (quenching effect, occurrence of spin-M1 for spin-saturated nucleus ^{40}Ca), multifragmentation of the giant resonances

The study of particle-phonon or particle-multiphonon coupling (description of the odd systems)

