AARHUS UNIVERSITET

JULY 6 2015

MAGNETISM AND DYNAMICS IN STRONGLY INTERACTING ONE-DIMENSIONAL SYSTEMS

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DEPARTMENT OF PHYSICS AND ASTRONOMY

595. WE-Heraeus-Seminar on
Cold atoms meets Quantum Field Theory
July 6th-9th 2015, PHB, Bad Honnef, Germany

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23rd European Few-Body Conference

Aarhus, Denmark, 8-12 August 2016

- Sub-atomic systems – including light nuclei, hadrons, few-nucleon physics and nuclear astrophysics
- Atomic and molecular systems, cold atoms and ions.
- Few-body methods
- Few-body physics in many-body systems
- New topics in few-body physics

Website: **conferences.au.dk/efb23**

Registration opening soon!

We look forward to seeing you in Aarhus in August 2016!

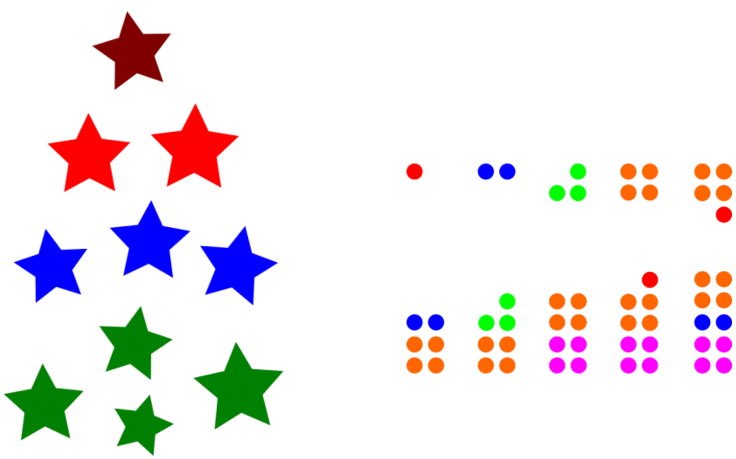
A 'SIMPLE' QUESTION

How many is 'many' really?

AN ANSWER?

Subitizing

Kaufman, E.L., Lord, M.W., Reese, T.W., & Volkmann, J. (1949). "The discrimination of visual number". *American Journal of Psychology* (The American Journal of Psychology) **62** (4): 498–525

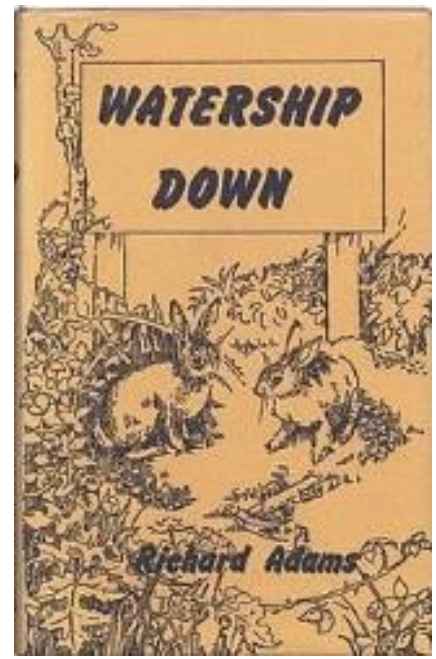


Counting or subitizing?

"Subitizing" by Nevit Dilmen, Wikipedia
(Thanks to Jose D'Incao for pointiing this out)

Source: Wikipedia

Richard Adams novel 1972



Rabbits cannot count beyond four so five is like a thousand ('infinity')

FIVE IS DIFFERENT!

Nathan L. Harshman

One-Dimensional Traps, Two-Body Interactions, Few-Body Symmetries: I. One, Two, and Three Particles

arXiv:1501.00215

One-Dimensional Traps, Two-Body Interactions, Few-Body Symmetries: II. N particles

arXiv:1505.00659

For more than four particles, the general case requires a solution of a degree five polynomial equation – no root formulas!

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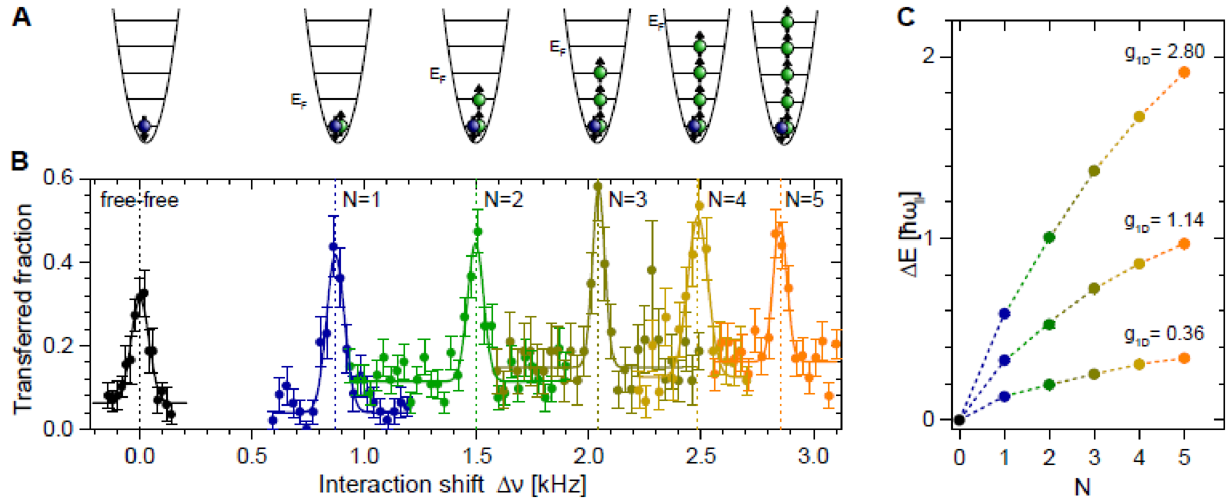
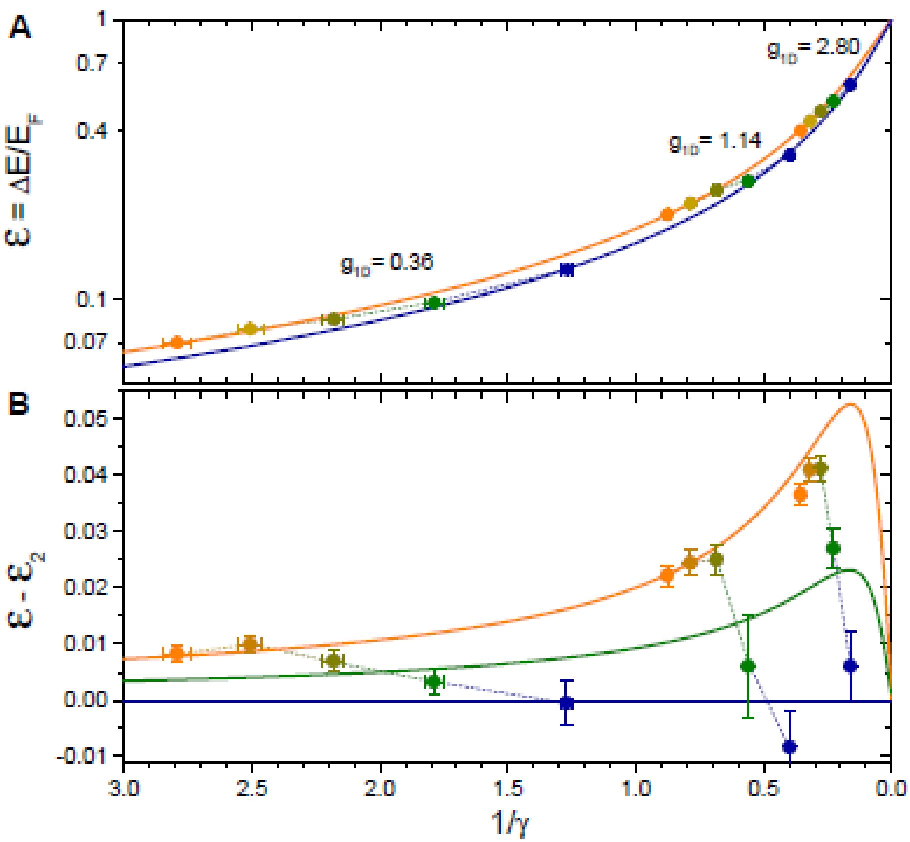
Selim Jochim experiments in Heidelberg.





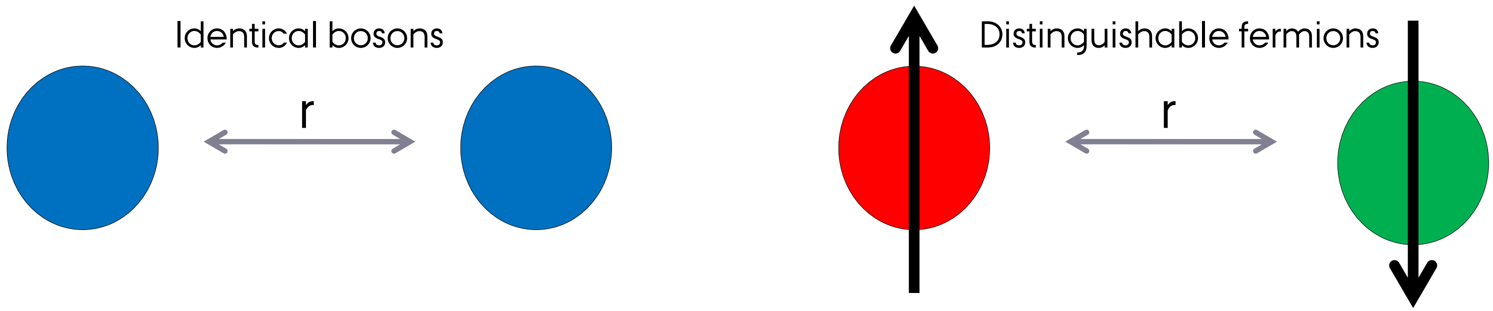
From Few to Many: Observing the Formation of a Fermi Sea One Atom at a Time

A. N. Wenz *et al.*, Science **342**, 457 (2013)

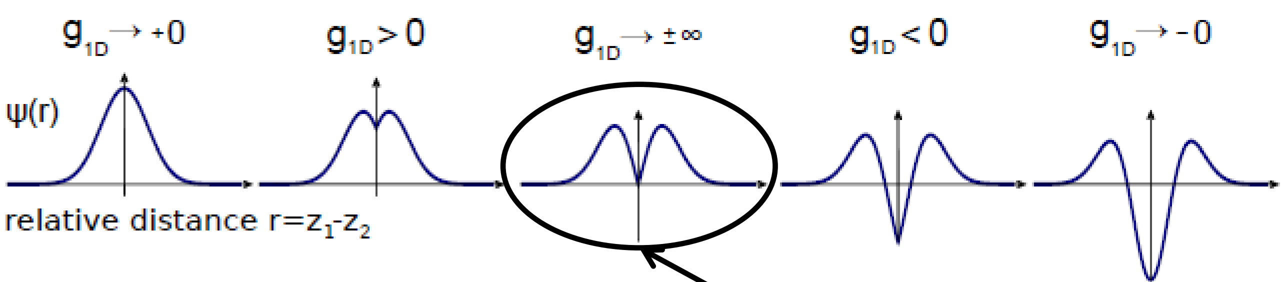
Green solid line from
S.E. Gharashi, K.M. Daily, and D. Blume, Phys. Rev. A **86**, 042702 (2012).

Orange 'many-body' line
J.B. McGuire, J. Math. Phys. **6**, 432 (1965).
G.E. Astrakharchik and I. Brouzos, Phys. Rev. A **88**, 021602 (2013).

A ONE DIMENSIONAL WORLD



Relative wave function



Interaction

$$g_{1D}\,\delta(r)$$

Source: G. Zürn, thesis

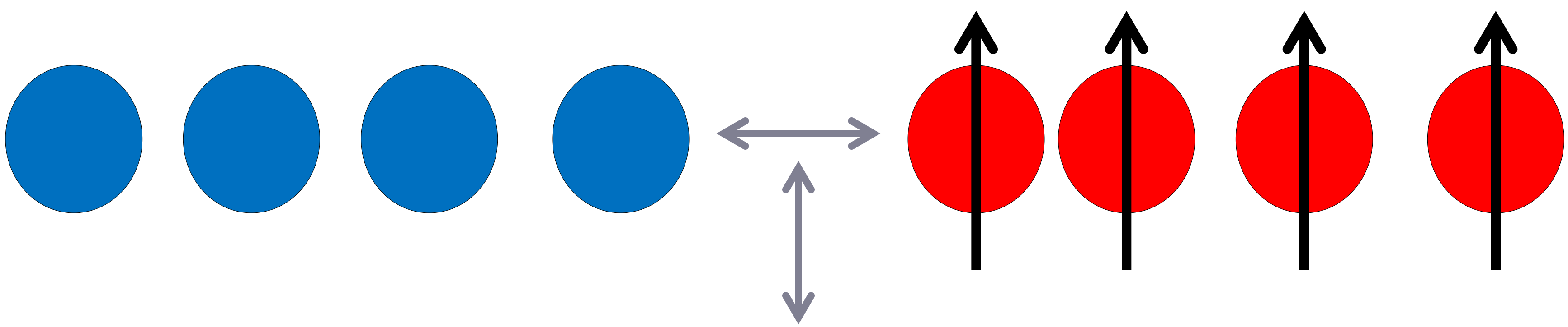
Strong interactions -> Impenetrability!

STRONGLY INTERACTING BOSONS

$|g_{1D}| \rightarrow \infty$ limit

Tonks (1936)-Girardeau (1960) gas
of impenetrable bosons

Mapping identical bosons to spin-polarized fermions. Girardeau (1960).



$$\Psi_B(z_1, \ldots, z_N) = |\Psi_F(z_1, \ldots, z_N)|$$

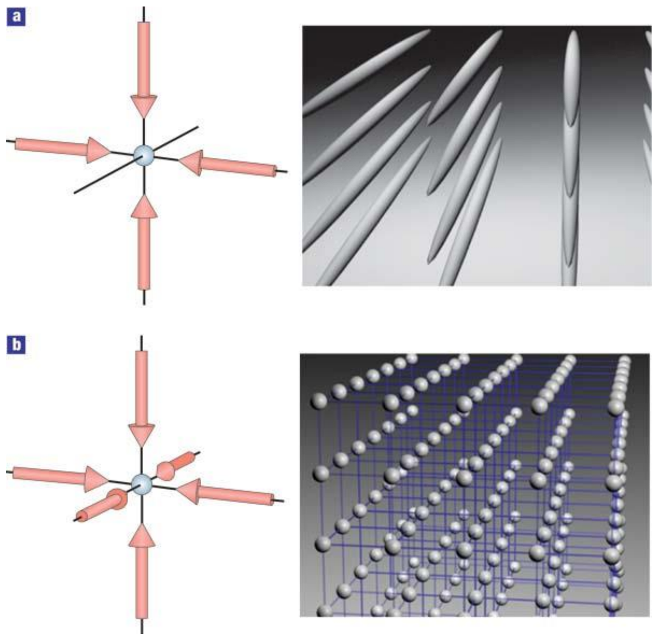
Impenetrable bosons

Antisymmetrized fermions

Lieb-Liniger (1963) used Bethe ansatz to
solve N boson problem for any g>0

EXPERIMENTAL REALIZATION

Optical lattices



I. Bloch, Nature Physics **1**, 23 (2005)

Confinement-induced resonances

Maxim Olshanii
Phys. Rev. Lett. **81**, 938 (1998)

$$g_{1D} = \frac{2\hbar^2 a_{3D}}{m a_\perp^2} \frac{1}{1 - C a_{3D}/a_\perp}$$

Divergent at specific point depending on lattice and 3D Feshbach resonance

EXPERIMENTAL REALIZATION

Tonks–Girardeau gas of ultracold atoms in an optical lattice

Belén Paredes[1], Artur Widera[1,2,3], Valentin Murg[1], Olaf Mandel[1,2,3], Simon Fölling[1,2,3], Ignacio Cirac[1], Gora V. Shlyapnikov[4], Theodor W. Hänsch[1,2] & Immanuel Bloch[1,2,3]

Nature **429**, 277 (2004)

Observation of a One-Dimensional Tonks-Girardeau Gas

Toshiya Kinoshita, Trevor Wenger, David S. Weiss*

Science **305**, 1125 (2004)

Experimentally produced and probed the Tonks-Girardeau gas on the repulsive side g>0

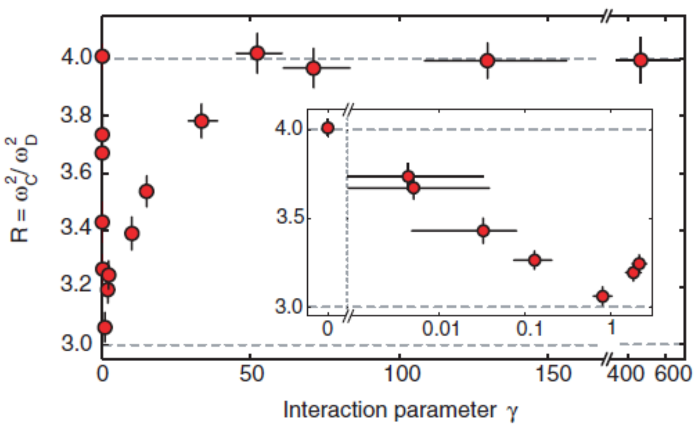
EXPERIMENTAL REALIZATION

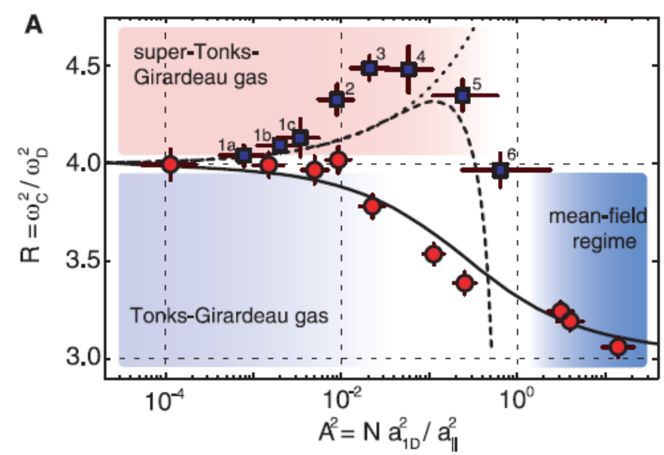
Realization of an Excited, Strongly Correlated Quantum Gas Phase

Elmar Haller,[1] Mattias Gustavsson,[1] Manfred J. Mark,[1] Johann G. Danzl,[1] Russell Hart,[1] Guido Pupillo,[2,3] Hanns-Christoph Nägerl[1]*

Science **325**, 1224 (2009)

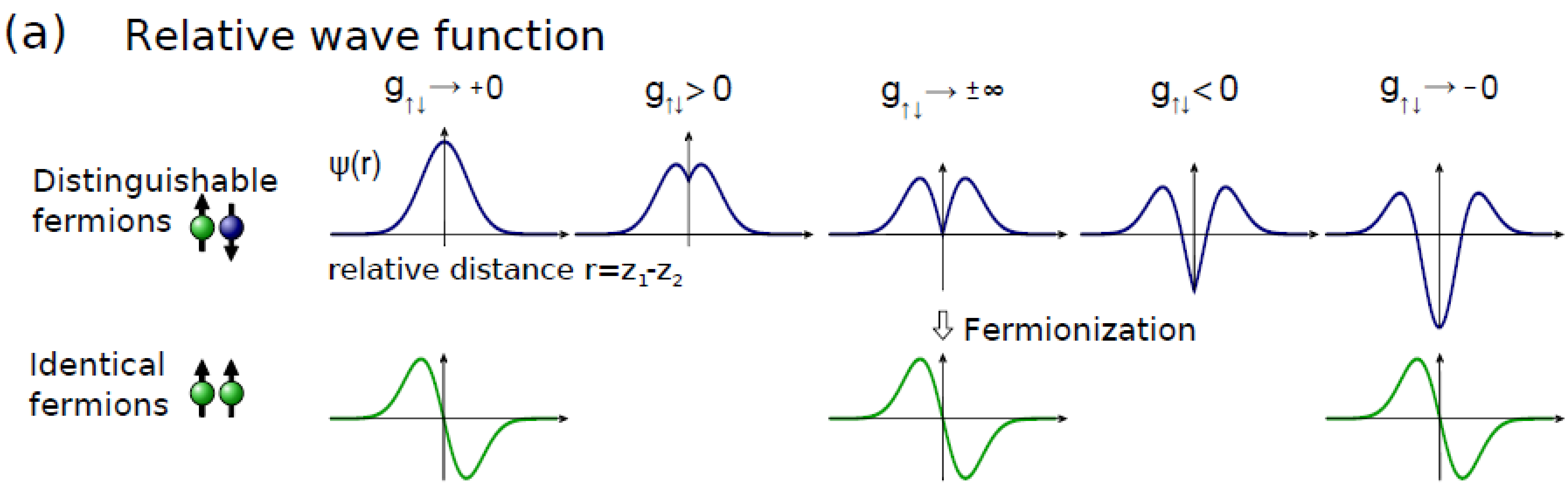
Study of the crossover from g>0 to g<0 in the strongly-interacting regime.





1D FERMIONS – A FRONTIER

Two kinds of relative motion for two-body states!

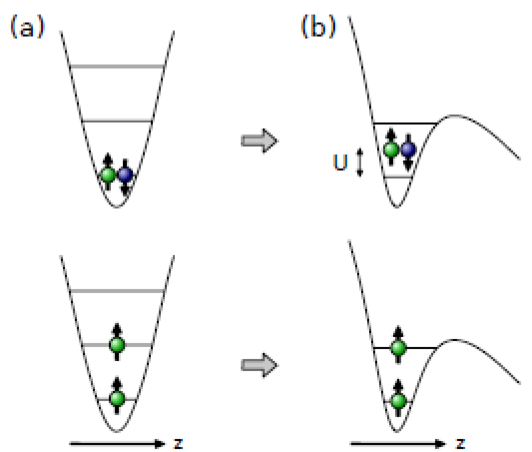


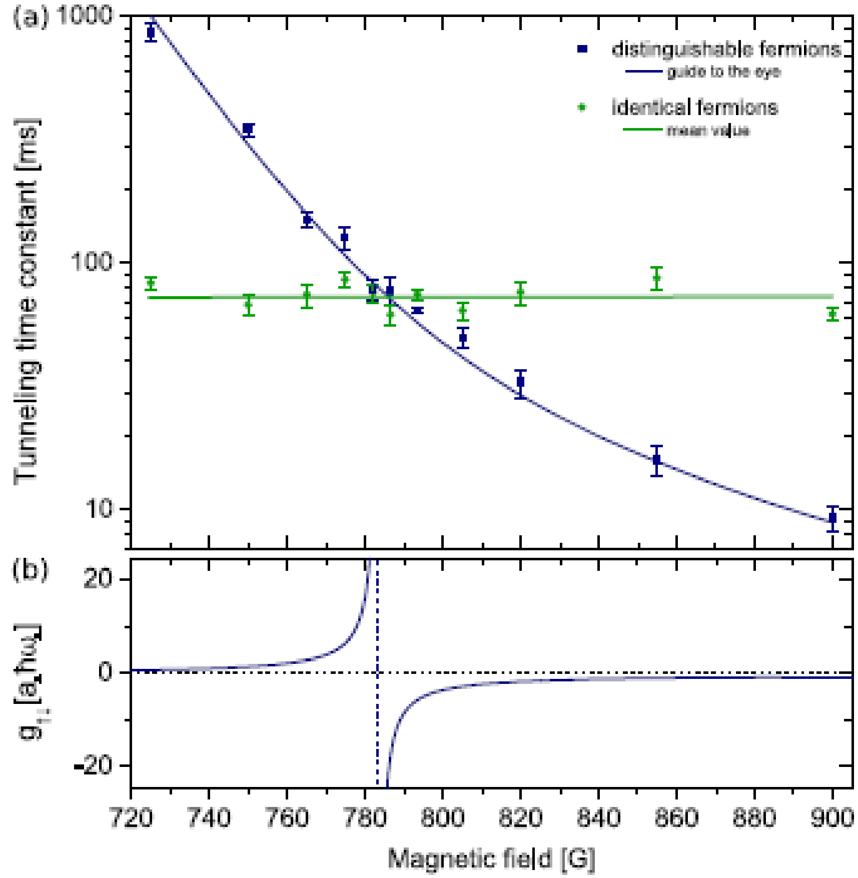
Source: G. Zürn, thesis

Fermionization of two fermions in a 1D harmonic trap:
G. Zürn *et al.*, Phys. Rev. Lett. **108**, 075303 (2012).

EXPERIMENTAL REALIZATION

Two-body tunneling experiments

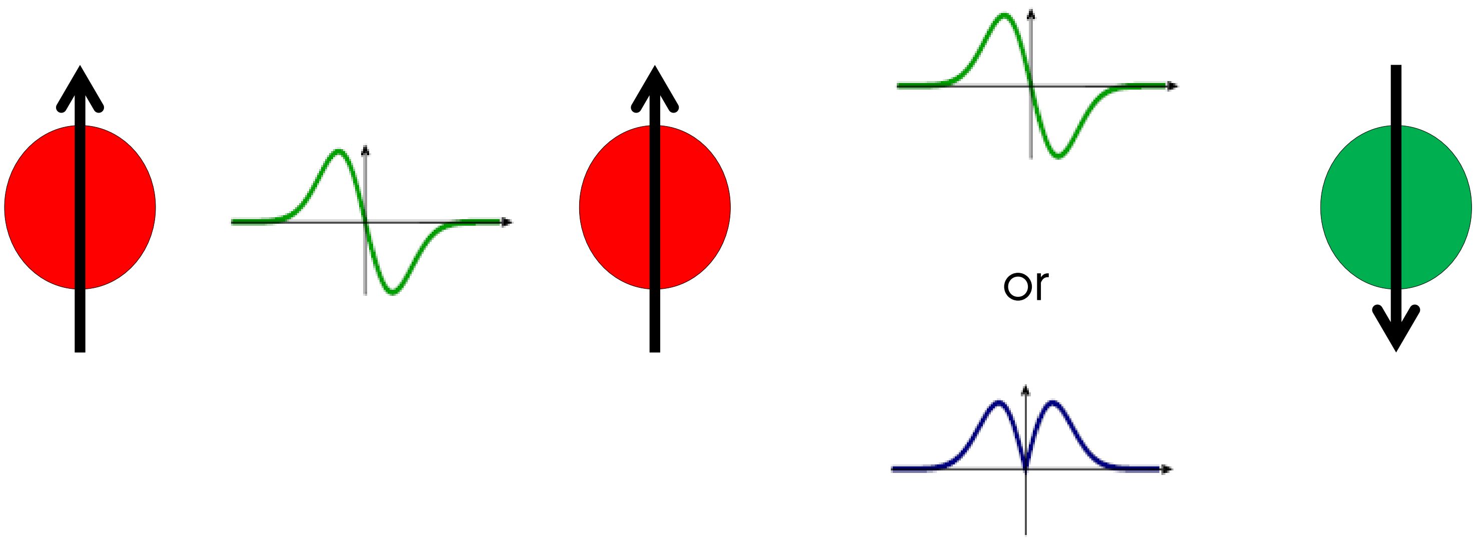




Fermionization of two fermions in a 1D harmonic trap:
G. Zürn *et al.*, Phys. Rev. Lett. **108**, 075303 (2012).

THREE FERMIONS

Relative wave functions. What should we take?



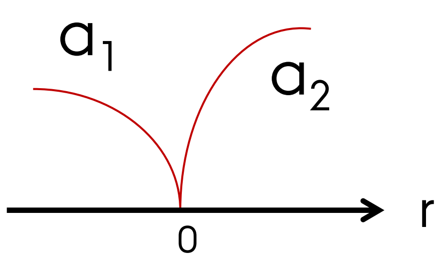
???

Conjecture: Use the symmetric choice for non-identical pairs for any N-body system

M.D. Girardeau, Phys. Rev. A **82**, 011607(R) (2010).

THREE FERMIONS

Let's keep an open mind!



Non-identical relative wave function

Two strict conditions:

1) In limit g_{1D}->∞, relative wave functions have not vanish at zero for identical and non-identical pairs!

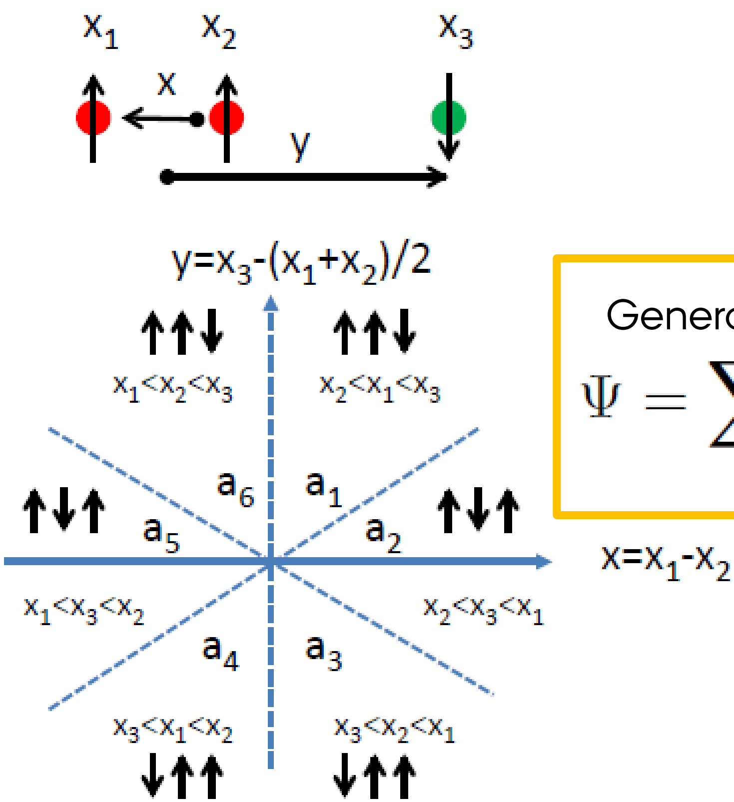
2) Identical fermions must have odd relative wave functions!

IDEA: Keep a_1 and a_2 as free parameters and do a variation!

A.G. Volosniev *et al.*, Nature Comm. **5**, 5300 (2014)

THREE FERMIONS - SOLUTION

Split space in patches

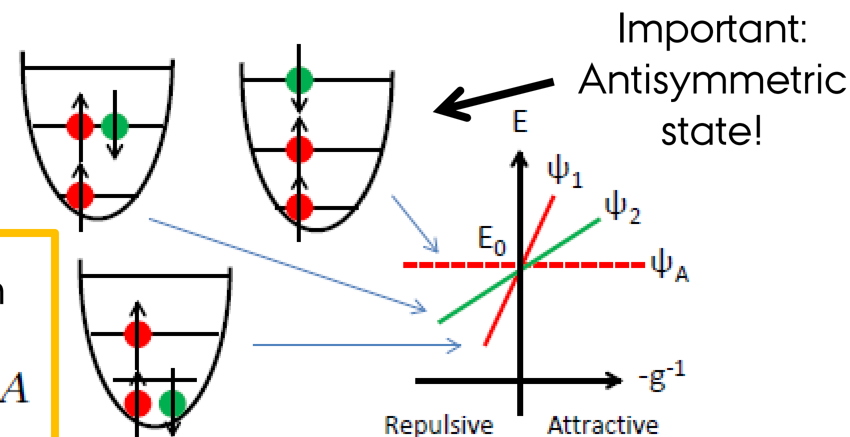


Pauli and parity reduces problem to a_1, a_2, and a_3.

General solution

$$\Psi = \sum_i a_i \Psi_A$$

Spectrum on resonance



Optimize derivative!

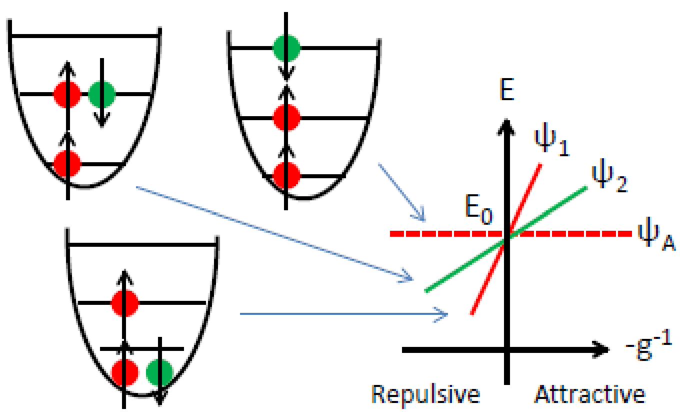
$$K = -\frac{\partial E}{\partial g^{-1}} = g^2 \frac{\sum_{ij} \int \prod_{k=1}^{N} dx_k |\Psi|^2 \delta(x_i - x_j)}{\langle \Psi | \Psi \rangle}$$

THREE FERMIONS - SOLUTION

$$K = \frac{27}{8\sqrt{2\pi}} \frac{(a_1 - a_2)^2 + (a_2 - a_3)^2}{a_1^2 + a_2^2 + a_3^2}$$

Extremizing solutions are:

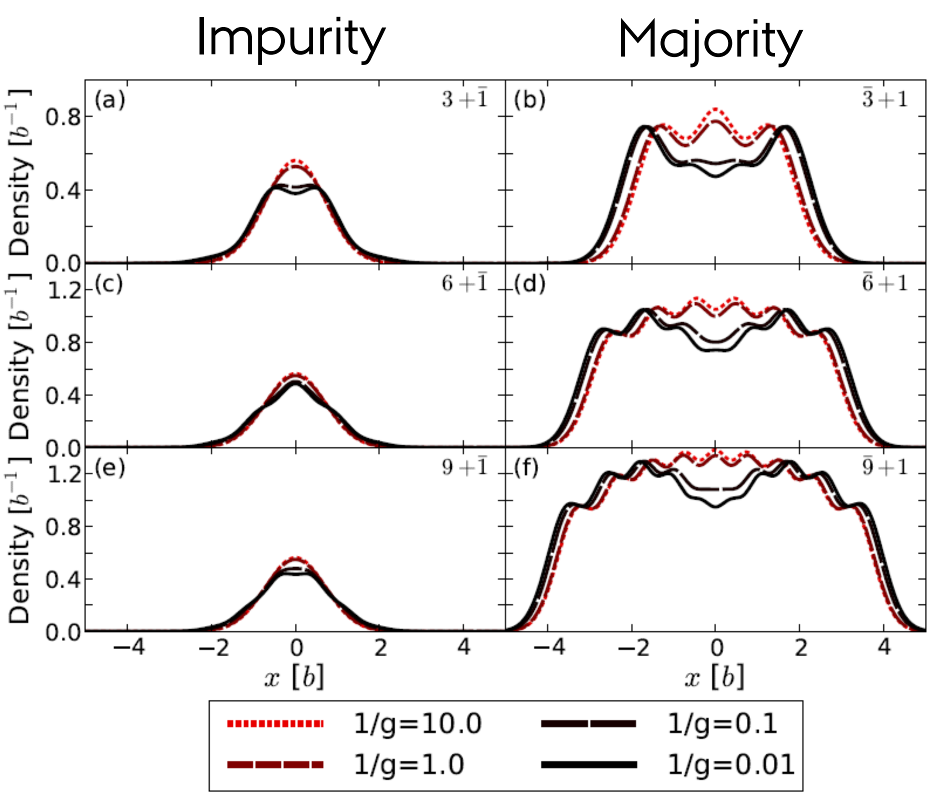
$a_1 = a_2 = a_3$	Non-interacting state
$a_1 = a_3$ and $a_2 = 0$	Excited state, even parity
$2a_1 = 2a_3 = -a_2$	Ground state, odd parity



IMPORTANT: Coefficients are generally NOT the same!

A.G. Volosniev *et al.*, Nature Comm. **5**, 5300 (2014).

TRAPPED 'POLARONS' IN 1D



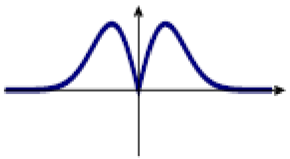
'Precursor' of magnetic structure! Phase-separation of spin up and spin down for strong interactions.

E.J. Lindgren *et al.*, New J. Phys. **16**, 063003 (2014)
S.E. Gharashi and D. Blume, Phys. Rev. Lett. **111**, 045302 (2013)
J. Levinsen *et al.*, arXiv:1408.7096
F. Deuretzbacher *et al.*, Phys. Rev. A **90**, 013611 (2014)

FERMIONIZATION OF FERMIONS

It is different from identical bosons and spin-polarized fermions!

The 'democratic' solution or trivial Bose-Fermi mapping uses:



between all non-identical pairs.

In the 2+1 case it is NOT a relevant eigenstate but rather a linear combination!

$$\Psi_{BF} = (8^{1/2}\Psi_{gs} + \Psi_{non})/3$$

BUT can we tell the difference in experiments?

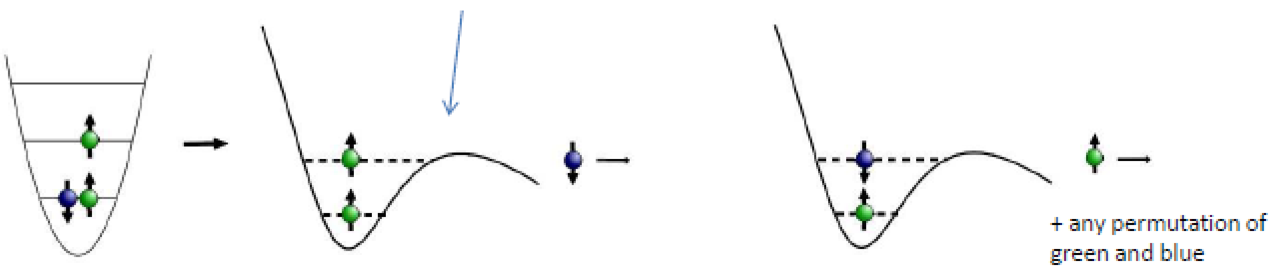
S.E. Gharashi and D. Blume, Phys. Rev. Lett. **111**, 045302 (2013)
A.G. Volosniev *et al.*, Nature Comm. **5**, 5300 (2014)

EXPERIMENTAL SIGNATURE

Do tunneling experiments!

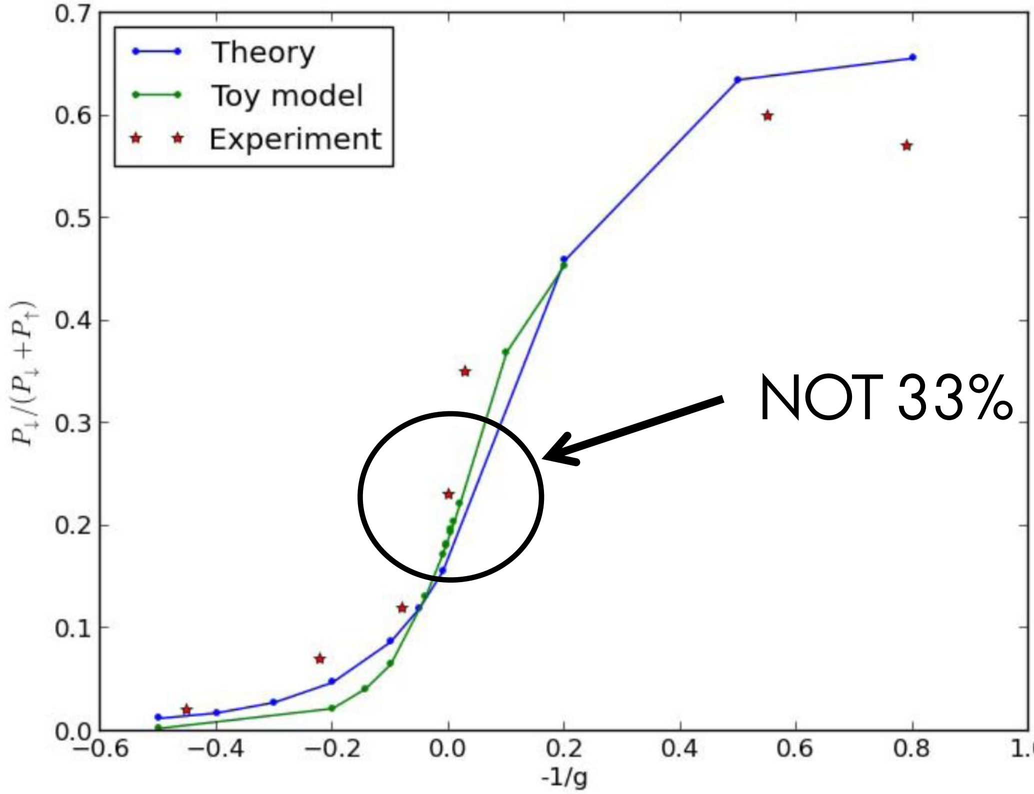
F. Serwane *et al.*, Science **332**, 336 (2011).

G. Zürn *et al.*, Phys. Rev. Lett. **108**, 075303 (2012).



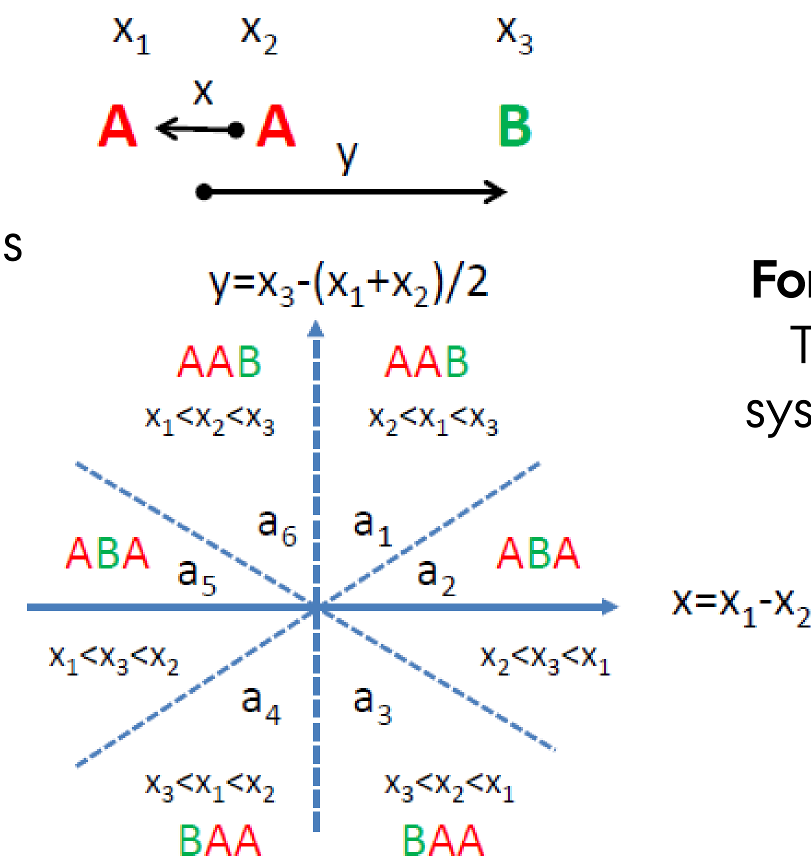
Source: G. Zürn

THEORY VS. EXPERIMENT



Theory by Lindgren and Volosniev.
Data from Jochim group (G. Zürn)

THREE TWO-COMPONENT BOSONS



Strong AB interactions

No AA interactions

No BB interactions

For more particles:
Two ideal Bose systems interacting strongly!

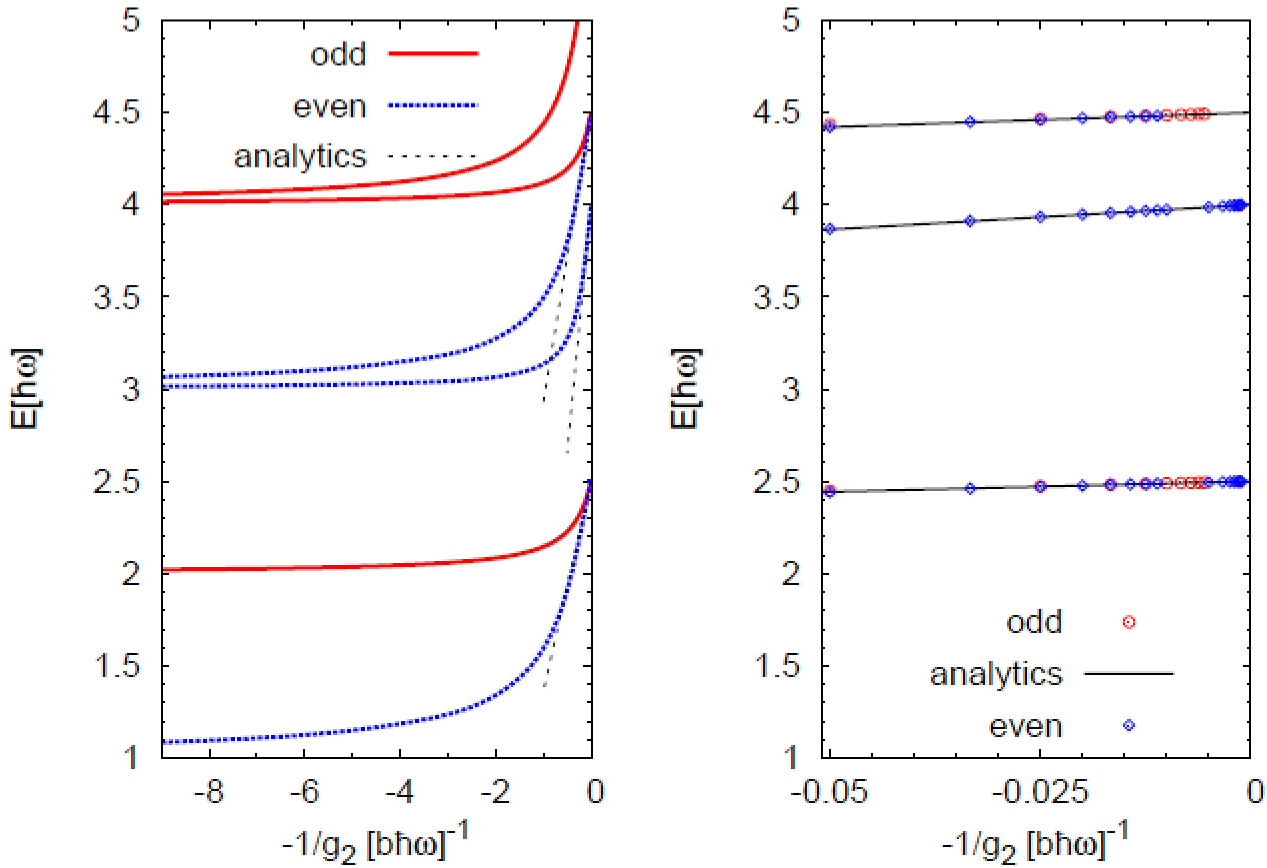
N.T. Zinner *et al.*, EPL **107**, 60003 (2014)

Fractional energies for strong interactions!

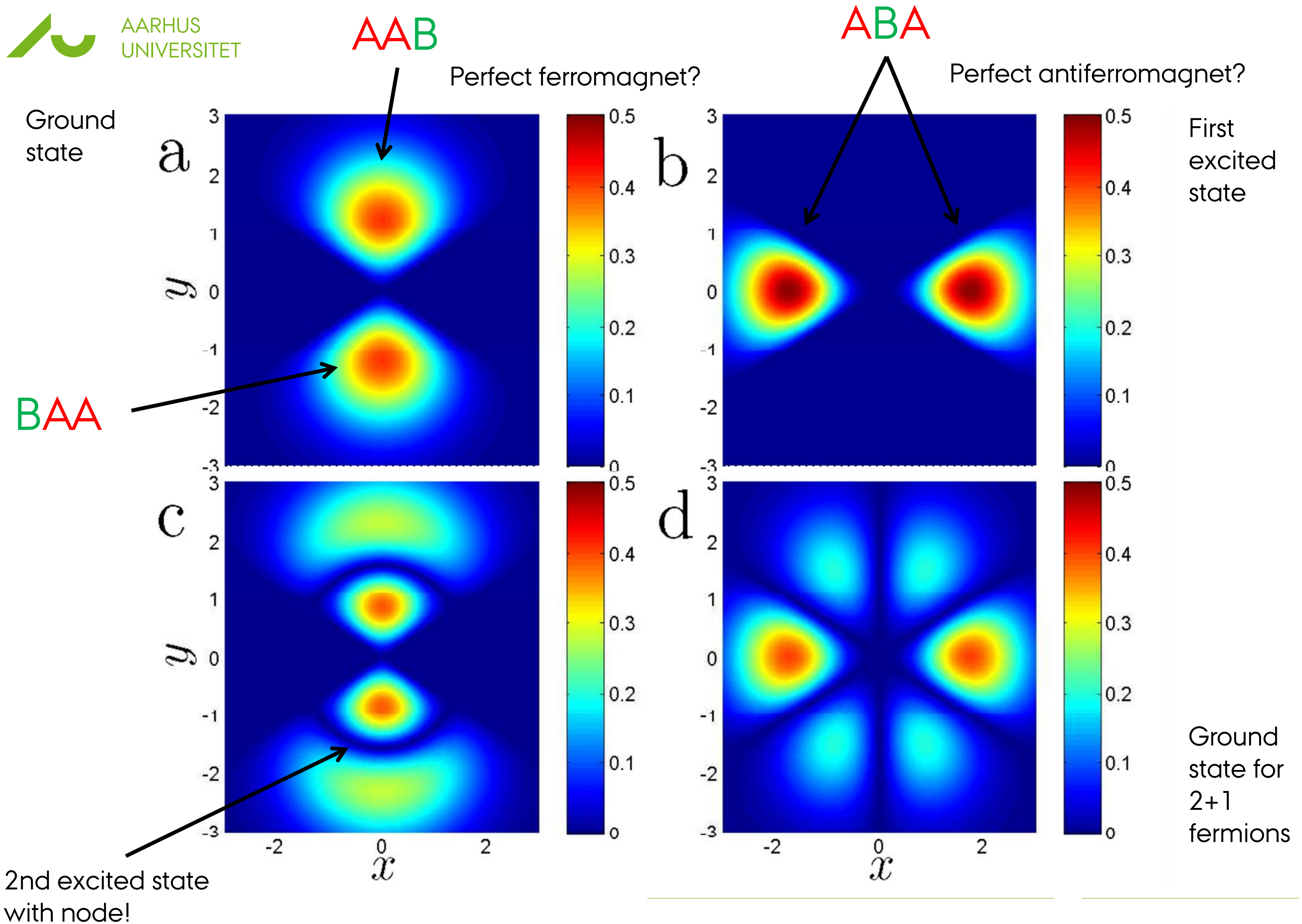
x_1 x_2 x_3

A A B

$$H=\sum_{i=1}^{3}\frac{p_i^2}{2m}+g_2\delta(x_1-x_3)+g_2\delta(x_2-x_3)$$



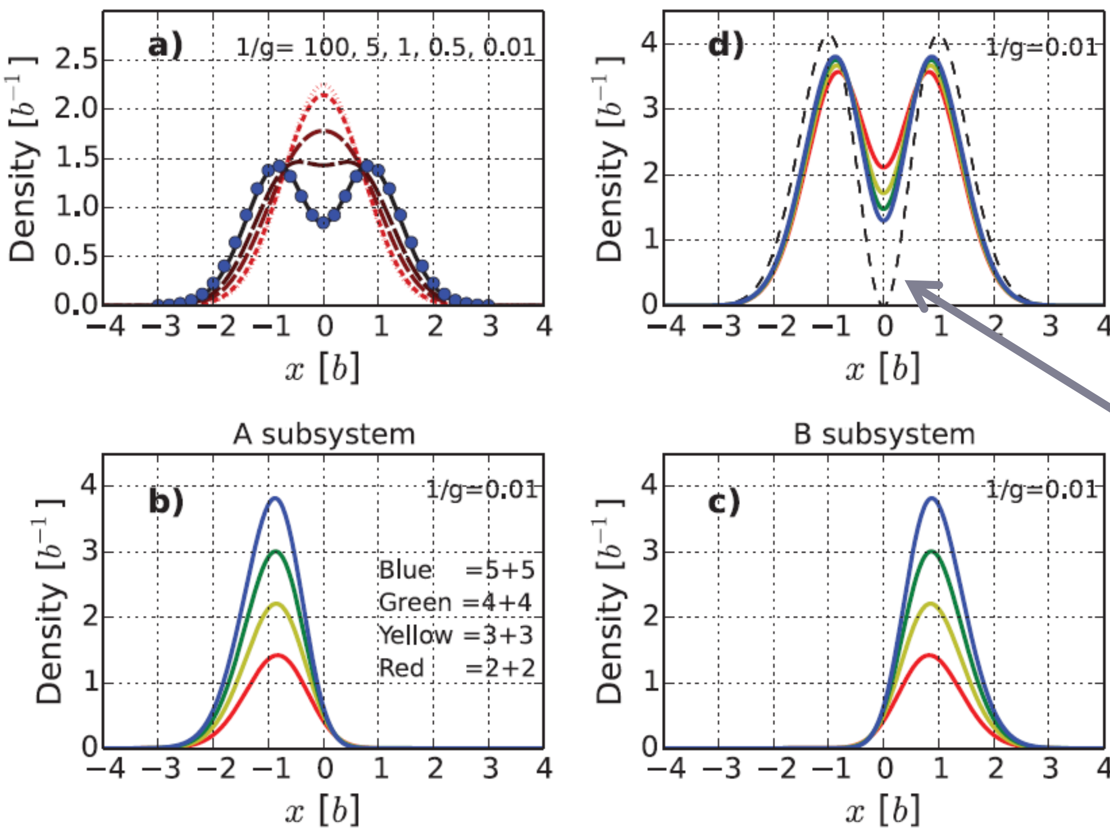
N.T. Zinner *et al.*, EPL **107**, 60003 (2014)



N.T. Zinner *et al.*, EPL **107**, 60003 (2014)

LARGER SYSTEMS

Energies are still fractional!



Ground state structure

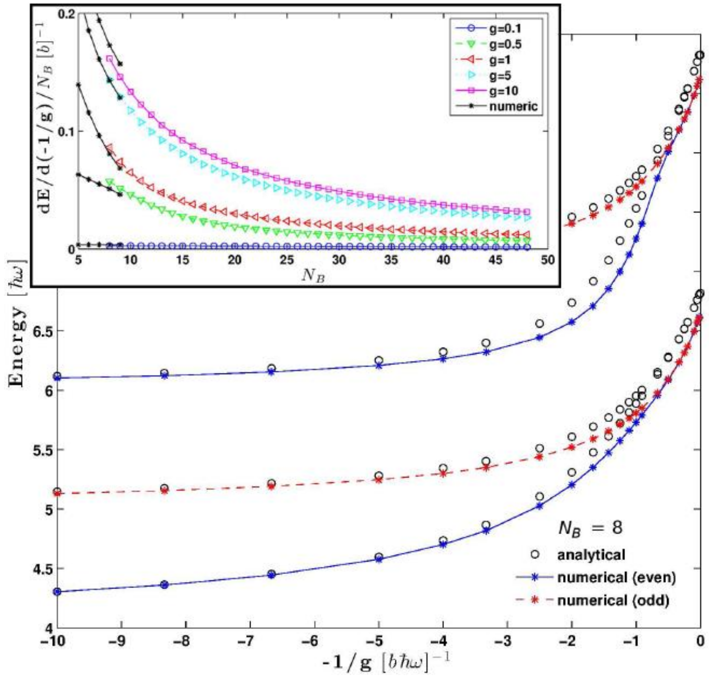
AAAAABBBBB+
BBBBBAAAAA

Perfect ferromagnetic ordering!

Many-body limit is approached quickly!

A.S. Dehkarghani *et al.*, Scientific Reports **5**, 10675 (2015).

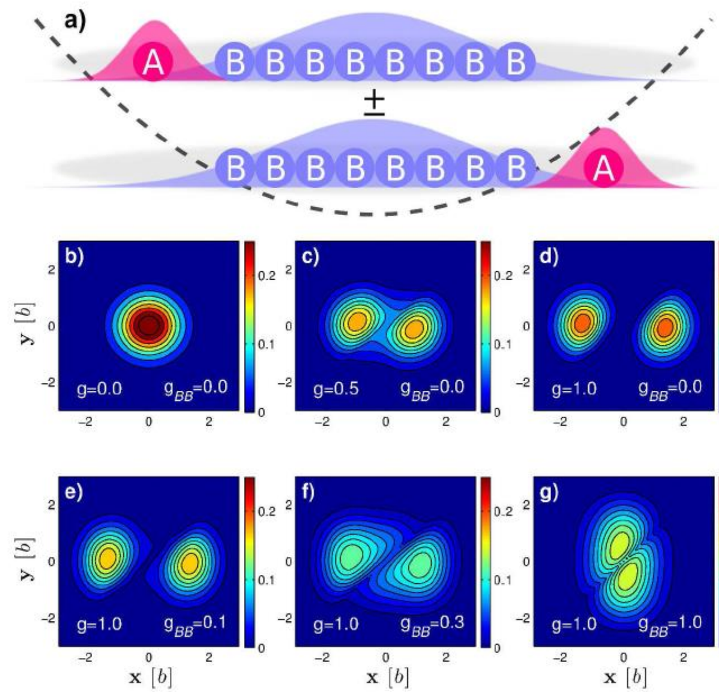
BOSE POLARONS



New semi-analytical approach to arbitrary particle numbers – a hyperspherical approach

Ground state structure

ABB….BB+BB….BBA



In the ground state, impurities will NEVER penetrate the majority component!

A.S. Dehkarghani *et al.*, arXiv:1503.03725

Poster by Amin S. Dehkharghani

SPIN MODELS

We can map strongly interacting two-component 1D systems in a trap to a spin model of XXZ type and do ENGINEERING!

Nearest-neighbor interactions are tunable via external trap!

$$H_s = \sum_{i=1}^{N-1} \left[J_i \left(\boldsymbol{S}(i) \cdot \boldsymbol{S}(i+1) \right) - \frac{2J_i}{\kappa} \left(S_z(i) S_z(i+1) \right) \right]$$

$$g_{\uparrow\uparrow} = g_{\downarrow\downarrow} = \kappa g_{\uparrow\downarrow}$$

Note: It is **not** a lattice index! It is a particle index.

TABLE I: Effective Heisenberg spin models for strongly interacting atoms in 1D traps.

Spin-$\frac{1}{2}$ model	Constituents	κ
XXZ	bosons	$0 < \kappa < \infty$
XXX	bosons or fermions	$\kappa \to \infty$
XX	bosons	$\kappa = 2$

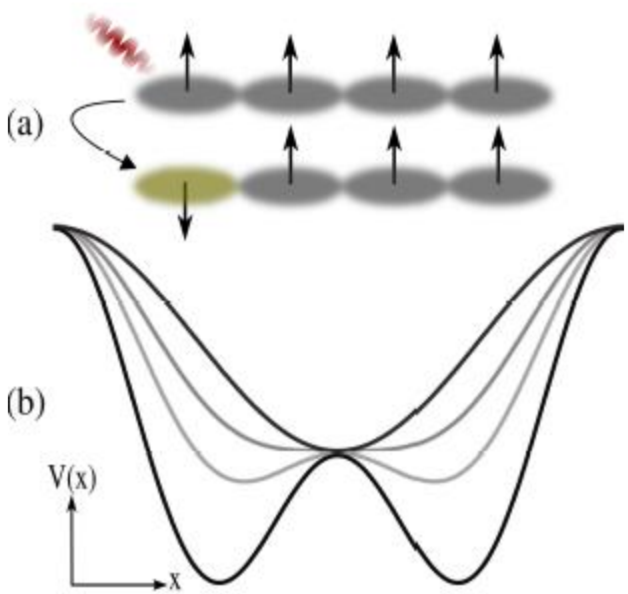
Confinement is taken into account **exactly**.

A.G. Volosniev *et al.*, Phys. Rev. A **91**, 023620 (2015)

See also F. Deuretzbacher *et al.*, and J. Levinsen *et al.*

STATE TRANSFER

Use trap to manipulate dynamics – example of quantum state transfer

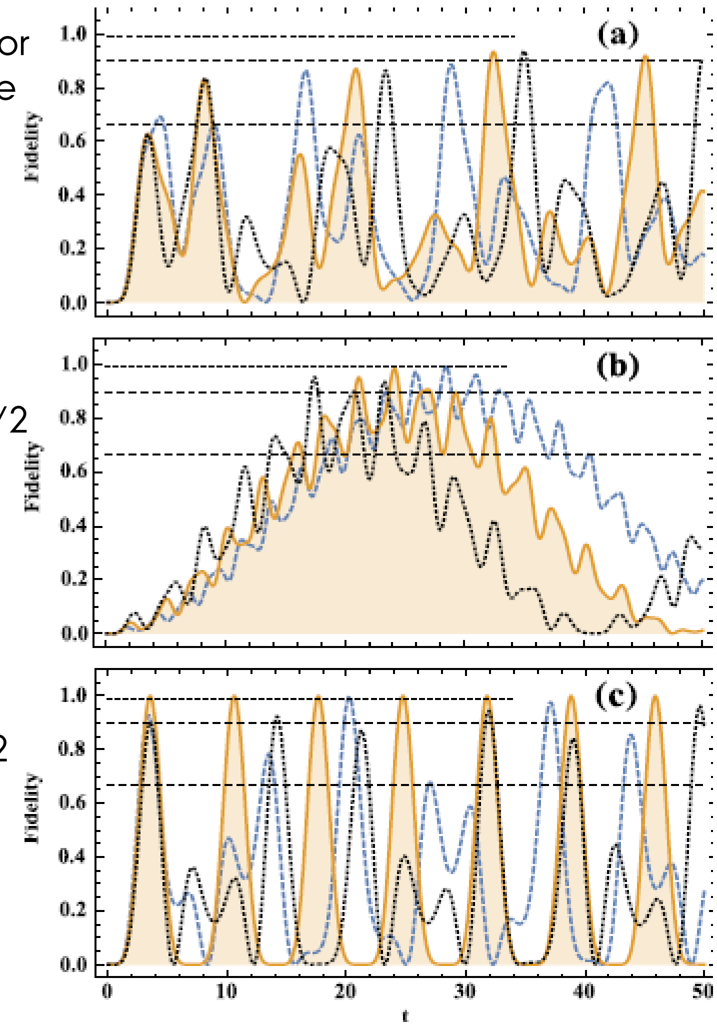


Fidelity of quantum state transfer

Fermions or hard-core bosons

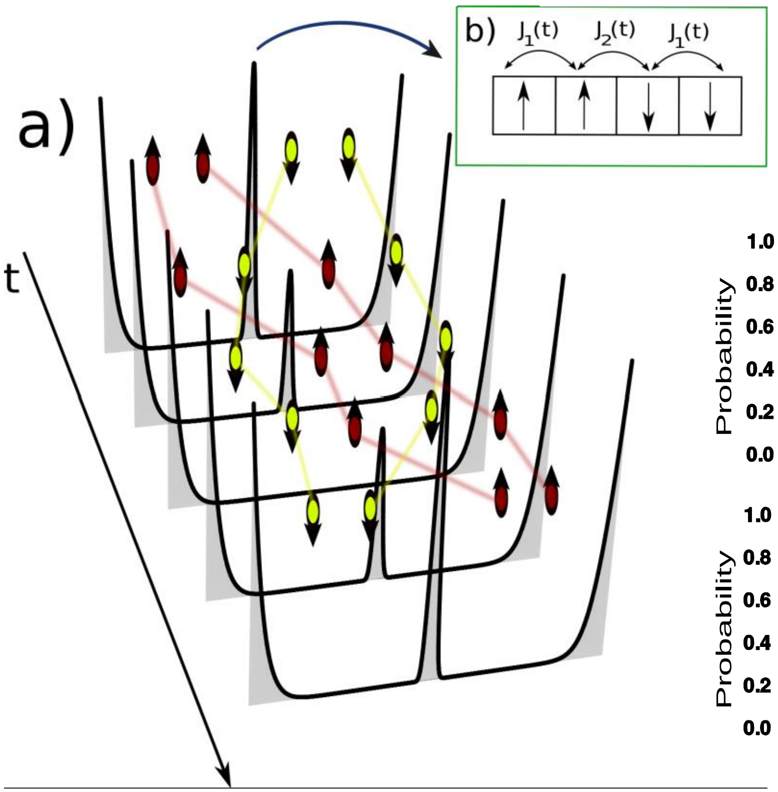
Bosons kappa=1/2

Bosons kappa=2



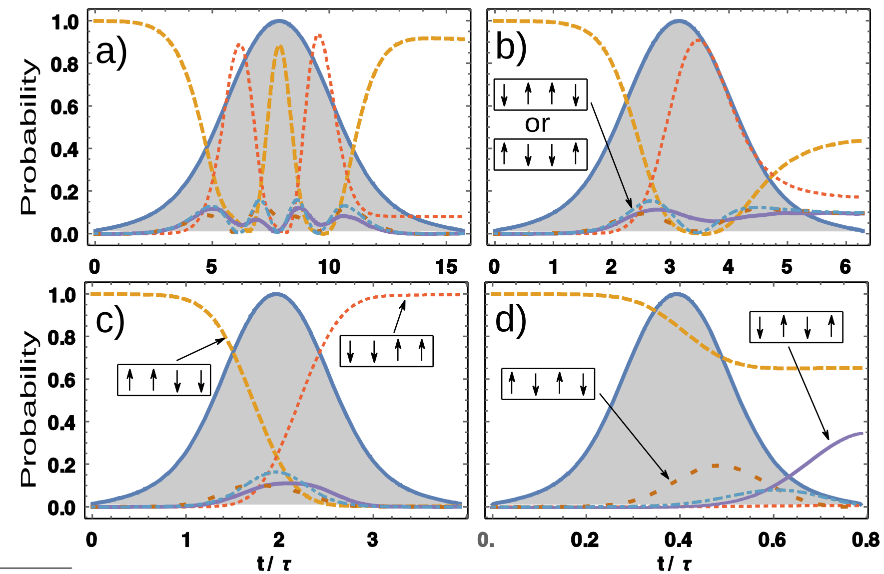
A.G. Volosniev*et al.*, Phys. Rev. A **91**, 023620 (2015)

TIME-DEPENDENT EXCHANGE



$$V(x,t) = \alpha f(t)\delta(x)$$

$$f(t) = (1 - \sqrt{\sin(At/\tau)})$$



A.G. Volosniev, H.-W. Hammer, NTZ, arXiv:1507.00186

SUMMARY

- Multi-component Bose and/or Fermi systems goes beyond Bose-Fermi mapping
- Magnetic correlations are accessible and we can engineer magnetic states, statically and dynamically!

ACKNOWLEDGEMENTS

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- Amin Dehkharghani, Oleksandr Marchukov, Dmitri Fedorov, Aksel Jensen (Aarhus)
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