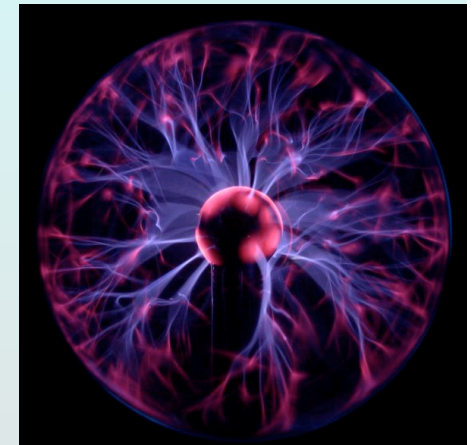


Measuring Properties of the Shortest Lived Plasma in the World



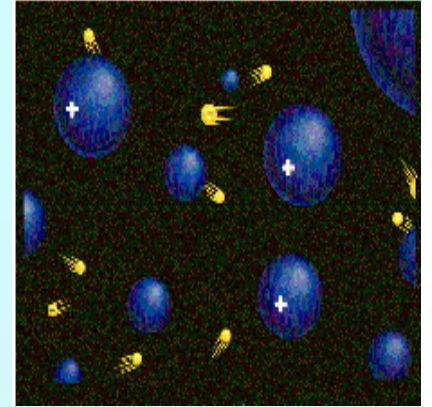
Barbara Jacak
Stony Brook

September 23,
2008

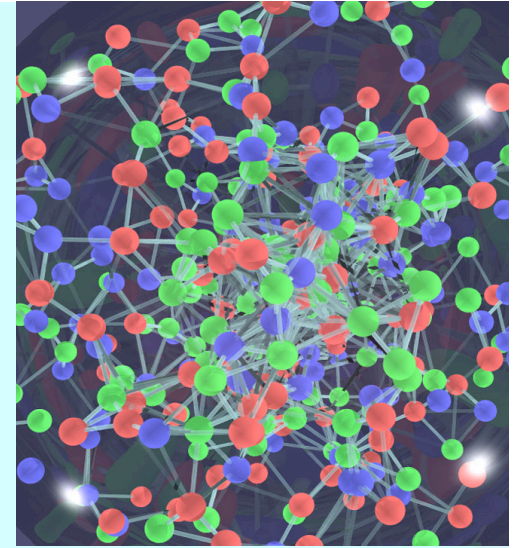


what is a plasma?

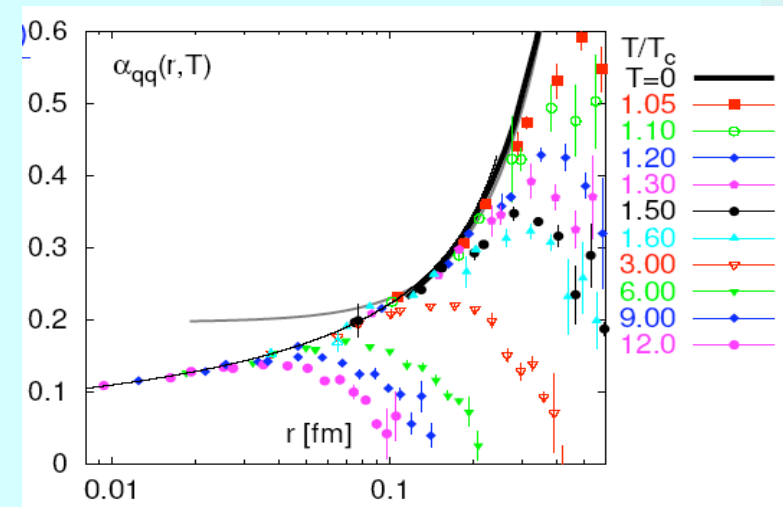
- 4th state of matter (after solid, liquid, gas)
- a plasma is:
 - ionized gas, macroscopically neutral
 - exhibits collective effects
- interactions among charges of multiple particles spreads charge into characteristic (Debye) length, λ_D
 - multiple particles inside this length
 - they screen each other
 - plasma size $> \lambda_D$
- “normal” plasmas are electromagnetic (e + ions)
 - quark-gluon plasma interacts via strong interaction



quark gluon plasma



- color forces rather than EM
exchanged particles: g instead of γ
gluons self-interacting
 \therefore theory is non-abelian
- “normal” plasmas
can vary ρ , T independently
strongly or weakly coupled
- QCD plasma
 T determines all properties
limits analogy to EM plasma
(heavy q mass sets new scale)
intrinsically strongly coupled for accessible T
 $\alpha_s \sim 0.35$; quarks & gluons NOT asymptotically free
...not your mother's plasma

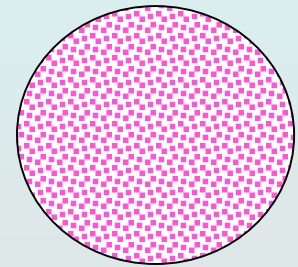


Debye screening in plasma

- Debye length: distance over which the influence of an individual charged particle is felt by the other particles in the plasma
- charged particles arrange themselves so as to effectively shield any electrostatic fields within a distance of order λ_D

- $\lambda_D = \left[\frac{\epsilon_0 kT}{n_e e^2} \right]^{1/2}$

n_e = number density
 e = charge



- Debye sphere = sphere with radius λ_D
- # of electrons inside Debye sphere is large

$$N_D = N/V_D = \rho V_D \quad V_D = \frac{4}{3} \pi \lambda_D^3$$

Debye screening in QCD: a tricky concept

- in leading order QCD (O. Philipsen, hep-ph/0010327)

$$V(r) \sim \frac{e^{-m_D^0 r}}{4\pi r}, \quad m_D^0 = \left(\frac{N}{3} + \frac{N_f}{6} \right)^{1/2} gT.$$

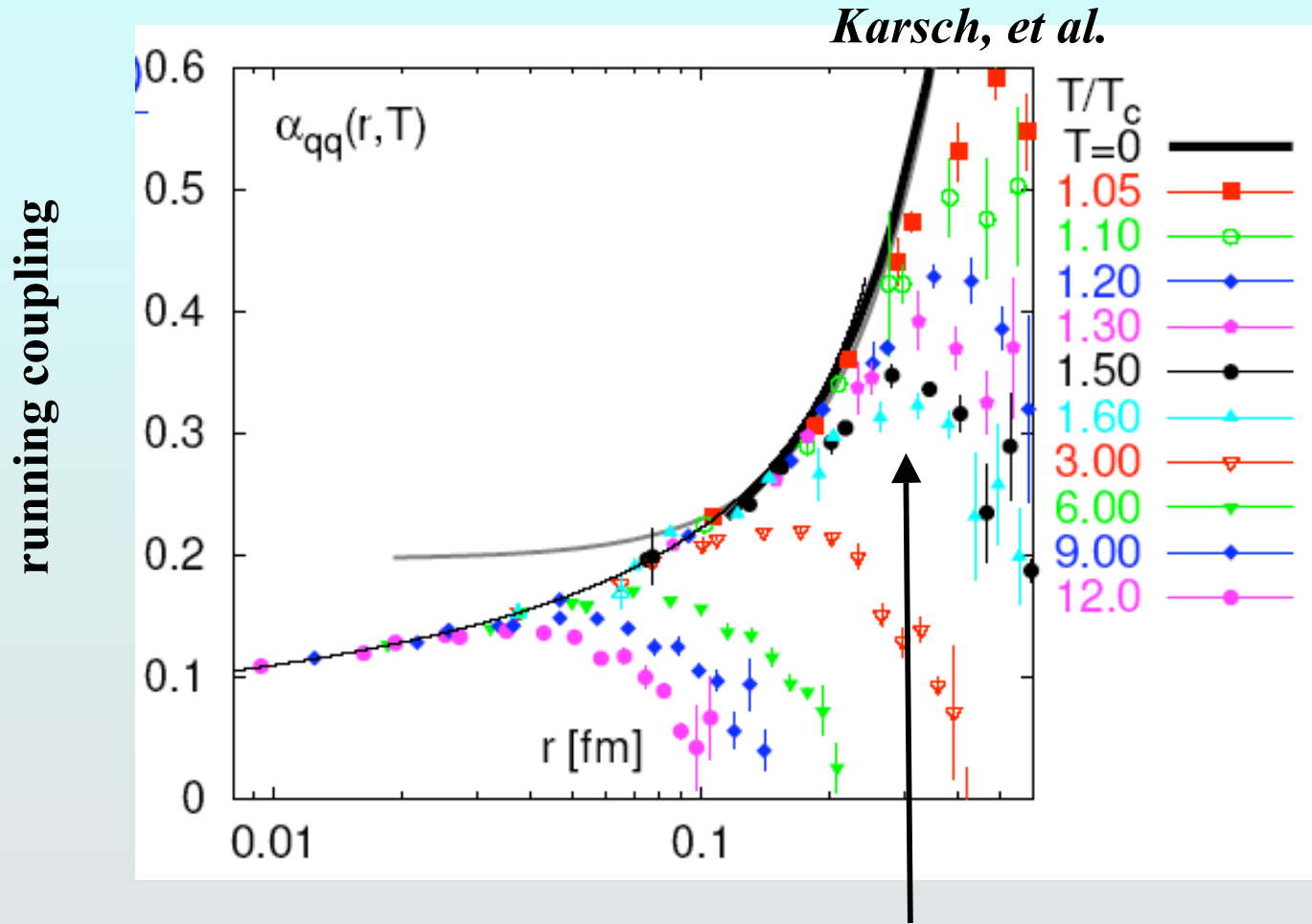
- However, at next-to-leading order the problem becomes non-perturbative. The general form of the series in g can be shown to be ²

$$m_D = m_D^0 + \frac{N}{4\pi} g^2 T \ln \frac{m_D^0}{g^2 T} + c_N g^2 T + \mathcal{O}(g^3 T), \quad (6)$$

which is non-analytic in the coupling constant. While the coefficient of the logarithm is fixed perturbatively, c_N is entirely non-perturbative. The reason is that, starting from this order in g , the non-abelian A_0 couples to the soft magnetic gluons $A_i \sim g^2 T$, and hence becomes sensitive to the non-abelian infrared divergencies in the magnetic sector for which there is no perturbative cure ³: $\Pi_{ii}(k_0 = 0, \mathbf{k} \rightarrow 0) \sim g^2 T \neq 0$, with contributions from all loop orders two and larger.

This raises the conceptual problem whether a perturbative definition of Debye screening in QCD is at all sensible.

don't give up! ask lattice QCD



coupling drops off for $r > 0.3$ fm

screening masses from gluon propagator

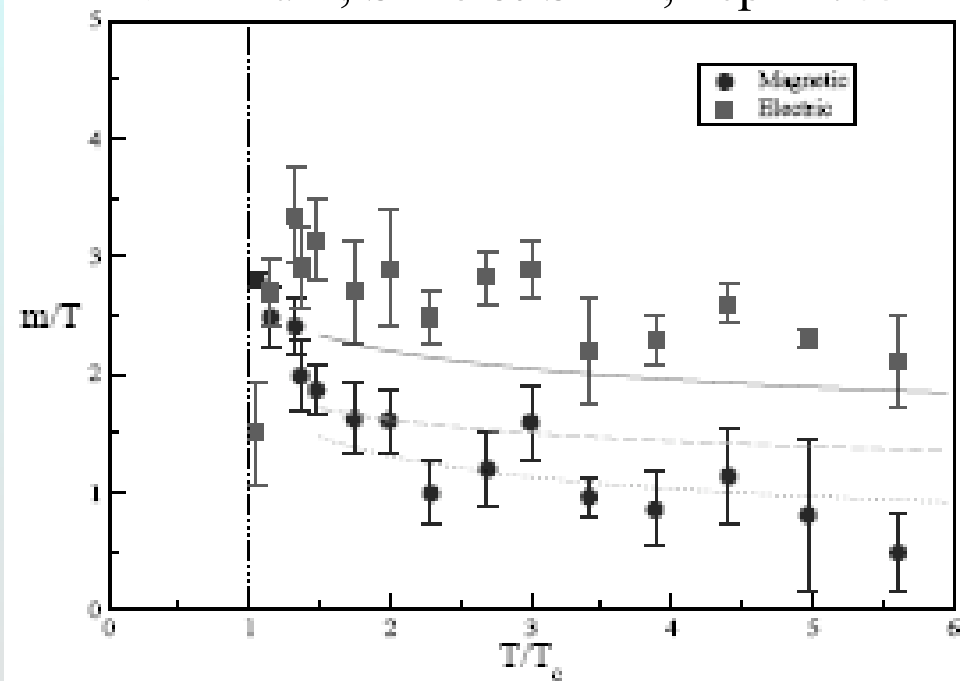
Screening mass, m_D , defines inverse length scale
Inside this distance, an equilibrated plasma is
sensitive to insertion of a static source
Outside it's not.

T dependence of electric & magnetic screening masses
Quenched lattice study
of gluon propagator

figure shows:

$$m_{D,m} = 3T_c, \quad m_{D,e} = 6T_c \text{ at } 2T_c$$
$$\therefore \lambda_D \sim 0.4, 0.2 \text{ fm}$$

Nakamura, Saito & Sakai, hep-lat/0311024



magnetic screening mass is non-zero
not very gauge-dependent, but DOES
grow w/ lattice size (long range is important)

Implications of $\lambda_D \sim 0.3$ fm?

- can use to estimate Coupling parameter, Γ

- $\Gamma = \langle PE \rangle / \langle KE \rangle$ but also given by $\Gamma = 1/N_D$

for $\lambda_D = 0.3$ fm and $\varepsilon = 15$ GeV/fm³

$$V_D = 4/3 \pi \lambda_D^3 = 0.113 \text{ fm}^3$$

$$E_D = 1.7 \text{ GeV}$$

to convert to number of quasiparticles:

use gT or g^2T for $T \sim 2T_c$ and $g^2 = 4$

get $N_D = 1.2 - 2.5 \quad \therefore \Gamma \sim 1$

- NB: for $\Gamma \sim 1$

plasma is NOT fully screened – strongly coupled!
affects interaction σ !

*strongly coupled EM plasmas behave as liquids,
can even make crystals for $\Gamma \geq 150$*

*dusty plasmas, cold atoms+ions, warm dense
matter*

calculating transport in QGP

weak coupling limit

perturbative QCD

kinetic theory, cascades

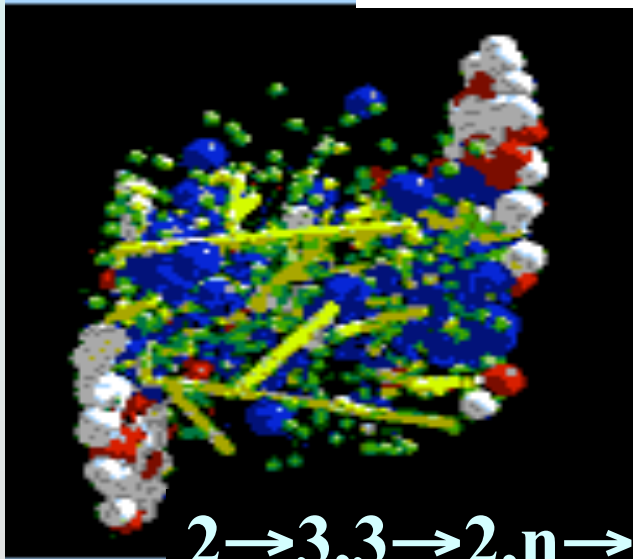
strong coupling limit

not so easy!

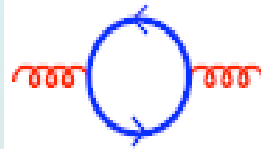
gravity \leftrightarrow supersym 4-d

AdS/CFT: QCD-like theory

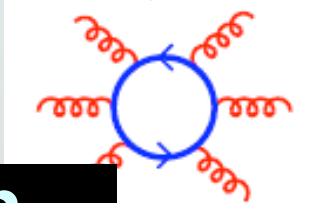
resummation of hard thermal loops



$2 \rightarrow 3, 3 \rightarrow 2, n \rightarrow 2 \dots$



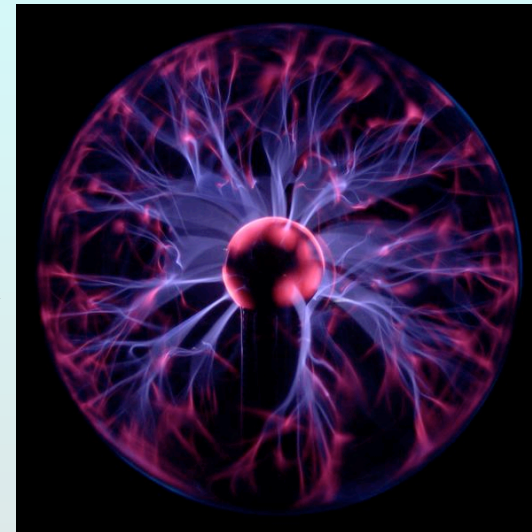
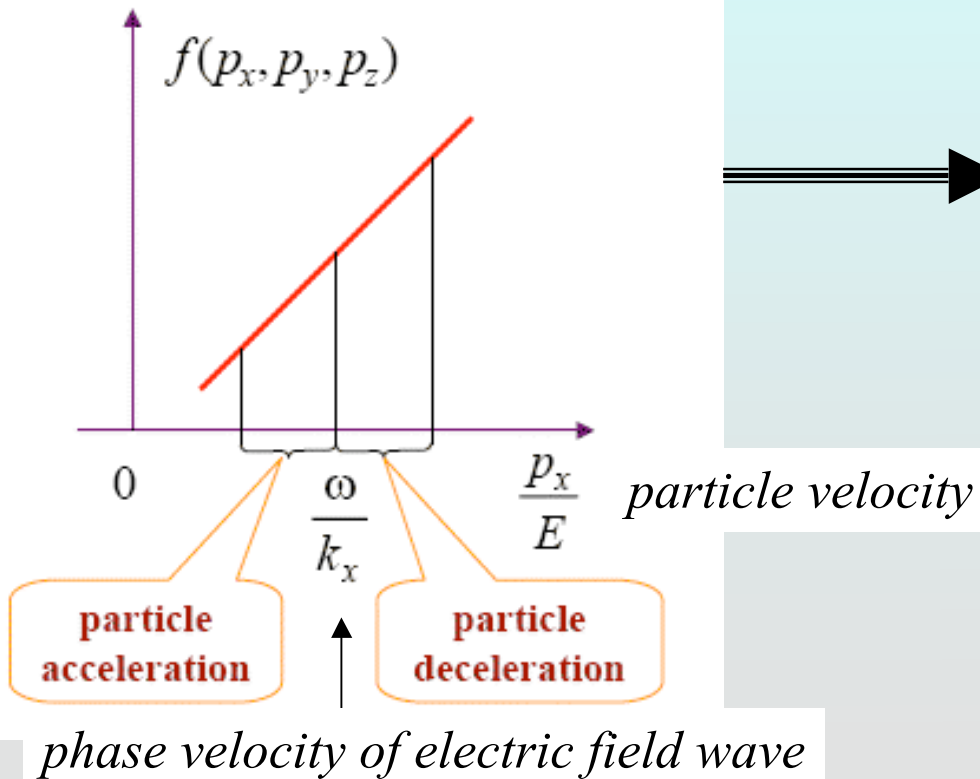
the 1-loop self-energy for gluons.



Plasma instabilities

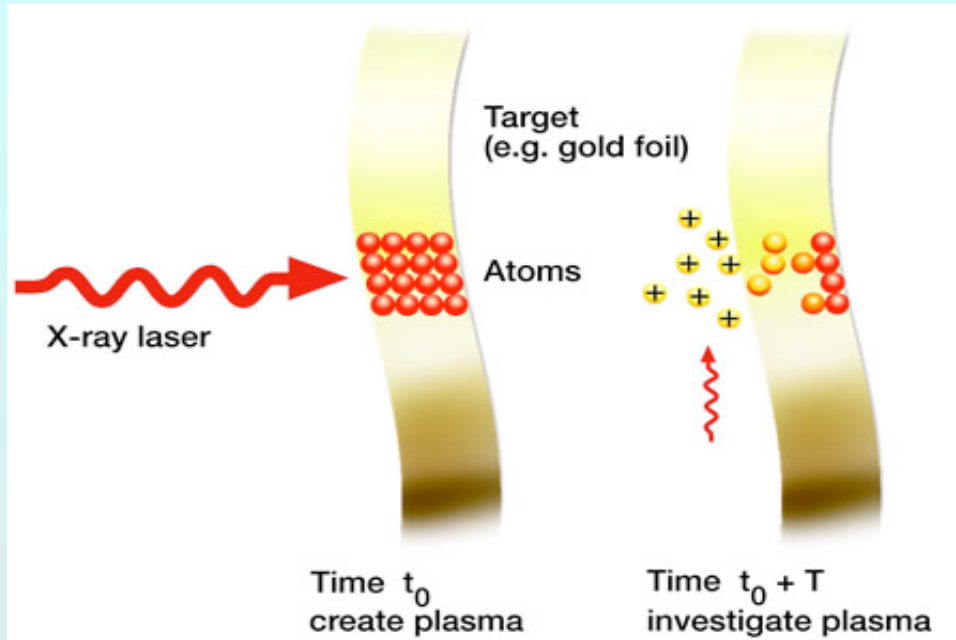
- a mode with growing amplitude
transfer energy from the plasma particles to EM field
- e.g. Weibel instability - causes beam filamentation

Electric field grows - **instability**

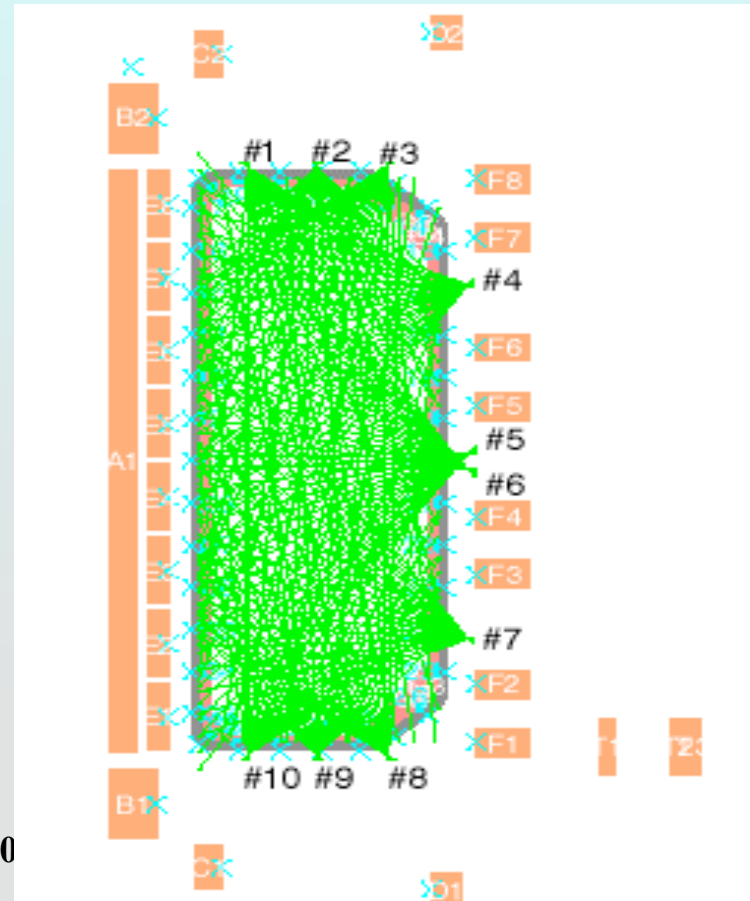


Mechanism to quickly isotropize momentum of colliding heavy ions?
Non-abelian plasma is self-limiting

Some typical plasma diagnostics



X-ray tomography: to get temporal and spatial evolution of plasma mode activity, we use 10 soft X-ray cameras with 20 channels each, at 70 kHz max.

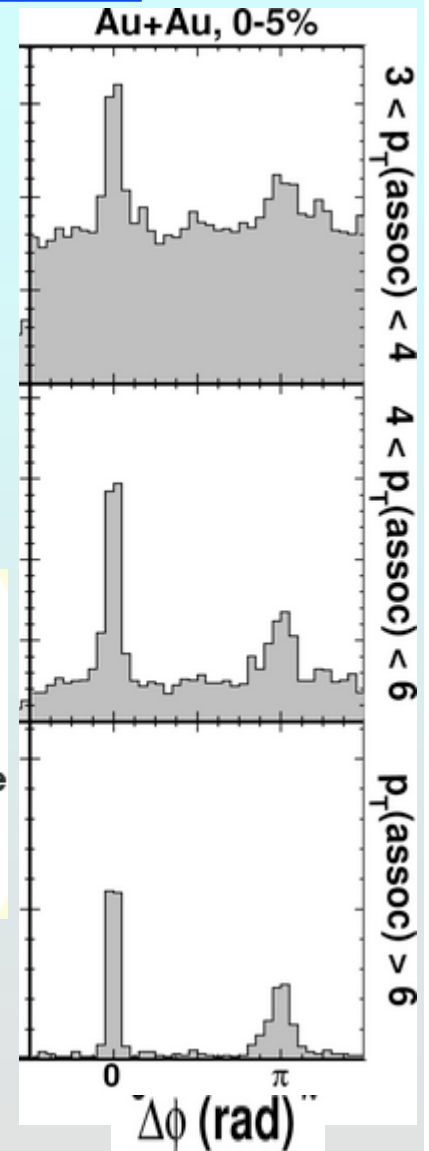
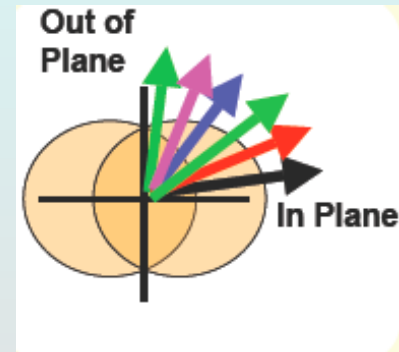


Parton energy loss: similar probe

- Instead of an external x-ray beam use partons from initial hard scattering “self-generated” probe → density

- R_{AA} integrates over all directions
Reaction plane dependent R_{AA} allows selecting a direction
Back-to-back jets control path length in plasma
di-hadrons work too

- Next step: “calibrate” the probe
 γ_{direct} - jet or γ -h, reconstructed jets



measure dynamic (transport) plasma properties

- measure dynamic properties - transport of:

particle number, energy, momentum, charge

diffusion *sound* *viscosity* *conductivity*

- emission from the bulk can reflect collective motion
- but other useful probes require
auto-generation in the heavy ion collision
large Q^2 processes to separate production & propagation
large E_{tot} (high p_T or M) to set scale other than $T(\text{plasma})$

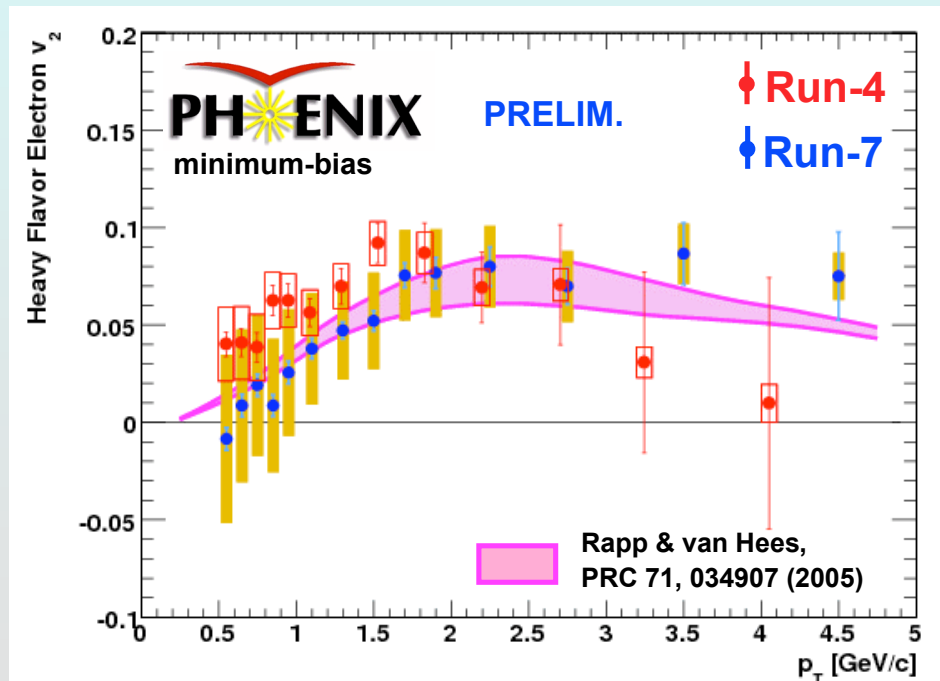
How to get at these

<u>property</u>	<u>measurement</u>	
T	γ, γ^* as fn. of ϵ, μ	
density	quark, gluon energy loss	
equation of state	particle flows as fn. of ϵ, μ critical point location	←
screening length	onium spectroscopy	←
$\rho(x, v)$	jet tomography	
diffusion	open C, B spectra & flow	←
viscosity	light & charmed hadron flows used to constrain 3d hydro	
energy transport i.e. sound	>2 particle correlations vs. T, p_T Medium response to jet energy	←

Diffusion (& viscosity, redux)

Consider heavy quark probes:

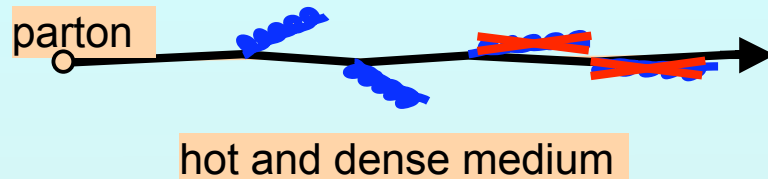
- Hard probe - production in initial NN collisions (maybe some thermal production at LHC...)
- Large mass: different scale for probe than medium



Charm actually flows along with the bulk medium!

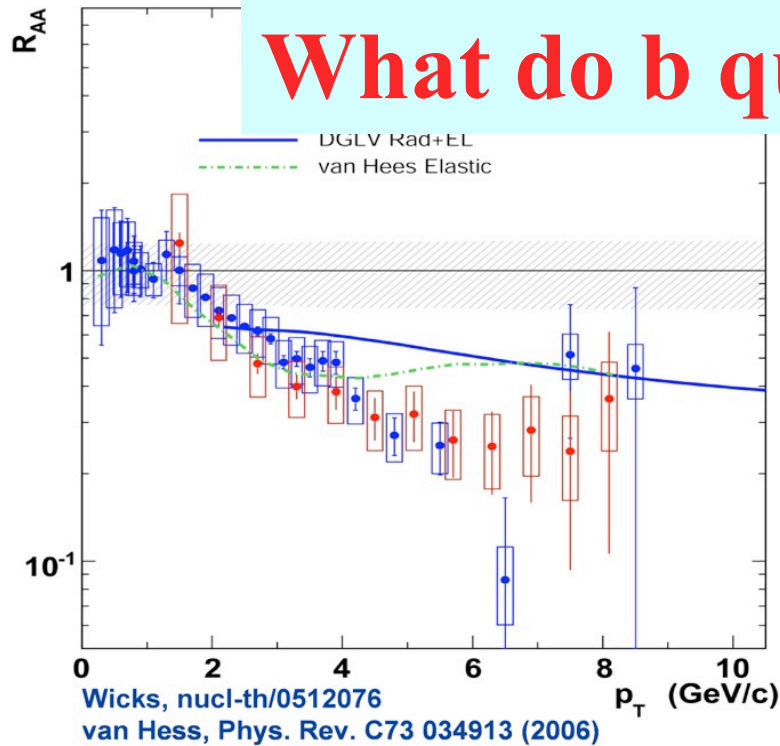
Consider how c quarks diffuse through the plasma

Heavy quarks radiate AND collide



* In strongly coupled medium: diffusion of charm quarks
 $D \propto$ collision time,
 determines relaxation time

STAR Phys. Rev. Lett. 98 (2007) 192301
 PHENIX Phys.Rev.Lett.98 (2007) 172301



Wicks, nucl-th/0512076
 van Hees, Phys. Rev. C73 034913 (2006)

$\langle v \rangle \lambda$

NB: $\eta = 1/3 \rho \langle v \rangle \lambda$

$$\therefore D = \eta/\rho \sim \eta/s$$

→ extract D get $\eta/s!$

$D \sim 4-6/(2\pi T)$ for charm

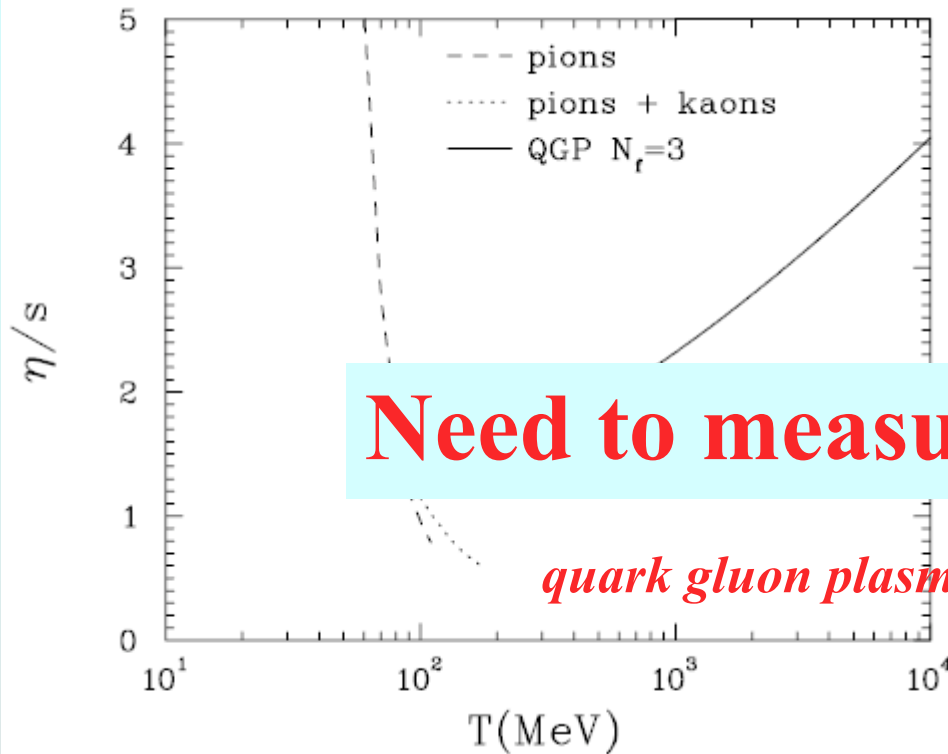
$$\eta/s = (1.3 - 2.0)/4\pi$$

(model dependent,
 but not on hydro)

minimum η/s at phase boundary?

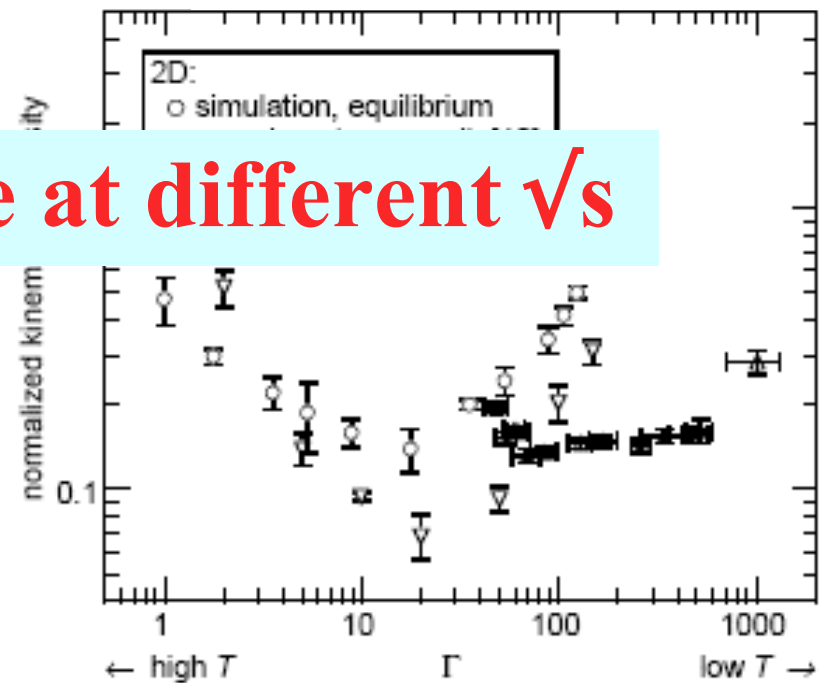
strongly coupled dusty plasma

B. Liu and J. Goree,
cond-mat/0502009



Need to measure at different \sqrt{s}

quark gluon plasma

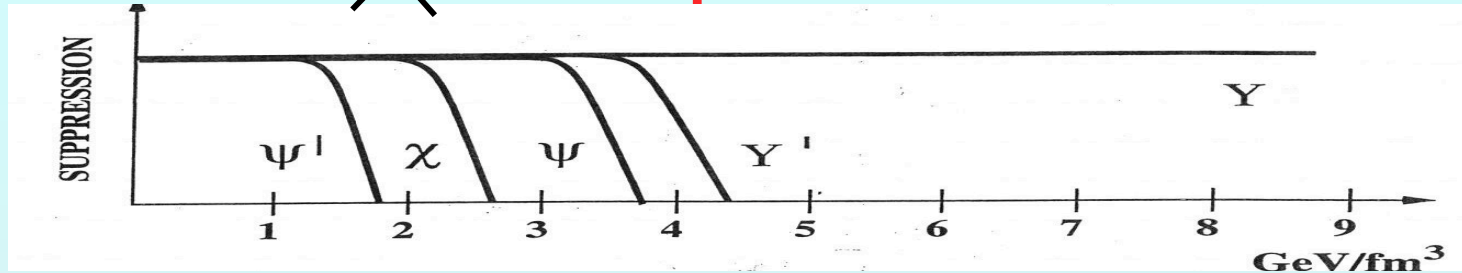


Csernai, Kapusta & McLerran
PRL97, 152303 (2006)

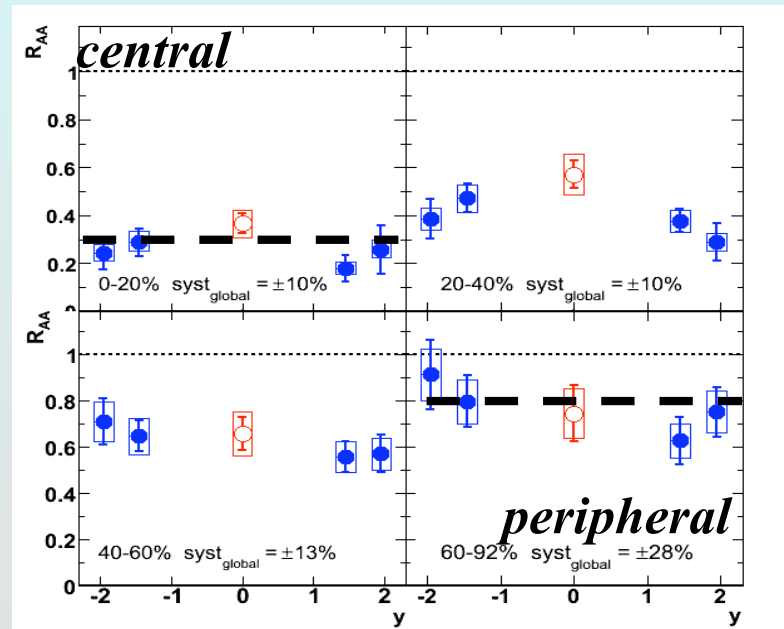
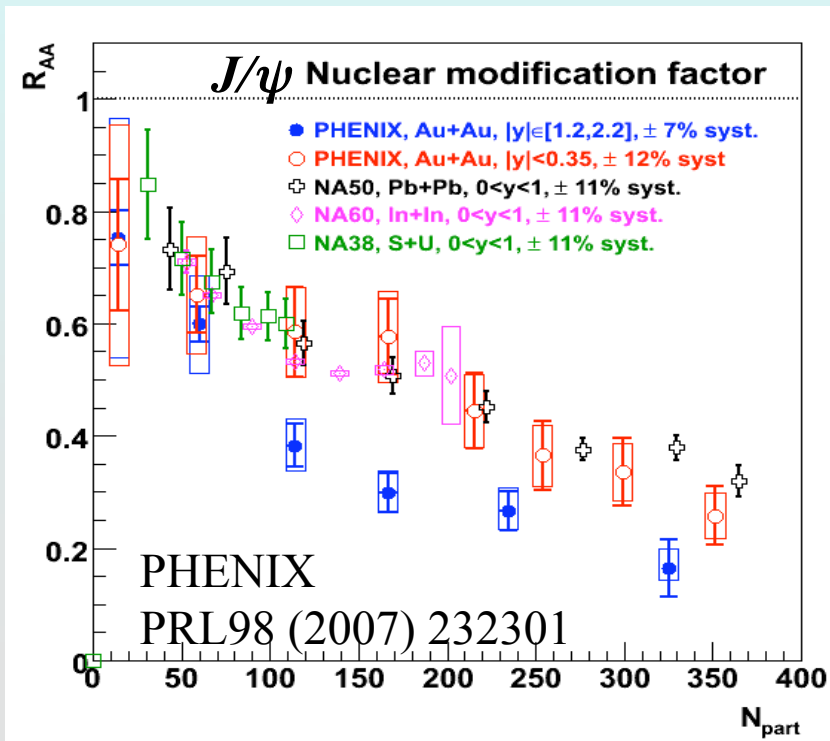
minimum observed in other strongly coupled systems –
kinetic part of η decreases with Γ while potential part increases

Screening length

- The idea ~~is~~ was simple



- Nature is not



*Disentangle final state coalescence
& feeding from plasma screening!*

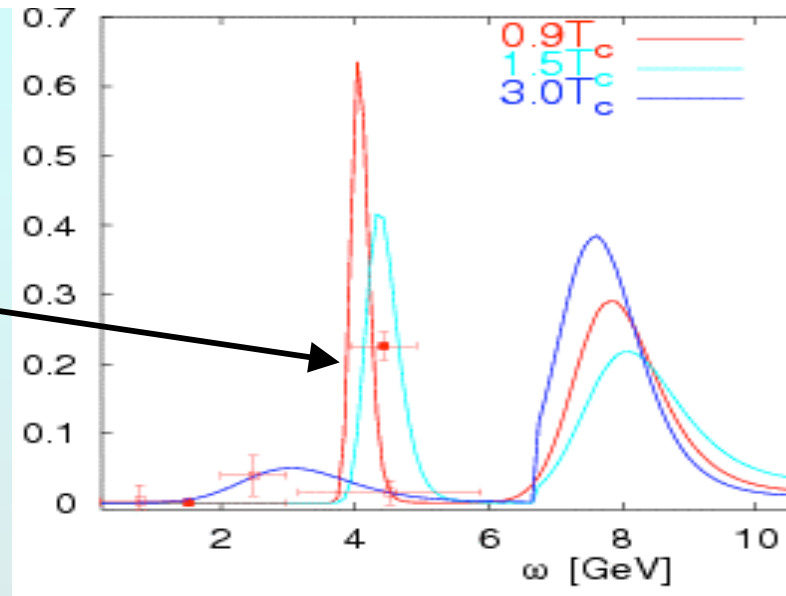
Lattice: J/ψ fully dissolved?

Lattice QCD shows $c\bar{c}$ correlations at $T > T_c$, implying that interactions are not zero

Expect in strongly coupled plasma?

Big debate: J/ψ all gone?
resonant states?
leftover correlations?

Asakawa & Hatsuda, PRL92, 012001 (2004)

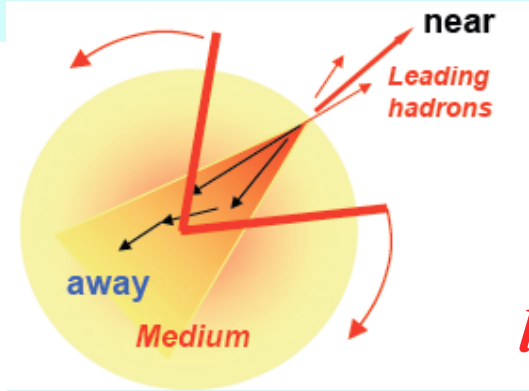
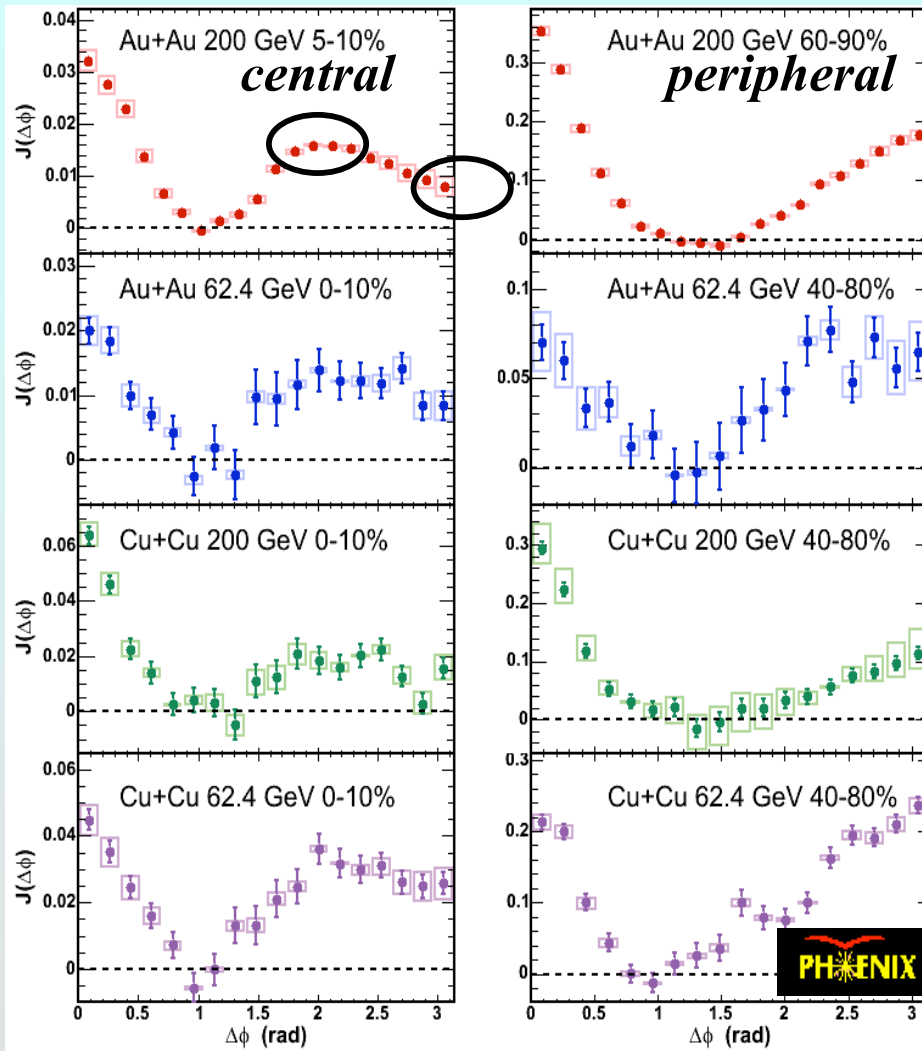


Mocsy & Petrecky, PRL99, 211602 (2007)

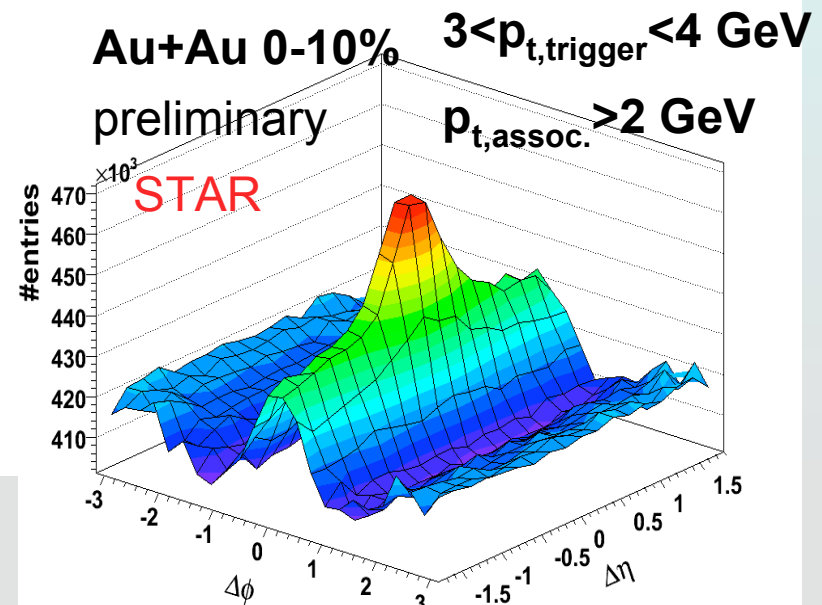
state	χ_c	ψ'	J/ψ	Υ'	χ_b	Υ
T_{dis}	$\leq T_c$	$\leq T_c$	$1.2T_c$	$1.2T_c$	$1.3T_c$	$2T_c$

TABLE I: Upper bound on dissociation temperatures.

Where does lost jet energy go?

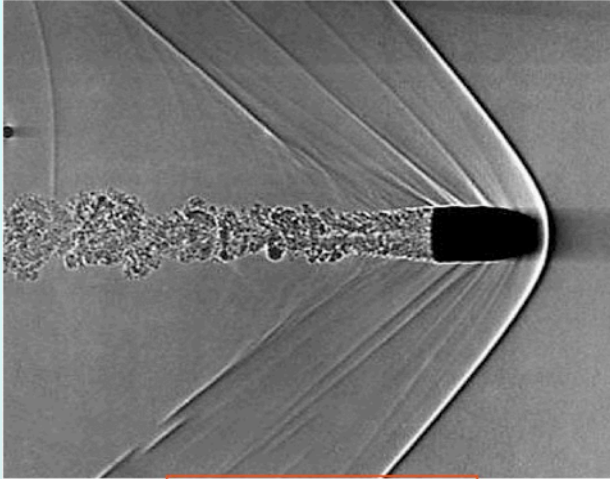


low p_T looks very modified medium response?



$1 < p_{T,a} < 2.5 < p_{T,t} < 4$ GeV/c

excites a sound (density) wave?



$$\cos(\theta_M) = c_s$$

Peak: $\phi = \pi - (1.2-1.38)$

→ speed of sound

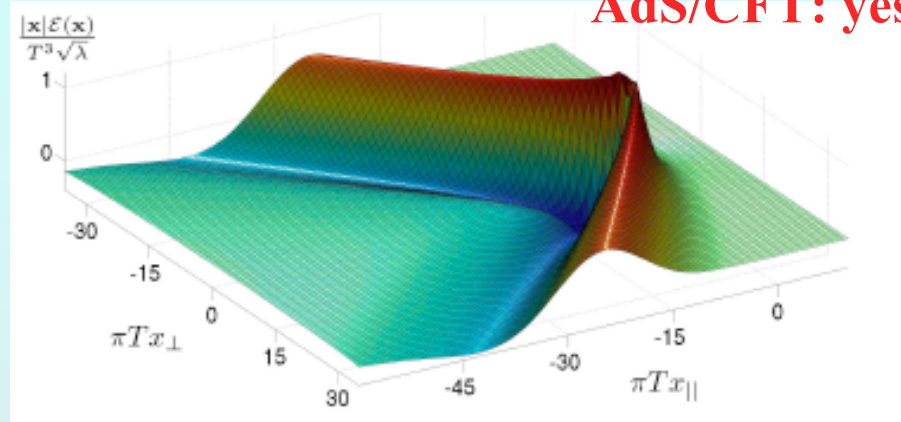
$c_s \sim 0.2-0.4$

($c_s^2 = 0.33$ in QGP,

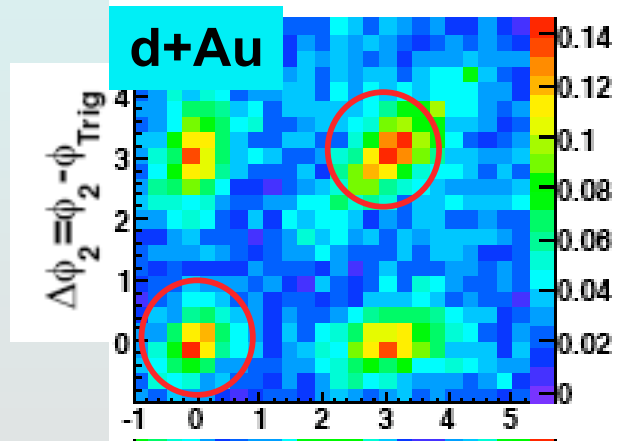
~ 0.19 in hadron gas)

Chesler & Yaffe, 0706.0368(hep-th)

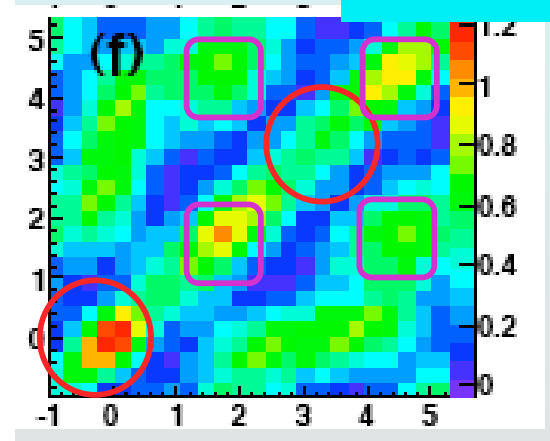
AdS/CFT: yes it does!



d+Au



Au+Au



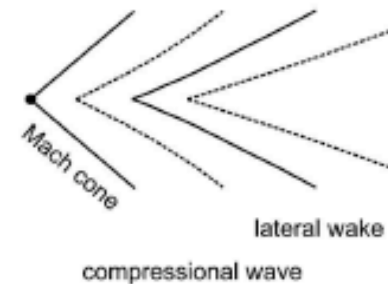
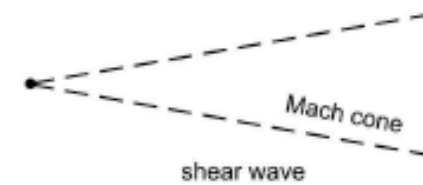
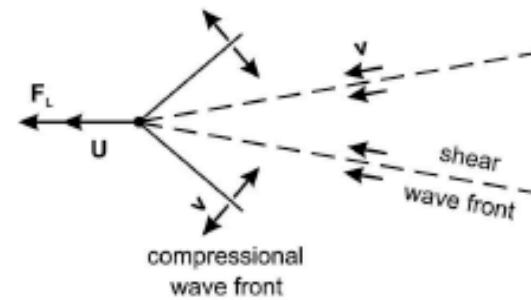
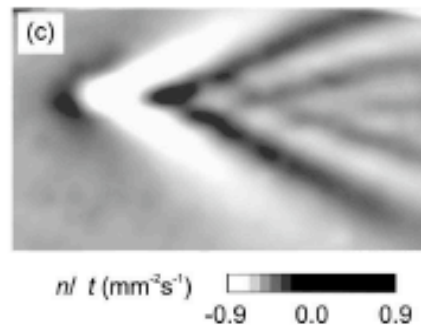
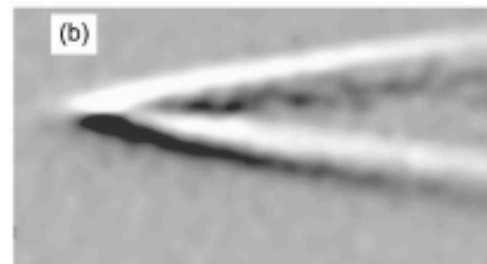
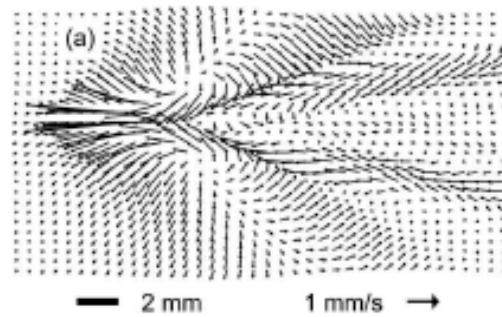
$$\Delta\phi_1 = \phi_1 - \phi_{\text{Trig}}$$

3 particle correlation STAR, 0805.0622

Compressional and shear wakes in a two-dimensional dusty plasma crystal

NOSENKO *et al.*

PHYSICAL REVIEW E 68, 056409 (2003)



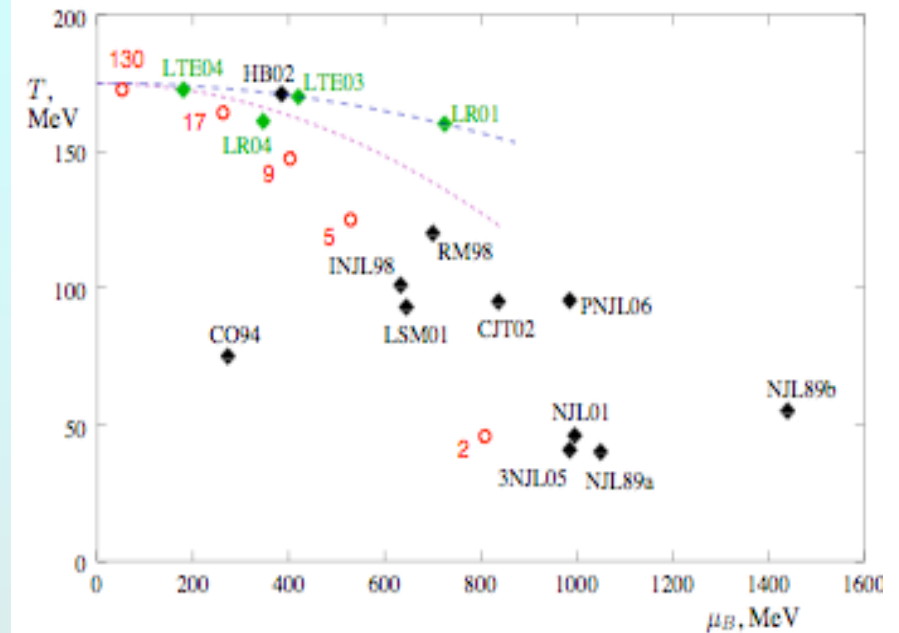
Are shocks in strongly coupled EM plasma. QGP also shown to support shockwaves

(shear generally a phenomenon in crystals but not liquids)

FIG. 4. The compressional- and shear-wave Mach cones, excited simultaneously. The scanning speed U is higher than the sound speed for both the compressional and shear waves. Maps are shown for (a) particle velocity \mathbf{v} , (b) vorticity $|\nabla \times \mathbf{v}|$, and (c) $\partial n / \partial t$, where n is the particle areal number density.

Critical endpoint → EOS

- *Where is it?*
Start at SPS & RHIC
Precision & rarer probes
at FAIR
- Bulk variables
Spectra, $\langle p_T \rangle$, yield ratios
many ~known from SPS
Fluctuations
2nd order phase transition - long range correlations
lots of work to do here (experiment AND theory)
Bulk collective flow
 v_2 scaling behavior near CEP? Need π, K, p v_2
HBT re-emerges as good probe of the space-time
evolution as \sqrt{s} decreases



Rare CEP probes

- **Multi-particle correlations**

In addition to fluctuations signatures*

- **Di-leptons and photons**

Use \sqrt{s} dependence to help disentangle hadron gas and partonic radiation*

Not only T_{init} , but also relative lifetimes change

Rate of partonic contribution increases with \sqrt{s}

Hadron gas always starts at T_c , lifetime decreases with \sqrt{s}

- **Charm approaching threshold**

How much does this add without confusion?*

*** Need help from theory!**

In closing:

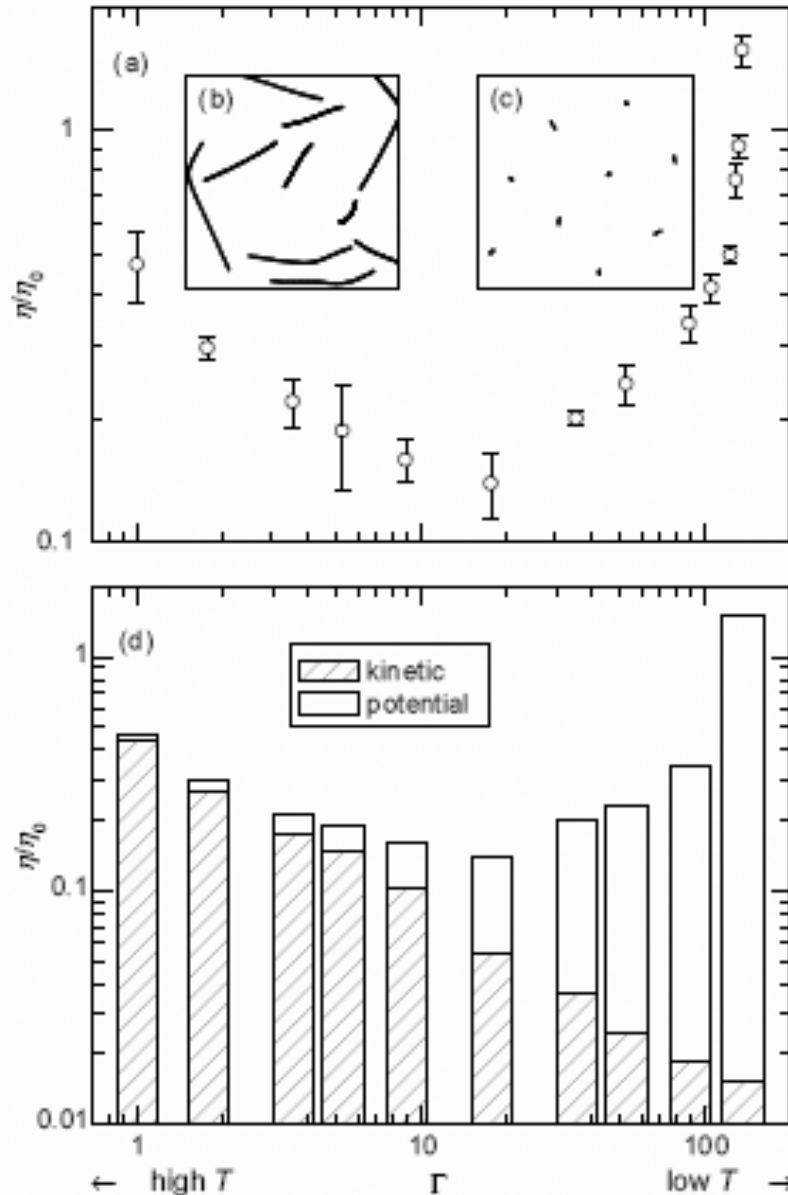
- **Nuclear physics with high energy heavy ion collisions connects to other branches of physics!**
- **Particle physics - obviously**
- **Condensed matter**
 - Strongly correlated systems
 - Condensates, cold atoms (Hatsuda's talk)
- **Plasma physics**
 - Theoretical ideas/tools in common
 - also molecular dynamics
 - Experimental approaches to learn matter properties
- **String theory**
 - Tool via AdS/CFT for ∞ strongly coupled limit
 - This connection was a big surprise!***

i'm in ur fizx lab



testn ur string therry

use this technique to measure viscosity



melt crystal with laser light
induce a shear flow (laminar)
image the dust to get velocity
study:

spatial profiles $v_x(y)$

moments, fluctuations $\rightarrow T(x,y)$

curvature of velocity profile

\rightarrow drag forces

viscous transport of drag in
 \perp direction from laser

compare to viscous hydro.

extract η/ρ

shear viscosity/mass density

PE vs. KE competition governs

coupling & phase of matter

Csernqi, Kapusta, McLerran nucl-th/0604032

important differences

- Medium created by colliding beams doesn't last long

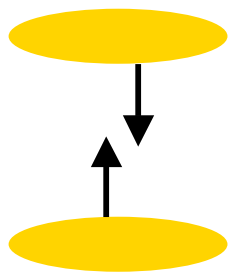
lifetime $\sim 10 \text{ fm}/c \sim 3 \times 10^{-23} \text{ sec}$ (3×10^{-5} attosecond)

- Relevant energy scale given by $T \sim 200 \text{ MeV}$
wavelength of a photon with energy of 200 MeV:
 $\sim 6 \text{ fm} = 6 \times 10^{-6} \text{ nm}$
1/1000th that of gamma rays!

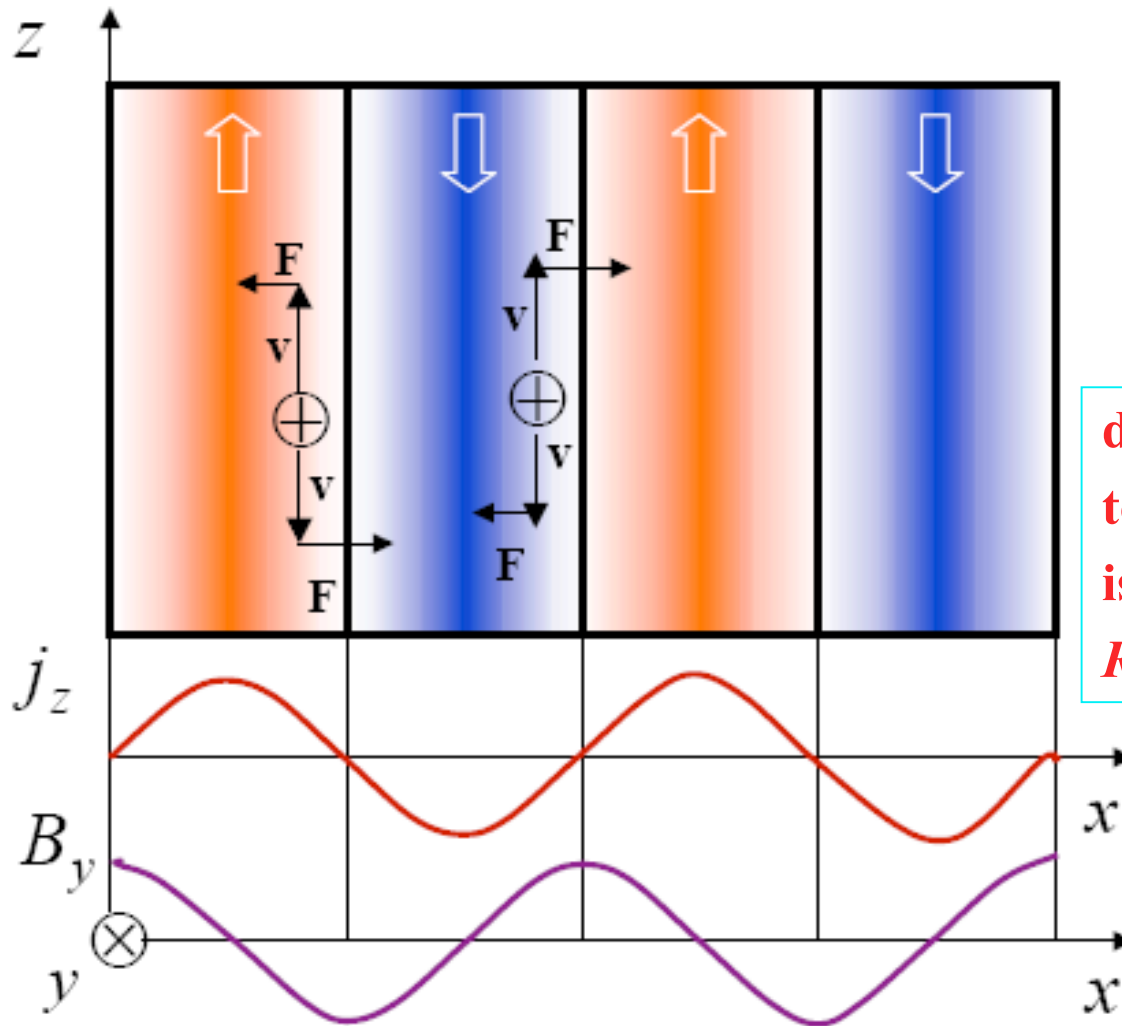
Not going to study this plasma with any laser known to man!

To study a plasma of gluons, want a probe sensitive to color not charge

Mechanism of filamentation



Ampere's law
 $\nabla \times \mathbf{B} = \mathbf{j}$



Lorentz force
 $\mathbf{F} = q \mathbf{v} \times \mathbf{B}$

drives system
toward fast
isotropization?
Rebhan

QGP viscosity is lower than that of liquid He!

η/s ~4 times smaller than of helium (but at $T \sim 10^{12}$ K)

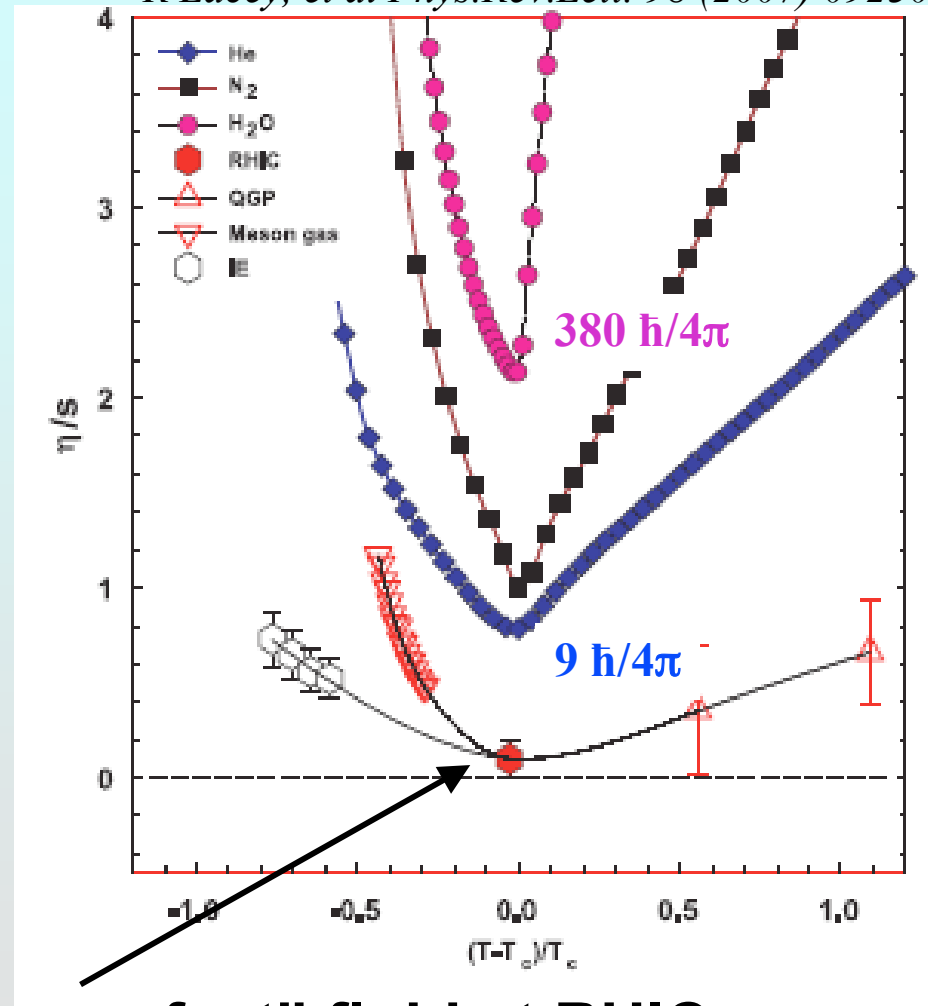
close to the conjectured limit $\eta/s = \hbar/4\pi$

P. Kovtun, D.T. Son, A.O. Starinets, Phys.Rev.Lett.94:111601, 2005

motivated by AdS/CFT (Anti-deSitter/Conformal Field Theory) correspondence

J. Maldacena: Adv. Theor. Math. Phys. 2, 231, 1998

R Lacey, et al Phys.Rev.Lett. 98 (2007) 092301



"near perfect" fluid at RHIC

expect low viscosity in strongly coupled plasma

Gelman, Shuryak, Zahed, nucl-th/0601029

◇ Shear viscosity η in (classical) quark gluon plasma

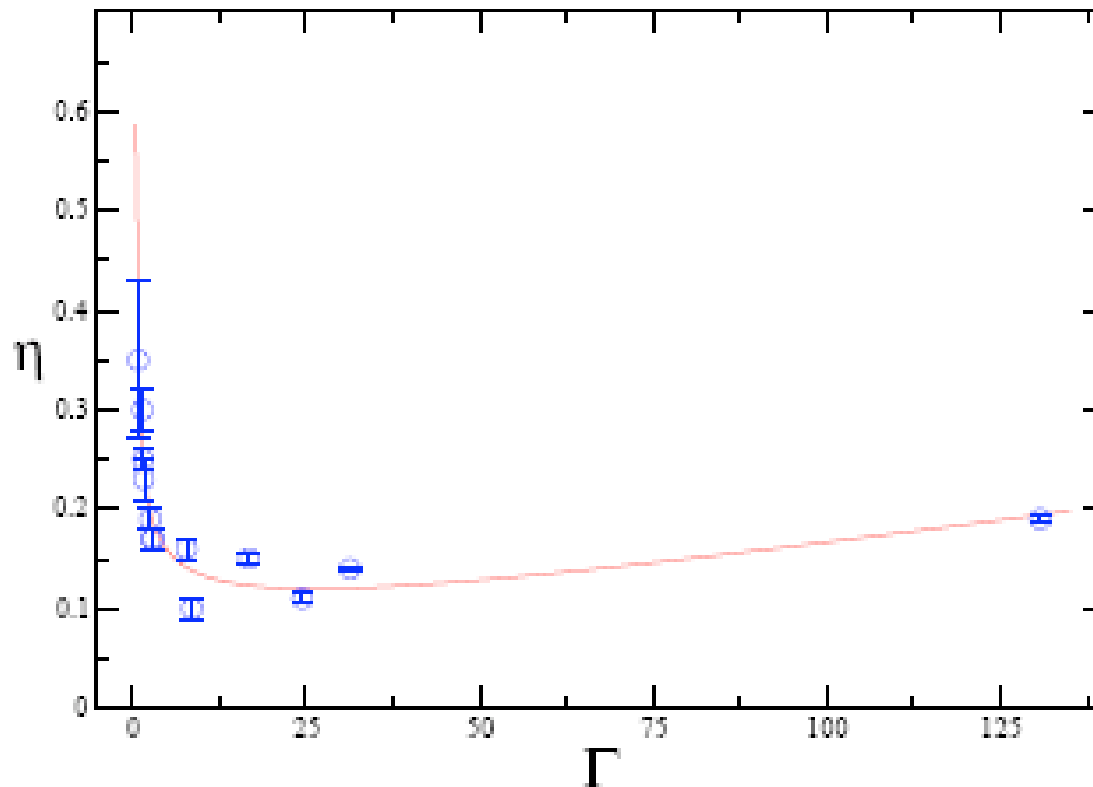


FIG. 5. Shear viscosity of one species cQGP as a function of the dimensionless coupling Γ . Blue points are the MD simulations; the red curve is the fit in eq. (32).