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## Chiral Partners and their Electromagnetic Radiation

Ingredients for a systematic in-medium calculation

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Erice School, Sicily, September 2008

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| Isospin partners |        |  |  |  |  |  |  |

- observe same spectra in different channels
- in particular: peaks (hadrons) at same position
- e.g. spectra of  $j^0_{\mu}$  (from  $e^+e^-$ ) and  $j^-_{\mu}$  (from  $\tau^-$  decay)

$$j^0_\mu = rac{1}{2} (ar u \gamma_\mu u - ar d \gamma_\mu d) \qquad o \qquad j^-_\mu = ar u \gamma_\mu d$$

 $\rightsquigarrow$  isospin partners  $\rho^0$  and  $\rho^-$ 



# Chiral partners Nature of chiral partners Chiral restoration VMD In-medium calculations Summary and outlook How do we know that chiral symmetry is broken?

study now instead of isospin transformation

$$j^0_\mu = rac{1}{2} (ar u \gamma_\mu u - ar d \gamma_\mu d) \qquad o \qquad j^-_\mu = ar u \gamma_\mu d$$

chiral transformation (from now on  $SU_F(3)$ )

$$j^0_\mu \qquad o \qquad j^b_\mu = ar q \, \lambda^b \gamma_5 \gamma_\mu q$$

- consequence of chirally symmetric world would be: same spectral information in vector and axial-vector current-current correlators (degeneracy)
- observable?  $\rightsquigarrow \tau$  decay

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#### Chiral symmetry breaking and $\tau$ decays

study decay  $\tau \rightarrow \nu_{\tau} + hadrons$ :

- couples to V-A (weak process)
- G parity: V/A couples to even/odd number of pions

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• are V and A spectra identical?



#### Chiral symmetry breaking and $\tau$ decays

study decay  $\tau \rightarrow \nu_{\tau}$ +hadrons:

- couples to V–A (weak process)
- G parity: V/A couples to even/odd number of pions
- are V and A spectra identical?
- → Phys. Rept. 421, 191 (2005) (ALEPH):





#### One of the clearest signs of chiral symmetry breaking



 $v_1: \tau \to \nu_\tau + m\pi$ (*m* even)  $a_1: \tau \to \nu_\tau + n\pi$ (*n* odd)



#### How to find the chiral partners?

- suggestive: get partners by relating lowest peaks/bumps in corresponding spectra
- before: isospin partners  $\rho^0$  and  $\rho^{\pm}$  (same mass)
- $\hookrightarrow$  now:  $\rho$  multiplet related to  $a_1$  multiplet
- not same mass due to spontaneous symmetry breaking!

→ conjecture: not even of same nature!

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#### Nature of chiral partners

- start with lowest states (of quark model): nucleon octet, pion nonet, Δ decuplet, ρ nonet
- → "LLH" = lowest-lying hadrons
- $\hookrightarrow$  LLH are quark-antiquark or three-quark states, respectively
  - conjecture:

chiral partners of LLH are dynamically generated states, i.e. hadron "molecules"

- N\*(1535) from coupled-channel dynamics of ηN, KΛ, ... Kaiser/Siegel/Weise, Phys. Lett. B362, 23, 1995
- σ meson from ππ ...
   Oller/Oset, Phys. Rev. D60, 074023, 1999
- $\Delta^*(1700)$  and  $N^*(1520)$  from  $\pi\Delta$  and flavor partners Lutz/Kolomeitsev, Phys. Lett. B585, 243, 2004
- a<sub>1</sub> (and b<sub>1</sub>) multiplets from πρ, (πω) and flavor partners Lutz/Kolomeitsev, Nucl. Phys. A 730, 392, 2004



#### General framework for dynamical generation

 study scattering of LLH state on Goldstone bosons in channel of interest (quantum numbers of chiral partner)

$$\rightarrow \text{ Bethe-Salpeter equation } T = K + K T$$

- interaction kernel always of same type:
  - lowest order chiral interaction

- $\rightsquigarrow$  strength fixed model independently  $\sim F_{\pi}^{-2}$
- → Weinberg-Tomozawa (WT) point interaction



 N.B.: renormalization point for loop fixed Lutz/Kolomeitsev, Nucl. Phys. A 730, 392, 2004 Hyodo/Jido/Hosaka, arXiv:0803.2550 [nucl-th]

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#### Spectra of chiral partners

- Low-energy parts of spectra:
  - *ρ*-meson in vector channel (left yellow)
  - *a*<sub>1</sub>-meson in axial-vector channel (right green)



→ How to understand spectra – resonances?

| Mature          |                           |                    |     |                        |                     |
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Nature of the *a*<sub>1</sub> meson

- experimental finding from τ decays (Dalitz plots): isovector–axial-vector current couples to π-ρ
- $\pi$ - $\rho$  system subject to final-state interactions (rescattering)
- experimental finding: resonant structure at  $\approx$  1250 MeV
- conjecture: emerges from final-state interaction of π-ρ Lutz/Kolomeitsev, Nucl. Phys. A 730, 392, 2004
   Roca/Oset/Singh, Phys. Rev. D 72, 014002, 2005
  - describe final-state interactions via Bethe-Salpeter eq., kernel from lowest order chiral interaction (Weinberg-Tomozawa - WT)
     parameter free

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#### Description of $a_1$ as final-state interaction effect

parameters for  $\tau \rightarrow \nu_{\tau} + 3\pi$ : renormalization points

• for loop for transition from W to hadrons



- ← renormalization point should be in reasonable range
  - for loop in Bethe-Salpeter equation (rescattering)

 → renormalization point fixed (cf. Lutz/Kolomeitsev, Nucl. Phys. A 730, 392 (2004); Hyodo/Jido/Hosaka, arXiv:0803.2550 [nucl-th])

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#### au decay



- reasonable description with one free parameter
- → indicates that  $a_1$  is  $\rho$ - $\pi$  "molecule" (Markus Wagner and S.L., PRD in press, arXiv:0801.0814 [hep-ph])



#### How does chiral restoration take place?

- typically spontaneous symmetry breaking lifted at some temperature/density (Ferro magnet: Curie temperature)
- → consequence at point of chiral restoration: same in-medium spectral information in vector and axial-vector channel
  - how does it look like?  $\rightarrow$  various (> 2) scenarios
- $\hookrightarrow$  scenario 1 (degenerate states):
  - $\rho$  meson is still (dominantly) single-particle state
  - $\rightarrow$  requires chiral partner which is also a single-particle state
  - → very high mass in vacuum (since  $\neq a_1$  meson) → ??
- $\hookrightarrow$  scenario 2 (melting):
  - *ρ* meson dissolves already in hadronic matter (precursor of deconfinement)
  - $\rightarrow a_1$  meson should also dissolve
  - → testable in our approach

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#### Dissolution of the $a_1$ meson?

- very simple model:
  - $\Gamma_{
    ho} \rightarrow 200, \ 250 \, {\rm MeV}$
  - no changes to pion
  - no momentum dep.
  - $\hookrightarrow$  can be improved







- broader  $\rho$  meson leads to broader  $a_1$  meson
- → no proof that melting scenario is correct, but at least consistent picture
- → problem of missing chiral partner (at single-particle level) solved by deconfinement

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#### Necessity for a systematic treatment

- for understanding of nature of resonances and for systematic in-medium calculations
- → need effective field theory for pion nonet,  $\rho$  nonet, nucleon octet,  $\Delta$  decuplet
- systematic calculations, i.e. with power counting, instead of models
  - suggested for meson sector in Lutz/Leupold, NPA in press, arXiv:0801.3821 [nucl-th]:
    - treat pseudoscalar and vector mesons as soft
    - allows for systematic inclusion of decays of vector mesons
    - yields clear statements about validity of vector-meson dominance (VMD)
    - vector mesons represented by antisymmetric tensor fields

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#### Extended VMD for elementary hadrons

example: decays of  $\omega$  meson





- both processes  $\omega \rightarrow \gamma \pi$  and  $\omega \rightarrow 3\pi$  in leading order given by VMD
- use first process to fix coupling of second one
- → prediction:  $\Gamma_{\omega \to 3\pi} = 7.3 \text{ MeV}$   $\Gamma^{exp}_{\omega \to 3\pi} = (7.57 \pm 0.13) \text{ MeV}$ Leupold/Lutz,

arXiv:0807.4686 [hep-ph]



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#### No VMD for hadron molecules

example: radiative decays of axial-vector mesons Lutz/Leupold, arXiv:0801.3821 [nucl-th]





#### No VMD for hadron molecules

example: radiative decays of axial-vector mesons

- formation of axial-vector meson dominated by Weinberg-Tomozawa (point) interaction
- radiative decay: VMD contribution and coupling of photon to constituents
- calculable from electromagnetic moments of constituents

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 uncertainty: dipole and quadrupole moments of vector mesons (lattice input?)

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### Towards self-consistent in-medium calculations

#### • ingredients:

- systematic vacuum input
- thermodynamically consistent resummation scheme ("Φ derivable")
- → to account self-consistently for "changes induced by changes"
  - current conservation in resummation scheme ensured by use of antisymmetric tensor fields (Leupold, Phys. Lett. B646, 155, 2007)
- $\hookrightarrow$  no proliferation of non-conserving part of current:

$$j^{\mu}\partial^{\nu}V_{\nu\mu} = j^{\mu}\left(\underbrace{g_{\mu\alpha} - \frac{\partial_{\mu}\partial_{\alpha}}{\partial^{2}}}_{P^{T}_{\mu\alpha}} + \underbrace{\frac{\partial_{\mu}\partial_{\alpha}}{\partial^{2}}}_{P^{L}_{\mu\alpha}}\right)\partial_{\nu}V^{\nu\alpha} = j^{\mu}P^{T}_{\mu\alpha}\partial_{\nu}V^{\nu\alpha}$$

→ check e.g. melting scenario at chiral restoration



#### Summary and outlook

- for the lowest-lying hadrons
  - (LLH = pion nonet,  $\rho$  nonet, nucleon octet,  $\Delta$  decuplet) the chiral partners can be understood as being dynamically generated ("hadron molecules")
- suggestion for systematic counting, i.e. effective field theory for LLH
- $\hookrightarrow\,$  prediction where VMD works and where not
  - opens the way for systematic and self-consistent in-medium calculations

Many thanks to my collaborators Markus Wagner and Matthias Lutz

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#### Backup: How to find the chiral partners?

general strategy:

- start with hadron
- Ind quark current with same quantum numbers
- apply chiral transformation
- determine quantum numbers of resulting current

Iook for state in PDG



#### Backup: How to find the chiral partners?

general strategy:

- start with hadron
- find quark current with same quantum numbers
- apply chiral transformation
- determine quantum numbers of resulting current

Iook for state in PDG

but:

- step 2 not unique for mesons and baryons
- step 4 ill-defined for baryons



- problem: corresponding quark current not unique
- e.g.  $\bar{q}\lambda_b\gamma^{\mu}q$  and  $\partial_{\nu}(\bar{q}\lambda_b\sigma^{\mu\nu}q)$  have quantum numbers of  $\rho$  multiplet

- chiral transformation on  $\bar{q}\lambda_b\gamma^\mu q$  leads to  $\bar{q}\lambda_c\gamma_5\gamma^\mu q$
- $\hookrightarrow$  quantum numbers of  $a_1$  multiplet
  - chiral transformation on  $\partial_{\nu}(\bar{q}\lambda_b\sigma^{\mu\nu}q)$  leads to  $\partial_{\nu}(\bar{q}\lambda_b\gamma_5\sigma^{\mu\nu}q)$
- $\hookrightarrow$  quantum numbers of  $b_1$  multiplet cf. Caldi/Pagels, Phys.Rev.D14, 809, 1976
  - N.B. even more complicated for baryons



- problem: corresponding quark current not unique
- e.g.  $\bar{q}\lambda_b\gamma^{\mu}q$  and  $\partial_{\nu}(\bar{q}\lambda_b\sigma^{\mu\nu}q)$  have quantum numbers of  $\rho$  multiplet

- chiral transformation on  $ar q\lambda_b\gamma^\mu q$  leads to  $ar q\lambda_c\gamma_5\gamma^\mu q$
- $\hookrightarrow$  quantum numbers of  $a_1$  multiplet
  - chiral transformation on  $\partial_{\nu}(\bar{q}\lambda_b\sigma^{\mu\nu}q)$  leads to  $\partial_{\nu}(\bar{q}\lambda_b\gamma_5\sigma^{\mu\nu}q)$
- $\hookrightarrow$  quantum numbers of  $b_1$  multiplet cf. Caldi/Pagels, Phys.Rev.D14, 809, 1976
  - N.B. even more complicated for baryons

in our framework: a1 and b1 dynamically generated

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### Example: chiral partner(s) of $\rho$ meson

dynamical generation of  $b_1$  (parameter free)



Lutz/Kolomeitsev, Nucl. Phys. A 730, 392 (2004)

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#### Pion form factor with final-state interaction only

#### parameters: renormalization points

• for loop for transition from photon to hadrons



- $\hookrightarrow$  renormalization point should be in reasonable range
  - for loop in Bethe-Salpeter equation (rescattering, final-state interaction)



 ← renormalization point fixed
 (cf. Lutz/Kolomeitsev, Nucl. Phys. A 730, 392 (2004);
 Hyodo/Jido/Hosaka, arXiv:0803.2550 [nucl-th])



#### Pion form factor with final-state interaction only



- resonance only for renormalization points in TeV range (same finding: Oller/Oset, Phys. Rev. D 60, 074023 (1999))
- no resonance for reasonable renormalization points