QCD phase diagram from effective models

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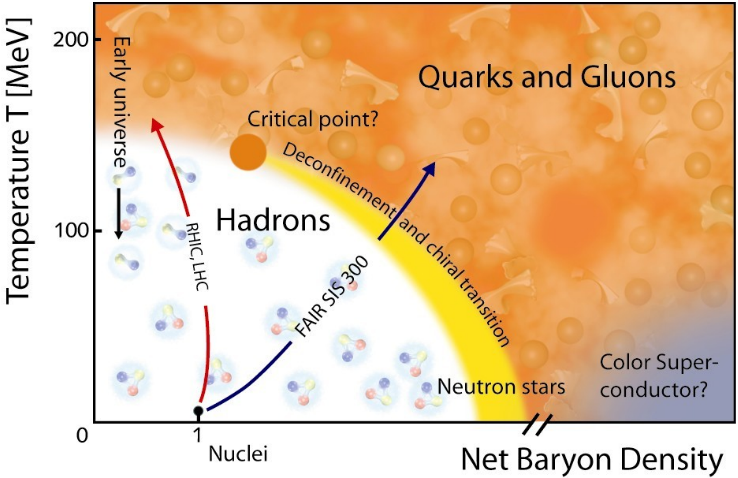
International School of Nuclear Physics

30th Course

Heavy-Ion Collisions from the Coulomb Barrier to the Quark-Gluon Plasma

Erice, Sicily

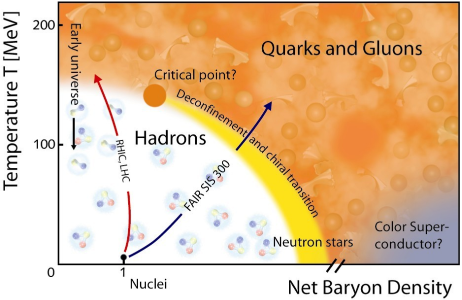
The conjectured QCD Phase Diagram



FAIR, Darmstadt

QCD: two phase transitions:

FAIR, Darmstadt



1. restoration of chiral symmetry

   $SU_{L+R}(N_f) \to SU_L(N_f) \times SU_R(N_f)$

   order parameter:

   $$\langle \bar{q}q \rangle \begin{cases} > 0 \Leftrightarrow \text{symmetry broken, } T < T_c \\ = 0 \Leftrightarrow \text{symmetric phase, } T > T_c \end{cases}$$

   associate limit: $m_q \to 0$

   chiral transition: spontaneous restoration of global $SU_L(N_f) \times SU_R(N_f)$ at high T

QCD Phase Transitions

QCD: two phase transitions:

1. restoration of chiral symmetry
2. de/confinement (center symmetry)

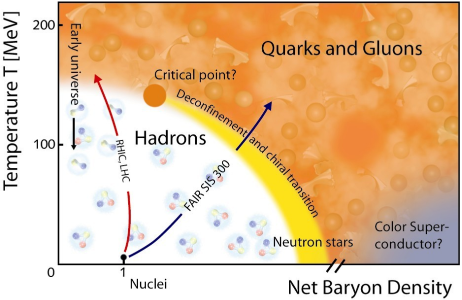
order parameter:

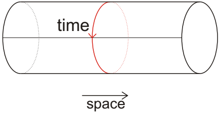
$$\Phi \begin{cases} = 0 \Leftrightarrow \text{confined phase}, & T < T_c \\ > 0 \Leftrightarrow \text{deconfined phase}, & T > T_c \end{cases}$$

$$\Phi = \frac{1}{N_c} \langle \mathrm{tr}_c \mathcal{P} e^{i \int_0^\beta d\tau A_0(\tau, \vec{x})} \rangle$$

associate limit: $m_q \to \infty$

FAIR, Darmstadt





➜ related to free energy of a static quark state: $\Phi = e^{-F_q}$

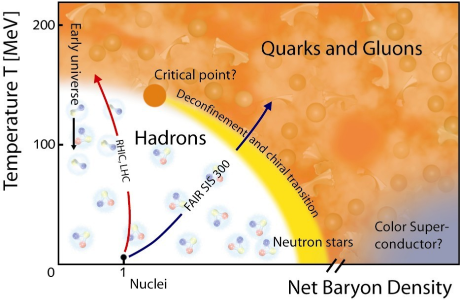
QCD: two phase transitions:

1. restoration of chiral symmetry
2. de/confinement

At densities/temperatures of interest

only model calculations available

FAIR, Darmstadt



effective models:

1. Quark-meson model — or other models e.g. NJL
2. Polyakov–quark-meson model — or PNJL models

Outline

- Two-Flavor Quark-Meson Model
  - ▷ Mean field approximation
  - ▷ Renormalization Group study

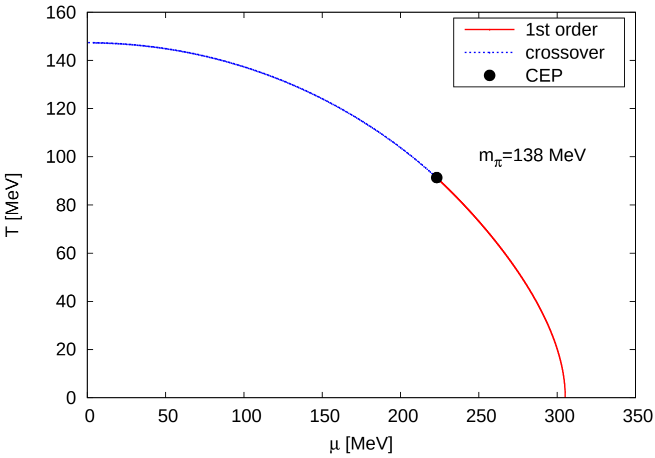
- Polyakov–Quark-Meson Model

- Three-Flavor Quark-Meson Model

- Lagrangian:

$$\mathcal{L}_{\text{qm}} = \bar{q}[i\gamma_\mu\partial^\mu - g(\sigma + i\vec{\tau}\vec{\pi}\gamma_5)]q + \frac{1}{2}(\partial_\mu\sigma)^2 + \frac{1}{2}(\partial_\mu\vec{\pi})^2 + \frac{\lambda}{4}(\sigma^2 + \vec{\pi}^2 - v^2)^2 - c\sigma$$

- Mean field analysis



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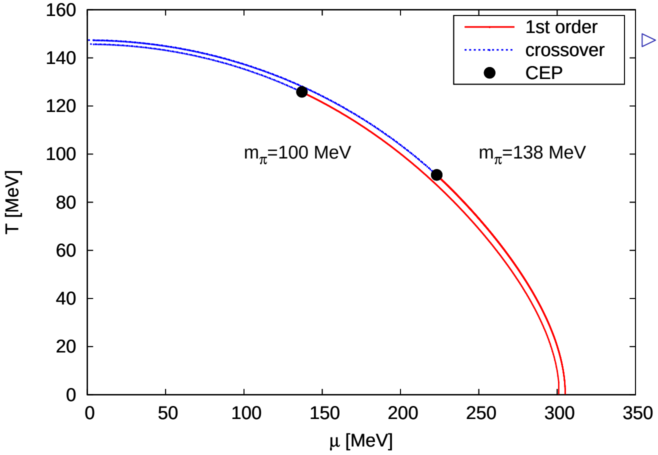
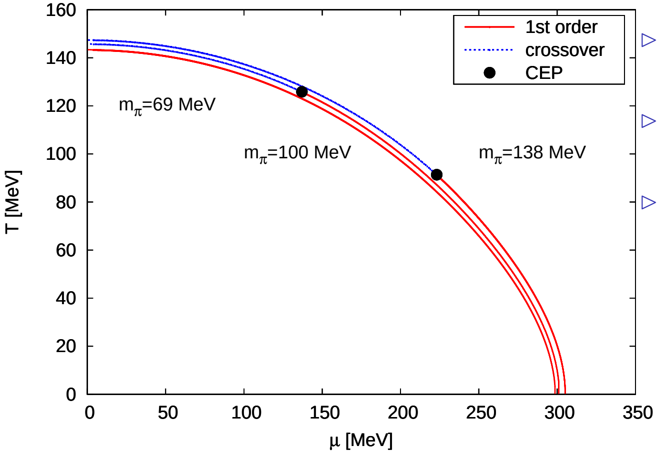


▷ as m_q decreased
CEP → T-axis

- Lagrangian:

$$\mathcal{L}_{\text{qm}} = \bar{q}[i\gamma_\mu\partial^\mu - g(\sigma + i\vec{\tau}\vec{\pi}\gamma_5)]q + \frac{1}{2}(\partial_\mu\sigma)^2 + \frac{1}{2}(\partial_\mu\vec{\pi})^2 + \frac{\lambda}{4}(\sigma^2+\vec{\pi}^2-v^2)^2 - c\sigma$$

- Mean field analysis



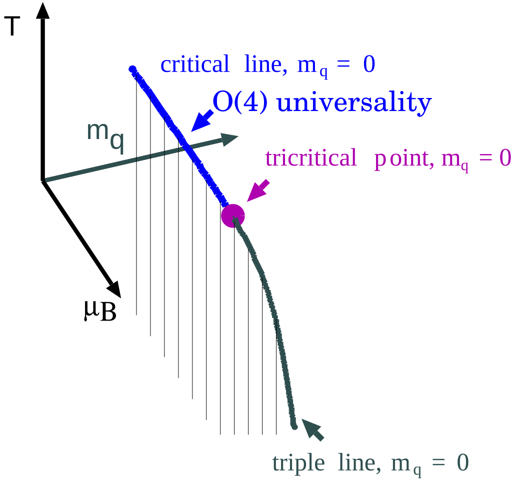
▷ as m_q decreased
CEP $\rightarrow$ T-axis

▷ chiral limit
no CEP

▷ at $\mu = 0$
$O(4)$ scaling expected
i.e. 2nd-order PT

$\rightarrow$ truncation effect

Chiral limit: $(m_q = 0)$ $SU(2) \times SU(2) \sim O(4)$-symmetry $\longrightarrow$ 4 modes critical $\sigma, \vec{\pi}$



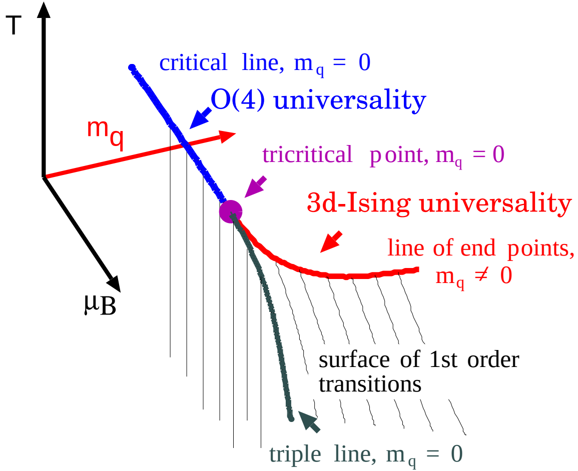
General properties

- chiral limit
  tricritical point (Gaussian fixed point)

Phase diagram in (T, μ_B, m_q)-space

Chiral limit: $(m_q = 0)$ $SU(2) \times SU(2) \sim O(4)$-symmetry $\longrightarrow$ 4 modes critical $\sigma, \vec{\pi}$

$m_q \neq 0$: no symmetry remains $\longrightarrow$ only one critical mode σ (Ising) ($\vec{\pi}$ massive)



General properties

- chiral limit
  tricritical point
  (Gaussian fixed point)
- finite m_q
  critical endpoints
  (3D-Ising class)

$\Gamma_k[\phi]$ scale dependent effective action ; $t = \ln(k/\Lambda)$; R_k regulators

FRG (average effective action) [Wetterich]

$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \mathrm{Tr}\, \partial_t R_k \left(\frac{1}{\Gamma_k^{(2)} + R_k} \right) \quad ; \quad \Gamma_k^{(2)} = \frac{\delta^2 \Gamma_k}{\delta\phi\delta\phi}$$

- Ansatz for Γ_k: (LO derivative expansion → arbitrary potential V_k)

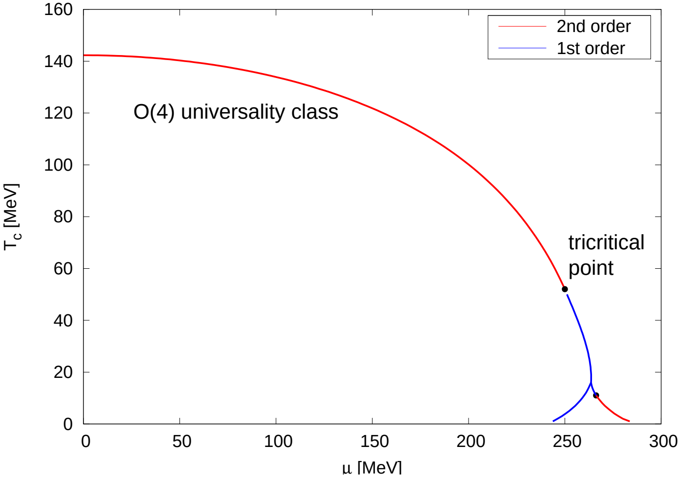
$$\Gamma_k = \int d^4x \bar{q}[i\gamma_\mu \partial^\mu - g(\sigma + i\vec{\tau}\vec{\pi}\gamma_5)]q + \frac{1}{2}(\partial_\mu \sigma)^2 + \frac{1}{2}(\partial_\mu \vec{\pi})^2 + V_k(\phi^2)$$

$$V_{k=\Lambda}(\phi^2) = \tfrac{\lambda}{4}(\sigma^2 + \vec{\pi}^2 - v^2)^2 - c\sigma$$

[BJS, J. Wambach, '05 & '06]

RG analysis:

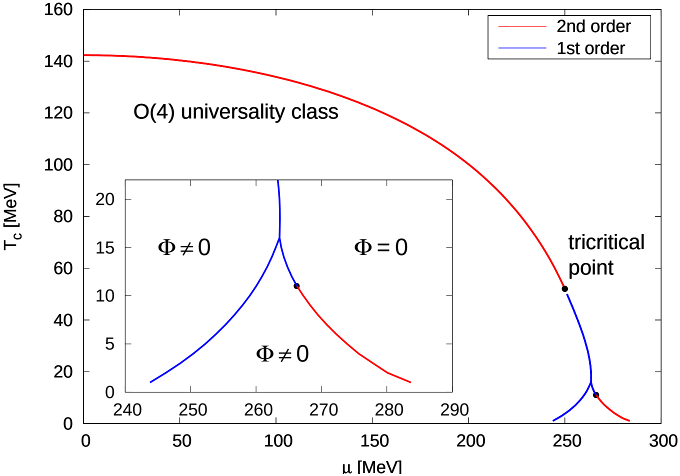
$O(4) \sim SU(2) \times SU(2)$ chiral limit



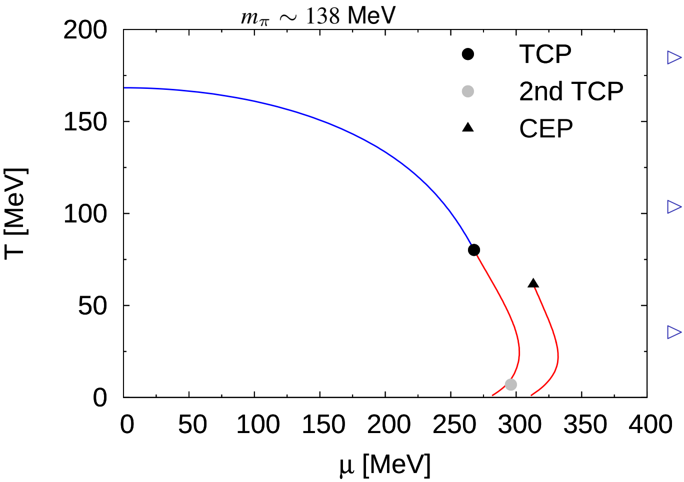
[BJS, J. Wambach, '05 & '06]

RG analysis:

$O(4) \sim SU(2) \times SU(2)$ chiral limit no spinodal lines!



RG Phase Diagram



- bending usual for RG

  Clausius-Clapeyron relation ok

- 2nd tricritical point

  in chiral limit

- features

  parameter independent

  but locations TCP/CEP

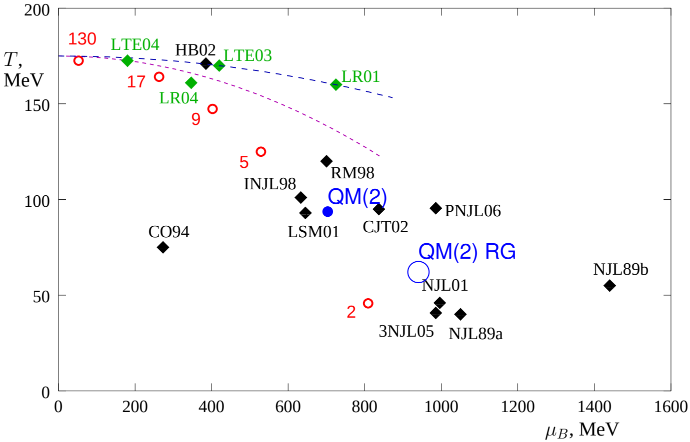
  parameter dependent

[BJS,Wambach ’05 & ’06]

Charts of QCD Critical End Points

model studies vs. lattice simulations

Black points: models   Lines & green points: lattice   Red points: Freezeout points for HIC



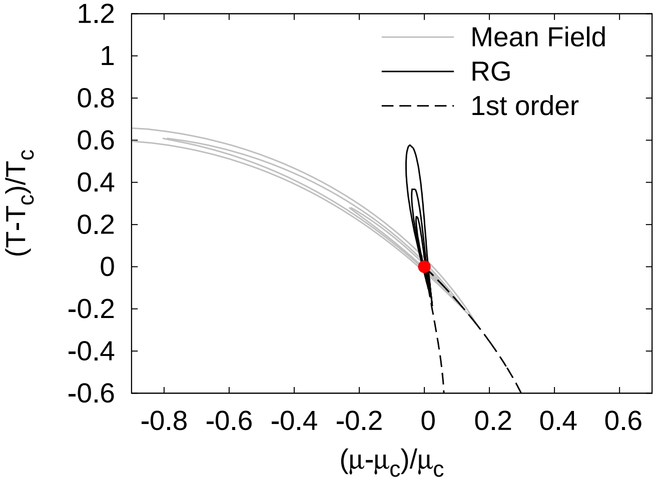
Stephanov '05 & '07

lattice methods:

- reweighting
- imaginary μ_B
- Taylor expansion around $\mu_B = 0$

Comparison with scalar χ_σ : MF↔RG

[BJS, J. Wambach '06]



Outline

- Two-Flavor Quark-Meson Model
  - ▷ Mean field approximation
  - ▷ Renormalization Group study
- Polyakov–Quark-Meson Model
- Three-Flavor Quark-Meson Model

- Lagrangian $\mathcal{L}_{\text{PQM}} = \mathcal{L}_{\text{qm}} + \mathcal{L}_{\text{pol}}$
- Polyakov loop potential:

Polyakov 1978
Meisinger 1996
Pisarski 2000

$$\mathcal{L}_{\text{pol}} = -\bar{q}\gamma_0 A_0 q - \mathcal{U}(\phi, \bar{\phi})$$

$$\frac{\mathcal{U}(\phi,\bar{\phi})}{T^4} = -\frac{b_2(T,T_0)}{2}\phi\bar{\phi} - \frac{b_3}{6}\left(\phi^3 + \bar{\phi}^3\right) + \frac{b_4}{16}\left(\phi\bar{\phi}\right)^2$$

Fukushima 2003
Ratti, Weise et al. 2004
Dumitru, Pisarski 2004
Friman, Redlich, Sasaki 2006

$\Rightarrow$ first-order transition at $T_0 = 270$ MeV

in presence of dynamical quarks: $T_0 = T_0(N_f)$

BJS, Pawlowski, Wambach, 2007

N_f	0	1	2	$2+1$	3
T_0 [MeV]	270	240	208	187	178

- Lagrangian $\mathcal{L}_{\mathrm{PQM}} = \mathcal{L}_{\mathrm{qm}} + \mathcal{L}_{\mathrm{pol}}$
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Friman, Redlich, Sasaki 2006

$\Rightarrow$ first-order transition at $T_0 = 270$ MeV

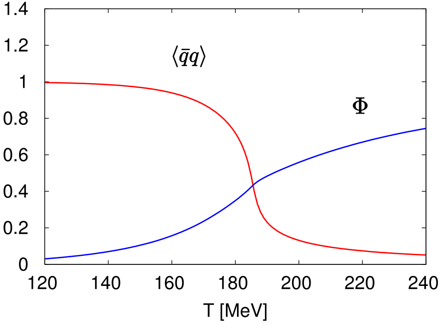
$\mu \neq 0: \quad T_0 = T_0(N_f, \mu)$ BJS, Pawlowski, Wambach, 2007

$\bar{\phi} \neq \phi^*$

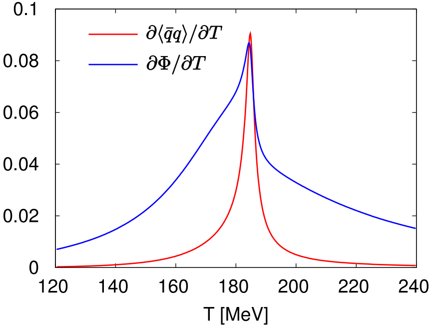
[BJS, Pawlowski, Wambach '07]

Numerical results:

order parameters



T-derivatives of order parameters

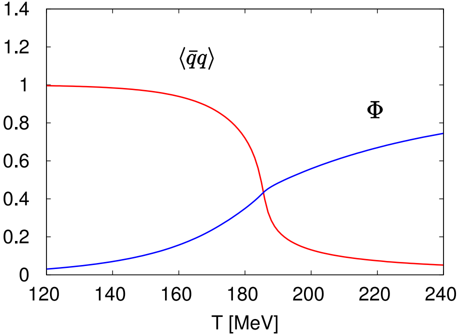


Finite temperature and $\mu = 0$

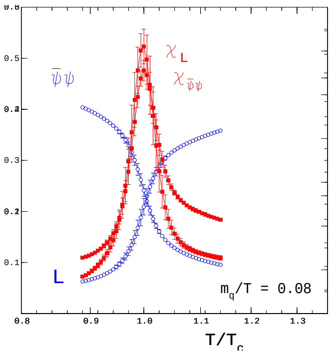
[Karsch et al.]

Numerical results:

order parameters



lattice

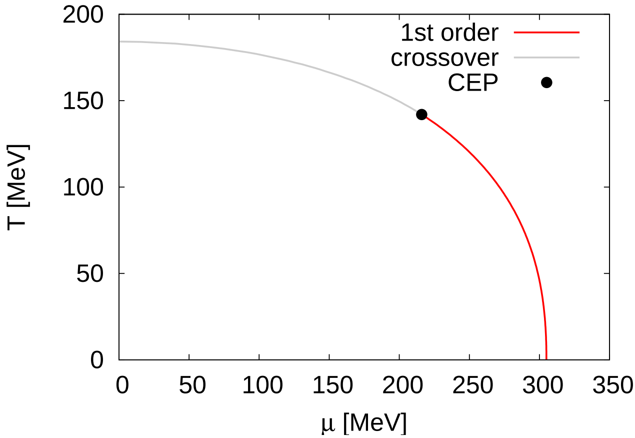


[BJS, Pawlowski, Wambach '07]

in mean field approximation

chiral transition and 'deconfinement' coincide

- for PQM model



[BJS, Pawlowski, Wambach '07]

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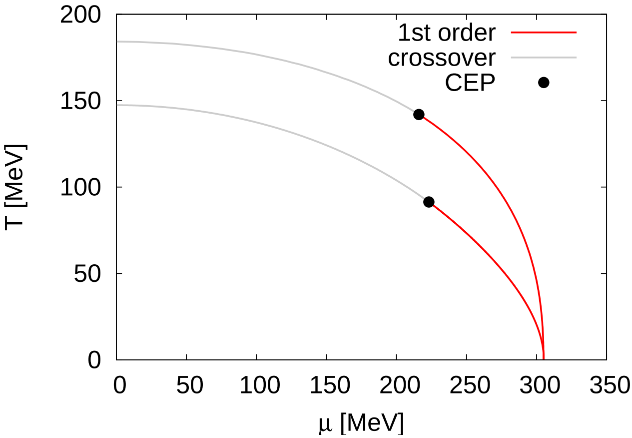
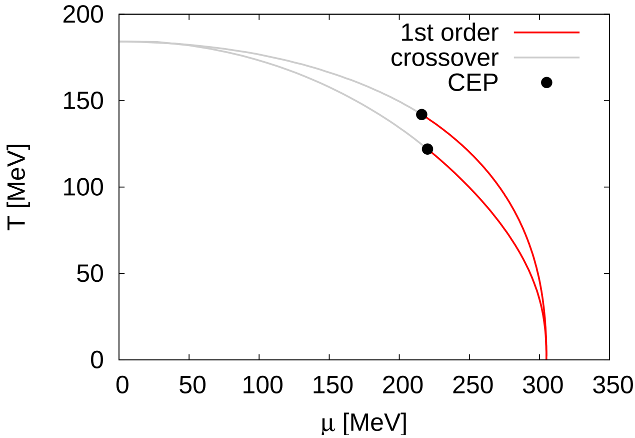


- for PQM model
- for QM model
  (lower lines)

[BJS, Pawlowski, Wambach '07]

in mean field approximation

chiral transition and 'deconfinement' coincide



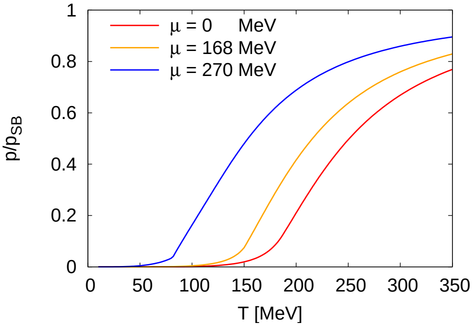
- for PQM model
- for PQM model with μ-modification in Polyakov loop potential (lower lines)

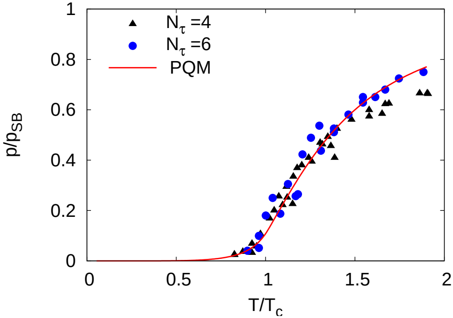
Pressure

- perturbative pressure of QCD with N_f massless quarks

$$\frac{p}{T^4} = (N_c^2 - 1)\frac{\pi^2}{45} + N_f\left[\frac{7\pi^2}{60} + \frac{1}{2}\left(\frac{\mu}{T}\right)^2 + \frac{1}{4\pi^2}\left(\frac{\mu}{T}\right)^4\right] .$$

- $N_f = 2$: lattice: $N_\tau = 4, 6$; $\mu = 0$





[Ali Khan et al. '01]

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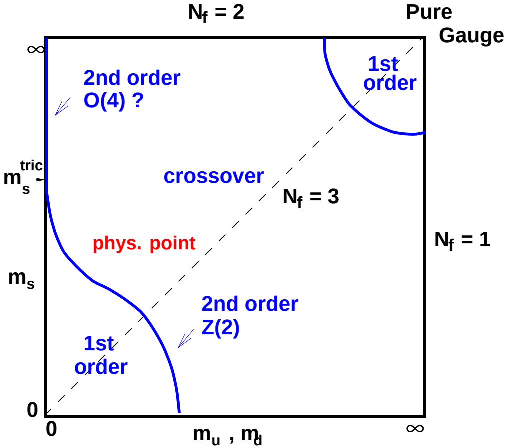
Columbia plot:

[Brown et al. '90]

$T_\chi^{N_f=2} \sim 175$ MeV

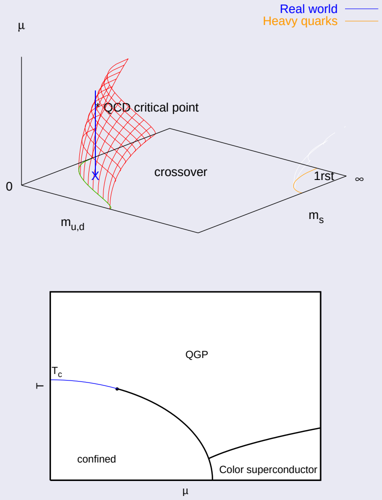
$T_d \sim 270$ MeV



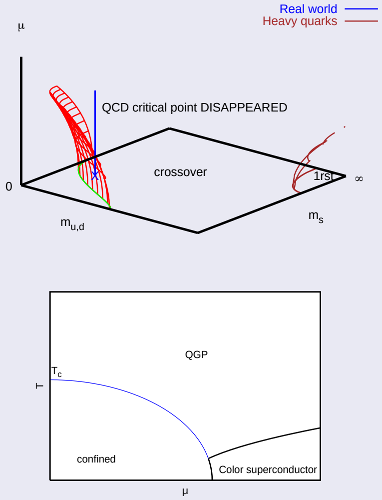
$T_\chi^{N_f=3} \sim 155$ MeV

Mass Sensitivity (lattice, $N_f = 3, \mu_B \neq 0$)

Standard scenario: $m_c(\mu)$ increasing



Nonstandard scenario: $m_c(\mu)$ decr.



[de Forcrand, Philipsen: hep-lat/0611027]

$N_f = 3$ quark-meson model

- Lagrangian $\mathcal{L} = \mathcal{L}_{\text{quark}} + \mathcal{L}_{\text{meson}}$

$$\mathcal{L}_{\text{quark}} = \bar{q}(i\partial\!\!\!/ - g\frac{\lambda_a}{2}(\sigma_a + i\gamma_5\pi_a))q$$

$$\begin{aligned}\mathcal{L}_{\text{meson}} = & \, tr(\partial_\mu\phi^\dagger\partial^\mu\phi) - m^2 tr(\phi^\dagger\phi) - \lambda_1(tr(\phi^\dagger\phi))^2 \\ & -\lambda_2 tr((\phi^\dagger\phi)^2) + c\,(\det(\phi) + \det(\phi^\dagger)) \\ & + trH(\phi + \phi^\dagger)\end{aligned}$$

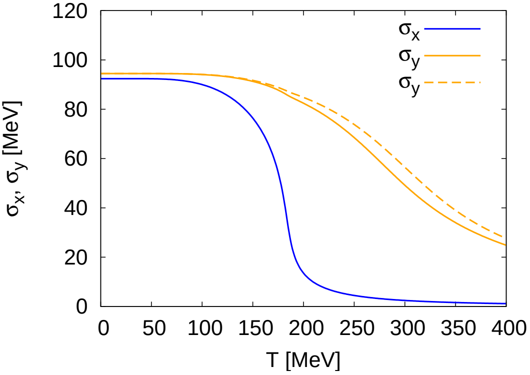
fields: $\phi = \sum_a \frac{\lambda_a}{2}(\sigma_a + i\pi_a)$ and $H = \sum_a \frac{\lambda_a}{2} h_a$

σ_a scalar and π_a pseudoscalar nonet

[BJS, M. Wagner, arXiv:0808.1491]

➜ two condesates: nonstrange $\sigma_x(T, \mu_f)$ and strange $\sigma_y(T, \mu_f)$

with (solid) and without (dashed) $U(1)_A$ anomaly



▷ almost no effect due $U(1)_A$ anomaly

▷ T_c depends on m_σ

Chiral symmetry restoration

[BJS, M. Wagner, arXiv:0808.1491]

➜ two condesates: nonstrange $\sigma_x(T, \mu_f)$ and strange $\sigma_y(T, \mu_f)$

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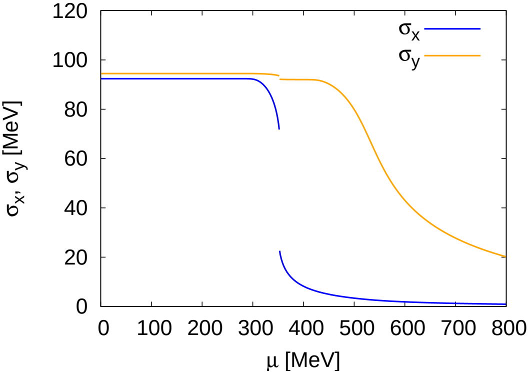


- ▷ almost no effect due $U(1)_A$ anomaly
- ▷ T_c depends on m_σ
- ▷ $\mu \equiv \mu_q = \mu_s$
- ▷ μ_c depends on m_σ

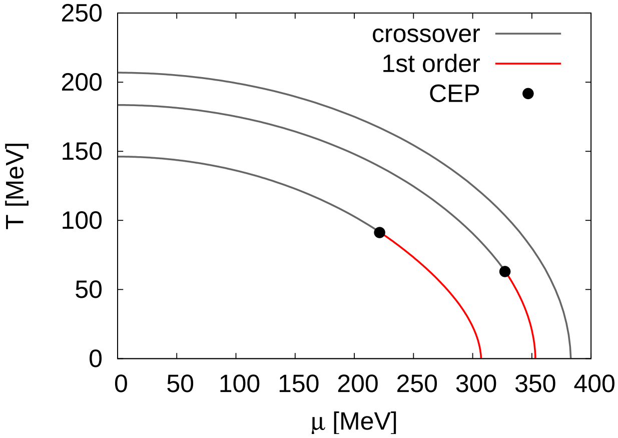
Phase diagram $N_f = 3$ $\quad (\mu \equiv \mu_q = \mu_s)$

[BJS, M. Wagner, arXiv:0808.1491]

$m_\sigma = 600$ MeV (lower lines), 800 and 900 MeV

PDG: $f_0(600)$ mass=$(400 \ldots 1200)$ MeV

➜ influence of existence of CEP!



In-medium meson masses

[BJS, M. Wagner, arXiv:0808.1491]

- ▷ genuine problem of linear sigma model w/o quarks at finite T

  $\rightarrow$ negative meson masses

- ▷ but not in this approximation

  ➜Ward identities, Goldstone theorem etc. are all valid in-medium e.g.

$$h_x = f_\pi m_\pi^2$$

  similar in strange sector

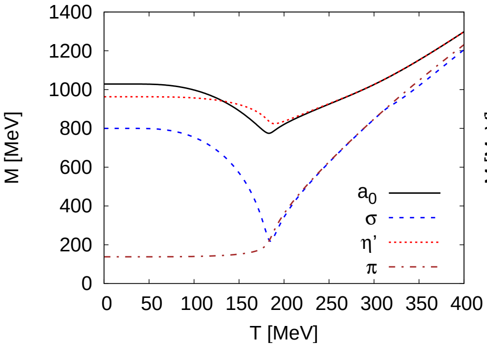
- ▷ At low temperatures: mesons dominate

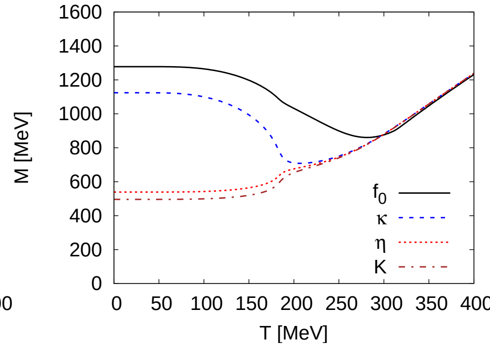
  At high temperatures: quarks dominate

In-medium meson masses

[BJS, M. Wagner, arXiv:0808.1491]

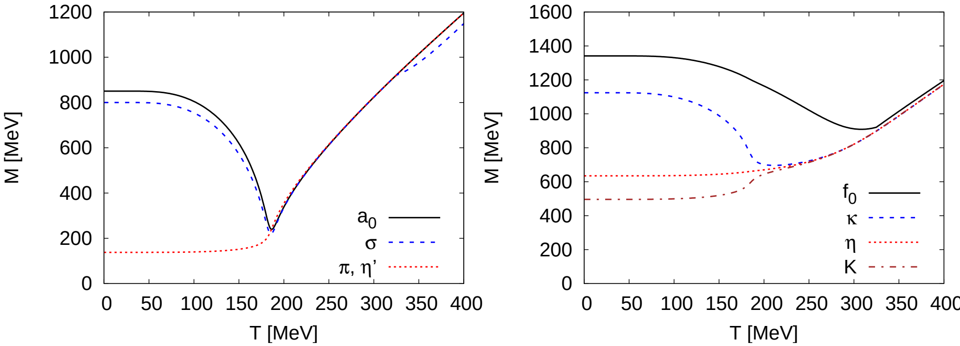
masses with $U(1)_A$ anomaly





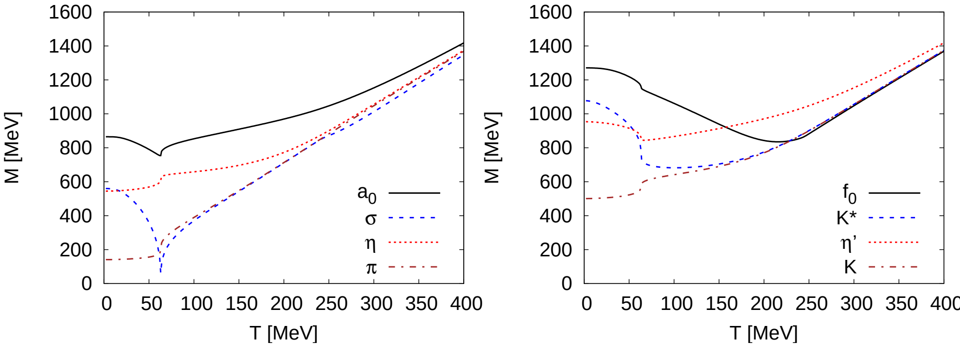
[BJS, M. Wagner, arXiv:0808.1491]

masses without $U(1)_A$ anomaly



[BJS, M. Wagner, arXiv:0808.1491]

masses with $U(1)_A$ anomaly

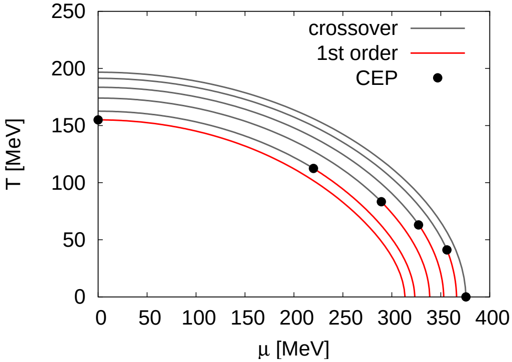


Mass sensitivity

RG arguments predict for $N_f = 3$ 1st-order in chiral limit

▷ $m_\pi / m_\pi^* = 0.49$ (lower line), $0.6, 0.8 \ldots, 1.36$ (upper line) $m_\pi^* = 138$ MeV

with $U(1)_A$, $m_\sigma = 800$ MeV

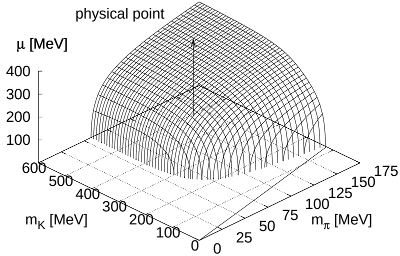


Chiral critical surface $(m_\sigma = 800$ MeV)

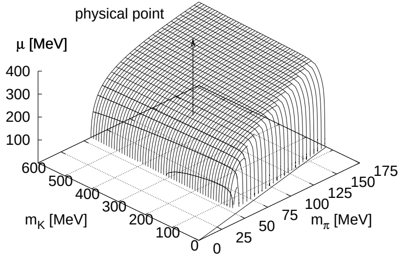
[BJS, M. Wagner, arXiv:0808.1491]

- chiral critical surface in (m_π, m_K) plane

with $U(1)_A$



without $U(1)_A$



Chiral critical surface for different m_σ

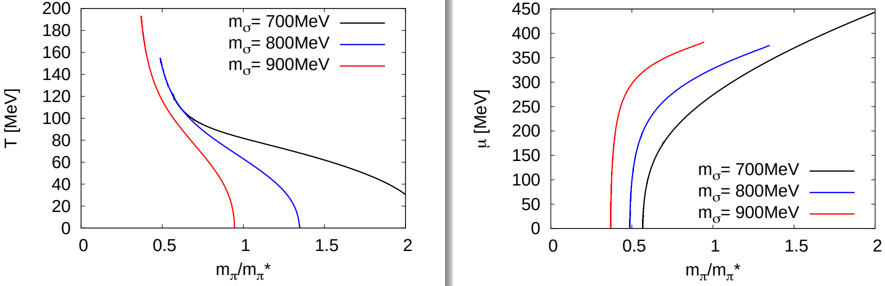
[BJS, M. Wagner, arXiv:0808.1491]

- chiral critical surface in (m_π, m_K) plane

➜ cuts along fixed m_π/m_K ratio through physical point

$m_\pi^* = 138$ MeV, $m_K^* = 496$MeV (physical point)

critical T_c / critical μ



Summary

Quark-meson model study for $N_F = 2$

➜ Mean field versus RG

Influence of fluctuations on phase diagram

Findings:

- ▷ MF phase diagram: no TCP (in chiral limit) found
- ▷ RG phase diagram: two TCP's (in chiral limit) & CEP found
- ▷ Size of critical region via susceptibilities: "compressed" with fluctuations

Quark-meson model study for $N_F = 3$

➜ Mean-field approximation

no need for Optimized Perturbation Theory

with and without axial anomaly

Polyakov–quark-meson model study for $N_F = 2$

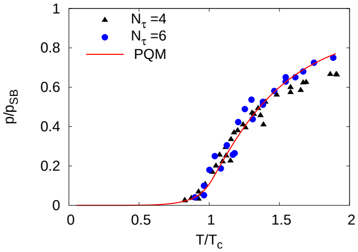
➜ mean-field approximation

Findings:

- ▷ Parameter in Polyakov loop potential: $T_0 \Rightarrow T_0(N_f, \mu)$
  - pure gauge: $T_0 \sim 270$ MeV
  - $N_f = 2$: $T_0 \sim 210$ MeV
- ▷ Chiral & deconfinement transition coincide
- ▷ Mean-field approximation encouraging

  Quark-meson model is renormalizable

  → no UV cutoff parameter (cf. PNJL model)



Outlook

- ▷ include quark-meson dynamics in PQM model and for $N_f = 3$ with FRG
- ▷ include glue dynamics with FRG → full QCD

  (step by step)