

# Equation of state for strongly interacting matter within a HTL quasiparticle model

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- QPM with  $Im \Pi_i \neq 0$ , plasmons and plasminos from HTL
- extrapolation of lattice QCD to large baryon densities, e.g. CBM@FAIR
- EOS for hydrodynamic phase of heavy-ion collisions (SPS, RHIC, LHC)

## From QCD to thermodynamics



### CJT formalism

• effective action

$$\Gamma[D, S] = I - \frac{1}{2} \left\{ \operatorname{Tr} \left[ \ln D^{-1} \right] + \operatorname{Tr} \left[ D_0^{-1} D - 1 \right] \right\} + \left\{ \operatorname{Tr} \left[ \ln S^{-1} \right] + \operatorname{Tr} \left[ S_0^{-1} S - 1 \right] \right\} + \Gamma_2[D, S]$$

• translation-invariant systems, no broken symmetries

$$\frac{\Omega}{V} = \operatorname{tr} \int \frac{\mathrm{d}^4 k}{(2\pi)^4} n_{\mathrm{B}}(\omega) \operatorname{Im} \left( \ln D^{-1} - \Pi D \right) + 2 \operatorname{tr} \int \frac{\mathrm{d}^4 k}{(2\pi)^4} n_{\mathrm{F}}(\omega) \operatorname{Im} \left( \ln S^{-1} - \Sigma S \right) - \frac{T}{V} \Gamma_2$$

### 2-loop QCD thermodynamics

• truncate  $\Gamma_2$  at 2-loop order

$$\Gamma_2 = \frac{1}{12} + \frac{1}{8} \left( - \frac{1}{2} \right) - \frac{1}{2} \left( - \frac{1}{2} \right)$$

 $\rightarrow$  self-energies of 1-loop order



• gauge invariance: HTL self-energies instead

### HTL self-energies

•  $Im \Pi \neq 0$  below the lightcone (solid lines)



### Entropy

• stationarity of  $\Omega$   $s := -\frac{1}{V} \left. \frac{\partial \Omega}{\partial T} \right|_{\mu} = -\frac{1}{V} \left( \left. \frac{\partial \Omega}{\partial T} \right|_{\text{expl.}} + \underbrace{\frac{\delta \Omega}{\delta D}}_{0} \left. \frac{\partial D}{\partial T} \right)_{\mu}$   $= s_{g,\text{T}} + s_{g,\text{L}} + s_{q,\text{Pt}} + s_{q,\text{Pl}} + s' \qquad s' = 0$ • e.g. gluons: • e.g. gluons:

$$s_{g,T} \sim \int_{d^{4}k} \frac{\partial n_{B}}{\partial T} \left\{ \underbrace{\pi \varepsilon(\omega) \Theta(-\text{Re}D_{T}^{-1})}_{\text{qp contribution}} + \underbrace{\text{Re}D_{T}\text{Im}\Pi_{T} - \arctan(\frac{\text{Im}\Pi_{T}}{\text{Re}D_{T}^{-1}})}_{\text{damping terms}} \right\}$$
• HTL QPM model
$$\iff \underbrace{\text{effective QPM}}_{\text{coll. modes, damping, asympt. disp. rel.}}_{\text{Blaizot, lancu, Rebhan: PRD'01}} \bigotimes \underbrace{\text{Reban, Romatschke: PRD'03}}_{\text{Bluhm, Kämpfer, RS, Seipt: EPJC'07}}$$

### Pressure

$$\begin{split} s &\sim \Big(\frac{\partial \Omega}{\partial T}\bigg|_{\exp \mathbf{l}} + \frac{\delta \Omega}{\delta D} \frac{\partial D}{\partial T}\Big)_{\mu} \\ s_i &\sim \int_{\mathrm{d}^4 k} \frac{\partial n_{B/F}}{\partial T} \left\{ \mathrm{qp} + \mathrm{damping} \right\} \end{split}$$

• self-consistent formulation of the pressure

$$p = -\frac{\Omega}{V} := \sum_{i} p_{i} - B \qquad p_{i} \sim \int_{\mathrm{d}^{4}k} n_{\mathrm{B/F}} \left\{ \mathrm{qp} + \mathrm{damping} \right\}$$
$$\frac{\partial B}{\partial T} := \sum_{i} \frac{\partial p_{i}}{\partial \Pi_{i}} \frac{\partial \Pi_{i}}{\partial T} \qquad \left( \frac{\partial B}{\partial \mu} = \sum_{i} \frac{\dots}{\dots} \frac{\partial \mu}{\partial \mu} \right)$$

• entropy density

$$s = \frac{\partial p}{\partial T} = \sum_i s_i + \frac{\partial p_i}{\partial \Pi_i} \frac{\partial \Pi_i}{\partial T} - \frac{\partial B}{\partial T} = \sum_i s_i$$

• net quark density

$$n\sim \frac{\partial\Omega}{\partial\mu}\bigg|_{\mathrm{expl.}} +\underbrace{\frac{\delta\Omega}{\delta D}}_{0}\frac{\partial D}{\partial\mu}$$

$$n_q = \frac{\partial p}{\partial \mu} \sim \int_{\mathrm{d}^4 k} \left( \frac{\partial n_{\mathrm{F}}}{\partial \mu} + \frac{\partial n_{\mathrm{F}}^A}{\partial \mu} \right) \left\{ \mathrm{qp} + \mathrm{damping} \right\}$$

### Effective coupling

• fundamental parameter

$$g^{2}(x^{2}) = \frac{16\pi^{2}}{\beta_{0}\ln(x^{2})} \left(1 - \frac{4\beta_{1}}{\beta_{0}^{2}} \frac{\ln[\ln(x^{2})]}{\ln(x^{2})}\right)$$

• running coupling  $g^2$   $x = \frac{\bar{\mu}}{\Lambda_{\rm QCD}}$   $\bar{\mu} \sim T$ 

$$T > T_c, \mu = 0$$

• effective coupling  $G^2 \qquad x = \frac{\lambda}{T_c}(T - T_s)$ 

### Adjustment @ $\mu\!=\!0$

• 
$$\mu = 0$$
: adjust to  $\ell \text{QCE}$   
 $\rightarrow T_s, \lambda \text{ fixed}$   
 $\rightarrow G^2(T, \mu = 0)$ 





### Influence of coll. modes + LD @ $\mu = 0$

individual entropy contributions



• Landau damping large close to  $T_{c}$ , decreases for higher temperatures



### Into the T- $\mu$ -plane

•  $\mu > 0$ : stationary potential, self-consistent model  $\rightarrow$  impose Maxwell's relation

$$\frac{\partial s}{\partial \mu} = \frac{\partial n}{\partial T} \qquad \longrightarrow \qquad a_T \frac{\partial G^2}{\partial T} + a_\mu \frac{\partial G^2}{\partial \mu} = b$$

Peshier, Kämpfer, Soff: PRC'00, PRD'02

- solve quasilinear PDE for  $G^2(T, \mu \neq 0)$  using method of characteristics
- test with lattice data for  $\mu \simeq 0$ : successful (eQPM)

### Larger chemical potential



• full HTL QPM: stable characteristics

entropy density





RS, Bluhm, Kämpfer: EPJ ST'08

# Results (1)

• entropy density



RS, Bluhm, Kämpfer: arXiv:0803.1571

• increases with temperature and chemical potential

# Results (2)

• net quark density



• small area of negative net quark density below transition line

# Results (3)

• pressure also increases with T and  $\mu$ 



- small area of negative pressure also above transition line
  - $\rightarrow$  no problems for EoS @ RHIC, LHC, SPS, FAIR

### EOS for RHIC and LHC



• energy density

 $e = -p + sT + \mu n$ 



 $e \, [\text{GeV/fm}^3]$ 

• EOS for LHC, RHIC

 $n_{\rm b}/s \approx 0$ 

### EOS for SPS

PRELIMINARY



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### Summary & Outlook

- 2-loop  $\Gamma_2$  + eff. coupling  $G^2 \rightarrow \mathsf{HTL} \ \mathsf{QPM}$
- Landau damping + collective modes  $\rightarrow$  large  $\mu$  accessible
- limitation: negative pressure for small T and  $\mu$   $\rightarrow$  however deconfinement region @ LHC, RHIC, FAIR, SPS accessible

• outlook: EOS for FAIR, critical endpoint

Kämpfer, Bluhm, RS, Seipt: NPA'06