

Nuclear EoS and 3-body forces in the T-matrix approach

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Motivation

➤ describe nuclear matter starting from realistic bare nucleon-nucleon potentials (CD-Bonn, Nijmegen, Argonne)

★ thermodynamic (macroscopic) properties

★ single-particle (microscopic) properties



thermodynamic Green's functions

□ (Hot) dense matter \longrightarrow supernovae explosions ($\sim 4 \rho_0$) and neutron stars ($\sim 8 \rho_0$)



Real-time Green's functions

► single-particle propagator on the contour ($1 = \mathbf{r}, t$)

$$i \mathbf{G}(1, 1') = \langle \mathcal{T} \psi(1) \psi^\dagger(1') \rangle$$

✦ matrix representation

$$i \mathbf{G}(1, 1') = i \begin{pmatrix} G^c(1, 1') & G^<(1, 1') \\ G^>(1, 1') & G^a(1, 1') \end{pmatrix}$$

✦ spectral representation

$$i G^>(\mathbf{p}, \omega) = [1 - f(\omega)] A(\mathbf{p}, \omega)$$

$$-i G^<(\mathbf{p}, \omega) = f(\omega) A(\mathbf{p}, \omega)$$

► Dyson equation

$$\mathbf{G}(1, 1')^{-1} = \mathbf{G}_0(1, 1')^{-1} - \mathbf{\Sigma}(1, 1')$$

T-matrix approximation

Approximation for the two-particle propagator

$$\overleftrightarrow{G}_2 = \overleftrightarrow{G}_0 - \overleftrightarrow{G}_0 \overleftrightarrow{X} + \overleftrightarrow{G}_0 \overleftrightarrow{T} - \overleftrightarrow{G}_0 \overleftrightarrow{T} \overleftrightarrow{X}$$

where the in-medium two-particle scattering matrix T is introduced:

$$\overleftrightarrow{T} = \overleftrightarrow{V} + \overleftrightarrow{V} \overleftrightarrow{T}$$

Iterative scheme involving the calculation of

- T-matrix
- Self-energy
- Spectral function (via Dyson equation)

Thermodynamic consistency

A class of approximation can be derived from a *generating functional* $\Phi[G, V]$

$$\Sigma(x, y) = \frac{\delta\Phi}{\delta G(x, y)}$$

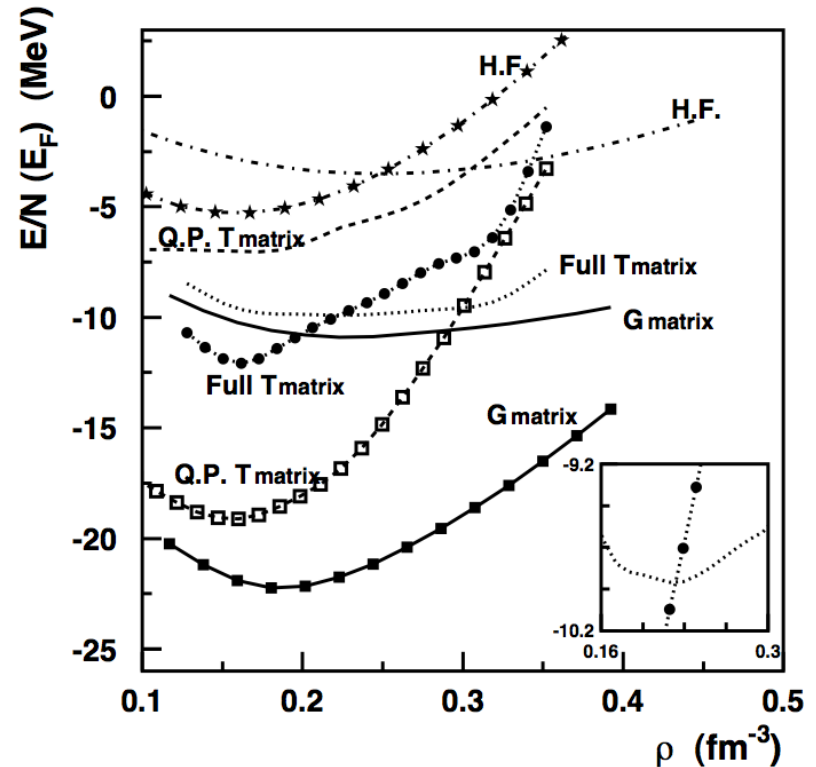
(Baym, Kadanoff 1961)

Example: pressure at $T = 0$

$$P = \rho^2 \frac{\partial(E/N)}{\partial\rho} = \rho(E_F - E/N)$$

→ Hugenholtz - Van Hove property at saturation

$$E_F = E/N$$



(P. Bożek and P. Czerski,
Eur.Phys.J. A11 (2001) 271)

Pressure and entropy

The pressure is related to the thermodynamic potential

$$\Omega(T, \mu, V) = -PV$$

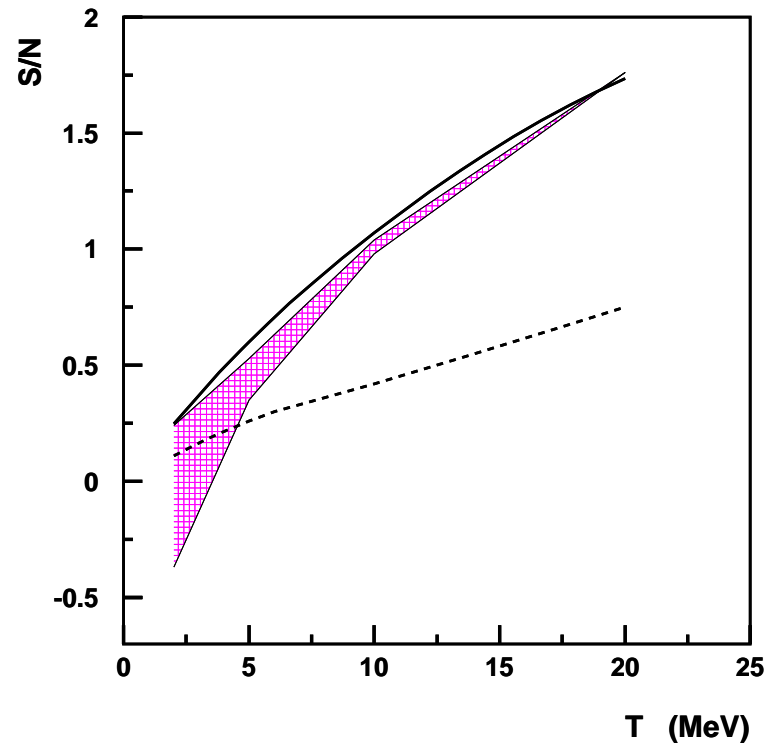
→ in a self-consistent approximation

$$\Omega = -\text{Tr}\{\ln[G^{-1}]\}$$

$$-\text{Tr}\{\Sigma G\} + \Phi$$

★ entropy

$$\frac{S}{N} = \frac{1}{T} \left[\frac{E}{N} + \frac{P_{tot}}{\rho} - \mu \right]$$

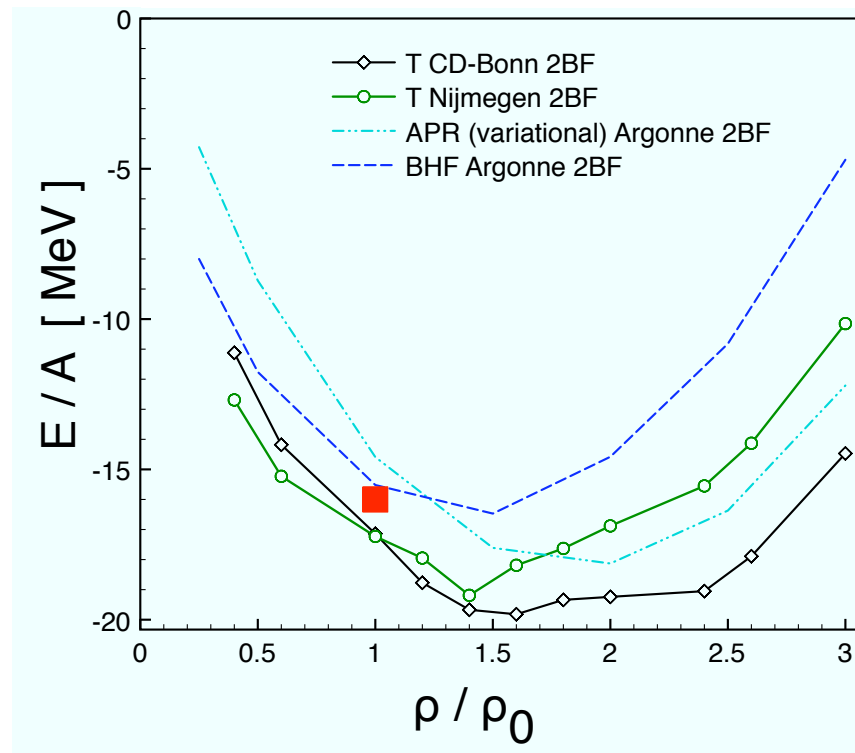


(V. S. and P. Božek, Phys.Rev. C74 (2006) 045809)

Three-body forces

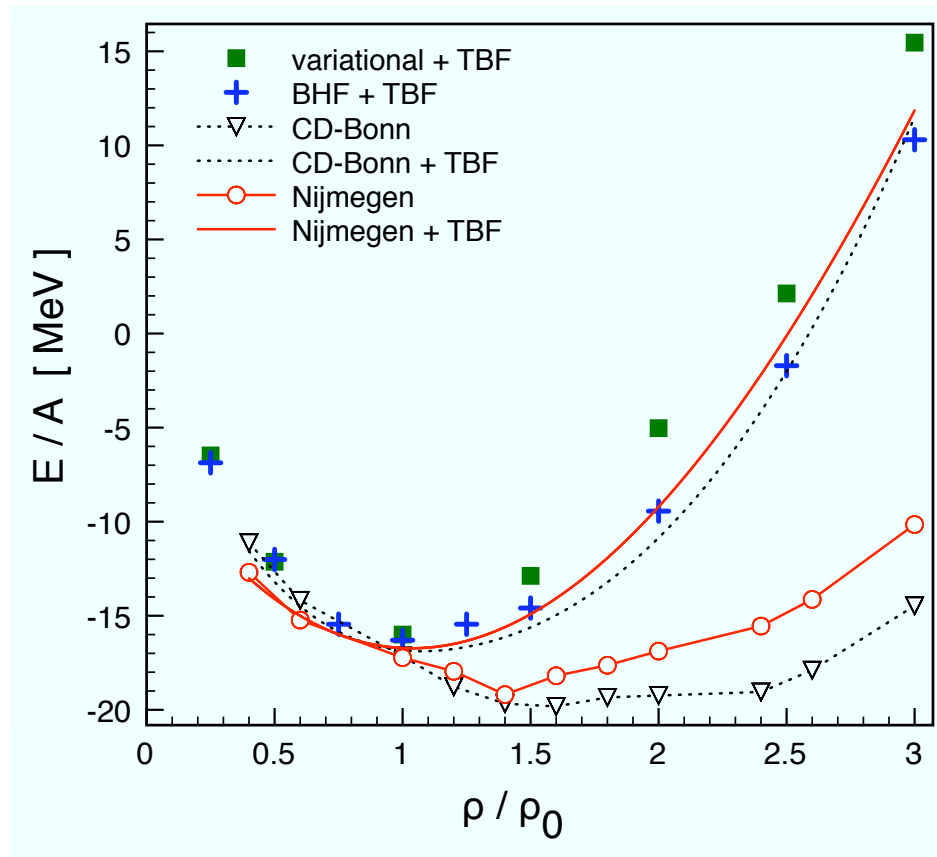
- Necessary to include three-body forces
- Urbana TBF

$$V_{ijk} = V_{ijk}^{2\pi} + V_{ijk}^R$$
- Two parameters ($A < 0$ and $U > 0$) to be tuned with the saturation properties
- Derive an effective two-particle potential



$$V_{eff}^3(\mathbf{q}, \mathbf{q}') = \sum_{\sigma\tau} \int \frac{d\mathbf{k}}{(2\pi)^3} n(\mathbf{k}) V^3(\mathbf{k}, \mathbf{q}, \mathbf{q}')$$

Energy per particle (symmetric matter)



$$\langle H_{\text{pot}} \rangle = \frac{1}{2} \sum_n \left[\text{Diagram 1} - \text{Diagram 2} \right]$$

$$= \frac{1}{2} \left[\text{Diagram 3} - \text{Diagram 4} \right]$$

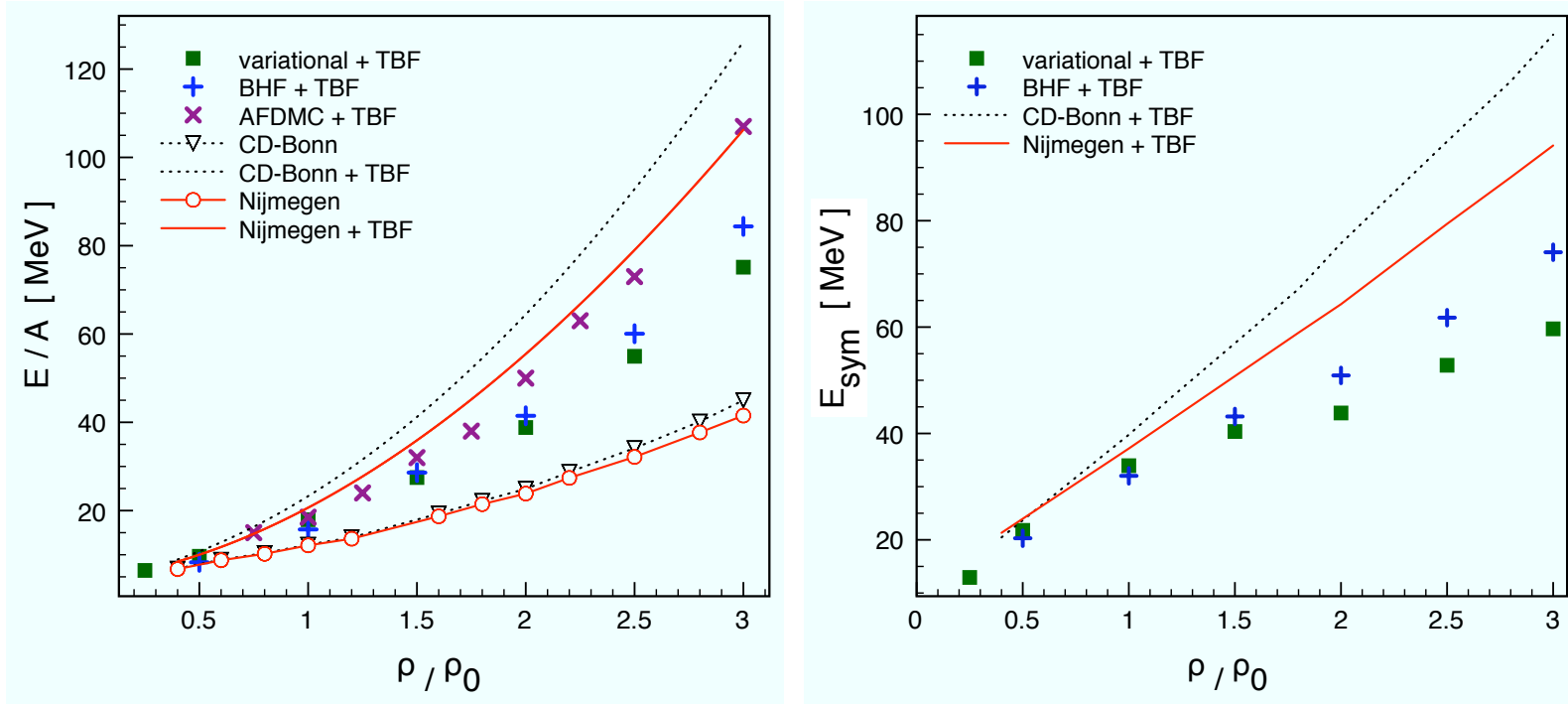
The diagrams represent various particle interactions: Diagram 1 and 2 are two-particle exchange diagrams with wavy lines; Diagram 3 and 4 are diagrams involving a box labeled 'T' representing a transition or interaction.

(V. S. and P. Božek, arXiv:0808.2929)

Alternatively through the Galitskii-Koltun sum rule

$$\frac{E}{N} = \frac{1}{\rho} \int \frac{d^3p}{(2\pi)^3} \frac{d\omega}{2\pi} \left[\frac{\mathbf{p}^2}{2m} + \omega \right] A(\mathbf{p}, \omega) f(\omega).$$

Neutron matter and symmetry energy

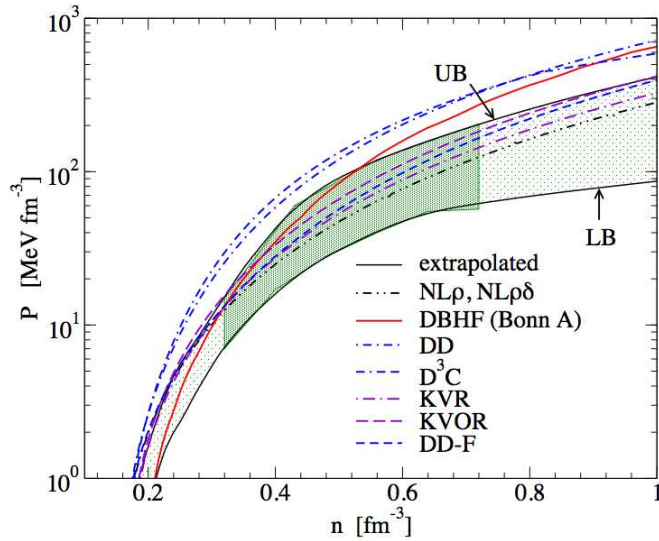


Symmetry energy computed with the parabolic approximation

$$\frac{E}{A}(n, \delta) = \frac{E}{A}(n, \delta = 0) + \delta^2 E_{sym}(n)$$

(Bombaci, Lombardo 1991)

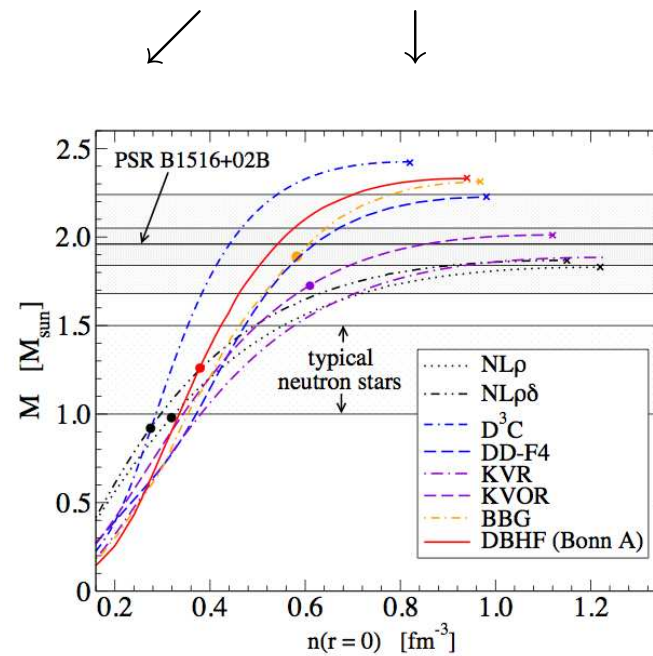
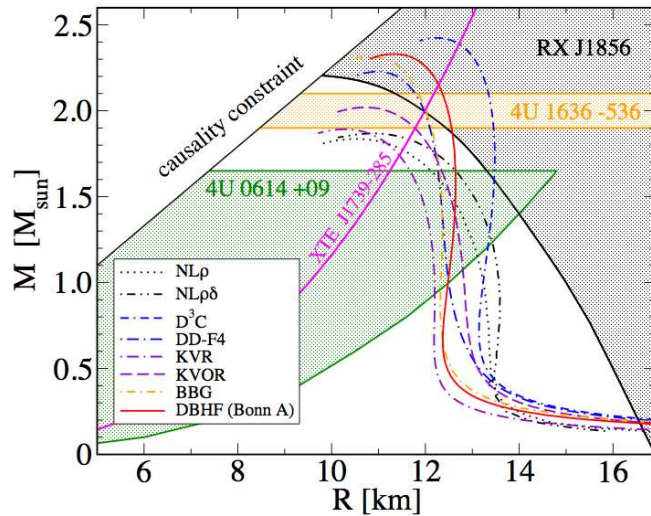
Constraints on the EoS



← Heavy-ion collisions

(Blaschke et al.
Phys.Rev. C74 (2006) 035802
arXiv:0808.1279)

Neutron stars



Single-particle properties

- Spectral function

$$G^<(\mathbf{p}, \omega) = f(\omega) A(\mathbf{p}, \omega)$$

- Self-energy

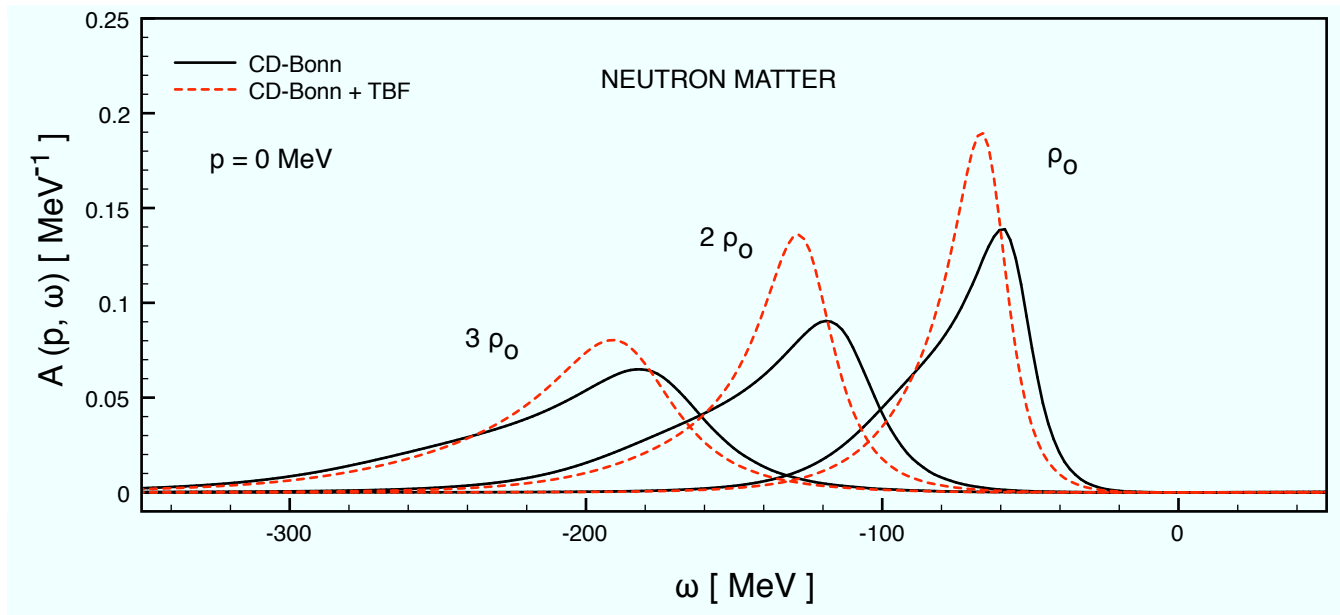
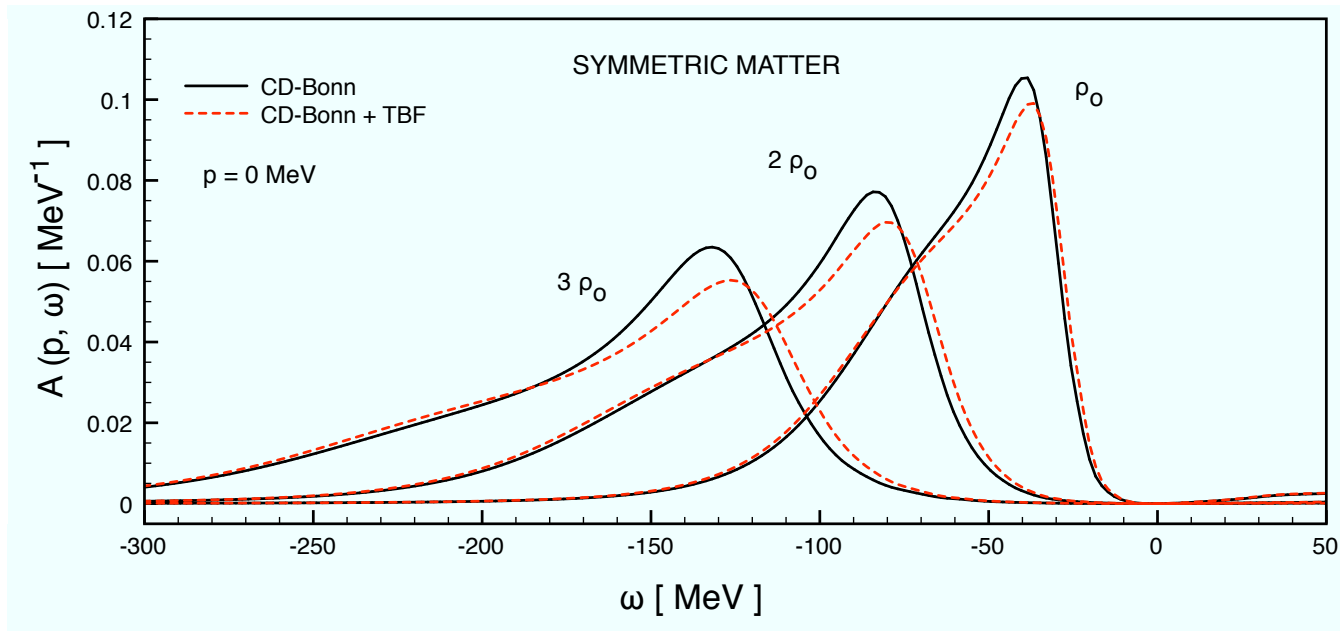
$$\text{Re } \Sigma(\mathbf{p}, \omega) = \Sigma_{HF}(\mathbf{p}) + \mathcal{P} \int \frac{d\omega'}{\pi} \frac{\text{Im } \Sigma(\mathbf{p}, \omega')}{\omega - \omega'}$$

dispersion relation $\omega_p = \frac{p^2}{2m} + \text{Re } \Sigma(\mathbf{p}, \omega_p)$

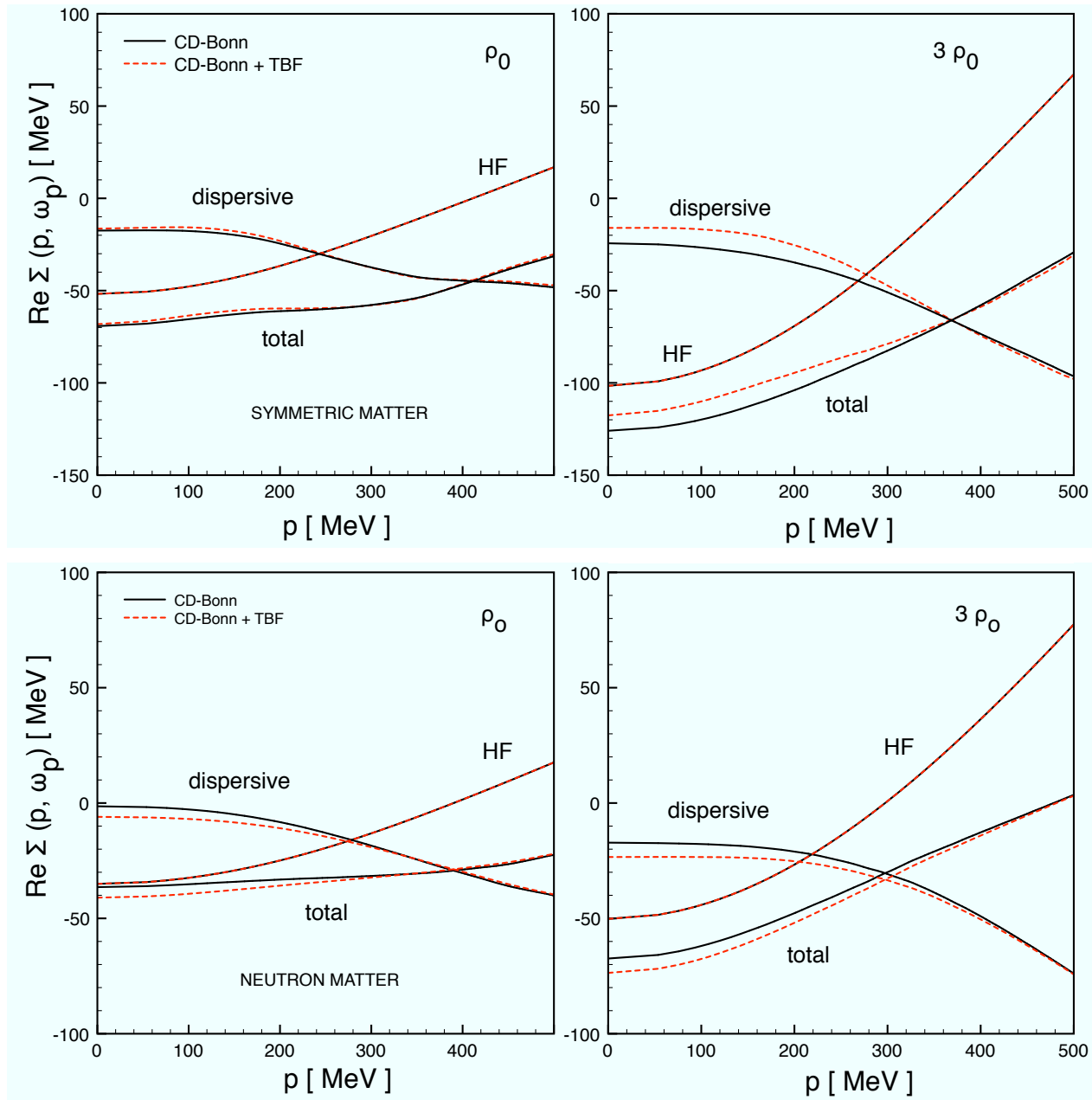
- Effective mass

$$\frac{\partial \omega_p}{\partial p^2} = \frac{1}{2m^*}$$

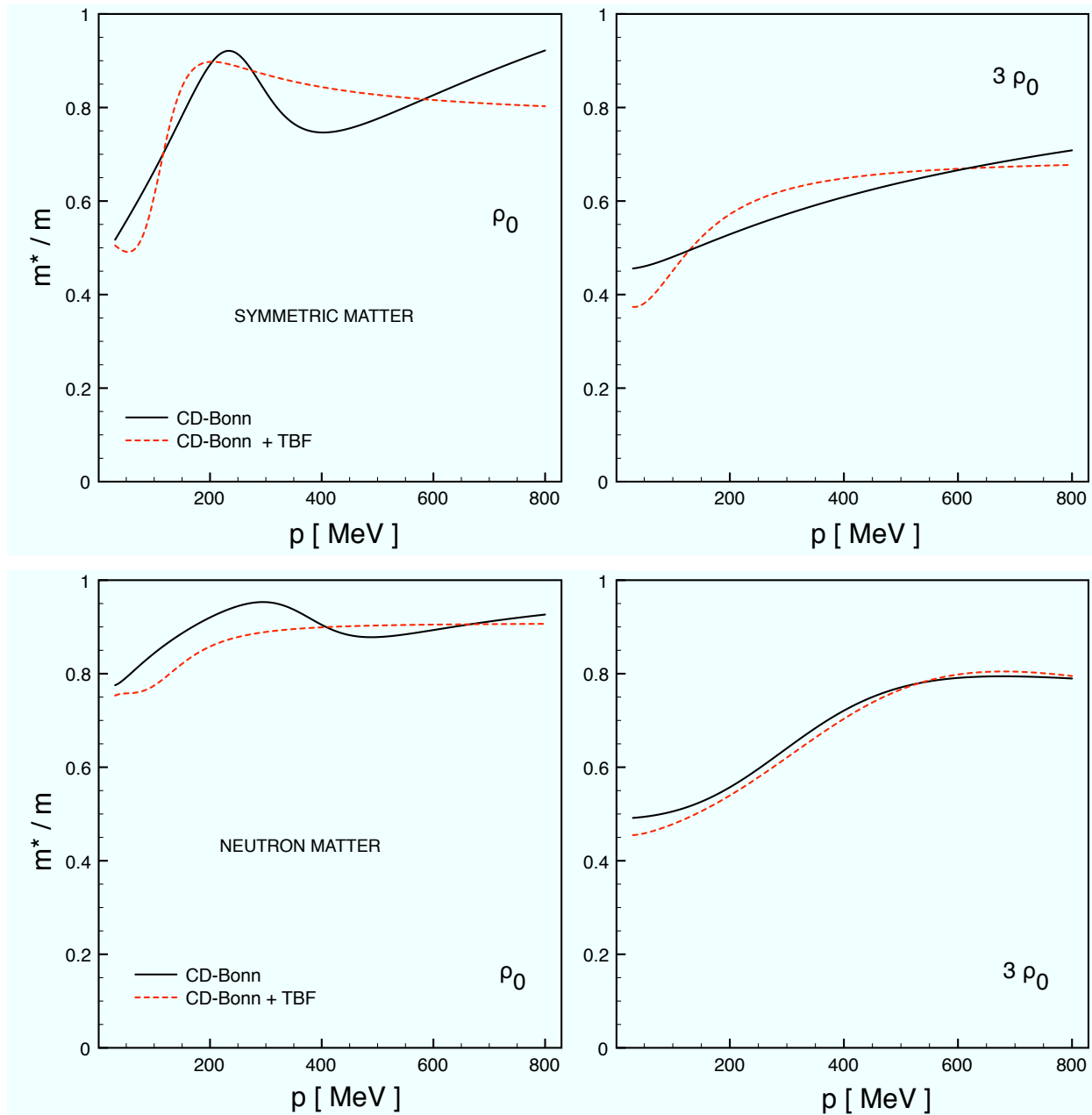
Modification of the spectral function



Self-energies



Effective mass



Summary

	$\rho_0 [\text{fm}^{-3}]$	E_0 [MeV]	K_0 [MeV]	S_0 [MeV]	m^*/m
exp	0.16 ± 0.01	-16 ± 1	210 ± 30	32 ± 6	≈ 0.8
CD-B 2	0.287	-19.9	70	32.2	0.90
CD-B 3	0.171	-16.3	148	39.7	0.87
Nijm 2	0.235	-18.4	76	30.5	0.87
Nijm 3	0.164	-16.4	158	37.1	0.90

★ on progress: complete EoS $\rightarrow \left\{ \frac{E}{N}, P, \frac{S}{N} \right\} (T, \rho)$

➔ (proto-) neutron star matter (T, ρ, x_p)

★ application of the Green's function formalism to the dynamics of central nuclear reactions (A. Rios, P. Danielewicz, arXiv:0801.4171)