International School of Nuclear Physic, 31st Course Neutrínos in Cosmology, in Astro-, Particle- and Nuclear Physics

Erice, Sicily, September 16-24, 2009

Supernova neutríno detectíon by terrestríal neutríno detectíon targets

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Outline

Introduction

- i. Neutrino Production Sources
- ii. Supernova Neutrino production
- iii. SN- v detection by terrestrial experiments (COBRA)
- iv. Neutrino-nucleus reaction cross sections at low energies

Nuclear Response to SN- v Spectra

- i. Response of Te isotopes to SN- v spectra
- ii. Convolution (folding) method for:
 - Differential cross sections <dσ/dω>
 - Double differential cross sections <d²σ/dΩdω>
- iii. Low energy beam neutrinos in SN- v searches
 - Reactor neutrino spectra
 - Beta beam neutrino spectra
- Summary Conclusions Outlook

Neutrino Sources

1) Astrophysical Neutrino Sources

- i. Solar Neutrinos
- ii. Supernova Neutrinos
- iii. Atmospheric Neutrinos
- iv. Cosmological Neutrinos

2) Laboratory Neutrino Sources

- i. Reactor neutrinos
 - Beta decay neutrinos
 - Slow pion and muon decay neutrinos
- ii. Accelerator ν (high energy neutrino beams)

Star evolution and v-production mechanisms

 $H \rightarrow He$

 $He \rightarrow C, O$

C→ Ne, Mg

O<u>→</u> S, Sì

S,Si→Fe

>At the end of hydrostatic burning a massive star ~ 8 Msun consists of concentric shells that are the relics of its previous burning phases

> When the mass of the iron core exceeds the $M_{Ch} = 1.4$ Msun, the gravitational pull > thermal pressure

> The core-collapse supernova starts!!!

Core-collapse simulation results

The main stages of stellar evolution (for massive stars) according to Janka et al., are:

>Initial phase of collapse

- Neutrino trapping
- Bounce and shock formation
- >Shock propagation and neutrino burst
- >Shock stagnation and neutrino heating
- >Neutrino cooling and neutrino driven wind

H. Th. Janka, K. Langanke et al, astro-ph/0612072 v1, 2006





Average energy of SN-v spectra

$$\left\langle E_{\nu}\right\rangle_{e} < \left\langle E_{\overline{\nu}}\right\rangle_{e} < \left\langle E_{\nu,\overline{\nu}}\right\rangle_{x}$$

Average energy of emitted neutrinos reflects the temperature of matter around the neutrinosphere.

$$T_{v_e} \approx 3,5 M e V \qquad T_{\overline{v_e}} \approx 5 M e V \qquad T_{v_x} \approx 8 M e V$$
flux
$$v_e \quad v_e \quad v_e \quad v_e \quad v_x \quad$$

11 16

25

E (MeV)

The Convolution Method in SN-v Searches

The differential v-nucleus cross-section $d\sigma(\varepsilon_v,\omega)/d\omega$ is folded by using the expression to study the nuclear response to SN-vspectra:

$$\left[\frac{d\sigma(\omega)}{d\omega}\right]_{fold} = \int_{\omega}^{\infty} \frac{d\sigma(\varepsilon_{\nu},\omega)}{d\omega} \mathcal{N}(\varepsilon_{\nu}) d\varepsilon_{\nu}$$

 $\omega = E_i - E_f = \varepsilon_i - \varepsilon_f$: excitation energy of the nucleus

The $n(\varepsilon_{v})$ is a specific v –energy distribution normalized to unity as:

$$\int \mathcal{N}(\varepsilon_{\nu}) d\varepsilon_{\nu} = 1$$

Energy distribution for SN-v

Fermi - Dirac
$$n_{FD}[T,n_{eff}](\varepsilon_v) = \frac{1}{F(n_{eff})T^3} \frac{\varepsilon_v^2}{Exp\left[\left(\frac{\varepsilon_v}{T}\right) - n_{eff}\right] + 1}$$

$$n_{\rm PL}[\langle \varepsilon_{\rm v} \rangle, a](\varepsilon_{\rm v}) = \frac{1}{c} \left(\frac{\varepsilon_{\rm v}}{\langle \varepsilon_{\rm v} \rangle}\right)^{\alpha} e^{-(a+1)\frac{\varepsilon_{\rm v}}{\langle \varepsilon_{\rm v} \rangle}}$$

Reactor neutrino spectrum

Power - law

$$\mathbf{n}_{v_{e}}(\varepsilon_{v_{e}}) = \frac{96\varepsilon_{v_{e}}^{2}}{m_{\mu}^{4}}(m_{\mu} - 2\varepsilon_{v_{e}})$$

Boosted beta-beam neutrino spectra

$$n_{y_{i}}(\varepsilon_{v}) = \frac{\ln 2}{m_{e} ft} F(\pm Z, E_{e}) E_{e} P_{e} \frac{\varepsilon_{v}^{2}}{\gamma^{2}(1+u^{2})} \frac{1}{2\gamma(1-u)}$$



Low-energy beta-beams in SN-V physics

Low – energy *beta-beams* can provide information about SN-v.



C. Volpe, J. Phys. G30, L1 (2004) C. Volpe, J. Phys. G34, R1 (2007) J. Serreau, C. Volpe: Phys.Rev. C70 (2004) 055502

Low-energy beta-beams in SN-V physics

We construct linear combinations of boosted beta beam spectra $n_{\nu i}$:

$$\mathcal{N}_{N_{\gamma}}(\varepsilon_{\nu}) = \sum_{i=1}^{N} a_{i} \mathcal{N}_{\gamma_{i}}(\varepsilon_{\nu})$$

The expansion coefficients $a_{i=1,2,...,N}$ for the boost factors $\gamma_{i=1,2,...,N}$ are obtained by minimizing the expression

$$\int_{\varepsilon_{\nu}} d\varepsilon_{\nu} \left| n_{N_{\gamma}}(\varepsilon_{\nu}) - n_{SN}(\varepsilon_{\nu}) \right|$$

C. Volpe, J. Phys. G30, L1 (2004)

N.Jachowicz, G.C.McLaughlin and C.Volpe, Phys. Review C 77,2008

Power-law fitting with 3-5 boost components



Nuclear response to SN-v of COBRA target

The aim of this work is to study the response to SN-v of the nuclear isotopes *Te*, contained in the *COBRA* and other detectors, through the neutral current reactions

 $^{128,130}Te(v_e,v_e')^{128,130}Te^*$



K. Zuber, Phys. Lett. B 519, 1-7 (2001); Prog. Part. Nucl. Phys. 57, 235-240 (2006);















ω (MeV)

ω (MeV)

Summary – Conclusions

>We study the response of Te isotopes to the SN- ν spectra by evaluating the folded:

1) differential cross-sections do/dw

2) double deferential cross-sections <d²σ/dΩdω>

>We used the convolution method and employed

 (i) Fermi-Dirac neutrino energy distribution
 (ii) Power-law neutrino energy distribution
 (iii) Reactor neutrino energy distribution
 (iv) Linear-combination of boosted beta-beam neutrinos

 They are appropriate for low energy neutrinos produced during Supernova explosions.

> We found that there are not dramatical differences between the above distributions.

Currently we are working on the charged-current neutrinos processes of these isotopes



Thank you!!!

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<u>Acknowledgments:</u>

I wish to acknowledge financial support from the ΠΕΝΕΔ-03/807, Hellenic GSRT Project