

Baryogenesis via Sterile Neutrino Oscillation and Neutrino Parameters

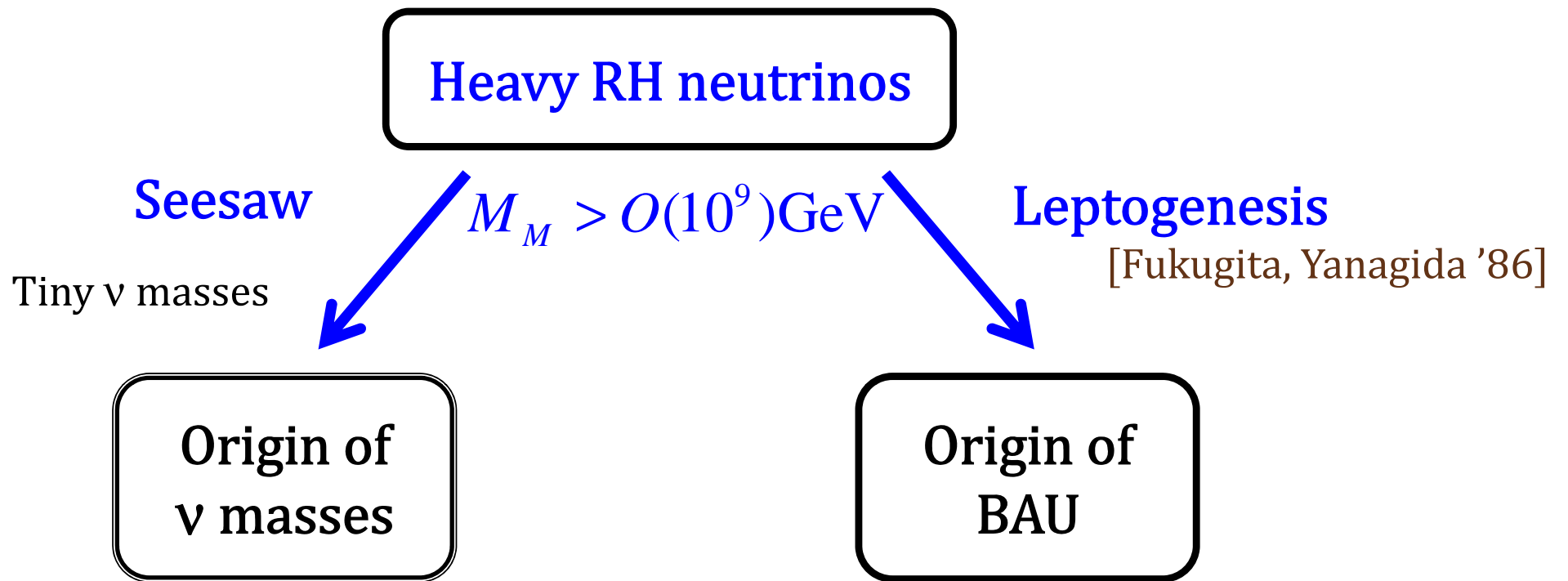
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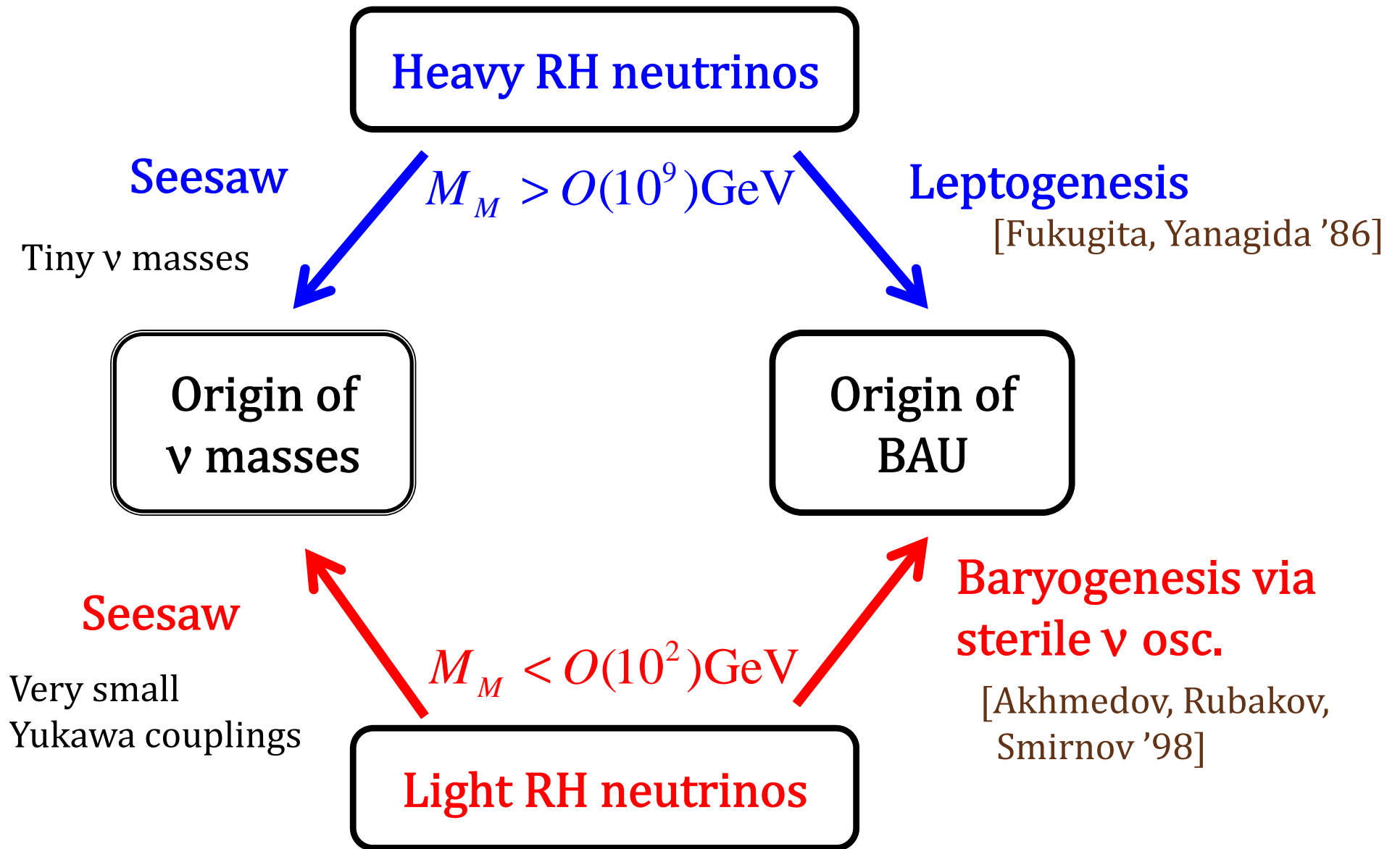
@ Erice-Sicily, 20 September, 2009

In collaboration with H. Ishida (Niigata University)

Origin of
 ν masses

Origin of
BAU





■ The MSM + Three RH neutrinos (N_1, N_2, N_3)

$$L = i\bar{N}_I \gamma^\mu \partial_\mu N_I - F_{\alpha I} \bar{L}_\alpha \Phi N_I - \frac{M_I}{2} \bar{N}_I^c N_I + h.c.$$

- Dirac and Majorana masses of neutrinos

$$[M_D]_{\alpha I} = F_{\alpha I} \langle \Phi \rangle \quad \text{and} \quad [M_M]_{II} = M_I \quad (\alpha = e, \mu, \tau; \quad I = 1, 2, 3)$$

■ Key Assumption

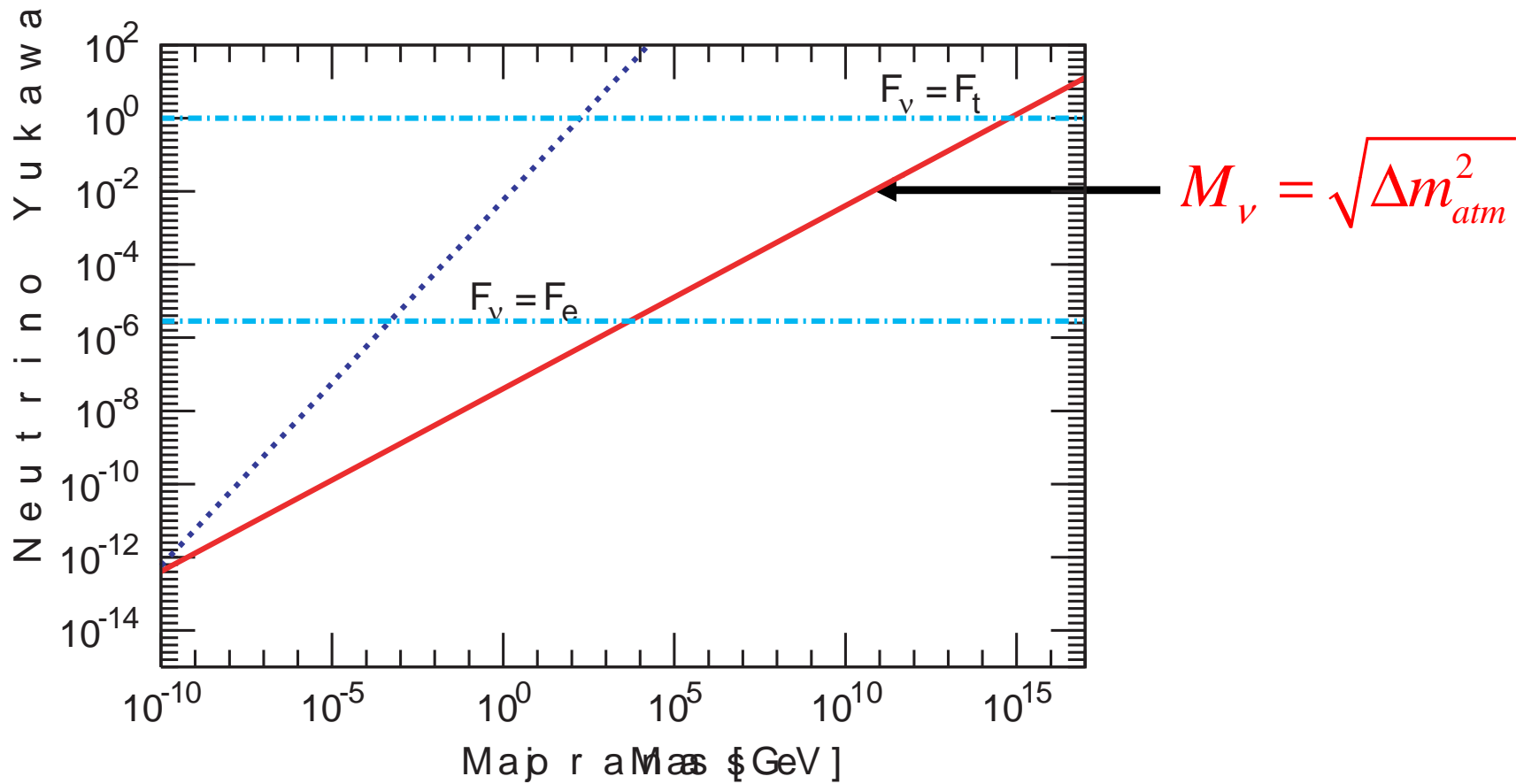
$$|M_D|_{\alpha I} \ll M_I < O(10^2) \text{ GeV}$$

- Light RH neutrinos could be tested!!
- Seesaw mechanism still works!! $M_\nu = -M_D M_M^{-1} M_D^T$

$$F \sim 4 \times 10^{-8} \left(\frac{M_I}{\text{GeV}} \right)^{1/2} \left(\frac{m_\nu^2}{2.5 \times 10^{-3} \text{ eV}^2} \right)^{1/4}$$

Scale of Majorana mass

$$M_\nu = -M_D^T \frac{1}{M_M} M_D \Rightarrow F^2 = M_M M_\nu / \langle H \rangle^2$$



Roles of three sterile (RH) neutrinos

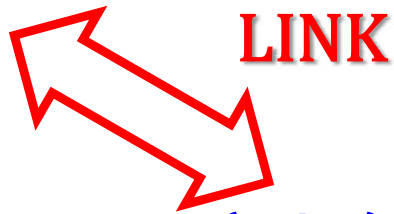
N_1 (“Dark” sterile neutrino)

- **Dark Matter Candidate** (Here, we do not specify its detail.)
- To avoid constraints, Yukawa’s should be suppressed

essentially, $F_{\alpha 1} \approx 0$

N_2 and N_3 (“Bright” and “Clear” sterile neutrinos)

- **Neutrino Oscillation data**
 - Masses and mixings
- **Baryon Asymmetry of the Universe (BAU)**
 - Mechanism via sterile neutrino oscillation



■ We will show

How CPV in neutrino sector relates with BAU?

■ Outline

- Baryogenesis via sterile neutrino oscillation
- BAU and low-energy CPV in neutrino sector
- Summary

Baryogenesis via Sterile Neutrino Oscillation

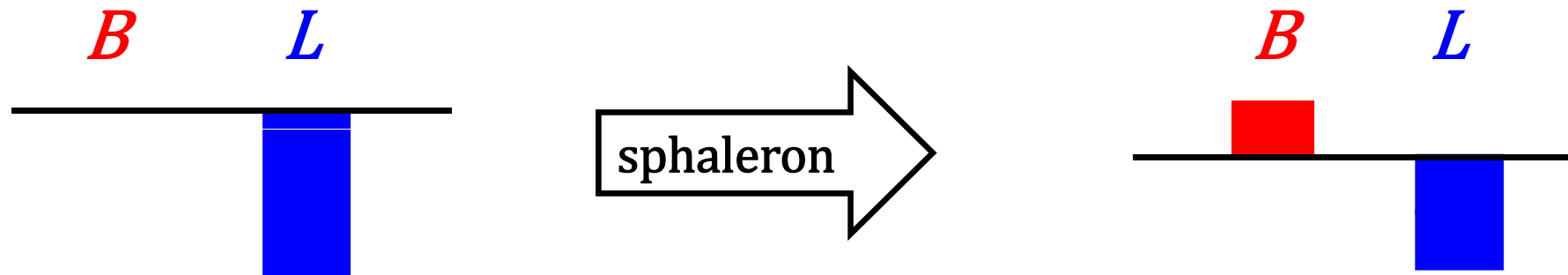
[Akhmedov, Rubakov, Smirnov '98]
[TA, Shaposhnikov '05]

Baryogenesis scenario
by sterile neutrinos
with $M_M < O(100) \text{ GeV}$

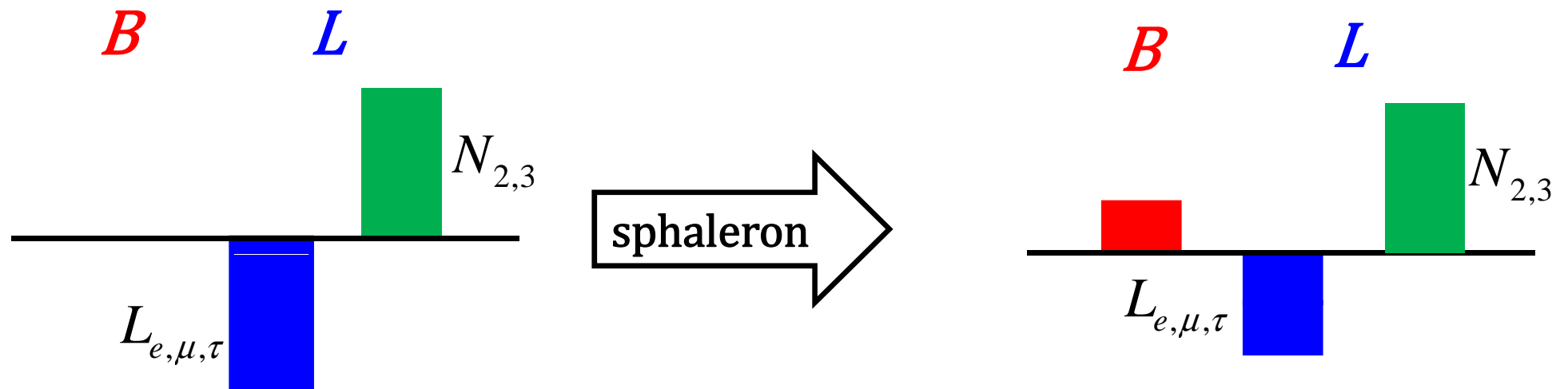
- The Sakharov's Conditions
 - CPV
 - CPV phases in neutrino sector
 - Out of equilibrium
 - No 1st order phase transition as in the SM
 - $N_{2,3}$ can be out of equilibrium due to small Yukawa couplings
 - B (L) violation
 - (B+L) violation due to sphaleron for $T > 100 \text{ GeV}$
 - Total lepton number violation is ineffective for $T > 100 \text{ GeV}$ since $M_M < 100 \text{ GeV}$

Key points of this baryogenesis

■ Baryogenesis via leptogenesis

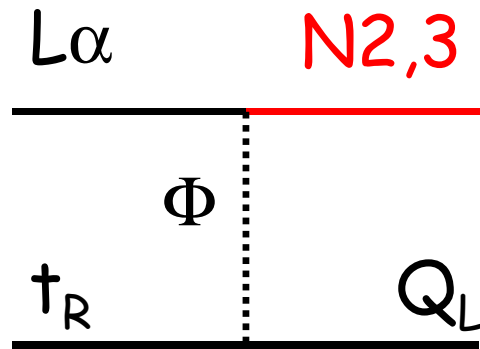


■ Baryogenesis via sterile neutrino oscillation

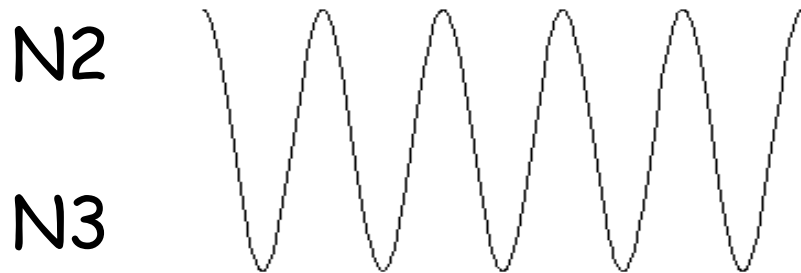


First step: at F^2 order

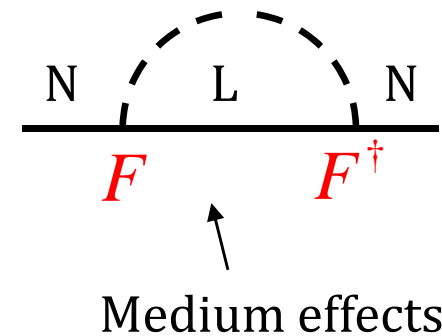
- N2 and N3 are produced via scatterings



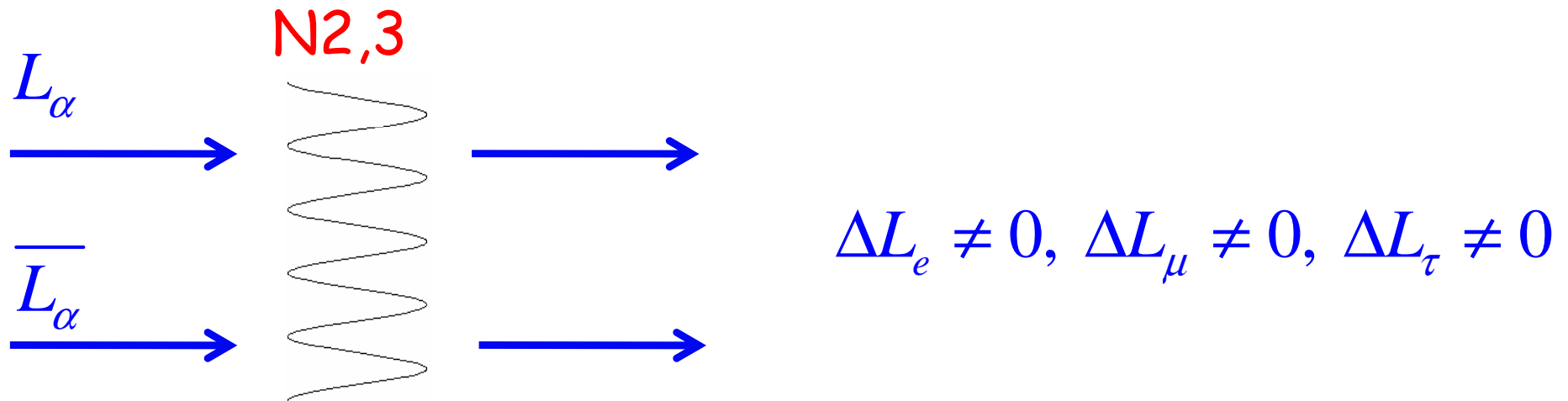
- N2 and N3 oscillate at $T = T_L \sim (M_{pl} \Delta M_{32}^2)^{1/3}$



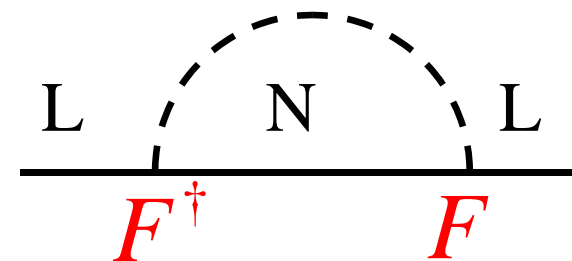
$$V_N = \frac{T^2}{8k} F^\dagger F$$



- Active flavor asymmetries are generated



- Evolution rates of L_α and \overline{L}_α are different due to CVP in

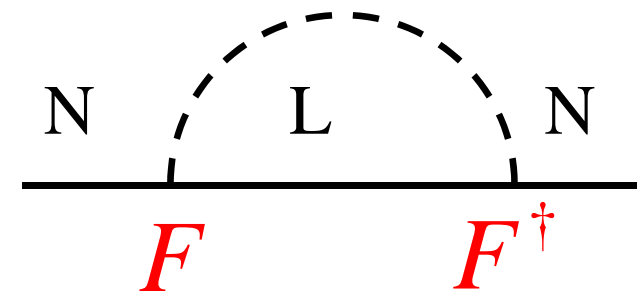


Final step: at F^6 order

- Total asymmetries in active and sterile sectors are generated.

$$\begin{array}{ccc}
 N_I & \longrightarrow & \Delta L_e \\
 \longrightarrow & & \longrightarrow \\
 \overline{N_I} & \longrightarrow & \Delta L_\mu \\
 \longrightarrow & & \longrightarrow \\
 \overline{N_I} & \longrightarrow & \Delta L_\tau \\
 \longrightarrow & & \longrightarrow
 \end{array}
 \quad
 \begin{array}{l}
 \Delta N_{\text{tot}} = \Delta N_2 + \Delta N_3 \neq 0 \\
 \Delta L_{\text{tot}} = \Delta L_e + \Delta L_\mu + \Delta L_\tau \neq 0
 \end{array}$$

- Evolution rates of N_I and $\overline{N_I}$ are different due to ΔL_α and CPV in



Evolution of each asymmetries

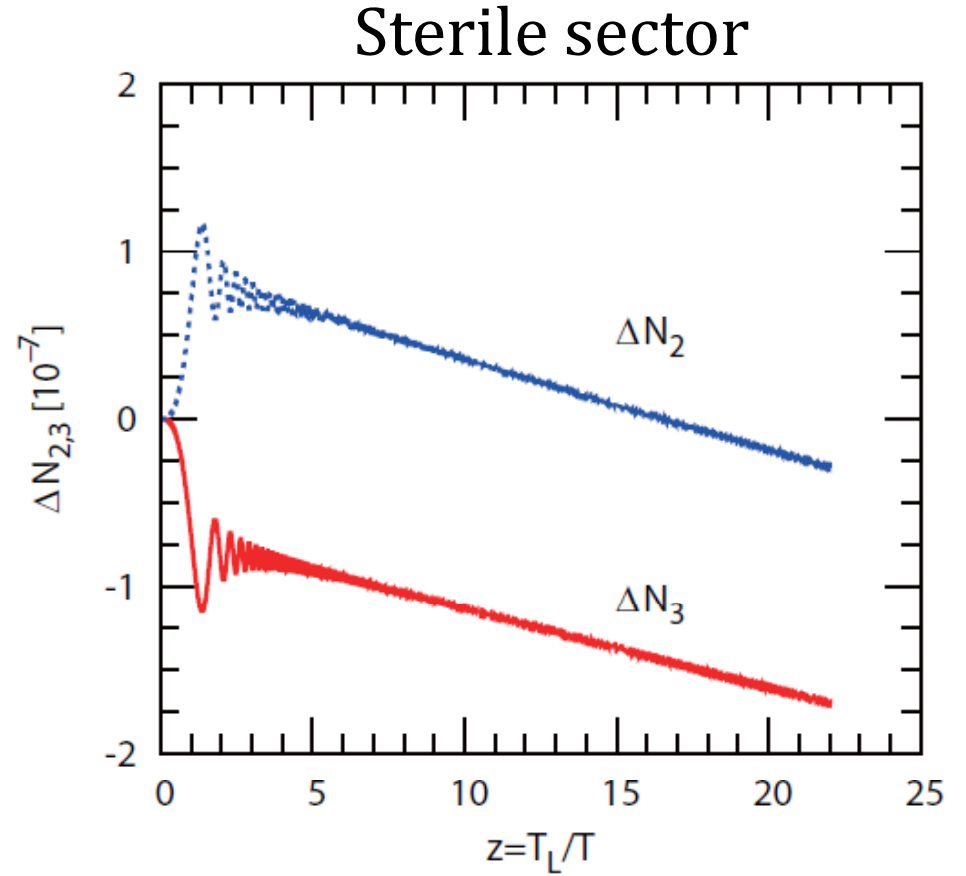
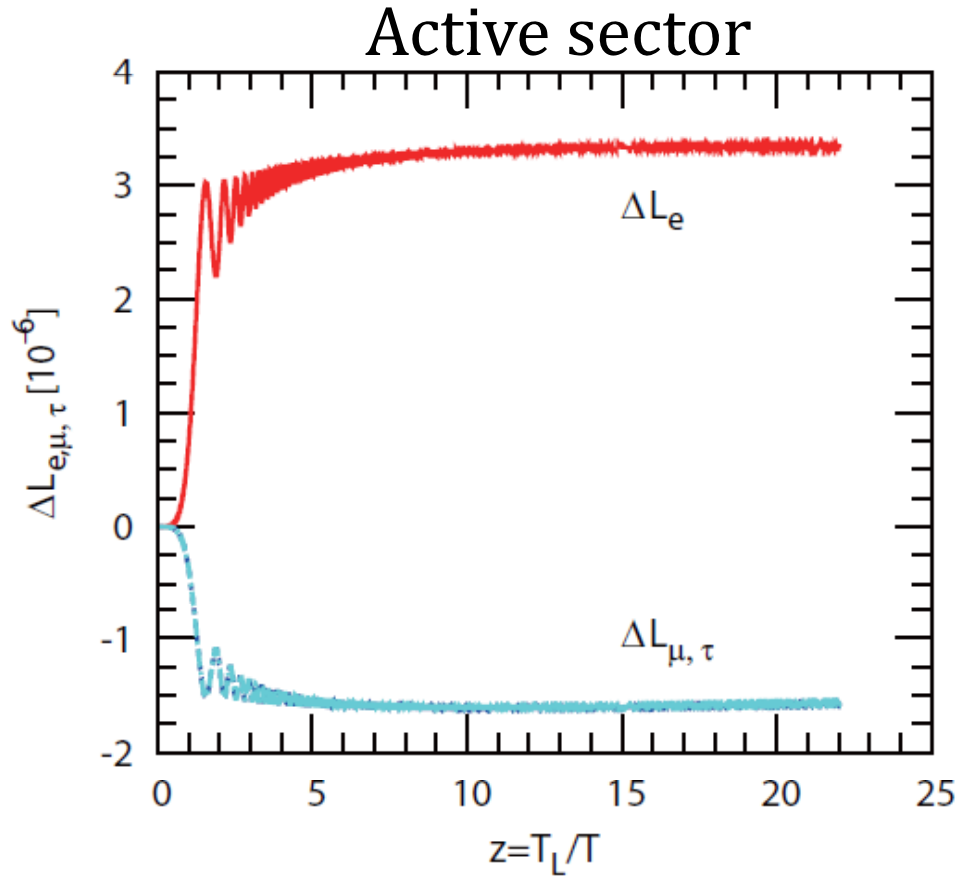


Figure 5: Evolution of asymmetries in terms of $z = T_L/T$. Here we take $M_3 = 3$ GeV, $\Delta M_{32}^2/M_3^2 = 10^{-8}$, $\xi = +1$, $\sin \theta_{13} = 0.2$, $\phi = 0$, $\omega = \pi/4$ and $\delta = 3\pi/2$.

Figure 6: Evolution of asymmetries in terms of $z = T_L/T$. Here we take $M_3 = 3$ GeV, $\Delta M_{32}^2/M_3^2 = 10^{-8}$, $\xi = +1$, $\sin \theta_{13} = 0.2$, $\phi = 0$, $\omega = \pi/4$ and $\delta = 3\pi/2$.

$$T_L \sim (M_{pl} \Delta M_{32}^2)^{1/3} = 2.2 \text{ TeV}$$

Evolution of asymmetries

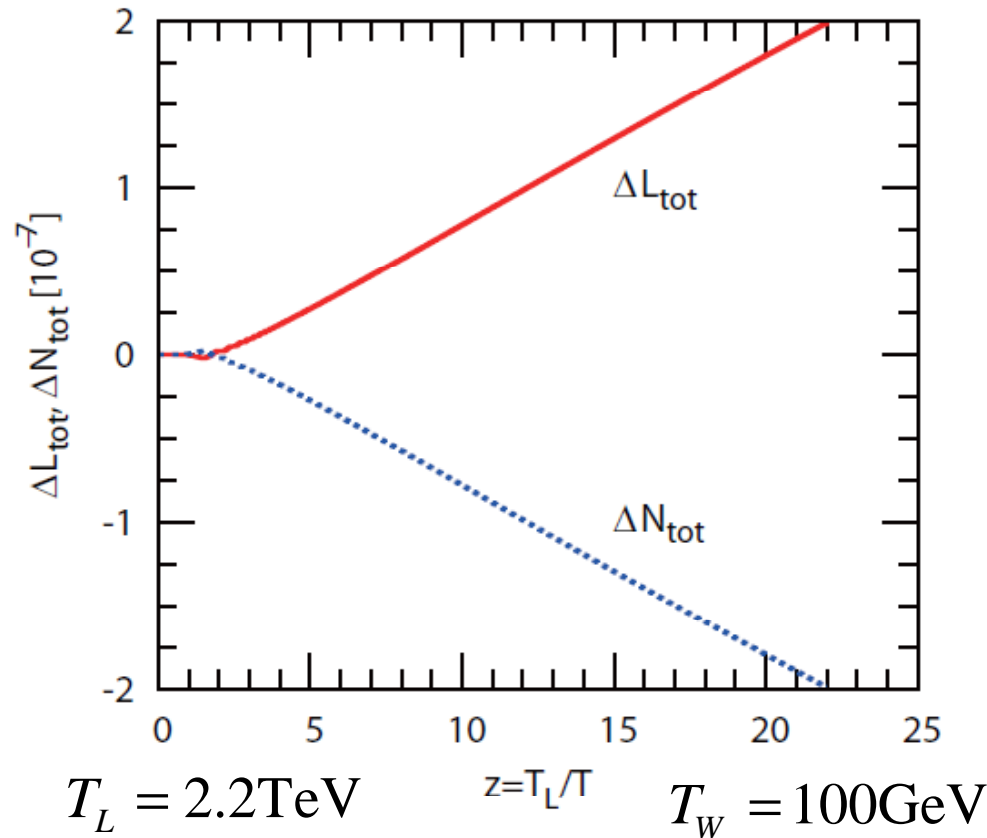


Figure 7: Evolution of asymmetries in terms of $z = T_L/T$. Here we take $M_3 = 3 \text{ GeV}$, $\Delta M_{32}^2/M_3^2 = 10^{-8}$, $\xi = +1$, $\sin \theta_{13} = 0.2$, $\phi = 0$, $\omega = \pi/4$ and $\delta = 3\pi/2$.

$$M_3 = 3\text{GeV} \quad M_2^2 = M_3^2 (1 - 10^{-8})$$

Degenerate masses are needed!

Sphaleron converts ΔL_{tot} partially into baryon asymmetry

[Kuzmin, Rubakov, Shaposhnikov '85]

$$\Delta B = -\frac{28}{79} \Delta L_{tot} \neq 0$$

BAU

$$\left. \frac{n_B}{s} \right|_0 = -7.0 \times 10^{-4} \left. \Delta L_{tot} \right|_{T=T_W}$$

$$\left. \frac{n_B}{s} \right|_{\text{obs}} = (8.81 \pm 0.23) \times 10^{-11}$$

Baryon Asymmetry of Univ. and Low-energy CPV in neutrino sector

[TA, Ishida in preparation]

Sterile Neutrinos: N_2 and N_3

Seesaw

$$[M_\nu]_{\alpha\beta} = - \sum_{I=2,3} \frac{[M_D]_{\alpha I} [M_D]_{\beta I}}{M_I}$$

BAU

$$\left. \frac{n_B}{s} \right|_{\text{obs}} = (8.81 \pm 0.23) \times 10^{-11}$$

Neutrino Yukawa couplings for N2,3

$$F = U_{\text{PMNS}} D_\nu^{1/2} \Omega D_N^{1/2} / \langle \Phi \rangle \quad [\text{Casas, Ibarra '01}]$$

(in NH)

Low energy parameters

$D_\nu^{1/2} = \text{diag}(\sqrt{m_1}, \sqrt{m_2}, \sqrt{m_3})$: active ν masses

$$U_{\text{PMNS}} = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & 0 \\ -c_{23}s_{12} - s_{23}c_{12}s_{13}e^{i\delta} & c_{23}c_{12} - s_{23}s_{12}s_{13}e^{i\delta} & s_{13}e^{-i\delta} \\ s_{23}s_{12} - c_{23}c_{12}s_{13}e^{i\delta} & -s_{23}c_{12} - c_{23}s_{12}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix}$$

Dirac phase δ

$$\begin{pmatrix} 1 & & \\ & e^{i\phi} & \\ & & 1 \end{pmatrix}$$

Majorana phase ϕ

High energy parameters

$D_N^{1/2} = \text{diag}(\sqrt{M_2}, \sqrt{M_3})$: sterile ν masses

$$\Omega = \begin{pmatrix} 0 & 0 \\ \cos \omega & -\sin \omega \\ \xi \sin \omega & \xi \cos \omega \end{pmatrix} \quad \begin{array}{l} \omega : \text{complex number} \\ \xi = \pm 1 \end{array}$$

High energy phase ω

- Sterile neutrinos: ρ_{NN} : density matrix for $N_{2,3}$

$$i \frac{d\rho_{NN}}{dt} = \left[H_{NN}^0 + V_N, \rho_{NN} \right] - \frac{i}{2} \left\{ \Gamma_{NN}^d, \rho_{NN} - \rho_{NN}^{eq} \right\}$$

Effective potential

Destruction rate

$$V_N = \frac{T^2}{8k} F^\dagger F$$

$$\Gamma_N^d \approx 0.04 V_N$$

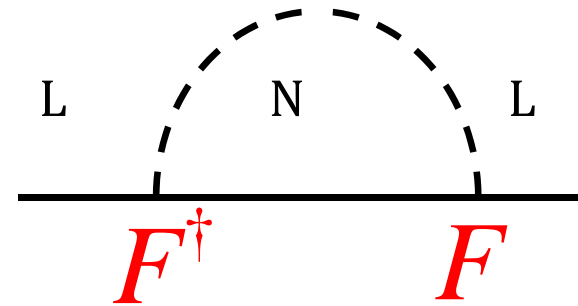
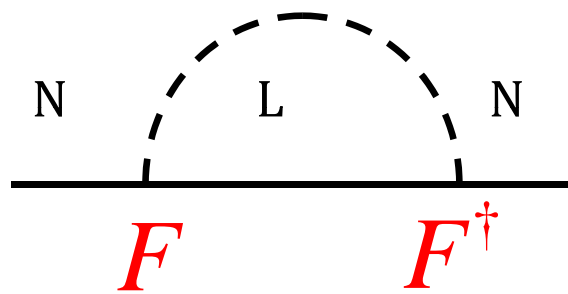
$$F^\dagger F = D_N^{1/2} \Omega^\dagger D_\nu \Omega D_N^{1/2}$$

Independent on neutrino mixing matrix U_{PMNS} !

→ insensitive to low-energy neutrino parameters !

Baryogenesis via sterile neutrino osc.

- Include the new effects [TA, Shaposhnikov '05]
 - Exchange of asymmetries between sterile neutrinos and active neutrinos (+ charged leptons)



■ Sterile neutrinos:

[TA, Shaposhnikov '05]

$$i \frac{d\rho_{NN}}{dt} = [H_{NN}^0 + V_N, \rho_{NN}] - \frac{i}{2} \{ \Gamma_{NN}^d, \rho_{NN} - \rho_{NN}^{eq} \} + \frac{i \sin \phi}{4} T \cdot F^\dagger (\rho_{LL} - \rho_{LL}^{eq}) F$$

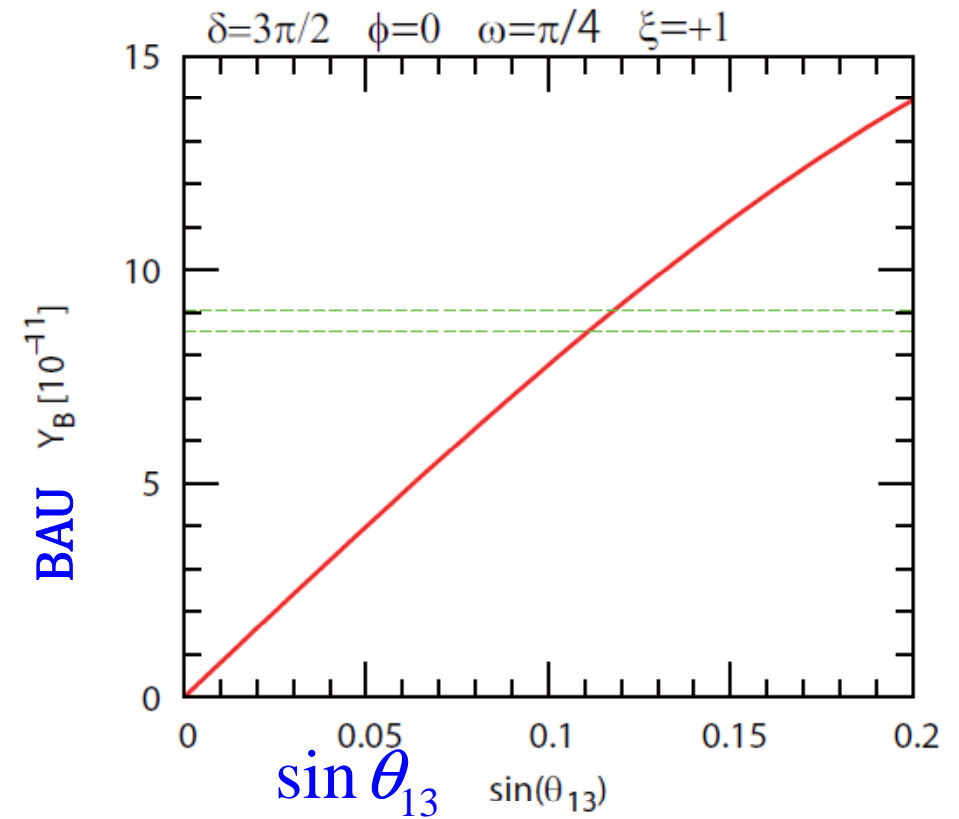
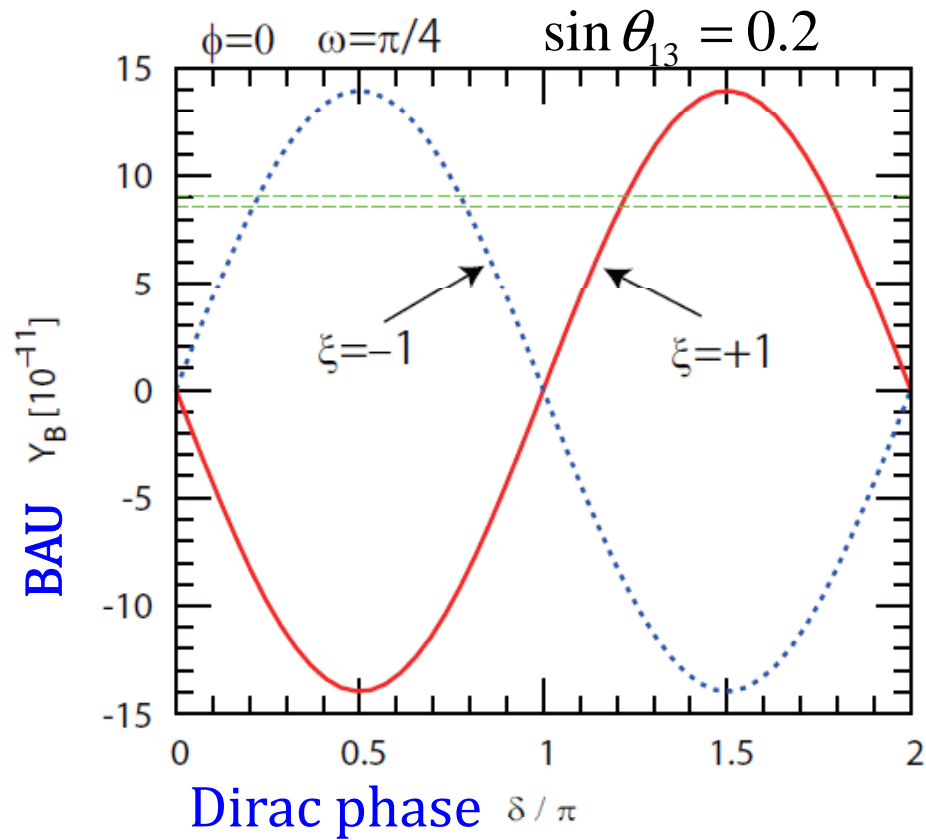
■ Active neutrinos:

$$i \frac{d\rho_{LL}^{diag}}{dt} = [H_{LL}^0 + V_L, \rho_{LL}^{diag}] - \frac{i}{2} \{ \Gamma_{LL}^d, \rho_{LL}^{diag} - \rho_{LL}^{eq} \} + \frac{i \sin \phi}{4} T \cdot F (\rho_{NN} - \rho_{NN}^{eq}) F^\dagger$$

Does depend on neutrino mixing matrix U_{PMNS} !
→ sensitive to low-energy neutrino parameters !

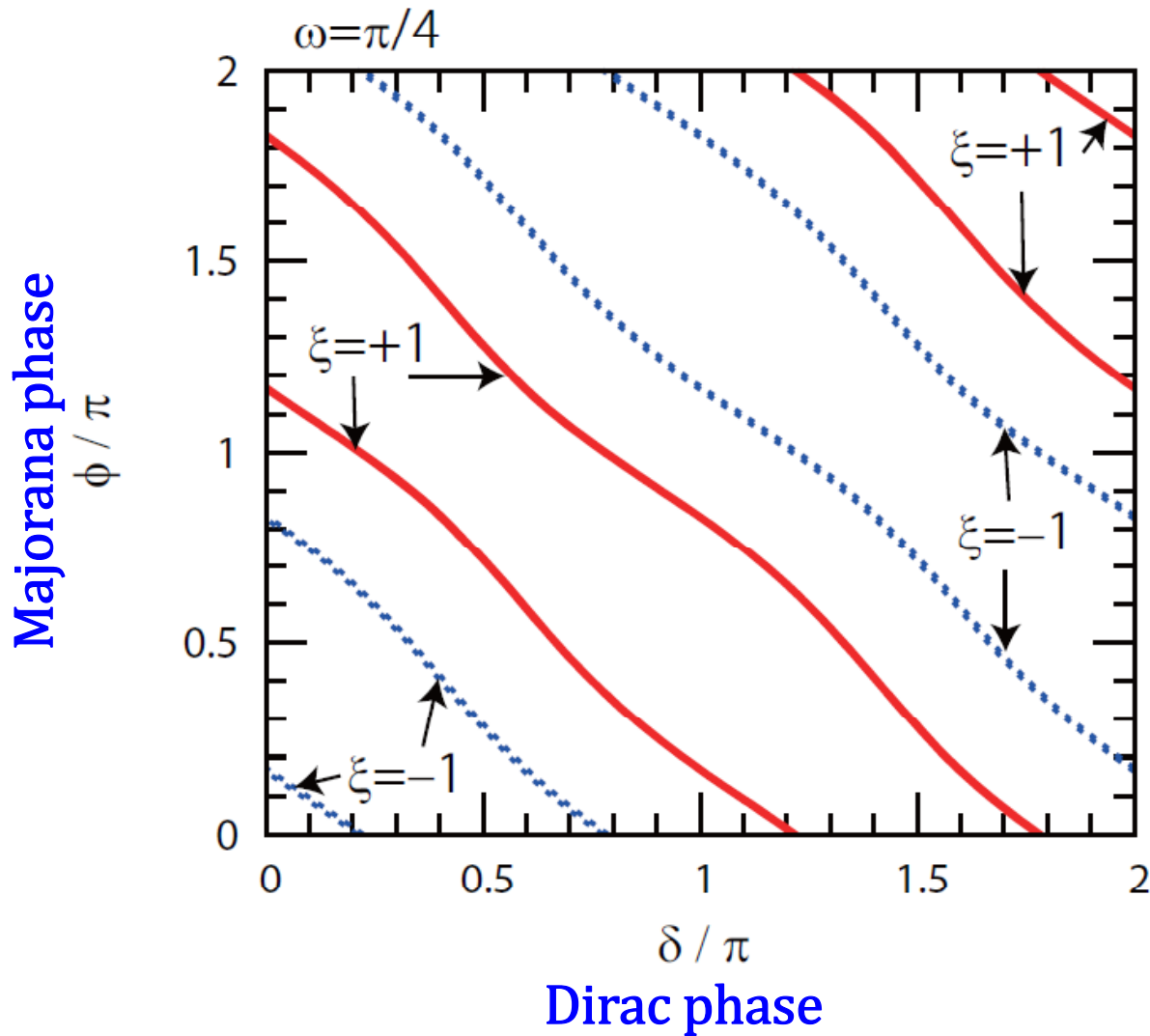
Let us see how BAU depends on

Dirac phase δ and Majorana phase ϕ !



$$U_{\text{PMNS}} = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -c_{23}s_{12} - s_{23}c_{12}s_{13}e^{i\delta} & c_{23}c_{12} - s_{23}s_{12}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{23}s_{12} - c_{23}c_{12}s_{13}e^{i\delta} & -s_{23}c_{12} - c_{23}s_{12}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix} \begin{pmatrix} 1 & & \\ & e^{i\phi} & \\ & & 1 \end{pmatrix} \quad \Omega = \begin{pmatrix} 0 & 0 \\ \cos \omega & -\sin \omega \\ \xi \sin \omega & \xi \cos \omega \end{pmatrix}$$

Regions accounting for BAU



$$\left. \frac{n_B}{S} \right|_{\text{obs}} = (8.81 \pm 0.23) \times 10^{-11}$$

Inputs:

$$M_3 = 3\text{GeV}$$

$$M_2^2 = M_3^2 (1 - 10^{-8})$$

$$\sin \theta_{13} = 0.2$$

$$\Omega = \begin{pmatrix} 0 & 0 \\ \cos \omega & -\sin \omega \\ \xi \sin \omega & \xi \cos \omega \end{pmatrix}$$

- Connection between neutrino masses and BAU is attractive and important idea
- Conventional seesaw scenario ($M_M > 10^9 \text{ GeV}$)
[Seesaw + Leptogenesis]
 - Natural framework of SUSY GUT ...
 - Exp. test for RH neutrinos is impossible
- Connection can be obtained even with $M_M < 10^2 \text{ GeV}$
[Seesaw + Baryogenesis via sterile neutrino osc.]
 - Such sterile neutrinos might be tested
 - Connection between BAU and CPV in neutrino oscillations

■ Normal hierarchy of (active) neutrino masses

$$m_3 = \sqrt{\Delta m_{atm}^2 + \Delta m_{sol}^2}, \quad m_2 = \sqrt{\Delta m_{sol}^2}, \quad m_1 = 0 \quad \Delta m_{atm}^2 = 2.5 \times 10^{-3} \text{ eV}^2$$

$$\theta_{23} = \pi / 4, \quad \theta_{12} = \pi / 5 \quad \sin \theta_{13} \leq 0.2 \quad \Delta m_{sol}^2 = 8.0 \times 10^{-5} \text{ eV}^2$$

$$\delta, \phi$$

$$M_3 = 3 \text{ GeV}, \quad M_2^2 = M_3^2 (1 - 10^{-8})$$

$$\omega = \text{Re } \omega + \text{Im } \omega$$

$$\xi = \pm 1$$

Sterile neutrino oscillation

- Flavor mixing of sterile neutrinos is induced from thermal potential

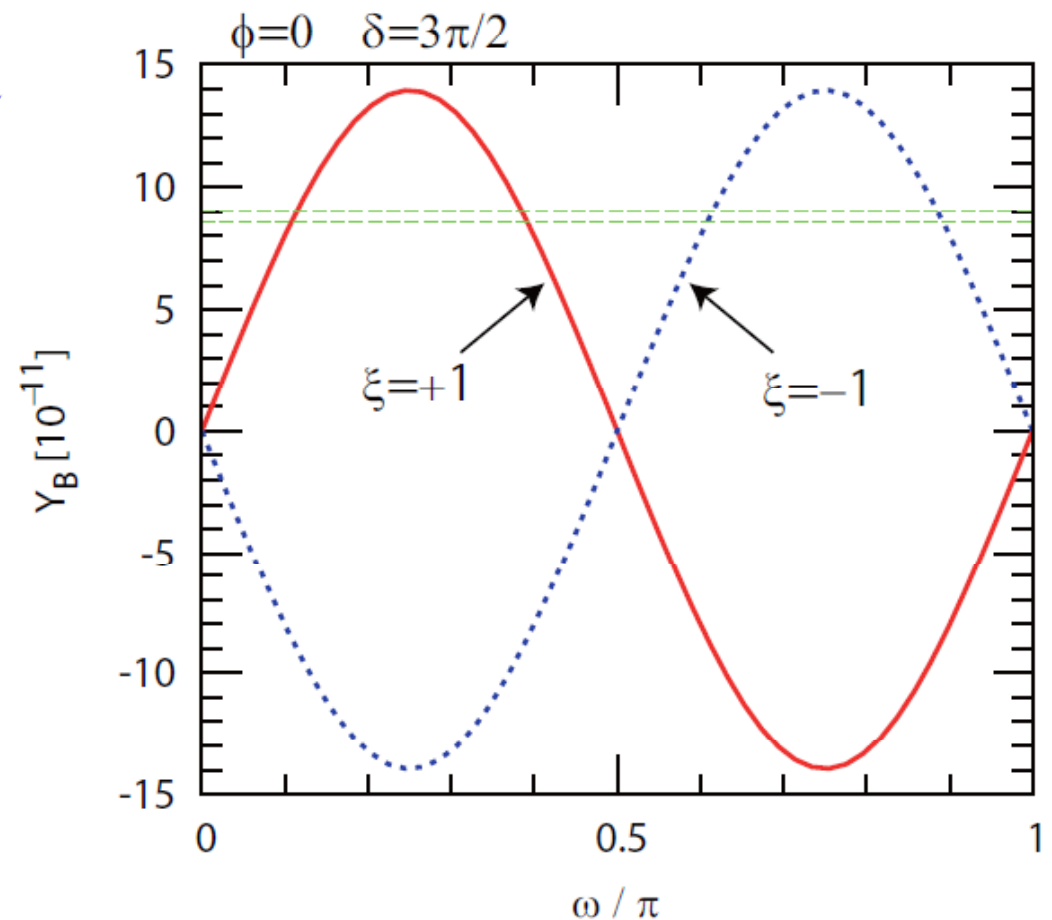
$$V_N \propto F^\dagger F = D_N^{1/2} \Omega^\dagger D_\nu \Omega D_N^{1/2}$$

$$\Omega = \begin{pmatrix} 0 & 0 \\ \cos \omega & -\sin \omega \\ \xi \sin \omega & \xi \cos \omega \end{pmatrix}$$

ω : complex number

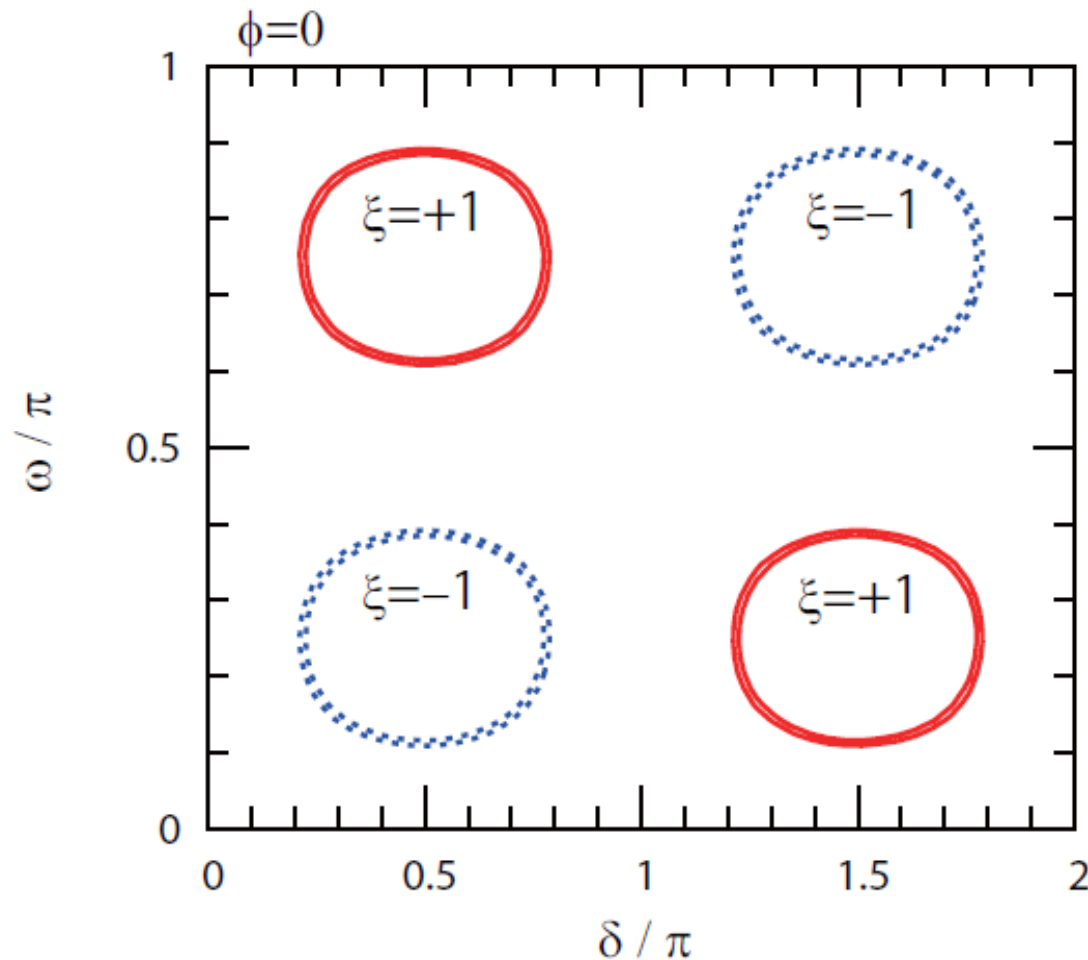
$$\xi = \pm 1$$

BAU vanishes when there is no sterile neutrino oscillation !



Regions accounting for BAU

■ $\left. \frac{n_B}{s} \right|_{\text{obs}} = (8.81 \pm 0.23) \times 10^{-11}$



Normal hierarchy

$$M_3 = 3 \text{ GeV}$$

$$M_2^2 = M_3^2(1 - 10^{-8})$$

$$\sin \theta_{13} = 0.2$$