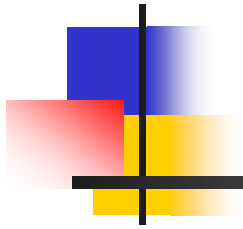


Neutrino Mass and Unification



R. N. Mohapatra



2009 Erice School on Neutrinos



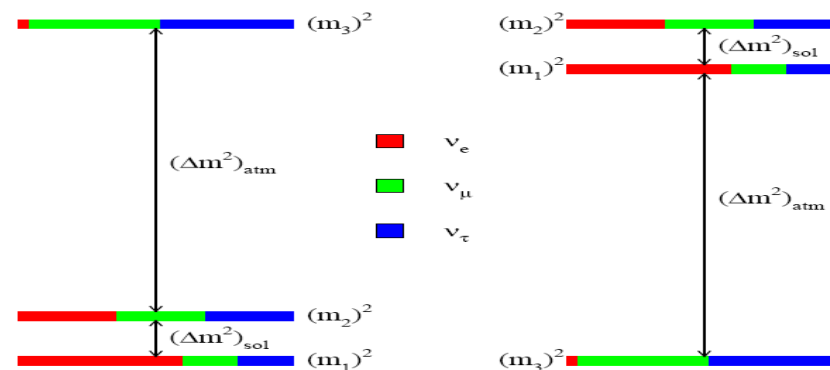
Implications of neutrino mass results for Theory

- What we may have learnt ?
- What we need to learn and how ?

A subjective overview

Present information: Masses and mixings

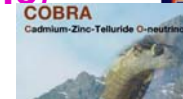
- **Masses:** $\Delta m_{sol}^2 \cong 7.67 \times 10^{-5} eV^2$; $\Delta m_{Atm}^2 \cong 2.39 \times 10^{-3} eV^2$
- **Mixings:** $\sin^2 \theta_{12} \cong .312$; $\sin^2 \theta_{23} \cong .466$; $\sin^2 \theta_{13} \leq .04$
- **Overall mass scale:** $< .1 - 1 eV$ (roughly) (WMAP,..)
- **Mass ordering not known:**



Need to know...

(i) Majorana or Dirac $\beta\beta_{0\nu}$

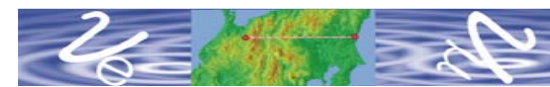
(Nucl matrix element : Fassler et al.; Talk by Simkovic, Suhonen ..)



(ii) Absolute mass scale:



(iii) Mass ordering:



(iv) Value of θ_{13}



(v) CP phase

Dirac vs Majorana

- Very important for determining the nature of new physics:

- **Can we tell experimentally?**

- Observe $\beta\beta_{0\nu} \rightarrow$ **Majorana**.

- If no signal till 20 meV \rightarrow could be Majorana with normal hierarchy;

- However, no $\beta\beta_{0\nu}$ signal till 20 meV + $\Delta m_{23}^2 < 0$ from long Baseline expts \rightarrow strong hint for

Dirac

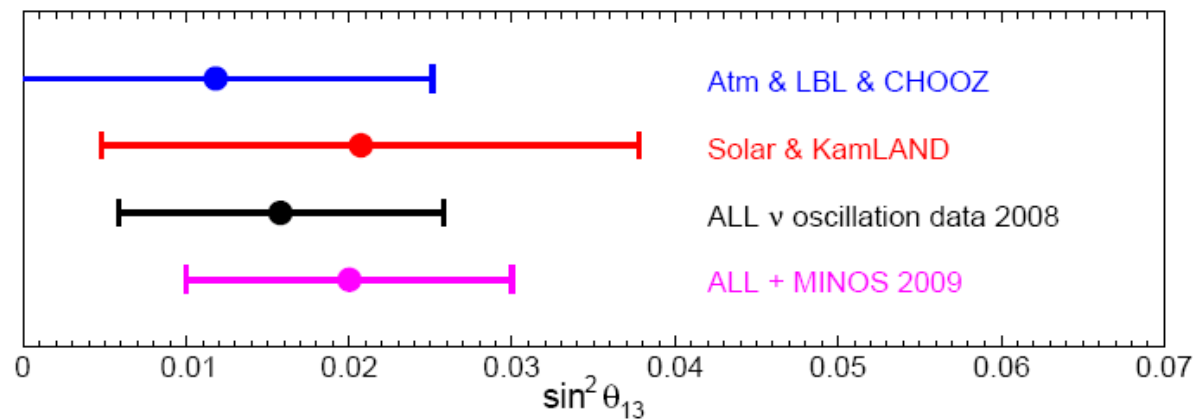
Rest of this talk assumes Majorana nu

What is value of

$$\theta_{13}$$

?

- Recent indications ? (Fogli, Lisi, Marrone, Pallazo and Rotunno'08)



Too early for definite conclusion--However

Value of θ_{13} significant for new physics



Goal of Theory

- Determining and understanding the Neutrino mass matrix :
- Two parts to the story:

$$M_\nu = m_\nu \times A_F$$

(i) Scale m_ν

(ii) Flavor structure A_F
(The neutrino matrix)



Challenges

(i) Scale issue: Why $m_\nu \ll m_{q,l}$? :

(ii) Flavor issues: A_F ?

A. Milder mass hierarchy compared to quarks and charged leptons:

$$\frac{m_{sol}}{m_{atm}} \approx \theta_c \gg \frac{m_\mu}{m_\tau}, \frac{m_s}{m_b}$$

B. Neutrino mixing angles much larger than quark mixings: e.g.

$$V_{23}^l \approx 0.7 \gg V_{23}^{CKM} \approx 0.04 \text{ etc.}$$

C. Quarks and leptons so different- are they unifiable ?

Why

$$m_\nu \ll m_{q,l}$$

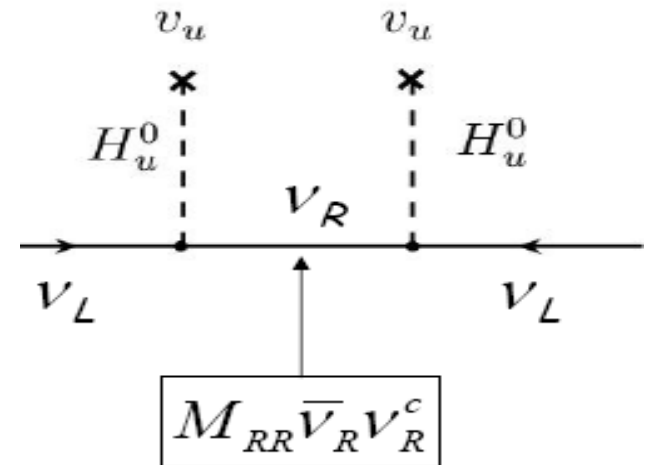
?

Seesaw Paradigm

- Add right handed neutrinos N_R to SM with Majorana mass:

$$L_Y = h_\nu \bar{L} H N_R + M_R N N$$

$$m_\nu \cong -\frac{h_\nu^2 v_{wk}^2}{M_R}$$



- **Standard seesaw (Type I):**

■ Minkowski, Gell-Mann, Ramond, Slansky, Yanagida, Mohapatra, Senjanovic, Glashow

Neutrino mass \rightarrow new symmetry of Nature: B-L

- Why B-L Symmetry ?
- Seesaw scale M_R breaks this symmetry
- The question is why $M_R \ll M_{\text{Pl}}$?
- Having a B-L symmetry explains this.
- Is it a global or local symmetry ?
- Most likely local since adding RH nu's to SM makes B-L local sym. $\rightarrow Z'$

Need to learn:

What is the B-L Scale M_R ?

- Two extreme cases: $m_D = h\nu_{wk}$
- $m_D \approx m_t$ (suggested by Q-L unif as in GUTs) leads to $M_R \approx 10^{14} GeV$ -SCALE CLOSE TO GUT SCALE-
- Corresponding theory of seesaw is: **SO(10)** :
- $m_D \approx m_e$ Scale much lower $M_R \approx TeV$;
- Theory is left-right model based on $SU(2)_L \otimes SU(2)_R \otimes U(1)_{B-L}$ motivated independently
Considerations e.g. parity, CP etc.
Physics accessible to LHC, $\beta\beta_{0\nu}$ decay etc.

Double (Inverse) Seesaw

- **Low scale seesaw with no small couplings:**
- RH neutrinos + 3 additional gauge singlet fermions \rightarrow 3x3 neutrino matrix:

$$\begin{pmatrix} 0 & hv_{wk} & 0 \\ hv_{wk} & 0 & M \\ 0 & M & \mu \end{pmatrix} \quad m_D = hv_{wk} \quad (\text{RNM, Valle, 86})$$
$$m_\nu \cong -m_D^T M^{-1} \mu M^{-1} m_D$$

- μ determines **both scale and flavor** structure;
- m_D decoupled from scale unlike type I; Seesaw scale can be in TeV range without small Yukawas

Second Challenge for theory:

Large lepton mixings:

- **Could they be hints of new symmetries for leptons:**
- **(i) Near maximal** θ_{23} : very suggestive of $\mu - \tau$ exchange sym. for neutrino matrix: **How maximal?**

(ii) Solar angle $\theta_{12} \approx 35^\circ$
suggests tribimaximal
scheme if $\theta_{13} = 0 \rightarrow$

Wolfenstein; Harrison, Perkins, Scott; Xing,

$$V = \begin{pmatrix} \frac{\sqrt{6}}{3} & \frac{\sqrt{3}}{3} & 0 \\ -\frac{\sqrt{6}}{6} & \frac{\sqrt{3}}{3} & \frac{\sqrt{2}}{2} \\ \frac{\sqrt{6}}{6} & -\frac{\sqrt{3}}{3} & \frac{\sqrt{2}}{2} \end{pmatrix}$$

The Neutrino Matrix:

- Flavor of the Neutrino flavor research

Find A_F

Generic mass matrix (NH) $\varepsilon_i \approx \lambda_{Cabibbo} \ll 1$

θ_{23} NEAR MAXIMAL

θ_{23} maximal

$$\begin{pmatrix} \varepsilon_5^{n \geq 1} & \varepsilon_4 & \varepsilon_3 \\ \varepsilon_4 & 1 + \varepsilon_1 & -1 \\ \varepsilon_3 & -1 & 1 + \varepsilon_2 \end{pmatrix} \rightarrow \begin{pmatrix} \varepsilon_5^{n \geq 1} & \varepsilon_3 & \varepsilon_3 \\ \varepsilon_3 & 1 + \varepsilon_1 & -1 \\ \varepsilon_3 & -1 & 1 + \varepsilon_1 \end{pmatrix} \xrightarrow{\text{TBM}} \begin{pmatrix} \varepsilon_1 & \varepsilon_3 & \varepsilon_3 \\ \varepsilon_3 & 1 + \varepsilon_1 & -1 + \varepsilon_3 \\ \varepsilon_3 & -1 + \varepsilon_3 & 1 + \varepsilon_1 \end{pmatrix}$$

5 parameters

$\rightarrow \mu \leftrightarrow \tau$ sym.

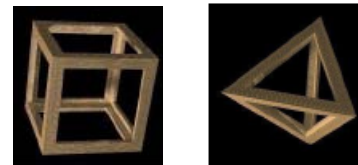
3 param.

2 param.

Understanding the pattern

- GUTs: $SO(10)$
- Family symmetries motivated by TBM:

$$S_{2(\mu-\tau)} \subset S_3, S_4, A_4, Z_2, \Delta(3n^2), \dots$$



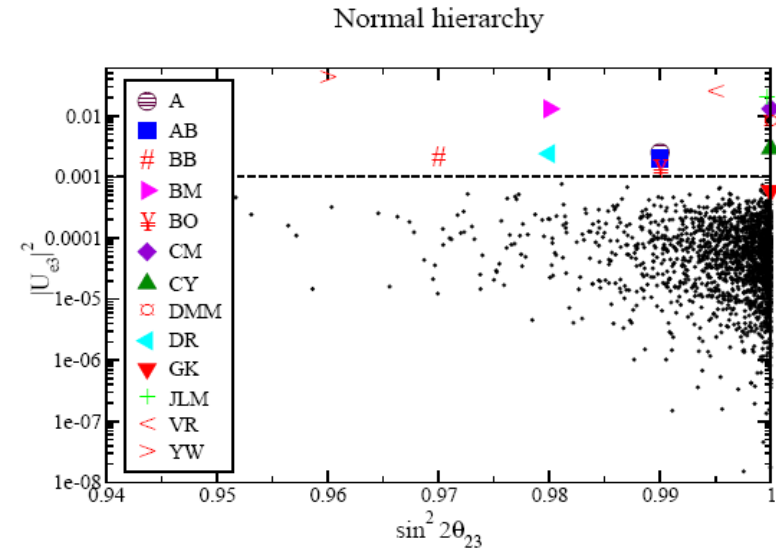
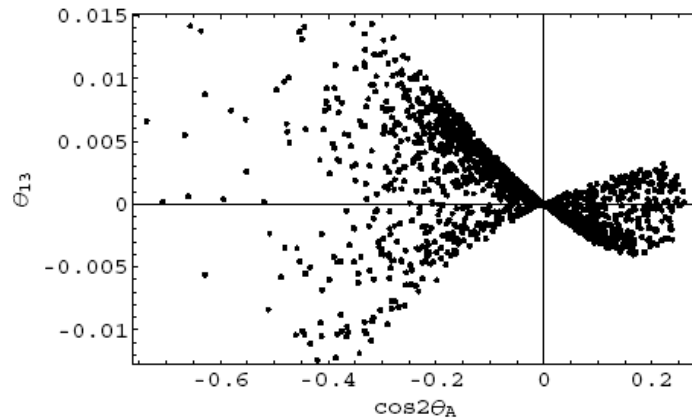
- Non-zero θ_{13} will provide important clue about new physics- is it symmetry + corrections or perhaps TBM an accident ?

(For extensive references, see G. Altarelli, Fermilab Neutrino summer school lectures;
Talks at this school by M. Chen, R. Volkas, F. Feruglio)

θ_{13} : Discriminator between Symmetry vs GUTs

■ **mu-tau sym**

TBM (Albright, Rodejohann)



- correlation with atm mixing –
- GUT predictions generally larger than 0.03.

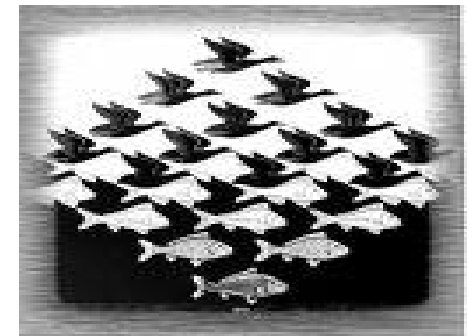
NEUTRINO MASS :A SIGNAL OF GRAND UNIFICATION ?

Grand unification hypothesis: all forces and all matter become one at high energies no matter how different they are at low energies.

Leptons →

quarks →

become same.



- Explains charge quantization;
- High scale goes well with ideas in cosmology ;
- Goes well with high scale version of seesaw.

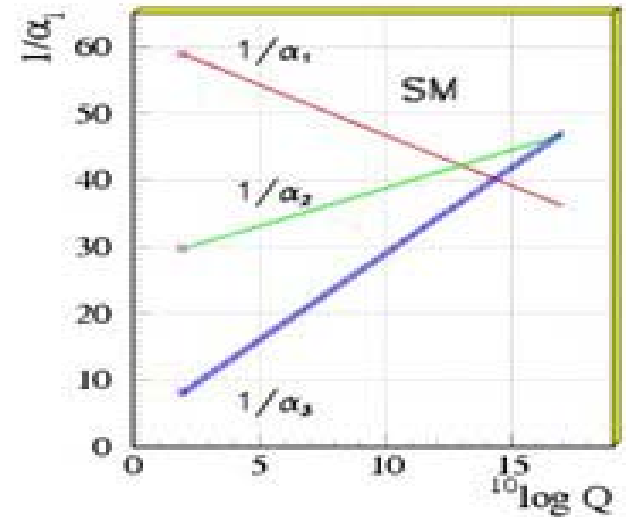
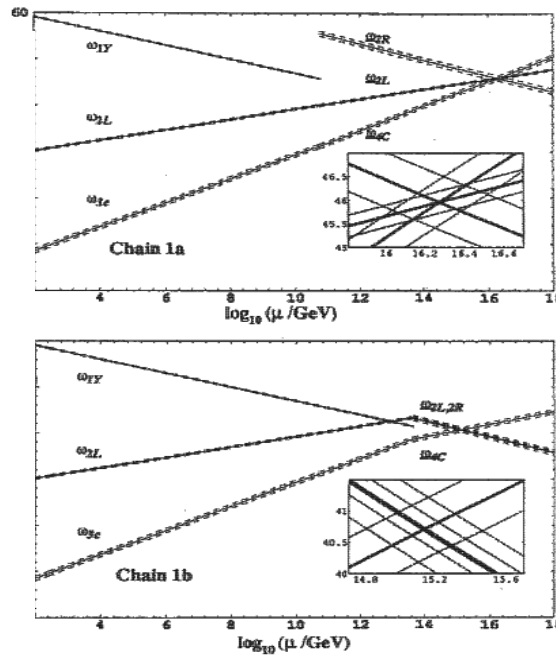
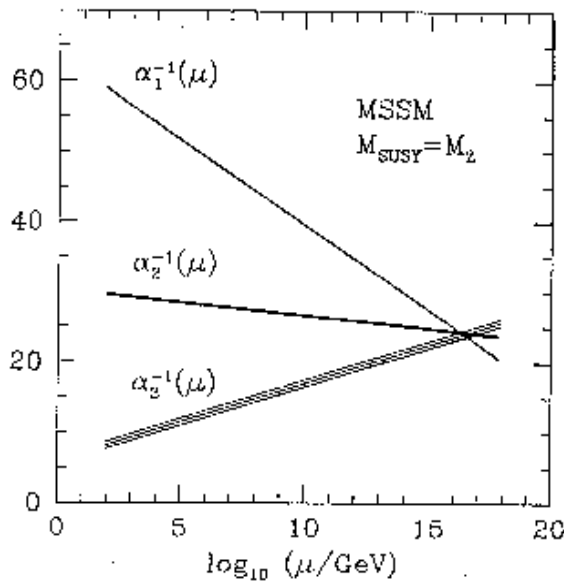
Some examples:

SUSY

Non-SUSY SO(10)

SM

with seesaw



Scale of grand unification $\sim 10^{16} \text{ GeV}$

Which GUT group ?

- Two simplest are: SU(5) and SO(10):

- (i) SU(5):**

minimal:

➤ Fermions: $5 = \begin{pmatrix} d^c \\ d^c \\ d^c \\ \nu \\ e^- \end{pmatrix}$ and $10 = \begin{pmatrix} 0 & u_3^c & -u_2^c & u_1 & d_1 \\ & 0 & u_1^c & u_2 & u_3 \\ & & 0 & u_3 & d_3 \\ & & & & e^+ \\ & & & & 0 \end{pmatrix}$

➤ : Higgs $5 \oplus \bar{5} \oplus 24$.

➤ Predicts: at M_U , $m_b = m_\tau$; very good prediction

Also predicts $m_s = m_\mu$; $m_d = m_e$; **VERY BAD PREDICTION!!**

➤ No explanation of neutrino mass:



Why SU(5) not satisfactory

- Minimal model ruled out by proton decay !
- Not predictive for neutrinos- so no advantage of GUTs except scale !
- However one nice feature: $m_b = m_\tau$

SO(10)-Just right for neutrinos

- Minimal GUT group with complete fermion unification (per family) is SO(10)-its spinor rep contains all 16 needed fermions (including RH nu) in single rep.

$$\begin{pmatrix} u & u & u & \nu \\ d & d & d & e \end{pmatrix}_{L,R}$$


- Georgi; Fritsch, Minkowski (74)
- **Contains B-L needed to understand why $M_R \ll M_{\text{Planck}}$.**
- **B-L if properly broken also allows a naturally stable dark matter in MSSM.**
- **Also helps proton decay problem.**

Appraising SO(10) as a theory of neutrinos

- Quark lepton unif. means:

$$M_d = M_l + \delta M \quad \text{and} \quad M_u = M_{\nu^D} + \delta M'$$

With δM small.

- This means quark and lepton mixings are similar – Disaster  !!
- Most models keep breaking symmetry till they get $\delta M_{\nu^D} \gg M_{u,d,l}$ and one gets large nu-mixings and a model. What trace is left of SO(10) ?

One exception !!



Minimal Predictive $SO(10)$

- Minimal model: $10 + 126 + \dots$ (Babu, Mohapatra, 93)
- Gives naturally stable dark matter without additional assumption.
- Relates RH neutrino spectrum to charged fermion spectrum reducing seesaw parameters; i.e. RH mass M_N has similar hierarchy as m_D
- Consequently type I inadequate: a new possibility emerges within the model.

How does it work ?

- {126}-Higgs relates nu matrix \mathbf{A}_F to quark-lepton flavor

$$M_\nu \cong c(M_d - M_l)$$

(Bajc, Senjanovic, Vissani'02)

- Even though quark and lepton masses are strongly hierarchical, due to $m_b \approx m_\tau$, M_ν becomes less so and gives

$$M_\nu = m_b c \lambda^2 \begin{pmatrix} \lambda^2 & \lambda^2 & \lambda \\ \lambda^2 & 1 + \lambda & 1 \\ \lambda & 1 & 1 \end{pmatrix}$$

$\lambda =$ Cabibbo angle.

(Goh, RNM, Ng'03)



Predictions

- Large solar and near maximal atmosph. mixing; diluted mass hierarchy, large θ_{13}
- Predictions: (qualitatively work very well.)
- ★ θ_{12}, θ_{23} large
- ★ $\theta_{13} \approx \lambda$
- ★ $\frac{m_{solar}}{m_{atmos}} \sim \lambda$ (Diluted hierarchy)

A quantitative model that works: Improved SO(10)

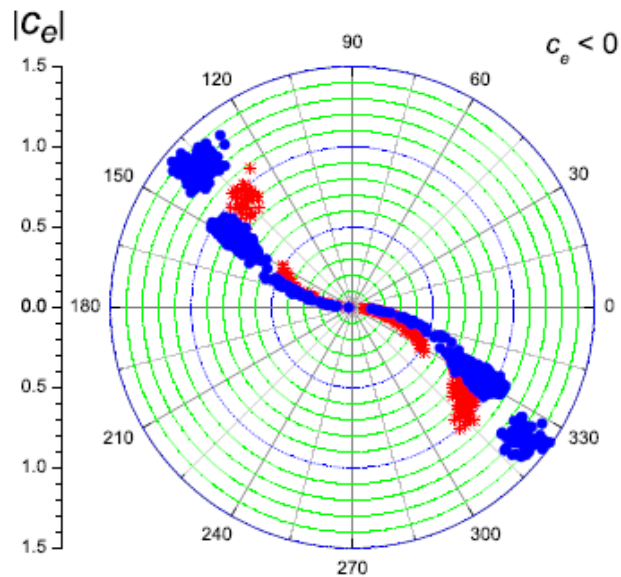
- **10+126** model-hard to include CKM CP violation !
- **Requires cancellation for proton decay !**
- **10+120+126** model with spontaneous CP **solves CP problem, proton decay problems** while keeping neutrino sector predictive; :(Dutta, Mimura, RNM,2005,06,07)
- **Solution to proton decay dictates flavor texture:**

$$h_{10} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}; \quad h_{126} = \begin{pmatrix} 0 & 0 & \lambda^3 \\ 0 & \lambda^2 & \lambda^2 \\ \lambda^3 & \lambda^2 & \lambda^2 \end{pmatrix};$$
$$h_{120} = \begin{pmatrix} 0 & \lambda^3 & \lambda^3 \\ -\lambda^3 & 0 & \lambda^2 \\ -\lambda^3 & -\lambda^2 & 0 \end{pmatrix};$$

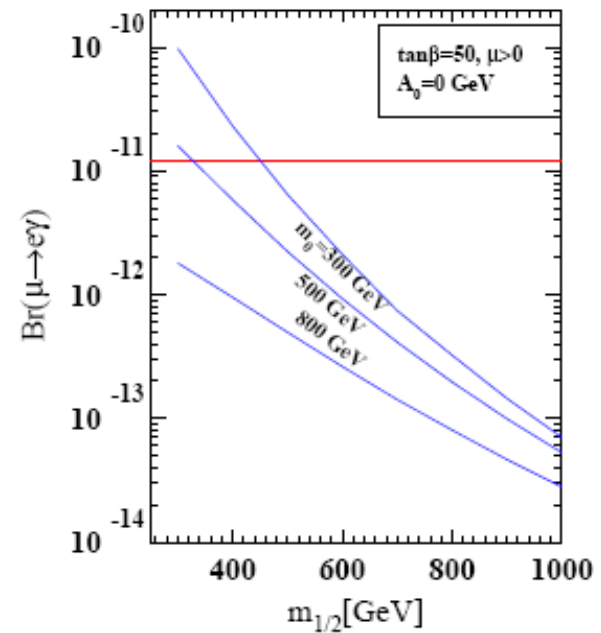
- **Predicts** $\theta_{13} \geq 0.06$

Other Predictions:

Dirac Phase:



$\mu \rightarrow e + \gamma$



$$B(\mu \rightarrow e + \gamma) > 10^{-14}$$

within range of





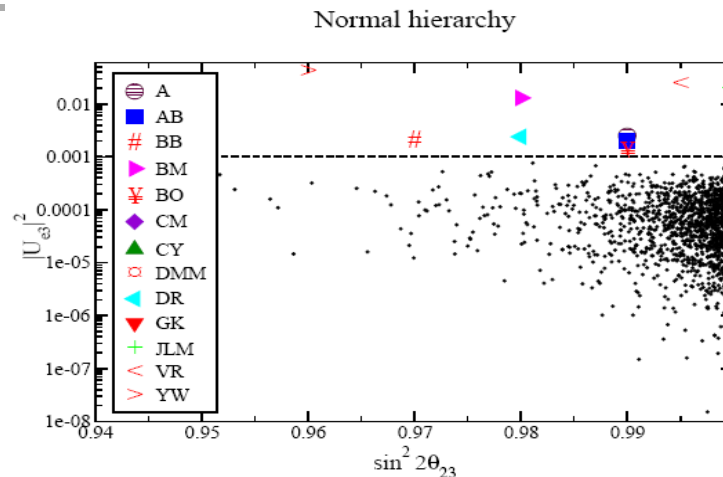
Further work on $SO(10) + 126$ models

- Fukuyama, Okada
- Goh, RNM, Ng
- Babu, Macesanu
- Bertolini, Malinsky, Frigerio
- Bertolini, Malinsky, Schwetz
- Aulakh, Bajc, Melfo, Senjanovic, Vissani
- Fukuyama, Okada, Kikuchi, Melajnac, Iljakovic
- Aulakh, Giridhar
- Dutta, Mimura, RNM
- Grimus, Kuhboch
- Aulakh, Garg
- Joshipura, Kodrani, Patel

Neutrino SUSY GUT summary:

PREDICTIONS

- Large θ_{13} -;



- $\mu \rightarrow e + \gamma$ within MEG range

ISSUES

- (i) True test of GUTs proton decay $\tau \approx 10^{35} - 10^{36}$ yrs
- (ii) How to suppress dim 5 planck induced operator for proton decay: $\{\psi_m(16)\}^4 / M_{Pl}$; strength has to be less than 10^{-7}



Reasons to consider TeV scale Seesaw:

- **Very hard to test high scale seesaw models !!**
- **Understanding the origin of matter within seesaw—**

Seesaw and Origin of Matter:

- One advantage of seesaw is the possibility to understand origin of matter, using RH neutrino seesaw couplings. Leptogenesis

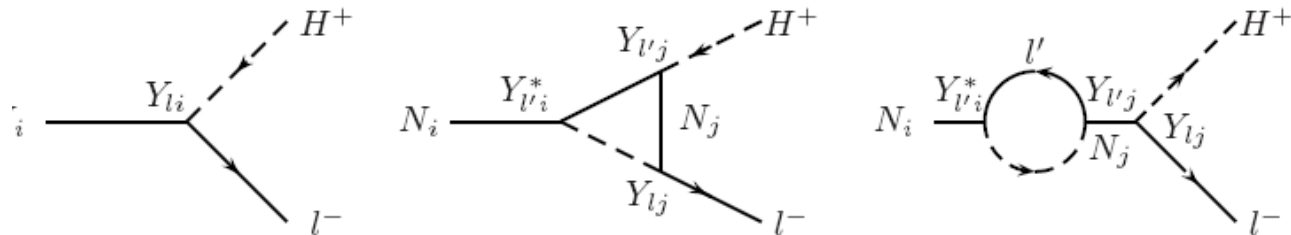
- Proposal: Heavy ν_R decays:

$$\begin{array}{l} \nu_R \rightarrow L + H \\ \nu_R \rightarrow \bar{L} + \bar{H} \end{array} \quad \begin{array}{l} R = (1 + \varepsilon) \\ \bar{R} = (1 - \varepsilon) \end{array}$$

- Generates lepton asymmetry which gets converted to baryons via sphaleron interactions; (Fukugita, Yanagida'86);

Two kinds of leptogenesis

Diagrams:



- Two classes of models depending on RH masses
- High Scale leptogenesis**: Adequate asymmetry; lightest RH nu $M \geq 10^9 \text{ GeV}$ for hierarchical RH nu's.
(Buchmuller, Plumacher, di Bari; Davidson, Ibarra)
- Resonant leptogenesis**: **degenerate** N 's, self energy diagram dominates: $\sim \frac{1}{M_i^2 - M_j^2 + M\Gamma}$
; Resonance $M_i \cong M_j$; **works for all B-L scales.**

(Liu, Segre'94; Covi et al. Flanz et al.'95; Pilaftsis'97)

ISSUES WITH HIGH SCALE SUSY LEPTOGENESIS

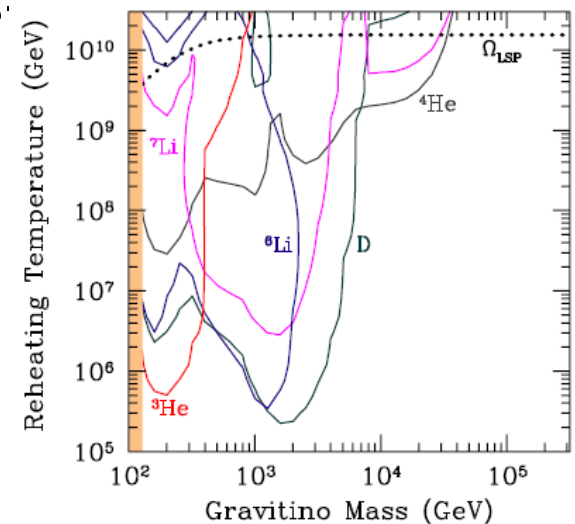
- Adequate baryogenesis requires the lightest RH neutrino mass $M_N \geq 3 \times 10^9$ GeV (Davidson, Ibarra)

- **Problem for supersymmetric models:**

they have gravitinos with TeV mass that are produced during inflation reheat along with all SM particles.

- **If stable Will overclose the universe for $T_R > 10^9$ GeV.**
- **If unstable, live too long -effect the success of BBN.**

(Kawasaki, Kohri, Moroi, Yatsuyanagi, 2008)



Could Seesaw be a TeV scale phenomenon ?

- If so it is very likely that there is a new gauge symmetry of Nature beyond SM that couples to RH neutrinos:
- The symmetry could involve either an extra Z' as in $SU(2)_L \times U(1)_{I_{3R}} \times U(1)_{B-L}$ or both W_R and Z' as in left-right models.
- If masses are in few TeV range, production and decays at LHC could provide evidence of their existence via Z' decays $Z' \rightarrow NN; N \rightarrow lH$ and W_R decays: $W_R^+ \rightarrow l^+ N ; N \rightarrow lH$

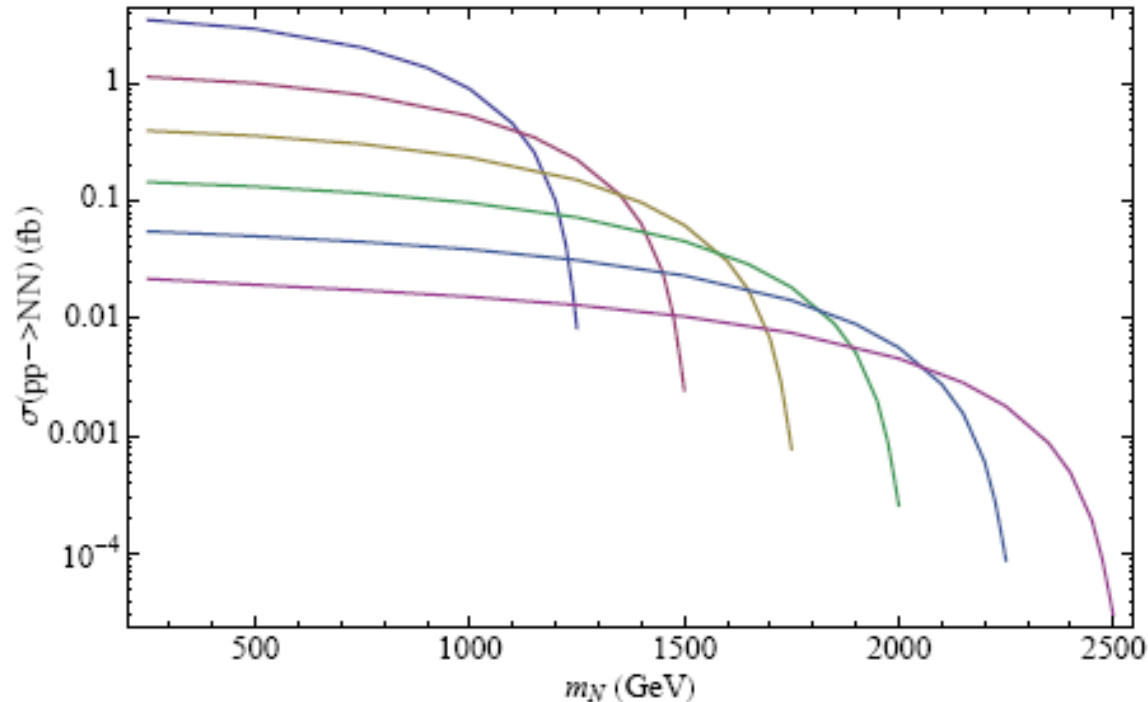


Current bounds on W_R , Z'

- **Collider limits on W_R :** around 780
- **Low energy limits:** K-K-bar, CPV, edm etc:
 W_R mass > 2.5 TeV (Zhang, An, Ji, RNM, 2008; adopted by PDG)
- **Limits from Neutrinoless double beta decay+ vacuum stability:**
 W_R mass > 1.5 TeV. (Paes, Simkovic talk)
- **Limits are lower for SUSYLR due to sparticle FCNC effects.** (Zhang, An, Ji 2008)
- **Z' mass bound: > 995 GeV** (Langacker, Erler, Munir, Pena)

TeV Z' cross section at LHC

- **LHC Z' reach - 4 TeV**
- Cross section for $pp \rightarrow Z' \rightarrow NN$ ($Z' \rightarrow NN$ branching ratio $\sim 20\%$)



2.5 TeV Z'

to

5 TeV

Testing seesaw with Z' decay

- $PP \rightarrow Z' + X$; xsection for a 3 TeV Z' \sim fb
- Seesaw signal: $N = \text{Majorana}$
- $N \rightarrow e^\pm W^\mp, \bar{\nu}^{(-)} + Z \quad W \rightarrow jj, l\nu$
- Di and Multi-lepton events: ($X = jjjj$)
 $pp \rightarrow l^\pm l^\pm X, l^\pm l^\pm l^\mp + \cancel{E}, l^\pm l^\pm l^\mp l^\mp + \cancel{E}$
- Important for signal to bg: very high p_T leptons coming from N -decay; inv mass reconstruction:
 (Del Aguila, Aguilar-Saavedra; P. Perez, Han, T. Li)

Does leptogenesis work with TeV Z' and WR ?

■ Conditions:

- (i) RH neutrinos must be degenerate in mass to the level of $M_1 - M_2 \sim 10^{-10} M$ since $h \sim 10^{-5}$;
- (ii) Since there are fast processes at that temperature, the net lepton asymmetry and primordial lepton asymmetry are related by

$$\eta_B \simeq 10^{-2} \sum_{i,\alpha} \epsilon_{i\alpha} \kappa_{i\alpha}$$

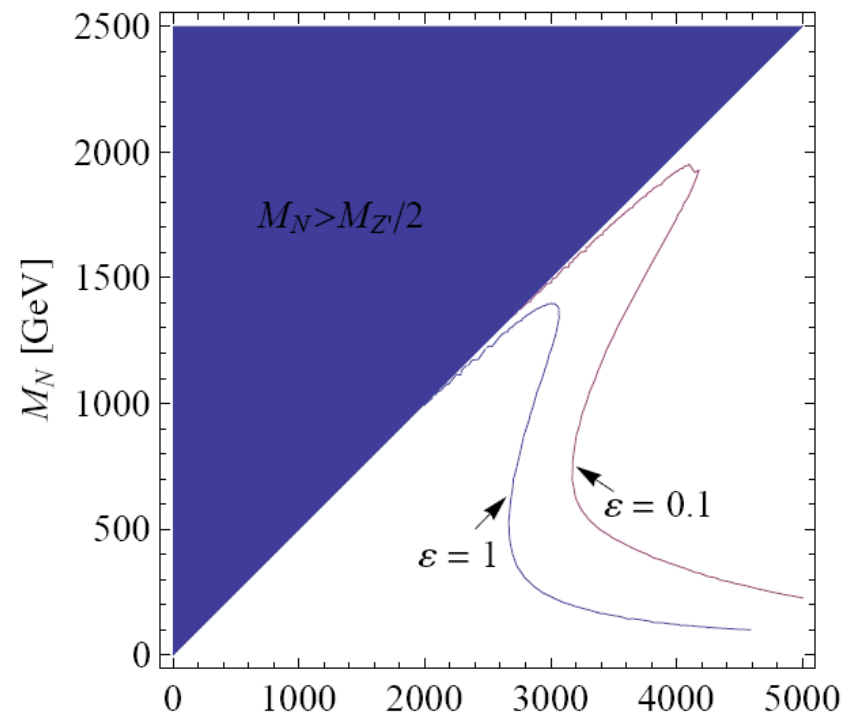
where $\kappa < 1$ - depends on Z' mediated $e^+ e^- \rightarrow NN$ and inverse decay $lH \rightarrow N$

Not clear that a TeV scale Z' is even allowed by baryogenesis due to rapid rates ?

Lower bound on Z' mass from leptogenesis

- Lower the Z' mass, faster the scattering and less the efficiency implying a lower limit on Z' mass !!

■ (BLANCHET, CHACKO, GRANOR, RNM: ARXIV:0904.2974)



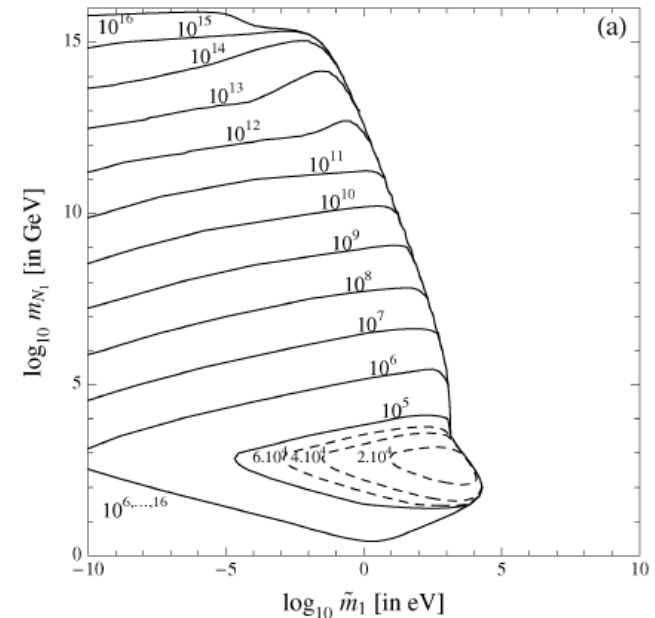
★ $M_{Z'} > 2.5 - 3.2 \text{ TeV}$ for $M_{Z'} > 2M_N$ (Accessible at LHC)

Limits on WR

- **Left-right Model:** $SU(2)_L \otimes SU(2)_R \otimes U(1)_{B-L}$
- New fast processes that erase the lepton asymmetry: $e_R + u_R \rightarrow N + d_R$

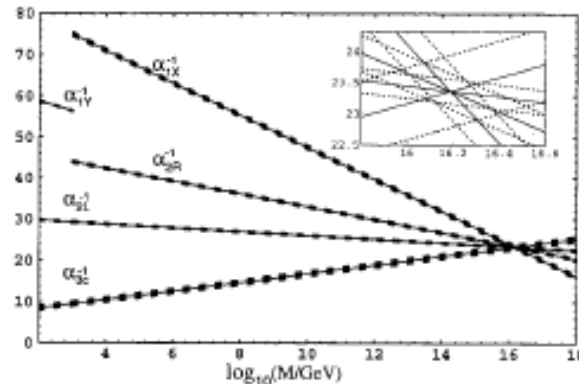
- Except when $M_{WR} > 18$ TeV,
(Frere, Hambye and Vertongen)

Sym br. to $U(1)_{I_{3R}} \times U(1)_{B-L}$
 then to SM at TeV-
 to do resonant leptogenesis.



Unification Prospects for TeV seesaw: An SO(10) possibility

- TeV scale Triplets with B-L=2 hard to unify to SUSY SO(10).
- Both for TeV Z' and WR, unification possible with B-L = 1 doublets breaking $U(1)_{B-L}$; (Deshpande, Keith and Rizzo; 93; Malinsky, Romao, Valle'05);



- This has implications for neutrino mixing:

Double seesaw for Neutrino masses

- B-L=1 breaking → inverse seesaw for neutrino masses

$$\begin{pmatrix} 0 & hv_{wk} & 0 \\ hv_{wk} & 0 & M \\ 0 & M & \mu \end{pmatrix} \quad m_\nu \cong -m_D^T M^{-1} \mu M^{-1} m_D$$

- **Unlike type I**, nu-N mixing m_D / M decoupled from neutrino mass- so can be large enhancing N-production at LHC.
- **Unlike type I**, Majorana character of RH N (amount of like sign dileptons) depends on how large $\frac{\mu}{M}$ is.
- **Unlike type I**, $\mu \rightarrow e + \gamma$ can be large without susy.
- Leptogenesis possible for $\frac{\mu}{M} > 10^{-6}$



Testing double seesaw

- Can lead to deviations from Unitarity for neutrino mixings (S.

[Antusch](#), [C. Biggio](#), [E. Fernandez-Martinez](#), [M.B. Gavela](#), [J. Lopez-Pavon](#) Goswami, Ota; Altarelli, Meloni; Malinsky, Ohlsson, Zhang, Xing;)

- **Mixing matrix**
$$N = (1 + \eta)U$$

- **-Current limits on $\eta_{\alpha\beta}$: $\eta_{\mu\tau} < 5 \times 10^{-3}$; $\eta_{e\mu} < 0.0001$ on 12-element from Lepton flavor violation.**

- **Observable oscillation effect at near detector in neutrino factories as well as far detector.**

$\eta_{\mu\tau} < 10^{-4}$ attainable from SBL with 50 GeV E.



Conclusion:

- **Need to know-Dirac vs Majorana:**

- ★ **What may we have learnt ?**

- Majorana nu \rightarrow seesaw good paradigm!!
- \rightarrow may explain the origin of matter

- ★ **What we need to know ?**

Scale of new physics (e.g. B-L sym.)- **GUT vs TeV scale** ? May give a hint as to whether large mixing is from dynamics or symmetry.

- ★ mass ordering, θ_{13} to understand \mathbf{A}_F

Large mixing from Dynamics:

Simple hierarchy to Double hierarchy

- A possibility within type I seesaw:

$$m_\nu = -m_D \frac{1}{M_N} m_D^T \quad \mathbf{m}_D \text{ hierarchical in most models:}$$

- $m_\nu \sim \begin{pmatrix} \varepsilon_5^{n \geq 1} & \varepsilon_4 & \varepsilon_3 \\ \varepsilon_4 & 1 + \varepsilon_1 & -1 \\ \varepsilon_3 & -1 & 1 + \varepsilon_2 \end{pmatrix}$ Not very hierarchical

- That means M_N could be "doubly" hierarchical:

(Altarelli, Feruglio, Masina; '03; He, Law, Volkas'08)

- A recent example: $\mathbf{m}_D = \text{diag}(m_e, m_\mu, m_\tau) \rightarrow$

$$(M_N)_{\text{NH}} = \frac{\lambda'}{6p^2 m_0} \begin{pmatrix} 4m_e^2 & -2m_e m_\mu & -2m_e m_\tau \\ -2m_e m_\mu & m_\mu^2 & m_\mu m_\tau \\ -2m_e m_\tau & m_\mu m_\tau & m_\tau^2 \end{pmatrix}$$

(Adulpravitchai, Lindner, Merle, RNM; arXiv:0908.0470)

Hierarchy Dilution by cancellation

- Suppose at very high scale, there is a sum-rule:

$$M_\nu \cong c(M_d - M_l)$$

- Since at high scale, $m_b \approx m_\tau$ most hierarchical term cancels out:

$$M_{d,l} = m_{b,\tau} \begin{pmatrix} \sim \lambda^4 & \sim \lambda^3 & \sim \lambda^3 \\ \sim \lambda^3 & \sim \lambda^2 & \sim \lambda^2 \\ \sim \lambda^4 & \sim \lambda^2 & 1 \end{pmatrix}$$

$\lambda =$ Cabibbo angle.

- Sum rule** \rightarrow

$$M_\nu = m_b c \lambda^2 \begin{pmatrix} \lambda^2 & \lambda^2 & \lambda \\ \lambda^2 & 1 + \lambda & 1 \\ \lambda & 1 & 1 \end{pmatrix}$$

- (Bajc, Senjanovic, Vissani; Goh, RNM, Ng)