



40 MCi Tritium Source for Experiments with Low Energy Neutrinos

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Neutrinos in Cosmology, in Astro-, Particle- and Nuclear Physics,

16-24 September 2009, Erice, Sicily

"Neutrinos, they are very small.
They have no charge and have no mass.
And do not interact at all..."

John Updike,
American writer, non-physicist

"In some sense, tritium in nuclear and weak
interaction physics is analogous to the
drosophila fly in genetic studies"

V.M. Lobashev,
Russian physicist,
Leader of Troitsk neutrino mass experiment



Why?

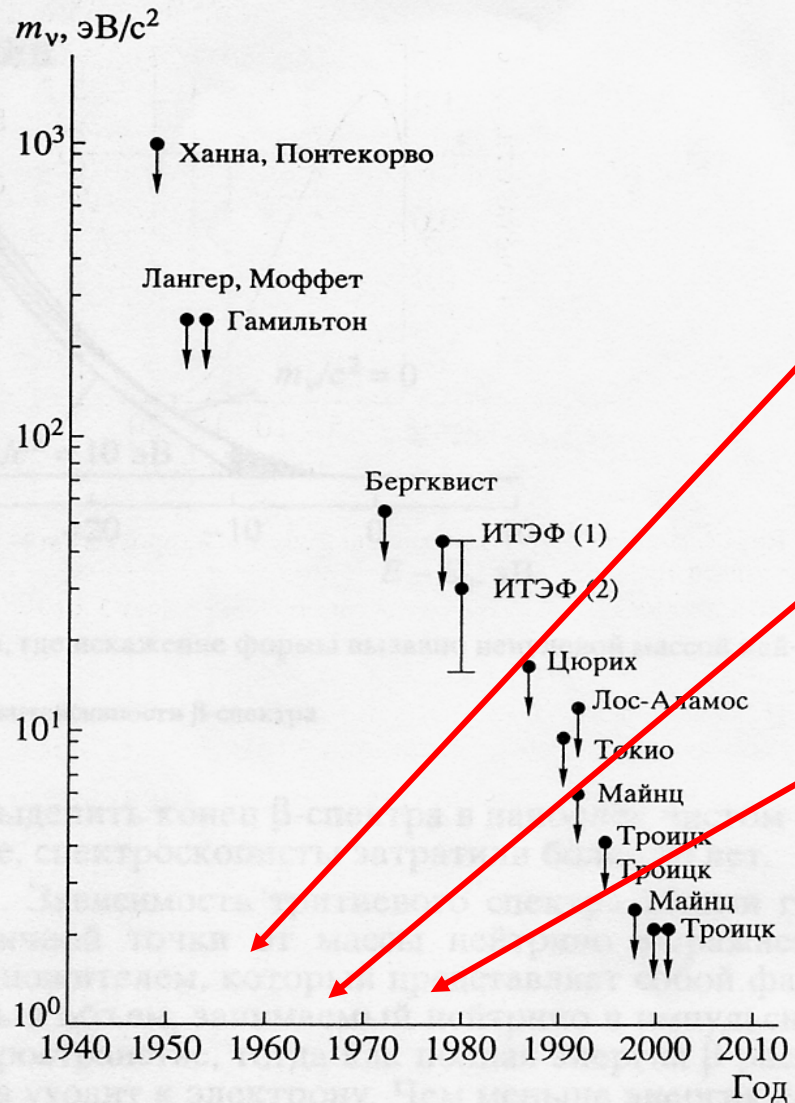
Triton is the lightest radioactive nucleus



The first estimation of the neutrino mass was done in 30-s, when a heavy hydrogen isotope – tritium – was discovered with a small decay energy. The fact of the existence of this allowed transition pointed at the extreme smallness of the neutrino mass:

$$m_\nu < 10 \text{ keV} < 0.02 m_e$$

Search of the Neutrino Mass in the Tritium β -Decay



Chronology

Theory of two-component neutrino (1958) - $m_\nu = 0$

Hypothesis of neutrino oscillations (1960-s) - $m_\nu \neq 0$

Hypothesis of grand unification of all interactions (1972) $m_\nu \neq 0$

$m_\nu < 2.2 \text{ eV}/c^2$

All experiments study the electron spectrum shape

Neutrino Magnetic Moment

Standard Model

Dirac neutrino

Majorana neutrino MM (some MSM extensions)

theory

$$\mu_\nu \cong 3(m_\nu/1\text{eV}) 10^{-19} \mu_B$$

$$\mu_\nu < 10^{-14} \mu_B$$

$$\mu_\nu \sim 10^{-(10\div 12)} \mu_B$$

Astrophysics & Cosmology (model dependent limit)

$$\mu_\nu \leq (1\div 3) \times 10^{-12} \mu_B$$

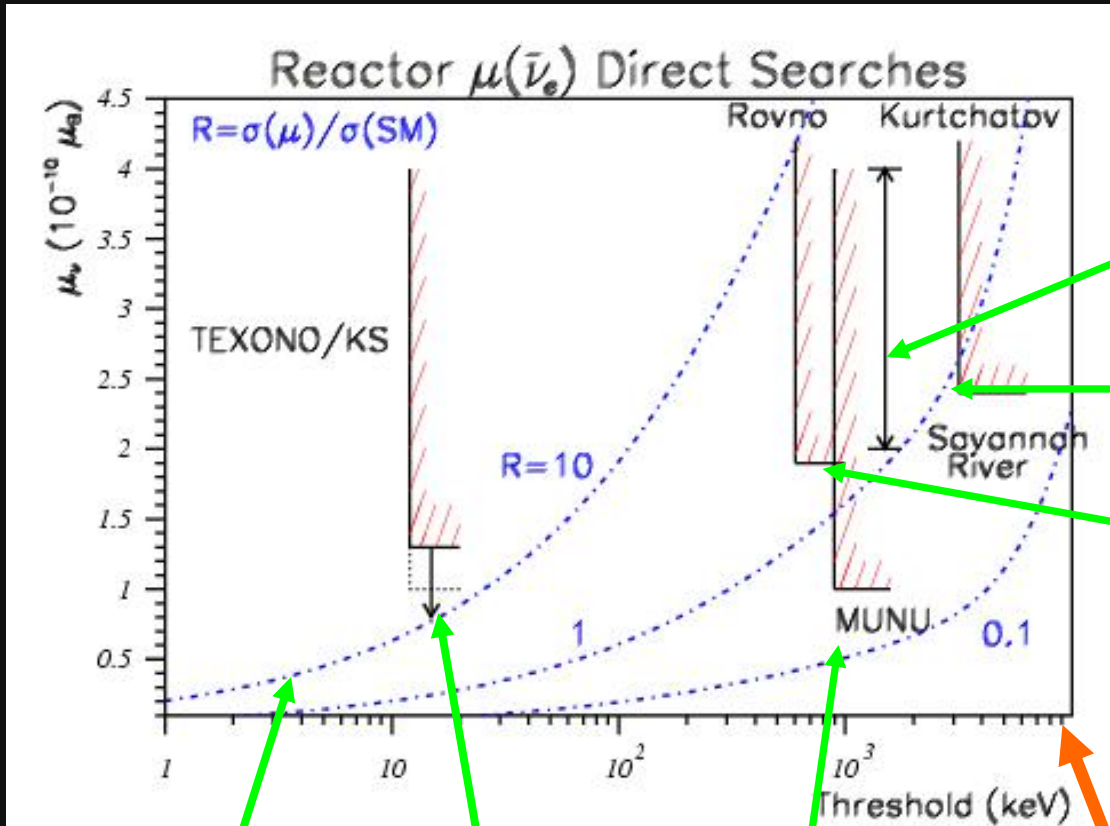
Experimental laboratory limit (reactor experiment 2009)

$$\mu_\nu \leq 3.2 \times 10^{-11} \mu_B \quad 90\% \text{CL}$$

A 40 MCi tritium source experiment

$$\mu_\nu \leq \underline{2.5 \times 10^{-12}} \mu_B$$

Neutrino Magnetic Moment Measurement



1976
Reines, et al.

1992

1993

June 2009,
GEMMA

2007,
TEXONO

2002

Antineutrino flux
 $F \approx (1 \div 5) \cdot 10^{13} \text{cm}^{-2} \cdot \text{s}^{-1}$

Detector threshold
(10 ÷ 3000) keV

Differential Cross Section of the ν -e Scattering for Reactor Antineutrino

$$\frac{d\sigma_{EM}}{dT} = \left(\frac{\mu_\nu}{\mu_B} \right)^2 \frac{\pi\alpha^2}{m_e^2} \left(\frac{1}{T_e} - \frac{1}{E_\nu} \right)$$

$$\mu_\nu = 6 \times 10^{-11} \mu_B$$

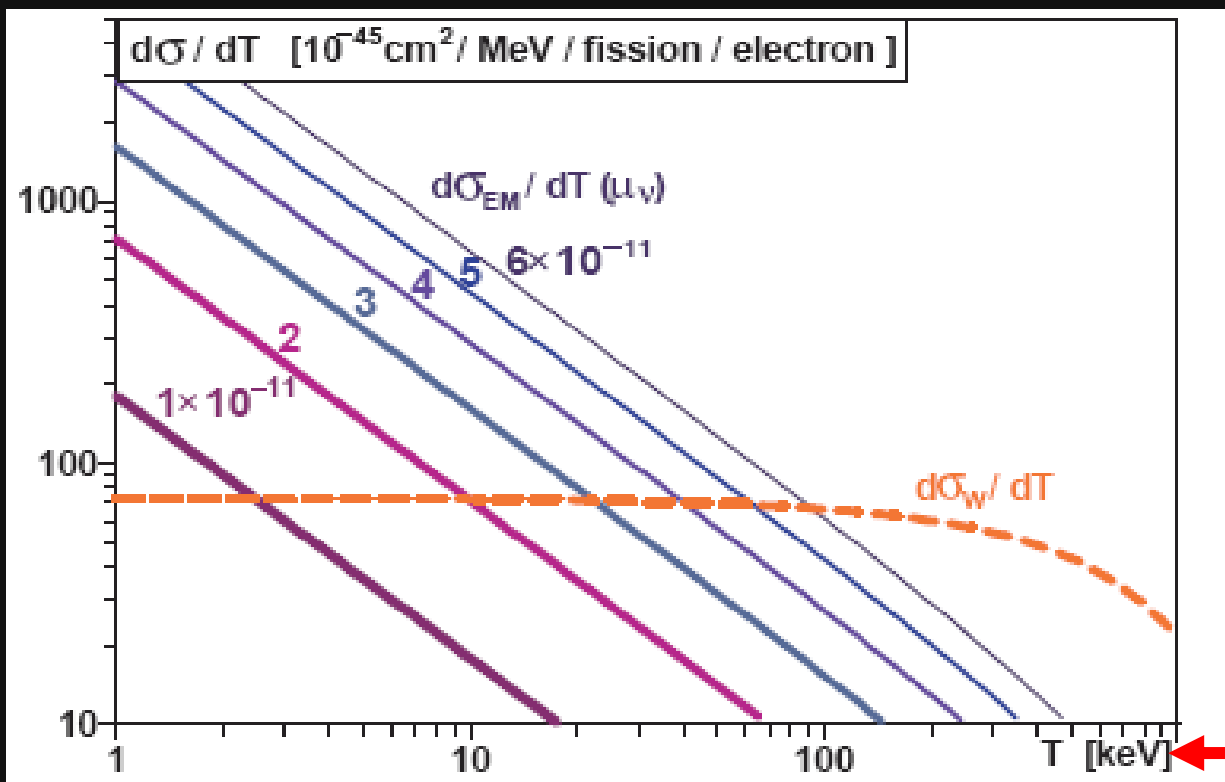
(2007)

$$\mu_\nu = 3 \times 10^{-11} \mu_B$$

(2009, GEMMA)

$$\mu_\nu = 1 \times 10^{-11} \mu_B$$

(planned)



T is the electron recoil energy

Beginning of the Story

In 1978 B. Neganov & V. Trofimov from JINR (Dubna) considered “**Possibility of producing a bulky supersensitive thermal detector at a temperature close to absolute zero**”

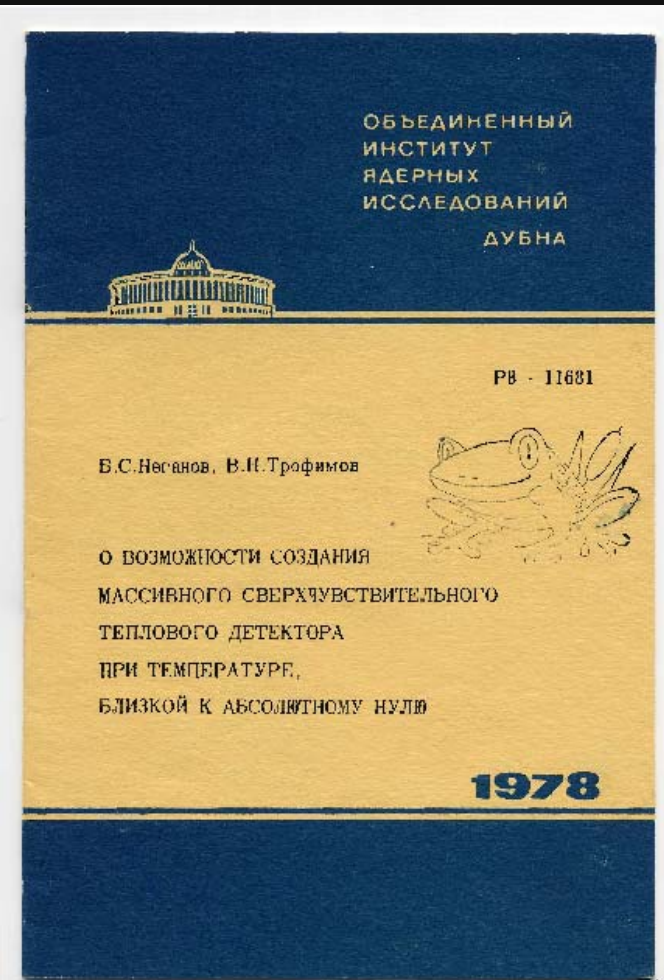


Photo of the first ^3He - ^4He dilution refrigerator $T \sim 70\text{mK}$ (B. Neganov, 1966)



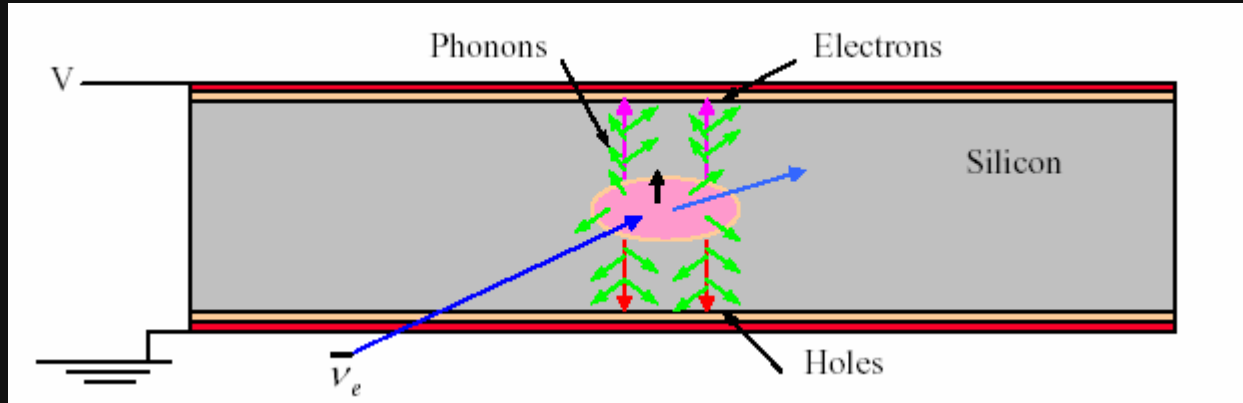
B. Neganov (center) with R. Gaitskell and D. Akerib from CDMS collaboration (JINR, 2001)

JETP Letters 28, 328 (1978)

Ionization-into-Heat Conversion (Neganov-Trofimov-Luke effect).

- ◆ Drift primary ionization in high electric field
- ◆ Measure resulting heat

Ideal gain for Silicon is $G_{th} = 1 + V_{applied}/3.8 \text{ Volts}$



Allows obtaining few keV heat release per one charge carrier in cooled semiconductor under irradiation - demonstrated by N&T at Dubna (1981) and P.Luke at LBL (1988)

Single particle detection in a bulky detector

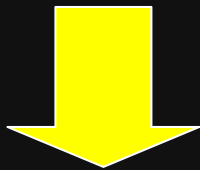
"We, experimentalists, are not like theorists: the originality of an idea is not for being printed in a paper, but for being shown in the implementation of an original experiment"

Patrick M.S. Blackett
(see at the entrance)

Proposal on Cryogenic Detection of the Neutrino Magnetic Moment with a Strong Tritium Source 1993 (J. of Low Temperature Physics)

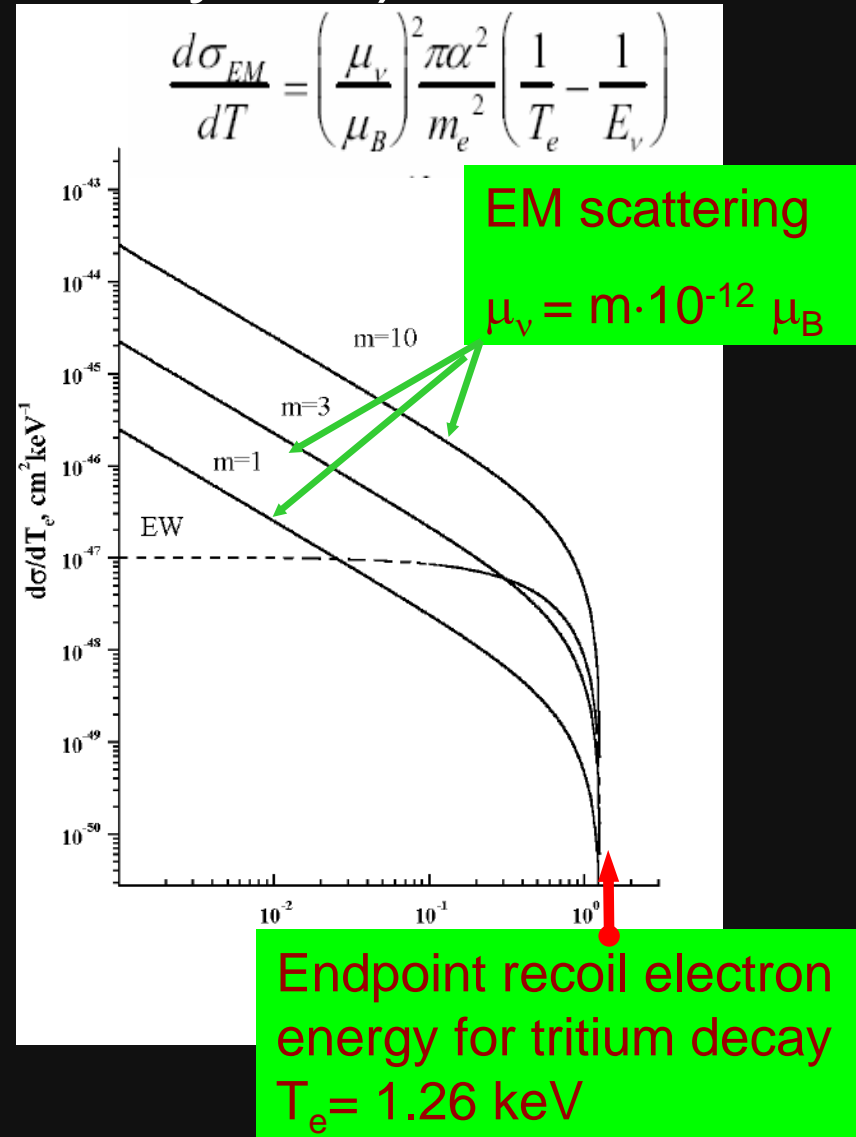
Why Tritium?

- Low decay energy
- High radioactive purity
- Long enough life-time (12.5 years)



- Low radiation background
- No passive shielding – closer to detector - bigger ν flux
- Enough time for data taking (~ 1 year)

9/27/2009



How "strong"?

Flux should exceed that from a power reactor by an order

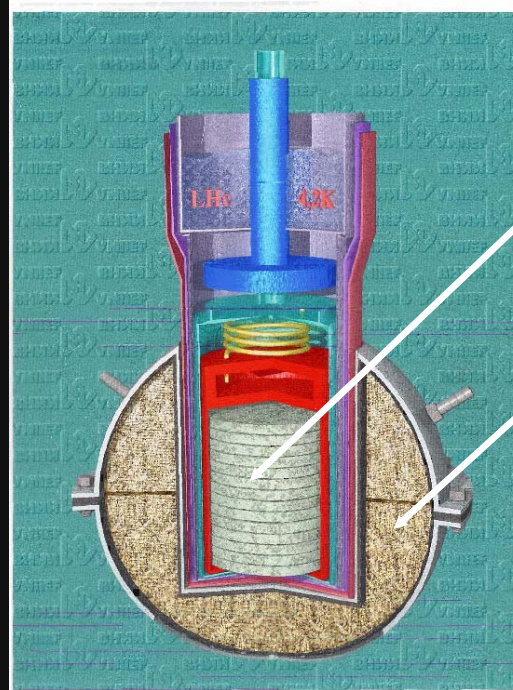
40 MCi (4 kg of ^3H)

Is it possible?

Problems to solve:

- Safety
- Long term stability
- Compactness

Original Design 1998



Cryodetectors

Tritium source



Rds-1



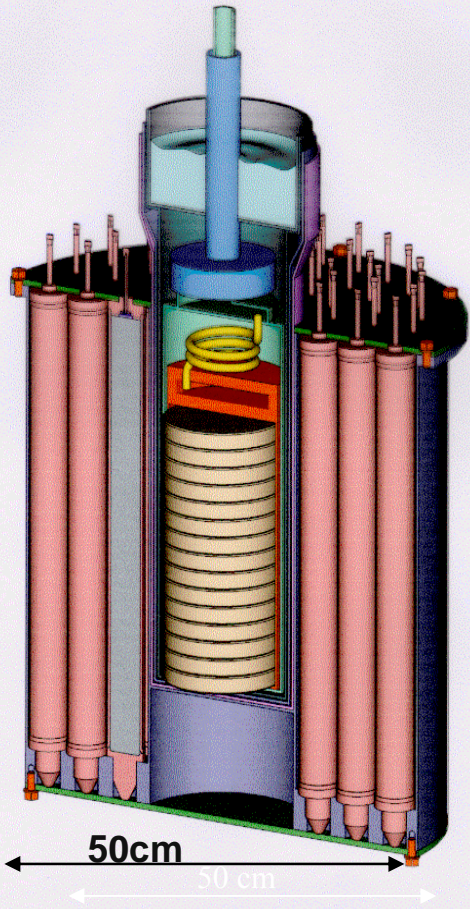
Sputnik 1
1957

9/27/2009

Metals and Alloys for Hydrogen Isotope Storage

| Materials | $P_{\text{dec.}}$ mm Hg at 25°C. | Decomposition temperature, °C | Sorption capacity H_2 , cm^3/g | Density in H_2 , cm^3/cm^3 (hydride) | Density, g/cm^3 | |
|---|--|-------------------------------------|---|---|------------------------------------|----------------|
| | | | | | metal | Hydride |
| U | 10^{-6} - 10^{-4} | 420-430 | ~141 | ~1570 | 18.7 | 11.9 |
| Ti | $\sim 10^{-9}$ | 550-620 | ~ 468 | ~ 1760 | 4.5 | 3.7-3.9 |
| Mg_2Ni | $\sim 10^{-2}$ | ~240 | ~418 | ~1074 | | |
| ZrCo | $\sim 10^{-5}$ | 340-350 | ~186 | ~1415 | | |
| LaNi_3Mn_2 | $\sim 10^{-2}$ | 270 | ~127 | | | |
| $\text{LaNi}_{5-x}\text{Al}_x$ ($x=1-1.3$) | $\sim 10^{-1}$ -10 | 180-250 | 80-100 | 550-690 | | |
| $\text{Zr}_{1-x}\text{Ti}_x\text{Co}_{0.5}\text{Ni}_{0.5}$ ($x=0.1-0.2$) | 10^{-2} - 10^{-1} | 200-300 | 160-180 | 1050-1180 | | |
| Pd | 30-50 | ~150 | ~105 | ~800 | 11.9 | |
| H liquid | | -253 | | 780 | | |
| T liquid | | -251 | | 1000 | | |

Search for the Neutrino Magnetic Moment



$\bar{\nu}_e - e$ scattering

$\bar{\nu}_e$ TRITIUM SOURCE of 40 MCi activity
(4 kg of ^3H) with a flux density of $6 \times 10^{14} \text{ cm}^{-2}\text{s}^{-1}$

ULTRA-LOW-THRESHOLD DETECTORS $E_{\text{th}} \sim 10 \text{ eV}$:

SILICON CRYODETECTOR @ $T=10\text{mK}$,
 $15 \times 100 \text{ cm}^3$, $M=3\text{kg}$,
ionization-into-heat conversion effect

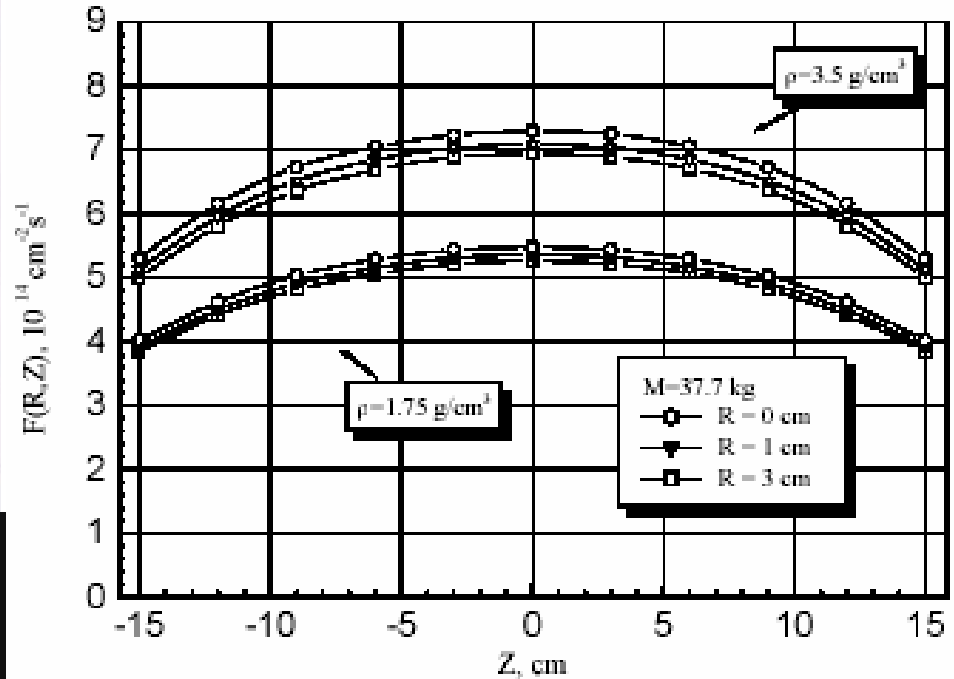
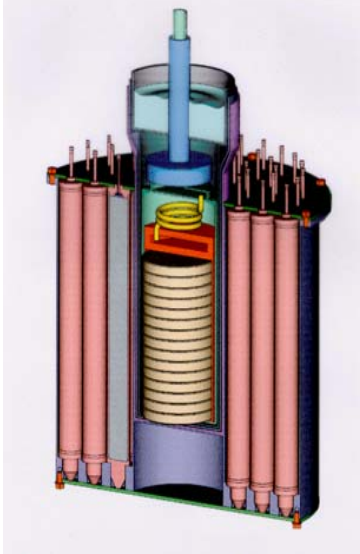
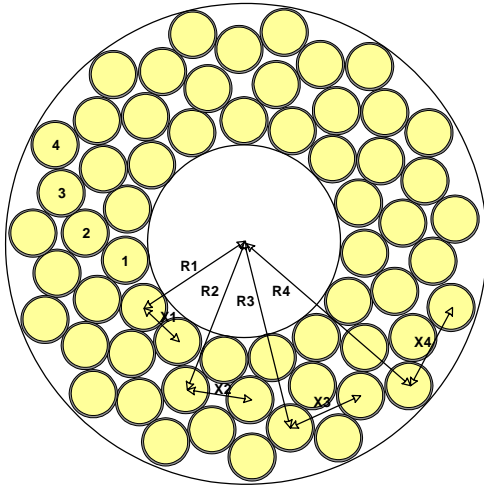
HIGH-PURITY-GERMANIUM DETECTOR

$6 \times 150 \text{ cm}^3$, $M=4.8\text{kg}$,
*internal proportional signal amplification by
avalanche multiplication in the electric field*

SENSITIVITY (95% C.L.): $\mu_{\nu} \leq 2.5 \times 10^{-12} \mu_{\text{B}}$

Conceptual layout of the
 ν -e scattering experiment
with 40 MCi tritium
source
9/27/2009

Neutrino Flux Inside the Source



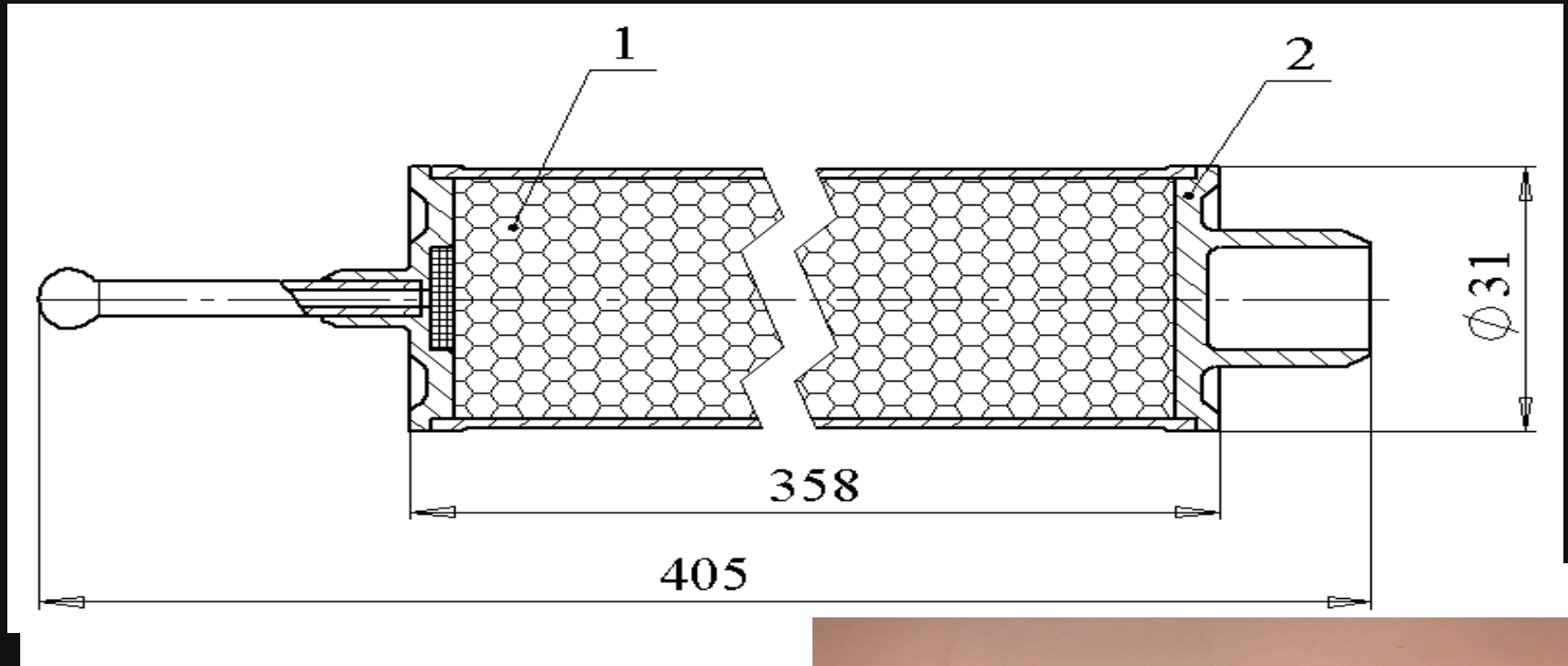
A cylinder-shaped source made of annular cells packed together and filled with titanium tritide



1/2 model

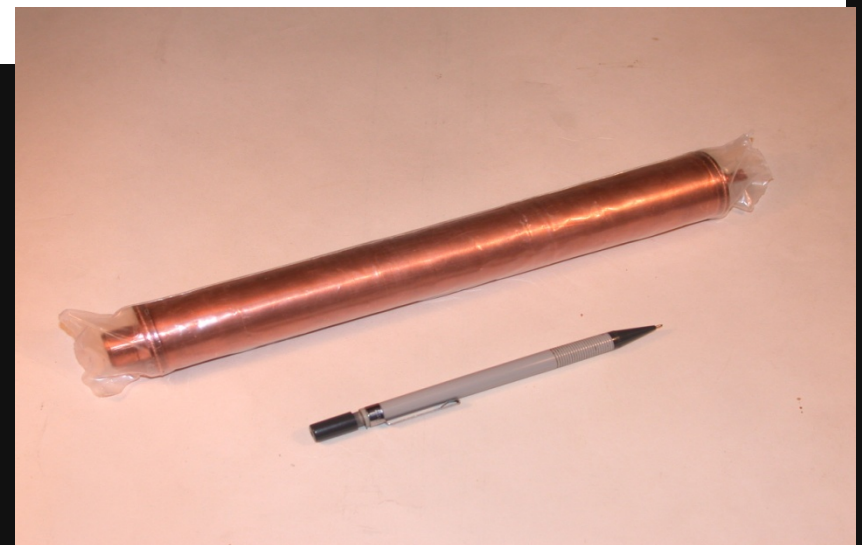
Neutrino flux density
 $F \approx 6 \cdot 10^{14} \text{ cm}^{-2} \cdot \text{s}^{-1}$ along the detector Z-axis for the multi-cylinder source design

Construction of the Tritium Cell



1 – titanium tritide

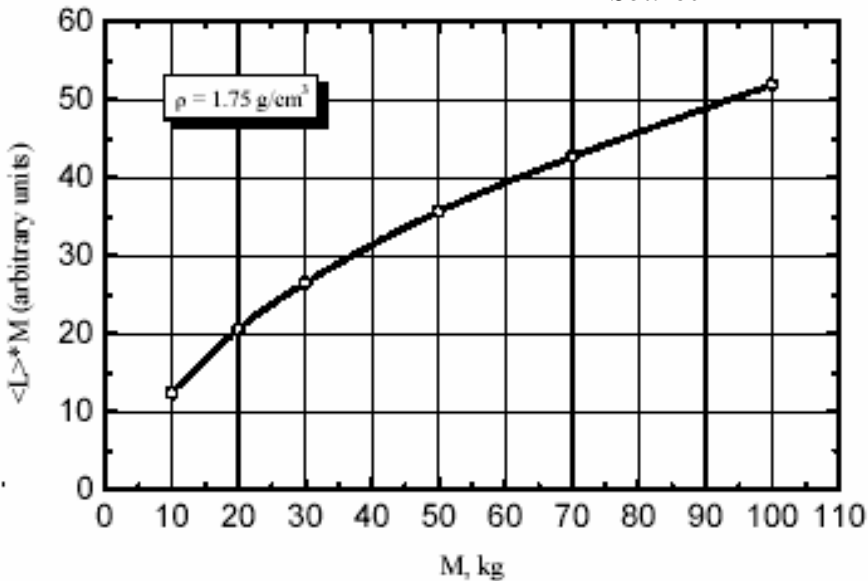
2 – copper body of the cell



How Much Tritium Can We Efficiently Add to?

Number of neutrino interactions can be increased by adding tritium mass. Source will get bigger, efficiency will be reduced.

efficiency $\langle L \rangle = \frac{\sum_{i=1}^{N_{Det}} L_i}{N_{Source}}$



- At tritium mass of 4 kg activity is $A=1.43 \cdot 10^{18}$ neutrino/s
- TiT_2 mass is 37.7 kg
- Maximum possible density of TiT_2 is $\rho_0=3.5 \text{ kg/cm}^3$
- Applicable value of the working density depends on the duration of an experiment and is limited by the allowed deformation of the cell.

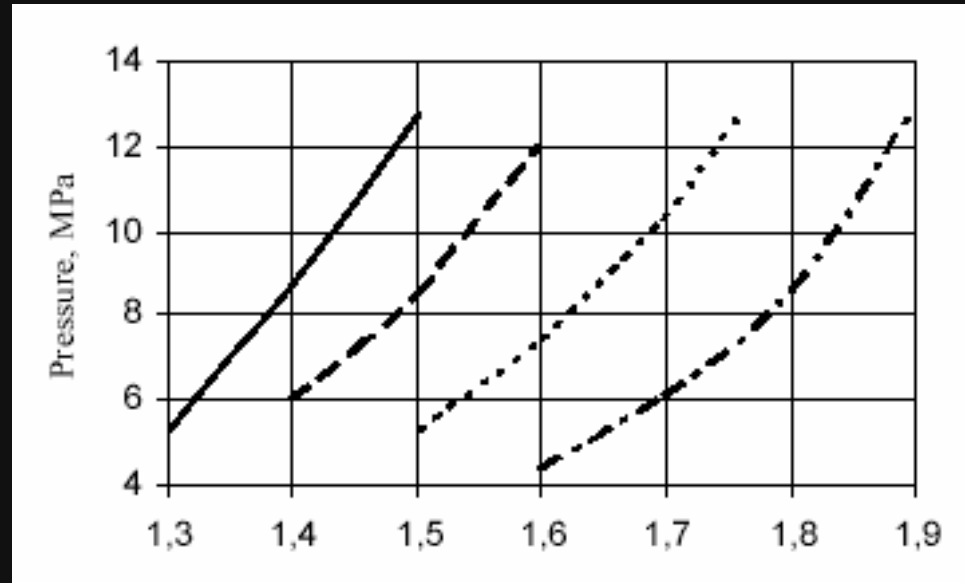
Number of neutrino interactions with detector electrons versus TiT_2 mass.

How Dense Titanium Tritide Can Be Packed?

The working substance TiT_2 swells inside the cell due to a deformation of pore structure and due to a storage of radiogenic helium.

At equal “swelling” of TiT_2 due to storage of radiogenic helium, the pressure from the titanium tritide will become smaller with increase of hydrogen content.

Maximum allowed pressure at the end of the exposure is 12 MPa – this enables safe dismounting of the cells (follows from the cell strength test and calculations).



Time dependency of the pressure in the cell versus the initial density (g/ccm) of titanium tritide:

solid – 4 years, dashed - 3 years, dotted - 2 years, dash-dotted - 1 year.

From the final design report

A cylinder-shaped multipart source construction seems to be most attractive technologically.

As the source design, a configuration of annular cells filled with TiT_2 that are stacked into a hollow cylinder is chosen.

Detector array will be mounted in the hole inside.

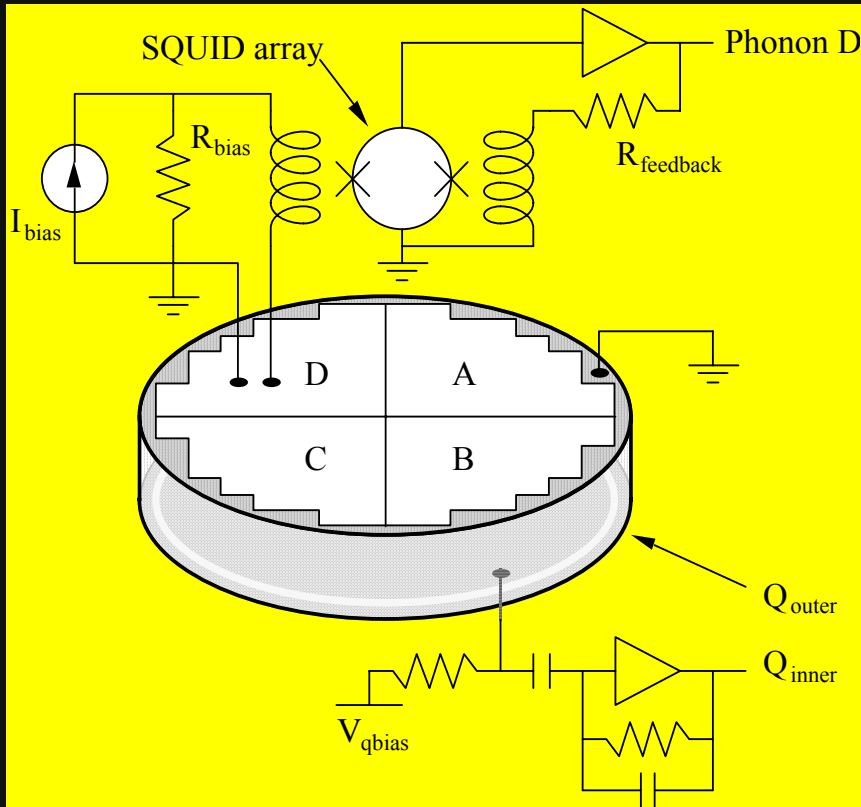
A multipart construction facilitates heat removal from and access to individual cells.

This design makes possible dealing with and transporting cells separately.

Special conditions are needed only when assembling the source of 40MCi activity to provide its safe operation.

Level of radiation purity of selected construction materials: Ti for T storage, electrolytic copper for cell body, satisfies the requirements of the physical experiment.

CDMS 'ZIP' Ionization & Phonon Detectors



- Tests conducted in CWRU detector test facility
- **S4 detector unsuited for CDMS**
 - Only one working phonon channel
 - **Poor energy resolution**
- Sufficient resolution to see features in ^{241}Am energy spectrum
 - Spot illumination on single working quadrant
- Measure collected phonon energy versus V_{qbias}
- Observe heat amplification

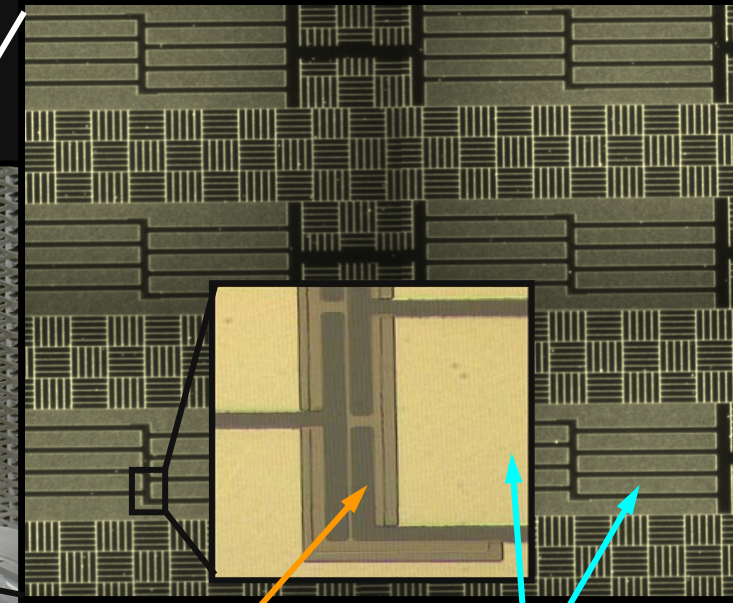
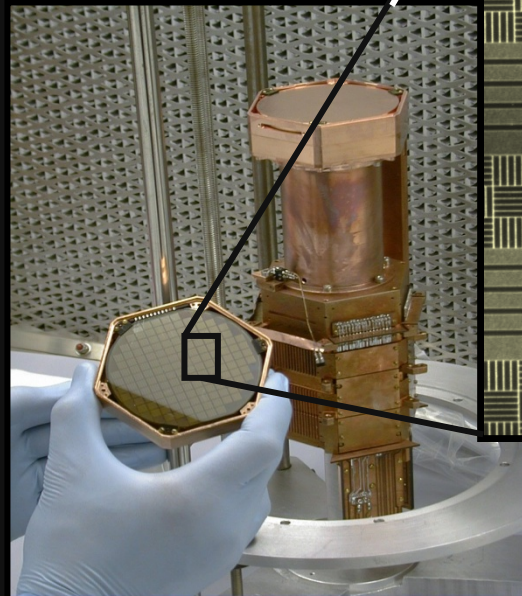
Fast athermal phonon technology

- Superconducting thin films of W, Al
- Stable Electrothermal Feedback
- Aluminum Quasiparticle Traps give area coverage

Low-Threshold Cryogenic Detectors

- First detector results are promising
 - ◆ Heat amplification observed to 200 Volts
 - ◆ Estimated intrinsic threshold is ~ 30 eV
- CDMS-style detectors already useful prototype as the basis for a neutrino magnetic moment search
 - ◆ Minimal manipulation of a maturing technology

100-g silicon
@T \sim 70 mK



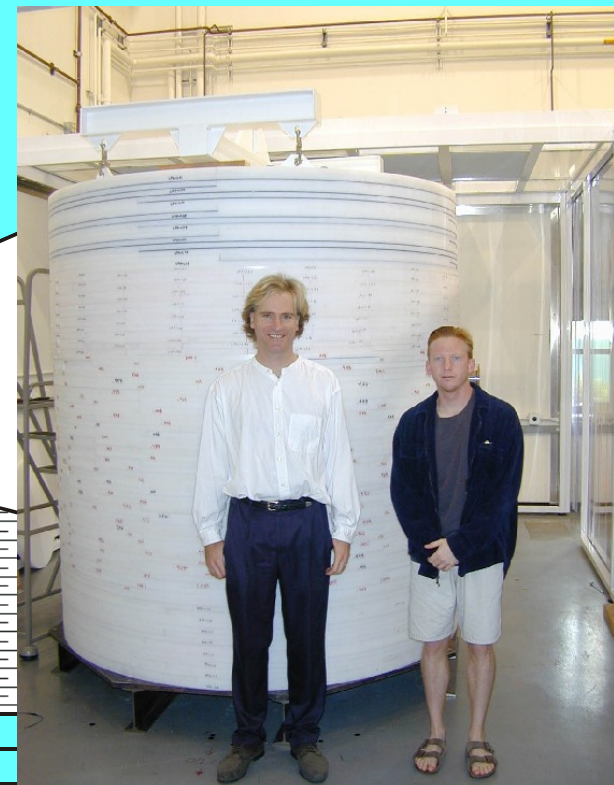
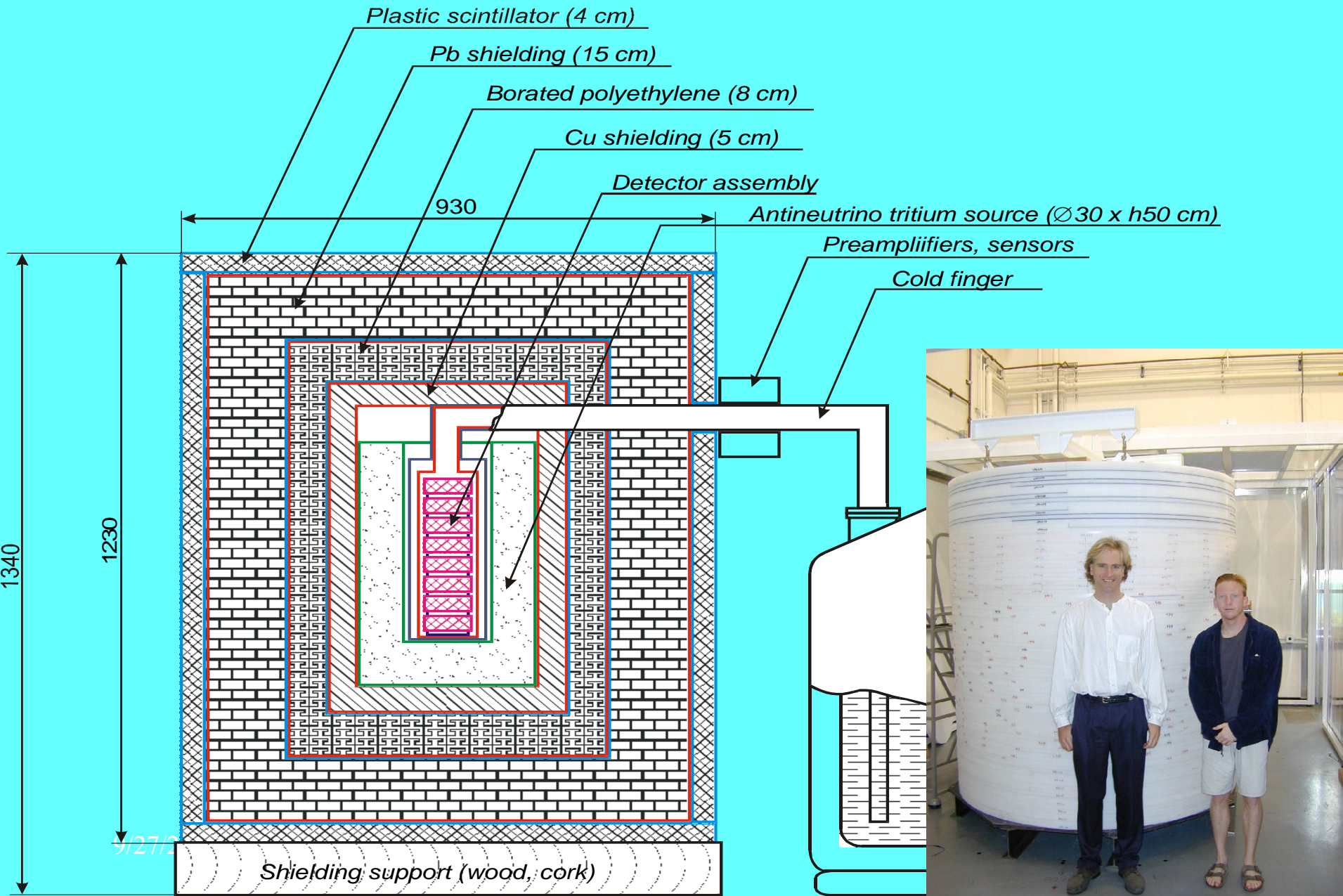
1 μ m tungsten

aluminum fins

Superconducting thin films of W, Al, Al

200-day tritium exposure x 3 kg at 30 eV threshold improve sensitivity to $3 \times 10^{-12} \mu_B$ or 10x current reactor-based limits

Principal Scheme of the Installation



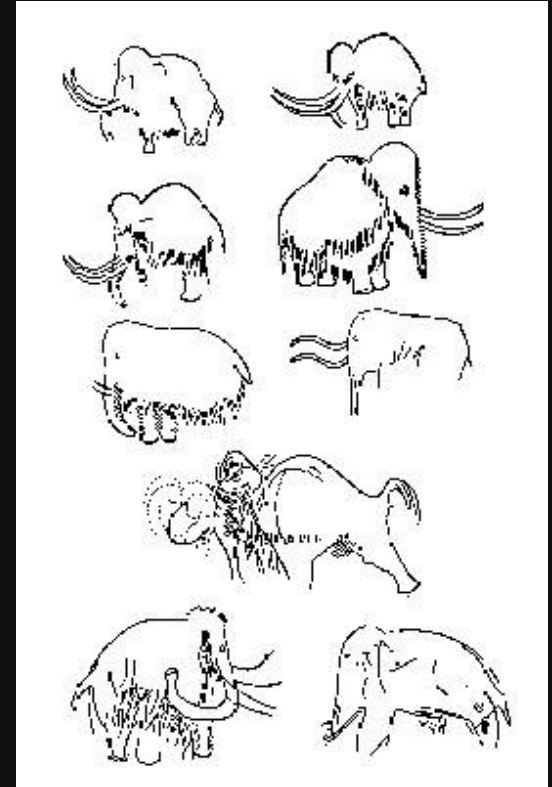
Sensitivity to the Neutrino Magnetic Moment

| Energy interval, eV | | 10 – 1260 | 10 – 200 |
|---|---------------------------------------|---|---|
| Magnetic scattering, N_M , events/day | $\mu_\nu \sim 1 \cdot 10^{-11} \mu_B$ | 2.4 | 1.4 |
| | $\mu_\nu \sim 3 \cdot 10^{-12} \mu_B$ | 0.22 | 0.13 |
| Weak scattering, N_W , events/day | | 0.15 | 0.04 |
| Background, N_B , (at the level 0.1 events/keV kg day) | | ~0.5 | ~0.1 |
| Sensitivity to the magnetic moment at 95% C.L. “ON”/“OFF”-200/200 days | | $\mu_\nu \leq 2.5 \cdot 10^{-12} \mu_B$ | $\mu_\nu \leq 2.2 \cdot 10^{-12} \mu_B$ |

MAMONT*) collaboration

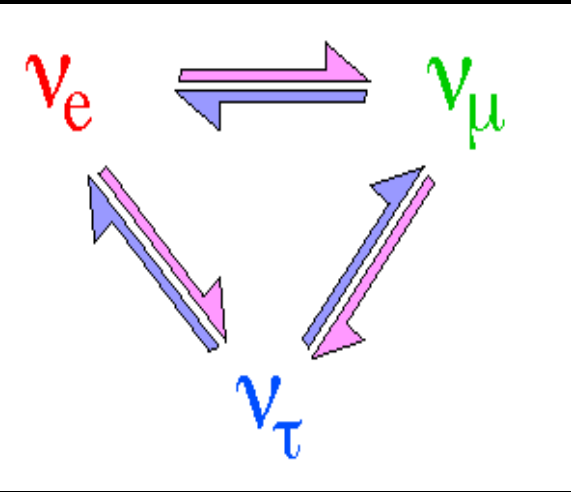
MAGnetic MOment Neutrino Tritium

- Brown University (Providence)
- Case Western Reserve University (Cleveland)
- Institute for Theoretical & Experimental Physics (Moscow)
- Institut für Kernphysik (Jülich)
- Joint Institute of Nuclear Research (Dubna)
- NIIMV (Zelenograd)
- Russian Federal Nuclear Center VNIIEF (Sarov)
- Stanford University (Stanford)



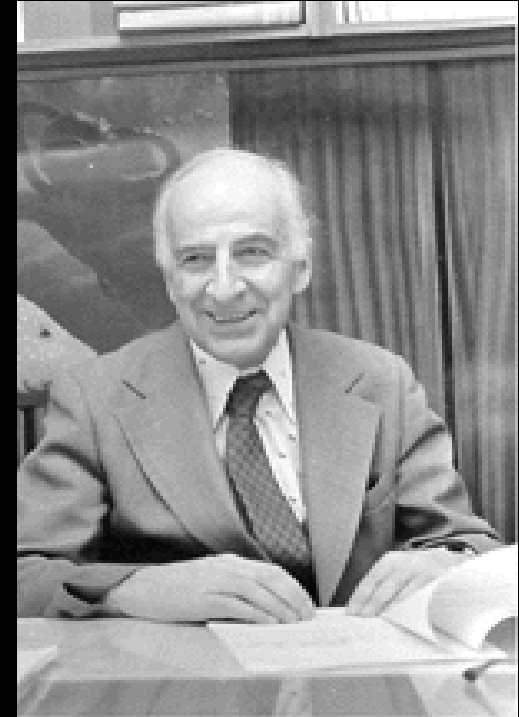
***)MAMONT (RUS) = MAMMOTH (ENG)**

Neutrino Oscillations



Neutrino oscillations were predicted by Bruno Pontecorvo in 1957.

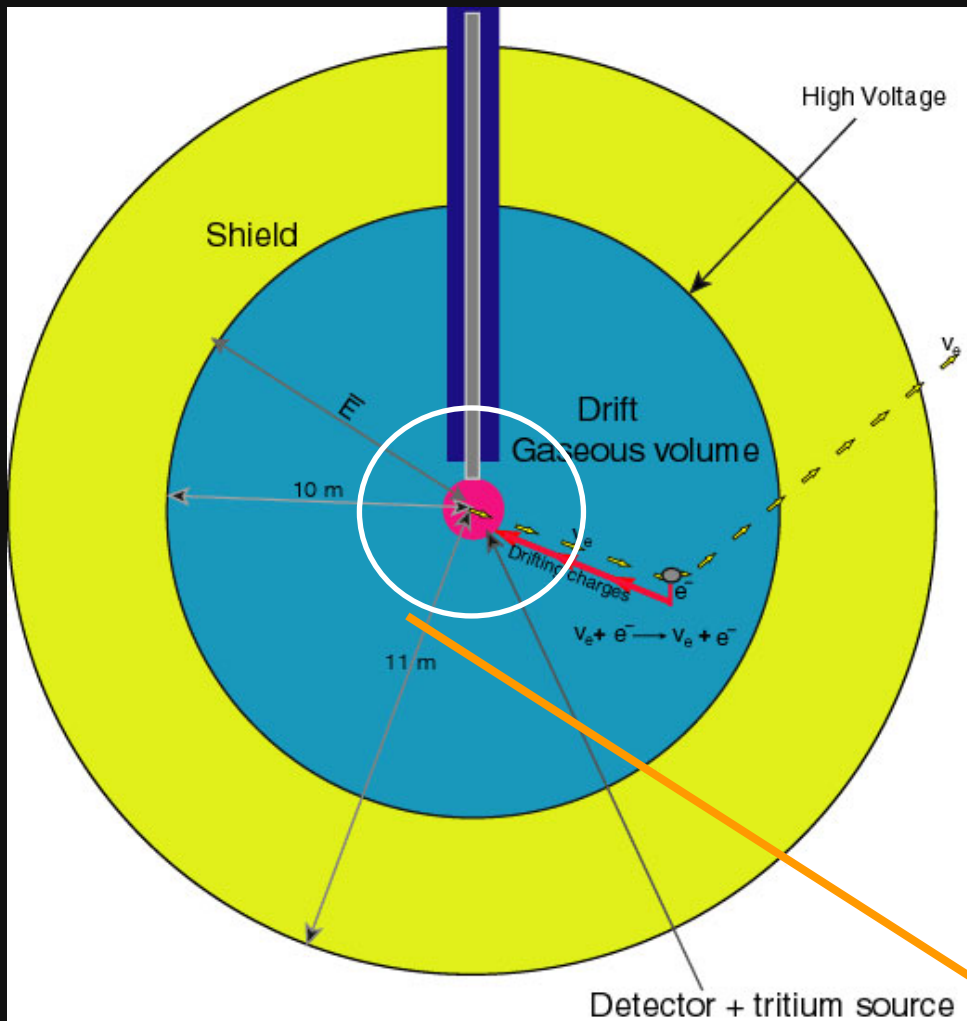
He discussed a possible solar neutrino deficit as a consequence of these oscillations.



Probability of oscillations between two neutrino flavors

$$P(\nu_1 \leftrightarrow \nu_2) = \sin^2 2\theta \cdot \sin^2 \left(\frac{\Delta m^2 \cdot L}{4E} \right)$$

Neutrino Properties Studied With a Triton Source Using Large TPC Detectors

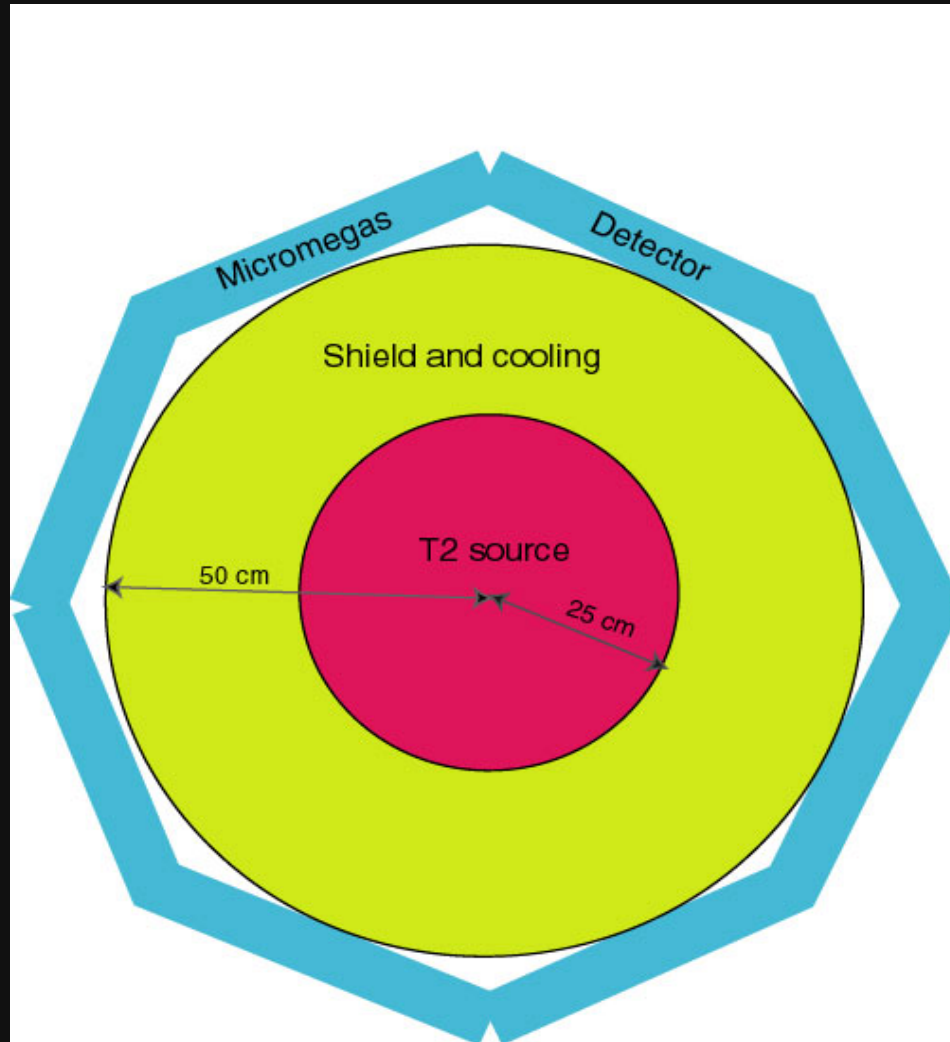


Proposal from CEA, Saclay, DAPNA, Gif-sur-Yvette, Cedex, France

Time Projection Chamber - large volume gaseous spherical vessel of 10-m radius that contains about 20 tons of gas (Xe).

Principal features of the proposed TPC - drift voltage of about ~ 100 kV
Next view (triton source)

A Schematic View of the Inner Part of the Vessel with the Detector and the Tritium Source



The 200 MCurie tritium source container is a sphere of ~25-cm radius.

The container of the source should fulfill the safety requirements and provide a flexible moving system for source on-off measurements.

“A unique sensitivity for the neutrino magnetic moment which is about two orders of magnitude beyond the current experimental limit”

Instead of Conclusion

1) 4 kg of tritium – not too big amount

How much tritium is needed for a thermonuclear reactor?



Constants: $\hbar = 1.054 \cdot 10^{-34} \text{ J}\cdot\text{s} = 6.582 \cdot 10^{-22} \text{ MeV}\cdot\text{s}$,
 $1 \text{ J} = 6.245 \cdot 10^{12} \text{ MeV}$, $1 \text{ year} = 3.15 \cdot 10^7 \text{ s}$

1 GW of thermal energy during a year requires

$$10^9 \text{ J/s} \cdot 3.15 \cdot 10^7 \text{ s} = 3.15 \cdot 10^{16} \text{ J} =$$

$$= 19.7 \cdot 10^{28} \text{ MeV} / (17.6 \text{ MeV/fusion}) = 1.12 \cdot 10^{28} \text{ fusions}$$

1 g of ${}^3\text{H}$ produces $\rightarrow 6.023 \cdot 10^{23} \text{ fusions} / 3 \sim 2 \cdot 10^{23} \text{ fusions}$

1 GW(thermal) requires $0.56 \cdot 10^5 \text{ g} = 56 \text{ kg of T per year}$

$\sim 1 \text{ kg/week!}$

2) Strong Physical Motivation

Boris Kayser
Neutrino 2008
May 28, 2008

New Physics can produce larger dipole moments than the $\sim 10^{-20} \mu_B$ SM ones.

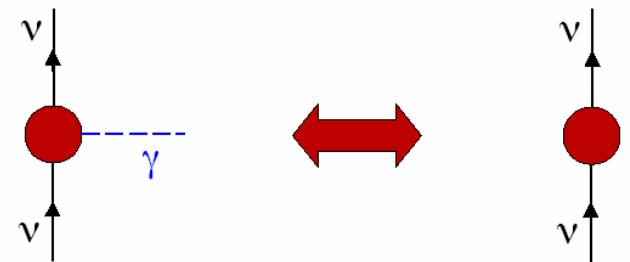
In the *Majorana* case, a *symmetry* suppresses the contribution of the dipole moment to the neutrino mass. So a bigger dipole moment is permissible. One finds —

For *Dirac* neutrinos, $\mu < 10^{-15} \mu_B$ for $\Lambda > 1$ TeV

For *Majorana* neutrinos, $\mu < \text{Present Bound}$

An observed μ below the present bound but well above $10^{-15} \mu_B$ would imply that neutrinos are *Majorana* particles.

A dipole moment that large requires L-violating new physics $\lesssim 1000$ TeV.



Dipole Moment

Mass Term

$$\mu_\nu \sim \frac{eX}{\Lambda} \left\{ \begin{array}{l} \text{Scale of} \\ \text{New Physics} \end{array} \right.$$

$$m_\nu \sim X\Lambda$$

R&D on 40 MCi Tritium Source

A.A. Yukhimchuk, Yu.I. Vinogradov,

A.N. Golubkov, S.K. Grishechkin,

R.I. Il'kaev, A.V. Kuryakin, B.L. Lebedev,

V.N. Lobanov, V.N. Mikhailov, D.P. Tumkin

Russian Federal Nuclear Center - VNIIEF,

Sarov

