

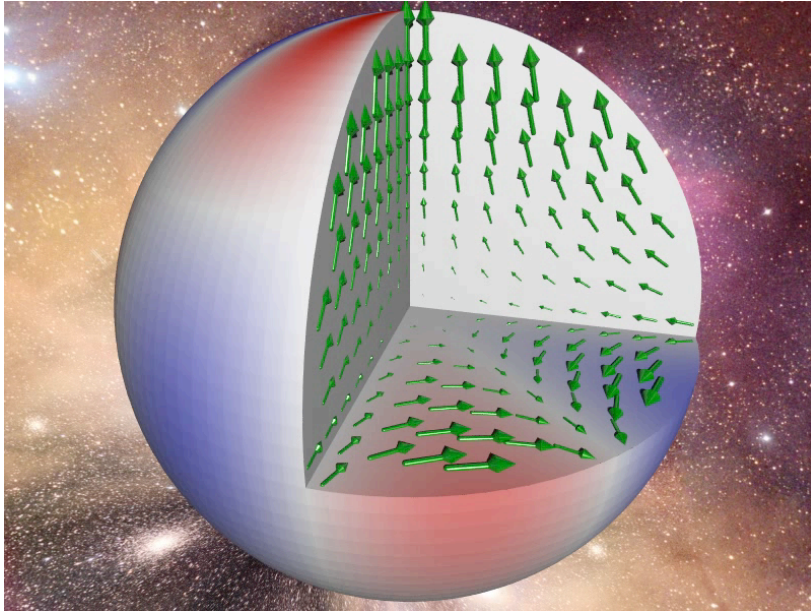


# NEUTRON STAR DYNAMICS

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**Theoretical Astrophysics, IAAT, Eberhard Karls University of Tübingen**

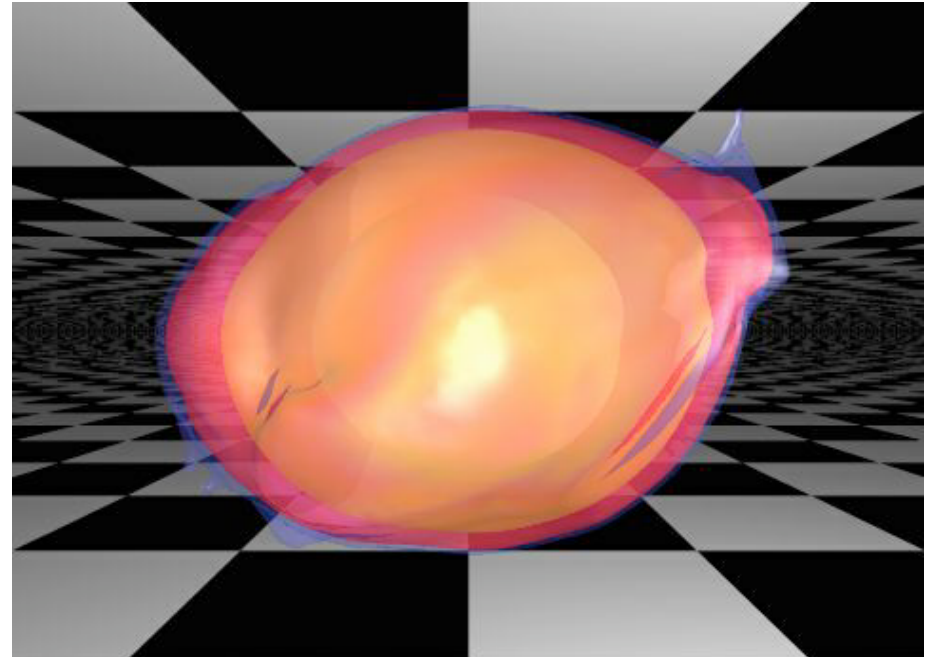
# Gravitational Wave Asteroseismology



We can estimate their **masses, radii, equations of state** by analysing the seismic data via the emitted gravitational waves

Neutron Stars oscillate wildly during the very first seconds of their life

Rotation is responsible for a number of instabilities which emit copious amounts of GWs



# A Laboratory for Theoretical Physics

**NS modelling involves the very extremes of physics:**

General Relativity

Alternative theories, Boson, Q-stars,...

Rotation

Slow, Fast, Differential  
**Instabilities**

Equation of State

Exotic nuclear physics, strange quarks,  
hyperons,... **Superfluidity**, Cold vs Warm

Magnetic Fields

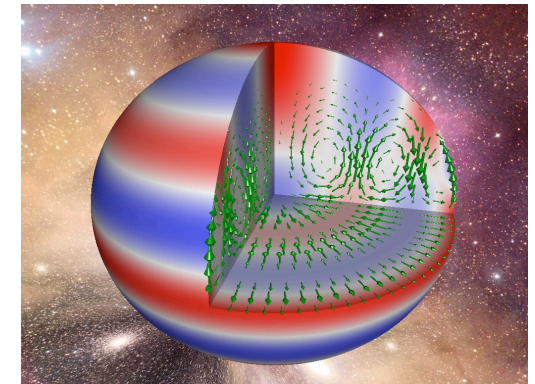
Magnetars, Slowdown, Suppression  
of Instabilities, **Uniform Rotation**

Crust – Core interface

Suppression of Instabilities,

**Can GW, x-ray,  $\gamma$ -ray observations constrain the theoretical models?**

# Neutron Star “ringing”



**p-modes:** main restoring force is the pressure (**f-mode**) ( $>1.5 \text{ kHz}$ )

**Inertial modes:** (**r-modes**) main restoring force is the **Coriolis force**

**w-modes:** pure **space-time modes** (only in GR) ( $>5 \text{ kHz}$ )

**Torsional modes** (t-modes) ( $>20 \text{ Hz}$ ) shear deformations. Restoring force, the weak **Coulomb force** of the crystal ions.

**... and many more**

$$\sigma \approx \sqrt{\frac{M}{R^3}}$$

$$\sigma \approx \Omega$$

$$\sigma \approx \frac{1}{R}$$

$$\sigma \approx \sqrt{\frac{v_s}{R}}$$

# GW Asteroseismology

Oscillation patterns can reveal the internal structure of neutron stars : **mass, radius, EoS, rotation, B-field, crust,...**

$$\sigma_f (kHz) \approx 0.78 + 1.64 \left( \frac{\bar{M}}{\bar{R}^3} \right)^{1/2}$$

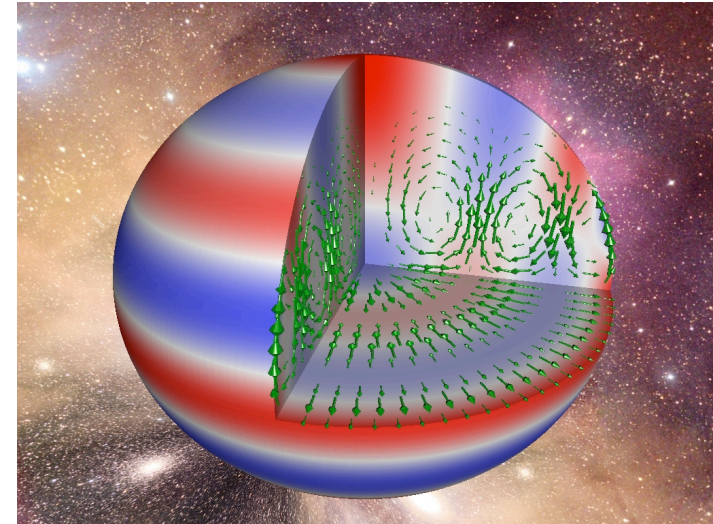
$$\frac{1}{\tau_f (s)} \approx \frac{\bar{M}^3}{\bar{R}^4} \left[ 22.85 - 14.65 \frac{\bar{M}}{\bar{R}} \right]$$

$$\sigma_w (kHz) \approx \frac{1}{\bar{R}} \left( 20.9 - 9.1 \frac{\bar{M}}{\bar{R}} \right)$$

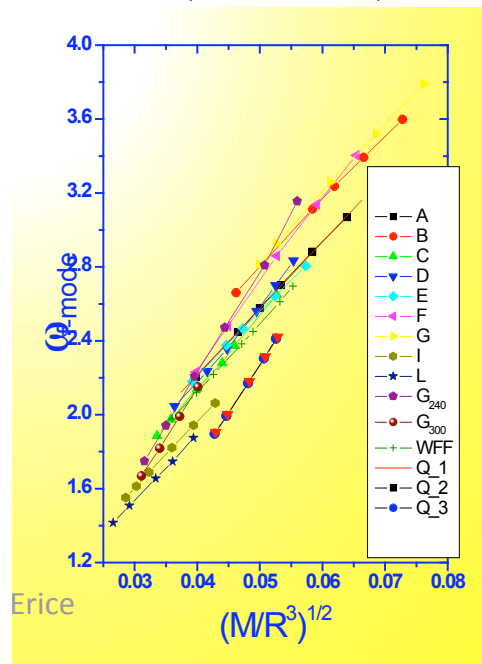
f-mode frequency

f-mode frequency

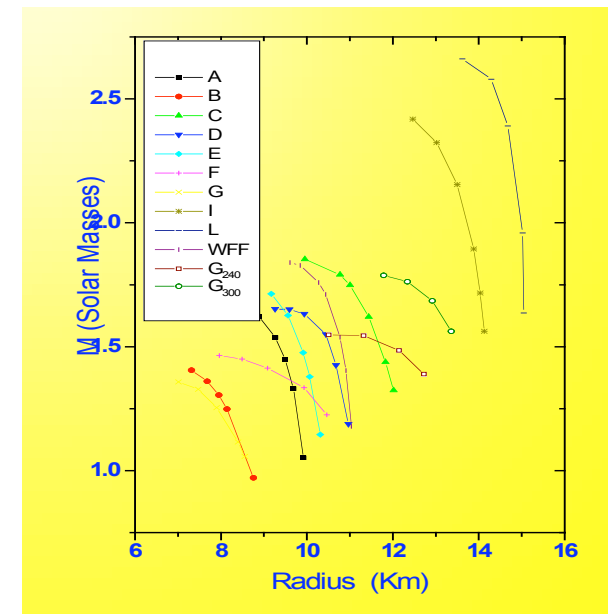
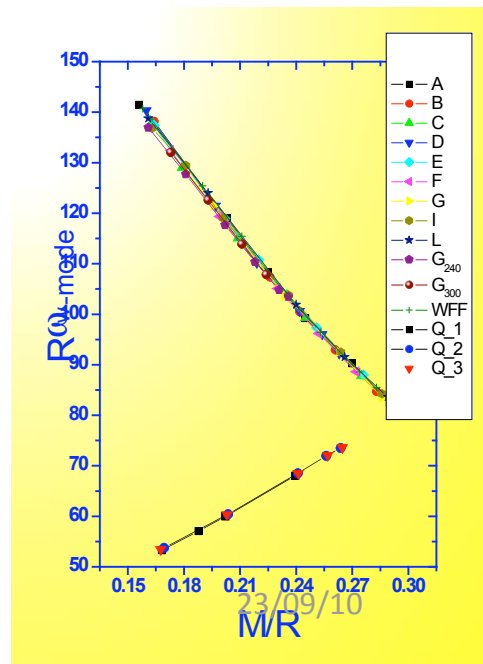
w-mode frequency



Andersson, KK 1996, 1998, 2001



+



# Effect of Rotation & Magnetic Fields

## ROTATION

- **Frame dragging**
- **Quadrupole deformation**
- **Rotational instabilities**
- The degeneracy in **m** is removed and the nonrotating mode of index **ℓ** is split into **2ℓ+1** different **(ℓ,m)** modes
- Shifting of the frequencies and damping times
- Coupling of polar **ℓ**-term to an axial **ℓ±1** term and v-v

## MAGNETIC FIELD

- **No significant effect** in the fluid frequencies and damping/growth times

$$\frac{\text{magnetic energy}}{\text{gravitational energy}} \sim \frac{B^2 R^3}{GM^2 / R} \sim 10^{-4} \left( \frac{B}{10^{16} G} \right)^2$$

- For **magnetars** we may observe **Alfvén oscillations**

# Stability of Rotating Stars

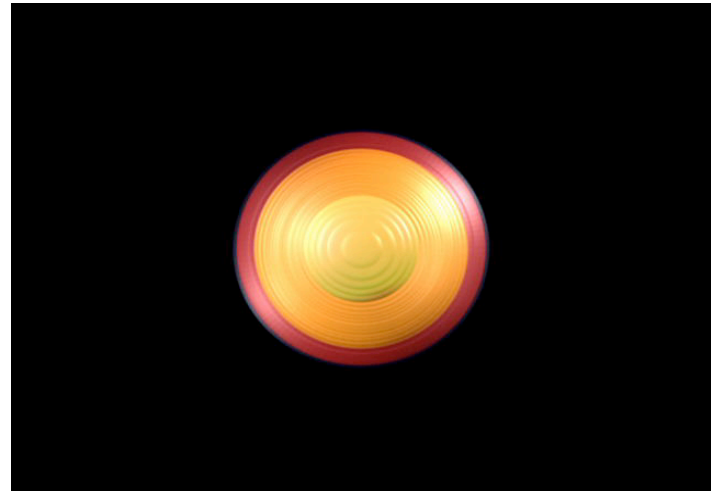
## Non-Axisymmetric Perturbations

A general criterion is:

**T** : rotational kinetic energy

**W** : gravitational binding energy

$$\beta = \frac{T}{|W|} \approx \frac{2}{15} e^2 + \dots$$



### Dynamical Instabilities

- Driven by **hydrodynamical forces** (bar-mode instability)
- Develop at a time scale of about **one rotation period**

$$e > 0.953 \quad \text{or} \quad \beta > 0.274$$

### Secular Instabilities

- Driven by **dissipative forces** (*viscosity, gravitational radiation*)
- Develop at a time scale of **several rotation periods**.  
Chandrasekhar-Friedman-Schutz (CFS)

$$e > 0.813 \quad \text{or} \quad \beta > 0.138$$

**GR predicts considerably lower  $\beta$**

$\beta \sim 0.24$  for the onset of the **dynamical instabilities**

$\beta \sim 0.07$  for the onset of the **secular instabilities**

# Bar-mode dynamical instability

- ✓ For rapidly (differentially!) rotating stars with:

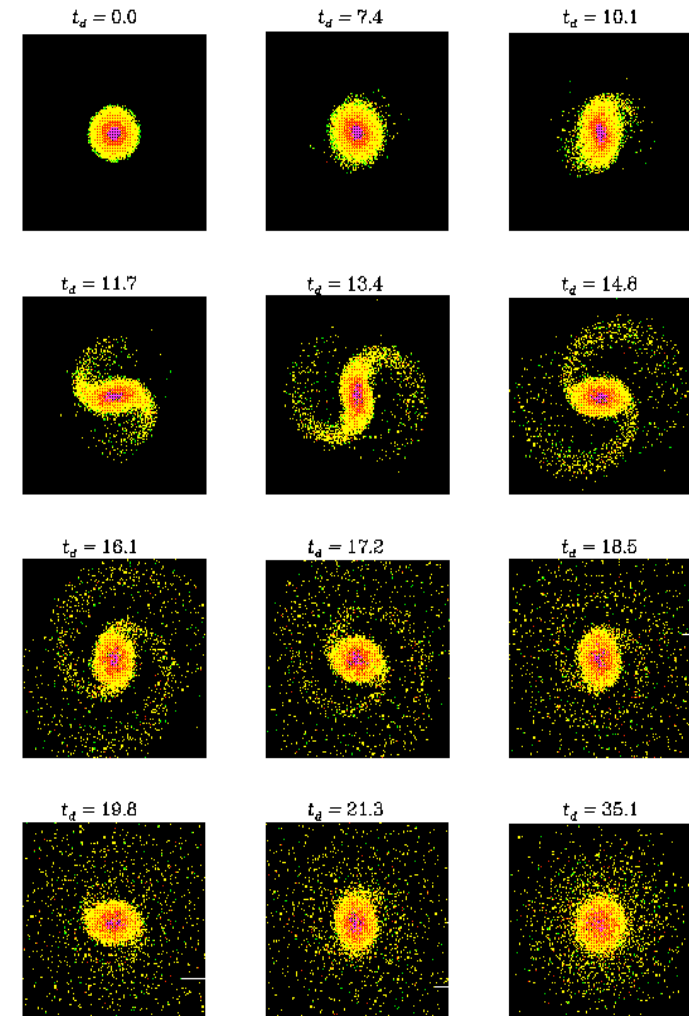
$$\beta = \frac{T}{|W|} \sim \frac{1}{R} > \beta_{\text{dyn}} \approx 0.27$$

- ✓ GR enhances the onset of the instability ( $\beta_{\text{dyn}} \gtrsim 0.24$ ) and  $\beta$  decreases with increasing  $M/R$ .

- ✓ The “bar-mode” grows on a dynamical timescale.

$$h \approx 10^{-22} \left( \frac{\varepsilon}{0.2} \right) \left( \frac{f}{3 \text{ kHz}} \right)^2 \left( \frac{15 \text{ Mpc}}{\text{d}} \right) M_{1.4} R_{10}^2$$

- ✓ If the bar persists for many ( $\sim 10-100$ ) rotation periods, the signal will be easily detectable from at least Virgo cluster.
- ✓ Typical Frequencies  $\sim 1.5-3.5 \text{ kHz}$





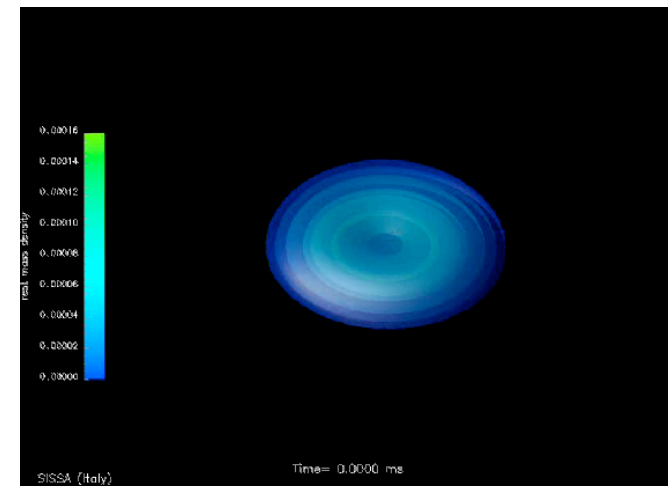
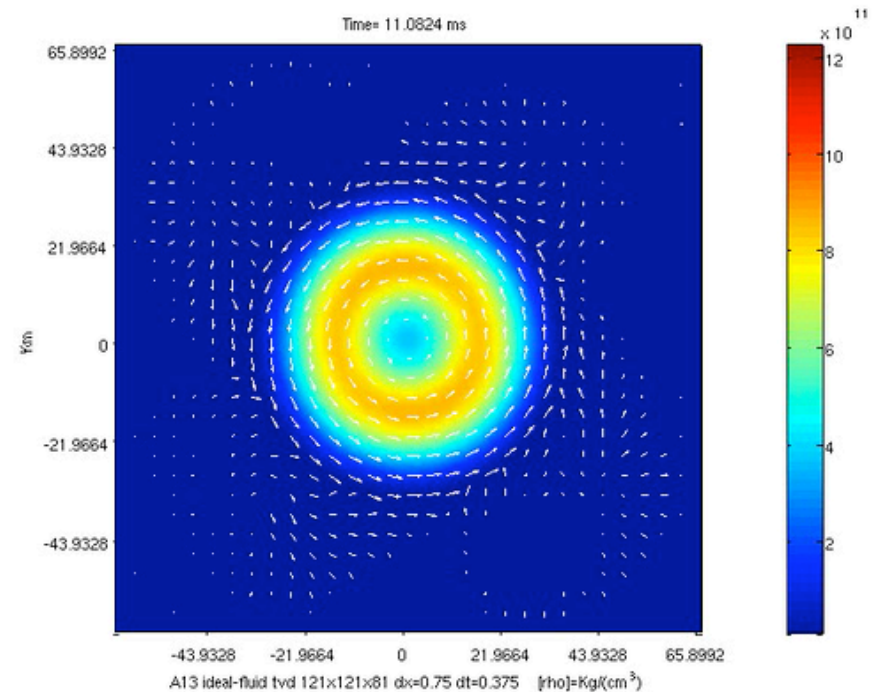
# Bar Mode Dynamical Instability

- Bars can be also created during the merging of **NS-NS**, **BH-NS**, **BH-WD** and **Collapsars (type II)**.
- Bar-mode **instability might happen for much smaller  $\beta$**  if centrifugal forces produce a peak in the density off the source's rotational center.

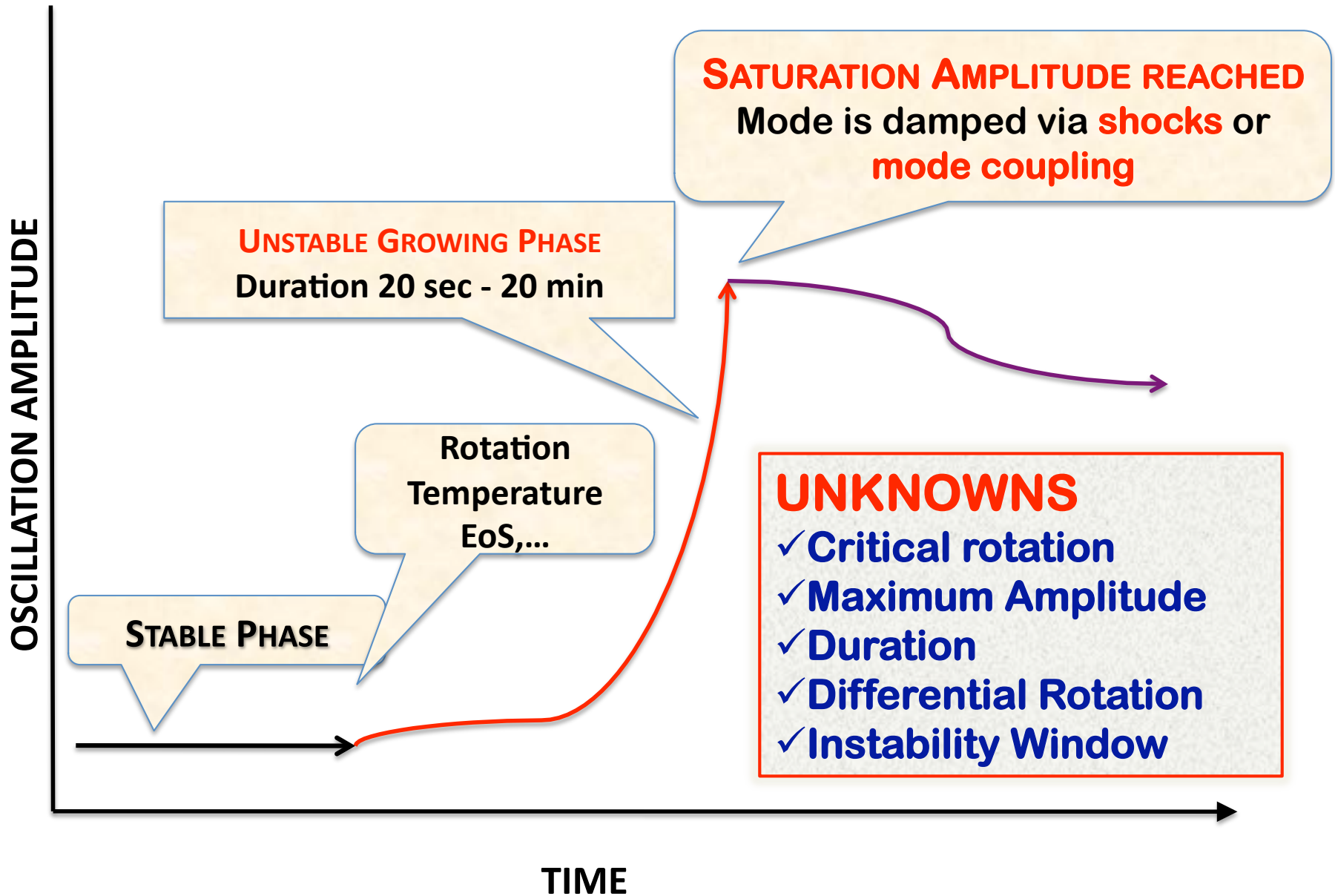
$$h_{\text{eff}} \approx 3 \times 10^{-22} \left( \frac{f}{800 \text{ Hz}} \right)^{1/2} \left( \frac{R_{\text{eq}}}{30 \text{ km}} \right) \left( \frac{M}{1.4 M_{\odot}} \right)^{1/2} \left( \frac{100 \text{ Mpc}}{d} \right)$$

## ➤ LOW $T/|W|$ Instability

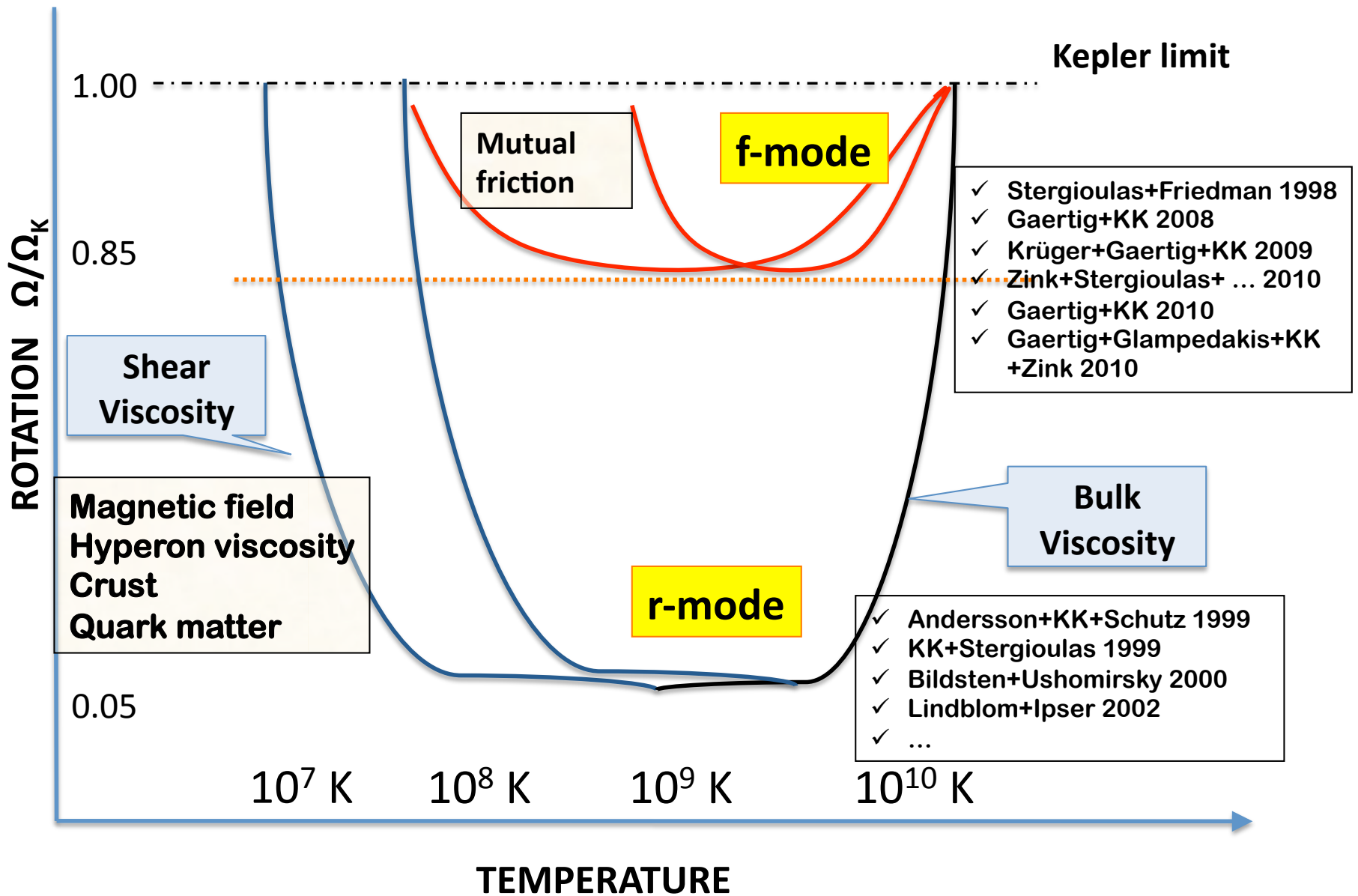
- **Highly differentially rotating stars** are shown to be dynamically unstable for significantly lower  $\beta$  (even when  $\beta \gtrsim 0.01$ ).
- Bars can be also create during the collapse of a SMS before the creation of a SMBH. **Ideal sources for LISA**.



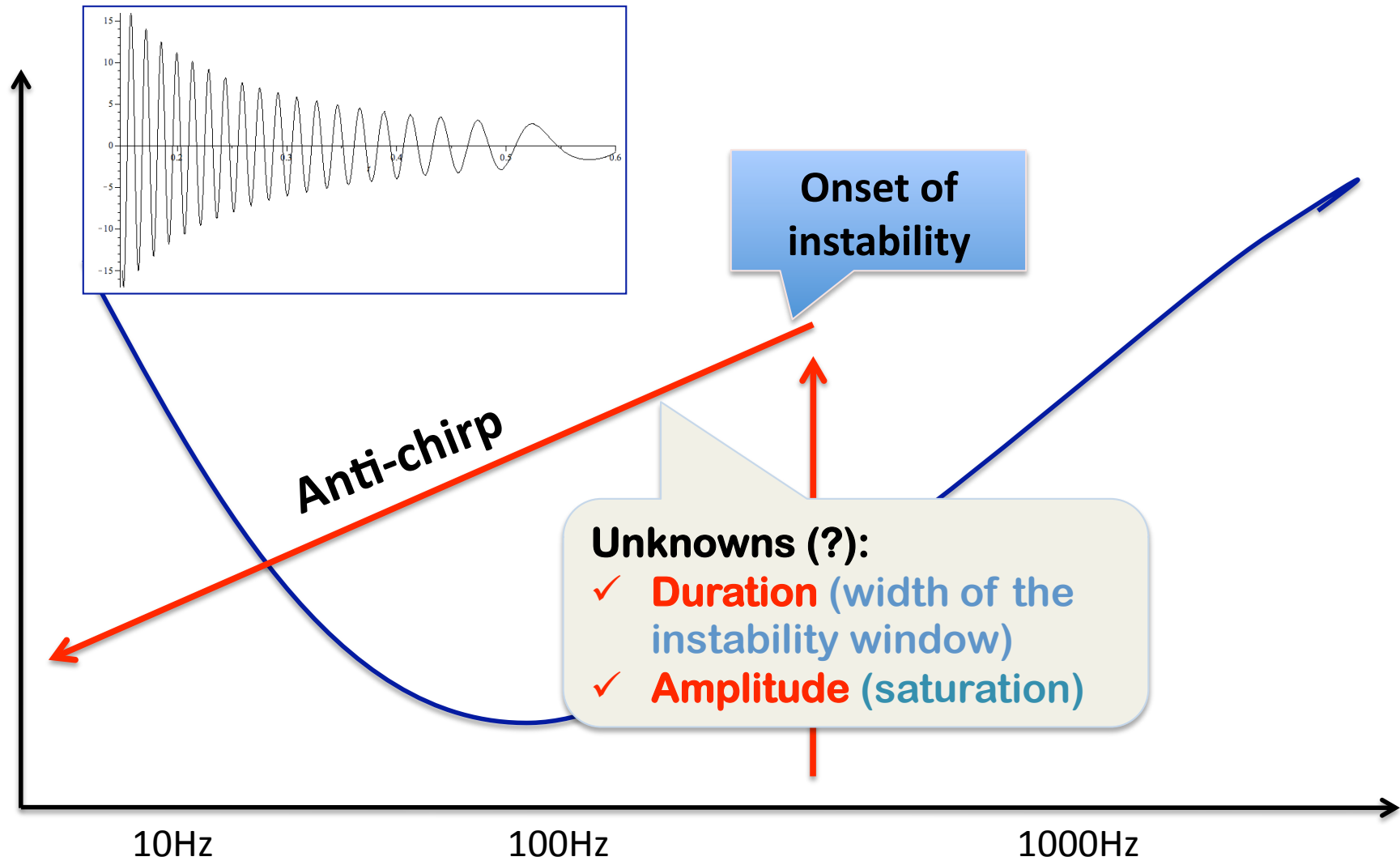
# The Excitation of Secular Instabilities



# INSTABILITY WINDOW

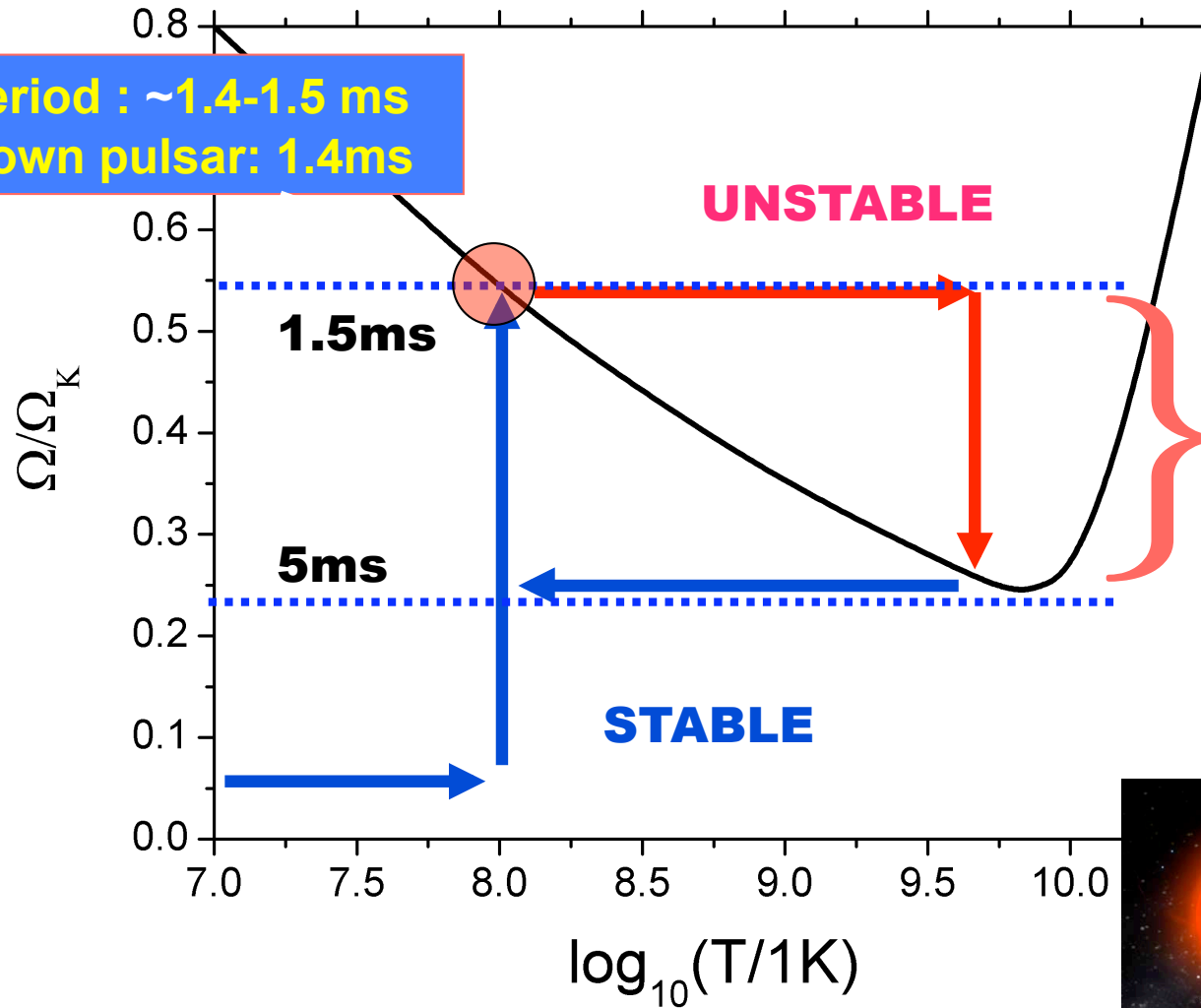


# f-mode Instability

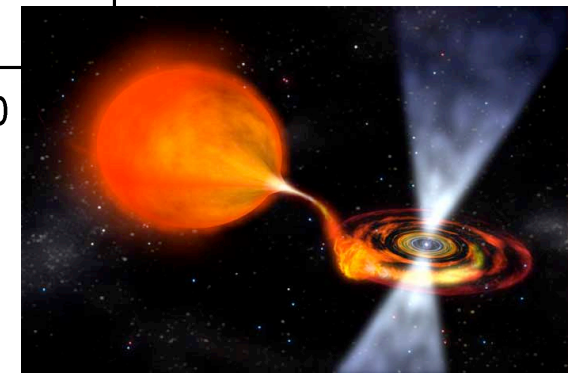


# LMXBs & r-modes

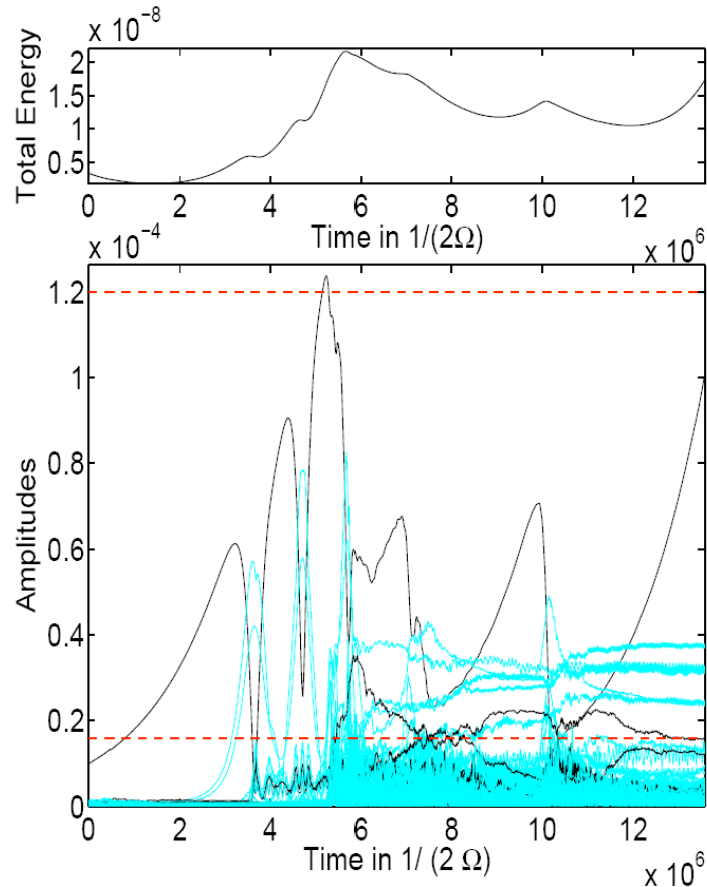
Limiting Period : ~1.4-1.5 ms  
Fastest known pulsar: 1.4ms



Andersson, KK, Stergioulas 1999, Levin 2000  
Andersson, Jones, KK, Stergioulas 2000, Heyl 2002  
Andersson, Jones, KK 2002



# R-modes



$$h(t) \approx 10^{-20} \alpha \left( \frac{\Omega}{1 \text{ kHz}} \right) \left( \frac{10 \text{ Kpc}}{d} \right)$$

$$\alpha \approx 10^{-3} - 10^{-4}$$

**GW amplitude depends on the saturation amplitude**

- ✓ **Mode coupling** might not allow the growth of instability to high amplitudes (Cornell group '04-'07)
- ✓ The existence of **crust**, **hyperons** in the core, **magnetic fields**, affects the efficiency of the instability.
- ✓ For newly born neutron stars might be **quite weak** ; unless we have the creation of a strange star
- ✓ **Old accreting neutron** (or strange) stars, probably the best source!

# Fast Rotating NS in GR: f-mode

## Frequency

$$\omega^2 = \frac{2l(l-1)}{2l+1} \left( \frac{GM}{R^3} \right)$$

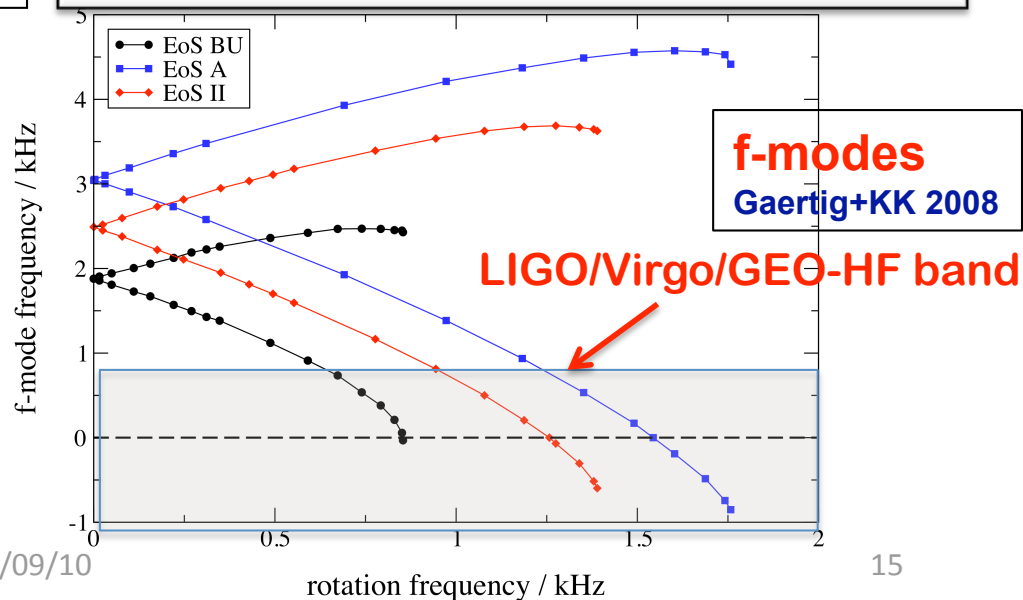
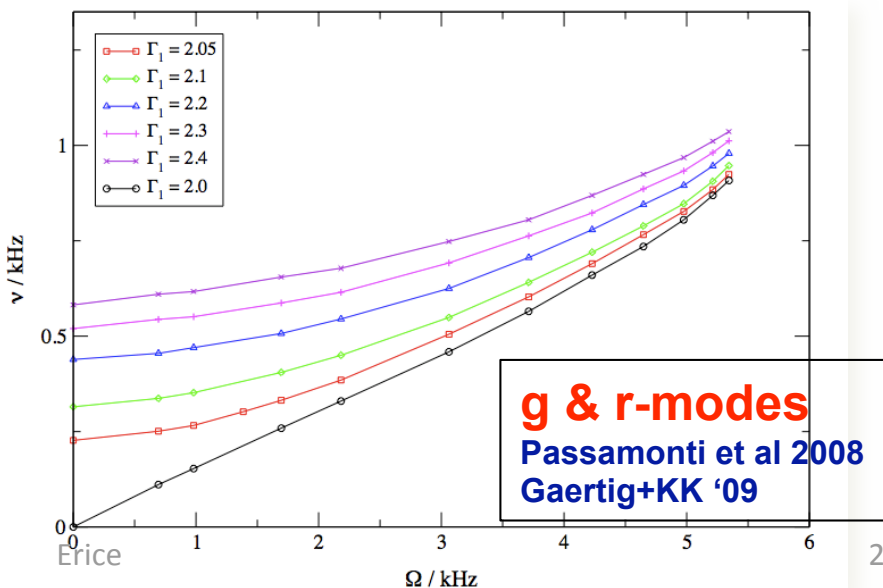
Gaertig+KK 2008,09,10  
 Krueger,,Gaertig, KK 2009  
 Zink etal 2010

## Damping/Growth time

$$t_{GW} \approx f(l)R \left( \frac{R}{M} \right)^{l+1} \sim 0.07 \left( \frac{1.4 M_{\odot}}{M} \right)^3 \left( \frac{R}{10km} \right)^4 \text{ sec}$$

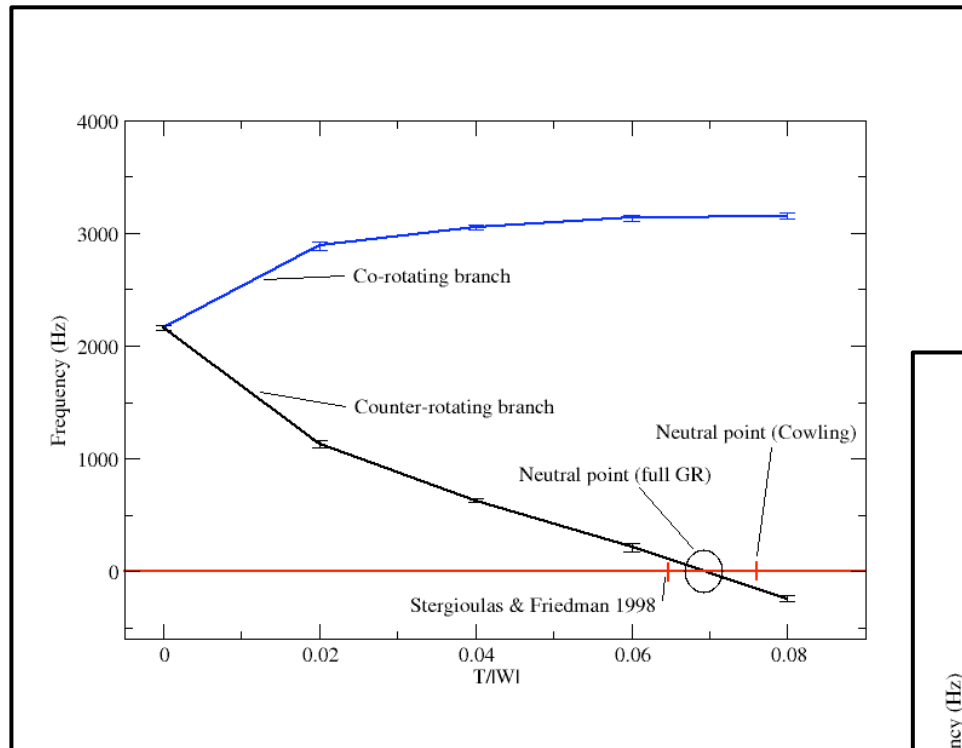
- ✓ In GR the  $m=2$  mode becomes unstable for  $\Omega > 0.85 \Omega_{Kepler}$
- ✓ Differential rotation affects the onset of the instability
- ✓ Up to 10% of energy and angular momentum will be dissipated by GWs.

- ### Major uncertainties:
1. Relativistic growth times ✓
  2. Nonlinear saturation ✓
  3. Initial rotation rates of proto-neutron stars ✗



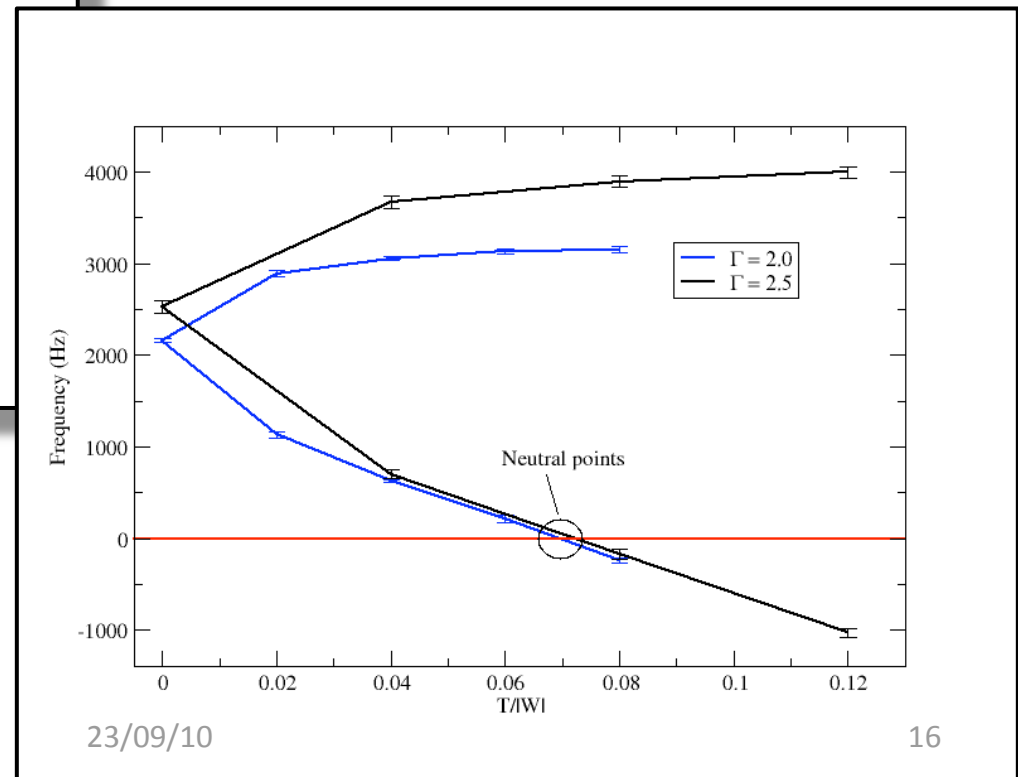
# f-mode : non-linear results

B. Zink, N. Stergioulas, O. Korobkin, P. Diener, E. Schnetter (2010)



**Top:** Frequencies of  $l = |m| = 2$  f-modes in rapidly rotating  $\Gamma = 2$  polytropes

**Bottom:** Comparison between  $\Gamma = 2$  and  $\Gamma = 2.5$





# f-modes: Asteroseismology

- ✓ We can trace the effect of rotation of the f-mode (and any p, i or g-mode) frequency and the onset of the CFS instability
- ✓ We can produce empirical relation relating the parameters of the neutron stars to the observed frequencies.

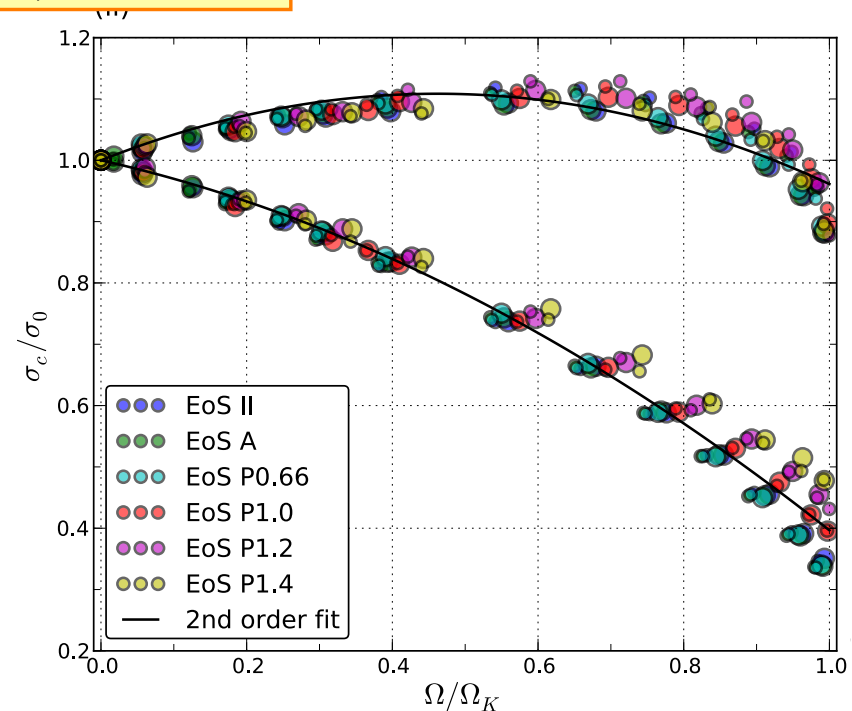
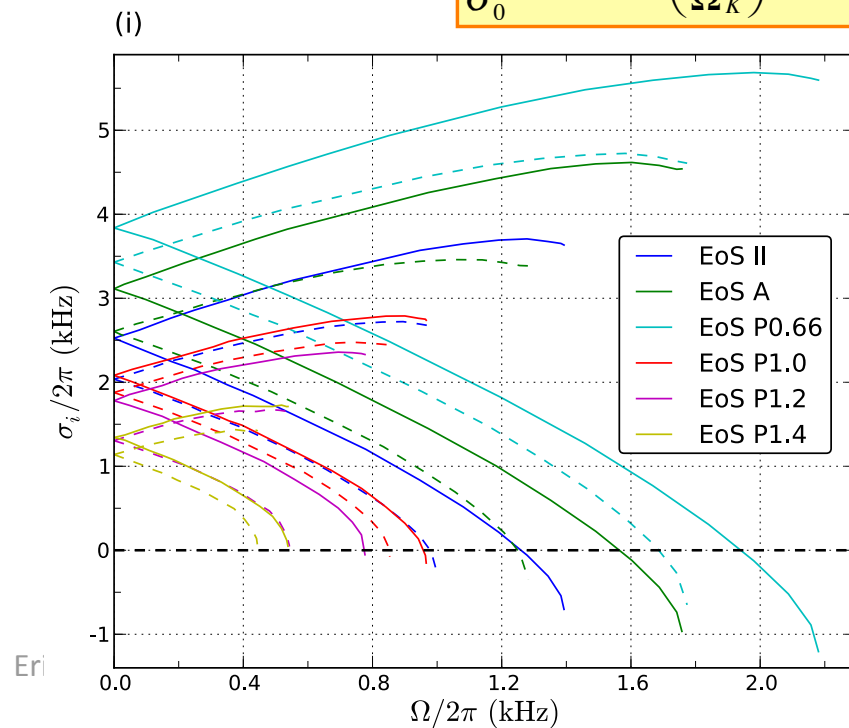
$$\frac{\sigma}{\sigma_0} \approx 1 + 0.63 \left( \frac{\Omega}{\Omega_K} \right) - 0.32 \left( \frac{\Omega}{\Omega_K} \right)^2 + \dots \quad (m=2)$$

$$\frac{\sigma}{\sigma_0} \approx 1 - 0.41 \left( \frac{\Omega}{\Omega_K} \right) - 0.53 \left( \frac{\Omega}{\Omega_K} \right)^2 + \dots \quad (m=-2)$$

$$\Omega_K \approx 0.67 \left( \frac{M}{R^3} \right)^{1/2} \quad \text{or} \quad \Omega_K \approx C(\chi) \left( \frac{M}{R^3} \right)^{1/2}$$

$$C(\chi) = 0.47 + 0.76 \frac{M}{R}$$

Gaertig-KK 2008, 2010



Eri

# f-mode: Damping/Growth time

$$E = \frac{1}{2} \int \left[ \rho \delta u^a \delta u_a^* + \frac{\delta p}{\rho} \delta \rho^* \right] d^3x \Rightarrow E \approx \sigma^2$$

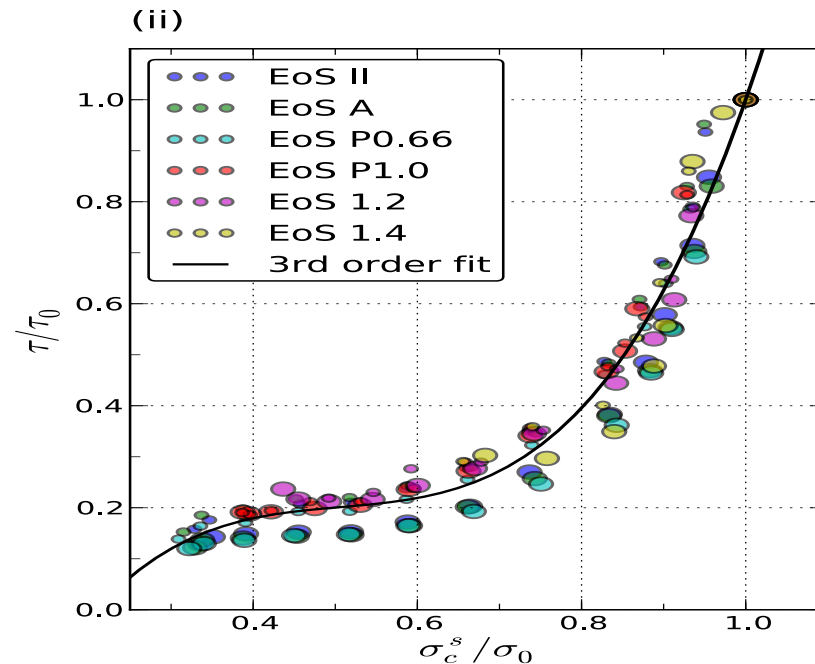
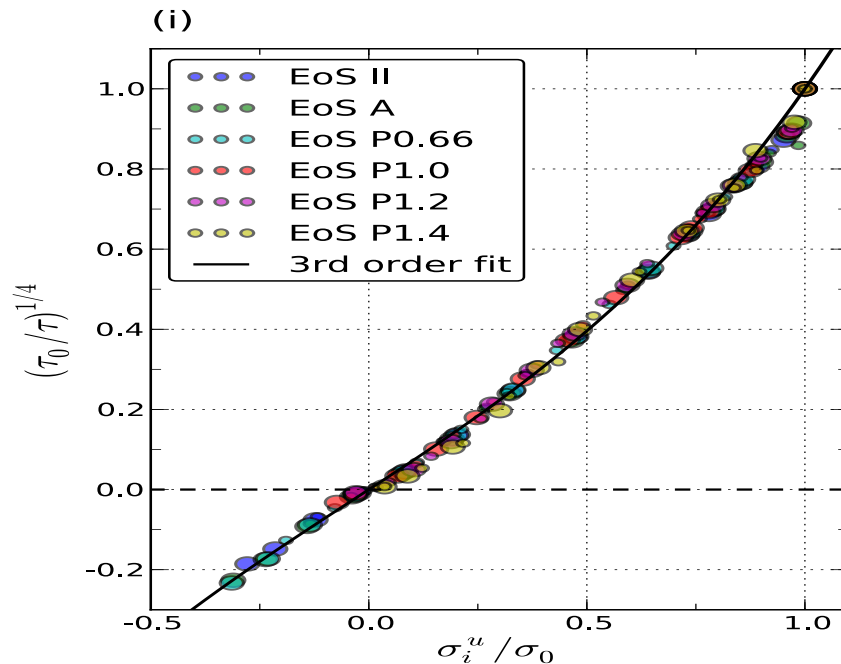
Gaertig-KK 2010

$$\frac{1}{\tau_{GR}} = -\frac{1}{2E} \left( \frac{dE}{dt} \right) \approx \sigma^4$$

$$\frac{dE}{dt} = -\sigma_i (\sigma_i + m\Omega) N_\ell |\delta D_{\ell m}| \sigma_i^4 \Rightarrow \frac{dE}{dt} \approx \sigma_i^6$$

$$\left[ \frac{\tau_0}{\tau} \right]^{1/4} \approx \text{sgn}(\sigma_i) 0.71 \left( \frac{\sigma_i}{\sigma_0} \right) \left[ 1 + 0.048 \left( \frac{\sigma_i}{\sigma_0} \right) + 0.35 \left( \frac{\sigma_i}{\sigma_0} \right)^2 \right]$$

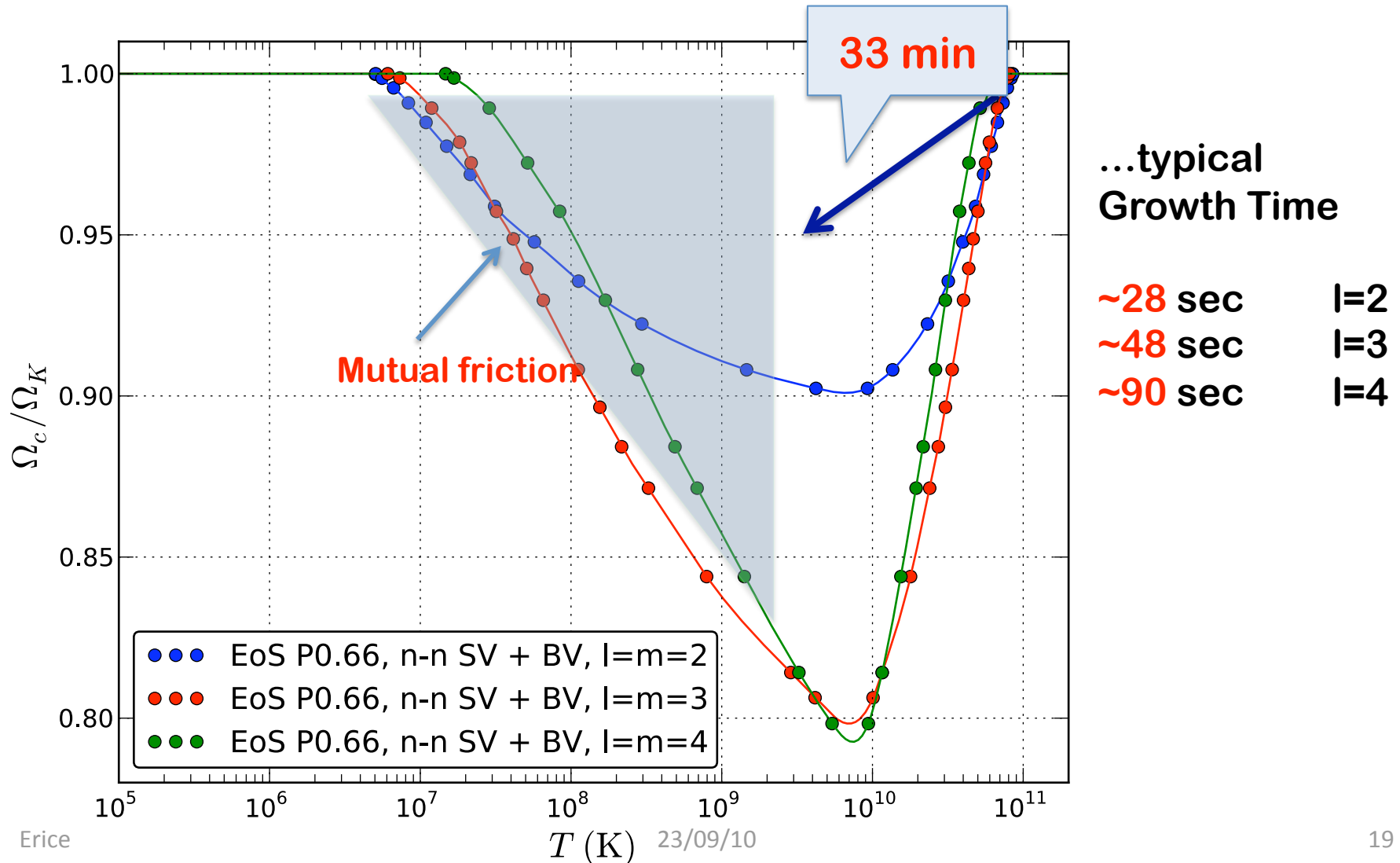
$$\left[ \frac{\tau_0}{\tau} \right] \approx -0.66 \left[ 1 - 7.33 \left( \frac{\sigma_c}{\sigma_0} \right) + 15.06 \left( \frac{\sigma_c}{\sigma_0} \right)^2 - 9.26 \left( \frac{\sigma_c}{\sigma_0} \right)^3 \right]$$



# Instability Window

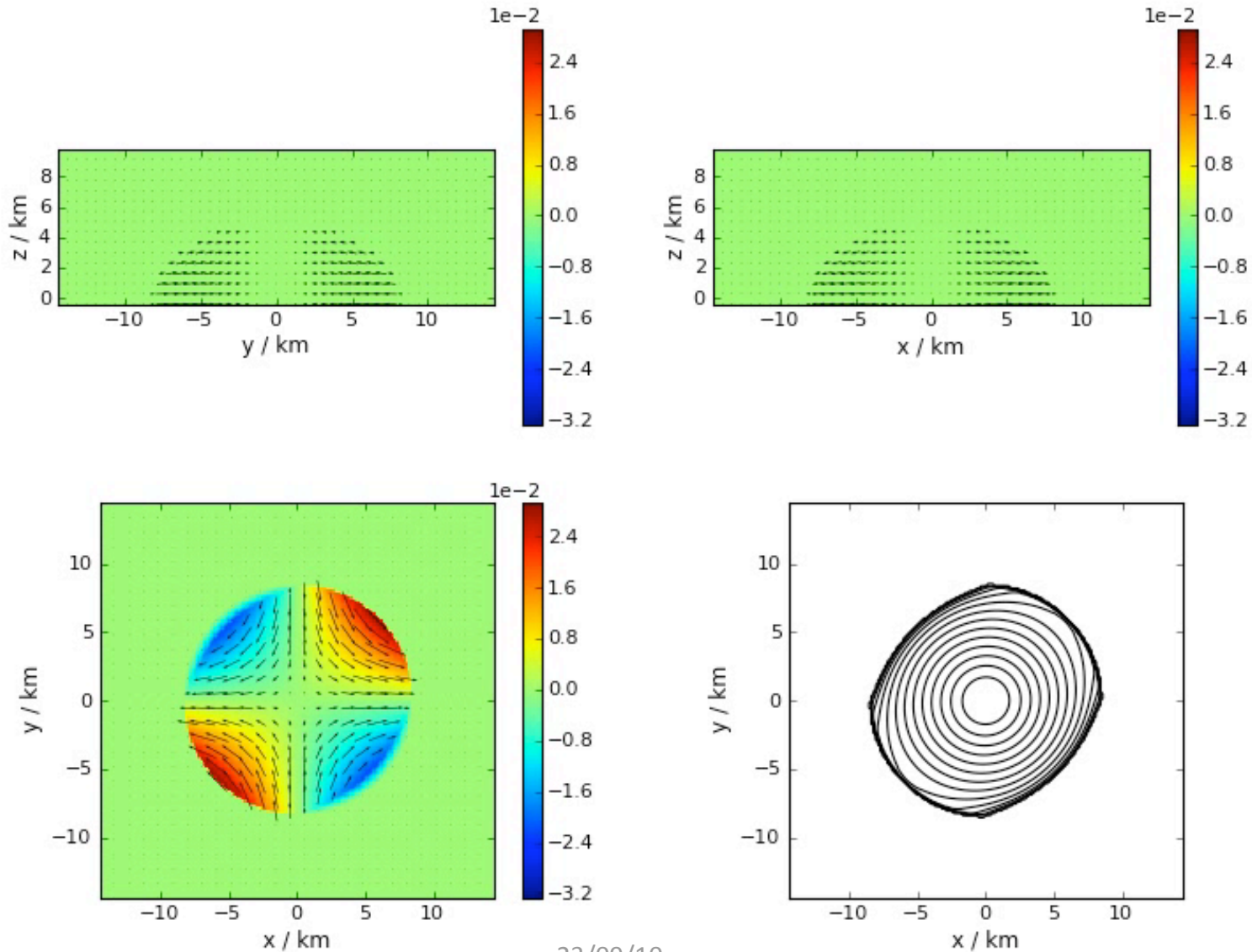
For the first time we have the window of f-mode instability in GR

Gaertig, Glampedakis, KK, Zink (2010)



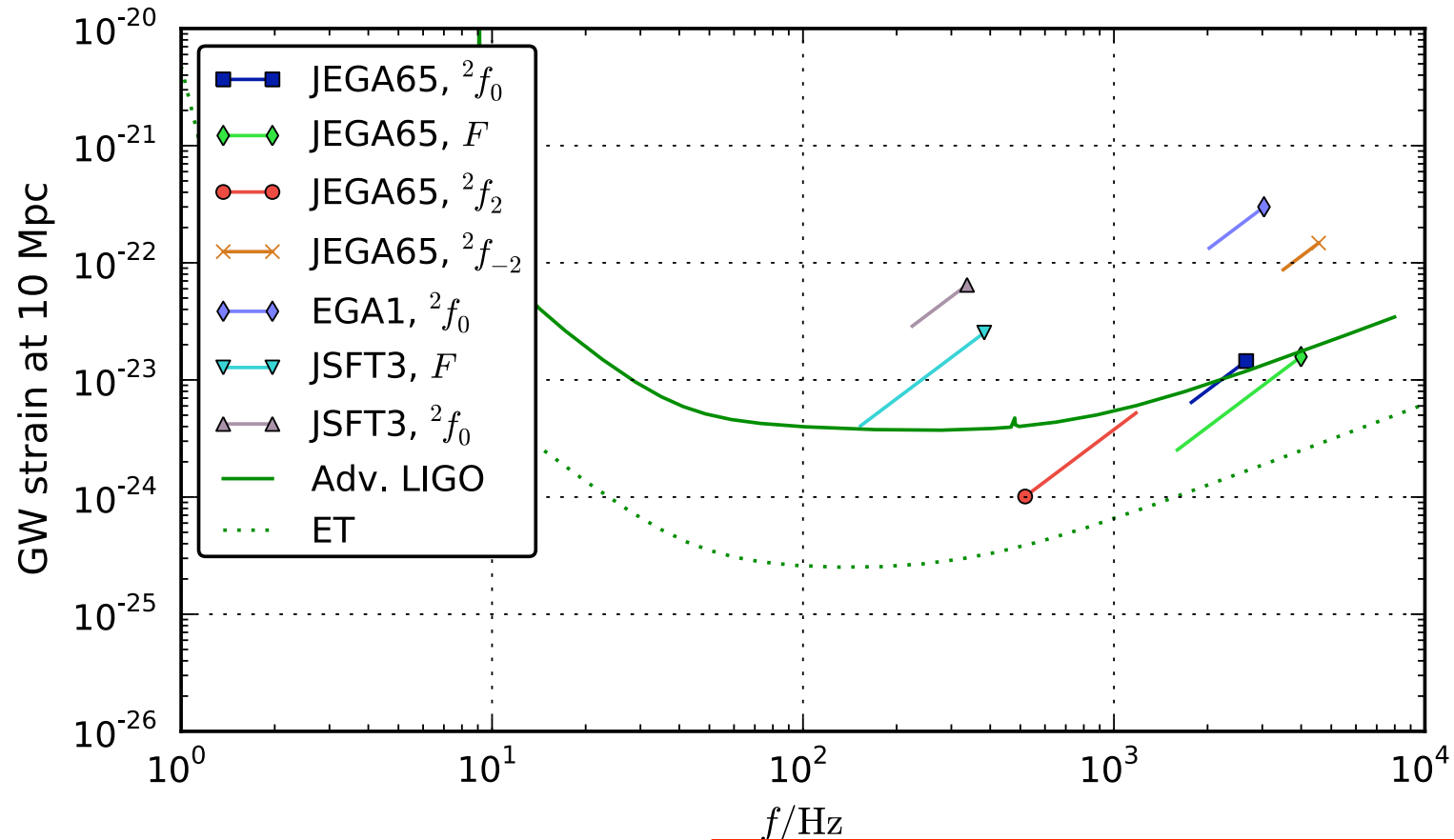
# Animation of the $l=m=2$ f-mode

Kastaun, Willburger, KK (2010)



# Detectability (10Mpc)

Kastaun, Willburger, KK (2010)



$$h \approx 10^{-23} - 10^{-24} \left( \frac{10\text{Mpc}}{r} \right)$$

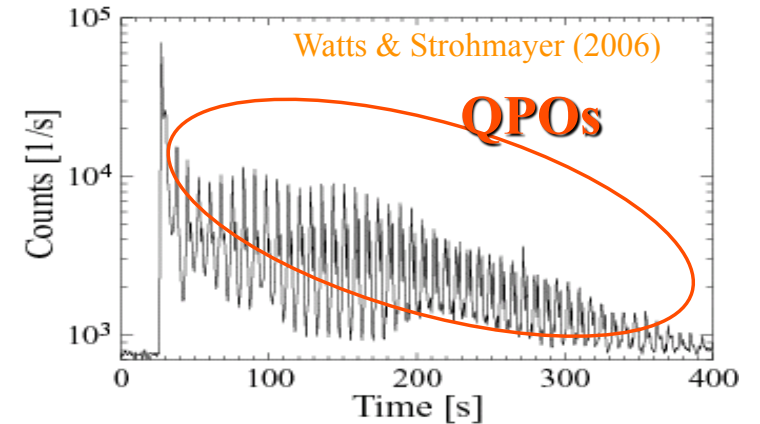


$$h_{\text{eff}} = h\sqrt{N} \approx 10^{-23} - 10^{-24} \left( \frac{10\text{Mpc}}{r} \right) \times \sqrt{N}$$

$$N \Rightarrow 10 - 10^8$$

# Magnetized Stars

## Observations



### ✓ Giant flares in SGRs

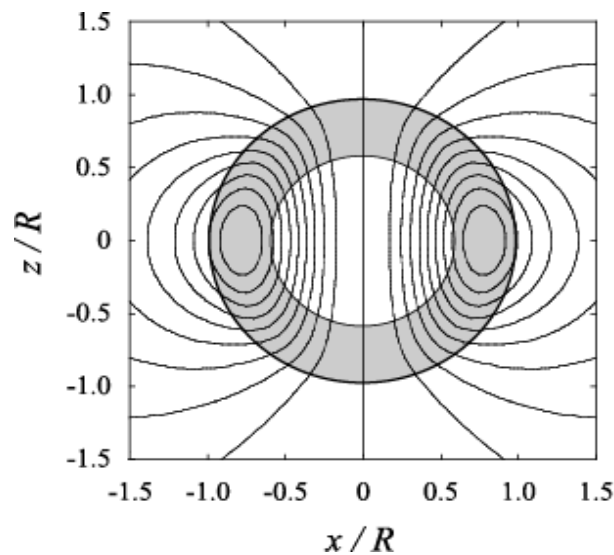
- Up to now, three giant flares have been detected.
  - *SGR 0526-66* in 1979,
  - *SGR 1900+14* in 1998,
  - *SGR 1806-20* in 2004
- Peak luminosities :  $10^{44} - 10^{46}$  erg/s
- A decaying tail for several hundred seconds follows the flare.

### ✓ QPOs in decaying tail (Israel *et al.* 2005; Watts & Strohmayer 2005, 2006)

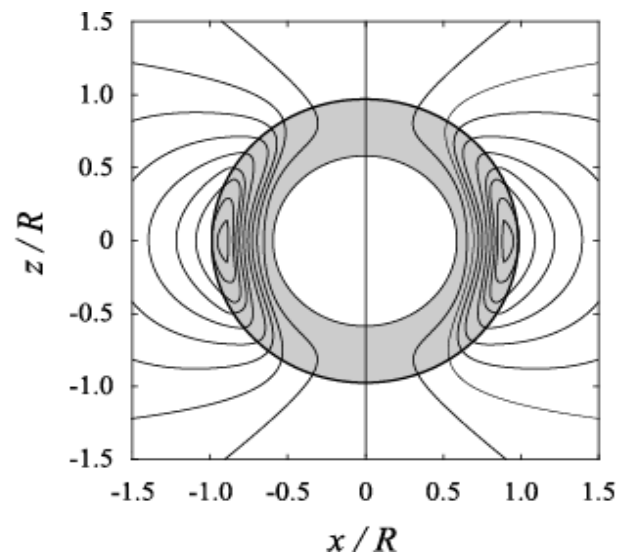
- *SGR 1900+14* : 28, 54, 84, and 155 Hz
- *SGR 1806-20* : 18, 26, 29, 92.5, 150, 626.5, & 1837 Hz  
(possible additional frequencies : 720 & 2384 Hz)

# B-field configuration

- ✓ Axisymmetric dipole magnetic fields with poloidal and toroidal components.
- ✓ Two types of B-field geometries;
  - I. B-fields permeating star.
  - II. Magnetic fields are confined in the crust region.



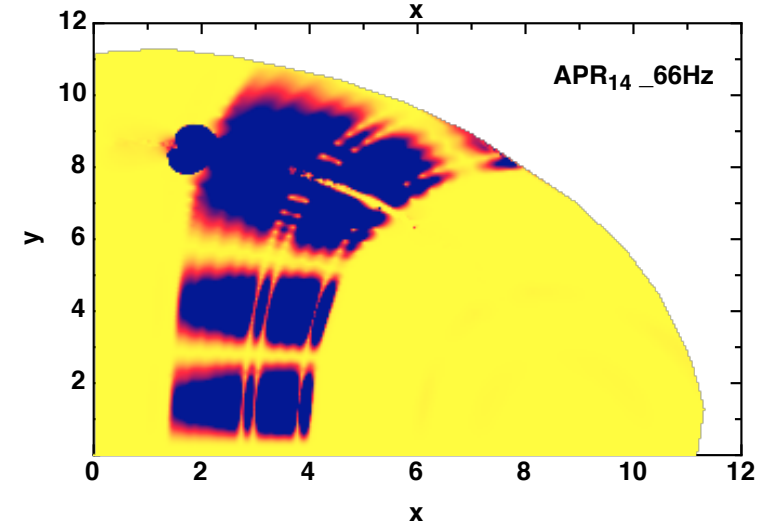
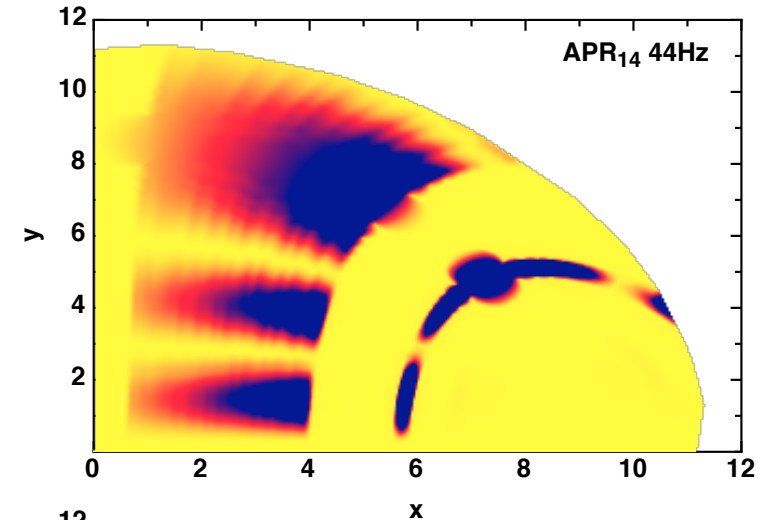
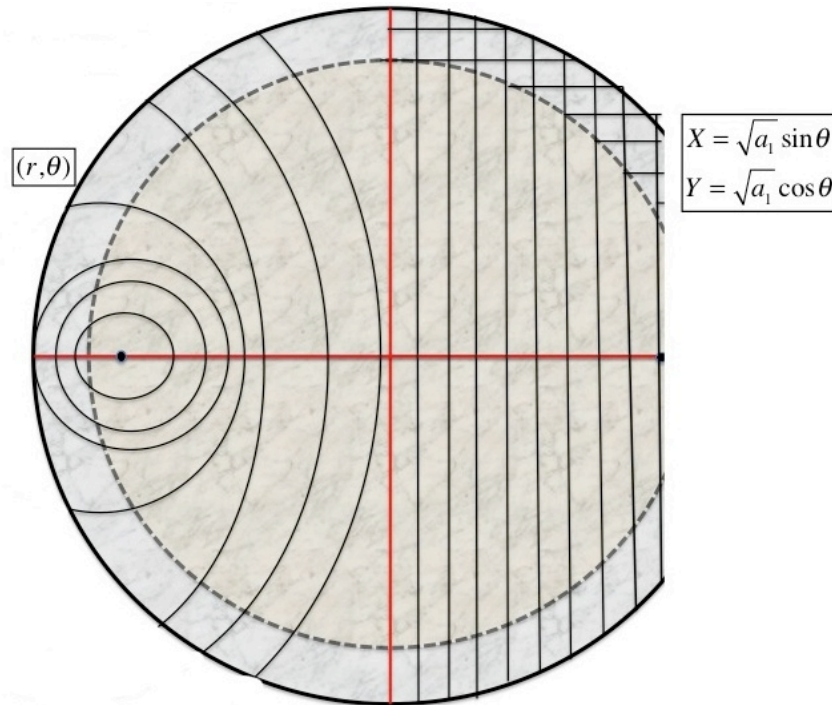
(i)



(ii)

- Piro 2005
- Glampedakis, Samuelsson, Andersson 2006
- Sotani, KK, Stergioulas 2007
- Samuelsson, Andersson 2007
- Levin 2007
- Sotani, KK, Stergioulas 2008
- Vavoulidis, Stavridis, KK 2008
- Sotani, Colaiuda, KK 2008
- Colaiuda, Beyer, KK 2009
- Cerde-Duran, Stergioulas, Font 2009
- Sotani, KK 2009
- Steiner, Watts 2009
- Lander, Jones, Passamonti 2009
- Van Hoven, Levin 2010
- Cerde-Duran et al 2010
- Colaiuda+ KK 2010

# Alfven Continuum + Discrete oscillations



## Without Crust

- Levin 2007
- Sotani, KK, Stergioulas 2008
- Colaiuda, Beyer, KK 2009
- Cerda-Duran, Stergioulas, Font 2009

## With Crust

- Van Hoven, Levin 2010
- Cerda-Duran, Stergioulas, Font 2010
- Colaiuda, KK 2010



# Attempt to Fit the Observational data

- ✓ **Axial type** of perturbations described by **discrete modes** (crust) and a **continuum** (core) and *they can explain the lower frequencies* observed
- ✓ **Polar type** of perturbations only **discrete modes** with higher frequencies

Alfvén modes

$$f_n^{\text{even}} \approx (2n + 1)f_0$$

$$f_n^{\text{odd}} \approx (n + 1)f_0$$

$$f_n^{\text{C}} \approx (n + 1)f_0^{\text{C}}$$

$${}_{\ell}t_n \approx \left[ 1 + {}_{\ell}a_n \left( \frac{B}{B_{\mu}} \right)^2 \right]^{1/2} \cdot {}_{\ell}t_0$$

Crust modes

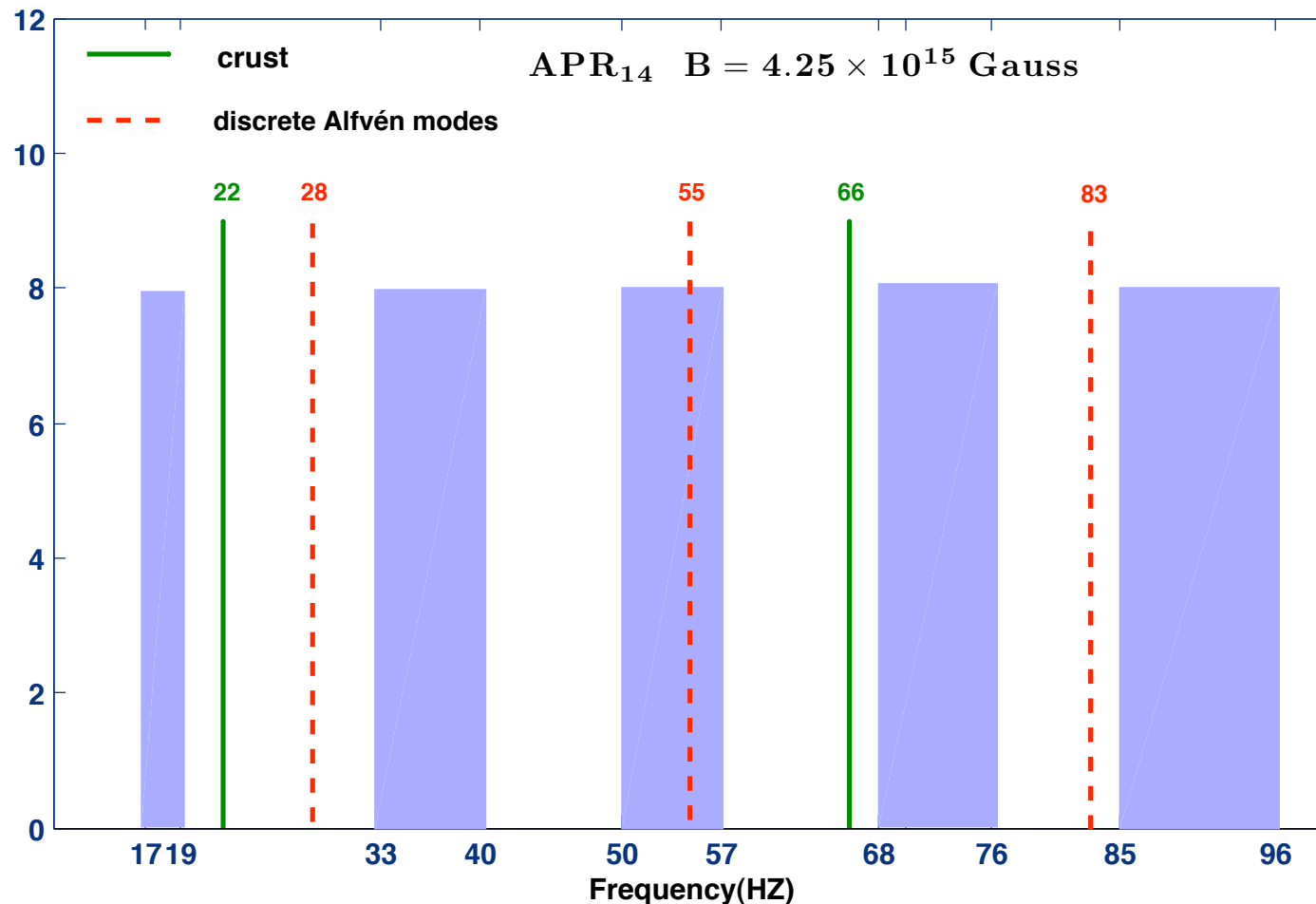
Sotani, KK, Stergioulas 2007

- On the other hand the observed frequencies of QPOs in SGRs
  - SGR 1900+14 : 28, 54, 84, 155 Hz
    - x2    x3    ← crust torsional oscillation ? or polar oscillation ?
  - SGR 1806-20 : 18, 26, 30, 92.5, 150 Hz
    - x3    x5
    - 0.6    ← crust torsional oscillation?

EoS => APR+DH,  
 M=1.4 M<sub>⊙</sub>, R=12.1Km  
 ΔR/R=0.93km,  
 B=4x10<sup>15</sup>Gauss

# SGR 1806-20

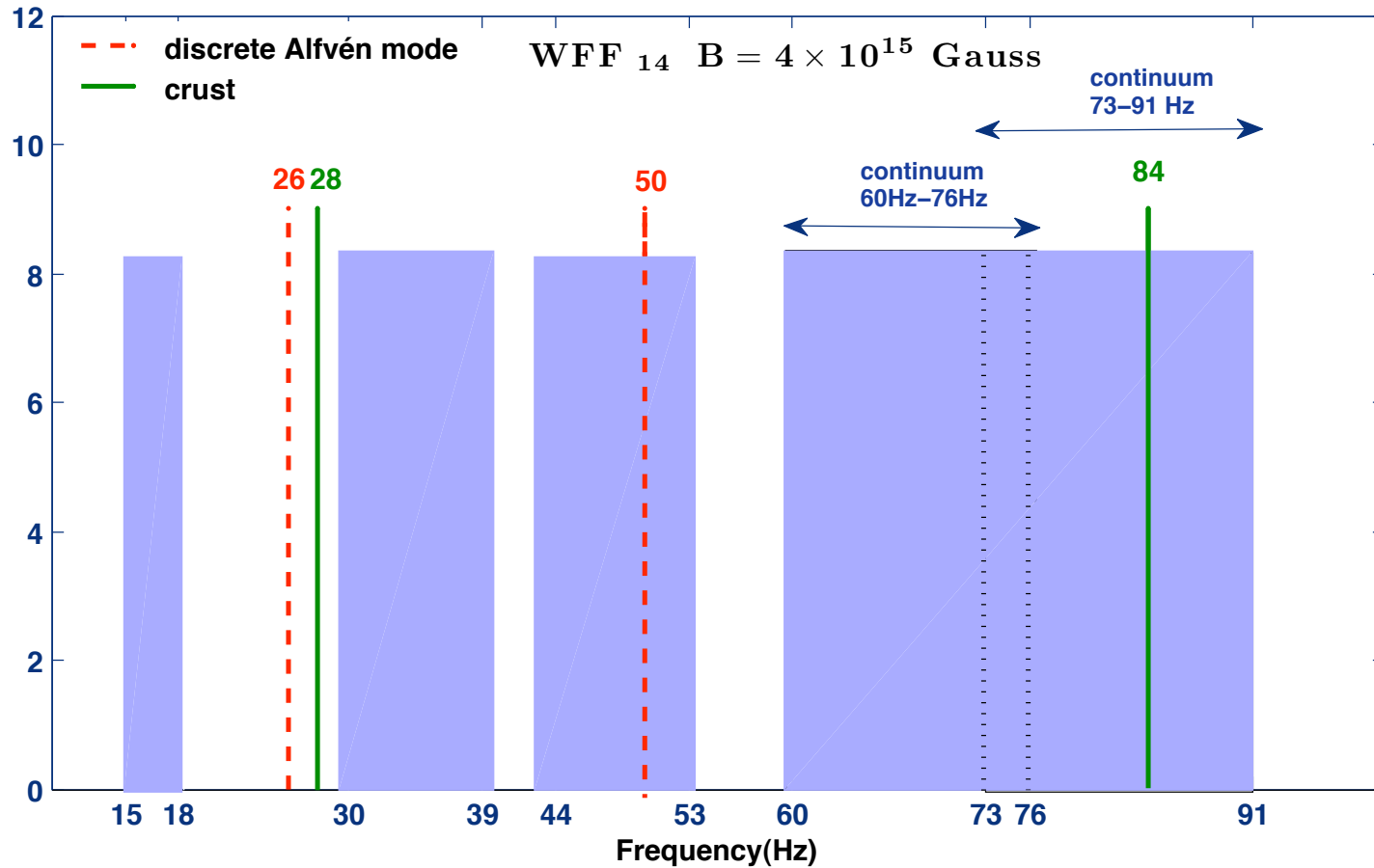
Colaiuda, KK 2010



Identification of the frequencies of **SGR 1806-20**. We show that a stellar model **APR** with mass  $M = 1.4M_{\odot}$  can explain all the frequencies observed.

# SGR 1900+14

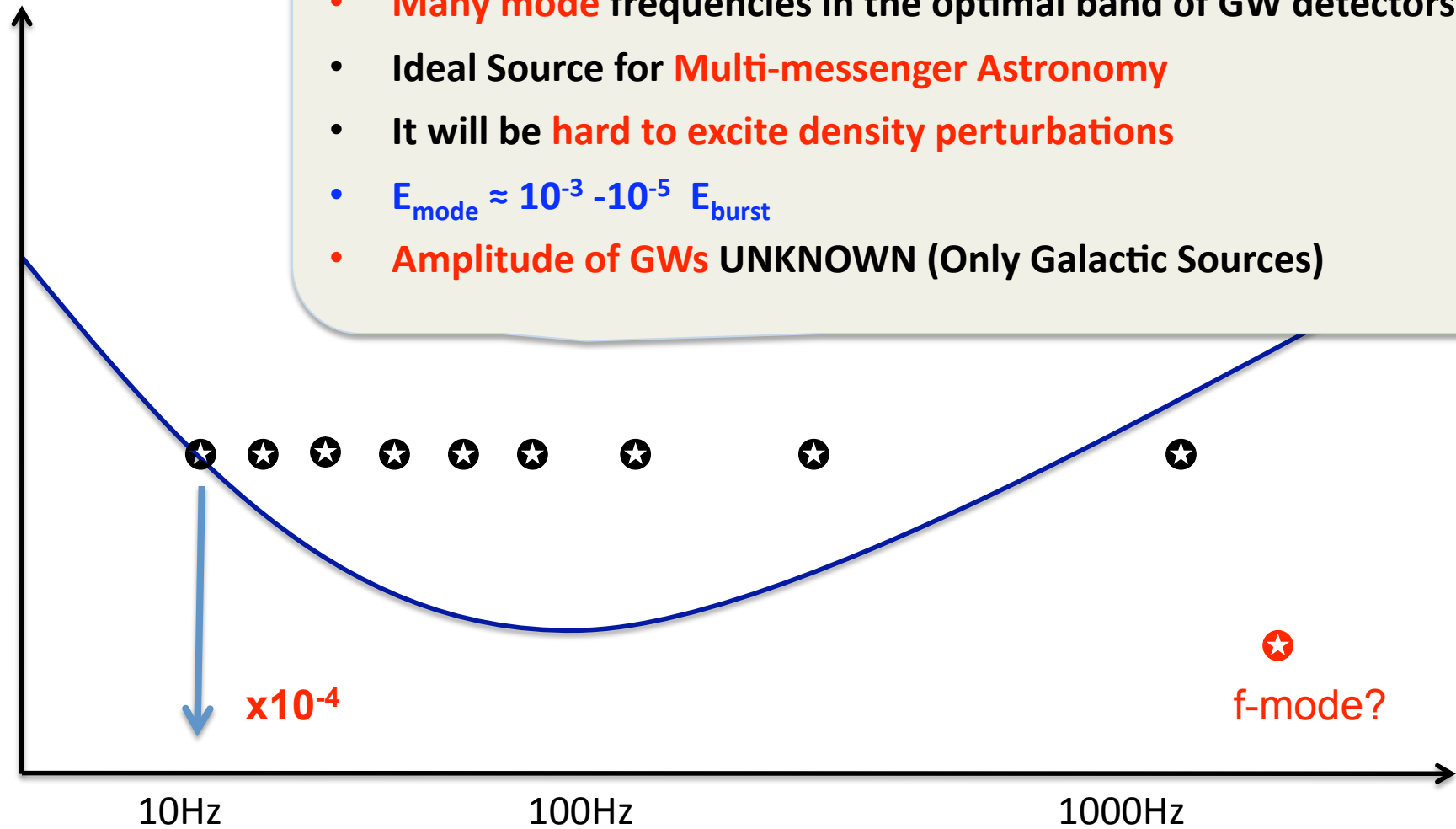
Colaiuda, KK 2010



Identification of the frequencies of **SGR 1900+14**. We show that a stellar model **WFF** with mass  $M = 1.4M_{\odot}$  can explain all the frequencies observed.

# Magnetars & GWs

- Many mode frequencies in the optimal band of GW detectors
- Ideal Source for Multi-messenger Astronomy
- It will be hard to excite density perturbations
- $E_{\text{mode}} \approx 10^{-3} - 10^{-5} E_{\text{burst}}$
- Amplitude of GWs UNKNOWN (Only Galactic Sources)



# Conclusions

- **Rotational Instabilities of Neutron Stars**
  - ✓ Are potential sources for GW **beyond** our galaxy
  - ✓ Many open issues (growth time, non-linear coupling,...) **are resolved one after the other.**
- **Dynamics of magnetars**
  - ✓ Offers the possibility to **understand their structure**
  - ✓ Most probably a **weak source** for GW with the present generation detectors