

Hamilton Approach to QCD in Coulomb Gauge

H. Reinhardt

Tübingen

Collaborators:

G. Burgio, D. Campagnari, M. Quandt, P. Watson

D. Epple, C. Feuchter, W. Schleifenbaum,
J. Heffner, M. Leder, M. Pak, C. Popovici ,

J. Pawlowski, A. Szczepaniak, A. Weber

Hamilton approach to QCD

- Introduction
- Yang-Mills theory at $T=0$
- Yang-Mills theory at finite T
- QCD
- Conclusions

Canonical Quantization of Yang-Mills theory

Weyl gauge: $A_0^a(x) = 0$ cartesian coordinates $A_i^a(x)$

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$$H\Psi[A] = E\Psi[A]$$

gauge invariant wave functionals: $\Psi[A]$ K. Johnson et al.....

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more convenient: gauge fixing

$$\partial A = 0$$

Hamiltonian approach to YMT in Coulomb gauge $\partial A = 0$

$$H = \frac{1}{2} \int (J^{-1} \Pi^\perp J \Pi^\perp + B^2) + H_C \quad \Pi^\perp = \delta / i \delta A^\perp$$

Christ and Lee

$$J(A^\perp) = \text{Det}(-D\partial) \quad D = \partial + gA$$

$$H_C = \frac{1}{2} \int J^{-1} \rho (-D\partial)^{-1} (-\partial^2) (-D\partial)^{-1} J \rho \quad \text{Coulomb term}$$

color charge density: $\rho = -A^\perp \Pi^\perp + \rho_m$

$$\langle \Phi | \dots | \Psi \rangle = \int_{\Lambda} D A J(A) \Phi^*(A) \dots \Psi(A)$$

$$H\Psi[A] = E\Psi[A]$$

Perturbation theory

D. Campagnari, H. R. & A. Weber, Phys. Rev D(2009)

■ Rayleigh-Schrödinger PT

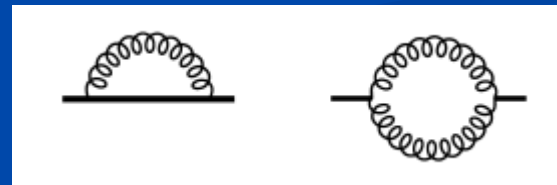
$$\tilde{H} = H_0 + g\tilde{H}_1 + g^2\tilde{H}_2 + \mathcal{O}(g^3),$$

$$\left[|0\rangle + g|0\rangle^{(1)} + g^2|0\rangle^{(2)} + \mathcal{O}(g^3) \right].$$

■ vacuum (QED)

$$\langle A|0\rangle = \mathcal{N} \exp \left\{ -\frac{1}{2} \int \bar{d}k A_\sigma^a(\mathbf{k}) |\mathbf{k}| A_\sigma^a(-\mathbf{k}) \right\}.$$

■ β -function



$$v(\mathbf{k}) = 1 + g^2 \frac{N_c}{(4\pi)^{2-\varepsilon}} \left\{ \frac{11}{3} \left[\frac{1}{\varepsilon} - \gamma - \ln \frac{\mathbf{k}^2}{\mu^2} \right] + \frac{31}{9} + \mathcal{O}(\varepsilon) \right\}.$$

$$\beta(g) = \frac{\partial g}{\partial \ln \mu} = \frac{1}{(4\pi)^2} \beta_0 g^3 + \mathcal{O}(g^5),$$

$$\beta_0 = -\frac{11}{3} N_C$$

Variational approach

- Gaussian ansatz,

$$\Psi(A) = \exp\left[-\frac{1}{2} \int dx dy A(x) \omega(x, y) A(y)\right]$$

D. Schütte 1984

.....

A. Szczepaniak & E. Swanson 2002

C. Feuchter & H. R. 2004

-ansatz

-FP determinant

-renormalization

- Greensite, Matevosyan, Olejnik, Quandt, Reinhardt, Szczepaniak, PRD83

Variational approach

- trial ansatz

C. Feuchter & H. R. PRD70(2004)

$$\Psi(A) = \frac{1}{\sqrt{\text{Det}(-D\partial)}} \exp\left[-\frac{1}{2} \int dx dy A(x) \omega(x, y) A(y)\right]$$

- : gluon propagator

$$\langle A(x) A(y) \rangle = (2\omega(x, y))^{-1}$$

variational kernel

$\omega(x, x')$

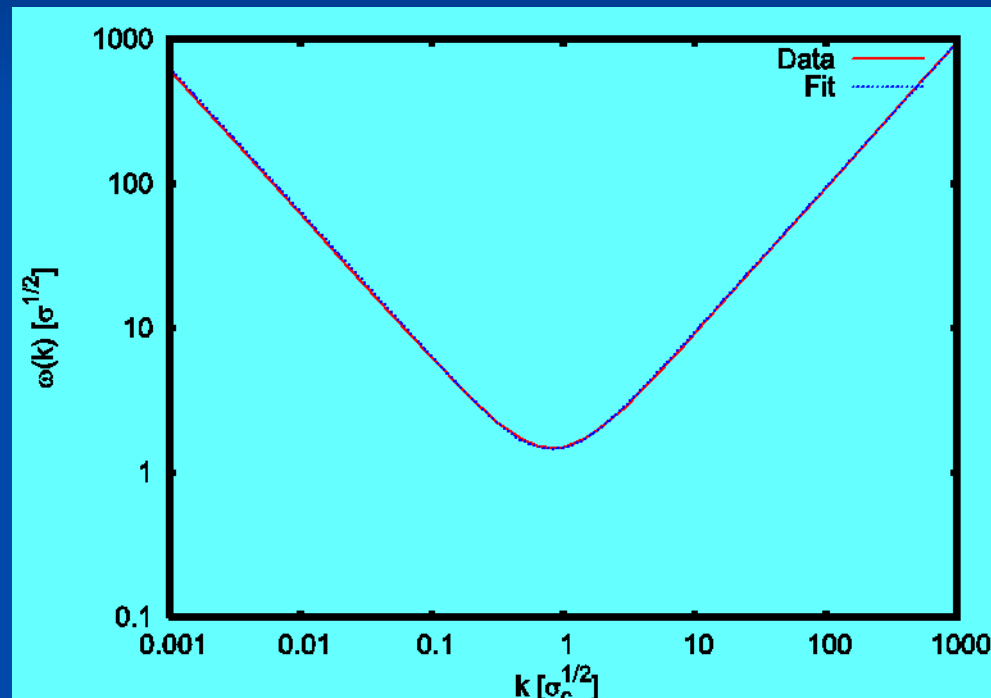
determined from

$$\langle \Psi | H | \Psi \rangle \rightarrow \min$$

Numerical results

D. Epple, H. R., W. Schleifenbaum,
PRD 75 (2007)

gluon energy



$$IR: \omega(k) \sim 1/k, \quad UV: \omega(k) \sim k,$$

Static gluon propagator in D=3+1

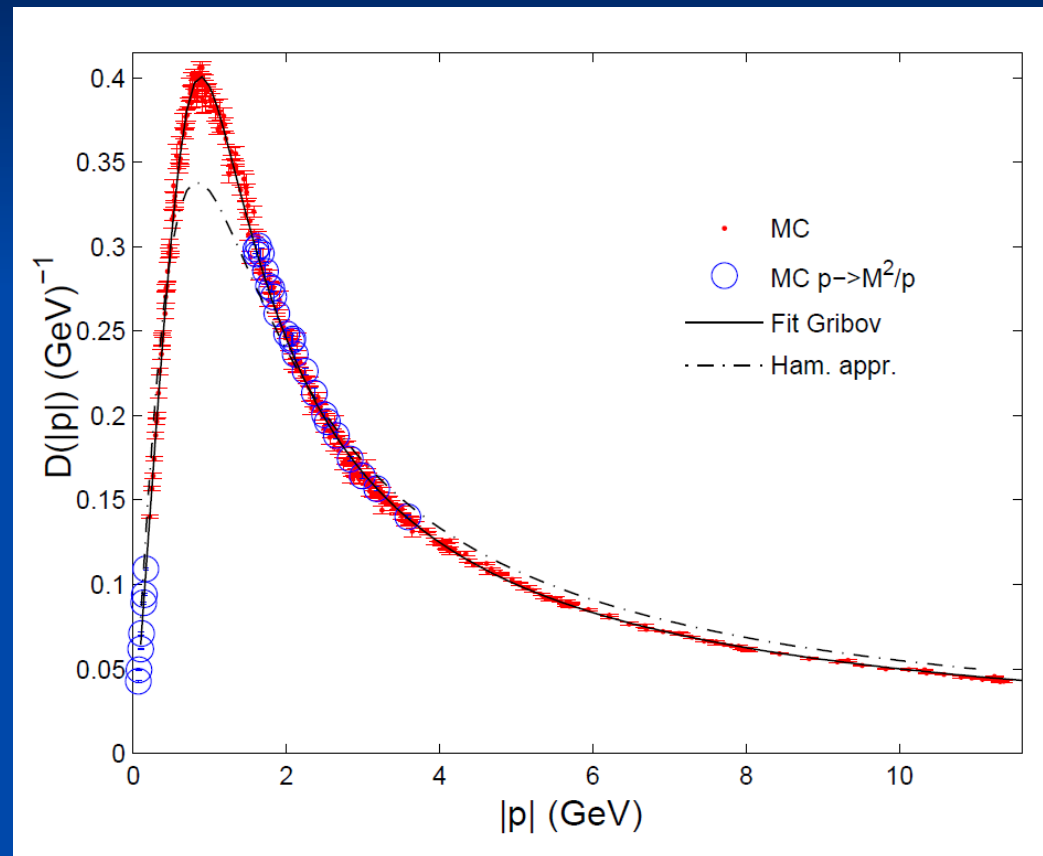
$$D(k) = (2\omega(k))^{-1}$$

Gribov's formula

$$\omega(k) = \sqrt{k^2 + \frac{M^4}{k^2}}$$

$$M = 0.88 \text{ GeV}$$

missing strength in
mid momentum regime:
missing gluon loop



G. Burgio, M.Quandt , H.R., PRL102(2009)

Variational approach to YMT with non-Gaussian wave functional

D. Campagnari & H.R,
Phys.Rev.D82(2010)
See talk by D. C.

wave functional

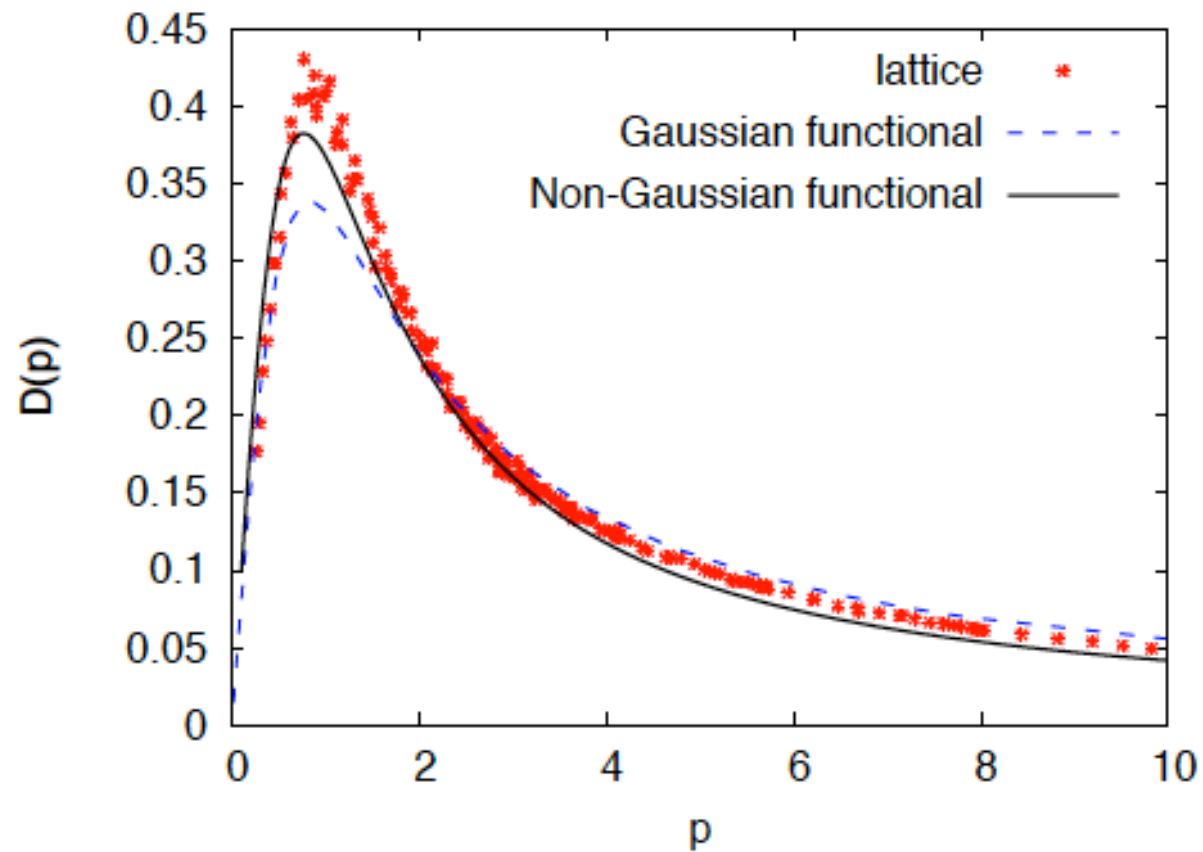
$$|\psi[A]|^2 = \exp(-S[A])$$

ansatz

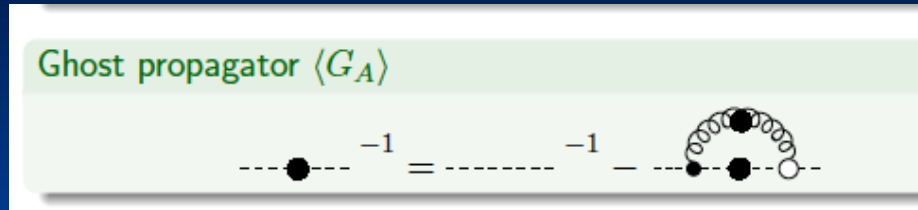
$$S[A] = \int \omega A^2 + \frac{1}{3!} \int \gamma^{(3)} A^3 + \frac{1}{4!} \int \gamma^{(4)} A^4$$

exploit DSE

Corrections to the gluon propagator

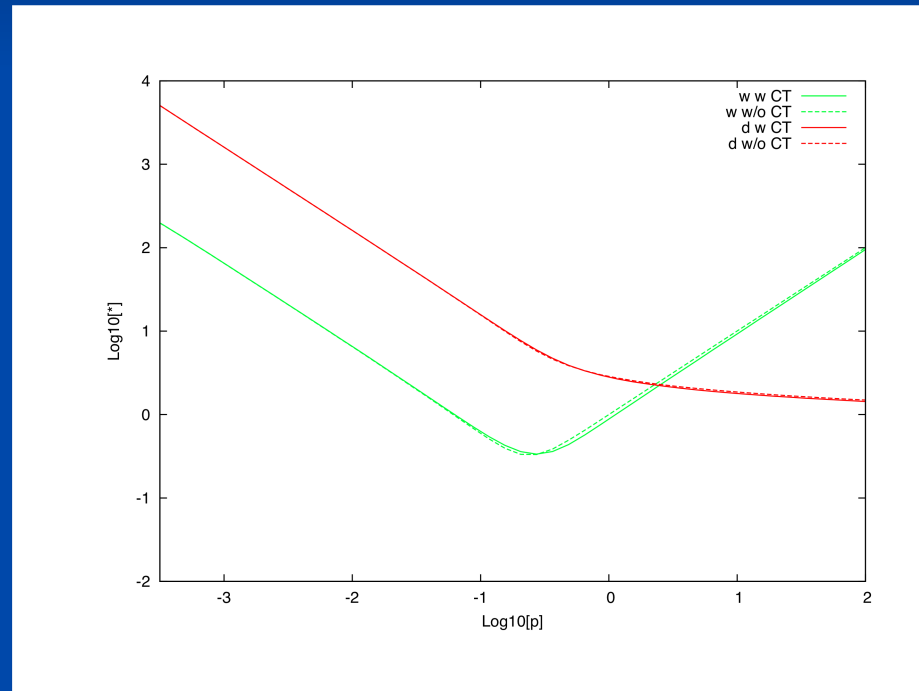


ghost propagator : $\langle (-D\partial)^{-1} \rangle = d(k) / k^2$



horizon condition $d^{-1}(0) = 0$

supported by the lattice



The color dielectric function of the QCD vacuum

- ghost propagator

$$\langle (-D\partial)^{-1} \rangle = d / (-\Delta)$$

- dielectric „constant“

$$\varepsilon = d^{-1}$$

H.R. PRL101 (2008)

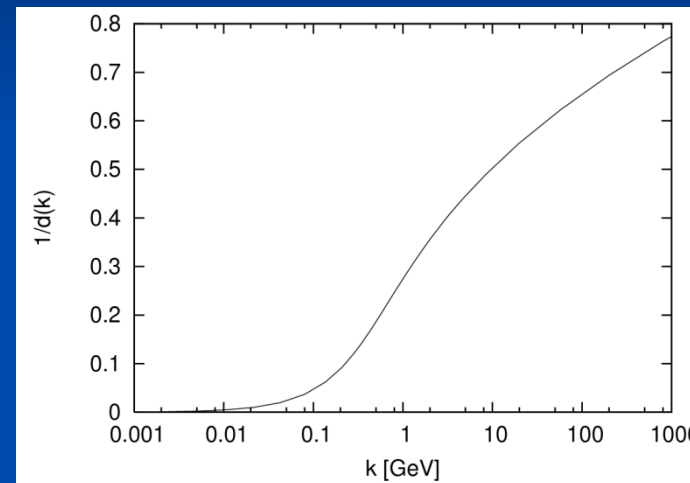
- horizon condition:

- : $d^{-1}(k=0) = 0$ $\varepsilon(k=0) = 0$

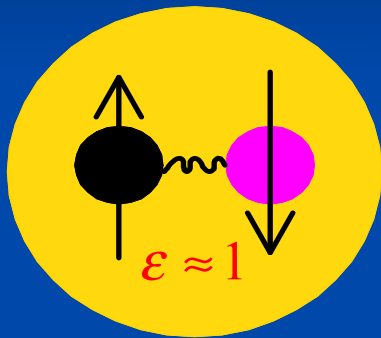
- QCD vacuum: perfect color dia-electricum

- dual superconductor

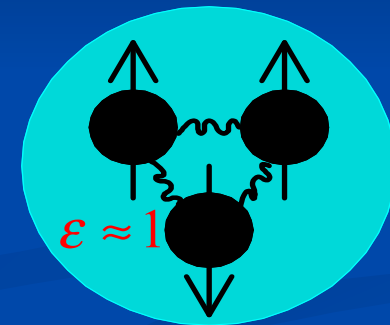
- : $\varepsilon(k) < 1$ anti-screening



$$D = \epsilon E \quad \partial D = \rho_{free}$$



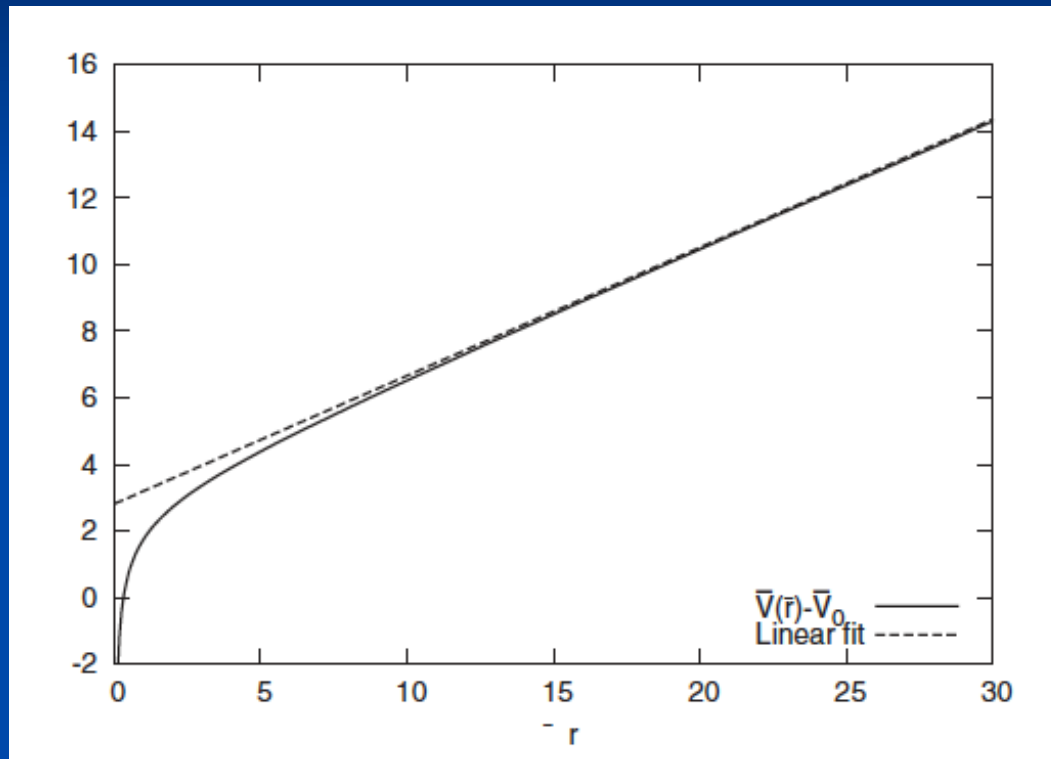
$$\epsilon = 0$$



no free color charges in the vacuum: confinement

Static Coulomb potential

$$V(|x-y|) = g^2 \langle \langle x | (-D\partial)^{-1} (-\partial^2) (-D\partial)^{-1} | y \rangle \rangle$$

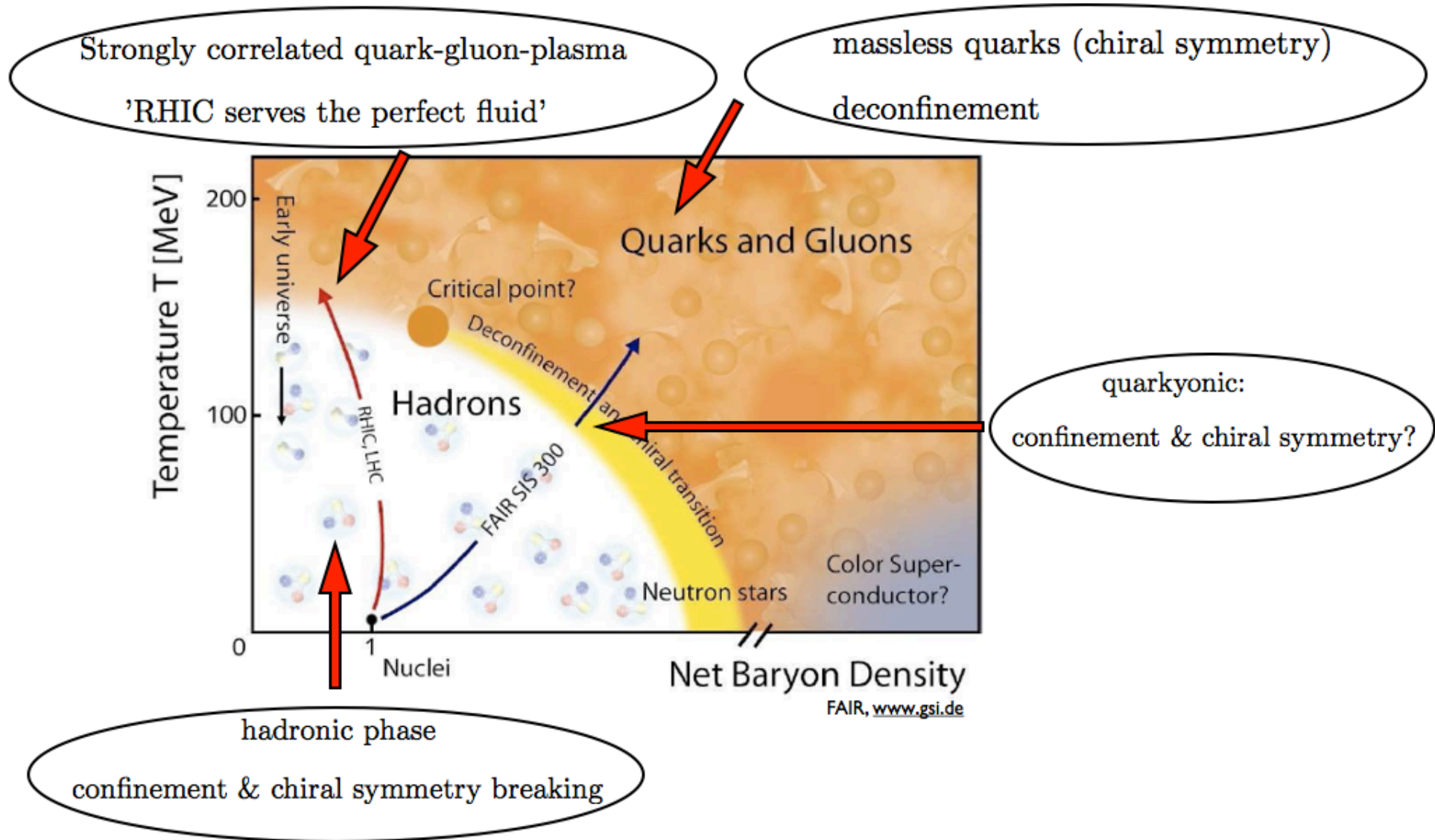


D. Epple, H. Reinhardt
W. Schleifenbaum,
PRD 75 (2007)

$$V(R) \xrightarrow{R \rightarrow \infty} \sigma_C R, \quad \text{lattice } \sigma_C = 2 \dots 3 \sigma_w$$

$$V(R) \xrightarrow{R \rightarrow 0} \sim 1/R$$

Phase diagram of QCD



Hamiltonian approach to YMT at finite T

Reinhardt, Campagnari, Szczepaniak, PRD,2011
Heffner, Reinhardt, Campagnari, to be published

- Grand canonical ensemble with $\mu = 0$
 - density matrix

$$D = \exp(-\beta H) \quad \beta = 1 / k_B T$$

- thermal averages

$$\langle \dots \rangle_T = \frac{\text{Tr}(D \dots)}{\text{Tr} D}, \quad \text{Tr}(\dots) = \sum_k \langle \tilde{k} | \dots | \tilde{k} \rangle$$

- complete basis $|\tilde{k}\rangle$

Grand canonical ensemble

- choice of basis

$$|\tilde{k}\rangle = \frac{1}{\sqrt{\text{Det}(-D\partial)}} |k\rangle$$

gauge field

$$A(k) = \frac{1}{\sqrt{2\omega(k)}} (a(k) + a^\dagger(-k)), \quad \omega(k) - \text{arbitrary}$$

vacuum

$$a(k)|0\rangle = 0, \quad \langle A|0\rangle = \exp\left(-\frac{1}{2} \int A\omega A\right)$$

complete basis

$$|0\rangle, \quad a^\dagger(k)|0\rangle, \quad a^\dagger(k)a^\dagger(k)|0\rangle, \dots$$

Grand canonical ensemble

- density operator $D = \exp(-\beta H)$

$$D = \exp(-\beta \tilde{H}), \quad \tilde{H} = \int dk \Omega(k) a^\dagger(k) a(k)$$

- occupation numbers

$$\langle a^\dagger a \rangle \equiv n(k) = [\exp(\beta \Omega(k)) - 1]^{-1}$$

- gluon propagator

$$\langle AA \rangle = (1 + 2n) / (2\omega)$$

- entropy $S = -k_B \text{Tr}(D \ln D)$

- free energy $F[\Omega, \omega] = \langle H \rangle_T - TS$

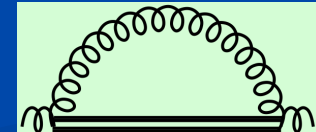
Finite T variational principle

$$n(k) = [\exp(\beta\Omega(k)) - 1]^{-1} \quad \langle A|0\rangle = \exp(-\frac{1}{2} \int A\omega A)$$

- free energy $F[\Omega, \omega] = E[\Omega, \omega] - TS[\Omega] \rightarrow \min$
- variation w.r.t. the density operator

$$\delta F[n, \omega] / \delta n(k) = 0 \quad \Rightarrow \quad \Omega(k) = \omega(k)(1 + I_C(k))$$

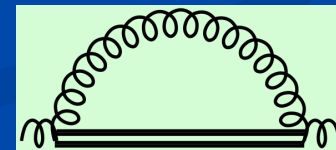
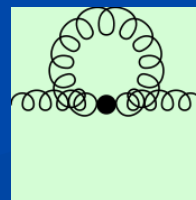
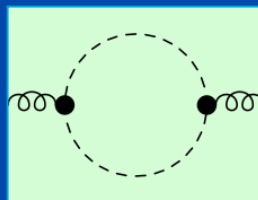
quasi-gluon energy



- variation w.r.t. the basis $\delta F[\Omega, \omega] / \delta \omega(k) = 0$

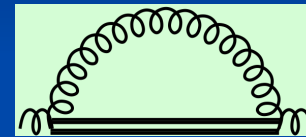
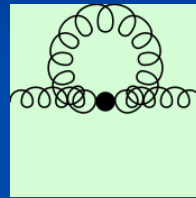
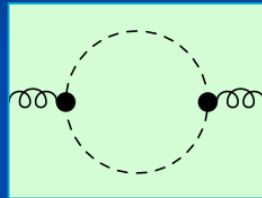
$$\omega^2(k) = k^2 + \chi^2(k) + I_{tad}[n] + I_C[n](k)$$

gap equation

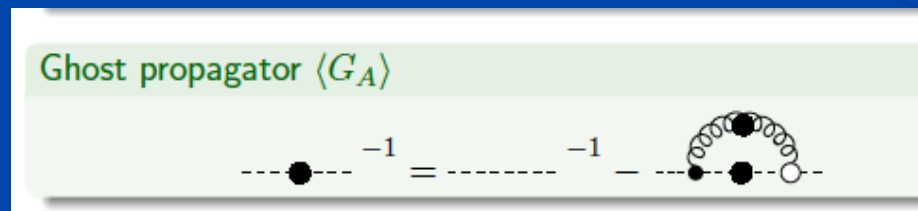


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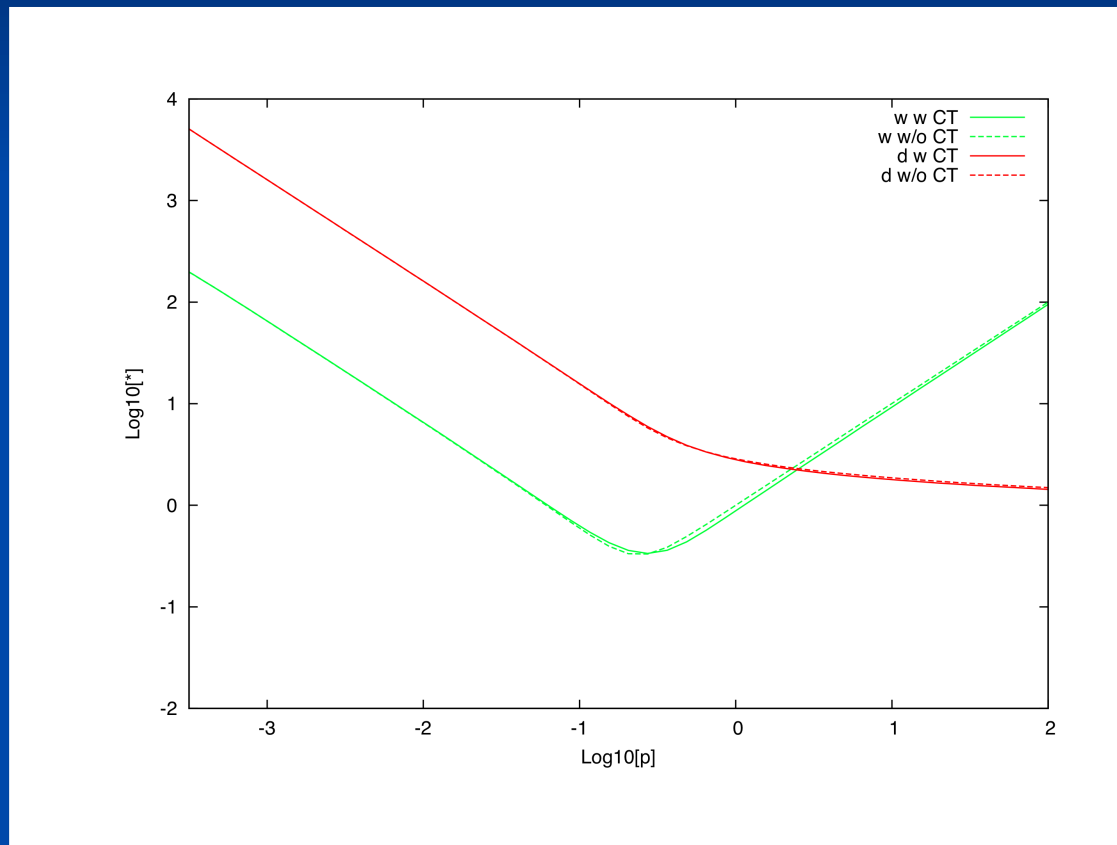
renormalization as at $T=0$



The influence of the Coulomb term

$$H_C = \frac{1}{2} \int J^{-1} \Pi J^{\parallel} \Pi^{\parallel} = \frac{1}{2} \int J^{-1} \rho (-D\partial)^{-1} (-\partial^2) (-D\partial)^{-1} J \rho$$

color charge density: $\rho = -A^{\perp} \Pi^{\perp} + \rho_m$

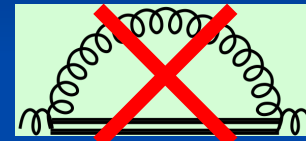
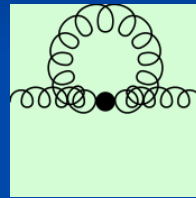
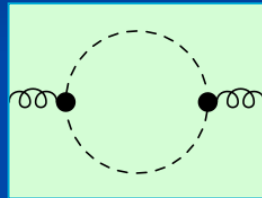


–irrelevant in the YM sector
–utterly important for quarks:
confining potential

Without Coulomb term

Heffner, Reinhardt & Campagnari

$$\omega^2(k) = k^2 + \chi^2(k) + I_{\text{tad}}[n] + I_C[n] \quad \text{X}$$



quasi-gluon energy

$$\Omega(k) = \omega(k)$$

vacuum $\langle A|0\rangle = \exp(-\frac{1}{2} \int A\omega A)$

occupation number $n(k) = [\exp(\beta\Omega(k)) - 1]^{-1}$

Infrared analysis

gluon energy

$$\omega(p) = A / p^\alpha$$

ghost form factor

$$d(p) = B / p^\beta$$

$T = 0$ *sum rule*

$$\alpha = 2\beta + 2 - d$$

$$d = 3$$

$$\beta = 1.0(0.99)$$

$$\beta = 0.796(0.85)$$

$$d = 2$$

$$\beta = 0.5(0.45)$$

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at arbitrarily finite T infrared analysis=impossible! $n(k) = [\exp(\beta\omega(k)) - 1]^{-1}$

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$$T \rightarrow \infty$$

$$n(k) \simeq 1 / \beta\omega(k)$$

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sum rule

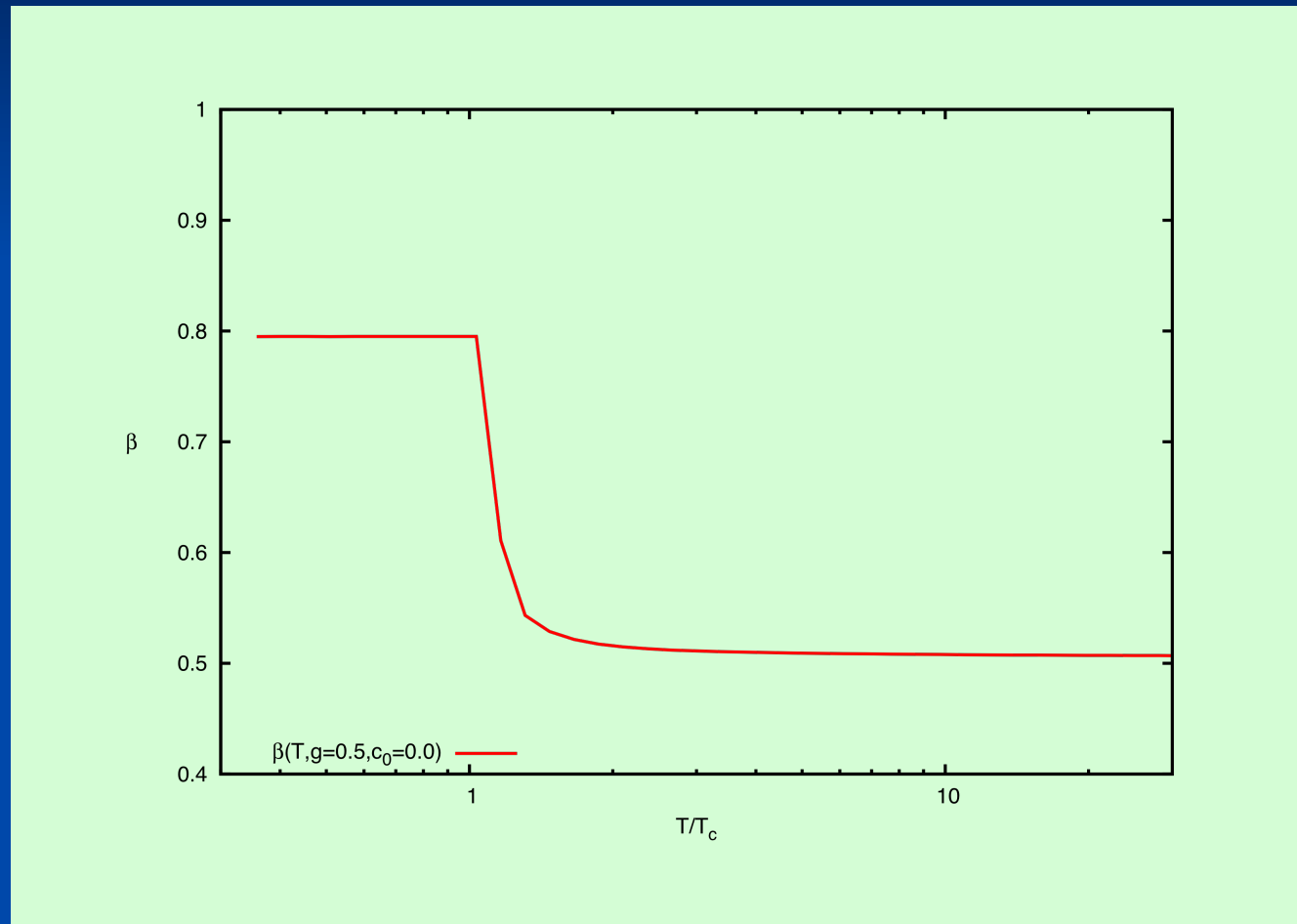
$$\alpha = 2\beta + 2 - d$$

$$d = 3$$

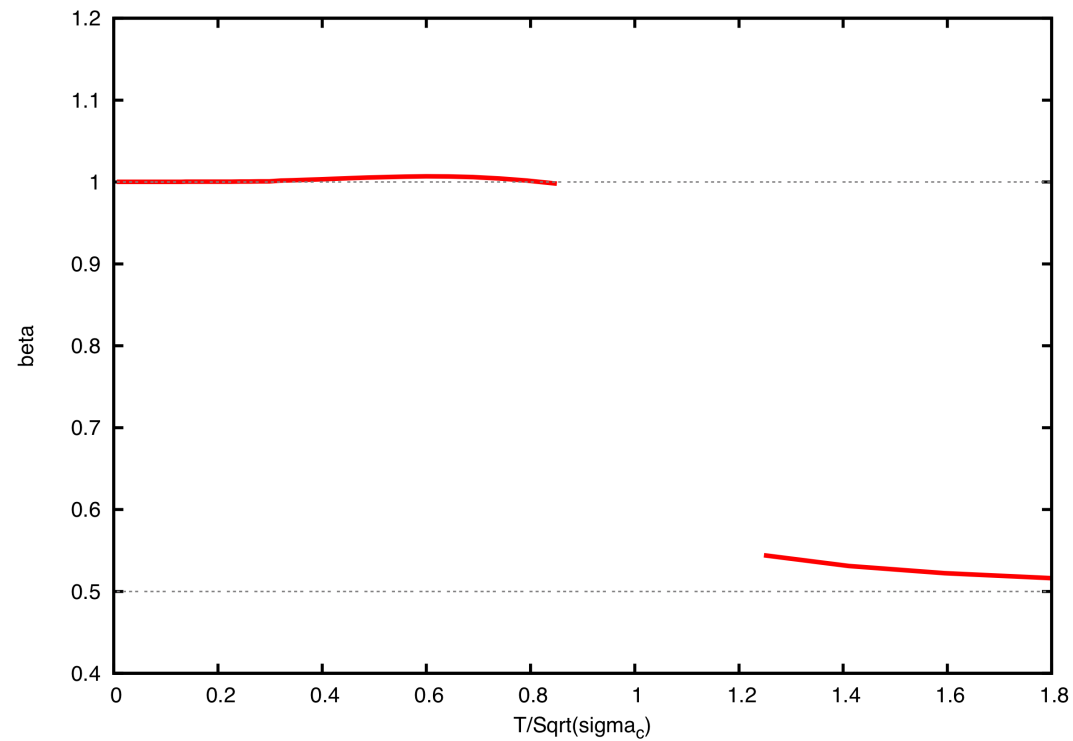
$$\beta = 0.5$$

$$\alpha = 0$$

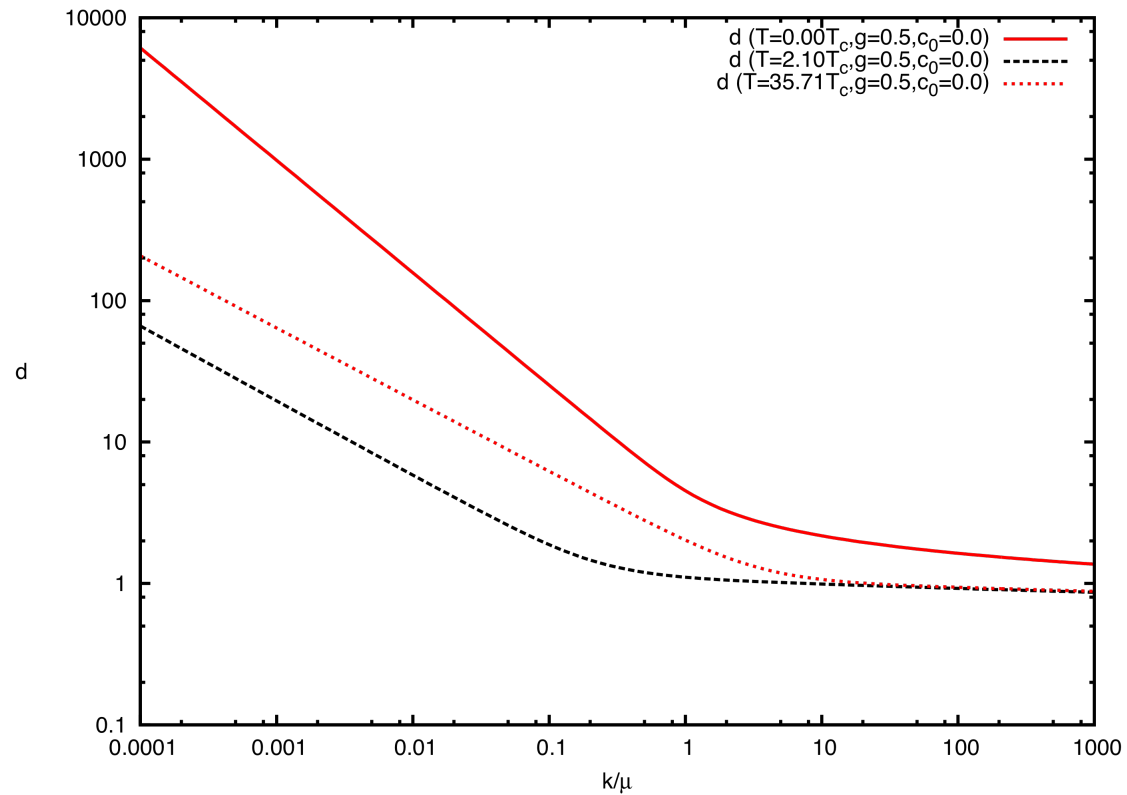
Infrared exponent of ghost



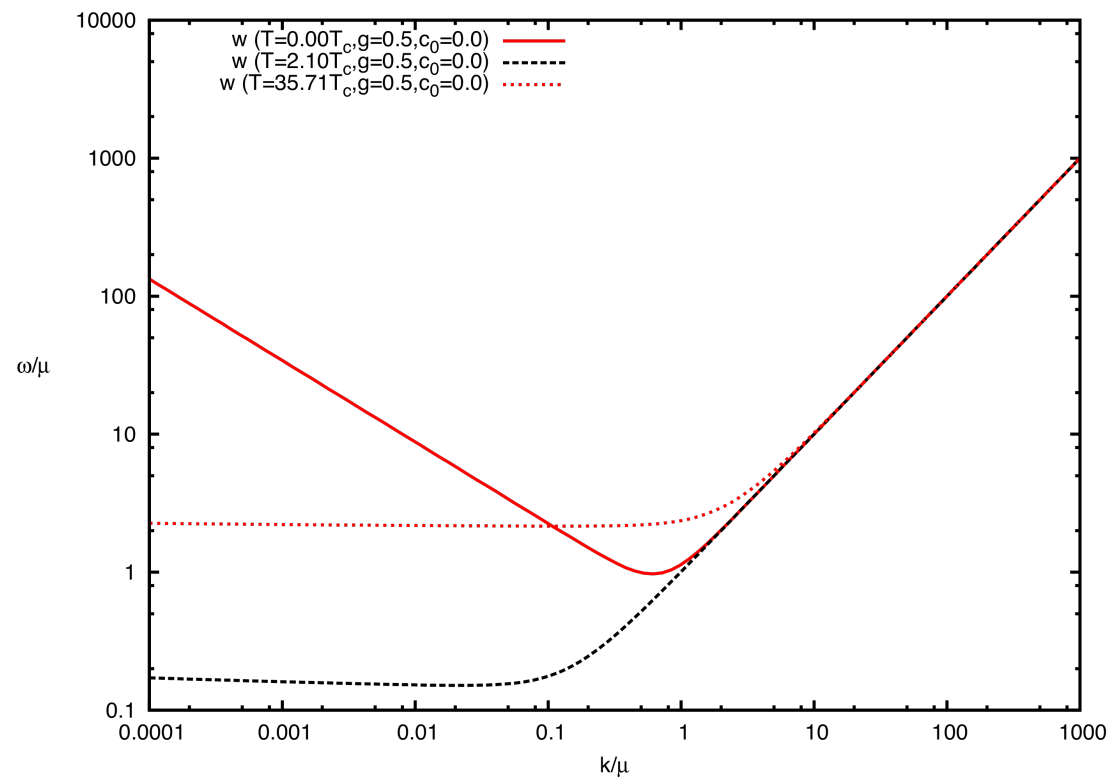
Infrared exponent of ghost



Ghost form factor

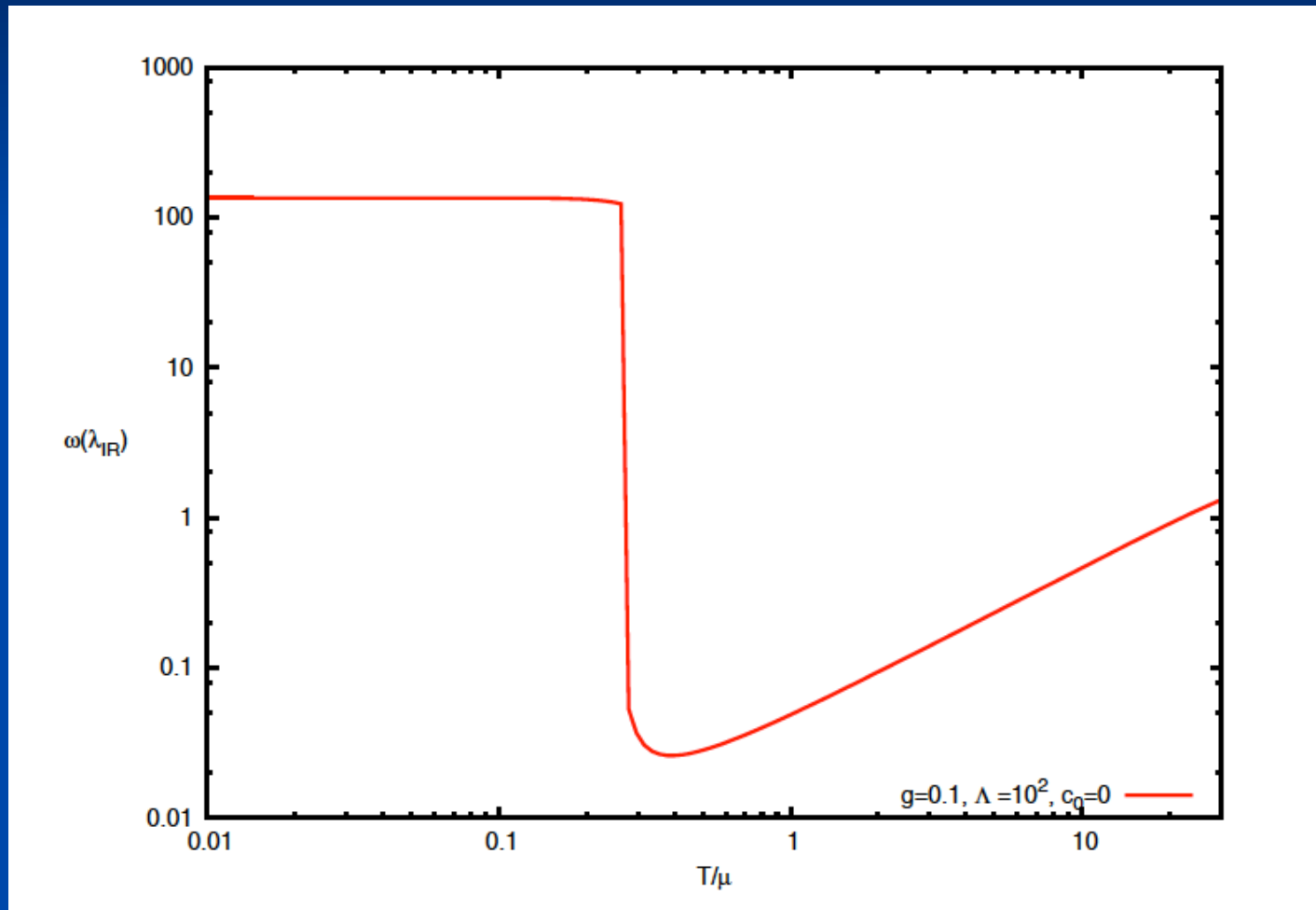


Gluon energy



Effective gluon mass (IR-value of gluon energy)

$$m(T) = \omega(k = \lambda_{IR}, T)$$



Critical temperature

$$T_C = 162 \text{ MeV} \sqrt{\sigma_C / \sigma_W} \quad \begin{array}{l} \sigma_C - \text{Coulomb string tension} \\ \sigma_W - \text{Wilson string tension} \end{array}$$

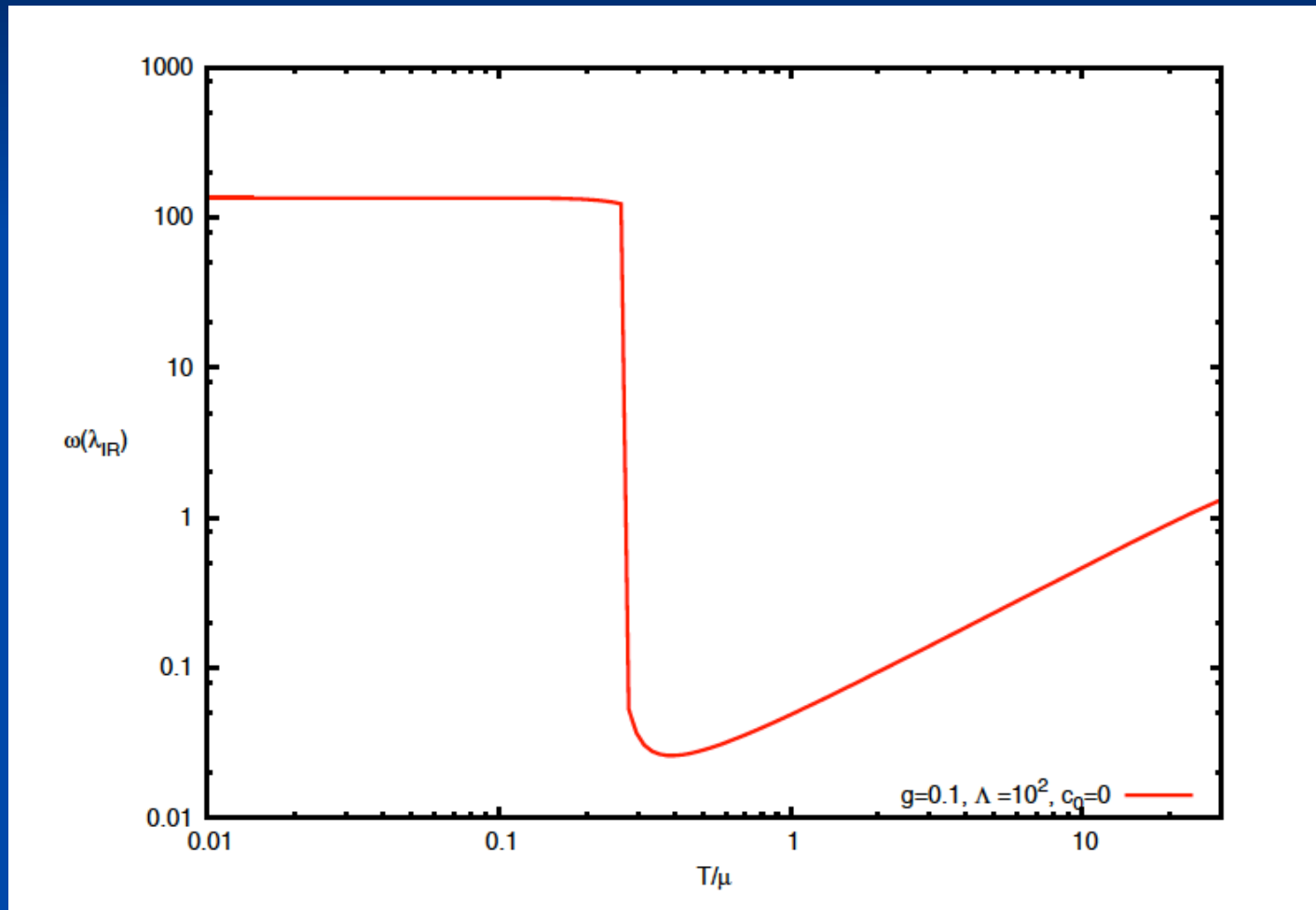
$$\text{SU}(2) - \text{lattice} : \sigma_C = 2 \dots 3 \sigma_W$$

$$\Rightarrow T_C = 230 \dots 280 \text{ MeV}$$

$$\text{SU}(2) - \text{lattice} : T_C = 295 \text{ MeV}$$

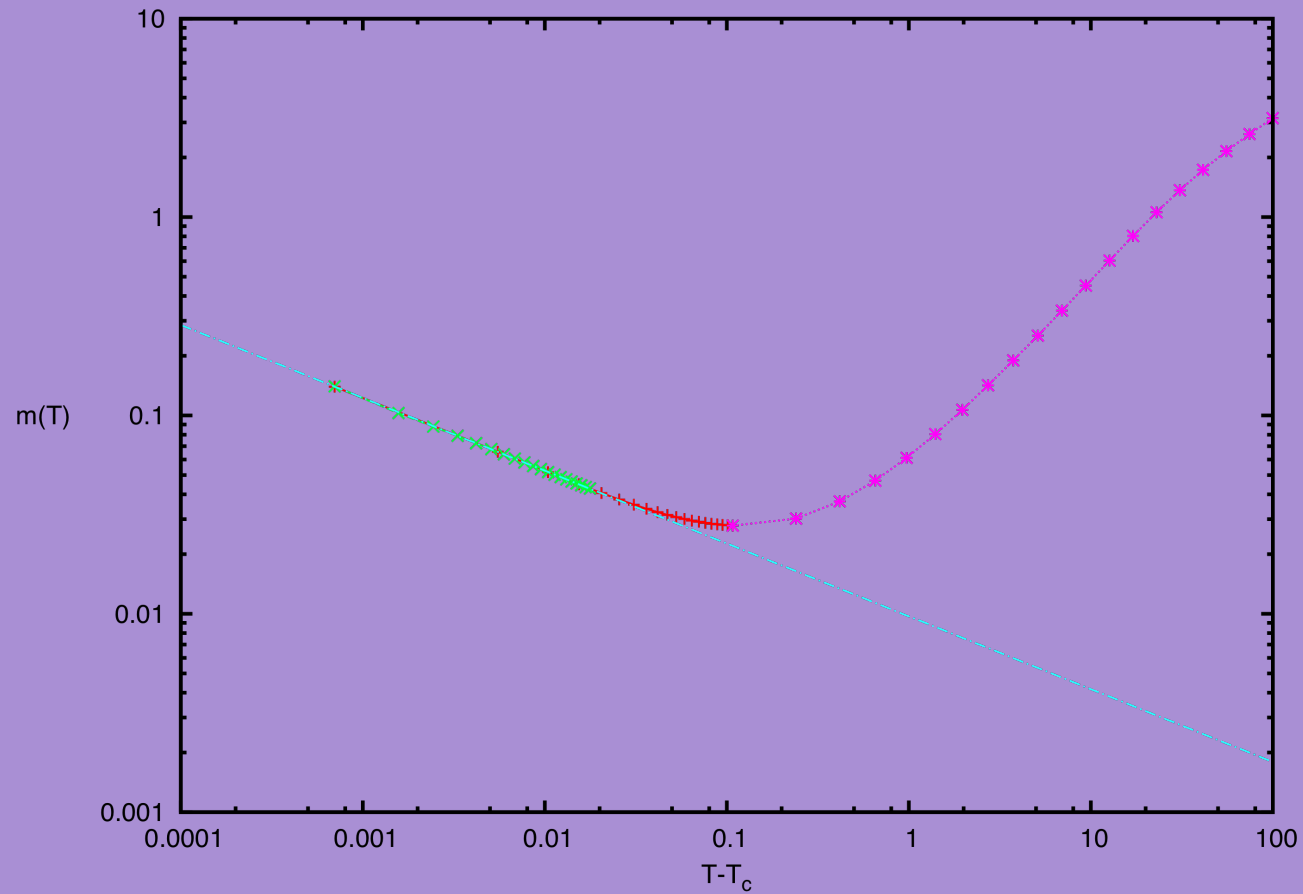
Effective gluon mass (IR-value of gluon energy)

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Critical mass

$$m(T) = \omega(k = \lambda_{IR}, T)$$



Critical exponent

critical behaviour of effective gluon mass near T_C :

$$m_{\text{crit}}(T) \sim (t-1)^{-\kappa} \quad t = T / T_C \quad \kappa \simeq 0.37$$

Critical exponent

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$$m_{crit}(T) \sim (t-1)^{-\kappa} \quad t = T / T_C \quad \kappa \simeq 0.37$$

quasi-gluon picture (Castorina, Miller, Satz 2011):

$$\kappa = (1 + \eta - \nu)$$

ν – critical exponent of the correlation length $\xi(T) \sim (t-1)^{-\nu}$

η – anomalous dimension exponent

$$\text{SU}(2) \hat{=} \text{d=3 Ising model: } \nu \simeq 0.63, \eta \simeq 0.04, \quad \kappa \simeq 0.41$$

Hamiltonian approach to QCD in Coulomb gauge $\partial A = 0$

$$H = H_{YM} + H_C + H_q$$

Christ and Lee

$$H_{YM} = \frac{1}{2} \int (J^{-1} \Pi^\perp J \Pi^\perp + B^2)$$

$$J(A^\perp) = \text{Det}(-D\partial) \quad D = \partial + gA$$

$$H_C = \frac{1}{2} \int J^{-1} \rho (-D\partial)^{-1} (-\partial^2) (-D\partial)^{-1} J \rho$$

color charge density: $\rho^a = (-\hat{A}^\perp \Pi^\perp)^a + \Psi^\dagger(\mathbf{x}) t^a \Psi(\mathbf{x})$

$$H_q = \int \Psi^\dagger(\mathbf{x}) [\alpha(\mathbf{p} + gA) + \beta m_0] \Psi(\mathbf{x})$$

Quark wave functional

M. Pak & H. R. 2011

quark w.f. (Thoules theorem)

$$\langle A | \Phi \rangle_q = \exp \left[\int \Psi^\dagger (\mathbf{s}\beta + \mathbf{v}\vec{\alpha} \cdot \vec{A}) \Psi \right] | 0 \rangle$$

$\mathbf{v} = 0$: *BCS – wave function (Adler & Davis)*

$\mathbf{v} \neq 0$: *quark – gluon coupling*

Quark condensate

without quark – gluon coupling ($v = 0$)

$$\langle \bar{q}q \rangle = (-113 \text{ MeV} \sqrt{\sigma_C / \sigma_W})^3$$

with quark – gluon coupling

$$\langle \bar{q}q \rangle = (-135 \text{ MeV} \sqrt{\sigma_C / \sigma_W})^3 \quad 20\% \text{ – increase}$$

lattice: $\sigma_C / \sigma_W = 2 \dots 3$

$$\langle \bar{q}q \rangle = (-191 \dots 234 \text{ MeV})^3$$

exp: $\langle \bar{q}q \rangle = (-230 \text{ MeV})^3$

Constituent quark mass

without quark – gluon coupling ($v = 0$)

$$m = 84 \text{ MeV} \sqrt{\sigma_C / \sigma_W}$$

with quark – gluon coupling

$$m = 132 \text{ MeV} \sqrt{\sigma_C / \sigma_W} \quad 57\% \text{ – increase}$$

lattice: $\sigma_C / \sigma_W = 2 \dots 3$

$$m = 187 \dots 230 \text{ MeV}$$

Hamiltonian approach to QCD in Coulomb gauge

- variational approach to YMT at $T=0$
 - decent description of IR sector
$$\omega(k) \sim k^{-\alpha} \quad d(k) \sim k^{-\beta} \quad \alpha = 2\beta - 1 \quad \beta = 1 \quad \beta = 0.8$$
 - dual superconductor
- FRG flow equation
 - similar results as in variational approach
- YMT at finite T
 - deconfinement phase transition $\beta = 1/2$
 - critical temperature $T_c = 230 \dots 280 \text{ MeV}$
- QCD
 - spontaneous breaking of chiral symmetry
 - quark condensate: $\langle \bar{q}q \rangle = (-191 \dots 234 \text{ MeV})^3$

Thanks for your attention