Strange quarks and the detection of dark matter



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International School of Nuclear Physics 33rd Course *From Quarks and Gluons to Hadrons and Nuclei* Erice, Sicily September 16-24, 2011





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The direct search for dark matter

 "Dark Matter Results from 100 Live Days of XENON100 Data" arXiv:1104.2549

The New York Eimes



Today... convince you this CMSSM blob is too "high"

CMSSM: Constrained Minimal Supersymmetric Standard Model

WIMP-Nucleon cross section

• The Constrained MSSM (CMSSM):

Eur. Phys. J. C (2011) 71: 1634 DOI 10.1140/epjc/s10052-011-1634-1 THE EUROPEAN PHYSICAL JOURNAL C

Regular Article - Theoretical Physics

Implications of initial LHC searches for supersymmetry

O. Buchmueller¹, R. Cavanaugh^{2,3}, D. Colling¹, A. De Roeck^{4,5}, M.J. Dolan⁶, J.R. Ellis^{4,7}, H. Flächer⁸, S. Heinemeyer⁹, G. Isidori¹⁰, K. Olive^{11,a}, S. Rogerson¹, F. Ronga¹², G. Weiglein¹³



Spin-independent neutralino cross section

see eg. Ellis, Olive & Savage, PRD77:065026(2008)

Scalar neutralino-quark contact interaction

q=u.d.s

 $\mathcal{L}_{SI} = \sum_{i} \alpha_{3i} \bar{\chi} \chi \bar{q}_{i} q_{i}$ depend on model (eg. CMSSM) α_{3i} we had down to low-energy sca evolved down to low-energy scale • Cross section $\sigma_{SI}^p \propto |f_p|^2$ $\frac{f_p}{M_p} = \sum_{q=u,d,s} \bar{\sigma}_{pq} \frac{\alpha_{3q}}{m_q} + \frac{2}{27} f_{TG}^p \sum_{q=c,b,t} \frac{\alpha_{3q}}{m_q}$ Nucleon scalar quark content $\bar{\sigma}_{pq} = \frac{m_q}{M_N} \langle N | \bar{q}q | N \rangle$ $\Sigma_{\pi N} = M_N(\bar{\sigma}_{pu} + \bar{\sigma}_{pd}) = \begin{cases} 45 \pm 8 \,\text{MeV} & \text{Gasser et al. (1991)} \\ 64 \pm 7 \,\text{MeV} & \text{GWU (2002)} \end{cases}$ $f_{TG}^p = 1 - \sum \bar{\sigma}_{pq}$ Trace anomaly:

Shifman, Vainstein & Zakharov, PLB(1978)

The missing ingredient

- Strangeness scalar content $\ ar{\sigma}_{ps} = m_s \langle N | ar{s}s | N
 angle / M_N$
- Commonly used quantity

$$\sigma_0 \equiv \hat{m} \langle N | \bar{u}u + \bar{d}d - 2\bar{s}s | N \rangle$$

some algebra

$$\Rightarrow \bar{\sigma}_{ps} = \frac{m_s}{2\hat{m}} \left(\Sigma_{\pi N} - \sigma_0 \right) / M_N$$

- Use Feynman-Hellmann relation $m_q \langle N | \bar{q}q | N \rangle = m_q \frac{\partial M_N}{\partial m_q}$
- First-order breaking in SU(3) baryon masses

$$\sigma_0 \simeq \hat{m} \frac{m_{\Xi} + m_{\Sigma} - 2m_N}{m_s - \hat{m}} = 26 \,\mathrm{MeV}$$

• With higher-order terms in chiral expansion

$$\sigma_0 \simeq 36 \pm 7 \,\mathrm{MeV} \qquad \Rightarrow \sigma_{ps} = \begin{cases} 110 \pm 130 \,\mathrm{MeV} & [\Sigma_{\pi N}(1)] \\ 350 \pm 120 \,\mathrm{MeV} & [\Sigma_{\pi N}(2)] \end{cases}$$

Lattice QCD can probe scalar matrix elements

- Can directly probe matrix elements in 3-pt correlation functions
 - See talks:
 - Alexandrou, Bali
- Here: Parameterise quark mass dependence and use Feyman–Hellmann relation

$$m_q \langle N | \bar{q}q | N \rangle = m_q \frac{\partial M_N}{\partial m_q}$$

2+1-flavour lattice results

Dynamical : $m_u = m_d \& m_s$

Octet baryon masses



 State-of-the-art lattice results approaching the physical domain

Chiral EFT: SU(3) expansion to $m_q^{3/2}$

- Chiral EFT is low-energy effective theory of QCD
- Only way to perform chiral extrapolation consistent with the chiral symmetries and symmetry breaking of QCD

Octet baryon masses

- 4 free parameters (at this order)
 - 1 Overall mass scale
 - 3 Linear perturbation in quark masses

Chiral nonanalytic contributions come with *model-independent* coefficients

$$\alpha_{M}, \beta_{M}, \sigma_{M}$$

$$= \frac{\pi, K, \eta}{+}$$

 M_0

nputs:
$$g_A = 1.267, D \simeq \frac{3}{5}g_A, F \simeq \frac{2}{5}g_A, C \simeq -2D, f_\pi \simeq 0.087 \,\text{GeV}$$

 $\pm 15\% \pm 15\% \pm 15\% \pm 15\% \pm 5\%$

Corrections to the linear expansion

 $\frac{2}{\pi}\int dk \frac{k^4}{k^2+m^2} \stackrel{R}{\longrightarrow} m^3$

Poorly converging



Finite Range Regularization (FRR)

- Suppress ultraviolet contributions to loop integrals from scale beyond the validity of the EFT
- Maintain renormalization such that scale dependence is removed to working order

Text book
$$\frac{2}{\pi} \int dk \frac{k^4}{k^2 + m^2} \xrightarrow{R} m^3$$
FRR
$$\frac{2}{\pi} \int dk \frac{k^4}{k^2 + m^2} \theta(\Lambda^2 - k^2) \xrightarrow{R} m^3 \frac{2}{\pi} \arctan \frac{\Lambda}{m}$$

Donoghue, Holstein & Borasoy, PRD59,036002(1999) Leinweber *et al.*, PRD61,074502(2000) Young, Leinweber & Thomas, PPNP 50,399(2003) Leinweber, Thomas & Young, PRL92,242002(2004)

Lattice results "choose" regularisation scale

• Lattice results prefer a regularisation scale of order 1 GeV (Dipole)



New development: preferred scale is **not** input from phenomenology

Young & Thomas, PRD(2010)

Fit to 8 LHPC points



Young & Thomas, PRD(2010)

Fit to 8 PACS-CS points

PACS-CS: 2+1-flavour simulation; different action discretization to LHPC



PACS-CS have an additional run with a different strange quark mass

Young & Thomas, PRD(2010)

Strange-quark mass dependence



Sigma terms from lattice QCD



πN Sigma Term (Expt): GL: Gasser & Leutwyler (1991) GW: Pavan et al. (2001)

Octet Masses & Breaking:

Gasser (1981) NK: Nelson & Kaplan (1987) BM: Borasoy & Meissner (1997)

3-flavour Lattice QCD: YT: Young & Thomas (2009) TF: Toussaint & Freeman (2009)

Updated cross sections for benchmark models



Ellis, Olive & Savage PRD(2008)

Strong dependence on sigma term from poorly known strangeness

$$\bar{\sigma}_{ps} = \frac{m_s}{2\hat{m}} \left(\Sigma_{\pi N} - \sigma_0 \right)$$

Giedt, Thomas & Young, PRL(2009)

Significant reduction in uncertainty

Cross-section reduced by order of magnitude from XENON100 figure

XENON100

• Shift the "blob" down



QCDSF *lattice results*



• "Prediction": curves are NOT a fit

Lattice QCD determination

	N	Λ	\sum	[I]
$\bar{\sigma}_{Bl}$	0.050(9)(1)(3)	0.028(4)(1)(2)	0.0212(27)(1)(17)	0.0100(10)(0)(4)
$\bar{\sigma}_{Bs}$	0.033(16)(4)(2)	0.144(15)(10)(2)	0.187(15)(3)(4)	0.244(15)(12)(2)



Young & Thomas, PRD(2010)

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> Strange scalar content is small