



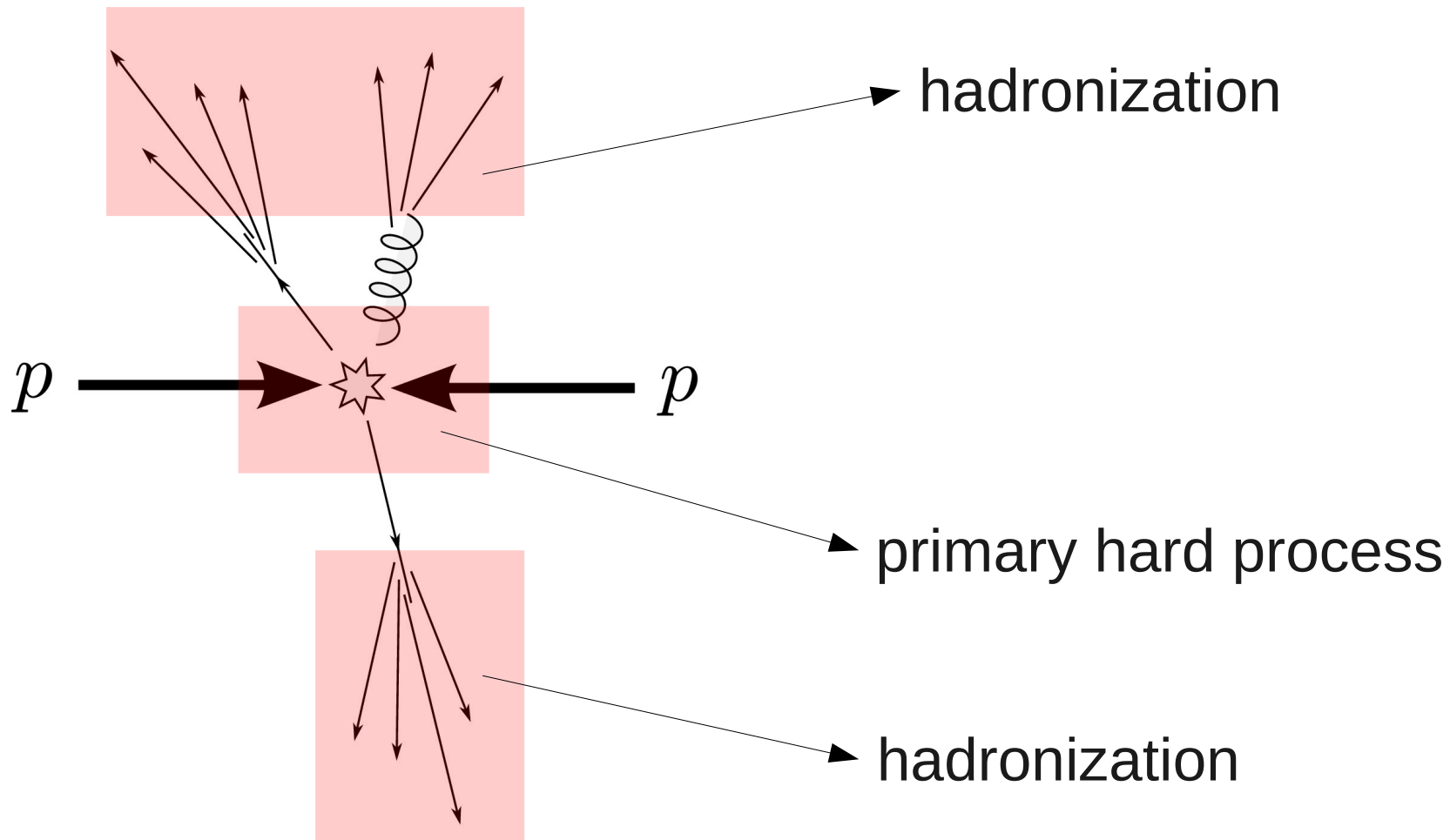
**“From Quarks and Gluons to Hadrons and Nuclei”  
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# Hadron Production in Jets at LHC

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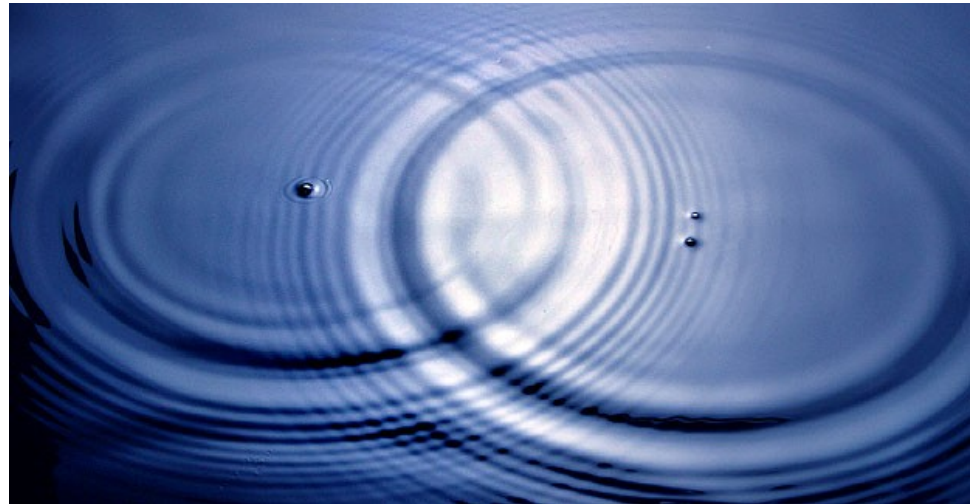
# A hard process at LHC



# Waves and obstacles: size is important

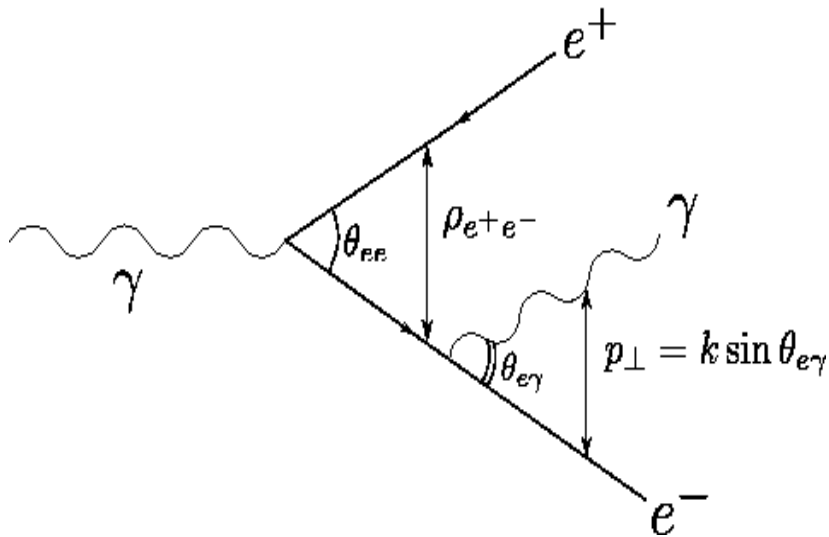


Wavelength is smaller or comparable with the size of an obstacle:  
**obstacle shape plays important role in dynamics**



Wavelength is larger than the size of an obstacle:  
**obstacle shape is irrelevant**

# QED: the essence of coherence



- To which extent  $e^+$  and  $e^-$  independently emit gamma's?

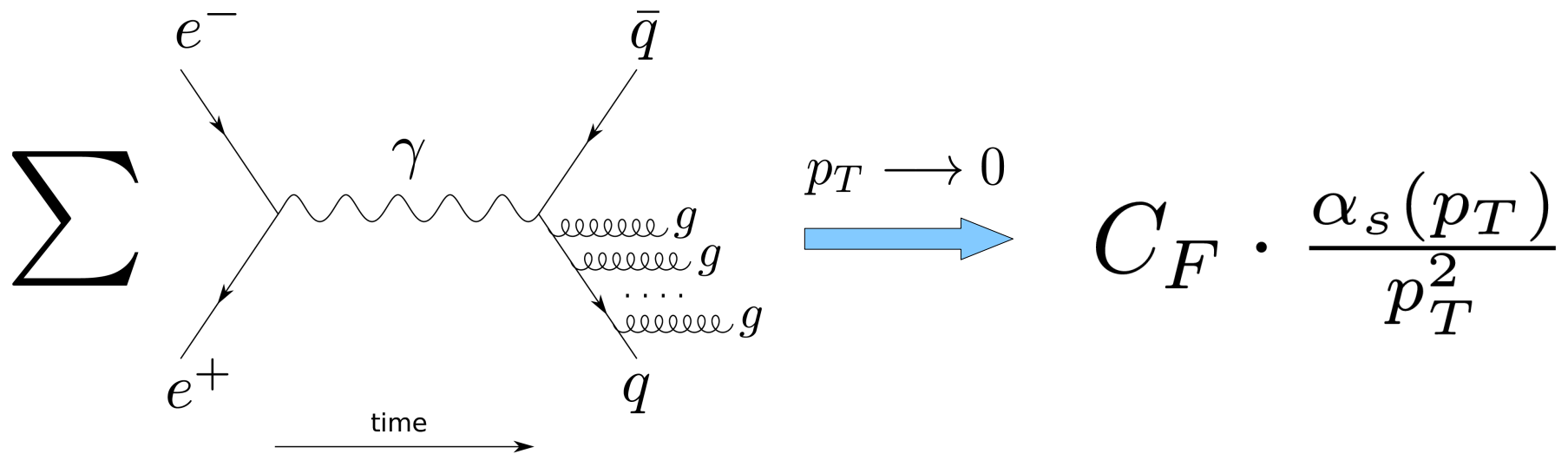
- If the wavelength of the photon is larger than the transverse separation of  $e^+ e^-$ , it cannot resolve the internal structure of the pair and probes only the electric charge, i.e. is effectively emitted by the chargeless object  $\rightarrow$  emission suppressed

$$\lambda_{\perp\gamma} = \frac{1}{p_{\perp}} = \frac{1}{k\Theta_{e\gamma}} ; \quad \rho_{e^+e^-} = \lambda_{\perp\gamma} \frac{\Theta_{ee}}{\Theta_{e\gamma}}$$

$$\rho_{e^+e^-} \ll \lambda_{\perp\gamma}, \quad \text{when} \quad \frac{\Theta_{ee}}{\Theta_{e\gamma}} \ll 1$$

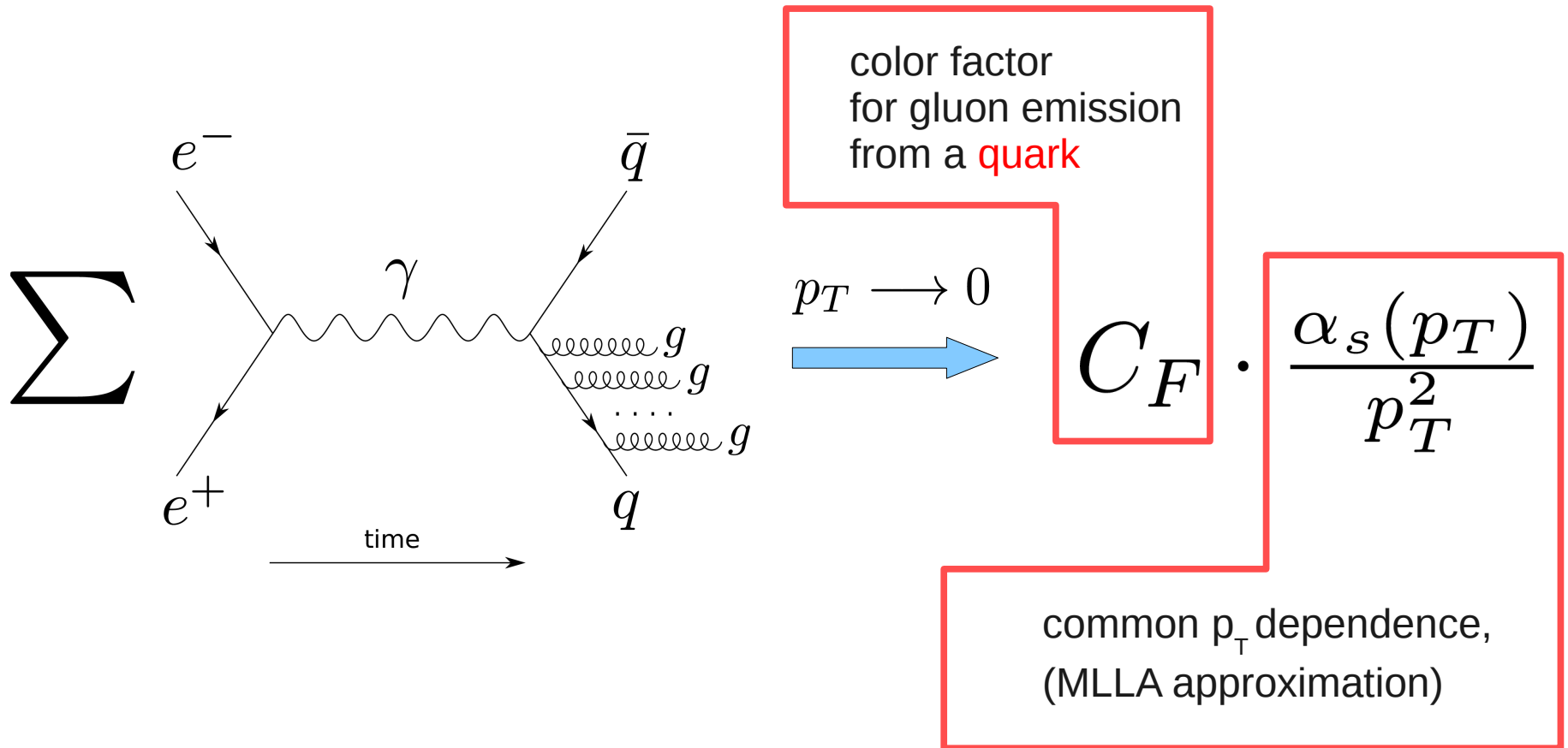
The emission at large angles is suppressed (Chudakov effect)

# QCD: soft hadron production in $e^+e^-$



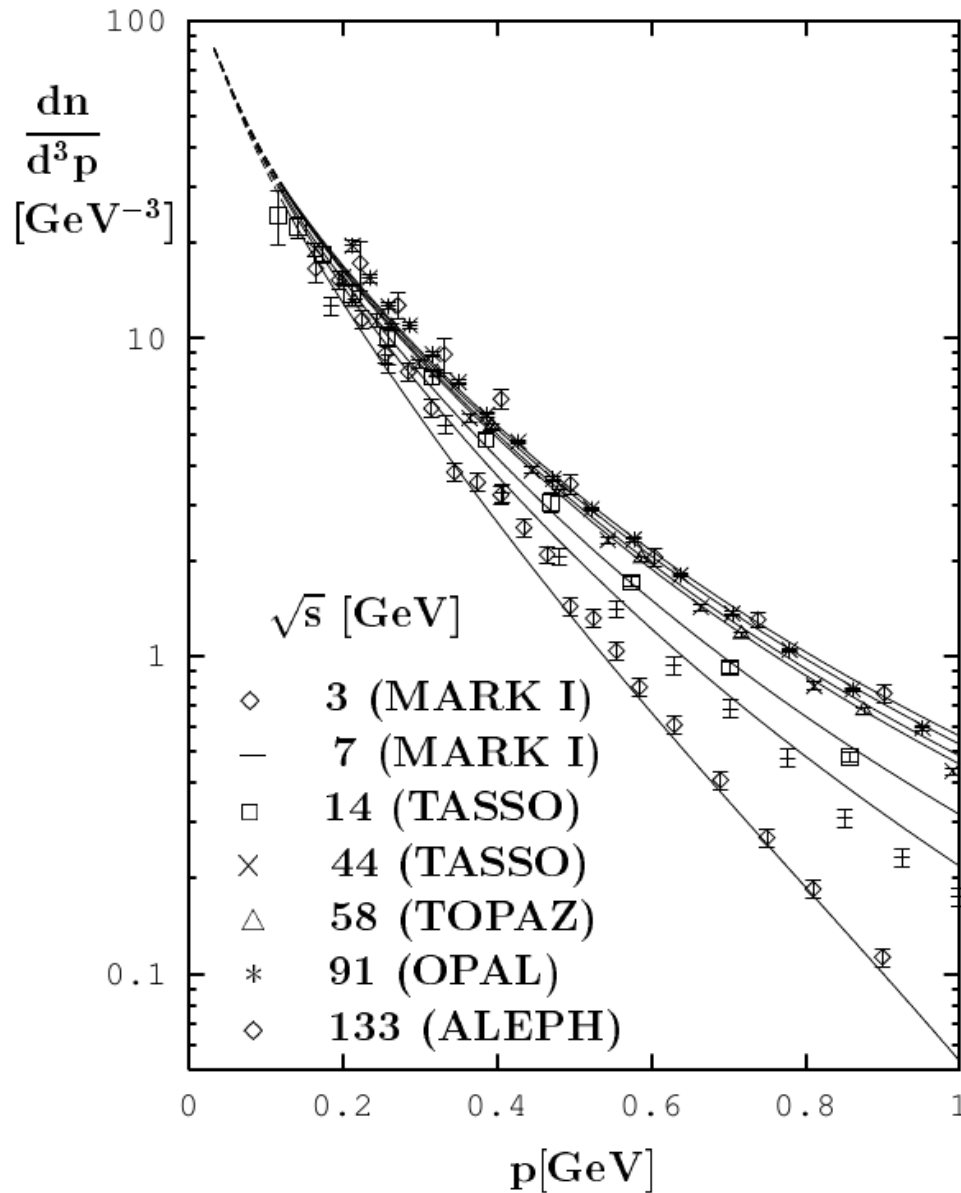
Soft gluons are emitted coherently. Due to their large wavelength soft gluons do not resolve the structure of  $q\bar{q}$  pair and see only the total color charge.

# QCD: soft hadron production in $e^+e^-$



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# Soft hadron production in $e^+e^-$



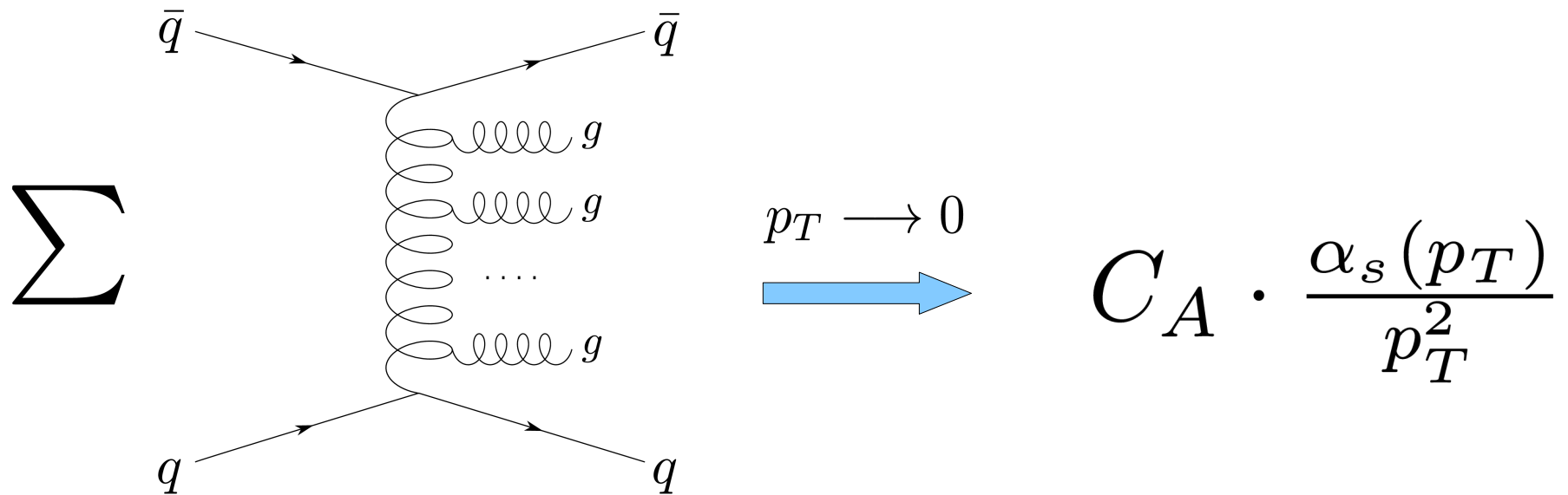
$n$  – inclusive particle densities of charged particles in  $e^+e^-$  annihilations for different collision energies

$$n \sim C_F \cdot \frac{\alpha_s(p_T)}{p_T^2}$$

$n$  does not depend on the collision energy in the limit  $p_T \rightarrow 0$

This prediction is consistent with experimental data

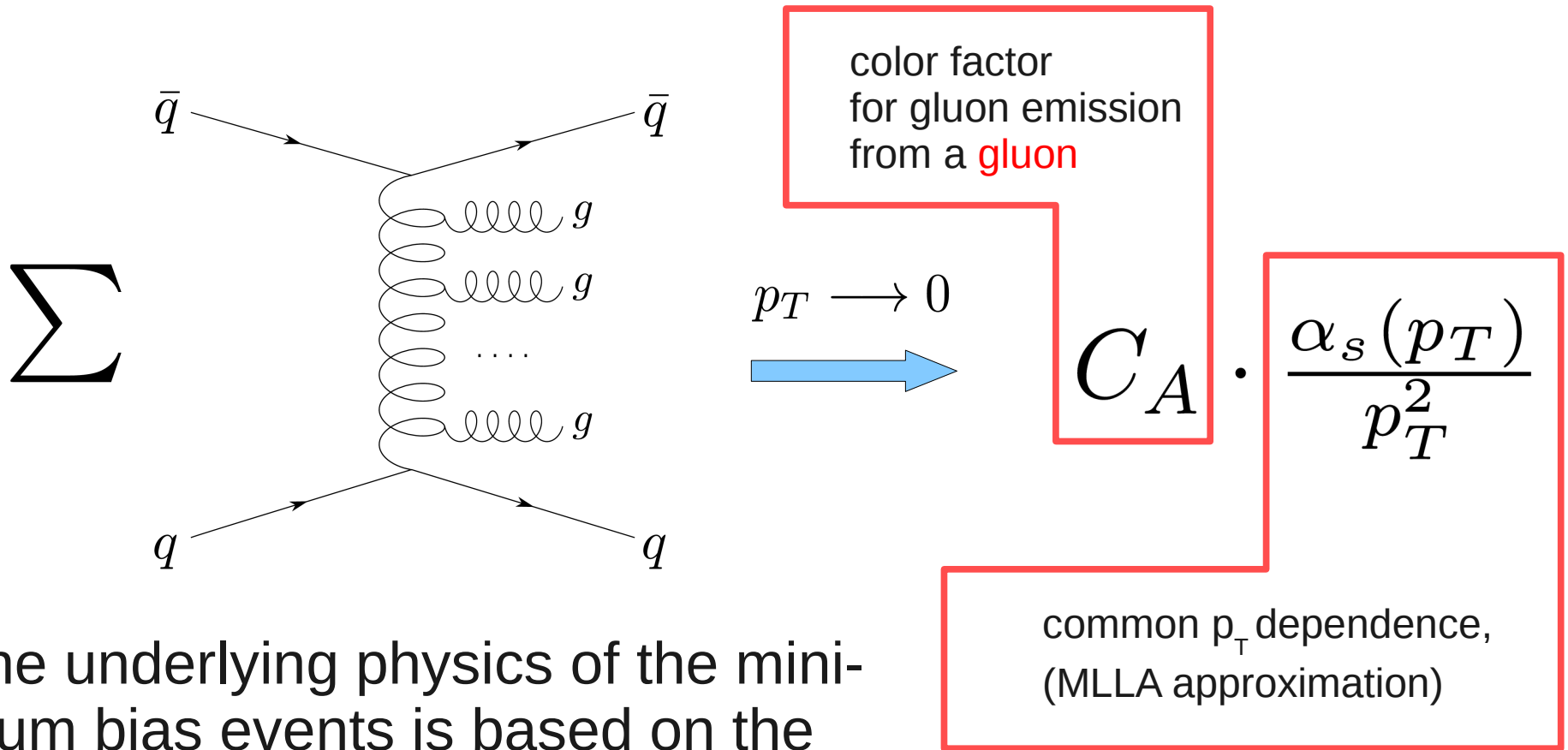
# Soft hadron production in pp



The underlying physics of the minimum bias events is based on the collisions of two partons within the protons. The exchange of the t-channel gluon rearranges the incoming colors and leads to the radiation of the soft gluons from the effective color **octet** dipole.



# Soft hadron production in pp



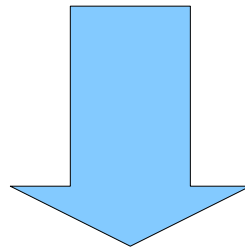
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# More on soft hadron production in $e^+e^-$ and $pp$ collisions

In the  $p_T \rightarrow 0$  limit:

$$I_0^{pp} = E \frac{dn^{pp}}{d^3p} = \text{const} \sim C_A$$

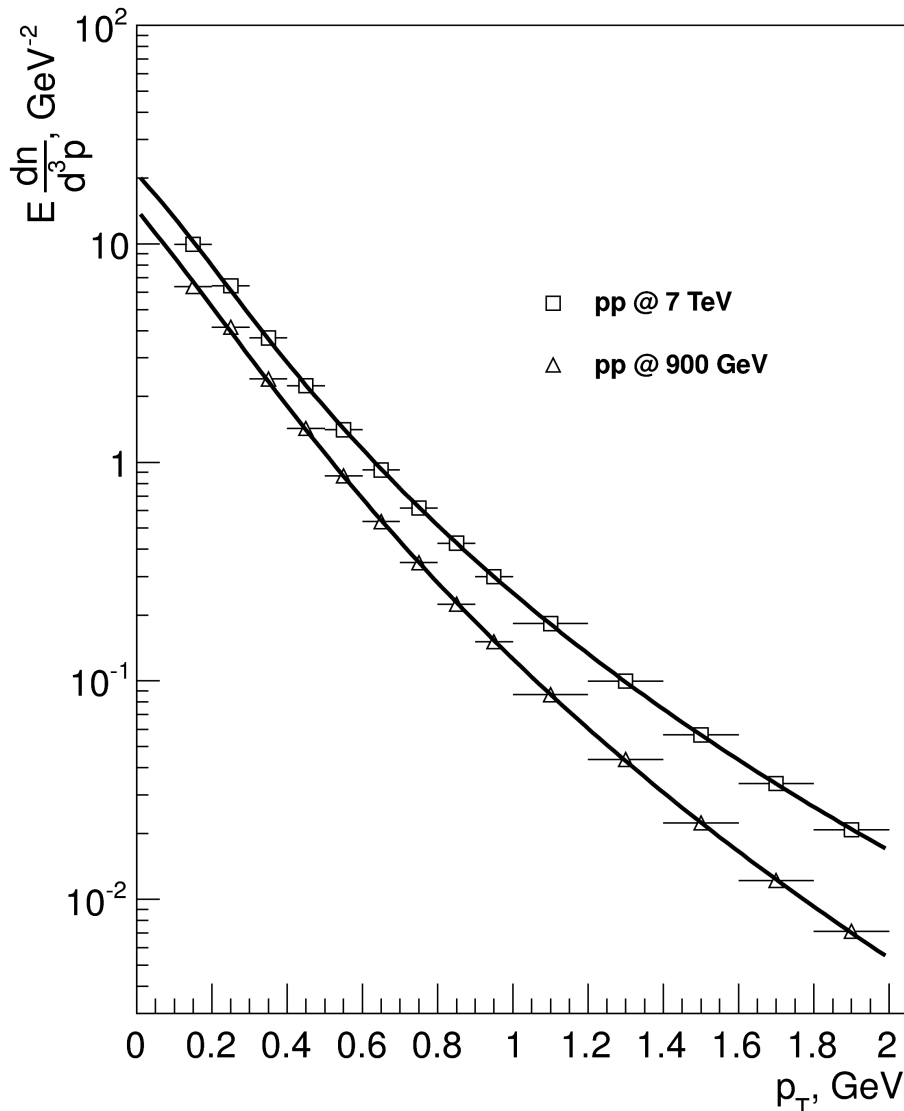
$$I_0^{ee} = E \frac{dn^{ee}}{d^3p} = \text{const} \sim C_F$$



$$\frac{I_0^{pp}}{I_0^{ee}} = \frac{C_A}{C_F} = \frac{9}{4}$$

This prediction can be tested experimentally

# CMS charged particle multiplicity



Fit extrapolation to  $p_T=0$  point shows  $\sim 30\%$  difference.

Errors for the extrapolation to  $p_T=0$  are not shown.

ALICE experiment is designed for low  $p_T$  measurements and may provide more precise results.

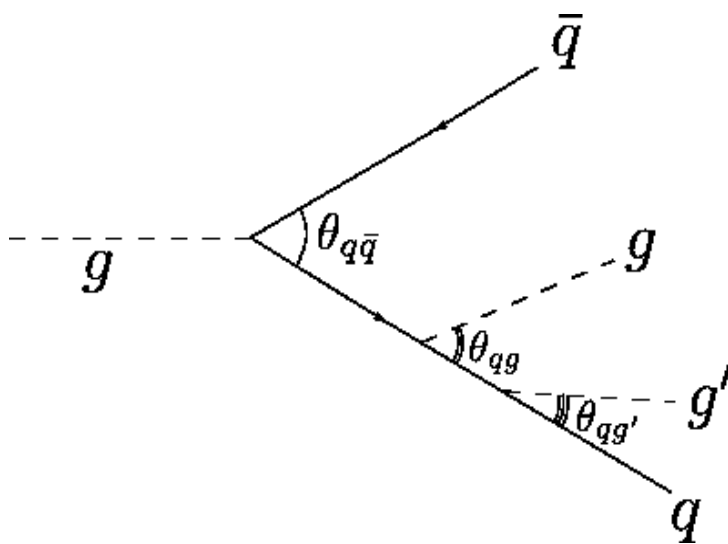
# Summary

- Inclusive soft particle spectra is independent of collision energy.
- The relative normalization of spectra in different processes is given by the color factors relevant for the minimal partonic process.
- Is there any new incoherent contributions in the new energy regime at LHC?

# Backup

# QCD: color coherence

- The same effect takes place in QCD where soft gluon radiation is governed by conserved color currents.



- **On the one hand**, the wavelength of the emitted gluon should be smaller than the hadronization scale  $R$ :

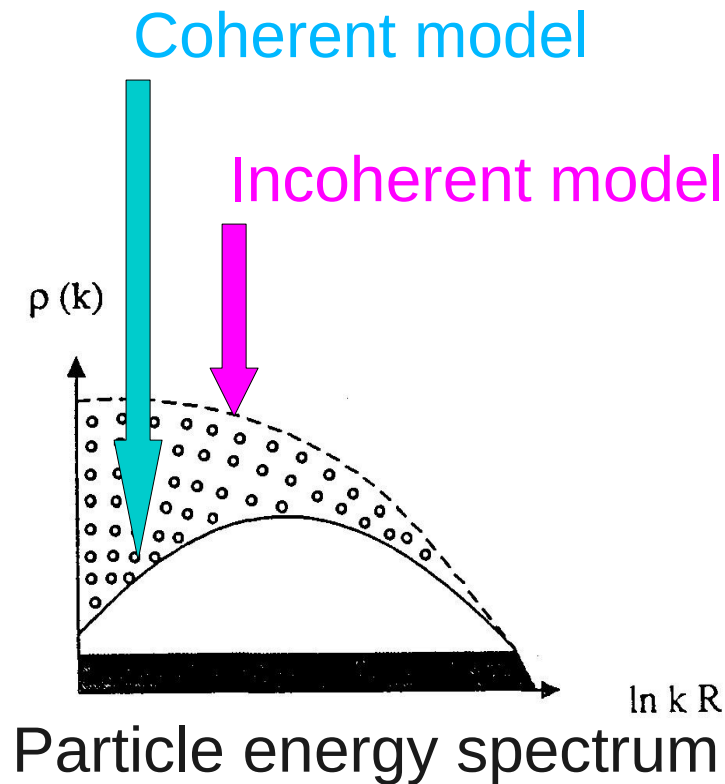
$$\lambda_{\perp g} = \frac{1}{k \sin \Theta_{qg}} > R \Rightarrow \Theta_{qg} > \frac{1}{kR}$$

- **On the other hand**, we have angular ordering:

$$\Theta_{qq} \ll \Theta_{qg} \ll \Theta_{qg'}$$

Yu.L. Dokshitzer, V.A. Khoze, A.H. Mueller, and S.I. Troyan. *Rev. Mod. Phys.*, 60:373, 1988.

# What can we observe with color coherence?



$$\rho(k) \equiv \frac{dn}{d \ln k}$$

R – hadronization scale,  
k – particle momentum

- Let us illustrate the influence of color coherence on particle spectra on a toy model
- **The suppression of soft radiation** follows from the angular ordering of partonic cascade and is a direct manifestation of the color coherence.
- This can be understood on kinematics ground **as a result of two conflicting tendencies**: due to the hadronization a slow particle is 'forced out' at large emission angle, on the contrary, the allowed decaying angle, after a few successive branching, is shrunk to small values.