

# Muon $g - 2$ and the pion form factor in lattice QCD

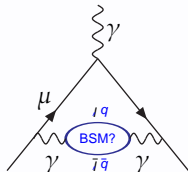
International School of Nuclear Physics  
33rd course  
From Quarks and Gluons to Hadrons and Nuclei  
Erice-Sicily  
September 16-24, 2011

Andreas Jüttner  
CERN Theory Division

# Introduction: elm. pion FF and $a_\mu^{\text{LO, had}}$

Some cartoons:

$\mu$  LO<sup>elm</sup> vacuum polarisation

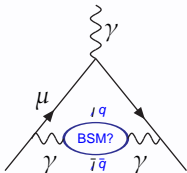


$$a_\mu^{\text{LHV}} \propto \int_0^\infty dQ^2 K(Q^2) (\Pi(Q^2) - \Pi(0))$$

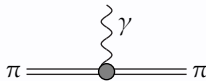
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pion form factor



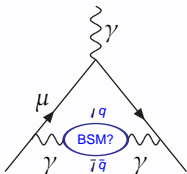
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$$f^{\pi\pi}(q^2)(p' + p) = \langle \pi(p') | j^{\text{elm}} | \pi(p) \rangle$$

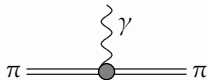
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non-perturbative QCD-contributions to both processes

- talks by Tobias Goecke and Christian Fischer for analytical approaches
- this talk: lattice QCD  
in principle a well defined problem in field Euclidean field theory  
in practice brute force neither feasible nor nice  
new ideas in the following
- Mainz group: Bastian Brandt, Michele Della Morte, Benjamin Jäger, A.J., Hartmut Wittig

# Motivation

- lattice QCD has seen continuous progress  
(e.g. FLAG summary for low energy physics results)  
*Eur.Phys.J. C71 (2011) 1695, arXiv:1011.4408*
- a number of yet unsolved obstacles, some of which affect  $a_\mu^{\text{LHV}}$  and  $f_{\pi\pi}$
- in particular for  $a_\mu^{\text{LHV}}$  the impact on particle physics potentially huge:  
 $(a_\mu^{\text{exp}} - a_\mu^{\text{th}}) \times 10^{-10} = 29(9)$  while  $a_\mu^{\text{LHV}} = 690(5)$   
*according to Jegerlehner and Nyffeler, Physics Reports 477(2009)1-110, arXiv:0902.3360*  
(of similar importance: light-by-light scattering but not covered here)

## What is ... ?

	QCD
$N_c$	3
$N_f$ , fundamental	1+1+1+1+1+1
$SU(2)$ iso-spin brk.	✓
$m_\pi$	135MeV
$V$	$\infty$
$a$	0

## What is ... ?

	QCD	Lattice QCD
$N_c$	3	3
$N_f$ , fundamental	1+1+1+1+1+1	0, 2, 2+1, 2+1+1
$SU(2)$ iso-spin brk.	✓	✗
$m_\pi$	135MeV	$\lesssim m_\pi^{\text{sim}}$
$V$	$\infty$	2-4fm
$a$	0	0.05-0.1fm

for  $a_\mu^{\text{LHV}}$  and  $f^{\pi\pi}(q^2)$  severe systematic effects due to

- finite volume
- unphysically heavy pion mass

## Difficulties (and solutions): finite volume

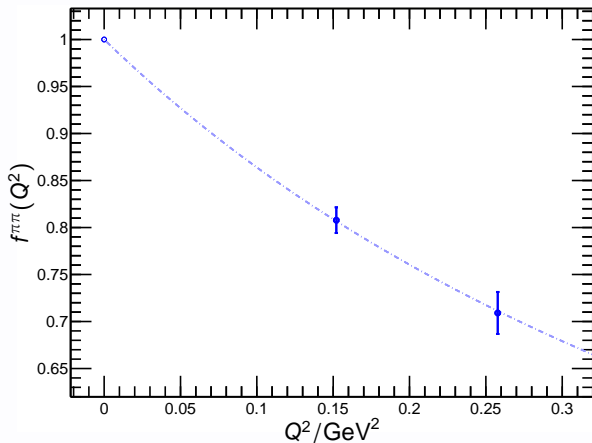
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→ hadron matrix elements can only be computed for these Fourier momenta:



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pion form factor in  $N_f = 2 + 1$  lattice QCD by RBC/UKQCD

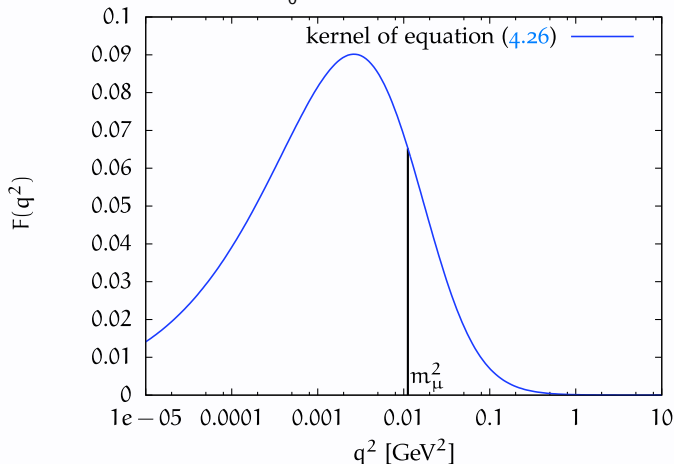


UKQCD JHEP 05(2007)016, JHEP 0807(2008)112

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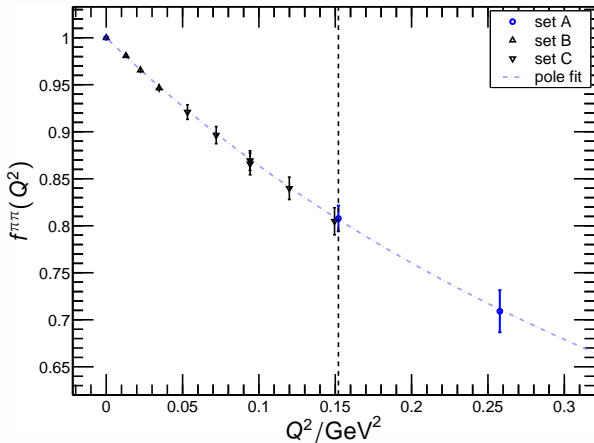
Tuning the fermion boundary conditions allows to induce arbitrary hadron momenta: (partially) twisted boundary conditions

*Bedaque PLB539(2004), Divitiis et al., PLB 595 (2004) 408, Bedaque, Chen, PLB 616:208-214,2005, Sachrajda, Viladoro, PLB 609:73-85,2005, ...*

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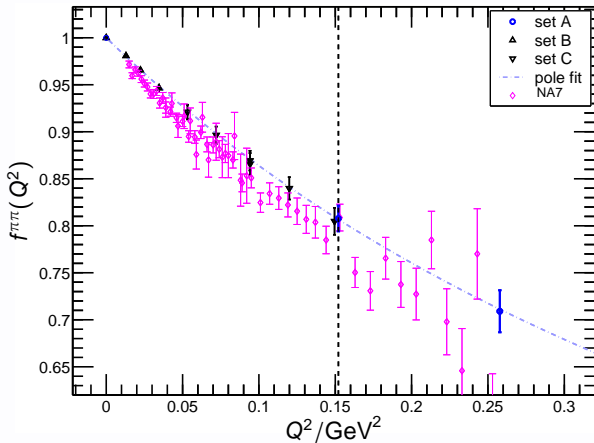


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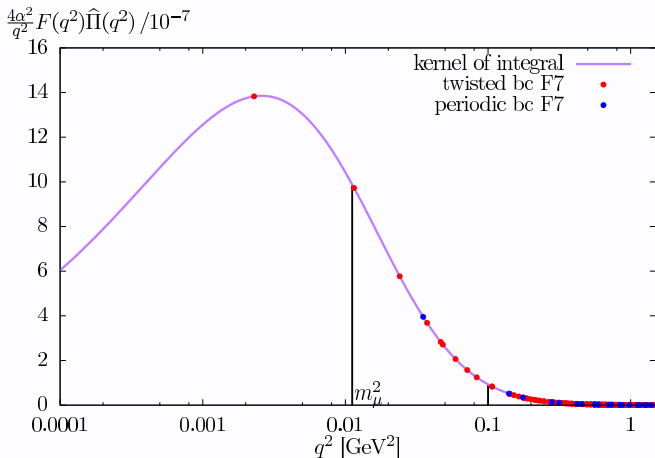


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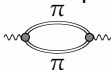


*CLS preliminary, plot by Benjamin Jäger, Mainz*

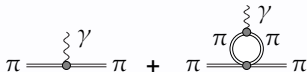
## Difficulties (and solutions): Quark mass dependence

- to date no single simulation of  $a_\mu^{\text{LHV}}$  and  $f^{\pi\pi}(q^2)$  with physical quark masses exists
- for making a *prediction* at the physical point in LQCD one resorts to extrapolations in the quark mass based on chiral perturbation theory

$a_\mu^{\text{LHV}}$



$f^{\pi\pi}(q^2)$



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## Difficulties (and solutions): disconnected diagrams

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- can be done, but numerically *expensive* and therefore neglected in all current simulations!!
- new technique allows their prediction in chiral perturbation theory; crude estimate at NLO: -10% shift in VP  $\Pi(q^2)$

*Della Morte & A.J. JHEP 1011 (2010) 154, arXiv:1009.3783*

## Summary of difficulties

- ✓ momentum resolution (twisted boundary conditions)
- ✗ mass dependence (extrapolation must be avoided)
- ? disconnected diagrams (they are small and we have a rough estimate)

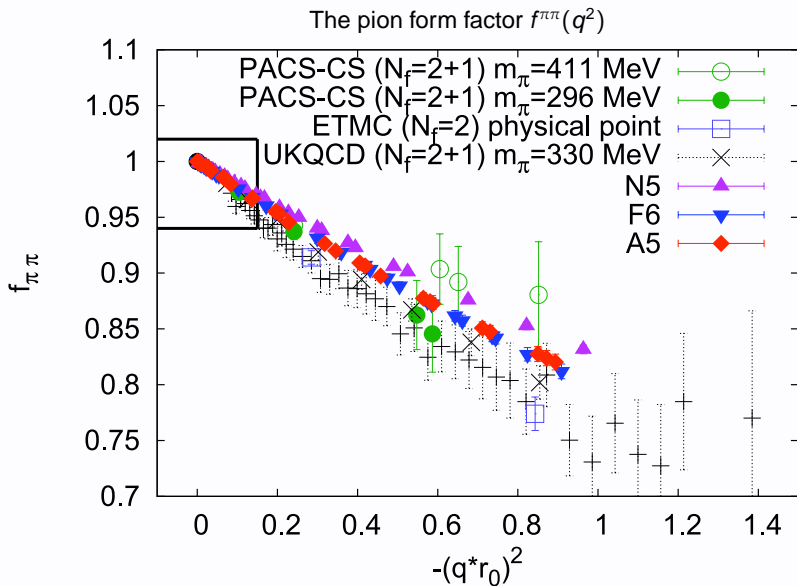
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Remainder: Status of ongoing calculations

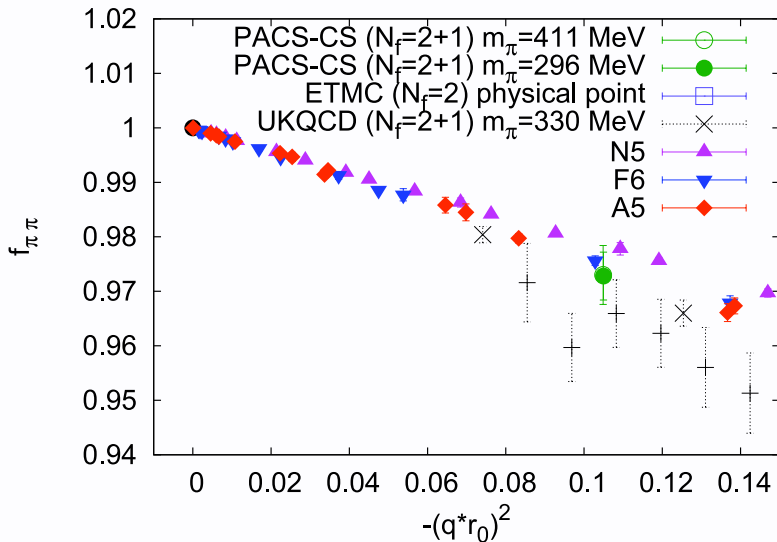


**Status:**  $f^{\pi\pi}(q^2)$



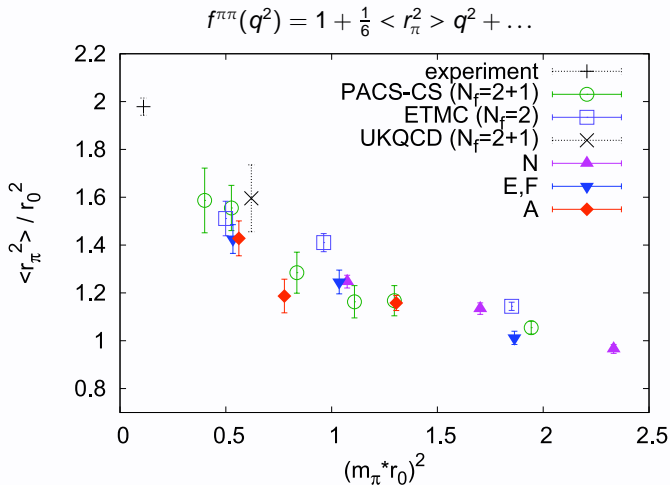
CLS preliminary, plot by Bastian Brandt, Mainz, arXiv:1109.0196

The pion form factor  $f_{\pi\pi}(q^2)$

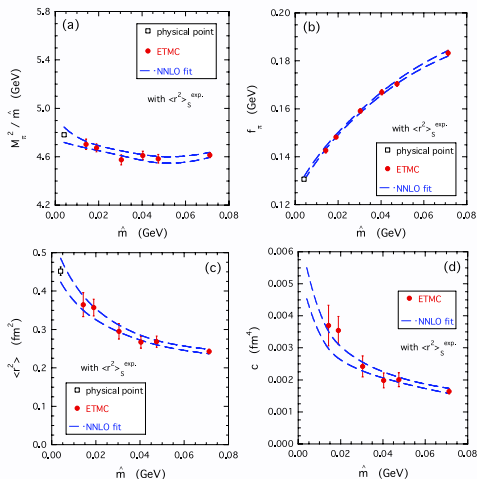


CLS preliminary, plot by Bastian Brandt, Mainz, arXiv:1109.0196

# Status: $f^{\pi\pi}(q^2)$



CLS preliminary, plot by Bastian Brandt, Mainz, arXiv:1109.0196



- all collaborations try to extrapolate the lattice data using  $\chi$ PT or some model

example here ETM [Phys.Rev. D79 \(2009\) 074506, arXiv:0812.4042](#)

# Summary $f^{\pi\pi}(q^2)$ - FLAG report

- G. Colangelo, S. Dürr, A. J., L. Lellouch, H. Leutwyler, V. Lubicz, S. Necco, C. Sachrajda, S. Simula, A. Vladikas, U. Wenger, H. Wittig  
[Eur.Phys.J. C71 \(2011\) 1695, arXiv:1011.4408](#)

→ quantities we consider:  $m_u, m_d, m_s, f_+(0), f_K/f_\pi, B_K, \text{NLO LEC's}$ , potentially more in the future

Collaboration	$N_f$	publication status	chiral extrapolation	continuum extrapolation	finite volume	$\langle r^2 \rangle_V^{\pi} [\text{fm}^2]$	$c_V (\text{GeV}^{-4})$	$\bar{\zeta}_6$
RBC/UKQCD 08A	2+1	A	●	■	★	0.418(31)	—	12.2(9)
LHP 04	2+1	A	●	■	●	0.310(46)	—	—
JLQCD/TWQCD 09	2	A	●	■	■	0.409(23)(37)	3.22(17)(36)	11.9(0.7)(1.0)
ETM 08	2	A	●	●	●	0.456(30)(24)	3.37(31)(27)	14.9(1.2)(0.7)
QCDSF/UKQCD 06A	2	A	★	●	●	0.441(19)(56)(29)	—	—
BCT 98						0.437(16)	3.85(60)	16.0(0.5)(0.7)
NA7 86						0.439(8)		
GL 84								16.5(1.1)

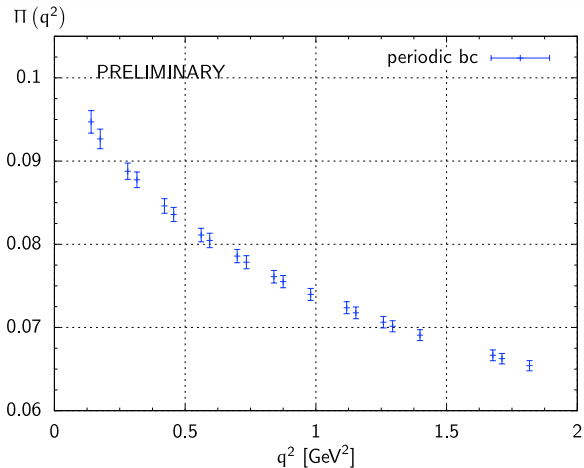
## Outlook $f^{\pi\pi}(q^2)$

- only twisted boundary conditions allow for model-independent extraction of the charge radius (derivative at  $q^2 = 0$ )
- a *final* word on the lattice on  $f^{\pi\pi}(q^2)$  musn't depend on  $\chi$ PT in the way current results do

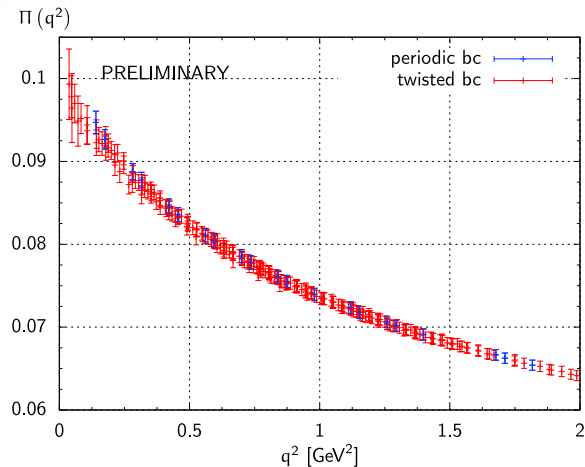




# Status $a_\mu^{\text{LHV}}$



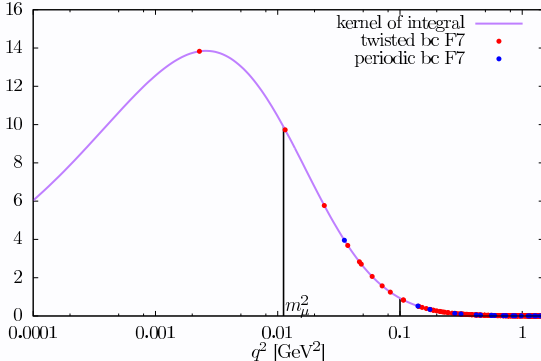
plots by Benjamin Jäger, Mainz, CLS-collaboration



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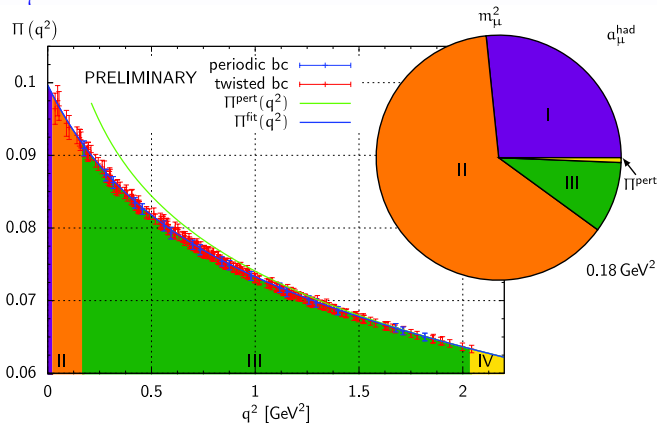
- twisting helps - many more data points in  $q^2$

$$\frac{4\alpha^2}{q^2} F(q^2) \hat{\Pi}(q^2) / 10^{-7}$$



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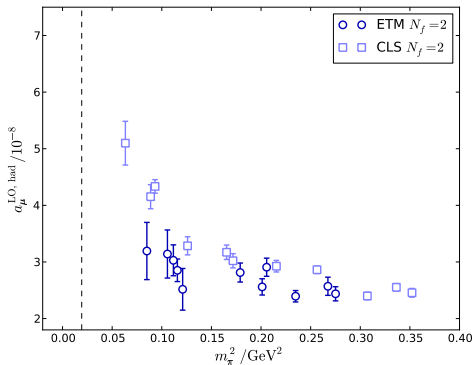


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$$a_\mu^{\text{LHV}} \propto \left(\frac{a}{\pi}\right)^2 \int_0^\infty dQ^2 \mathcal{K}(Q^2) (\Pi(Q^2) - \Pi(0))$$

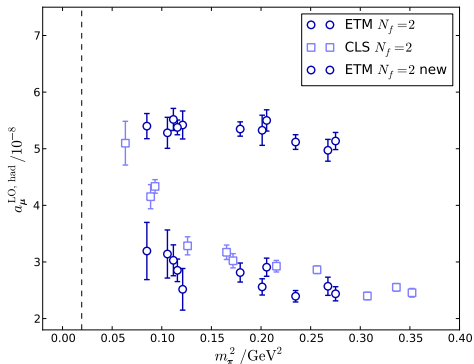
- twisting helps - many more data points in  $q^2$
- fit raw lattice data to various ansätze (pole, polynomial, Padé)
- match with PT (vary matching point)
- integrate: fold with QED PT

## $a_\mu^{\text{LHV}}$ in $N_f = 2$ flavour QCD



(thanks to Dru Renner for providing ETM's data)

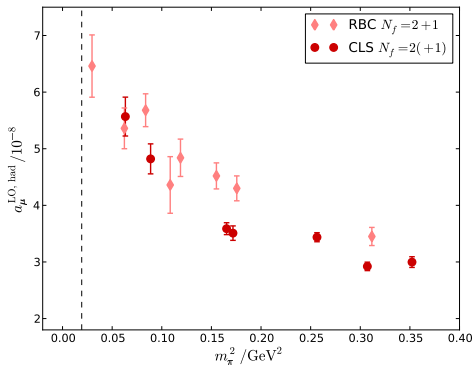
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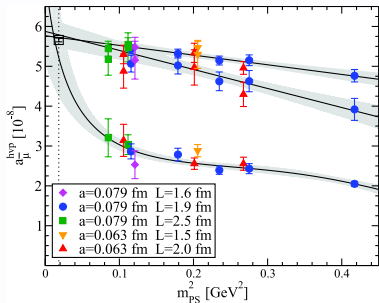
- $m_\pi \neq m_\pi^{\text{phys}}$  in all simulations (ETM divide out  $\rho$ -pole and they then extrapolate linearly in  $m_\pi^2$ , [arXiv:1103.4818](https://arxiv.org/abs/1103.4818))

## $a_\mu^{\text{LHV}}$ in $N_f = 2 + 1$ flavour QCD

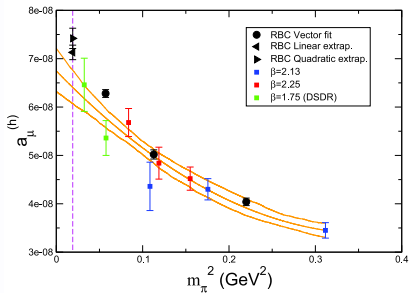


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## parameterising the mass dependence



ETM arXiv:1103.4818



UKQCD arXiv:1107.1497



# Numbers

Collaboration	$N_f$	$a_\mu^{\text{LHV}}/10^{-10}$
UKQCD <a href="https://arxiv.org/abs/1107.1497">arXiv:1107.1497</a>	2+1	641(33)(32)
ETM <a href="https://arxiv.org/abs/1103.4818">arXiv:1103.4818</a>	2	572(16)
exp. <i>e.g. Jegerlehner &amp; Nyffeler</i> <i>Physics Reports 477(2009)1-110</i>		690(5)

(for older efforts see: *Aubin and Blum, Phys.Rev. D75 (2007) 114502, hep-lat/0608011*  
*QCDSF, Nucl.Phys. B688 (2004) 135-164 hep-lat/0312032*)

## TODO:

- stabilise momentum dependence by using twisting
- compute contribution of disconnected diagrams
- simulate closer/at the physical pion
- iso-spin breaking effects?

# Summary

- recent summaries provide an impressive account of the advanced status of lattice QCD (*FLAG and Laiho, Lunghi, van de Water Phys.Rev. D81 (2010) 034503, arXiv:0910.2928*)
- efforts to predict  $f^{\pi\pi}(q^2)$  and  $a_{\mu}^{\text{LHV}}$  are still struggling to meet phenomenologists/experimentalist expectations/needs
- there has been progress:
  - twisted boundary conditions
  - estimates for the disconnected diagrams
  - simulations closer to the physical point
- for  $a_{\mu}^{\text{LHV}}$  ETM and RBC have recently published status reports and the Mainz group will soon present their results
- for  $f^{\pi\pi}(q^2)$  Mainz will also soon provide new results