

Chiral crossover effects on the shear viscosity of a pion gas



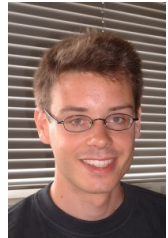
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Michael Buballa

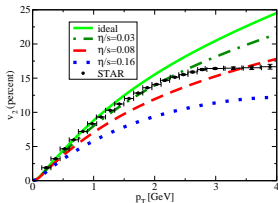
International School of Nuclear Physics
33rd Course
“From Quarks and Gluons to Hadrons and Nuclei”
Erice, Sicily, September 16 – 24, 2011

▶ based on

- ▶ **Klaus Heckmann**, Ph.D. thesis, TU Darmstadt, 2011
- ▶ K. Heckmann, M.B., J. Wambach, work in progress



► flow data at RHIC

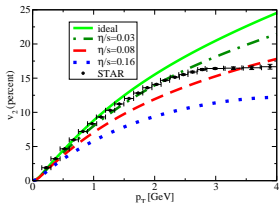


[P. & U. Romatschke, PRL (2007)]

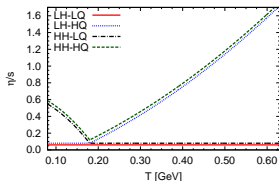
► conventional interpretation: QGP = “nearly perfect fluid” ($\eta/s \sim \frac{1}{4\pi}$)

Motivation

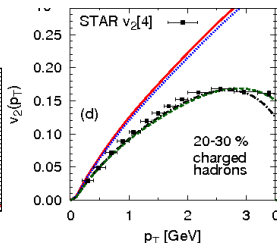
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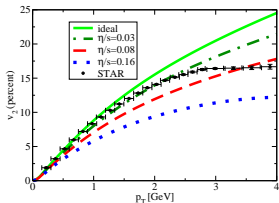


[H. Niemi et al., PRL (2011)]

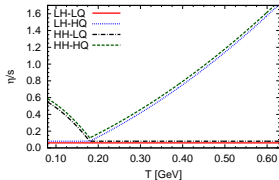


- conventional interpretation: QGP = “nearly perfect fluid” ($\eta/s \sim \frac{1}{4\pi}$)
- more recent: hydrodynamics more sensitive to hadronic phase

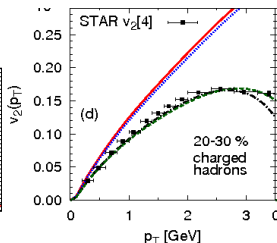
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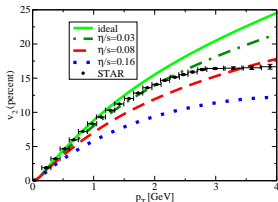
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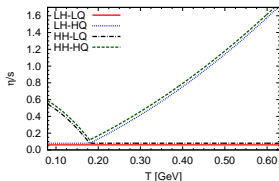
➔ **aim:** microscopic understanding of the shear viscosity in the hadronic phase

Motivation

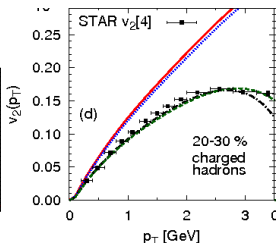
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▶ **conventional interpretation:** QGP = “nearly perfect fluid” ($\eta/s \sim \frac{1}{4\pi}$)

▶ **more recent:** hydrodynamics more sensitive to hadronic phase

➔ **aim:** microscopic understanding of the shear viscosity in the hadronic phase

▶ **here:** BUU approach to $\pi\pi$ -scattering in the NJL model

▶ correct low-temperature limit

▶ imprints of the chiral crossover and the compositeness of the pions



▶ basic ingredients and conservation laws:

▶ fluid 4-velocity $u^\mu(x)$, $u^\mu(x)u_\mu(x) = 1$

▶ energy-momentum tensor $T^{\mu\nu}(x)$,

$$\partial_\mu T^{\mu\nu}(x) = 0$$

▶ particle current $J^\mu(x) = n(x)u^\mu(x)$,

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▶ additional assumption:

EoS + local thermal equilibrium $\rightarrow \epsilon(x) = \epsilon(p(x))$

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▶ gradient expansion:

$T^{\mu\nu} = T^{(0)\mu\nu} + T^{(1)\mu\nu} + \dots$ (and similar for J^μ)

▶ ideal fluid:

$T^{(0)\mu\nu} = (\epsilon + p) u^\mu u^\nu - p g^{\mu\nu}$

▶ 1st-order viscous correction:

$T^{(1)\mu\nu} = \eta (\partial^\mu u^\nu + \partial^\nu u^\mu + u^\mu u^\lambda \partial_\lambda u^\nu + u^\nu u^\lambda \partial_\lambda u^\mu)$
 $+ (\zeta - \frac{2}{3}\eta) (g^{\mu\nu} - u^\mu u^\nu) \partial_\lambda u^\lambda$

Quantum relativistic kinetic theory



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- ▶ underlying assumption: mean free path $\lambda \gg$ interaction range r



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- ▶ particle phase-space distribution function: $f_a(\vec{x}, \vec{p}, t)$

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- ▶ Boltzmann-Uehling-Uhlenbeck (BUU) equation ($2 \rightarrow 2$)

$$\begin{aligned} \frac{d}{dt} f_a(\vec{x}, \vec{p}, t) = & \sum_b \frac{g_b}{1+\delta_{ab}} \int \frac{d^3p'}{(2\pi)^3} \int \frac{d^3p_1}{(2\pi)^3} \int \frac{d^3p_1'}{(2\pi)^3} \left\{ |\mathcal{M}_{ab}|^2 \frac{(2\pi)^4 \delta^4(p+p_1-p'-p_1')}{16E E_1 E' E_1'} \right. \\ & \left. \times [f'_a f'_{1b} (1+f_a)(1+f_{1b}) - f_a f_{1b} (1+f'_a)(1+f'_{1b})] \right\} \end{aligned}$$



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- ▶ linearization (2^{nd} -order Chapman-Enskog expansion)

$$f_a = f_a^{(0)} + f_a^{(1)} + \dots, \quad \text{local equilibrium: } f_a^{(0)}(x, p) = \frac{1}{\exp[(p^\mu u_\mu(x) - \mu_\pi(x))/T(x) - 1]}$$

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- ▶ linear integral equation for η



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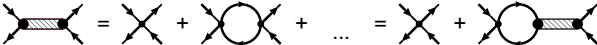
- ▶ linear integral equation for η

- ▶ physics input: scattering matrix element \mathcal{M}_{ab} here: $\pi\pi \rightarrow \pi\pi$ in NJL

Mesons in the NJL model

► Lagrangian: $\mathcal{L} = \bar{\psi}(i\partial - m)\psi + g [(\bar{\psi}\psi)^2 + (\bar{\psi}i\gamma_5\vec{\tau}\psi)^2]$

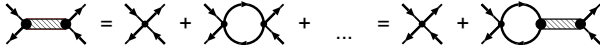
► gap equation:  \rightarrow dynamical quark masses

► mesons (RPA): 

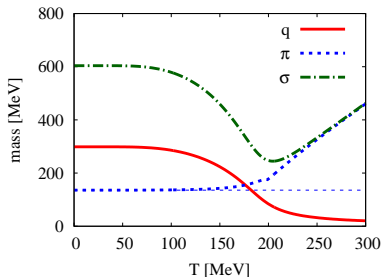
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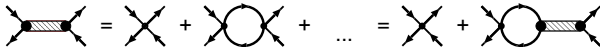
► in-medium masses:



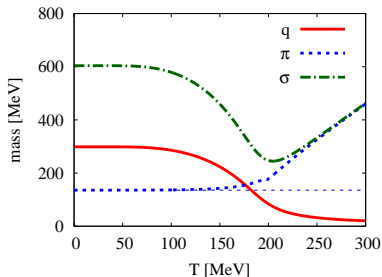
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▶ mesons (RPA): 

▶ in-medium masses:



▶ characteristic temperatures:

[Quack et al., PLB (1995)]

▶ σ -dissociation temperature:

$$m_{\sigma}(T_{diss}) = 2m_{\pi}(T_{diss})$$

here: $T_{diss} = 180$ MeV

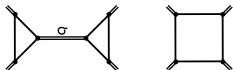
▶ Mott temperature:

$$m_{\pi}(T_{Mott}) = 2m_q(T_{Mott})$$

here: $T_{Mott} = 199$ MeV

► scattering amplitude

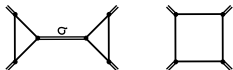
[Bernard et al., PLB (1991), Schulze, JPG (1995)]



- leading order $1/N_c$
- to be taken in s -, t -, and u -channel
- consistent with chiral low-energy theorems

► scattering amplitude

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- ## ► scattering length

$$a^l = \frac{1}{32\pi m_\pi} \mathcal{M}_{\pi\pi}^l(s = 4m_\pi^2, t = u = 0)$$

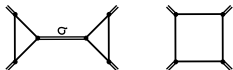
- ## ► chiral expansion [Weinberg, PRL (1966)]

$$a_W^0 = \frac{7m_\pi}{32\pi f_\pi^2}, \quad a_W^2 = -\frac{2m_\pi}{32\pi f_\pi^2}$$

In-medium $\pi\pi$ -scattering

▶ scattering amplitude

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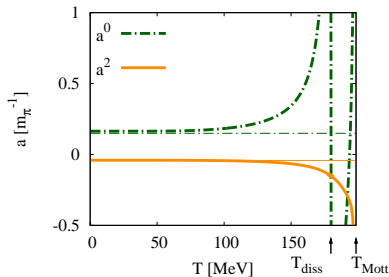
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▶ numerical results

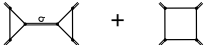
cf. [Quack et al. PLB (1995)]



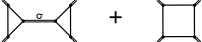
- ▶ Weinberg values at low T
- ▶ “Feshbach resonances” at T_{diss} and T_{Mott}

In-medium cross section

► isospin averaged cross section: $\left(\frac{d\sigma}{d\Omega}\right)_{cm} = \frac{1}{9} \sum_{l=0}^2 (2l+1) \frac{|\mathcal{M}'_{\pi\pi}|^2}{64\pi^2 s}$

► $i\mathcal{M}_{\pi\pi} =$ 

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▶ approximations:

1. Weinberg amplitude

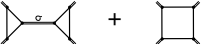
$$\mathcal{M}'_{\pi\pi} = 32\pi m_\pi a'_W$$

(T and momentum independent)

2. evaluate \mathcal{M} at threshold

$$\mathcal{M}'_{\pi\pi} = 32\pi m_\pi a^l(T)$$

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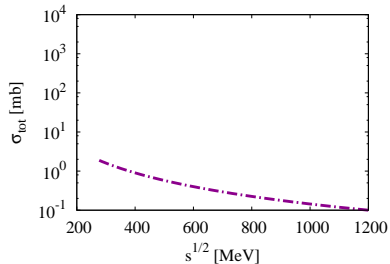
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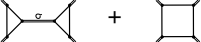
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total cross section ($T = 0$)



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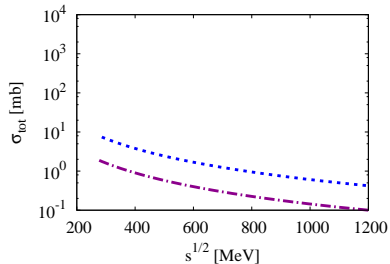
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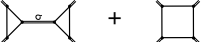
2. evaluate \mathcal{M} at threshold

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total cross section ($T = 150$ MeV)



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► $i\mathcal{M}_{\pi\pi} =$ 

► approximations:

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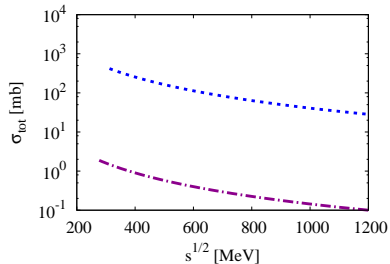
$$\mathcal{M}'_{\pi\pi} = 32\pi m_\pi a'_W$$

(T and momentum independent)

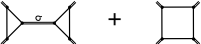
2. evaluate \mathcal{M} at threshold

$$\mathcal{M}'_{\pi\pi} = 32\pi m_\pi a'(T)$$

total cross section ($T = 177$ MeV)



► isospin averaged cross section:
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► $i\mathcal{M}_{\pi\pi} =$ 

► approximations:

1. Weinberg amplitude

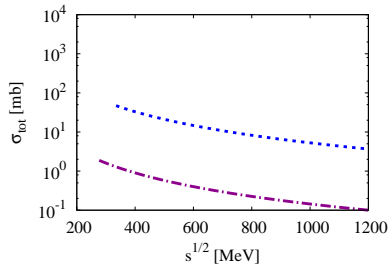
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(T and momentum independent)

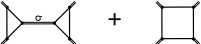
2. evaluate \mathcal{M} at threshold

$$\mathcal{M}'_{\pi\pi} = 32\pi m_\pi a^l(T)$$

total cross section ($T = 188$ MeV)



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► $i\mathcal{M}_{\pi\pi} =$ 

► approximations:

1. Weinberg amplitude

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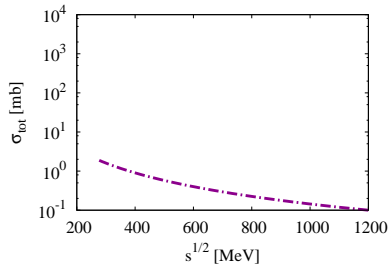
(T and momentum independent)

2. evaluate \mathcal{M} at threshold

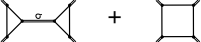
$$\mathcal{M}'_{\pi\pi} = 32\pi m_\pi a'(T)$$

3. keep momentum dependence of the σ exchange
(but still evaluate quark triangles and boxes at threshold)

total cross section ($T = 0$)



► isospin averaged cross section:
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► $i\mathcal{M}_{\pi\pi} =$ 

► approximations:

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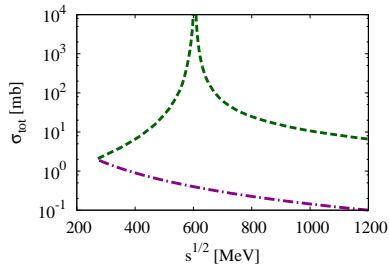
2. evaluate \mathcal{M} at threshold

$$\mathcal{M}'_{\pi\pi} = 32\pi m_\pi a'(T)$$

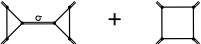
3. keep momentum dependence of the σ exchange

(but still evaluate quark triangles and boxes at threshold)

total cross section ($T = 0$)



- ▶ isospin averaged cross section:
$$\left(\frac{d\sigma}{d\Omega}\right)_{cm} = \frac{1}{9} \sum_{l=0}^2 (2l+1) \frac{|\mathcal{M}'_{\pi\pi}|^2}{64\pi^2 s}$$

▶ $i\mathcal{M}_{\pi\pi} =$ 

- ▶ approximations:

1. Weinberg amplitude

$$\mathcal{M}'_{\pi\pi} = 32\pi m_\pi a_W^l$$

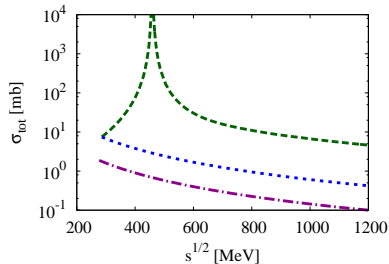
(T and momentum independent)

2. evaluate \mathcal{M} at threshold

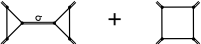

$$\mathcal{M}'_{\pi\pi} = 32\pi m_\pi a^l(T)$$

3. keep momentum dependence of the σ exchange
(but still evaluate quark triangles and boxes at threshold)

total cross section ($T = 150$ MeV)



- isospin averaged cross section:
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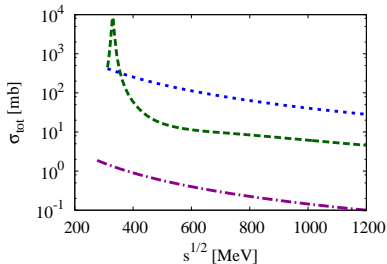
(T and momentum independent)

2. evaluate \mathcal{M} at threshold

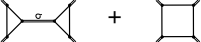
$$\mathcal{M}'_{\pi\pi} = 32\pi m_\pi a'(T)$$

3. keep momentum dependence of the σ exchange
(but still evaluate quark triangles and boxes at threshold)

total cross section ($T = 177$ MeV)



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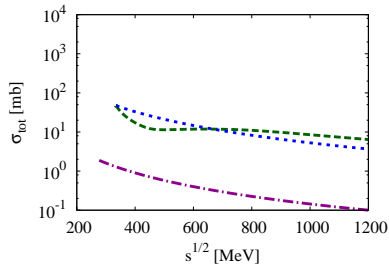
2. evaluate \mathcal{M} at threshold

$$\mathcal{M}'_{\pi\pi} = 32\pi m_\pi a'(T)$$

3. keep momentum dependence of the σ exchange

(but still evaluate quark triangles and boxes at threshold)

total cross section ($T = 188$ MeV)



Including the sigma-decay width



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DARMSTADT

► physical inconsistency:


$\sigma \leftrightarrow \pi\pi$ considered in scattering,
but not in RPA sigma propagator
 \Rightarrow width strongly underestimated

Including the sigma-decay width

▶ **physical inconsistency:**

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→ include 

▶ $1/N_c$ -correction term

▶ but there are many more


⇒ inconsistency with chiral symmetry

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→ include only imaginary part:

$$\Pi^{\text{dressed}} = \Pi_{\sigma}^{\text{RPA}} + \text{Im} \Pi_{\sigma}^{\pi\pi}$$


$$D_{\sigma}^{\text{dressed}} = \frac{-2g}{1 - 2g\Pi_{\sigma}^{\text{dressed}}}$$

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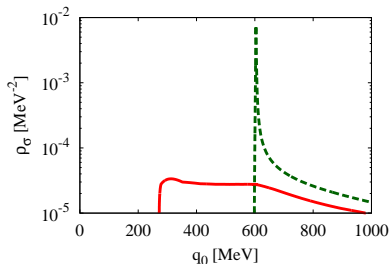
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▶ (unnormalized) spectral function

$$\rho_{\sigma}(q) = -2 \text{Im} D_{\sigma}(q)$$

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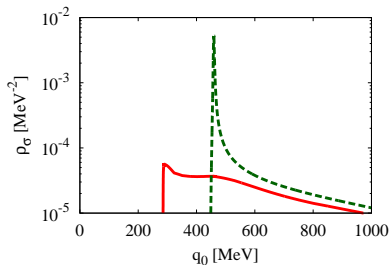
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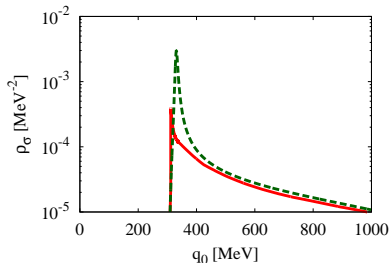
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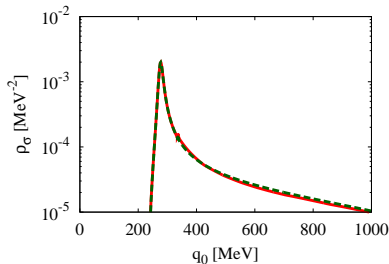
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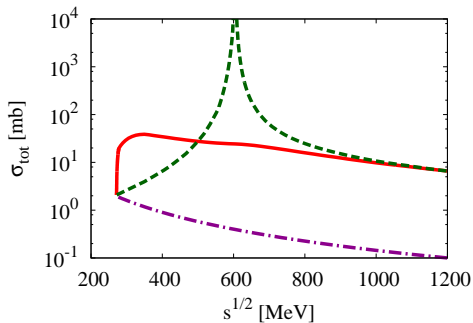
$$\rho_{\sigma}(q) = -2 \text{Im} D_{\sigma}(q)$$

$T = 188 \text{ MeV}$



Cross section with sigma-decay width

total cross section ($T = 0$)

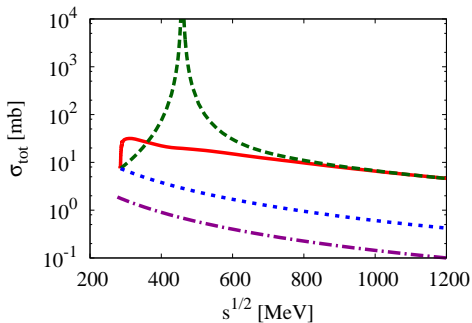


► $T = 0$:

$\sigma = \sigma_{\text{Weinberg}}$ at threshold

Cross section with sigma-decay width

total cross section ($T = 150$ MeV)



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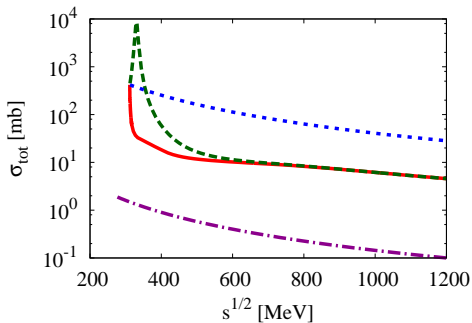
$\sigma = \sigma_{\text{Weinberg}}$ at threshold

▶ small and intermediate T :

$\sigma \gg \sigma(\mathcal{M})_{\text{thresh}} \gg \sigma_{\text{Weinberg}}$

Cross section with sigma-decay width

total cross section ($T = 177$ MeV)



▶ $T = 0$:

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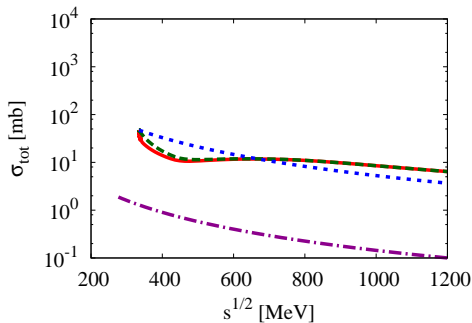
$\sigma \gg \sigma(\mathcal{M})_{thresh} \gg \sigma_{Weinberg}$

▶ $T \approx T_{diss}$:

$\sigma(\mathcal{M})_{thresh} \gg \sigma \gg \sigma_{Weinberg}$

Cross section with sigma-decay width

total cross section ($T = 188 \text{ MeV}$)



▶ $T = 0$:

$\sigma = \sigma_{Weinberg}$ at threshold

▶ small and intermediate T :

$\sigma \gg \sigma(\mathcal{M})_{thresh} \gg \sigma_{Weinberg}$

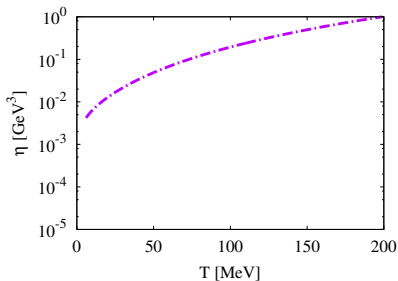
▶ $T \approx T_{diss}$:

$\sigma(\mathcal{M})_{thresh} \gg \sigma \gg \sigma_{Weinberg}$

▶ $T > T_{diss}$:

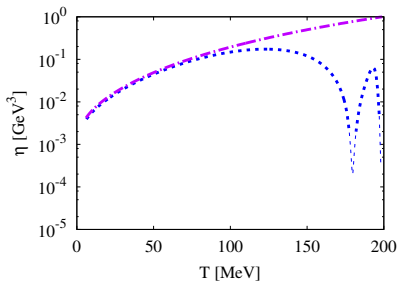
$\sigma \rightarrow \pi\pi$ irrelevant

Shear viscosity: numerical results



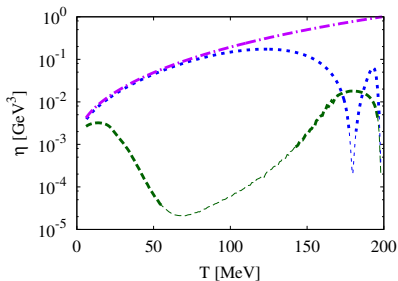
► Weinberg

Shear viscosity: numerical results



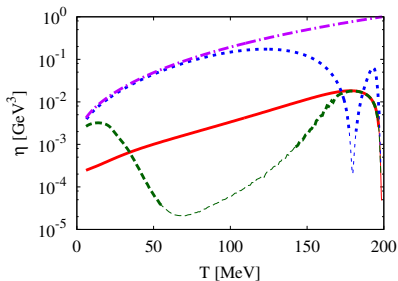
- ▶ Weinberg
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Shear viscosity: numerical results



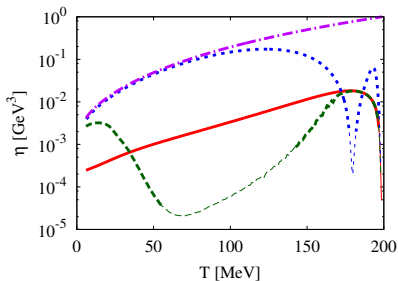
- ▶ Weinberg
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Shear viscosity: numerical results



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- ▶ RPA σ -propagator
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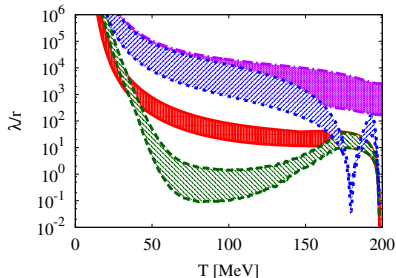
Shear viscosity: numerical results



- ▶ Weinberg
- ▶ $\mathcal{M} = \mathcal{M}_{threshold}$
- ▶ RPA σ -propagator
- ▶ dressed σ -propagator

▶ validity of the kinetic approach:

- ▶ criterion: $\frac{\lambda}{r} \gg 1$ (dilute gas)
- ▶ $\lambda = \frac{1}{n\sigma}$ mean free path
- ▶ r = interaction range
(e.g., $1/m_\sigma$, $1/m_\pi$, hard sphere: $\sqrt{\frac{\sigma}{\pi}}$)



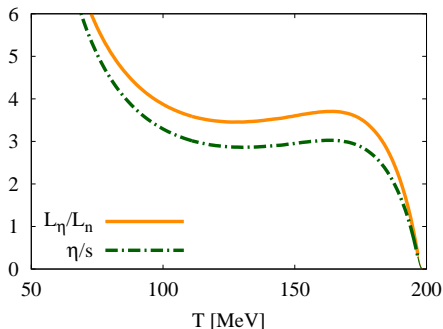
► most popular measure: η/s

- η from our “best model” (dressed σ -meson)
- entropy density of an ideal pion gas

► alternative measure: L_η/L_n

[Liao & Koch, PRC (2010)]

- $L_\eta = \frac{\eta}{hc_s}$, $L_n = n^{-1/3}$



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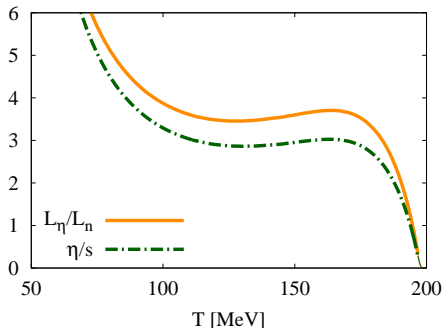
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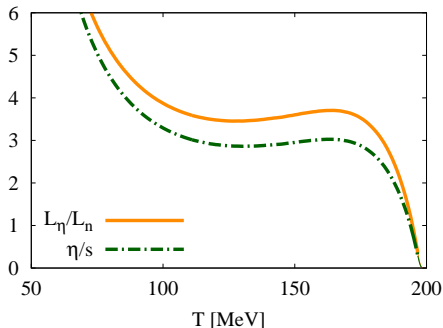
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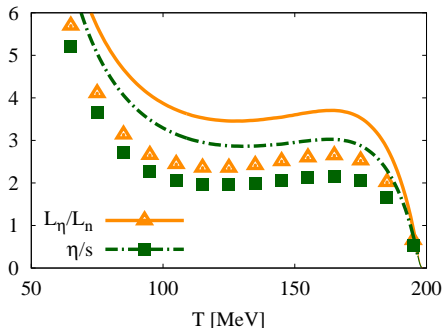
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▶ simple estimate: $\eta \approx \frac{1}{3} n \bar{p} \lambda = \frac{\bar{p}}{3\sigma(\bar{p})}$ works quite well ...





▶ summary:

- ▶ shear viscosity from $\pi\pi$ -scattering in the NJL model in kinetic theory
- ▶ agreement with lowest-order χ PT (Weinberg) at low T , much lower values when approaching the crossover
- ▶ quantitative results very sensitive to details of the model

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- ▶ quantitative results very sensitive to details of the model

▶ outlook:

- ▶ better description of p -wave $\pi\pi$ scattering (include ρ -meson)
- ▶ further scattering channels, e.g., kaons
- ▶ ...