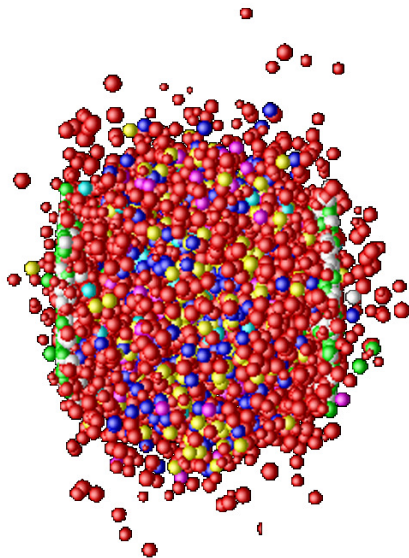


**Institut für
Theoretische Physik I**



Parton dynamics in heavy-ion collisions from FAIR to LHC

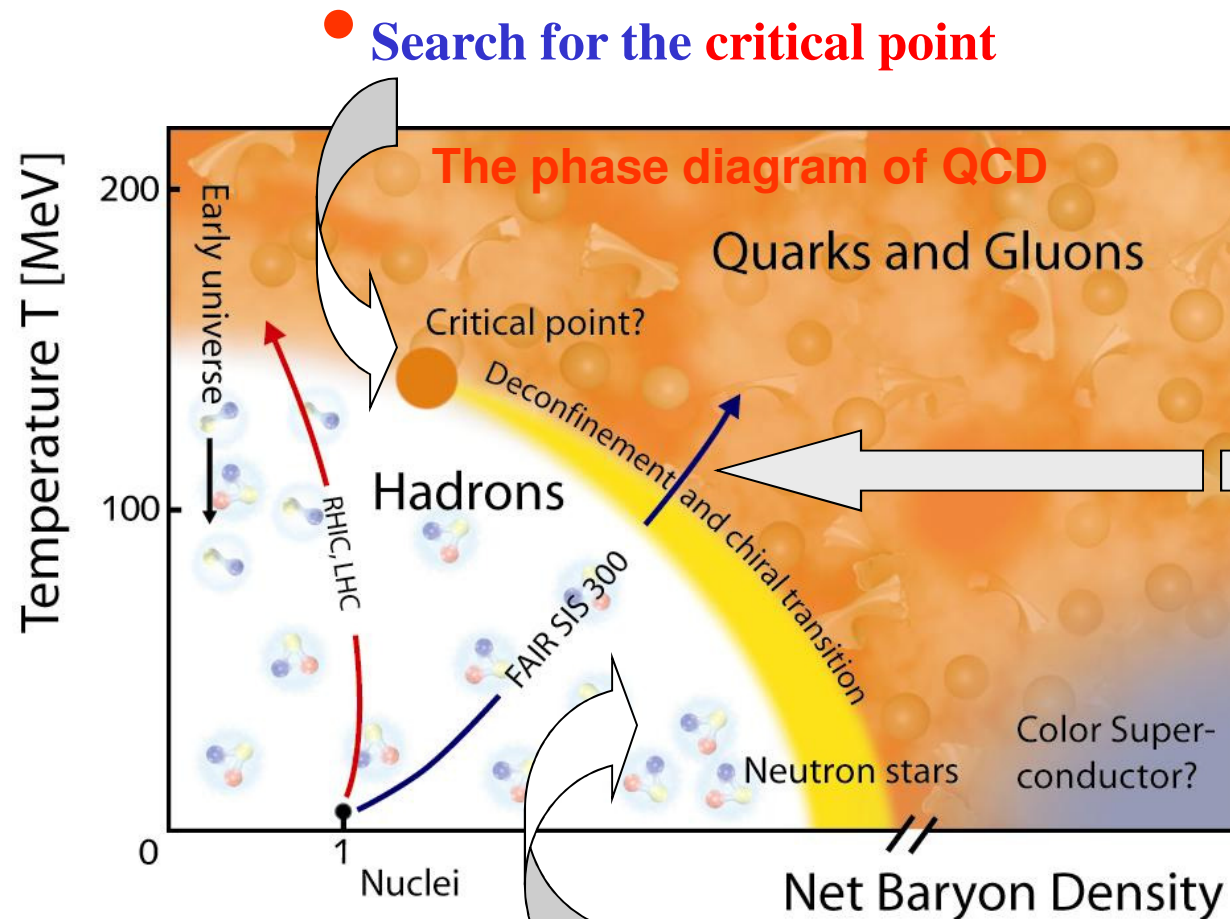


Wolfgang Cassing

Erice, 21.09.2012



The holy grail:



- Study of the **phase transition** from hadronic to partonic matter – **Quark-Gluon-Plasma**

- Study of the **in-medium** properties of hadrons at high baryon density and temperature
- Study of the partonic medium beyond the phase boundary

Signals of the phase transition:

- Multi-strange particle enhancement in A+A
- Charm suppression
- Collective flow (v_1, v_2, v_3, v_4)
- Thermal dileptons
- Jet quenching and angular correlations
- High p_T suppression of hadrons
- Nonstatistical event by event fluctuations and correlations
- Chiral Magnetic Effect ?

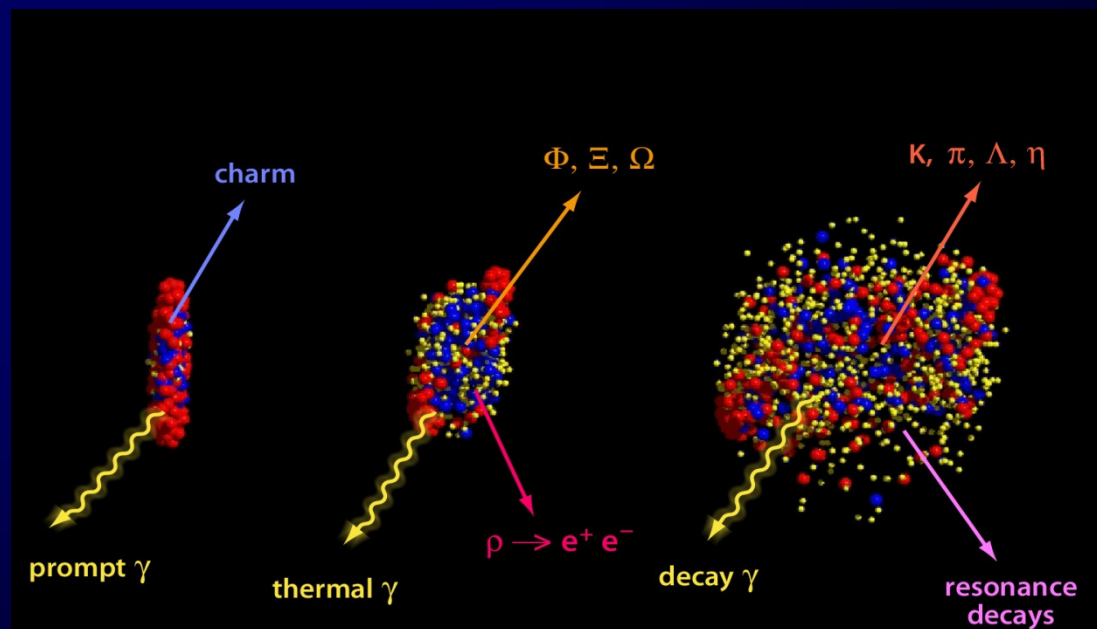
Experiment: measures final hadrons and leptons

How to learn about physics from data?

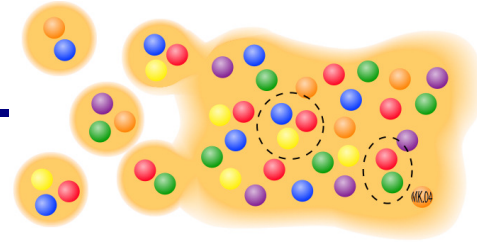
Compare with theory!



Microscopic transport models provide a unique dynamical description of nonequilibrium effects in heavy-ion collisions !



From hadrons to partons



In order to study the **phase transition** from hadronic to partonic matter – **Quark-Gluon-Plasma** –

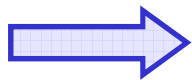
we need a **consistent non-equilibrium (transport) model with**

➤ **explicit parton-parton interactions** (i.e. between quarks and gluons) beyond strings!

➤ **explicit phase transition** from hadronic to partonic degrees of freedom

➤ **IQCD EoS** for partonic phase

Transport theory: off-shell Kadanoff-Baym equations for the Green-functions $S_h^<(x,p)$ in phase-space representation for the **partonic and hadronic phase**



Parton-Hadron-String-Dynamics (PHSD)



QGP phase described by

Dynamical QuasiParticle Model (DQPM)

W. Cassing, E. Bratkovskaya, PRC 78 (2008) 034919;
NPA831 (2009) 215;
W. Cassing, EPJ ST 168 (2009) 3

A. Peshier, W. Cassing, PRL 94 (2005) 172301;
Cassing, NPA 791 (2007) 365; NPA 793 (2007)

The Dynamical QuasiParticle Model (DQPM)

Basic idea: Interacting quasiparticles

- massive quarks and gluons (g, q, q_{bar}) with spectral functions :

$$\rho(\omega) = \frac{\gamma}{E} \left(\frac{1}{(\omega - E)^2 + \gamma^2} - \frac{1}{(\omega + E)^2 + \gamma^2} \right) \quad E^2 = \mathbf{p}^2 + M^2 - \gamma^2$$

■ quarks

mass: $m^2(T) = \frac{N_c^2 - 1}{8N_c} g^2 \left(T^2 + \frac{\mu_q^2}{\pi^2} \right)$

width: $\gamma_q(T) = \frac{N_c^2 - 1}{2N_c} \frac{g^2 T}{4\pi} \ln \frac{c}{g^2}$

running coupling: $\alpha_s(T) = g^2(T)/(4\pi)$

$$g^2(T/T_c) = \frac{48\pi^2}{(11N_c - 2N_f) \ln(\lambda^2(T/T_c - T_s/T_c)^2)}$$

➤ **fit to lattice (IQCD) results** (e.g. entropy density)

with 3 parameters: $T_s/T_c=0.46$; $c=28.8$; $\lambda=2.42$

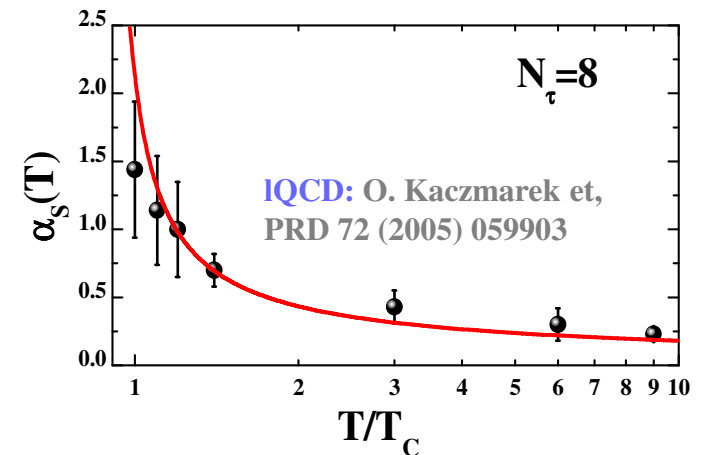
➔ **quasiparticle properties** (mass, width)

■ gluons:

A. Peshier, PRD 70 (2004) 034016

$$M^2(T) = \frac{g^2}{6} \left((N_c + \frac{1}{2}N_f) T^2 + \frac{N_c}{2} \sum_q \frac{\mu_q^2}{\pi^2} \right) \quad N_c = 3, N_f = 3$$

$$\gamma_g(T) = N_c \frac{g^2 T}{4\pi} \ln \frac{c}{g^2}$$



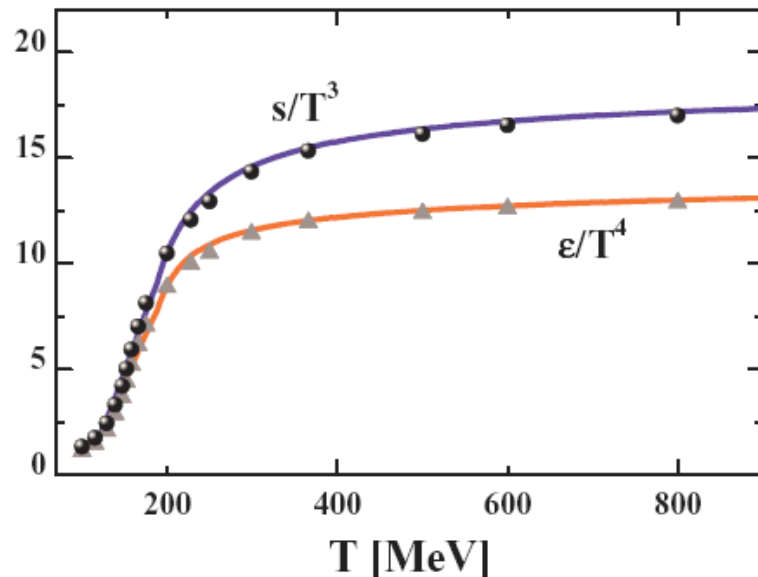
DQPM: Peshier, Cassing, PRL 94 (2005) 172301;
Cassing, NPA 791 (2007) 365; NPA 793 (2007)

DQPM thermodynamics ($N_f=3$) and IQCD

entropy $s = \frac{\partial P}{\partial T} \rightarrow$ pressure P

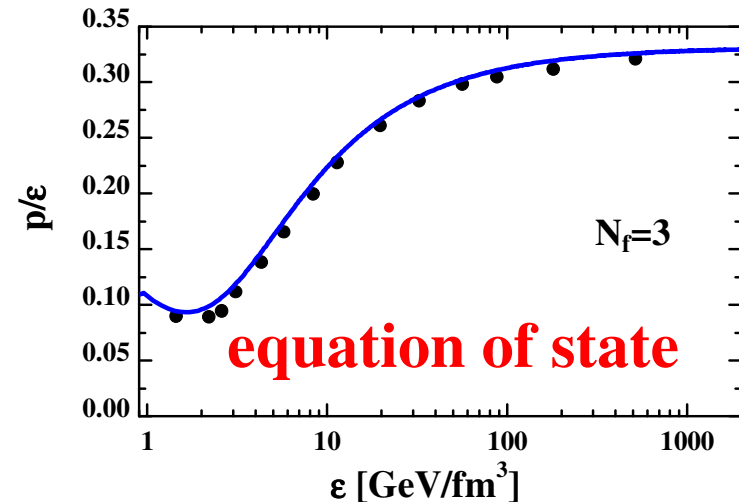
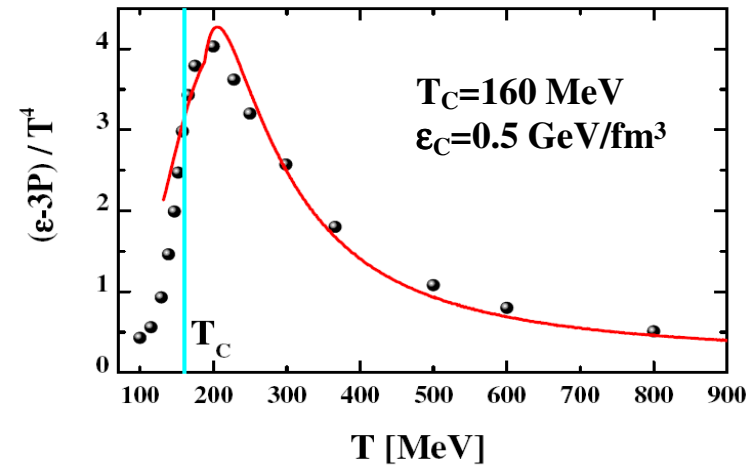
energy density: $\epsilon = Ts - P$

IQCD: Wuppertal-Budapest group
Y. Aoki et al., JHEP 0906 (2009) 088.



interaction measure:

$$W(T) := \epsilon(T) - 3P(T) = Ts - 4P$$

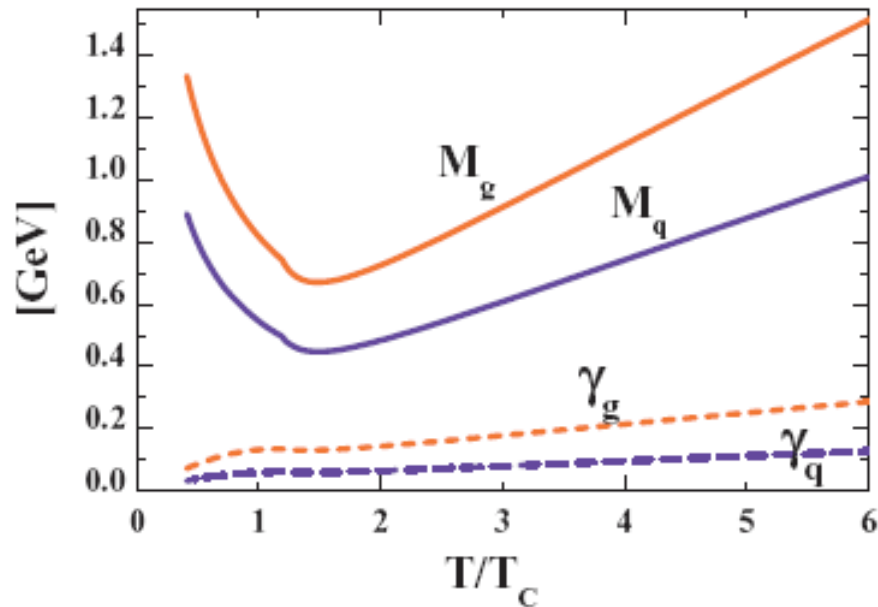


DQPM gives a good description of IQCD results !

The Dynamical QuasiParticle Model (DQPM)

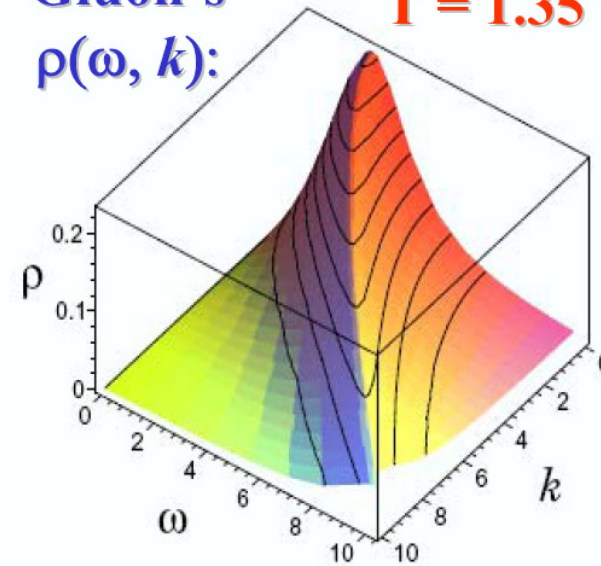
→ Quasiparticle properties:

- large width and mass for gluons and quarks

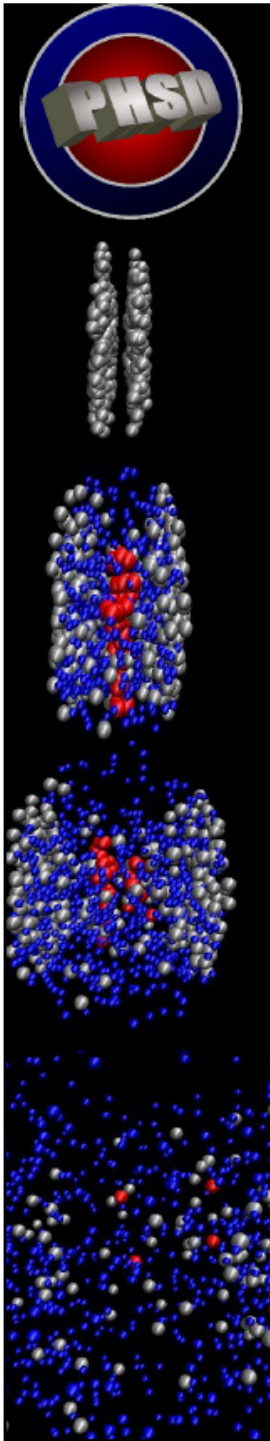


→ Broad spectral function :

Gluon's
 $\rho(\omega, k)$: $T = 1.35 T_c$



- DQPM matches well lattice QCD
- DQPM provides mean-fields (1PI) for gluons and quarks as well as effective 2-body interactions (2PI)
- DQPM gives transition rates for the formation of hadrons → PHSD



PHSD - basic concept

Initial A+A collisions – HSD: string formation and decay to pre-hadrons

Fragmentation of pre-hadrons into quarks: using the quark spectral functions from the **Dynamical QuasiParticle Model (DQPM)** - approximation to QCD

Partonic phase: quarks and gluons (= ,dynamical quasiparticles‘) with **off-shell spectral functions** (width, mass) defined by the DQPM

elastic and inelastic parton-parton interactions: using the effective cross sections from the DQPM

✓ $q + qbar$ (flavor neutral) \Leftrightarrow **gluon** (colored)

✓ **gluon + gluon** \Leftrightarrow **gluon** (possible due to large spectral width)

✓ $q + qbar$ (color neutral) \Leftrightarrow **hadron resonances**

self-generated mean-field potential for quarks and gluons !

Hadronization: based on DQPM - **massive, off-shell quarks and gluons** with broad spectral functions hadronize to **off-shell mesons and baryons:**

gluons \rightarrow $q + qbar$; $q + qbar \rightarrow$ **meson (or string);**

$q + q + q \rightarrow$ **baryon (or string)** (strings act as ,doorway states‘ for hadrons)

Hadronic phase: hadron-string interactions – **off-shell HSD**

W. Cassing, E. Bratkovskaya, PRC 78 (2008) 034919;
NPA831 (2009) 215; EPJ ST 168 (2009) 3; NPA856 (2011) 162.

PHSD: Hadronization details

**Local covariant off-shell transition rate for q+qbar fusion
=> meson formation**

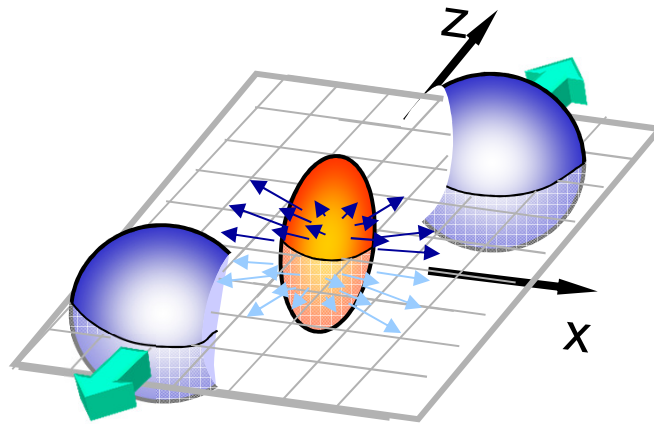
$$\frac{dN_m(x, p)}{d^4x d^4p} = Tr_q Tr_{\bar{q}} \delta^4(p - p_q - p_{\bar{q}}) \delta^4\left(\frac{x_q + x_{\bar{q}}}{2} - x\right) \\ \times \omega_q \rho_q(p_q) \omega_{\bar{q}} \rho_{\bar{q}}(p_{\bar{q}}) |v_{q\bar{q}}|^2 W_m(x_q - x_{\bar{q}}, p_q - p_{\bar{q}}) \\ \times N_q(x_q, p_q) N_{\bar{q}}(x_{\bar{q}}, p_{\bar{q}}) \delta(\text{flavor, color}).$$



using $Tr_j = \sum_j \int d^4x_j d^4p_j / (2\pi)^4$

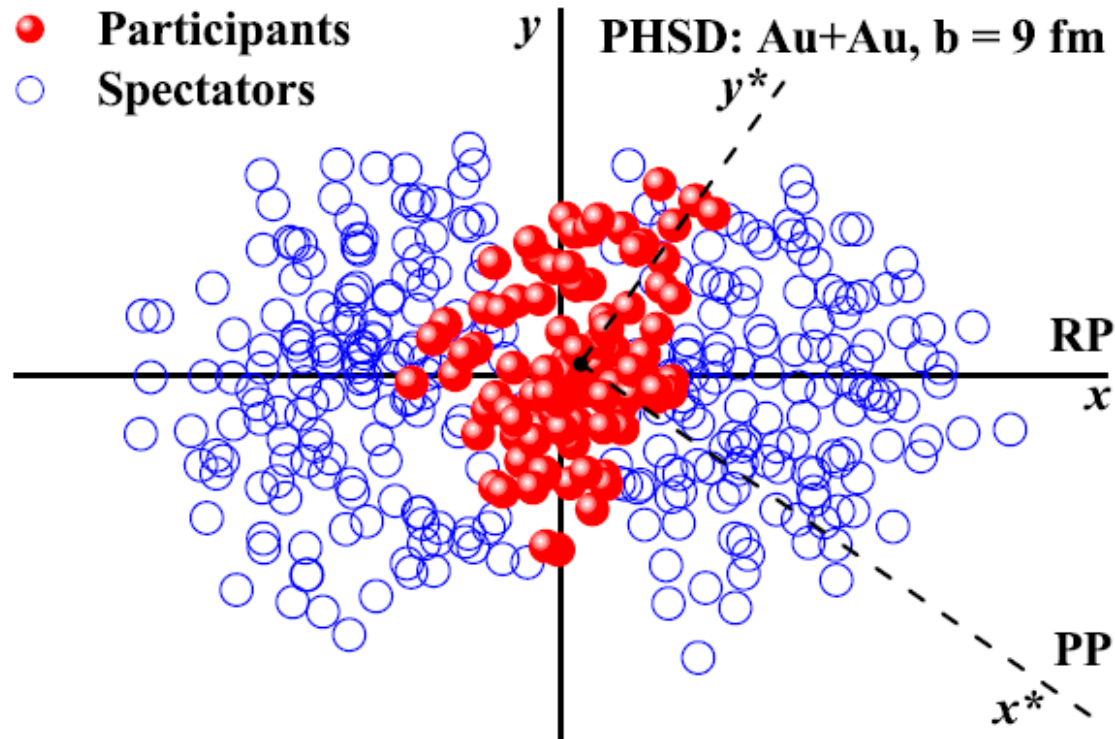
- $N_j(x, p)$ is the **phase-space density of parton j** at space-time position x and 4-momentum p
- W_m is the **phase-space distribution of the formed ,pre-hadrons‘**:
(Gaussian in phase space) $\sqrt{\langle r^2 \rangle} = 0.66 \text{ fm}$
- $v_{q\bar{q}}$ is the **effective quark-antiquark interaction from the DQPM**

**Collective flow:
anisotropy coefficients (v_1, v_2, v_3, v_4)
in A+A**



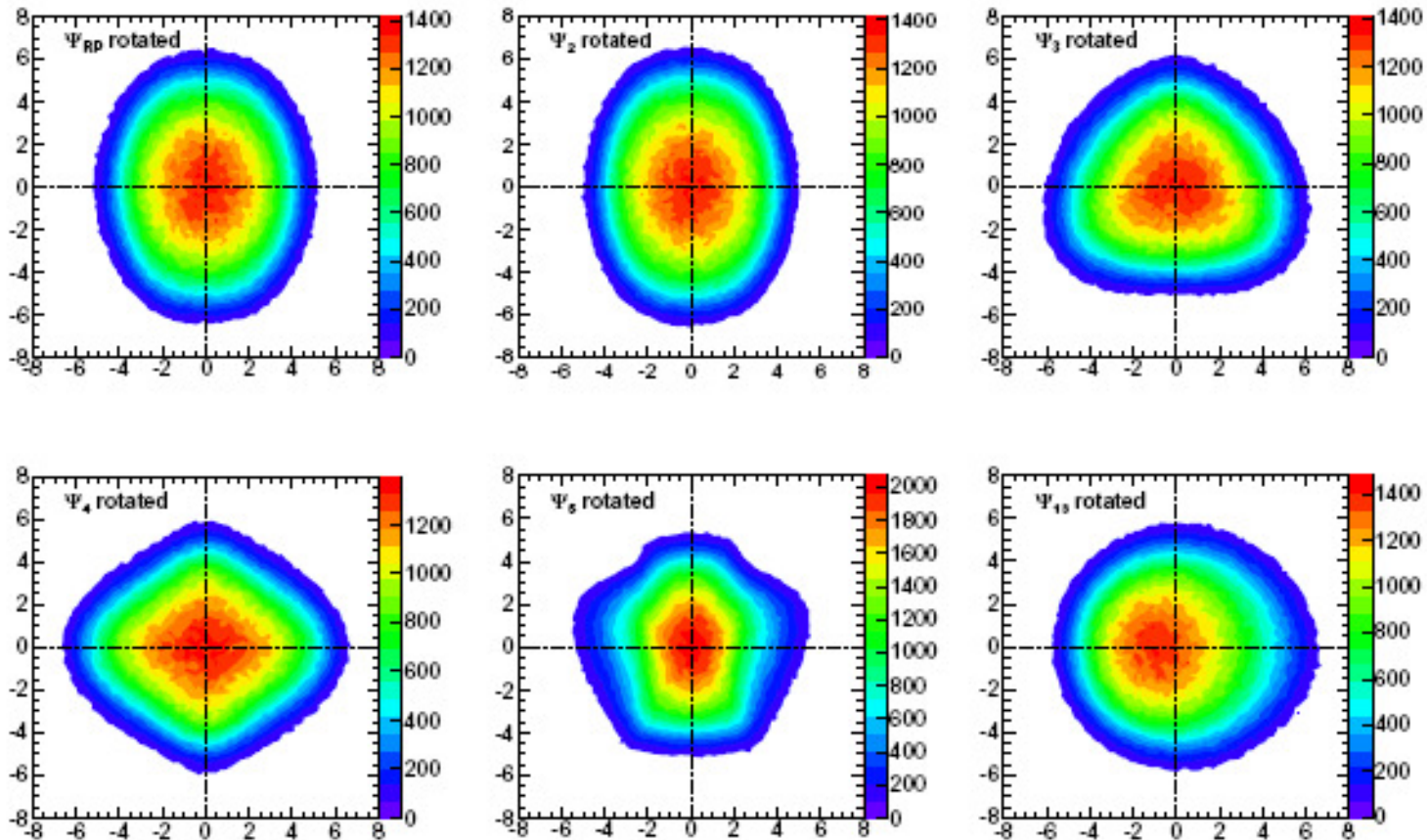
almond shaped interaction zone?

Initial participants and spectators in event by event calculations



In each event the reaction plane is tilted in coordinate space
and of irregular shape \rightarrow higher order harmonics!

Final angular distributions of hadrons



$$E \frac{d^3 N}{d^3 p} =$$

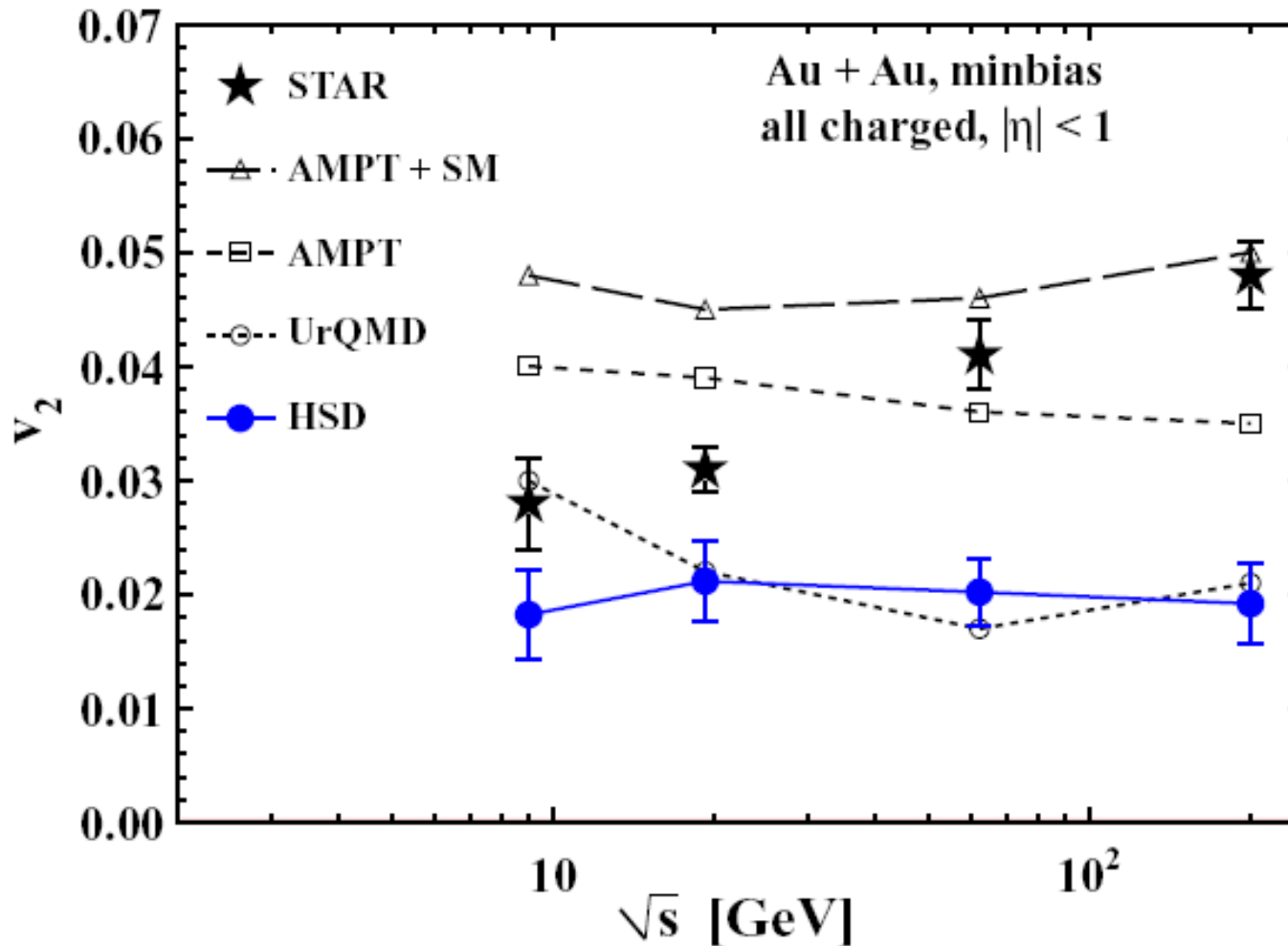
S. A. Voloshin

$$\frac{d^2 N}{2\pi p_T dp_T dy} \left(1 + \sum_{n=1}^{\infty} 2v_n(p_T) \cos(n(\psi - \Psi_n)) \right)$$

show higher order

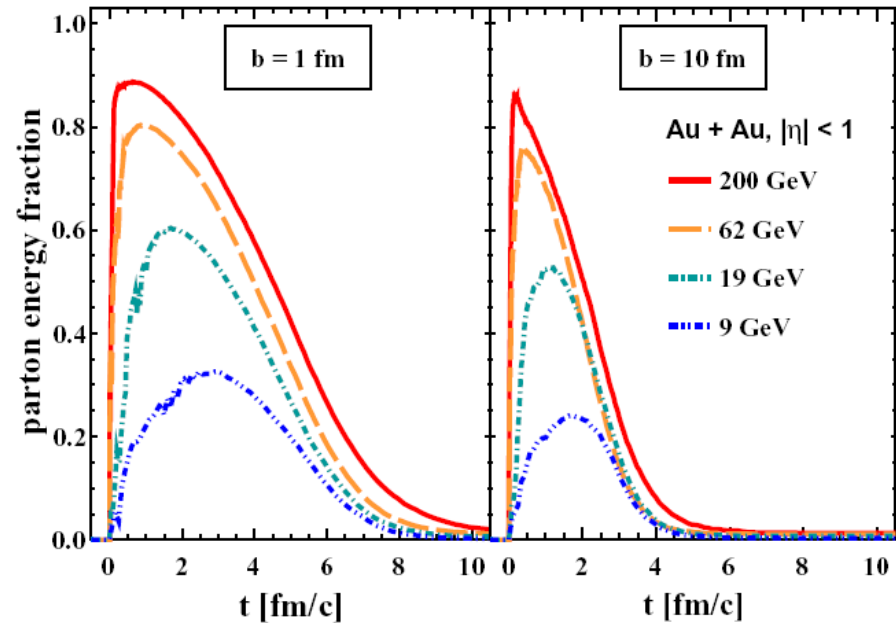
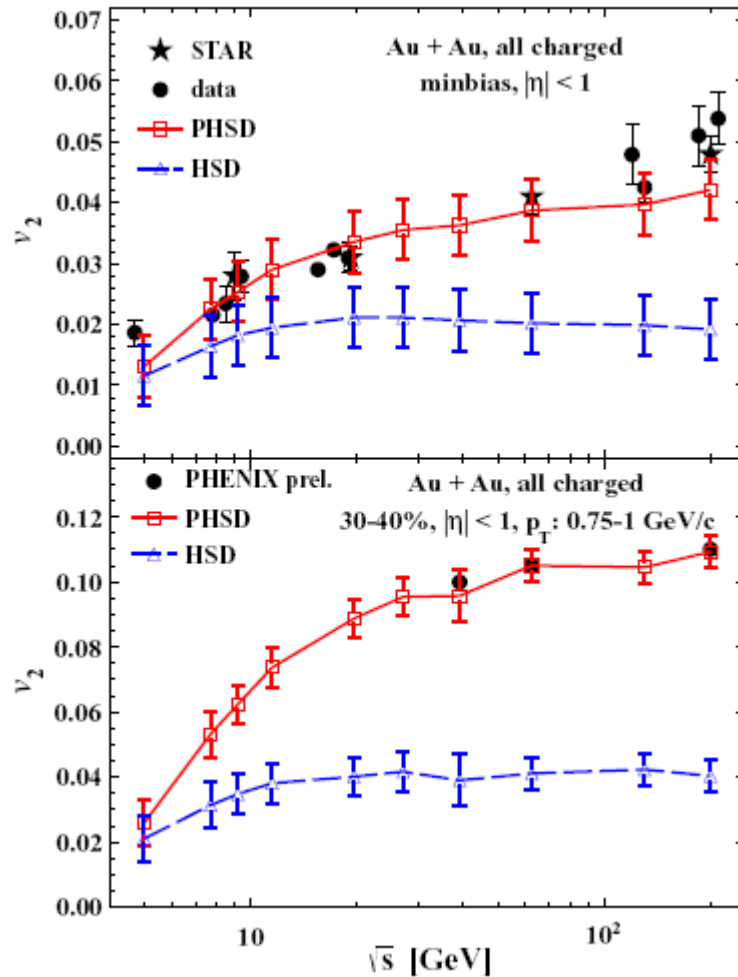
harmonics v_n

Excitation function of elliptic flow



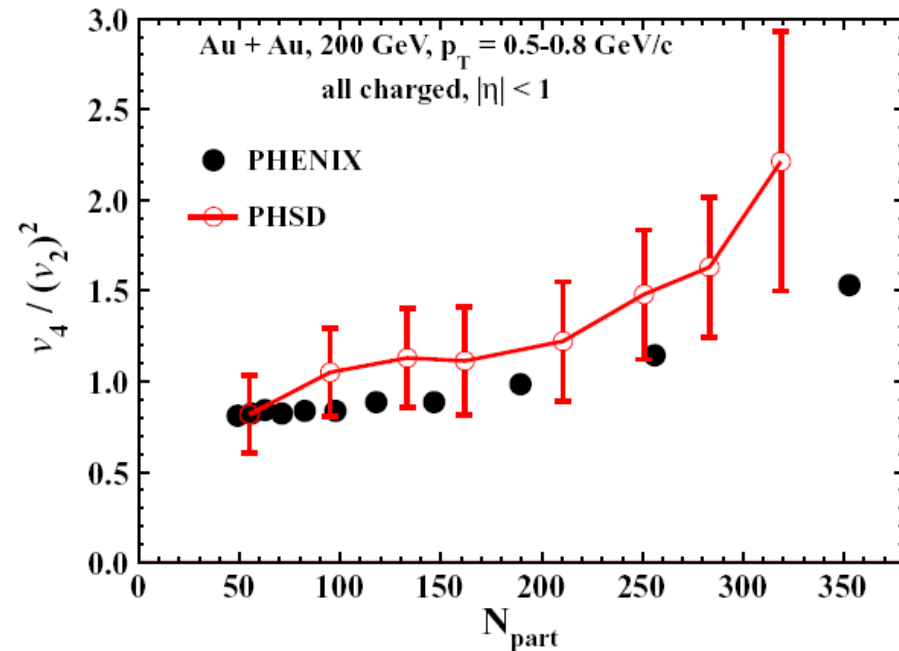
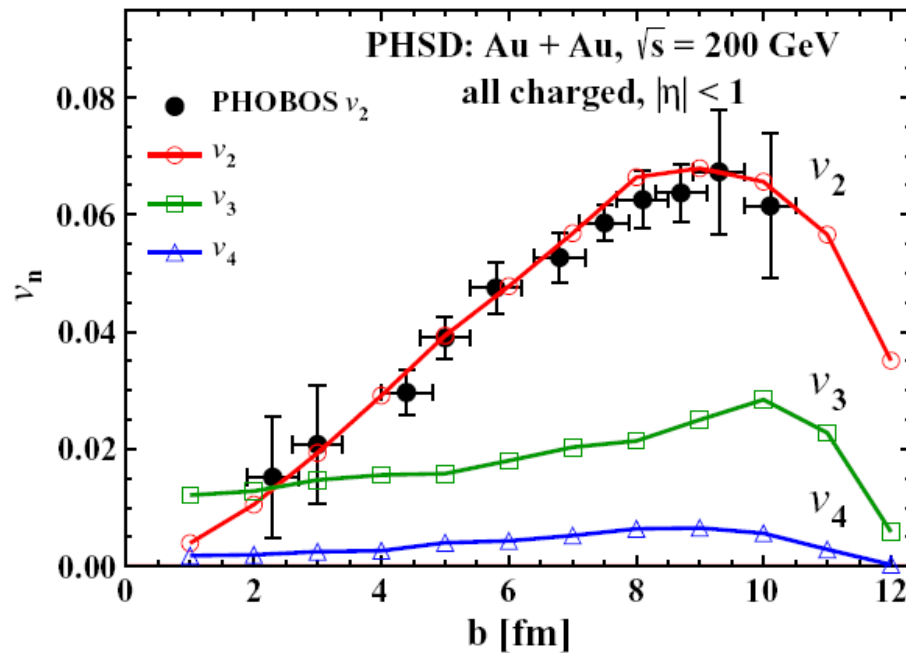
is not described by hadron-string or purely partonic models !

Elliptic flow from PHSD versus STAR/PHENIX



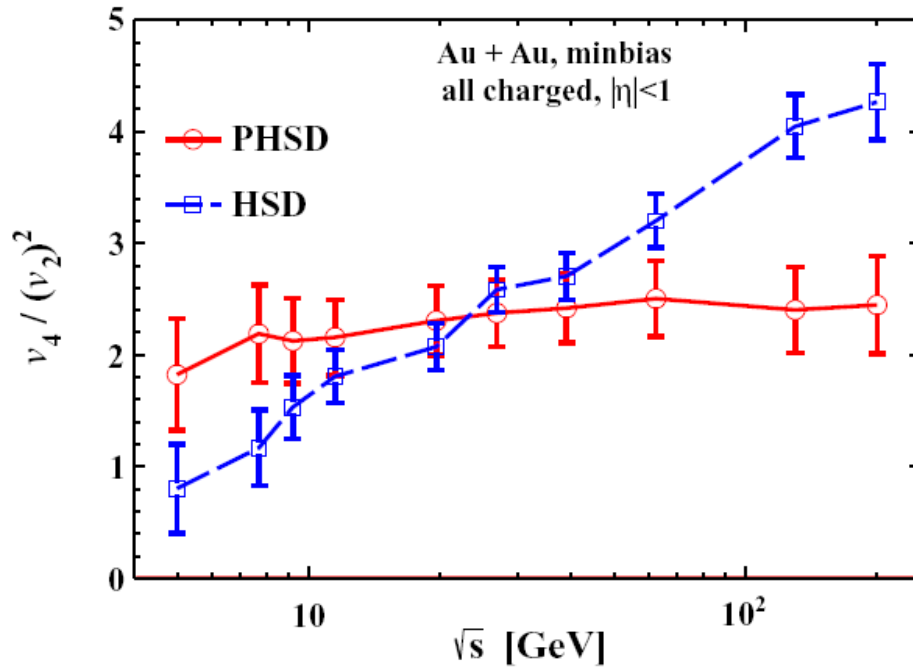
is reasonably described by PHSD due to an increasing fraction of partonic degrees-of-freedom !

Flow coefficients versus centrality



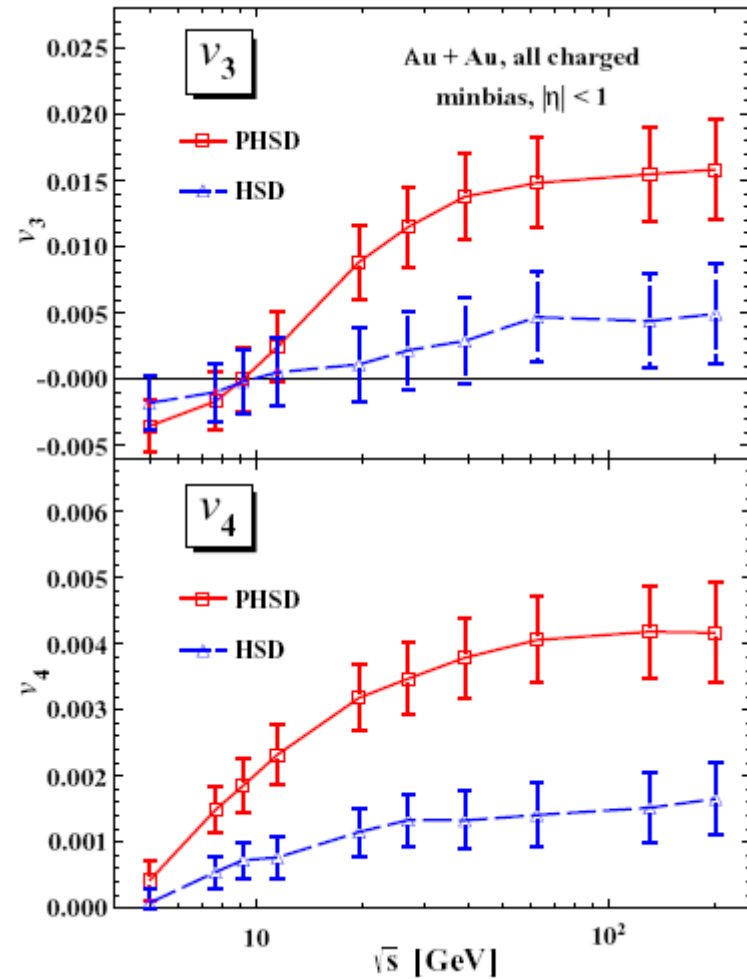
increase of v_2 with impact parameter but rather flat v_3 and v_4

More excitation functions



Very low v_3 and v_4 at FAIR energies

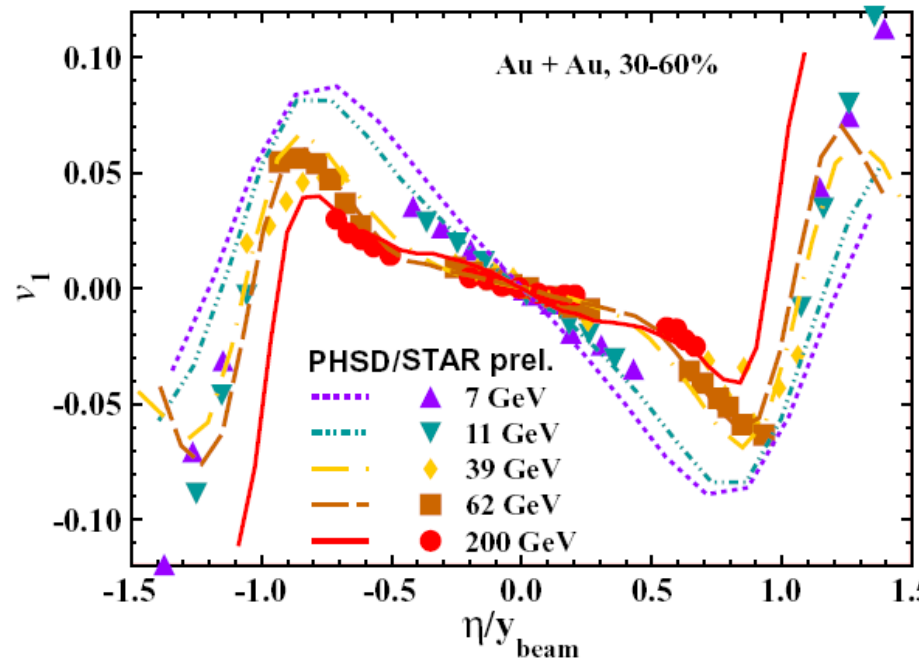
Almost constant $v_4 / (v_2)^2$ for PHSD



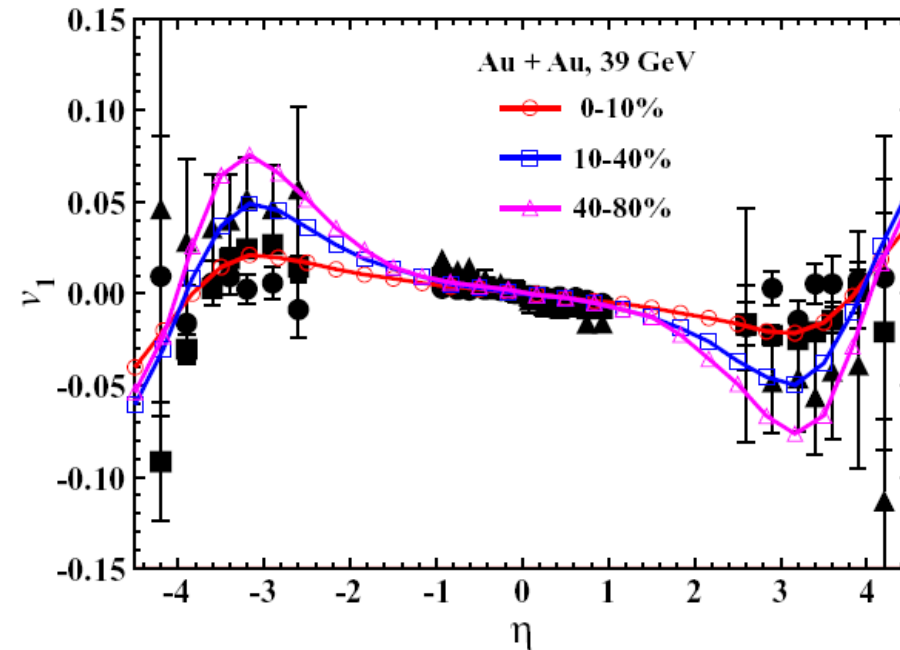
In-plane flow v_1



versus beam energy



versus centrality

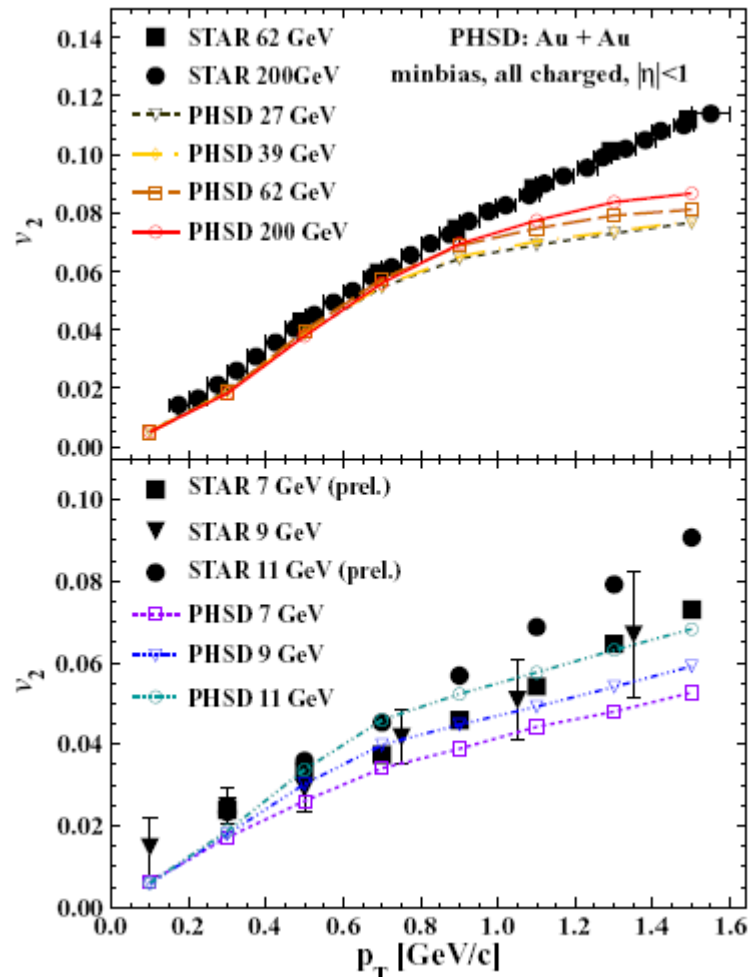


- **PHSD: v_1 vs. pseudo-rapidity follows an approximate scaling for high invariant energies $s^{1/2}=39, 62, 200$ GeV - in line with experimental data – whereas at low energies the scaling is violated!**

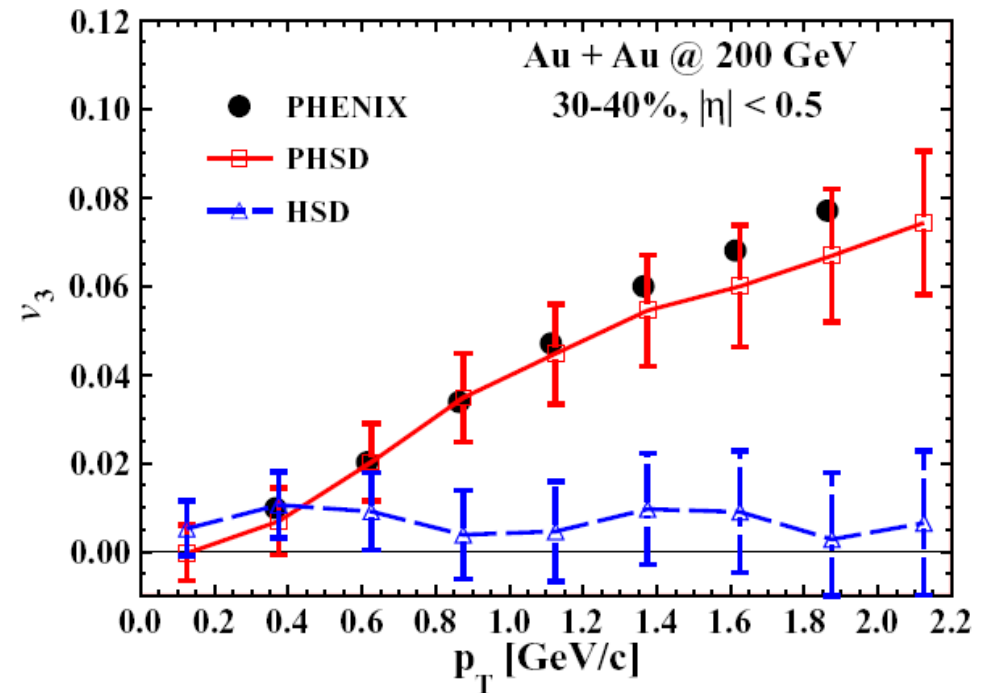
Transverse momentum dependence



elliptic flow



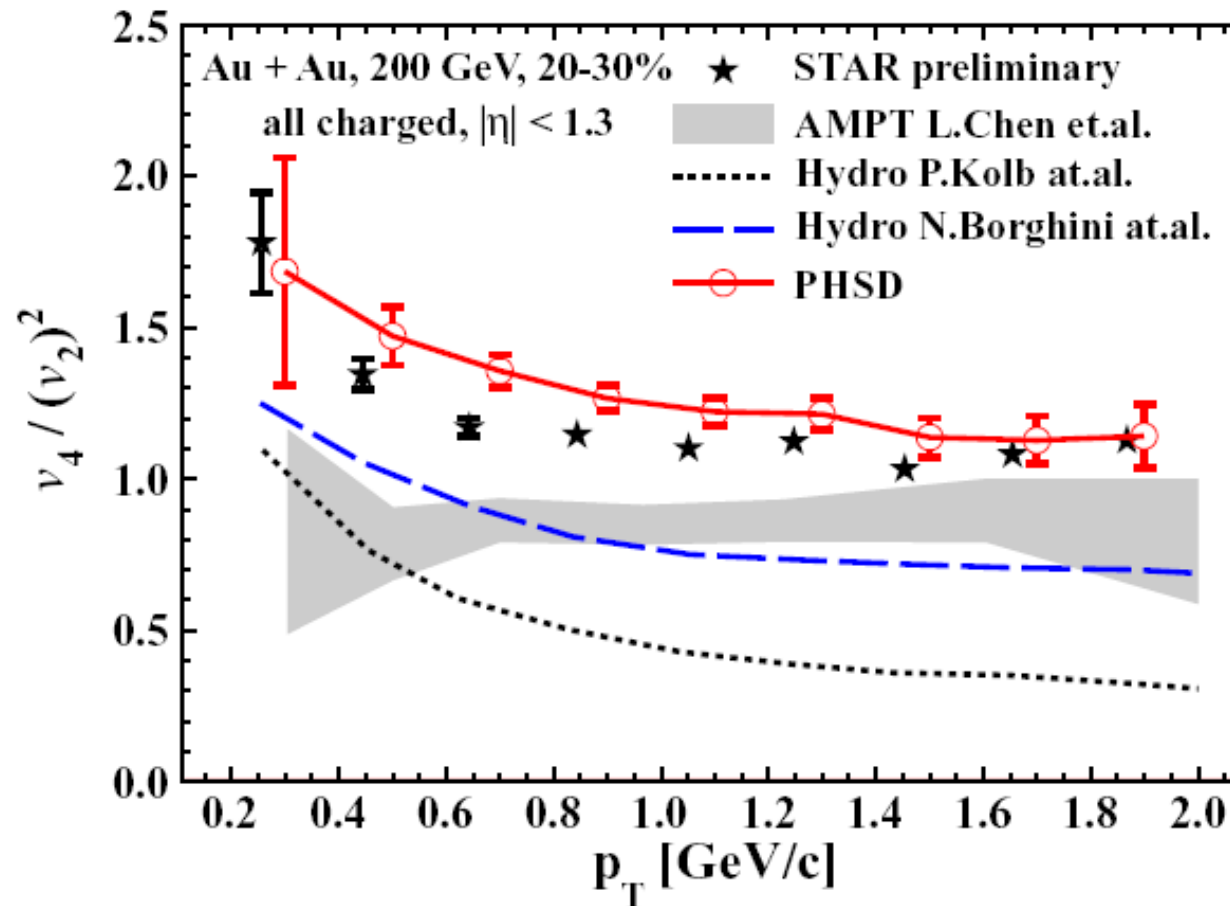
triangular flow



needs partonic degrees-of-freedom !

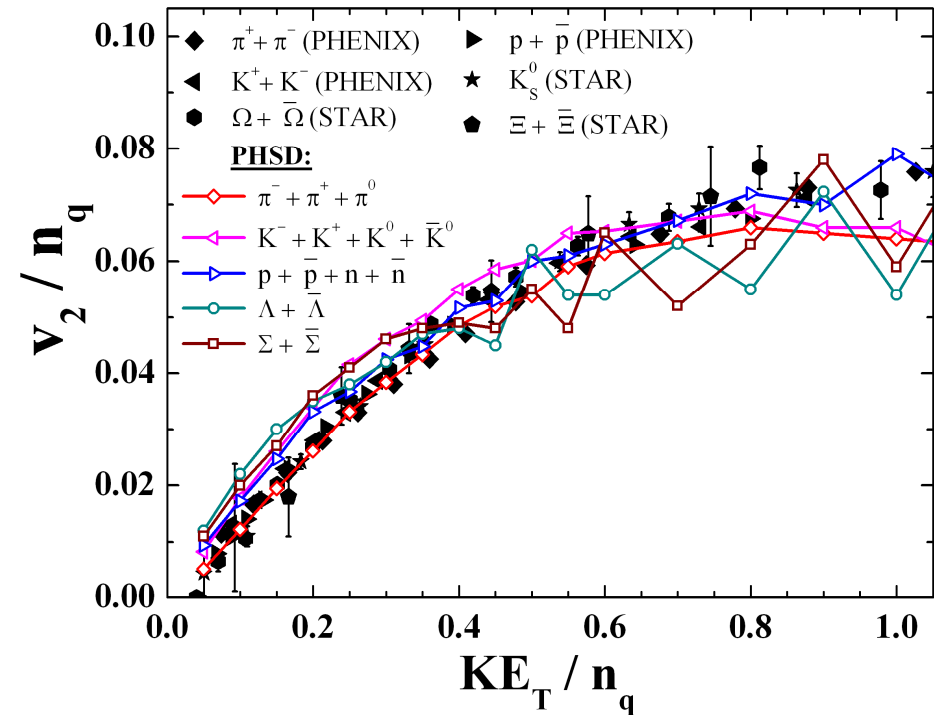
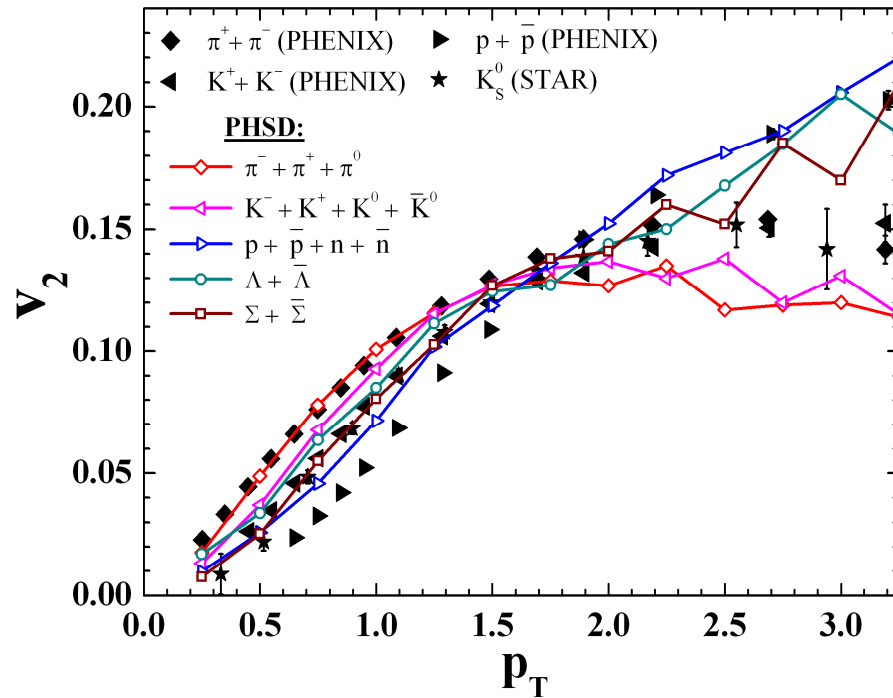


Ratio $v_4/(v_2)^2$ vs. p_T



is very sensitive to the microscopic dynamics !

Elliptic flow scaling at RHIC

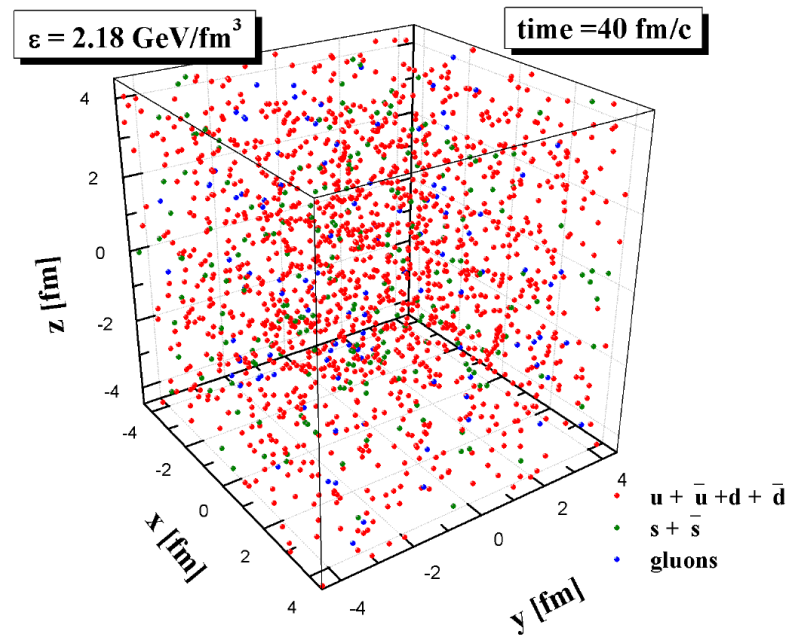


- The mass splitting at low p_T is approximately reproduced as well as the meson-baryon splitting for $p_T > 2$ GeV/c !
- The scaling of v_2 with the number of constituent quarks n_q is roughly in line with the data .

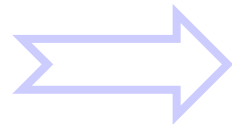
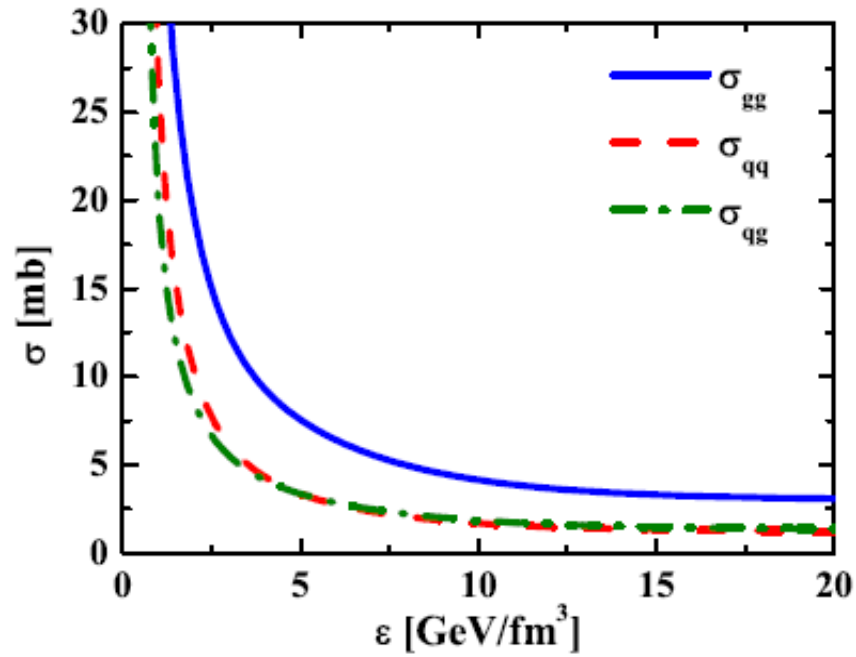
E. Bratkovskaya, W. Cassing, V. Konchakovski, O. Linnyk, NPA856 (2011) 162

Shear viscosity from PHSD in a box

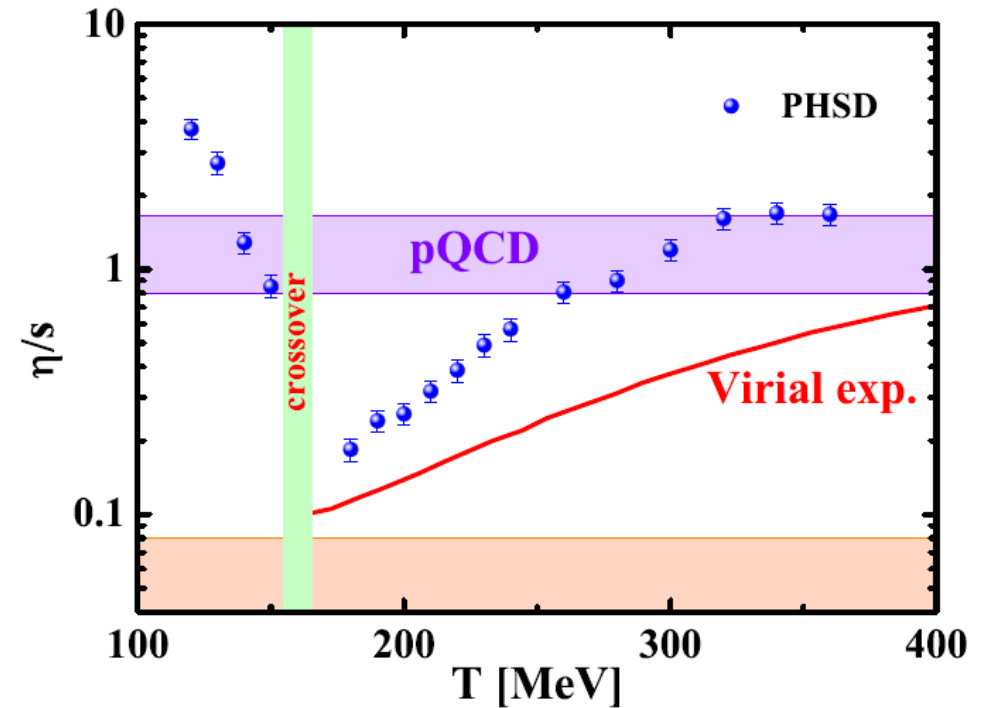
$$\eta = \frac{1}{T} \int d^3r \int_0^{\infty} dt \langle \pi^{xy}(\mathbf{0}, 0) \pi^{xy}(\mathbf{r}, t) \rangle_{\text{equil}}$$



Transport coefficients



Cross sections in PHSD



Summary



● **PHSD** provides a consistent description of **off-shell parton dynamics** in line with the lattice QCD equation of state

● **PHSD** versus **experimental observables**:

enhancement of meson m_T slopes (at top SPS and RHIC)

strange antibaryon enhancement (at SPS)

partonic emission of high mass dileptons at SPS and RHIC

enhancement of collective flow v_2 with increasing energy

quark number scaling of v_2 (at RHIC)

jet suppression

near-side ridge ...

⇒ **evidence for strong partonic interactions in the early phase of relativistic heavy-ion reactions from FAIR to LHC energies**

⇒ **initial fluctuations generate odd flow components!**

⇒ **Shear viscosity to entropy density has a minimum close to T_c .**

⇒ **Collective flow is also generated by the scalar partonic mean-field!**



PHSD group



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Mark Gorenstein

Barcelona Univ.

Laura Tolos, Angel Ramos



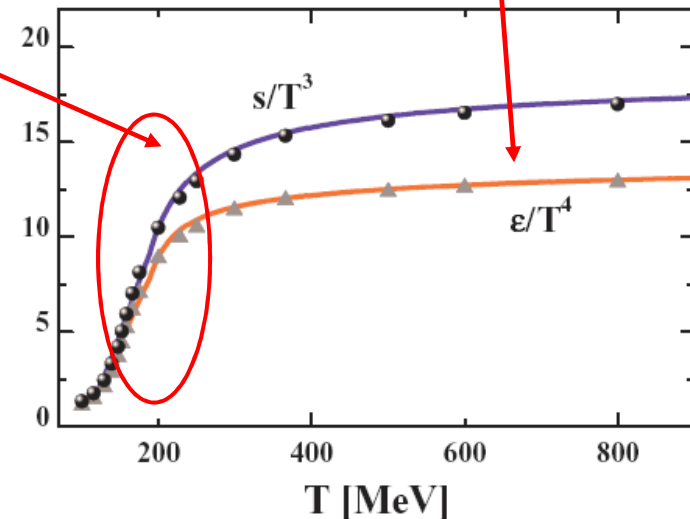
Outlook - Perspectives



What is the stage of matter close to T_c at finite quark chemical potential:

- ❑ 1st order phase transition?
- ❑ ,Mixed‘ phase = interaction of partonic and hadronic degrees of freedom?

Lattice EQS \rightarrow ,crossover‘ , $T > T_c$



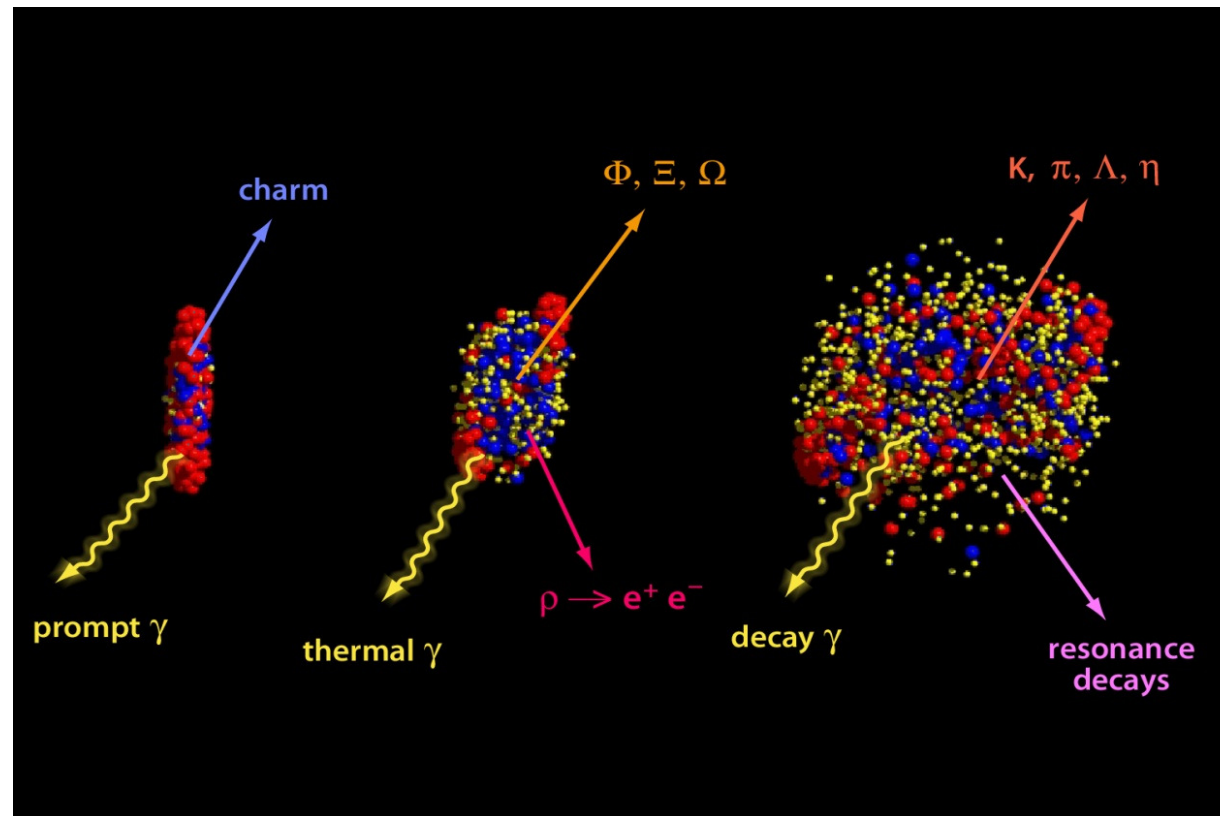
Open problems:

- How to describe a **first-order phase transition** in transport models?
- How to describe parton-hadron interactions in a **,mixed‘ phase**?

A possible solution (?):

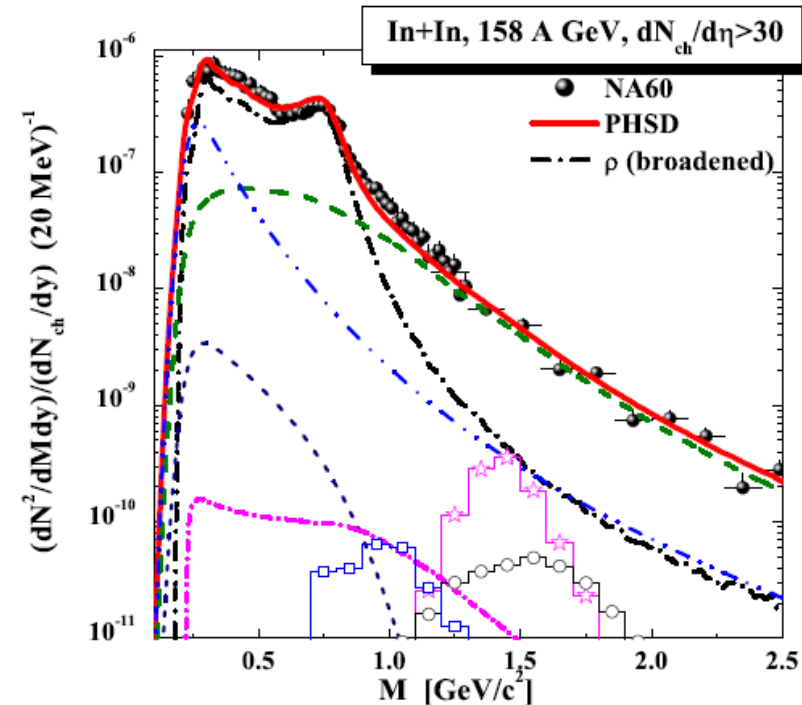
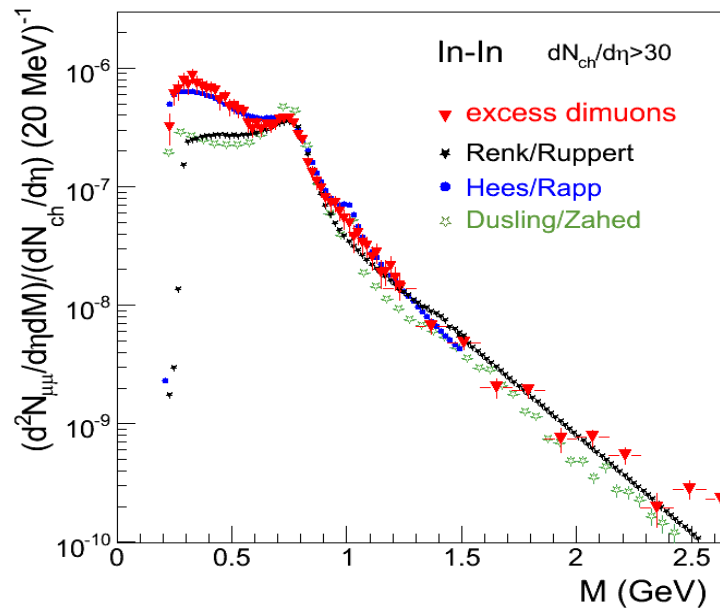
- **,mixed‘ phase description within molecular dynamics?**

Dileptons



are produced throughout all reaction phases!

Dileptons at SPS: NA60



Fireball model – Renk/Ruppert

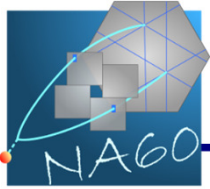
Fireball model – Rapp/vanHees

Hydro model – Dusling/Zahed

Parton-Hadron-String-Dynamics
microscopic transport model - PHSD

→ good agreement with data in
shape and absolute yield

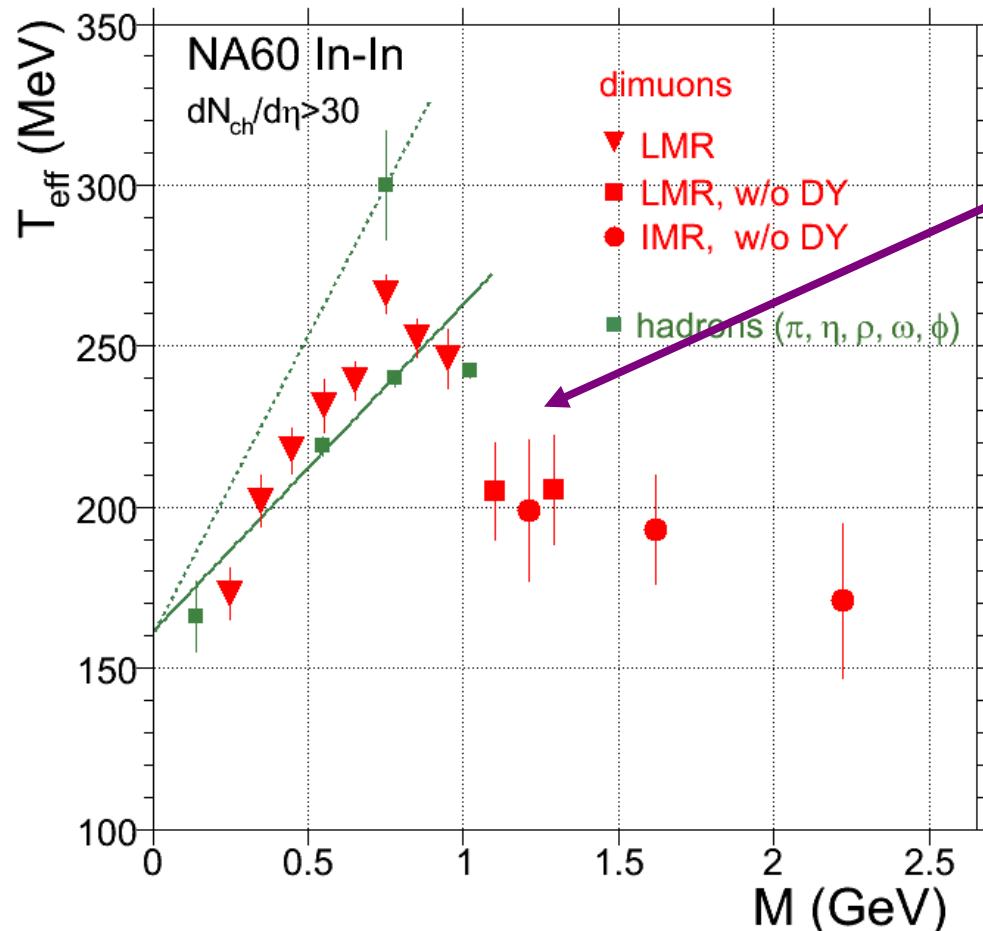
- Mass region above 1 GeV is dominated by partonic radiation
- Contributions of “4π” channels (radiation from multi-meson reactions) are small



The rise and fall of radial flow of thermal dimuons

▪ NA60: *Phys. Rev. Lett.* 100 (2008) 022302

▪ S. Damjanovic, Trento 2010:



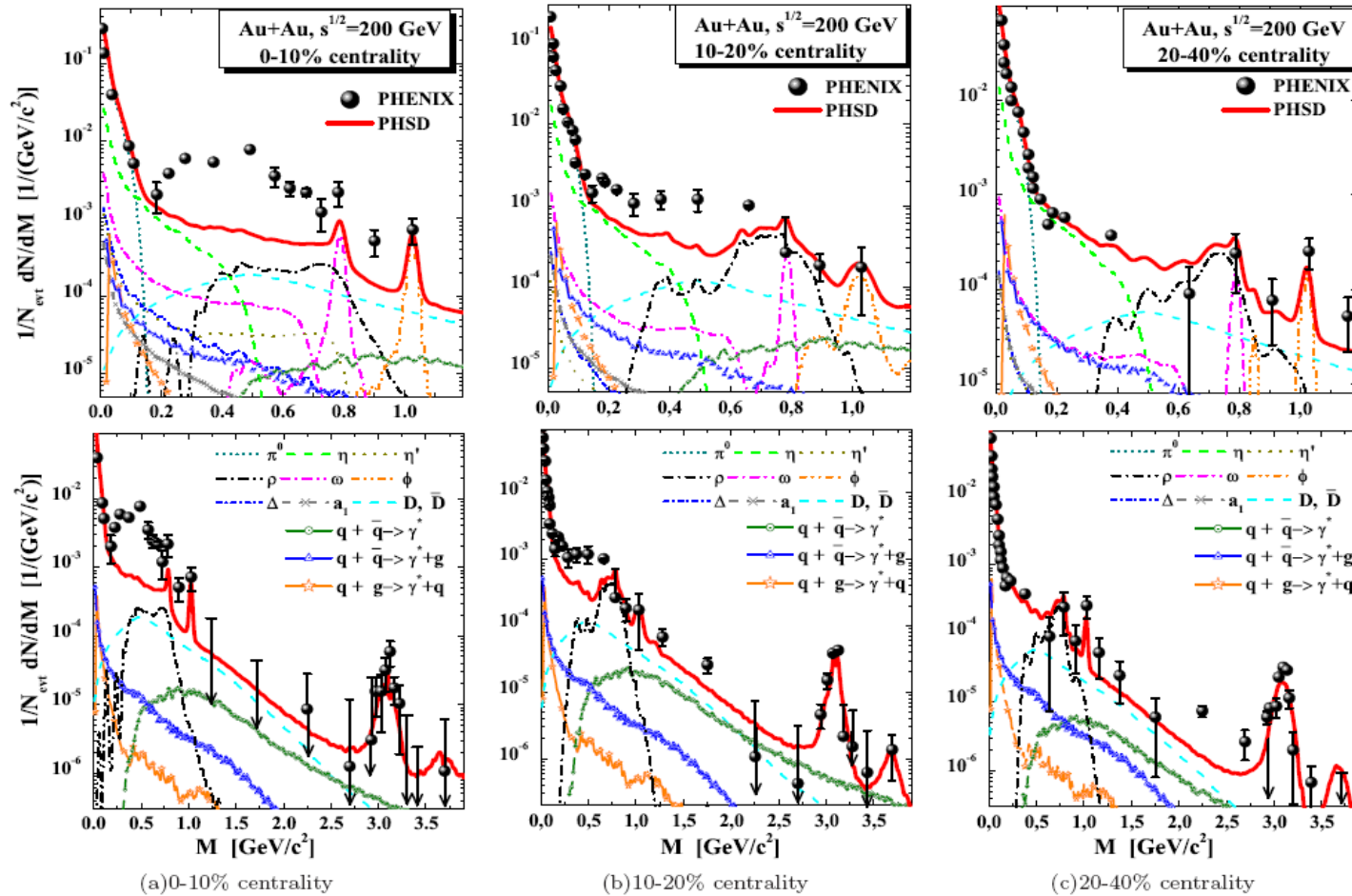
● Strong **rise** of T_{eff} with dimuon mass, followed by a sudden **drop** for $M > 1$ GeV:

■ **Rise** consistent with **radial flow of a hadronic source** (here $\pi^+\pi^- \rightarrow \rho \rightarrow \mu^+\mu^-$), taking the freeze-out \rangle as the reference (from a separate analysis of the ρ peak and the continuum)

■ **The drop of T_{eff}** signals a sudden transition to a **low-flow source**, i.e. a **source of partonic origin** (here $q\bar{q} \rightarrow \mu^+\mu^-$)

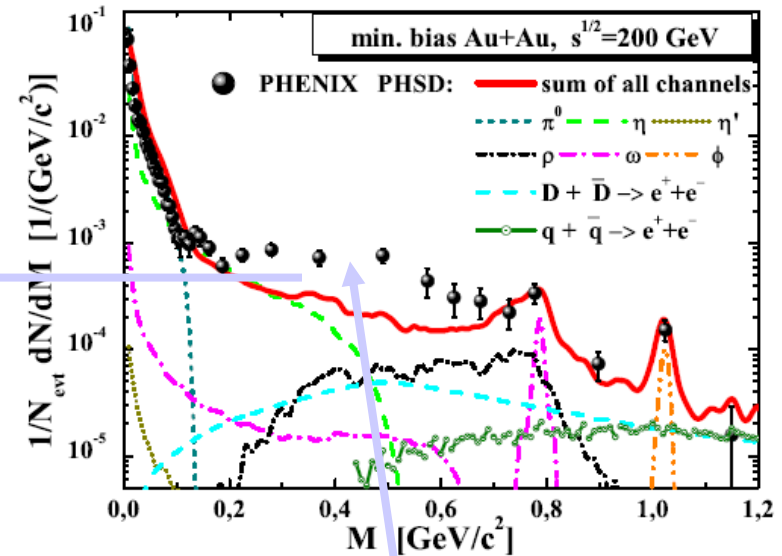
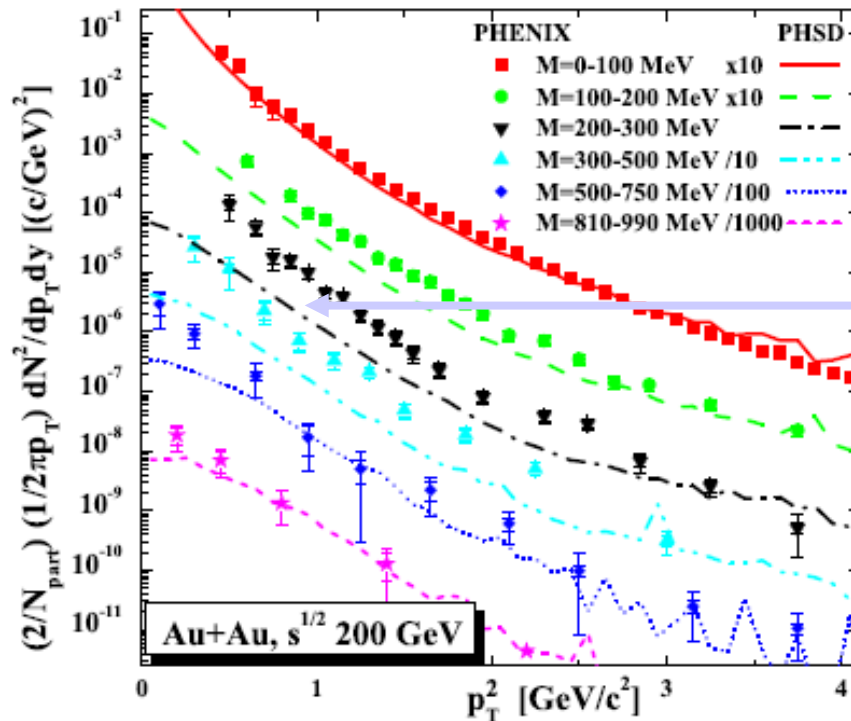
▪ NA60: Dileptons for $M > 1$ GeV are of dominantly partonic origin !

PHENIX: mass spectra



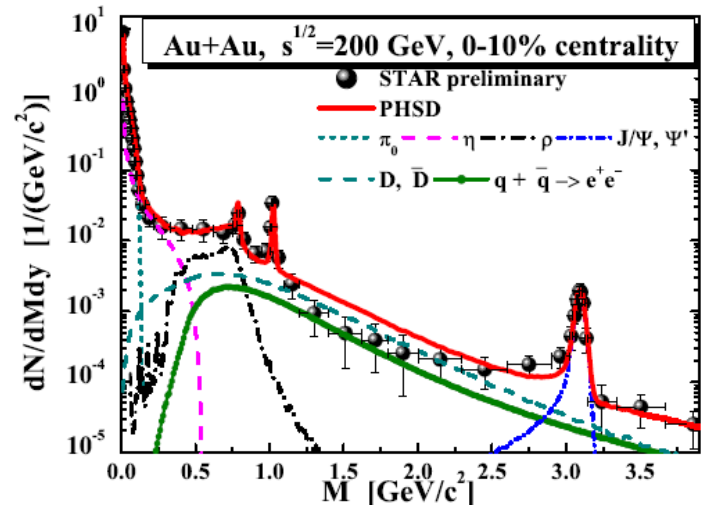
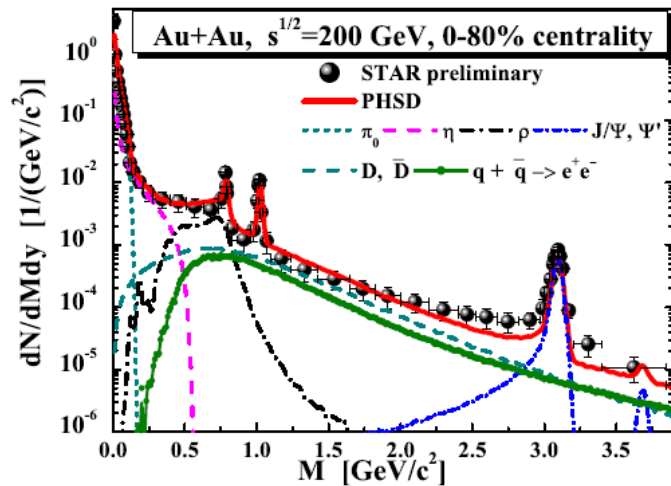
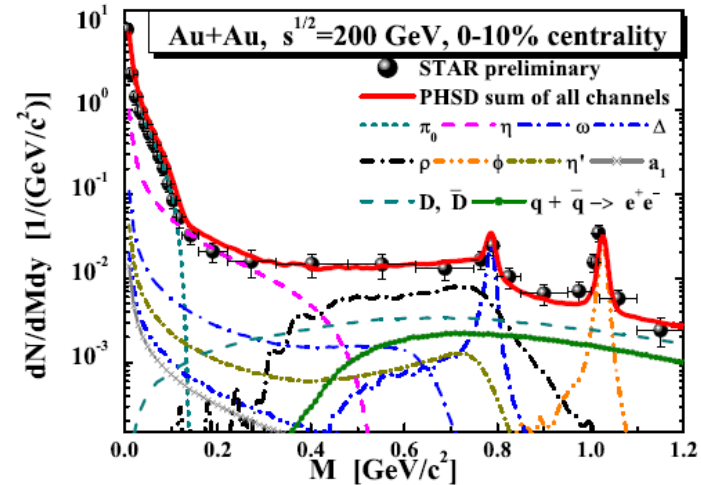
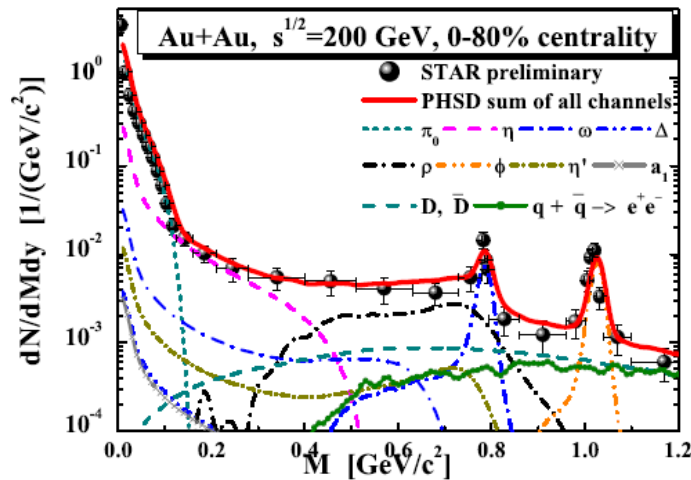
- Peripheral collisions (and pp) are well described, however, central fail!

PHENIX: p_T spectra



- The lowest and highest mass bins are described very well
- Underestimation of p_T data for $100 < M < 750$ MeV bins consistent with dN/dM
- The ‘missing source’(?) is located at low p_T !

STAR: dilepton mass spectra



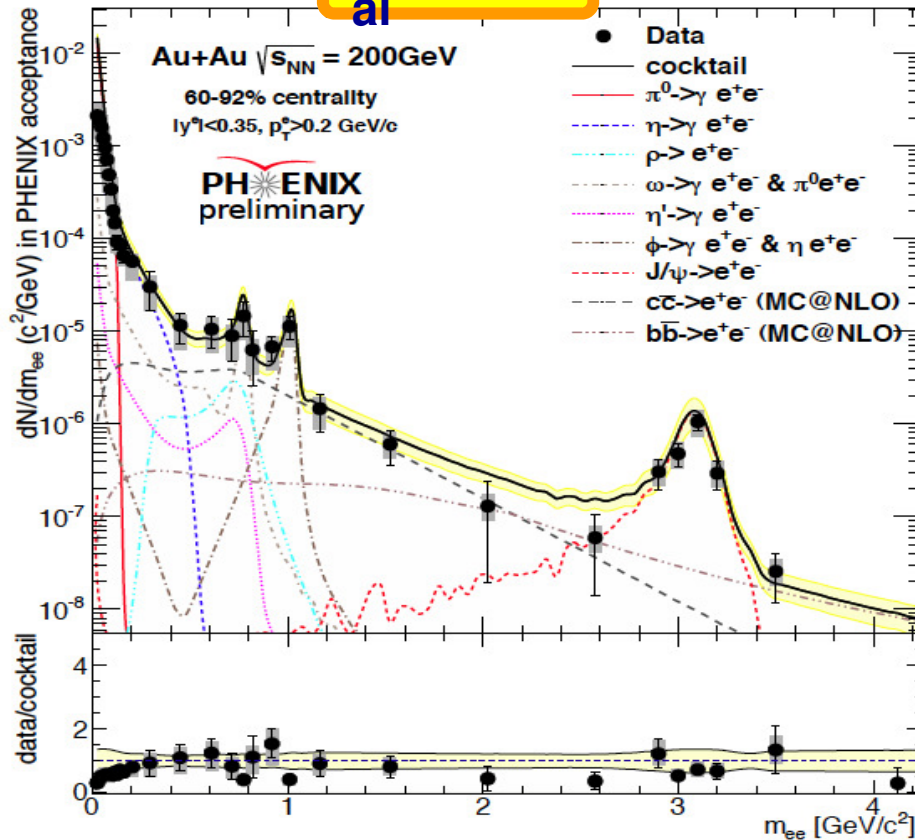
▪ STAR data are well described

▪ Confirmed by the extended data set at QM2012

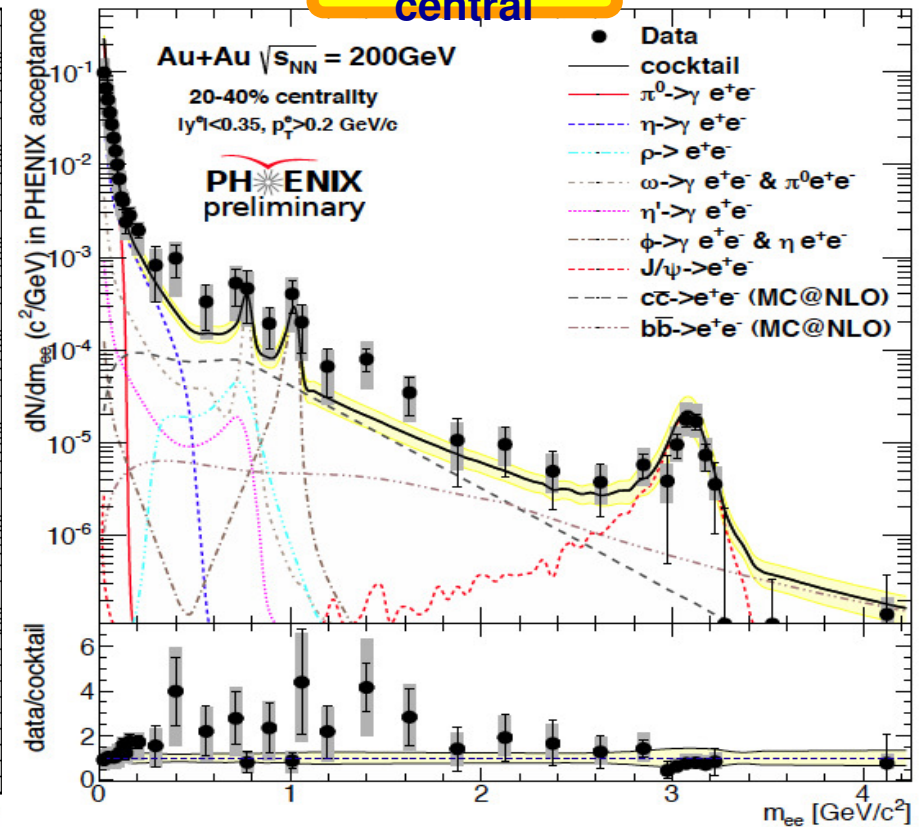
O. Linnyk, W. Cassing, J. Manninen, E. Bratkovskaya and C.M. Ko, PRC 85 (2012) 024910

PHENIX: mass spectra with HBD

Peripheral

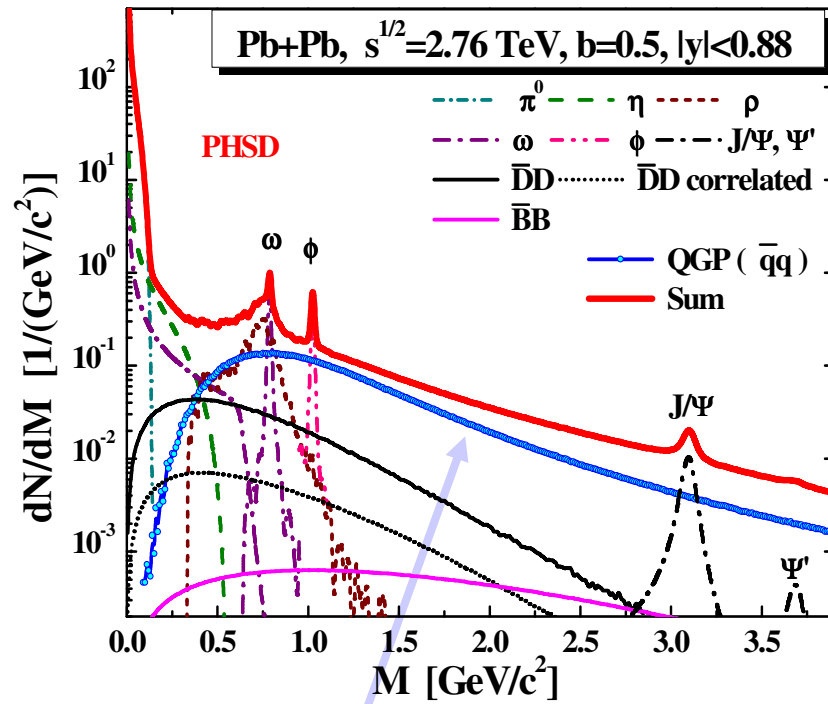


Semi-central

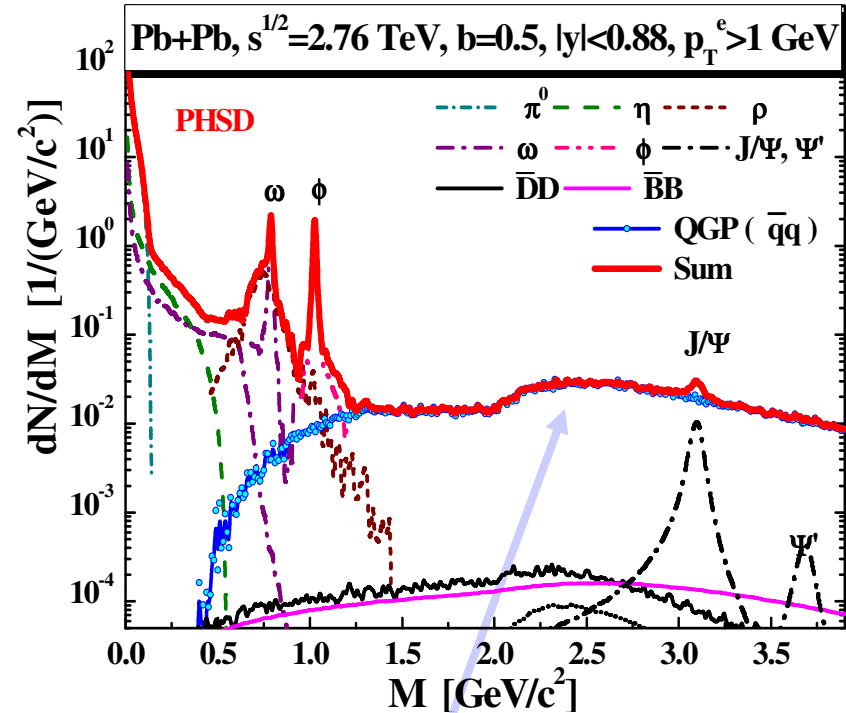


- Preliminary PHENIX data with Hadron-Blind Detector (HBD) presented at QM 2012
- Room for the QGP yield at $M > 1 \text{ GeV}$ and for the in-medium modification of rho at low mass

Predictions for LHC



QGP($\bar{q}q$) dominates at $M > 1.2$ GeV



p_T cut enhances the signal of QGP($\bar{q}q$)

D-, B-mesons energy loss from Pol-Bernard Gossiaux and Jörg Aichelin

JPsi and Psi' nuclear modification from Che-Ming Ko and Taesoo Song