Chiral Symmetry Breaking and Linear Confinement in a Dynamic Holography QCD Model

Danning Li



Institute of High Energy Physics, CAS

In collaboration with Mei Huang and Qi-Shu Yan

Erice, SEP, 23, 2012

Outlines

- I.A brief introduction to holography QCD
- II.Graviton-Dilaton system and pure gluon system
- III. Accommodate chiral symmetry breaking and linear confinement in a dynamic holography QCD model
- IV.Conclusion

I.A brief introduction to holography QCD

Why holography QCD?



Perturbation theory fails in both zero temperature case and finite temperature case

Non-perturbative methods: Lattice simulation, DSE, Effective field theory...

> Holographic QCD method: analytic understanding, easier calculation

AdS/CFT Correspondence

Juan Maldacena, Adv. Theor. Math. Phys. 2:231-252, 1998



Conjecture: Duality of the two theories

Equivalence of the partition function:

$$< e^{\int \phi_0^i O_i} >_{CFT} = Z_{string} [\phi^i(x, z \mid_{boundary}) = \phi_0^i(x)] \sim e^{-S_{Gravity}}$$

Towards Realistic Models

Further Conjecture

AdS/CFT Correspondence

General Gauge/Gravity

Task: find the proper non-conformal gravity background for QCD:

Top-Down Approach:

derive the background from fundamental string theory

Bottom-Up Approach:

find the background in a phenomenological way leave the origin of the background to the Top-down approach

II.Graviton-Dilaton system and pure gluon system

Confinement in QCD vacuum and deconfinement at finite temperature



Graviton-Dilaton System

$$S_{GD} = \frac{1}{2\kappa^2} \int dx^5 \sqrt{-g} e^{-2\phi} (R + 4\partial_m \phi \partial^m \phi + V_s(\phi))$$

Dynamical dilaton

deform the background metric: the warp factor

.

Zero temperature:

$$dS^{2} = g_{mn}dx^{m}dx^{n} = \frac{e^{2A_{s}}}{z^{2}}(-dt^{2} + dz^{2} + dx_{i}dx^{i})$$

Andreev-Zakharov

Quadratic warp factor

$$A_s = \frac{\sigma z^2}{2}$$



9

Study the finite T behavior: Black hole solution

$$dS^{2} = g_{mn}dx^{m}dx^{n} = \frac{e^{2A_{s}}}{z^{2}}(-f(z)dt^{2} + \frac{1}{f(z)}dz^{2} + dx_{i}dx^{i})$$
$$T = \frac{|f'(z_{h})|}{4\pi}.$$

 $A_s = \frac{\sigma z^2}{2}$

Dilaton profile and potential can be self-consistently solved from the equation of motion

Finite temperature behavior in the gravity background



D.N. Li, S. He, M.Huang, Q. S. Yan, arXiv:1103.5389, JHEP2011

O.Kaczmarek, F.Karsch, and F.Zantow, Phys.Rev.D.70.074505



EOS V.S. Quenched Lattice Results



 $T_c = 201 \text{MeV}$

D.N. Li, S. He, M.Huang, Q. S. Yan, arXiv:1103.5389, JHEP2011

EOS V.S. Quenched Lattice Results

$$c_s^2 = \frac{d\log T}{d\log s} = \frac{s}{Tds/dT},$$

Dots from G. Boyd et al., Nucl. Phys. B 469 (1996) 419



D.N. Li, S. He, M.Huang, Q. S. Yan, arXiv:1103.5389, JHEP2011

 At zero temperature and finite temperature, such a graviton-dilaton coupled system can describe the heavy quark potential behavior very well.

 The equation of states fits the quenched lattice data well (above Tc), which indicate that such a system can describe the pure gluon dynamics. III. Accommodate chiral symmetry breaking and linear confinement in a dynamic holography QCD model

Chiral symmetry breaking and confinement from meson spectra

Chiral symmetry breaking:

Chiral condensate spontaneously breaks the chiral symmetry, and provide quarks dynamical masses.

Mass split in chiral partner

Confinement:

Linear spectra in highly excited states



Data from C.Amsler et.al, Phys.Lett.B667(2008)1

KKSS model Karch, Katz, Son, Stephonov, Phys. Rev. D74:015005, 2006 $S_{kkss} = \int dx^5 \sqrt{-g} e^{-\phi} [|DX|^2 + 3|X|^2 - \frac{1}{4g_{\epsilon}^2} (F_L^2 + F_R^2)] \quad \phi = \mu^2 Z^2 \qquad m_{\rho,n}^2 = 4\mu^2 n$ Chiral symmetry is broken by m_n^2 , GeV² the expectation value of X p(2150) o(1900) p(1700) $X = \frac{\chi}{2} I_{2\times 2}, \quad \chi = m_q \ z + \sigma \ z^3 + o(z^4)$ 0(1450) + p(770) $V_{\rho} = \frac{A_{s}^{*} - \phi^{*}}{2} + \frac{(A_{s}^{*} - \phi^{*})^{2}}{4}$ 0.5 1.0 -0.1 $V_{a_1} = \frac{A_s^{"} - \phi^{"}}{2} + \frac{(A_s^{"} - \phi^{"})^2}{4} + g_5^2 \chi^2 e^{2A_s}$ -0.2 AdS metric -0.3Only coulomb part, -0.4 no linear part Chiral Symmetry breaking 17 -05

A successful holography QCD model

- should incoperate the linear potential and linear spectra (two aspects of confinement) simultaneously
- should realize chiral symmetry breaking and confinement simultaneously

Accommodate the two aspects at zero temperature

$$S = S_{GD} + \frac{N_f}{N_c} S_{KKSS}$$

$$N_f = 2, N_c = 3$$

Zero temperature

Dilaton profile

 $dS^{2} = e^{2A_{s}(z)}(-dt^{2} + dz^{2} + dx_{i}dx^{i})$

Cover chiral dynamics and

gluon dynamics

$$\phi = \mu^2 z^2$$

UV behavior of χ

$$\chi = m_q z + \sigma z^3 + o(z^4)$$

The Background Solution



D.Li, M.Huang, and Q.S.Yan, arxiv:1206.2824

Rho-a masses



D.Li, M.Huang, and Q.S.Yan, arxiv:1206.2824

Dots from C.Amsler et.al, Phys.Lett.B667(2008)1

pi-f0 masses



The quark antiquark potential

The linear potential is produced.

The slope of the linear part is: $\sigma_s \approx 4\mu^2$

$$\phi = \mu^2 z^2$$
$$m_{\phi}^2 L^2 = -4 \int_{V} \Delta_{\phi} = 2$$

Some dimension-2 operatror related to gluon dynamics ?

Need further study



D.Li, M.Huang, and Q.S.Yan, arxiv:1206.2824

IV.Conclusion

- The graviton-dilaton system describes the pure gluon dynamics above Tc quite well.
- Graviton-dilaton-scalar system can cover the gluon and chiral dynamics well. In this system, we accommodate chiral symmetry breaking, linear spectra and linear potential selfconsistently for the first time in holographic QCD frame. The slope of the linear potential and the linear spectra are the same and related to dilaton, which is consistent with the consideration that the two phenomena have the same origin---the gluon dynamics.
- Studying the confinement-deconfinement phase transition and chiral phase transition simultaneously is in progress.

Thanks for your attention !