

# The Need for an Early Anti-neutrino Run in NO $\nu$ A

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(Prakash, UR, Sankar, arXiv:1306.4125)



"ONE HUNDRED MILLION NEUTRINOS ARE PASSING THROUGH OUR BODIES EVERY SECOND AND WE'RE WORRIED ABOUT THE PRICE OF COFFEE."

[<http://www.sciencecartoonsplus.com/gallery/physics/galphys2h.php>]

# The Neutrino Mixing Parameters

- $|\nu_e\rangle, |\nu_\mu\rangle, |\nu_\tau\rangle$  are flavor states / interaction states
- $|\nu_i\rangle$ s are the mass eigen states / propagation states, where  $i=1, 2, 3$
- $|\nu_\alpha\rangle = U|\nu_i\rangle$
- $$U = \begin{bmatrix} c_{13}c_{12} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - s_{13}c_{12}c_{23}e^{i\delta} & -c_{12}s_{23} - s_{13}c_{23}s_{12}e^{i\delta} & c_{23}c_{13} \end{bmatrix}$$
- Oscillation probabilities depend on 3 mixing angles, two independent mass-squared differences  $\Delta_{21} = m_2^2 - m_1^2$ ,  $\Delta_{31} = m_3^2 - m_1^2$  and one CP violating phase  $\delta$ .

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# The Parameter Space Before Reactor Neutrino Results

( $bfp \pm 1\sigma$ )

- $\sin^2 \theta_{12} = 0.30 \pm 0.013$
- $\sin^2 \theta_{23} = 0.41_{-0.025}^{+0.037} \oplus 0.59_{-0.022}^{+0.021}$
- $\Delta_{21} = (7.50 \pm 0.185) \times 10^{-5} eV^2$
- $|\Delta_{31}| = (2.47_{-0.067}^{+0.069}) \times 10^{-3} eV^2$

[Table 1 of Gonzalez-Garcia, Maltoni, Salvado, Schwetz, arXiv:1209.3023]

- $\alpha = \frac{\Delta_{21}}{\Delta_{31}}$  and  $\theta_{13}$  very small
- $P_{ee}$  and  $P_{\mu\mu}$  generally can be written in effective two-flavor form.  
 $P_{\bar{e}\bar{e}} = 1 - \sin^2 2\theta_{12} \sin^2 \frac{\Delta_{21}L}{4E}$  (KamLAND)  
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# $|\nu_\mu\rangle \rightarrow |\nu_e\rangle$ Oscillation with Matter Effect in Long Baseline Experiments



$$\begin{aligned} P_{\mu e} &= \sin^2 2\theta_{13} \sin^2 \theta_{23} \frac{\sin^2 \hat{\Delta}(1 - \hat{A})}{(1 - \hat{A})^2} \\ &+ \alpha \cos \theta_{13} \sin 2\theta_{12} \sin 2\theta_{13} \sin 2\theta_{23} \cos(\hat{\Delta} + \delta_{CP}) \\ &\quad \frac{\sin \hat{\Delta} \hat{A}}{\hat{A}} \frac{\sin \hat{\Delta}(1 - \hat{A})}{1 - \hat{A}} \\ &+ \alpha^2 \sin^2 2\theta_{12} \cos^2 \theta_{13} \cos^2 \theta_{23} \frac{\sin^2 \hat{\Delta} \hat{A}}{\hat{A}^2} \end{aligned} \quad (0.1)$$

[Cervera et al., arXiv: hep-ph/0002108]

- $\hat{\Delta} = \Delta_{31}L/4E$ ,  $\hat{A} = A/\Delta_{31}$ ,  $\alpha = \Delta_{21}/\Delta_{31}$ ,  $A = 0.000076\rho E$  has energy dependence and hence can create matter effect.

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# $|\nu_\mu\rangle \rightarrow |\nu_e\rangle$ Oscillation with Matter Effect in Long Baseline Experiments

- $\Delta_{31}$  +ve for NH and -ve for IH
- A +ve for  $\nu$  and -ve for  $\bar{\nu}$
- For  $\nu$ ,  $\hat{A}$  +ve for NH and -ve for IH
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- $P_{\mu e}$  SENSITIVE to hierarchy
- $P_{\mu e}$  dependent of  $\theta_{13}$ , hierarchy, octant of  $\theta_{23}$ ,  $\delta_{CP} \rightarrow$  EIGHT-FOLD DEGENERACY

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# Eight -fold Degeneracy

[Barger et al., arXiv: hep-ph/0112119]

- The (hierarchy- $\delta_{CP}$ ) degeneracy:  $P_{\mu e}(\theta_{13}, \text{NH}, \delta_{CP})(\theta_{13}, \text{NH}, \delta_{CP}) = P_{\mu e}(\theta_{13}, \text{IH}, \delta_{CP}')$
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# Recent Reactor Neutrino Experiments

- $\sin^2 2\theta_{13} = 0.089$   
[An et al., arXiv: 1203.1669; Ahn et al., arXiv: 1204.0626; Abe et al., arXiv: 1207.6632]
- $\sin^2 \theta_{13} = 0.023 \pm 0.0023$   
[Table 1 of Gonzalez-Garcia, Maltoni, Salvado, Schwetz, arXiv:1209.3023]
- Daya Bay will reduce the error in  $\sin^2 2\theta_{13}$  from 10% to 5% at the end of its running.  
[Dwyer (Daya Bay collaboration), talk given at the Neutrino 2012 conference, June 3-9, 2012, Kyoto, japan, <http://neu2012.kek.jp>]
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- $\sin 2\theta_{13} \simeq 0.3$ , 10 times to  $\alpha = 0.03$ , neglect  $\alpha^2$

# Recent Reactor Neutrino Experiments

$$\begin{aligned} P_{\mu e} &= \sin^2 2\theta_{13} \sin^2 \theta_{23} \frac{\sin^2 \hat{\Delta}(1 - \hat{A})}{(1 - \hat{A})^2} \\ &+ \alpha \cos \theta_{13} \sin 2\theta_{12} \sin 2\theta_{13} \sin 2\theta_{23} \cos(\hat{\Delta} + \delta_{CP}) \\ &\quad \frac{\sin \hat{\Delta} \hat{A}}{\hat{A}} \frac{\sin \hat{\Delta}(1 - \hat{A})}{1 - \hat{A}} \end{aligned} \tag{0.2}$$

# The Unknown Ones

- $\delta_{CP}$ , Sign of  $\Delta_{31}$ , octant of  $\theta_{23}$
- Reactor experiments result on  $\theta_{13}$  allows mass hierarchy to be determined by current experiment  
[An et al., arXiv: 1203.1669; Ahn et al., arXiv: 1204.0626; Abe et al., arXiv: 1207.6632]
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# The Hierarchy- $\delta_{CP}$ Degeneracy

[Barger et al., arXiv: hep-ph/0112119; Mena, Parke, arXiv: hep-ph/040870]



$$P_{\mu e} = \sin^2 2\theta_{13} \sin^2 \theta_{23} \frac{\sin^2 \hat{\Delta}(1 - \hat{A})}{(1 - \hat{A})^2} + \alpha \cos \theta_{13} \sin 2\theta_{12} \sin 2\theta_{13} \sin 2\theta_{23} \cos(\hat{\Delta} + \delta_{CP}) \frac{\sin \hat{\Delta} \hat{A}}{\hat{A}} \frac{\sin \hat{\Delta}(1 - \hat{A})}{1 - \hat{A}} \quad (0.3)$$

- $P_{\mu e}$  (NH) >  $P_{\mu e}$  (IH), for  $\nu$ : consequences of  $\hat{A}$  dependence
- At oscillation maxima  $\hat{\Delta} \simeq 90^\circ$
- $\cos(\hat{\Delta} + \delta_{CP})$  is 1 for  $\delta_{CP} = -90^\circ$  and -1 for  $\delta_{CP} = 90^\circ$

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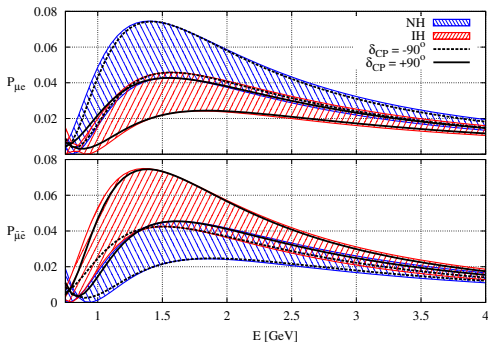
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# The Hierarchy- $\delta_{CP}$ Degeneracy in $\text{NO}\nu\text{A}$



**Figure:**  $P_{\mu e}$  (top panel) and  $P_{\bar{\mu} e}$  (bottom panel) vs. energy for  $\text{NO}\nu\text{A}$ . Variation of  $\delta_{CP}$  leads to the blue (red) bands for NH (IH). The plots are drawn for maximal  $\theta_{23}$  and other neutrino parameters given as the central value in slide 3.

[Prakash, Raut, Sankar, arXiv: 1201.6485v3]

# Favorable and Unfavorable Combinations

- $P_{\mu e}(NH, -180^\circ < \delta_{CP} < 0) > P_{\mu e}(IH, \text{any } \delta_{CP})$
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DEGENERATE solutions- (true hierarchy,  $\delta_{CP}$ ) and (wrong hierarchy,  $\delta_{CP}'$ ) for each measurement, no hierarchy determination



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# Favorable and Unfavorable Combinations

- (NH,  $\delta_{CP}$  in LHP) and (IH,  $\delta_{CP}$  in UHP) are **favorable** combinations for hierarchy determination in  $NO\nu A$ .
- (NH,  $\delta_{CP}$  in UHP) and (IH,  $\delta_{CP}$  in LHP) are **unfavorable** combinations for hierarchy determination in  $NO\nu A$ .
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# Potential for $\text{NO}\nu\text{A}$

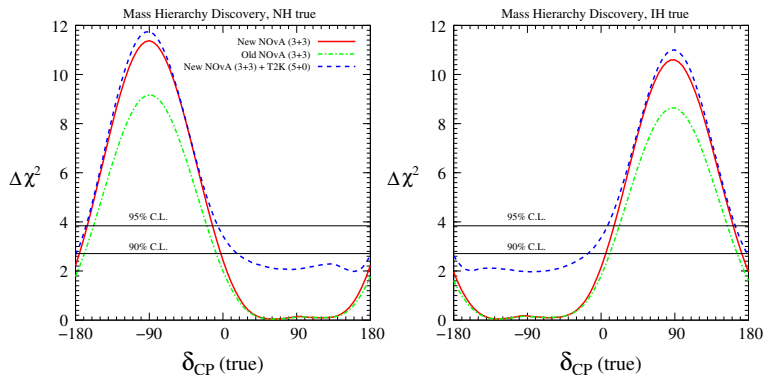


Figure: Hierarchy sensitivity for  $\text{NO}\nu\text{A}$  after complete run. In the left (right) panel, the true hierarchy is taken to be NH (IH).

[Agarwalla, Prakash, Raut, Sankar, arXiv: 1208.3644v2]

What can we learn in first 3 years of  $\text{NO}\nu\text{A}$ , if we have **favorable** hierarchy- $\delta_{CP}$  combinations as true combinations?

Consider two possibilities.

- 3 year of  $\nu$  run
- 1.5 year  $\nu$  run + 1.5 year  $\bar{\nu}$  run

WHY?

Originally first 3 year  $\nu$  run was considered to discover non-zero  $\theta_{13}$ , in case it was small. But now  $\theta_{13} \simeq 8^\circ$ , other possibilities can be considered.

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## NO $\nu$ A Experiment

[Ayres et al., NO $\nu$ A, Tech. Rep. (2007), Fermilab-Design-2007-01]

- 14 kiloton T ASD
- 810 km away from Fermilab
- Detector locaton: 0.8° off axis from the NuMI beam
- $\nu$  flux peaks sharply at 2 GeV, oscillation maximum energy 1.5 GeV
- Equal  $\nu$  and  $\bar{\nu}$  run of 3 years each
- NuMI beam power 700 kW, corresponding to  $6 \times 10^{20}$  protons on target per year
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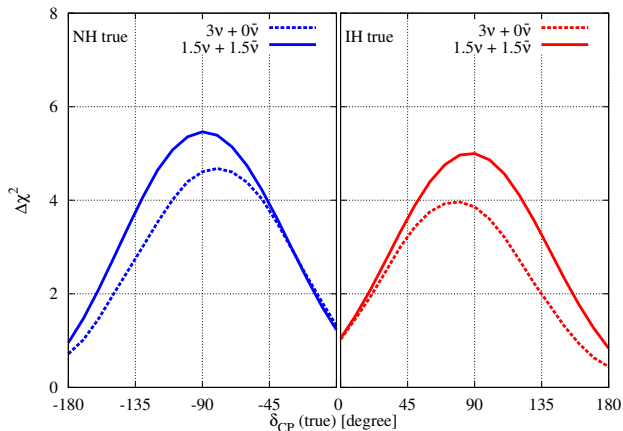
Event number simulations and the  $\Delta\chi^2$  calculations are done by using **GLOBES**

[Huber et al., arXiv: hep-ph/0407333, Huber et al., arXiv: hep-ph/0701187]

Minimum  $\Delta\chi^2$  is calculated by doing a **marginalization** over the above mentioned parameters.



# Effect of Precision of $\sin^2 2\theta_{13}$



**Figure:** Hierarchy sensitivity assuming 10% uncertainty in  $\sin^2 2\theta_{13}$  and maximal  $\theta_{23}$ . In the left (right) panel, the true hierarchy is taken to be NH (IH).

# Effect of Precision of $\sin^2 2\theta_{13}$

10% uncertainty in  $\sin^2 2\theta_{13}$

- $2\sigma$  hierarchy determination is possible in 1.5 year  $\nu$  run + 1.5 year  $\bar{\nu}$  for about 50% of the favorable half plane.
- $2\sigma$  hierarchy determination is possible over a very small range (no  $\delta_{CP}$ ) if NH,LHP (IH,UHP) is true in  $3\nu$  run

WHY?

- The lower sensitivity of  $3\nu$  run is due to hierarchy- $\sin^2 2\theta_{13}$  degeneracy.
- If the value of  $\sin^2 2\theta_{13}$  is not precisely known it is possible to fake the probability value with a wrong hierarchy and with a wrong  $\sin^2 2\theta_{13}$ .
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- Pure  $\nu$  is sensitive to this degeneracy but a combination of  $\nu$  and  $\bar{\nu}$  not.
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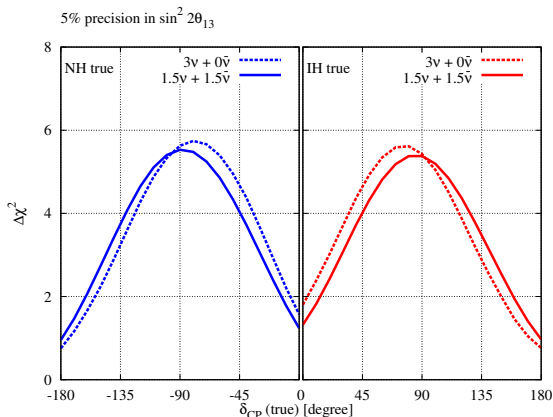
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# Effect of Precision of $\sin^2 2\theta_{13}$

If the uncertainty in  $\sin^2 2\theta_{13}$  is reduced to 5%, the hierarchy reach for  $3\nu$  does improve and becomes equal to that of  $1.5\nu + 1.5\bar{\nu}$  run.



**Figure:** Hierarchy sensitivity assuming 5% uncertainty in  $\sin^2 2\theta_{13}$  and maximal  $\theta_{23}$ . In the left (right) panel, the true hierarchy is taken to be NH (IH).

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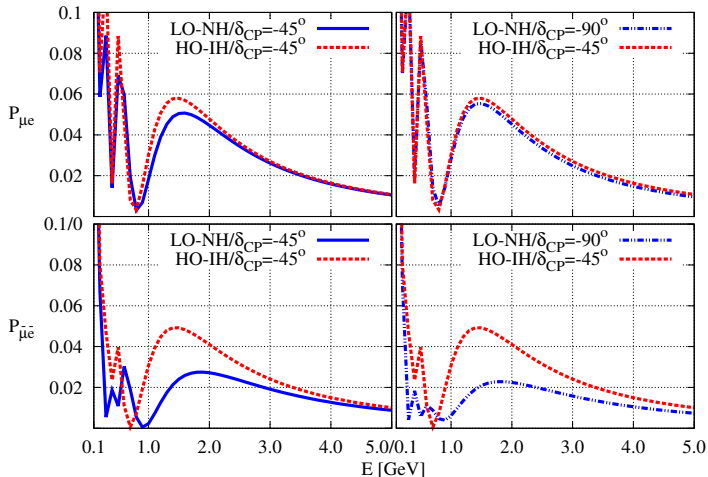
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**Figure:** Illustration of degenerate  $P_{\mu e}$  and non-degenerate  $P_{\bar{\mu} \bar{e}}$  for the following two cases. Left: (LO-NH,  $\delta_{CP} = -45^\circ$ ) and (HO-IH,  $\delta_{CP}' = -45^\circ$ ), Right: (LO-NH,  $\delta_{CP} = -90^\circ$ ) and (HO-IH,  $\delta_{CP}' = -45^\circ$ ).



- Dominant term in  $P_{\mu e}$  proportional to  $\sin^2 2\theta_{13} \sin^2 \theta_{23}$
- Matter effect in NH makes it larger and HO makes it even larger.
- $P_{\mu e}(\text{HO-NH}, \delta_{CP} \text{ in LHP}) \gg P_{\mu e}(\text{IH})$ , for any values of neutrino parameters  $\Rightarrow$  statistics for HO-NH will be quite large.
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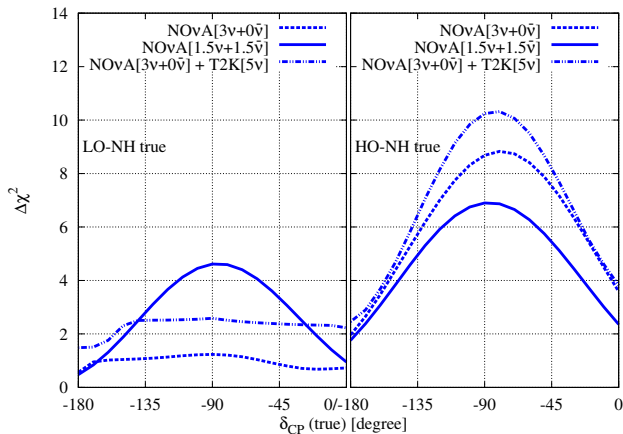


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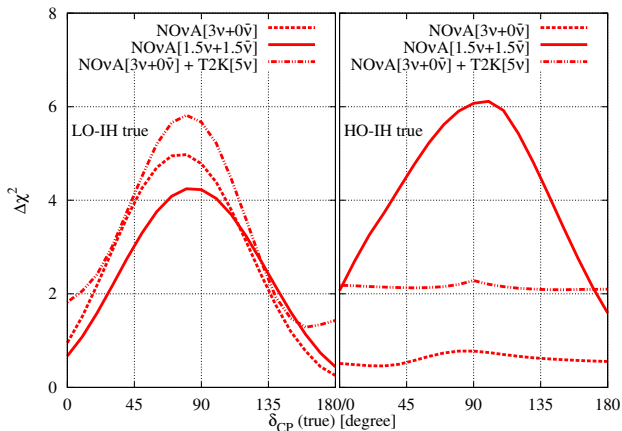
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$\nu$  and  $\bar{\nu}$  data have **different** dependence on hierarchy-octant degeneracy  $\Rightarrow$  a combination of **both** will help to solve this degeneracy in case of **LO-NH** and **HO-IH**.



**Figure:** Hierarchy sensitivity assuming 5% uncertainty in  $\sin^2 2\theta_{13}$  for NH and LHP. In the left (right) panel, the true  $\sin^2 \theta_{23}$  is taken to be 0.41 (0.59).

# Non-maximal $\theta_{23}$



**Figure:** Hierarchy sensitivity assuming 5% uncertainty in  $\sin^2 2\theta_{13}$  for IH and UHP. In the left (right) panel, the true  $\sin^2 \theta_{23}$  is taken to be 0.41 (0.59).

- **HO-NH** combination has a  $2\sigma$  hierarchy discrimination for **80%** (**70%**) of the favorable half plane for  $3\nu$  ( $1.5\nu + 1.5\bar{\nu}$ )
- **LO-IH** combination has a  $2\sigma$  hierarchy discrimination for 40% (20%) of the favorable half plane for  $3\nu$  ( $1.5\nu + 1.5\bar{\nu}$ )
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- After Daya Bay reduces  $\sigma(\sin^2 2\theta_{13})$  to  $5\%$ ,  $1.5\nu + 1.5\bar{\nu}$  has  $2\sigma$  hierarchy sensitivity in all possible hierarchy-octant combinations for a satisfying range of  $\delta_{CP}$  in favorable half plane, where as  $3\nu$  has no sensitivity at all for LO-NH and HO-IH
- Addition of T2K data does not help much, because we are confined in favorable region.
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# Re-optimizing $\text{NO}\nu\text{A}$

- Relaxed energy cut allowing more signal events  $\Rightarrow$  more backgrounds as well
- Signal efficiencies in new  $\text{NO}\nu\text{A}$  (45% for both neutrino and anti-neutrino) is 2 times the old one (26% for neutrino and 40% for anti-neutrino) for neutrino
- NC background  $\Rightarrow$  about 7 times (2% vs 0.3%) for  $\nu$  and 3 times (3% vs 0.9%) for  $\bar{\nu}$  than the old one.
- Misidentified muons  $\Rightarrow$  6 times (0.83% vs 0.13%) for  $\nu$  and 2 times (0.22% vs 0.13%) than the old one
- Earlier the number of NC background events were moderate. For that case assuming a Gaussian energy resolution function to smear the background events was a good approximation. But at present the NC backgrounds are higher by a factor 5. But their measured energy range in is much below the range of large flux. The NC spectrum shift to measured energies is implemented through migration matrices.



- We have considered the neutrino contamination in the anti-neutrino beam in both appearance and disappearance channels. While the anti- neutrino contamination in the neutrino beam can be ignored, the reverse is not true.
- The above optimization criteria was developed in the case of  $\nu_\mu \longrightarrow \nu_e$  vacuum oscillation with  $\delta_{cp} = 0$  and maximizing signal events while keeping the background events relatively small.