

Problems with single pion production

International School of Nuclear Physics 35th Course Neutrino Physics: Present and Future Erice-Sicily

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Jakub Żmuda Neutrino interactions...

Motivation





 Single pion production (SPP): second to quasielastic.



• T2K and other: huge no oscillation/ best fit difference $\rightarrow \Theta_{ij}$, Δm_{ij}^2 , ΔCP ?, ...?

Predictions of # events with/without oscillation

Based on Monte Carlo. What you put is what you get \rightarrow dependency on theoretical lepton-nucleus interaction modeling.

Single pion production



• At T2K mean energy $\nu_{\mu}p \rightarrow \mu^{-}p\pi^{+}$ channel: intermediate $\Delta(1232)$ baryon. $\Gamma^{\mu\alpha} = \begin{bmatrix} \frac{C_{3}^{V}}{M} (g^{\alpha\mu}q - q^{\alpha}\gamma^{\mu}) + \frac{C_{4}^{V}}{M^{2}} (g^{\alpha\mu}q \cdot p_{\Delta} - q^{\alpha}p_{\Delta}^{\mu}) + \frac{C_{5}^{V}}{M^{2}} (g^{\alpha\mu}q \cdot p - q^{\alpha}p^{\mu}) + g^{\alpha\mu}C_{6}^{V} \end{bmatrix} \gamma^{5} + \begin{bmatrix} \frac{C_{3}^{A}}{M} (g^{\alpha\mu}q - q^{\alpha}\gamma^{\mu}) + \frac{C_{4}^{A}}{M^{2}} (g^{\alpha\mu}q \cdot p_{\Delta} - q^{\alpha}p_{\Delta}^{\mu}) + C_{5}^{A}g^{\alpha\mu} + \frac{C_{6}^{A}}{M^{2}}q^{\alpha}q^{\mu} \end{bmatrix}$

• Vector part: (rather) well-known from pion photo- and electroproduction data.

• Axial part: dominated by
$$C_5^A(Q^2) = \frac{C_5^A(0)}{(1+Q^2/M_{A\Delta}^2)^2}$$
.
 $C_{3/4}^A$ - Adler's relations, C_6^A -PCAC. $C_5^A(0) \approx 1.2 \propto f_{\pi N\Delta}$
(Goldberger-Treiman). $M_{A\Delta}$: fits to ANL/BNL data.



The HNV model

- Other isospin channels: large nonresonant background (but in $\nu_{\mu}p \rightarrow \mu^{-}p\pi^{+}$ also non-negligible).
- Hernandez, Nieves, Valverde (HNV) model from Phys. Rev. D 76, 033005 (2007):



- Alltogether 7 currents: 2 from Δ resonance (a) and b)), rest from chiral effective field theory.
- Fit to ANL/BNL data with background: E. Hernandez et al. Phys. Rev. D 81, 085046 (2010): $C_5^A(0) = 1.0 \pm 0.11$, $M_{A\Delta} = 0.93 \pm 0.07$ GeV.



The HNV model

• Amplitudes in HNV $\mathcal{J}^{\mu}_{hadr.} = \langle N'\pi \, | s^{\mu} | \, N \rangle$. Example:

$$\begin{split} s^{\mu}_{\Delta P} &= iC_{\Delta P}\frac{f^{*}}{m_{\pi}}k^{\alpha}G_{\alpha\beta}(p+q)\Gamma^{\beta\mu}(p,q) \\ s^{\mu}_{NP} &= -iC_{NP}\frac{g_{A}}{\sqrt{2}f_{\pi}}k^{\prime}\gamma^{5}\frac{(p^{\prime}+q^{\prime}+M)}{(p+q)^{2}-M^{2}+i\epsilon}\Gamma^{\mu}(q)F_{\pi}(k-q) \end{split}$$

• $\Gamma^{\beta\mu}$: Δ excitation vertex, Γ^{μ} -nucleon weak vertex, $G_{\alpha\beta}$: Rarita-Schwinger propagator, $\Gamma_{\Delta}(p_{\Delta}^2)$ - $\Delta \rightarrow \pi N$ decay width.

$$G_{\alpha\beta}(p_{\Delta}) = \frac{-(\not\!\!\!p_{\Delta} + M_{\Delta})}{p_{\Delta}^2 - M_{\Delta}^2 + iM_{\Delta}\Gamma_{\Delta}(p_{\Delta}^2)} \left(g_{\alpha\beta} - \frac{1}{3}\gamma^{\alpha}\gamma^{\beta} - \frac{2}{3}\frac{p_{\Delta}^{\alpha}p_{\Delta}^{\beta}}{M_{\Delta}^2} + \frac{1}{3}\frac{p_{\Delta}^{\alpha}\gamma^{\beta} - p_{\Delta}^{\beta}\gamma^{\alpha}}{M_{\Delta}}\right)$$

• High complexity- Dirac and Lorentz algebra in numerical C++ code.



Tests: inclusive electron data

- O. Lalakulich, E. A. Paschos, and G. Piranishvili, Phys. Rev. D , 014009 (2006) electromagnetic Delta form factors.
- Test: inclusive electron-proton scattering data (J. S. O'Connell *et al.*, Phys. Rev. Lett. 53, 1627 (1984) (left), M. Christy and P. E. Bosted, Phys. Rev. C 81, 055213 (2010) (right)).



- Background contribution in interesting energy range (730 MeV)-large.
- Some imperfections for low energies. Above $W\approx 1.25$ GeV- 2π channel, then heavy resonances and DIS



• Double-differential cross section w.r.t. final lepton variables (angle Ω' and energy E'):

$$\frac{d^3\sigma}{d\Omega' dE'} = \frac{G_F^2 \cos^2 \theta_C}{4\pi^2} \frac{|\vec{l'}|}{\sqrt{(l \cdot P_i)^2}} \left(-\frac{1}{\pi} \Im L_{\mu\nu} \Pi^{\mu\nu} \right)$$
$$L_{\mu\nu} = \frac{1}{8} \text{Tr} [\ell' + m_x) \gamma_\mu (1 \mp \gamma^5) (\ell + m_{\nu_x}) \gamma_\nu (1 \mp \gamma^5)]$$

• Information about nuclear system excitations in polarization tensor, $\Pi^{\mu\nu}$:

$$\Pi^{\mu\nu}(q) \equiv i\Omega \int d^4x e^{iqx} \sum_i \left\langle i \left| T \left\{ \hat{J}_I^{\nu\dagger}(x) \hat{J}_I^{\mu}(0) \exp\left(-i \int d^4y \hat{H}_{int.}(y)\right) \right\} \right| i \right\rangle E_{ij}$$

- \hat{J}^{μ} nuclear current operator, $\hat{H}_{int.}$ interaction Hamiltonian (here-from chiral field theory).
- $\Pi^{\mu\nu}$: gauge boson self-energy in nuclear matter.



SPP on atomic nuclei

 Single pion production contribution to polarization tensor (following IFIC group Phys. Rev. C 83, 045501 (2011)):





Black circle- sum of 7 HNV diagrams

- Nuclear effects in local density approximation:
 - Fermi motion: local Fermi gas
 - Pauli blocking
 - Δ self-energy in nuclear matter Σ_{Δ} (*npnh* excitations, part of MEC) $\Gamma_{\Delta} \rightarrow \Gamma_{\Delta}^{PB} - 2\Im \Sigma_{\Delta}$, parameterization from E. Oset Nucl. Phys. A 468, 631 (1987).



SPP on atomic nuclei

• Basic formula with sum over isospins N, N', t_{π} , integration over nuclear volume d^3r , nucleon d^3p and pion d^3k momenta:

$$\begin{split} \frac{d^3\sigma}{d\Omega' dE'} &= \frac{G_F^2 \cos^2 \theta_C}{4\pi^2} \frac{|\vec{l}'|}{E} \sum_{N,N',t_\pi} \int\!\!\!d^3r \int\!\!\frac{d^3p}{(2\pi)^3} \int\!\!\frac{d^3k}{(2\pi)^3} \frac{1}{8E_\pi(k)E(p)E(p')} \\ & \left[n_N(p)(1\!-\!n_{N'}(p')) + n_{N'}(p')(1\!-\!n_N(p)) \right] \\ & \delta(E(p') - \tilde{q}^0 \!+\!E_\pi \!-\!E(p)) \text{Tr} \Big[A_{1p1h1\pi}^{\mu\nu}(p,q,k) \Big] \, L_{\mu\nu}. \end{split}$$

•
$$A_{1p1h1\pi}^{\mu\nu} \equiv \operatorname{Tr} \left[\mathcal{J}_{hadr.}^{\nu} \mathcal{J}_{hadr.}^{\mu*} \right], \ n_N(p) = \Theta(k_F^N(r) - p).$$

Separate piece for Δ-in-medium excitations:

$$\frac{d^{3}\sigma}{dE'd\Omega'} \approx \frac{G_{F}^{2}\cos^{2}(\Theta_{C})|\boldsymbol{l}'|}{16\pi^{5}|\boldsymbol{l}|} \int drr^{2} \sum_{iso} C_{iso} \int d^{3}p \frac{n_{N}(p)(\frac{1}{2}\tilde{\Gamma}-\Im\Sigma_{\Delta})}{E(p)(M_{\Delta}+W)} \cdot \\
\cdot \frac{\operatorname{Tr}\left[\gamma^{0}\Gamma^{\alpha\mu^{\dagger}}\gamma^{0}P_{\alpha\beta}^{3/2}(p_{\Delta})\Gamma^{\beta\nu}(\not{p}+M)\right]L_{\mu\nu}}{(W-(M_{\Delta}+\Re\Sigma_{\Delta}))^{2}+(\frac{1}{2}\tilde{\Gamma}-\Im\Sigma_{\Delta})^{2}}.$$
(1)

• Problem: total $\sigma \Delta$ +background integration in 8 dimensions.

Set $< |\vec{p}| > \text{in } A^{\mu\nu} \rightarrow \text{integral factorization (literal Phys. Rev. C 83, 045501 (2011) solution)}$

2 Monte Carlo integration with Vegas algorithm from GSL → "exact" solution (J. Sobczyk and J. Żmuda PRC87, 065503 (2013)).



"Approximate" and "exact" integration

• Difference for double-differential cross sections (rather large):



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Comparison with MiniBooNE

- MiniBooNE: Phys. Rev. D83, 052007 and 052009 (2011). All 1π final state included!
- Strong final state interactions, no MC simulation yet. Approximate ¹⁶O at 1 GeV (T. Golan et al., Acta Phys. Polon. B 40, 2519 (2009)):

$$P(\pi^{0} \to \pi^{0}) = 67\%, \quad P(\pi^{0} \to \pi^{+}) = 5\%$$

$$P(\pi^{+} \to \pi^{+}) = 69\%, \quad P(\pi^{+} \to \pi^{0}) = 5\%.$$
(2)

• In use:
$$C_5^A(0) = 1.0$$
 and $M_{A\Delta} = 0.93$ GeV.



- Data underestimated, but:
- Lack of coherent process, no 2nd resonance region or DIS (multiple mesons) + absorption (final state interactions) →1π final state.

Comparison with MiniBooNE

- (Almost) the same model with same $C_5^A(Q^2)$ E. Hernández, J. Nieves, M. J. Vicente Vacas, Phys. Rev. D 87, 113009 (2013).
- Additional coherent process, D_{13} resonance. FSI in cascade model.



• Two extra mechanisms for SPP, actual FSI and still not enough.



Comparison with MiniBooNE

- Different model O. Lalakulich, U. Mosel Phys.Rev. C87 (2013) 014602.
- Different background model, more resonances, migration from other channels through FSI (GiBUU transport model)



experiments. Again, theory below data.

Possible explanations



- Atomic nucleus: more mechanisms possible
- Example: SPP on top of meson exchange current (MEC 2p2h1π-type).
- Requirements: good understanding of SPP on free targets and MEC itself.

Other possibilities: need for consistent treatment of ∆ off-shell propagator. In use: on-shell propagator (solution to free Rarita-schwinger equation), width put "by hand":

 Another possible statement: MiniBooNE (carbon) "prefers" higher C₅^A, than ANL/BNL (deuteron)": medium modification of Δ excitation vertex Uniwersytet

Possible explanations



- Full model $C_5^A(0) = 1.0$, $M_{A\Delta} = 0.93$ GeV, no cut in invariant mass W.
- Deuteron effects $\sigma^{^{2}H} = \int d^{3}p \frac{f(p)}{v_{rel.}} \sigma^{free}(p).$
- Norm of deuteron wave function f(p), binding energy $B = 2\sqrt{p^2 + M^2} M_{^2H}.$
- Toying around with Delta-background interference sign.
- Interference phase- possibly large effect in neutron channels.
- Actual solution: unitarization with Watson's theorem (J. Nieves, L. Alvarez-Ruso)→ dynamical Delta-background complex phase (not a sign flip!).



- SPP problem leaves us with more questions, than definite answers. Still no completely satisfying SPP model.
- MiniBooNE CH₂ data: underestimated, problem with simultaneous ANL/BNL data reproduction in all isospin channels.
 - $\bullet\,$ Proper off-shell Δ treatment; possible in-medium excitation vertex modifications?
 - Proper weak Delta-background phases?
 - Extra channels coming from MEC $2p2h1\pi$ and multi-pion+FSI $\rightarrow 1\pi$?
- ANL-BNL-MiniBooNE data: seemingly different, large normalization errors. Need for more precise experiments.



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