

Observation of Low- and High-Energy Gamow-Teller Phonon Excitations in Nuclei

GT : weak response caused by simple $\sigma\tau$ operator



International School
of Nuclear Physics
@Erice

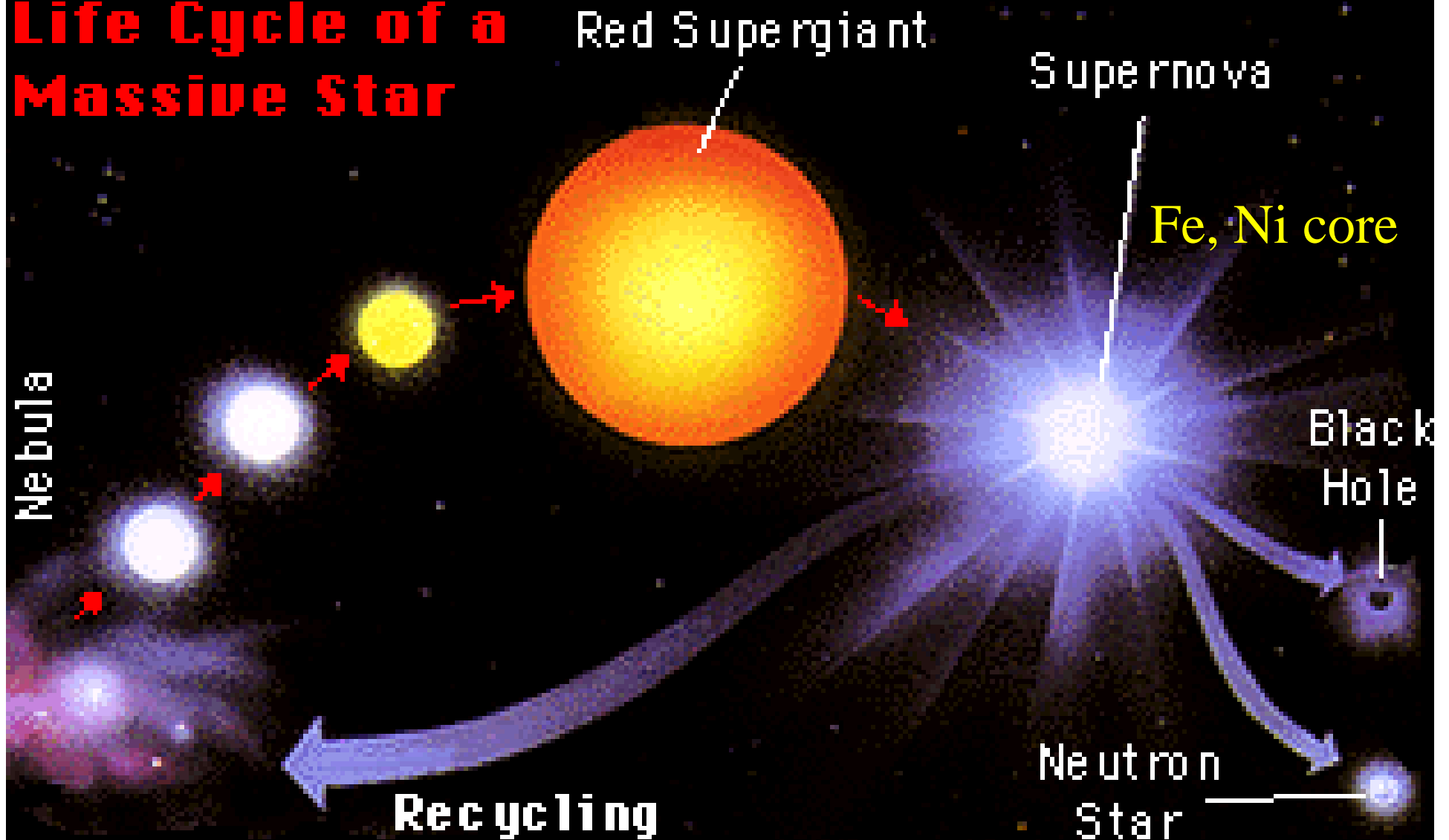
Sep. 16 – Sep. 24, 2014

Yoshitaka FUJITA

RCNP & Dept. Phys.,
Osaka Univ.

Stellar Evolution & Supernova Cycle

Life Cycle of a Massive Star



Gamow-Teller transitions

Mediated by $\sigma\tau$ operator

$$\Delta S = -1, 0, +1 \quad \text{and} \quad \Delta T = -1, 0, +1$$

($\Delta L = 0$, no change in radial w.f.)

→ no change in spatial w.f.

Accordingly, transitions among $j_>$ and $j_<$ configurations

$$j_> \rightarrow j_>, \quad j_< \rightarrow j_<,\quad j_> \leftrightarrow j_<$$

example $f_{7/2} \rightarrow f_{7/2}, \quad f_{5/2} \rightarrow f_{5/2}, \quad f_{7/2} \leftrightarrow f_{5/2}$

Note that Spin and Isospin are
unique quantum numbers in atomic nuclei !

→ GT transitions are sensitive to Nuclear Structure !

→ GT transitions in each nucleus are UNIQUE !

**Basic common understanding of β -decay and Charge-Exchange reaction

β decays :

Absolute $B(\text{GT})$ values,

but usually the study is limited to low-lying states

(p,n), (^3He ,t) reaction at 0° :

Relative $B(\text{GT})$ values, but **Highly Excited States**

** Both are important for the study of GT transitions!

β -decay & CE Nuclear Reaction

* β -decay GT tra. rate = $\frac{1}{t_{1/2}} = f \frac{\lambda^2}{K} B(\text{GT})$

$B(\text{GT})$: reduced GT transition strength
 $\propto (\text{matrix element})^2 = |\langle f | \sigma \tau | i \rangle|^2$

*Nuclear (CE) reaction rate (cross-section)
= reaction mechanism

⊗ operator

⊗ structure

$= (\text{matrix element})^2$

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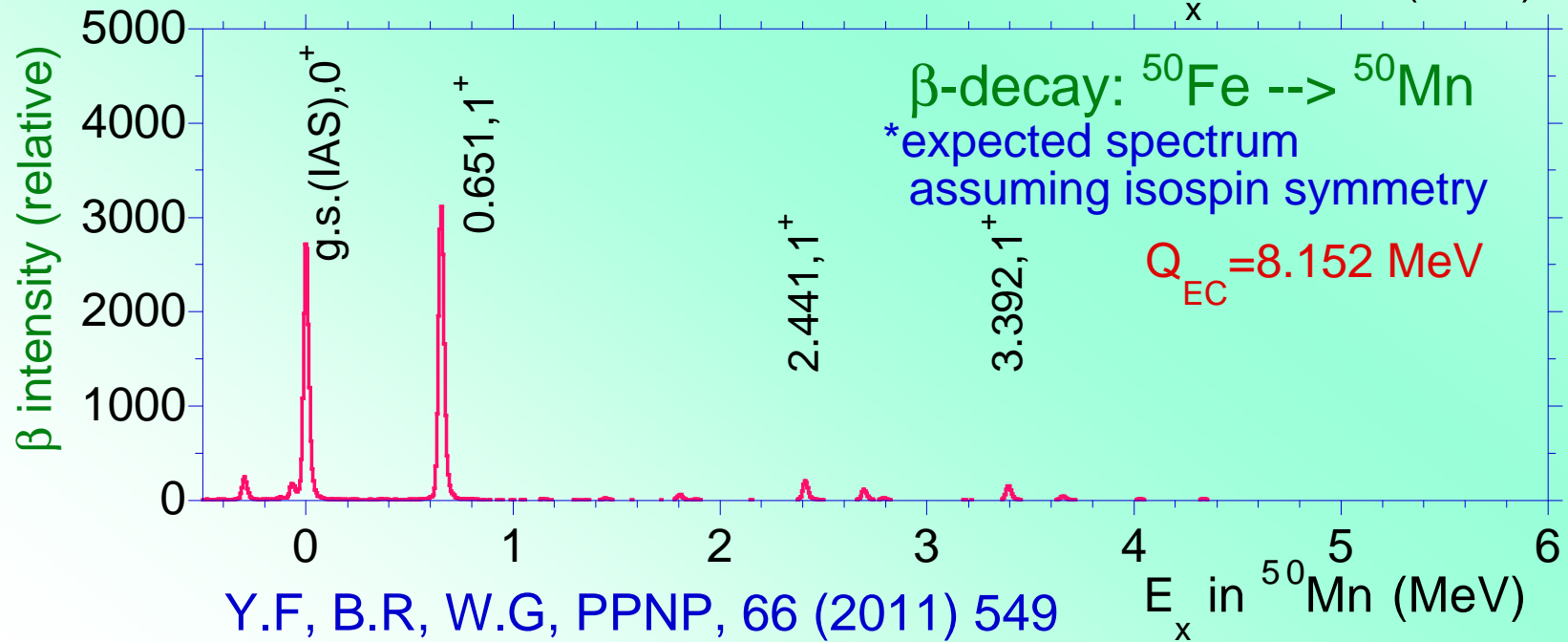
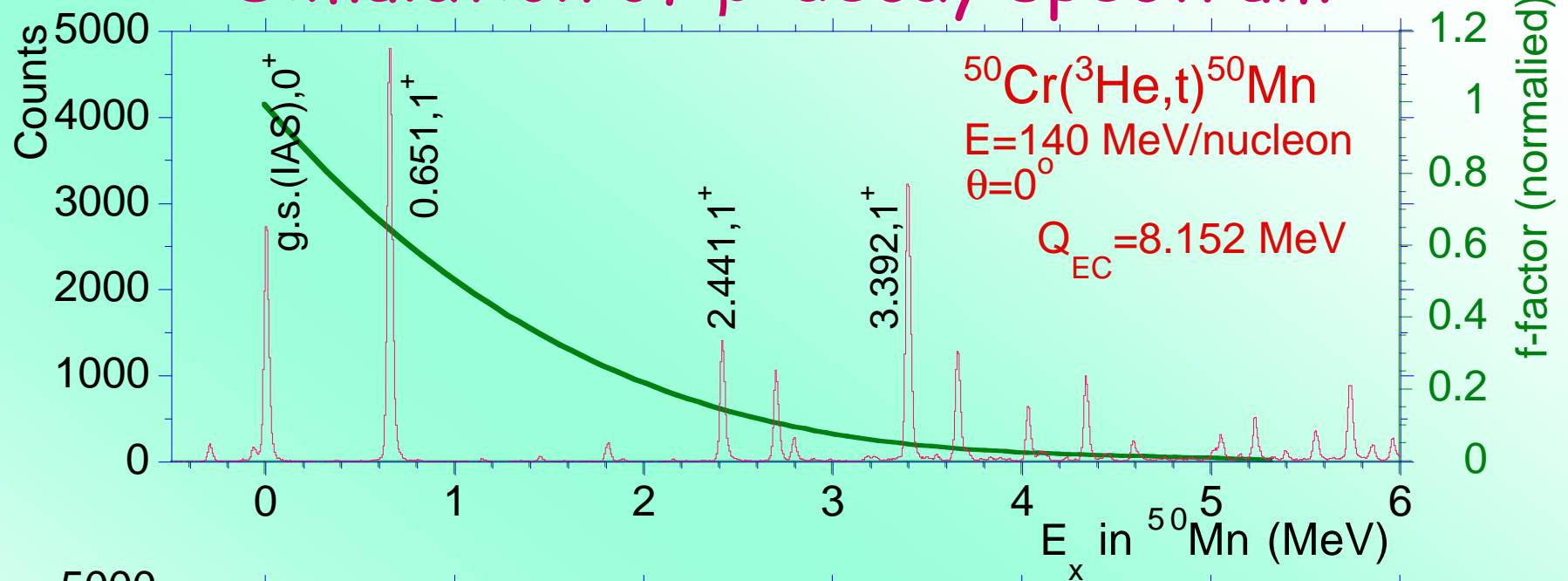
*Nuclear (CE) reaction rate (cross-section)
= reaction mechanism

\otimes operator

\otimes structure

$= (\text{matrix element})^2$

Simulation of β -decay spectrum



Y.F, B.R, W.G, PPNP, 66 (2011) 549

β -decay & Nuclear Reaction

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(\otimes) operator

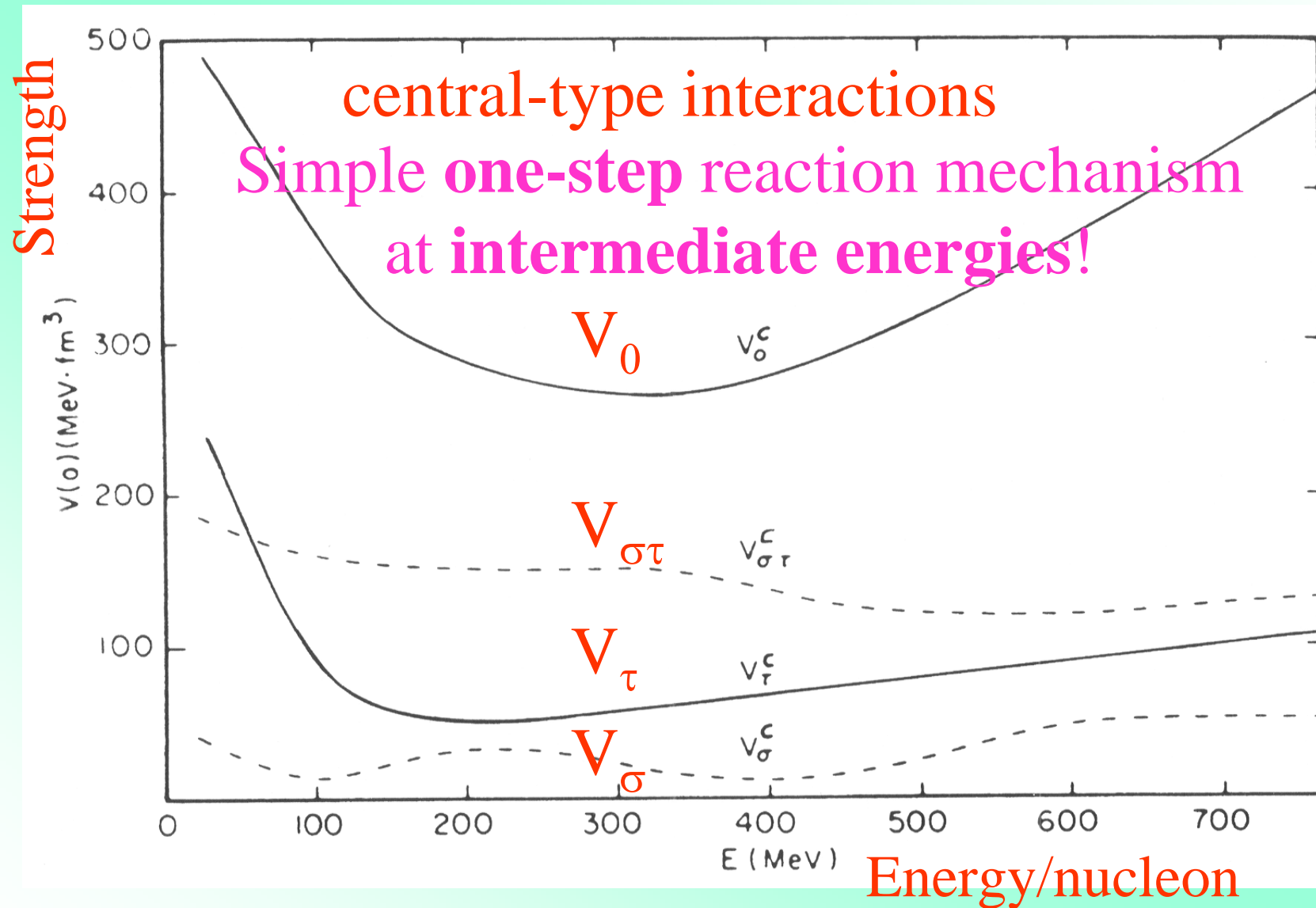
(\otimes) structure

=(matrix element)²

*At intermediate energies ($100 < E_{\text{in}} < 500$ MeV)

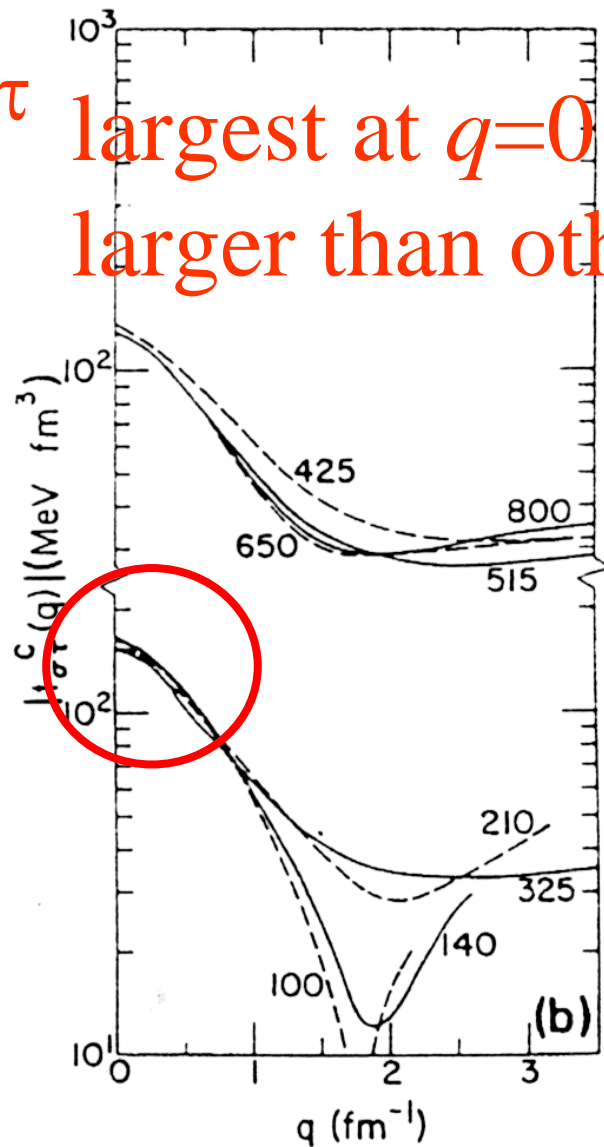
→ $d\sigma/d\omega(q=0)$: proportional to $B(\text{GT})$

Nucleon-Nucleon Int. : E_{in} dependence at $q=0$

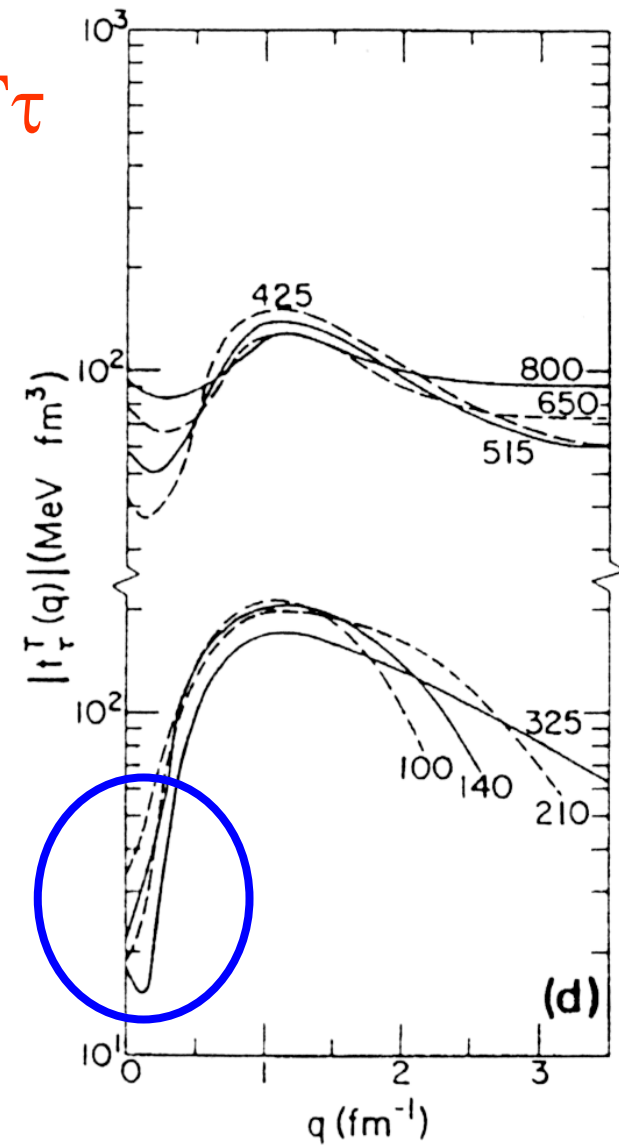


N.-N. Int. : $\sigma\tau$ & Tensor- τ q -dependence

$\sigma\tau$ largest at $q=0$!
larger than others !



$T\tau$



Love & Franey PRC 24 ('81) 1073

β -decay & Nuclear Reaction

$$*\beta\text{-decay GT tra. rate} = \frac{1}{t_{1/2}} = f \frac{\lambda^2}{K} B(\text{GT})$$

$B(\text{GT})$: reduced GT transition strength
 $\propto (\text{matrix element})^2 = |\langle f | \sigma \tau | i \rangle|^2$

*Nuclear (CE) reaction rate (cross-section)
= reaction mechanism

⊗ operator

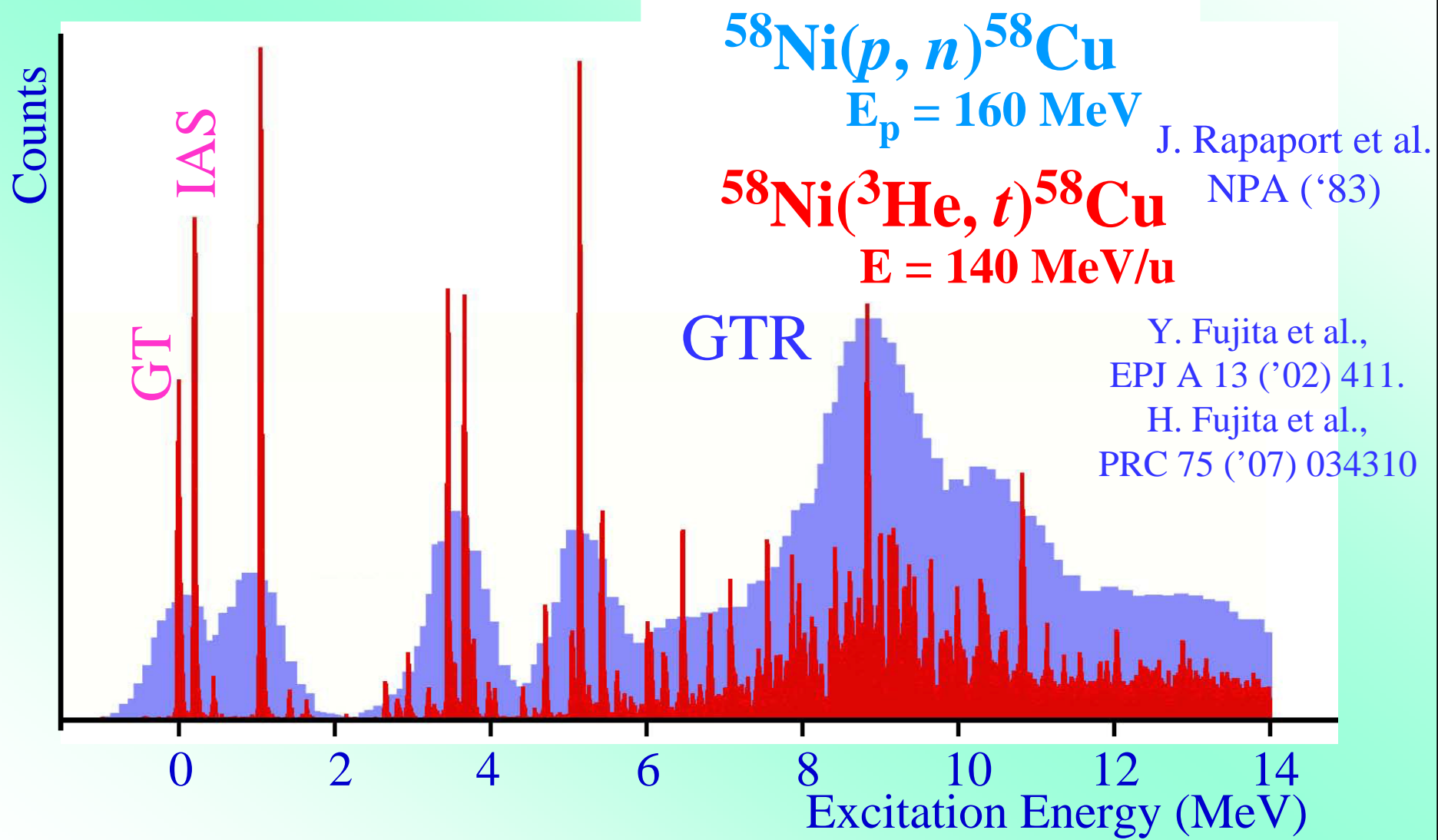
⊗ structure

=(matrix element)²

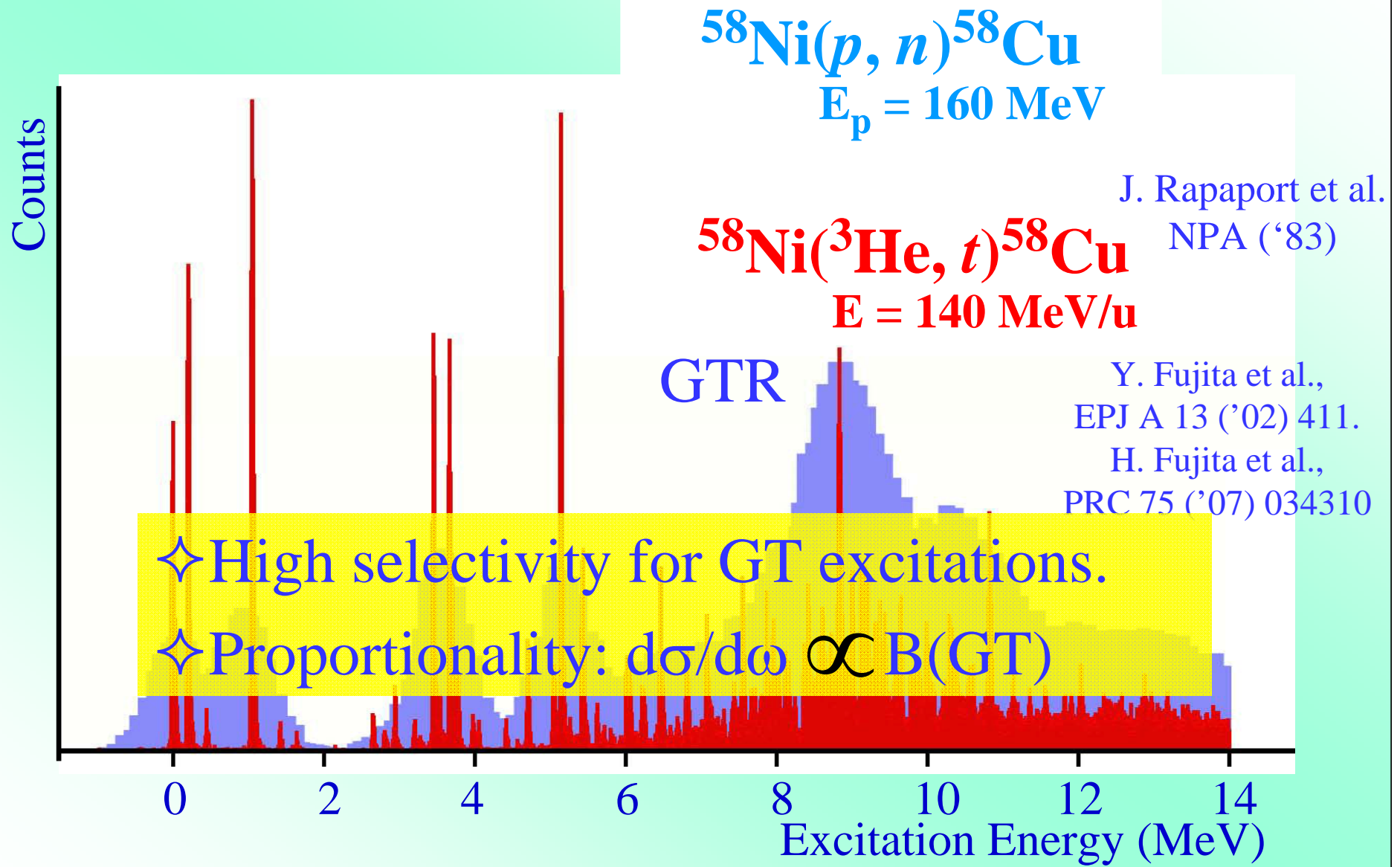
*At intermediate energies ($100 < E_{\text{in}} < 500$ MeV)

→ $d\sigma/d\omega(q=0)$: proportional to $B(\text{GT})$

Comparison of (p, n) and (³He, t) 0° spectra



Comparison of (p, n) and (³He, t) 0° spectra



β -decay & Nuclear Reaction

* β -decay GT tra. rate = $\frac{1}{t_{1/2}} = \underbrace{f}_{\text{operator}} \frac{\lambda^2}{K} \underbrace{B(\text{GT})}_{\text{structure}}$

$B(\text{GT})$: reduced GT transition strength
 $\propto (\text{matrix element})^2$
 Study of Weak Response of Nuclei
 by means of

* Nuclear (GE) reaction rate (cross-section)
 = reaction mechanism
 using β -decay as a reference

\otimes operator
 \otimes structure

$= (\text{matrix element})^2$

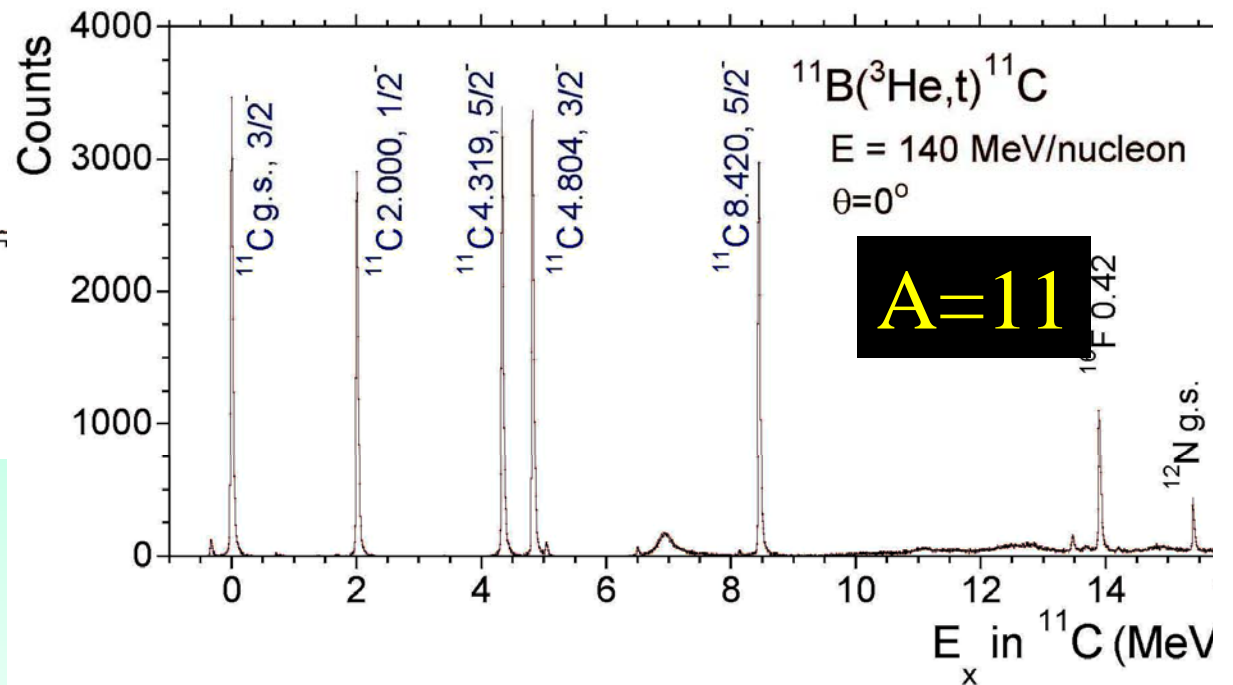
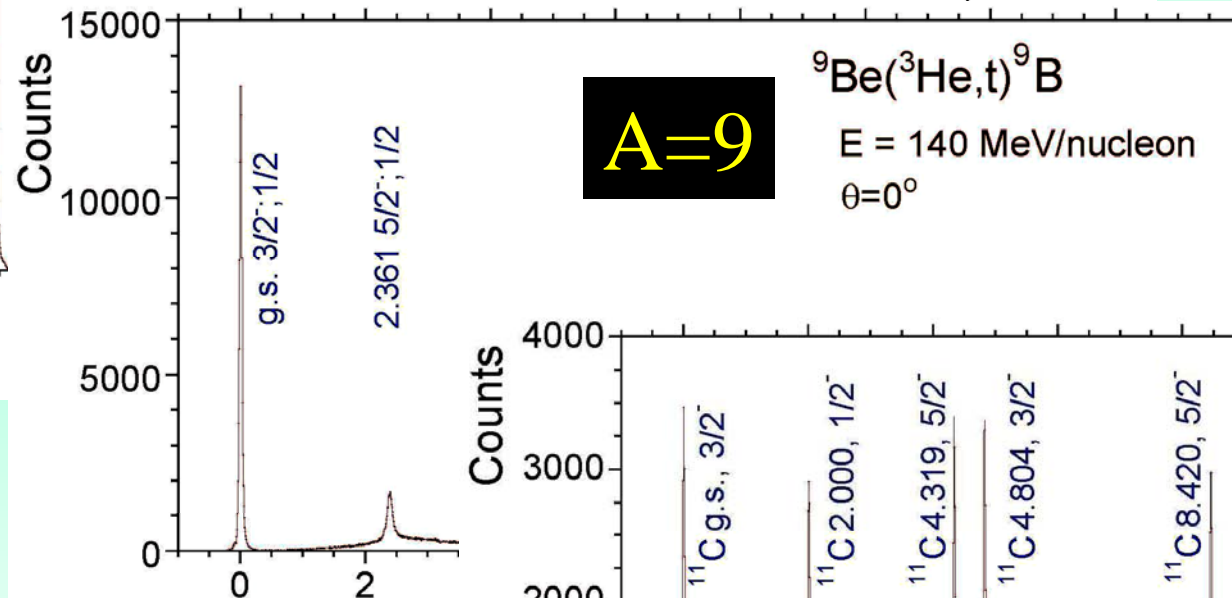
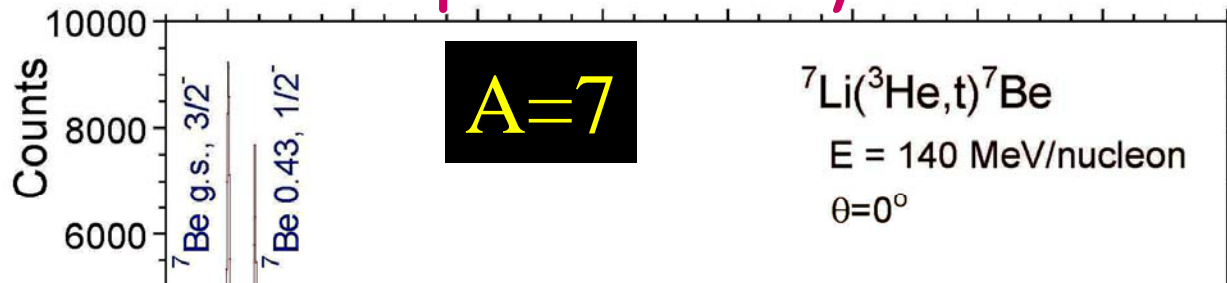
A simple reaction mechanism should be achieved !

\rightarrow we have to go to high incoming energy

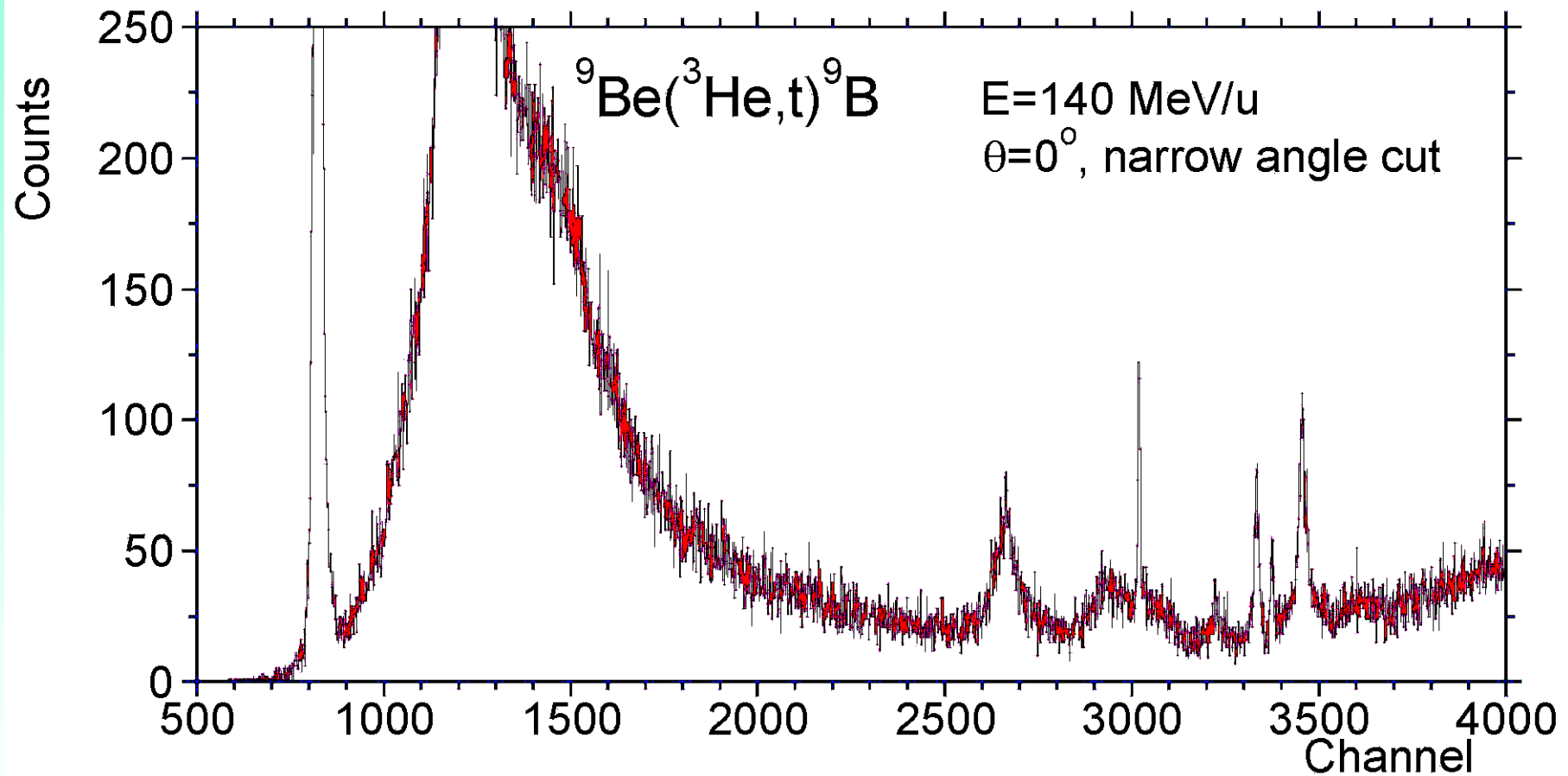
****GT transitions in each nucleus are
UNIQUE!**

***($^3\text{He}, t$): high resolution and sensitivity !**

Spectra of p -shell $T_z=1/2$ Nuclei



${}^9\text{Be}({}^3\text{He},t){}^9\text{B}$ spectrum (various scales)



Relationship: Decay and Width

Heisenberg's Uncertainty Principle

$$\Delta x \cdot \Delta p \approx \hbar$$

$$\Delta t \cdot \Delta E \approx \hbar$$

Width $\Gamma = \Delta E$

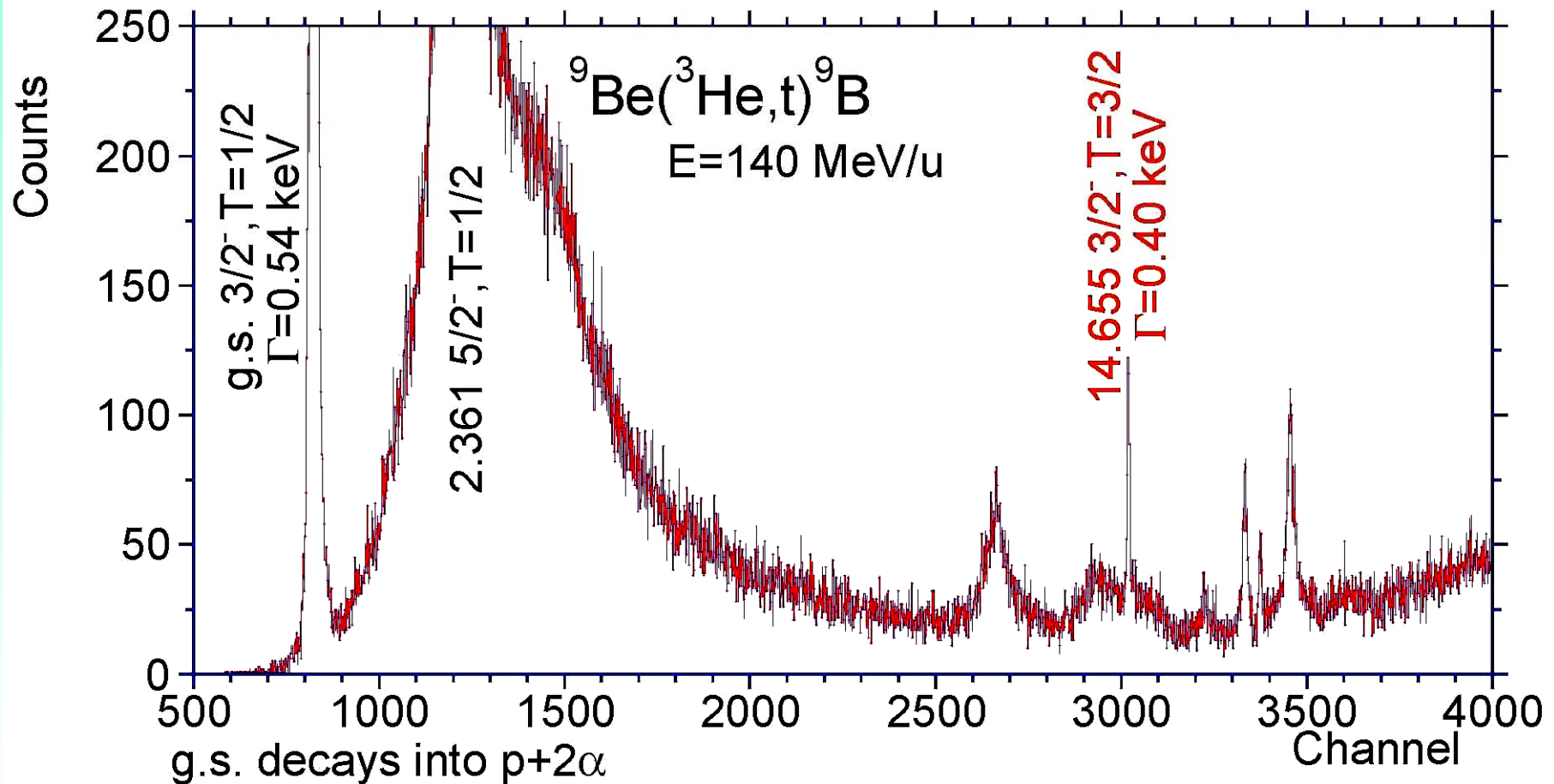
*if: Decay is Fast,

then: Width of a State is Wider !

*if $\Delta t = 10^{-20}$ sec $\rightarrow \Delta E \sim 100$ keV (particle decay)

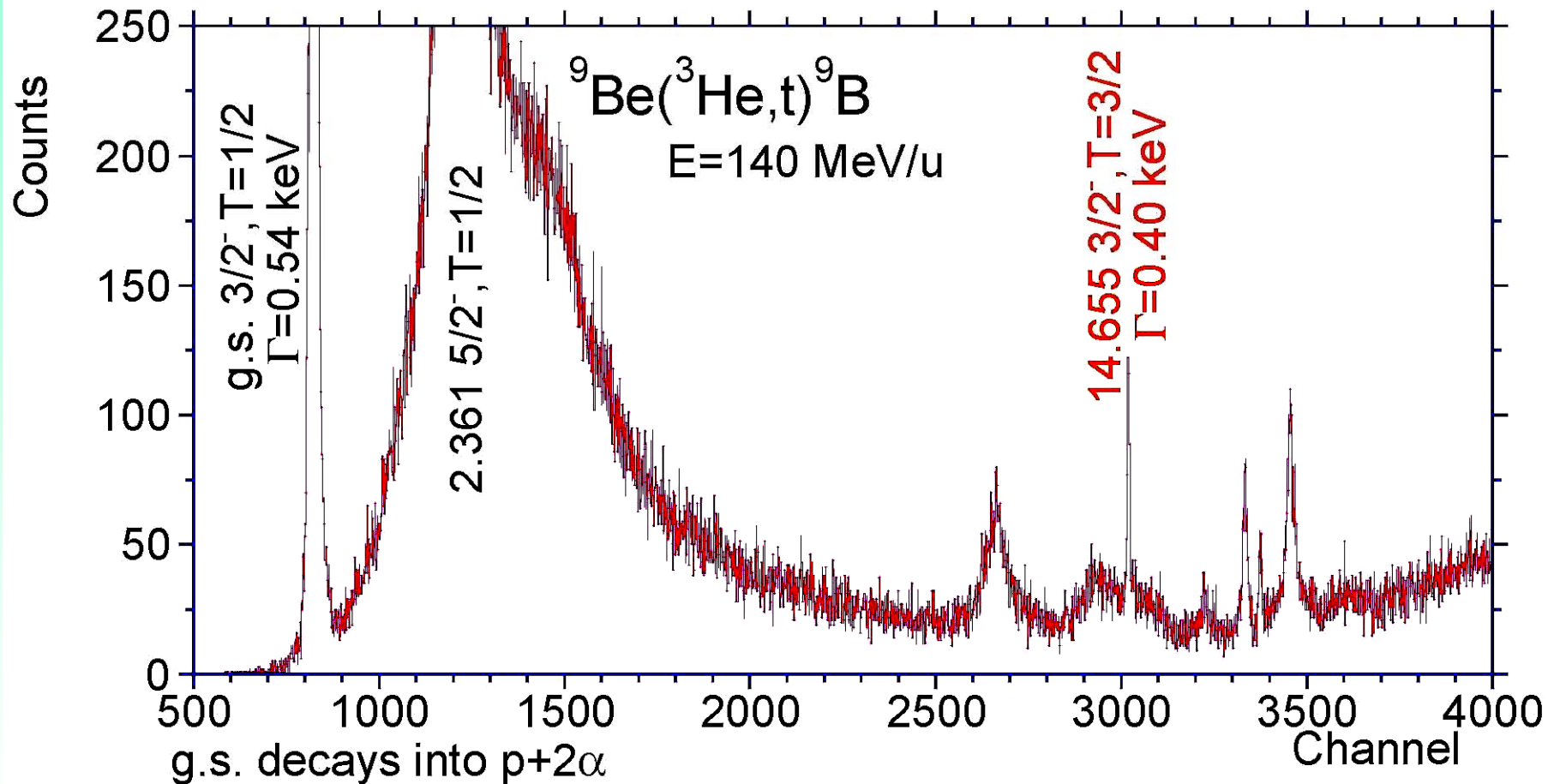
$\Delta t = 10^{-15}$ sec $\rightarrow \Delta E \sim 1$ eV (fast γ decay)

${}^9\text{Be}({}^3\text{He},t){}^9\text{B}$ spectrum (II)



**Isospin selection rule prohibits
proton decay of $T=3/2$ state!**

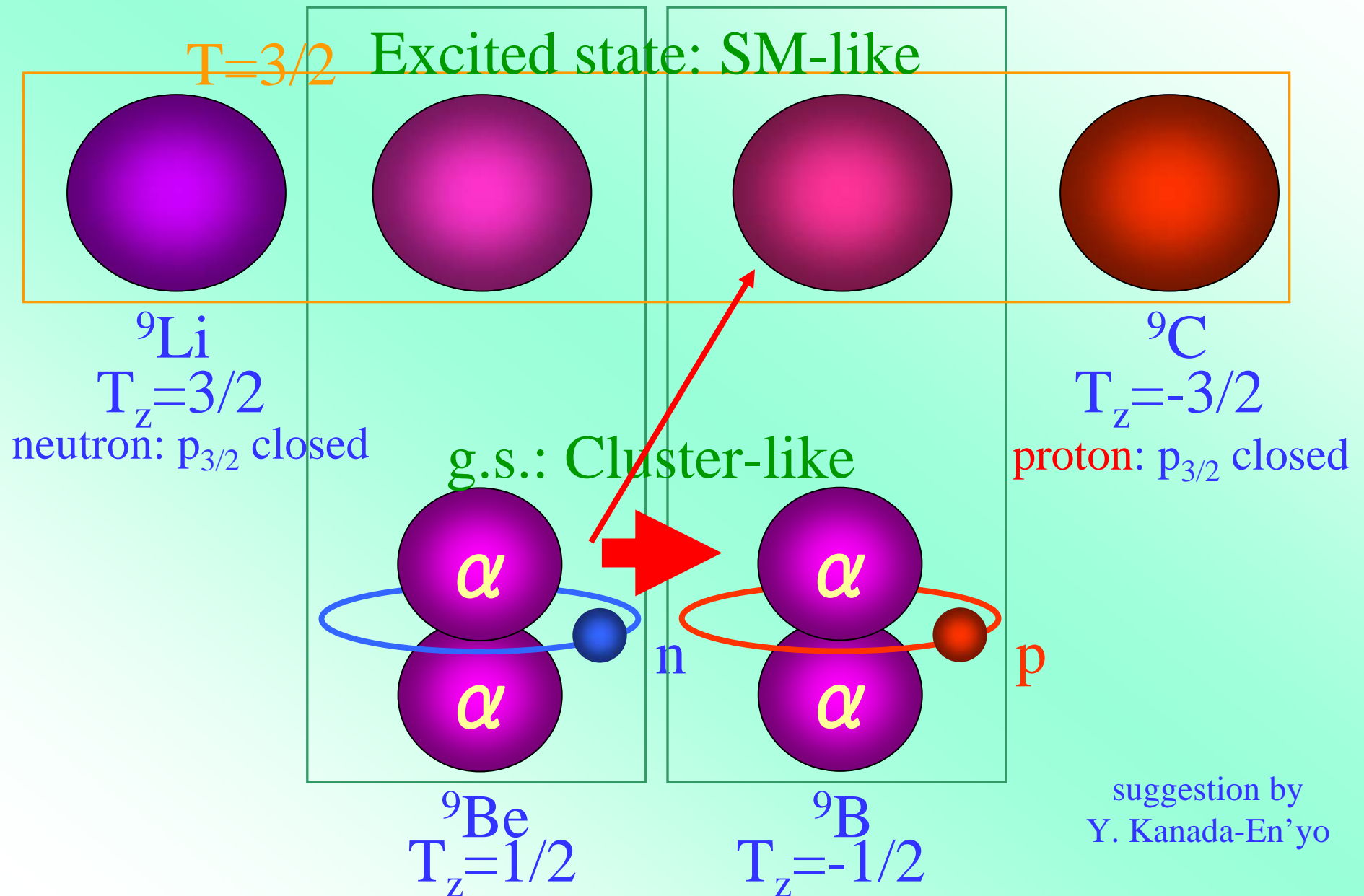
${}^9\text{Be}({}^3\text{He},t){}^9\text{B}$ spectrum (III)



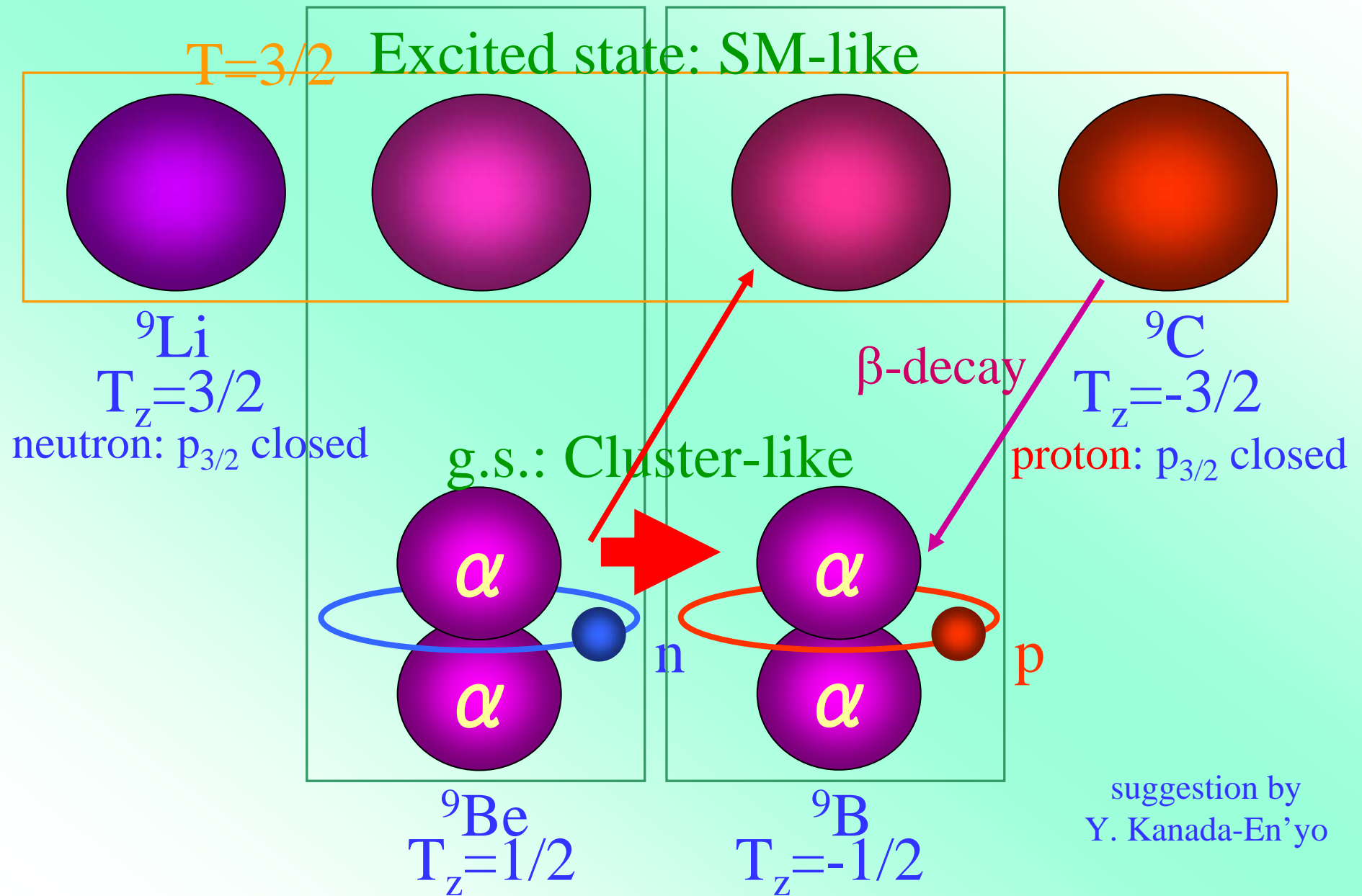
14.7 MeV $T=3/2$ state is very weak!

Strength ratio of g.s. & 14.7 MeV $3/2^-$ states: **140:1**

Shell Structure and Cluster Structure



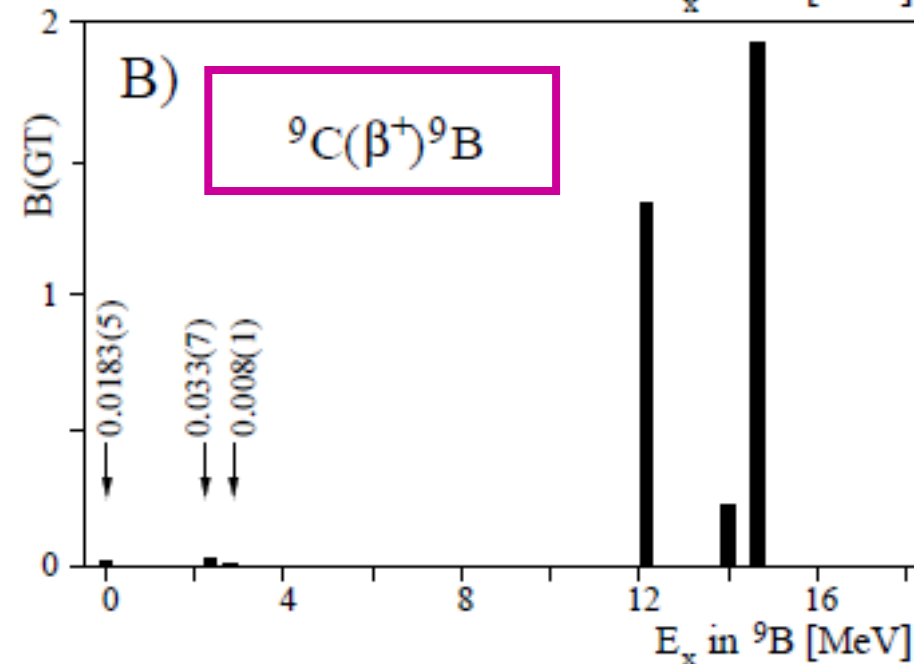
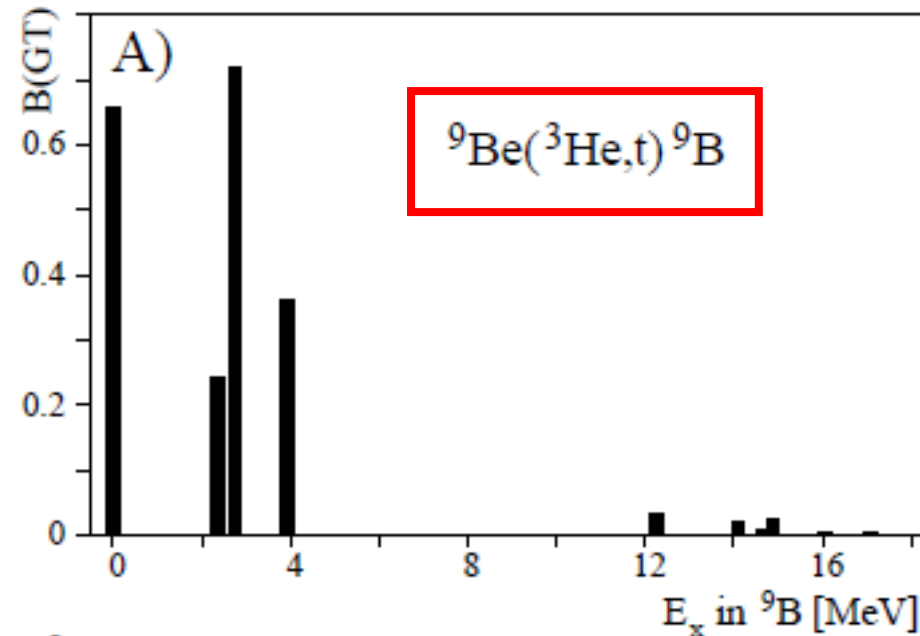
Shell Structure and Cluster Structure



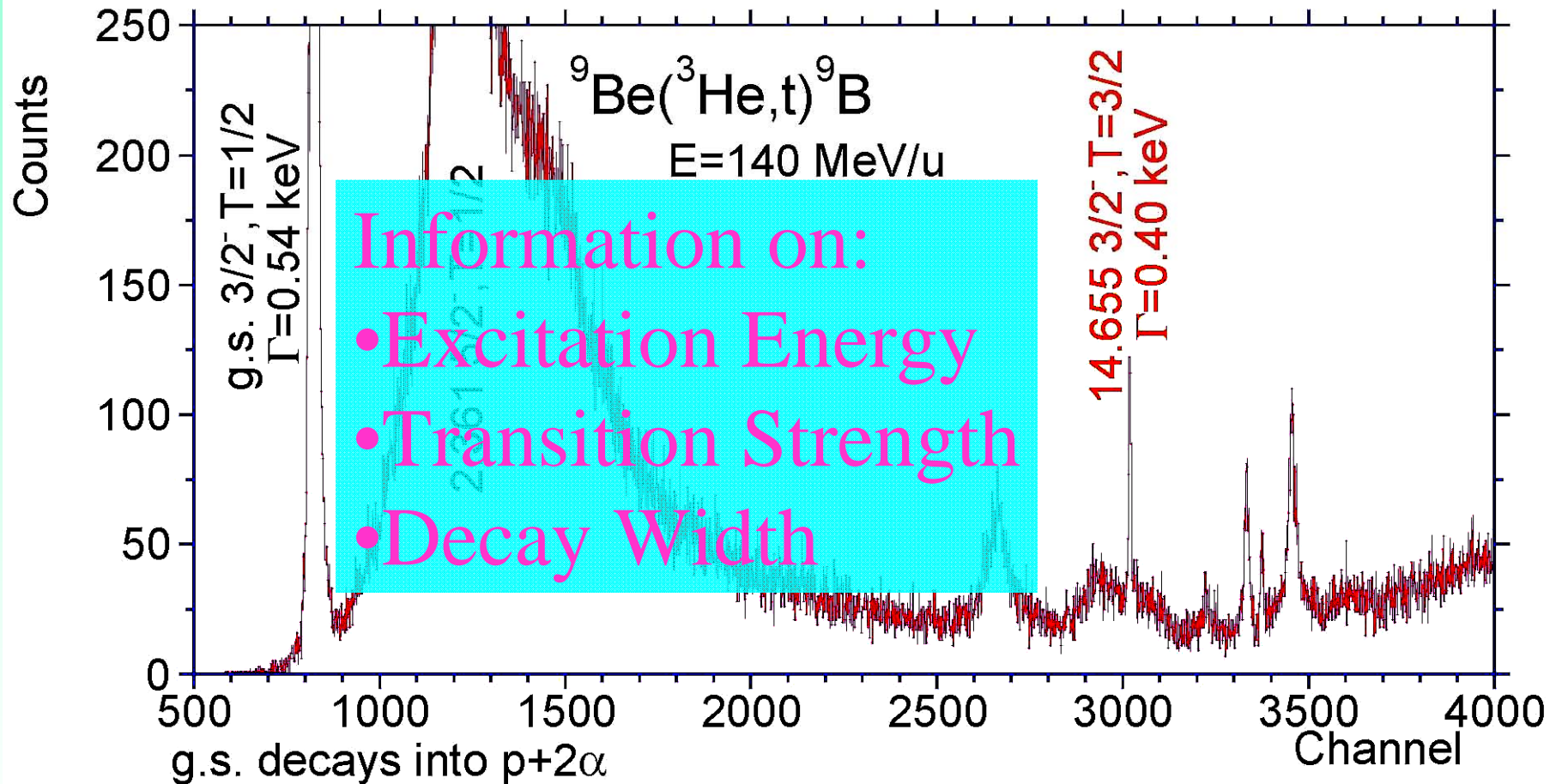
β -decay and $(^3\text{He}, t)$ results

C. Scholl et al,
PRC 84, 014308 (2011)

L. Buchmann et al.,
PRC 63 (2001) 034303.
U.C. Bergmann et al.,
Nucl. Phys. A 692 (2001) 427.



${}^9\text{Be}({}^3\text{He},t){}^9\text{B}$ spectrum (III)



14.7 MeV $T=3/2^-$ state is very weak!

Strength ratio of g.s. & 14.7 MeV $3/2^-$ states: **140:1**

****Connection between
 β -decay and ($^3\text{He},t$) reaction****

**by means of
Isospin Symmetry**

***Isospin Symmetry

an important idea to see the connection of
decays and excitations caused
by Strong, EM and Weak interactions !

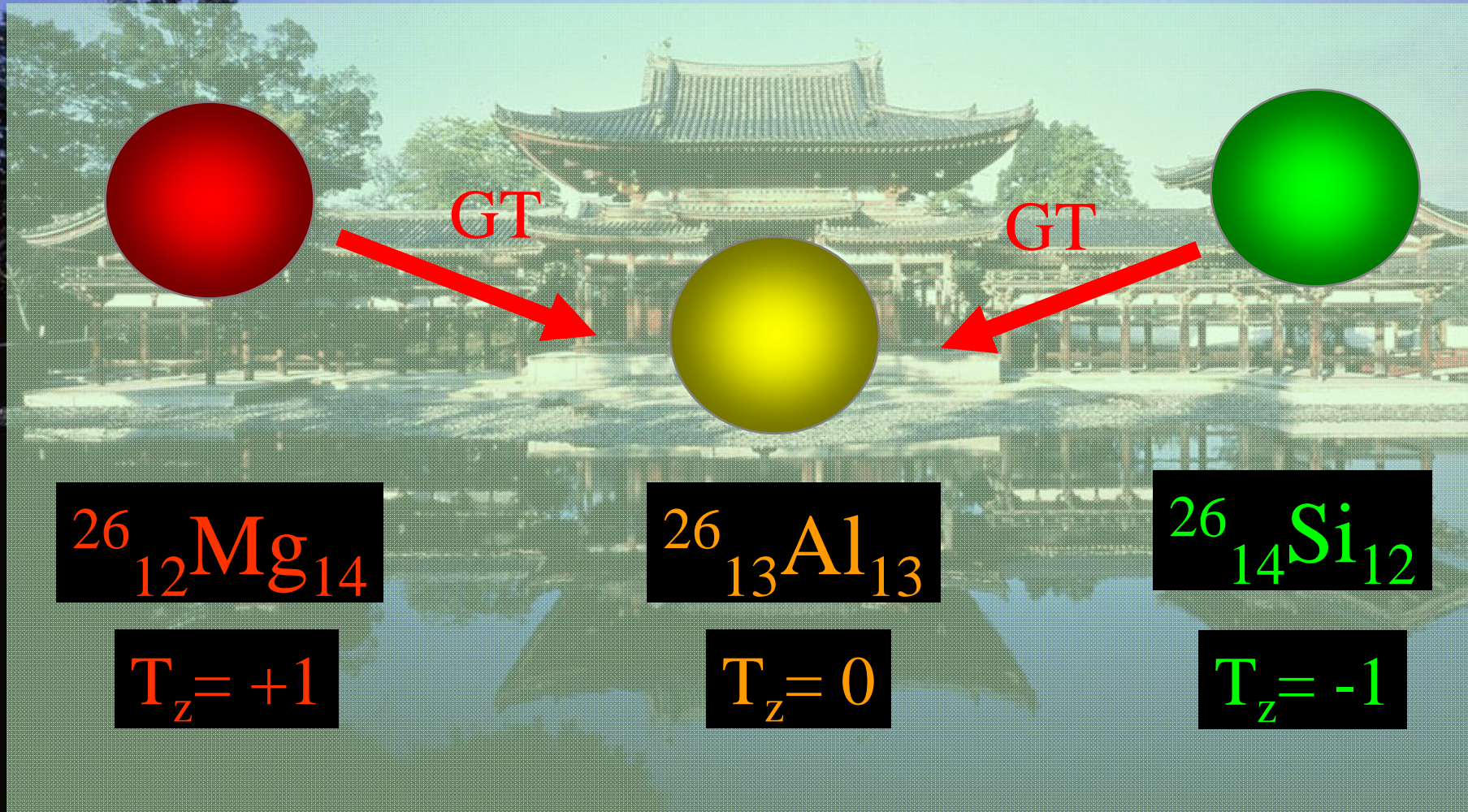
There are many cases that the “operators” are the same
in transitions caused by “strong,” “EM” and “weak” int.

T=1 Isospin Symmetry



Byodoin-temple,
Uji, Kyoto

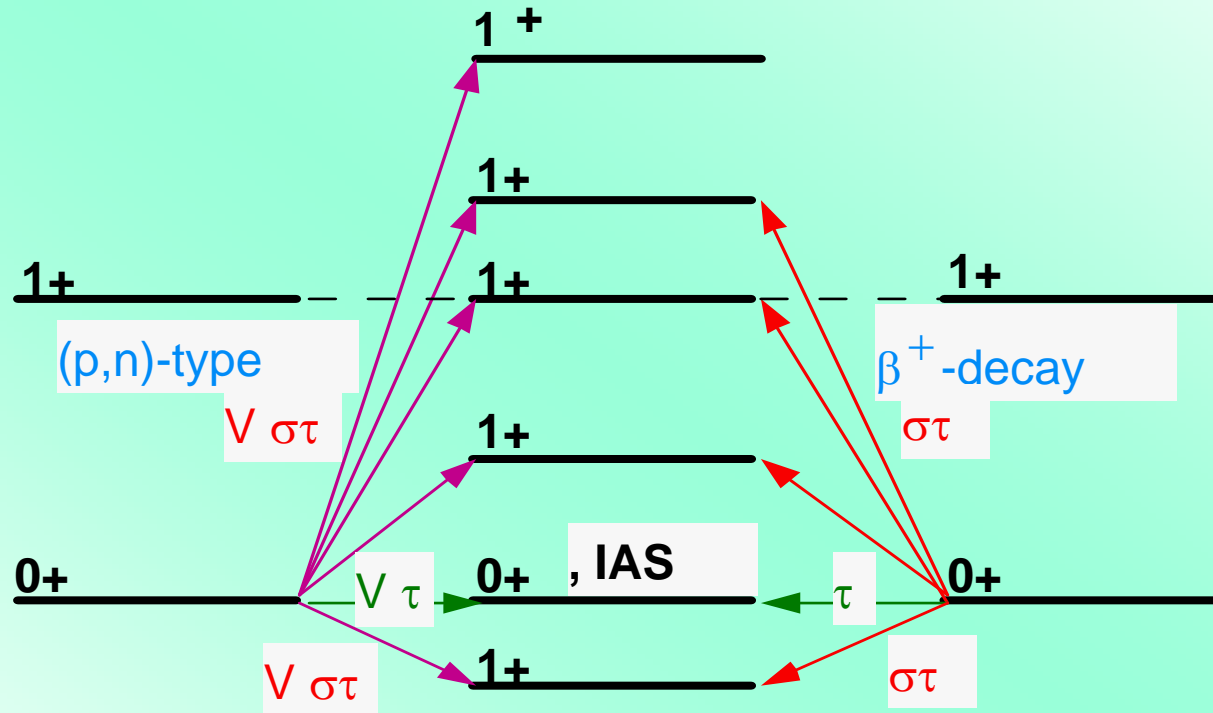
T=1 Isospin Symmetry



T=1 symmetry : Structures & Transitions

$T_z=+1$ \longrightarrow $T_z=0$ \longleftarrow $T_z=-1$

(in isospin symmetry space*)



$T_z=+1$

^{26}Mg

Z=12, N=14

$T_z=0$

^{26}Al

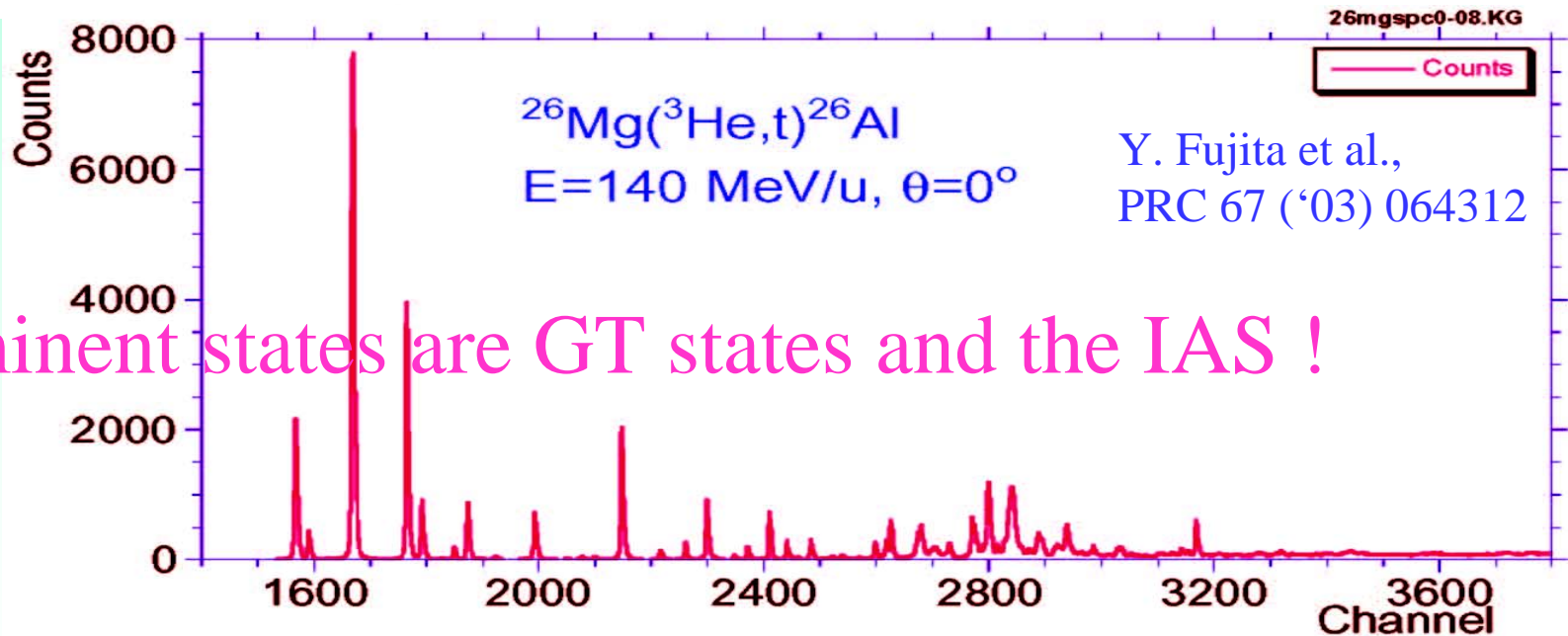
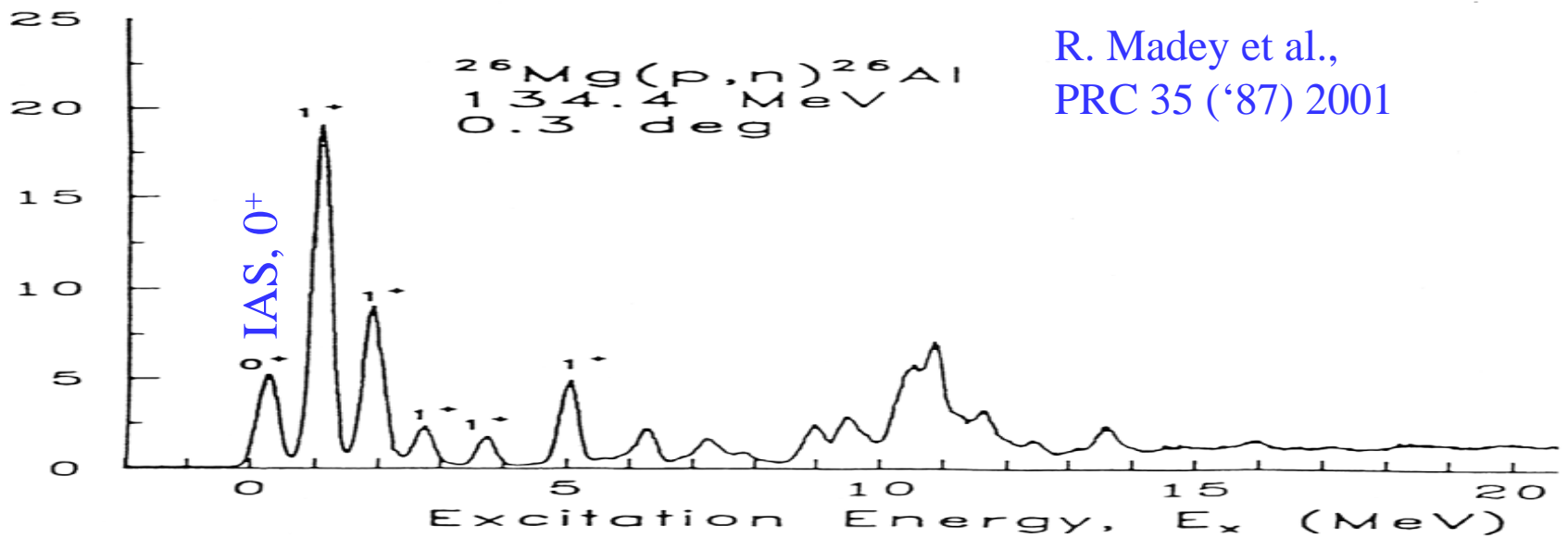
Z=13, N=13

$T_z=-1$

^{26}Si

Z=14, N=12

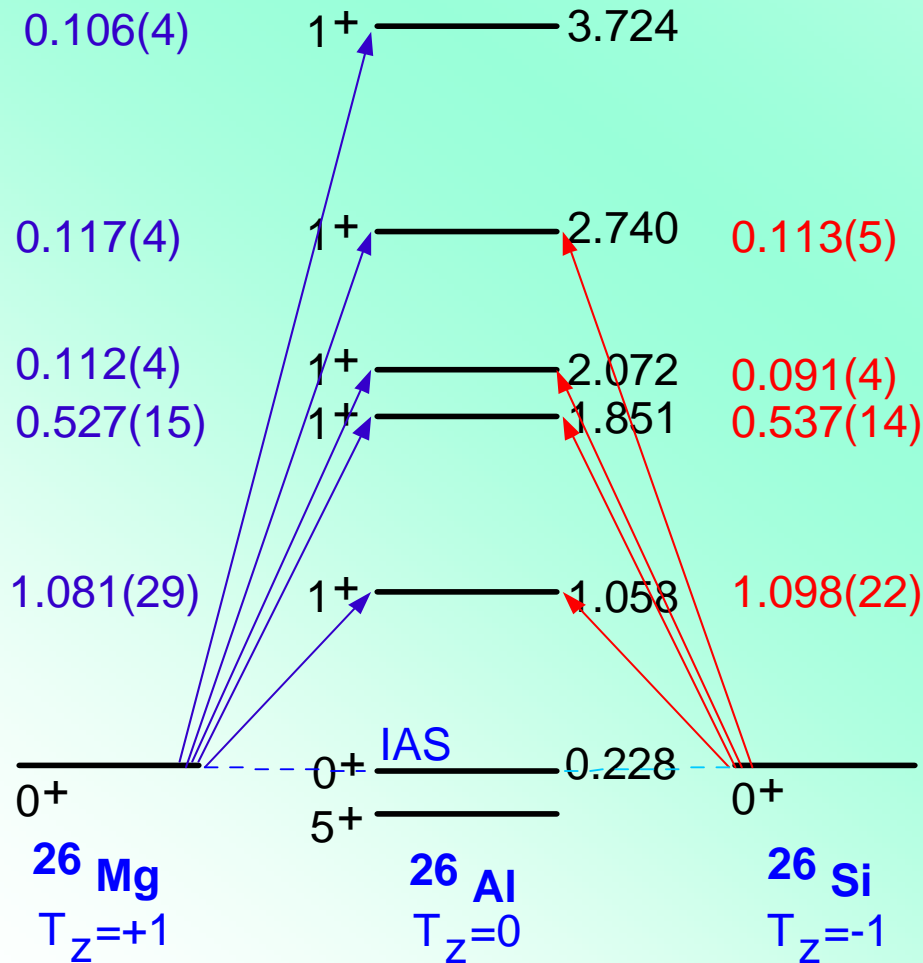
$^{26}\text{Mg}(p, n)^{26}\text{Al}$ & $^{26}\text{Mg}(^3\text{He}, t)^{26}\text{Al}$ spectra



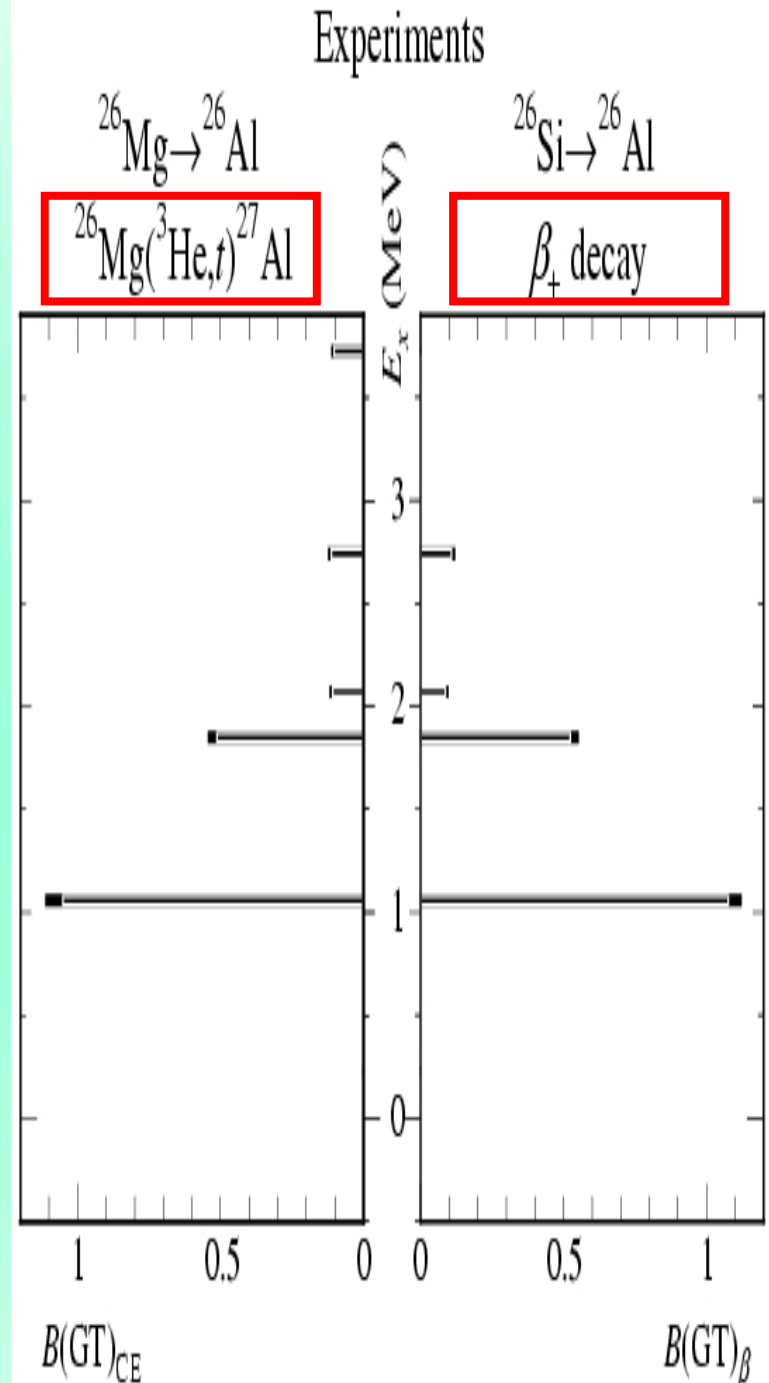
Prominent states are GT states and the IAS !

B(GT) values from Symmetry Transitions (A=26)

$(^3\text{He}, t)$ **β -decay**
 B(GT) B(GT)



Y. Fujita et al., PRC 67 ('03) 064312



RCNP (Osaka) Ring Cyclotron

Good quality ^3He beam (140 MeV/nucleon)

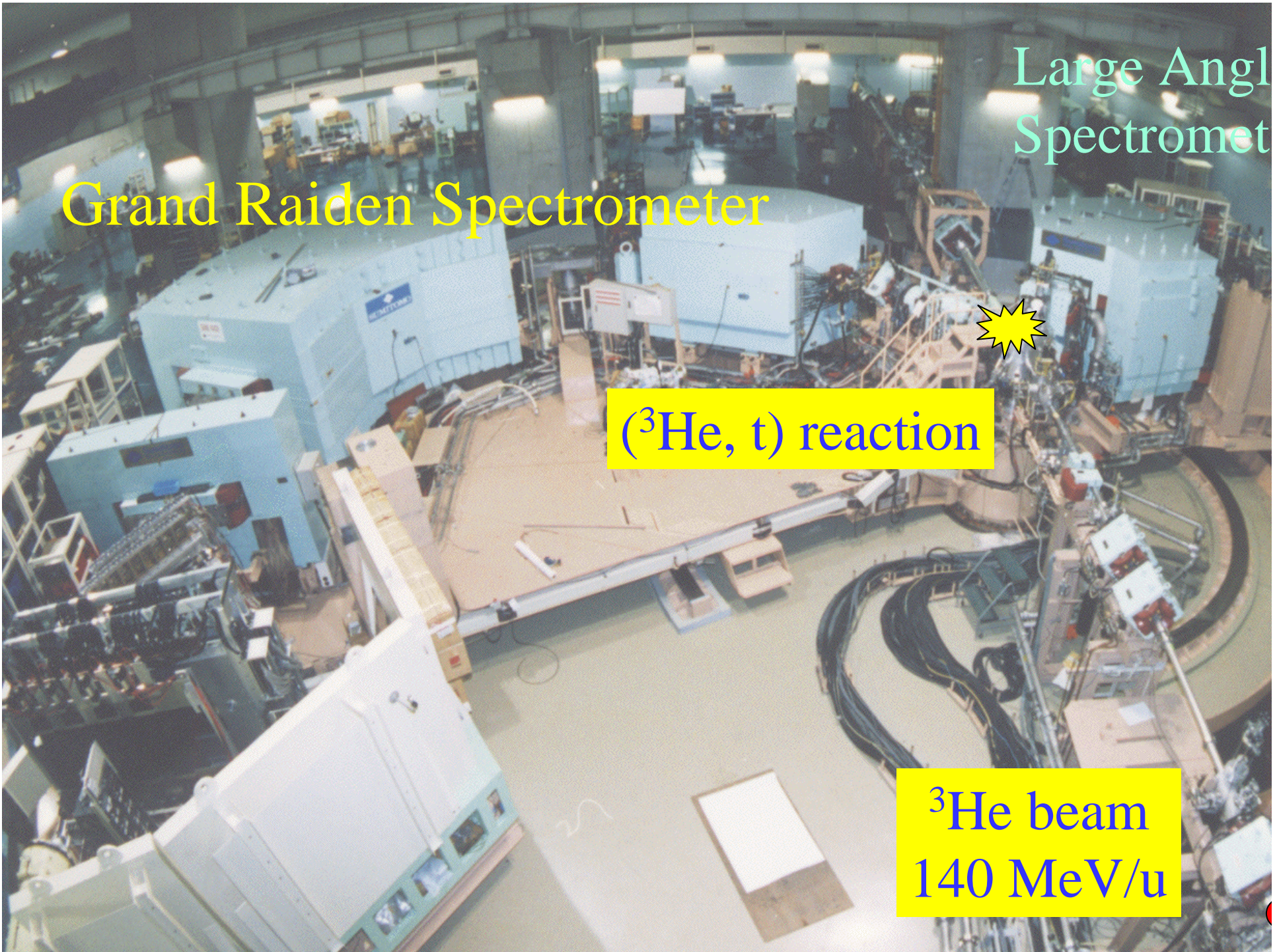


Large Angle
Spectrometer

Grand Raiden Spectrometer

$(^3\text{He}, t)$ reaction

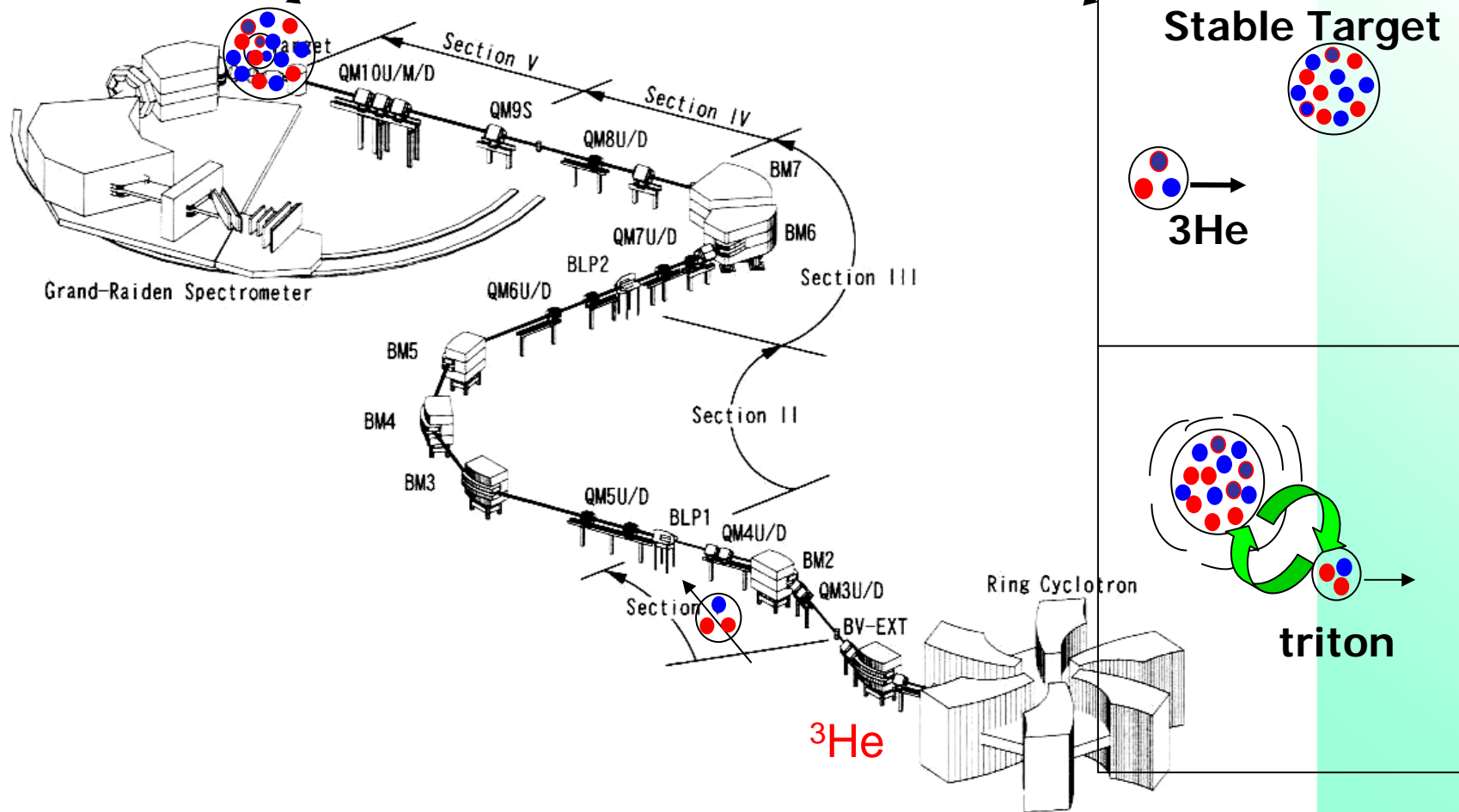
^3He beam
140 MeV/u



$(^3\text{He}, t)$ CE Reactions @ RCNP (Osaka)

$\theta_{\text{lab}} = 0^\circ$

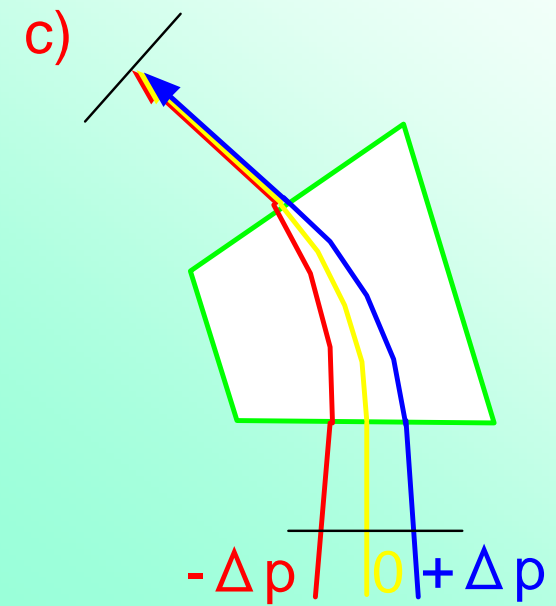
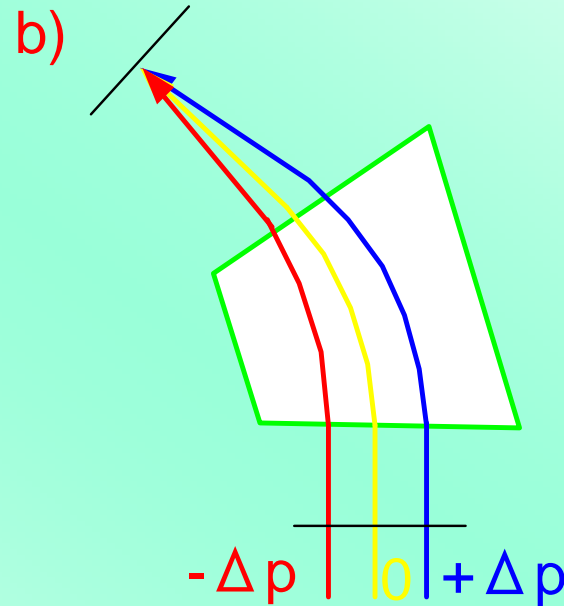
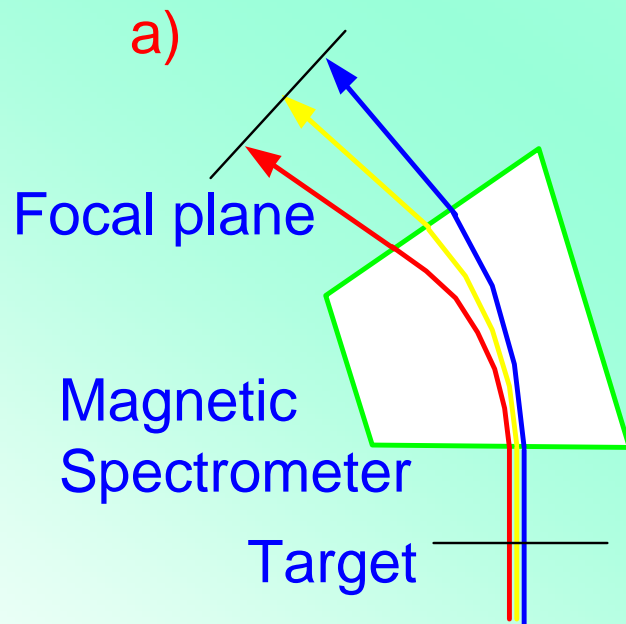
$(^3\text{He}, t)$ CE reaction



Matching Techniques

Y. Fujita et al., N.I.M. B 126 (1997) 274.

H. Fujita et al., N.I.M. A 484 (2002) 17.



*Achromatic beam
transportation*

$\Delta E \sim 200$ keV
for 140MeV/u ^3He beam

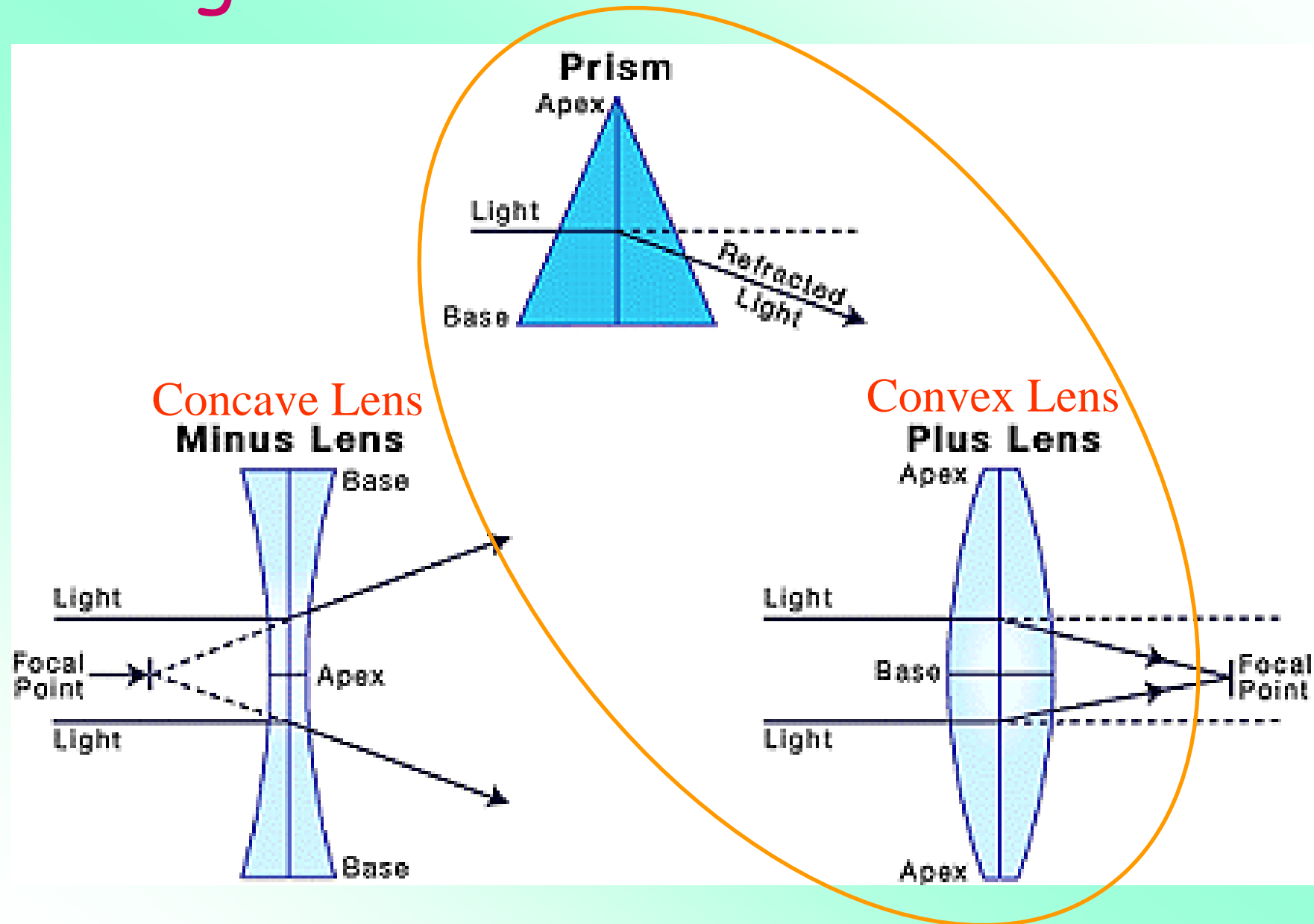
*Lateral dispersion
matching*

$\Delta E \sim 35$ keV
Horiz. angle resolution
 $\Delta\theta_{\text{sc}} > 15$ mrad

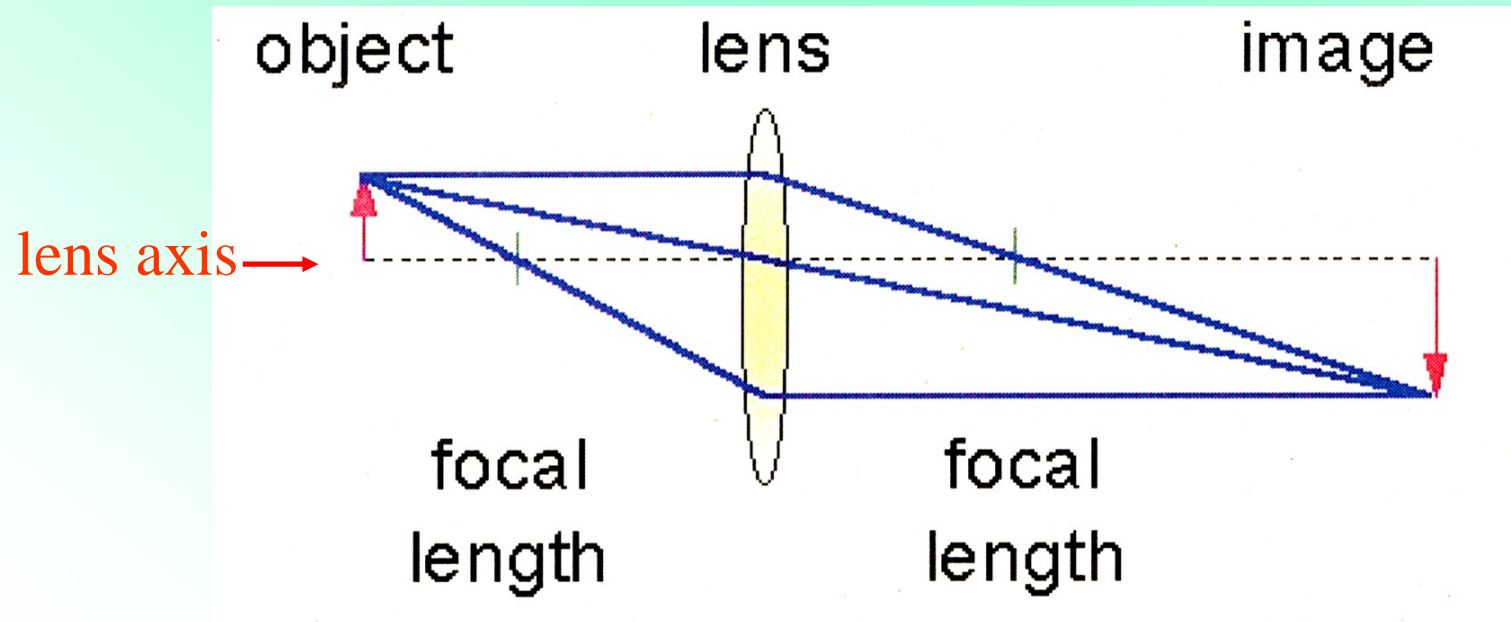
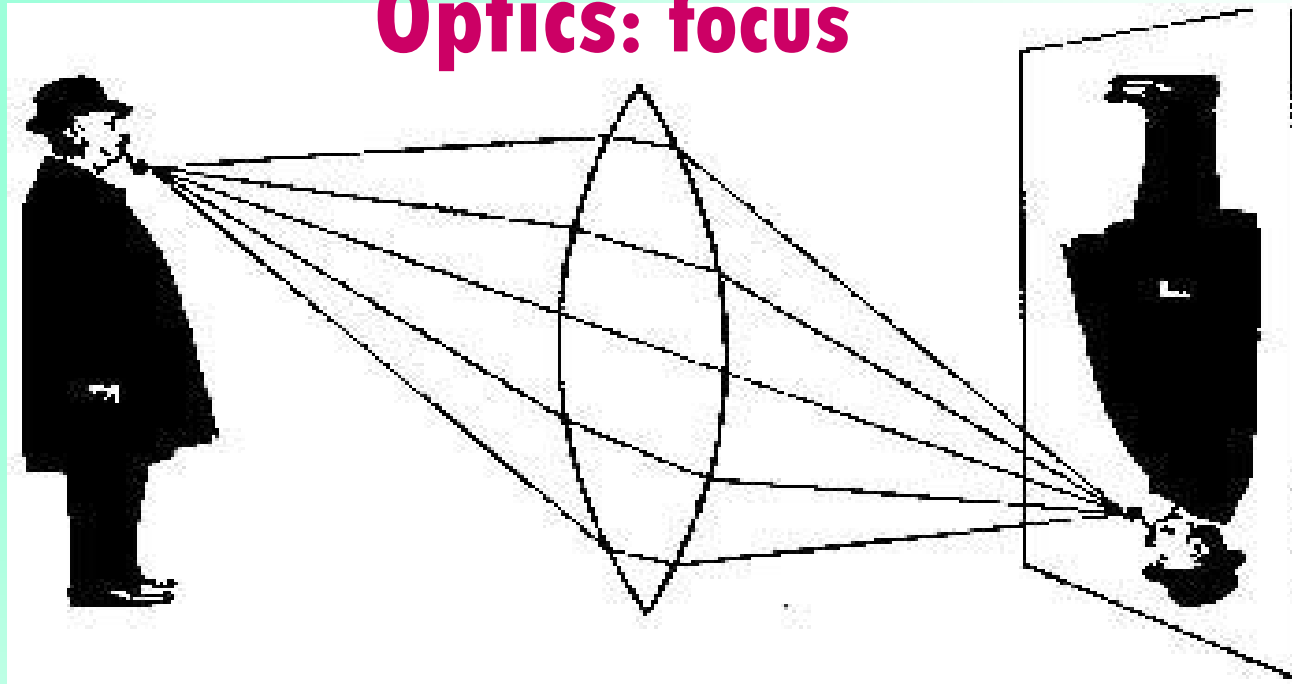
*Angular dispersion
matching*

$\Delta\theta_{\text{sc}} \sim 5$ mrad

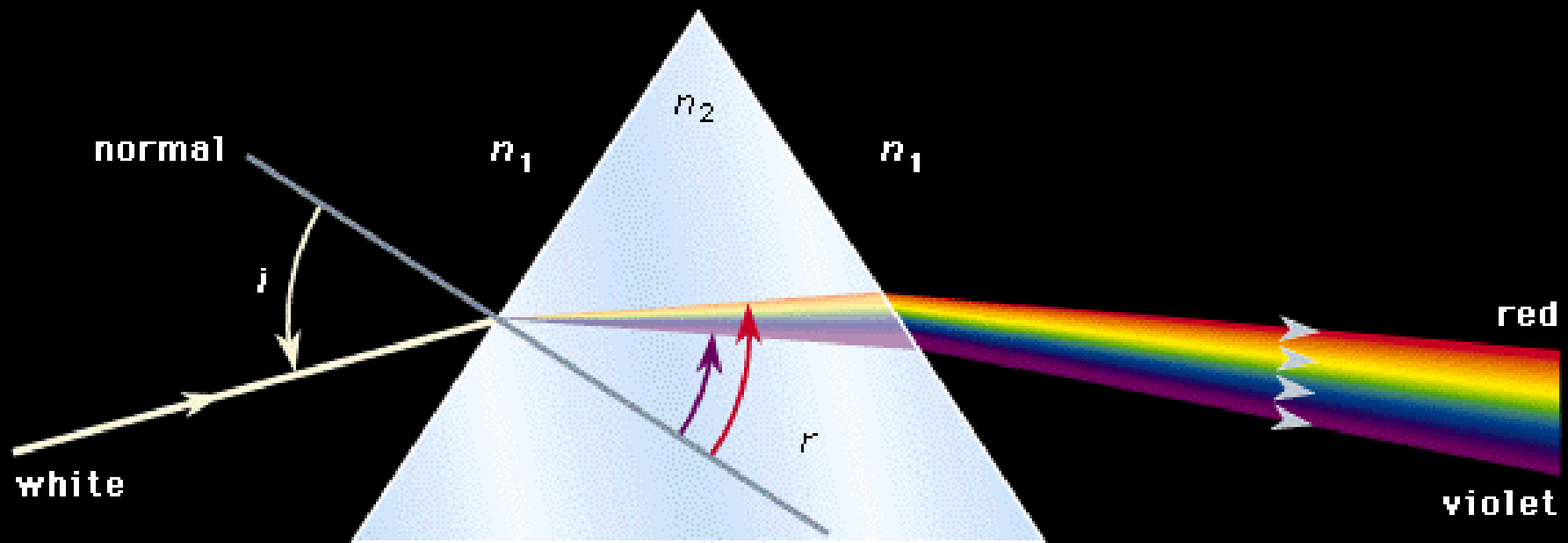
Magnet = Convex Lens + Prism



Optics: focus



Prism

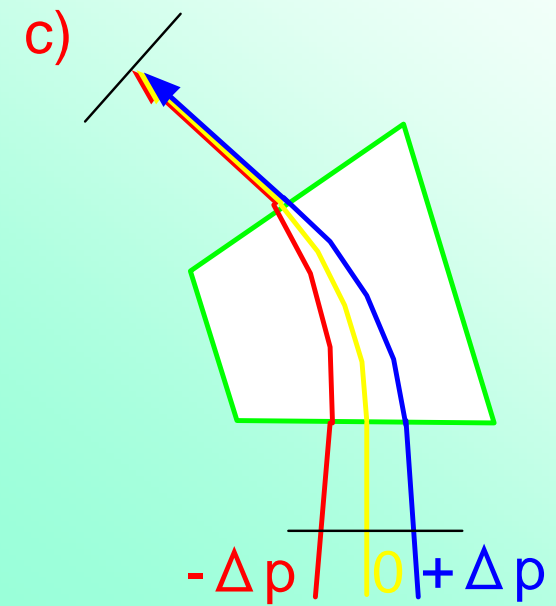
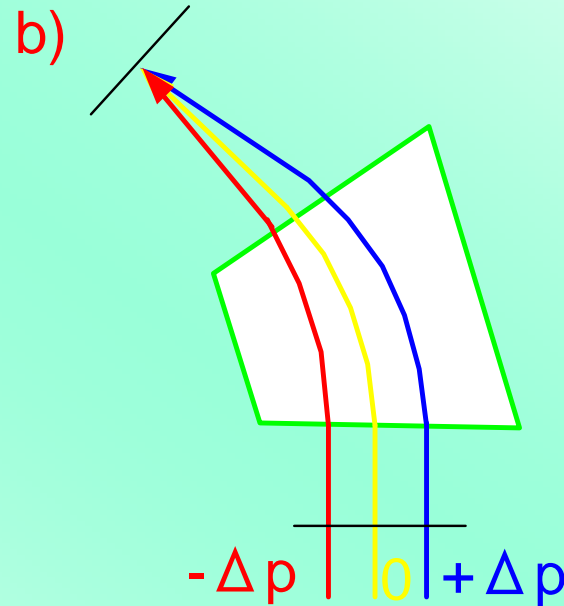
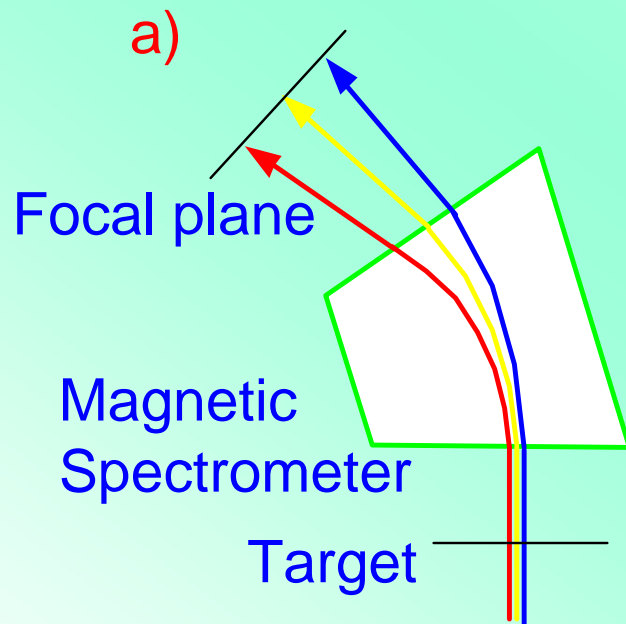


The angles i and r that the rays make with the normal are the angles of incidence and refraction. Because n_2 depends upon wavelength, the incident white ray separates into its constituent colours upon refraction, with deviation of the red ray the least and the violet ray the most.

Matching Techniques

Y. Fujita et al., N.I.M. B 126 (1997) 274.

H. Fujita et al., N.I.M. A 484 (2002) 17.



*Achromatic beam
transportation*

$\Delta E \sim 200$ keV
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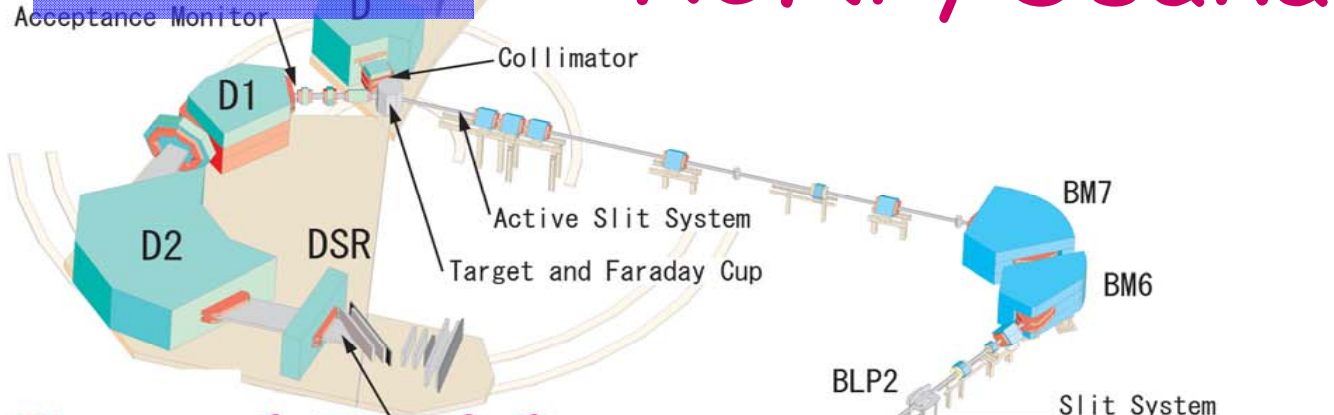
$\Delta E \sim 35$ keV
Horiz. angle resolution
 $\Delta\theta_{\text{sc}} > 15$ mrad

*Angular dispersion
matching*

$\Delta\theta_{\text{sc}} \sim 5$ mrad

$\Delta E = 30 \text{ keV}$

RCNP, Osaka Univ.



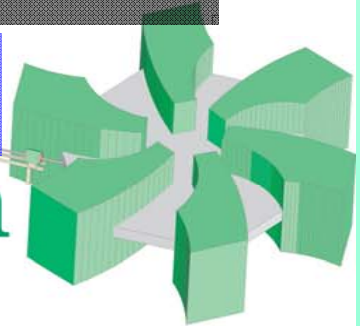
Grand Raiden

WS Beam Line

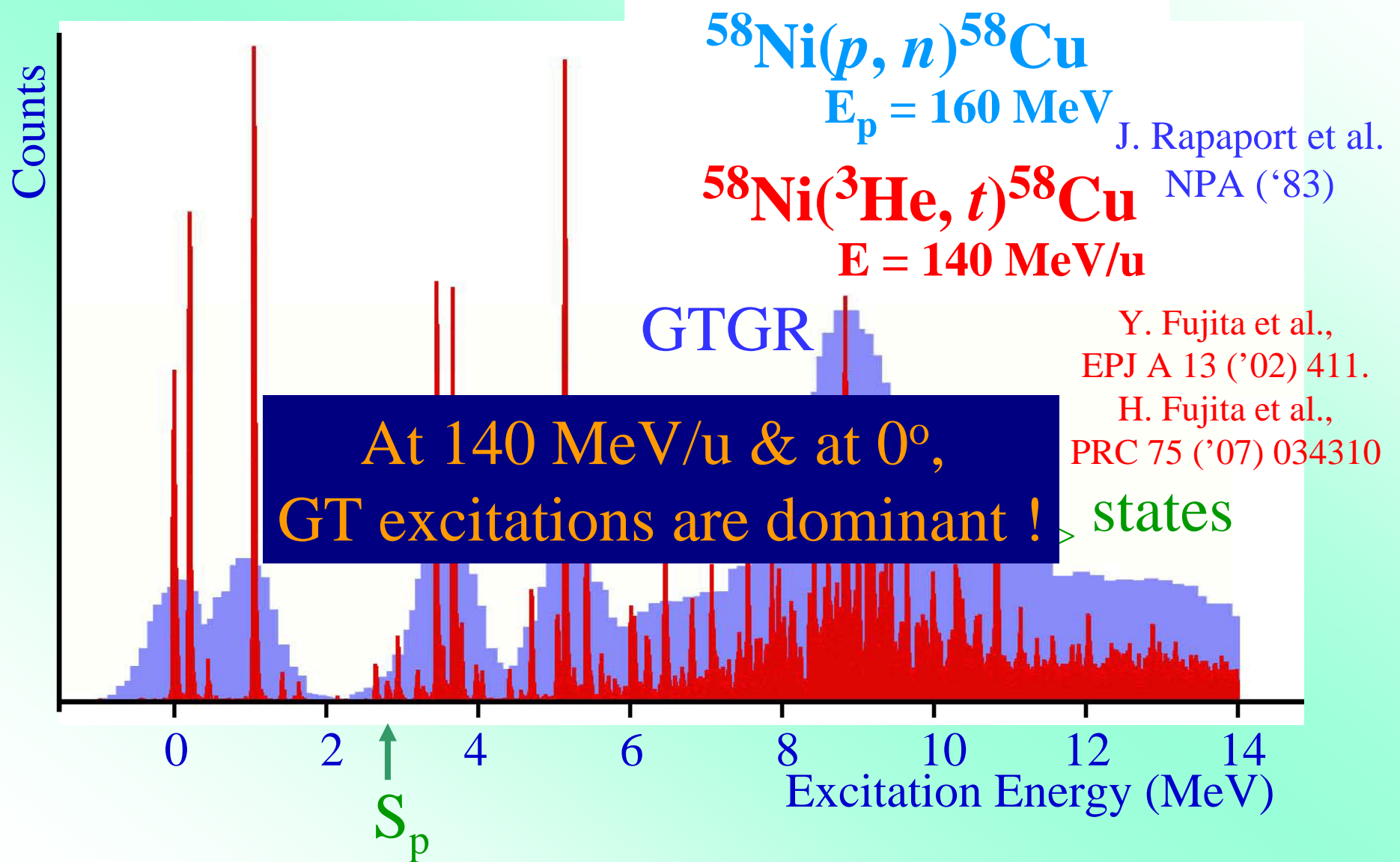
Dispersion Matching Techniques were applied!

$\Delta E = 150 \text{ keV}$

Ring Cyclotron

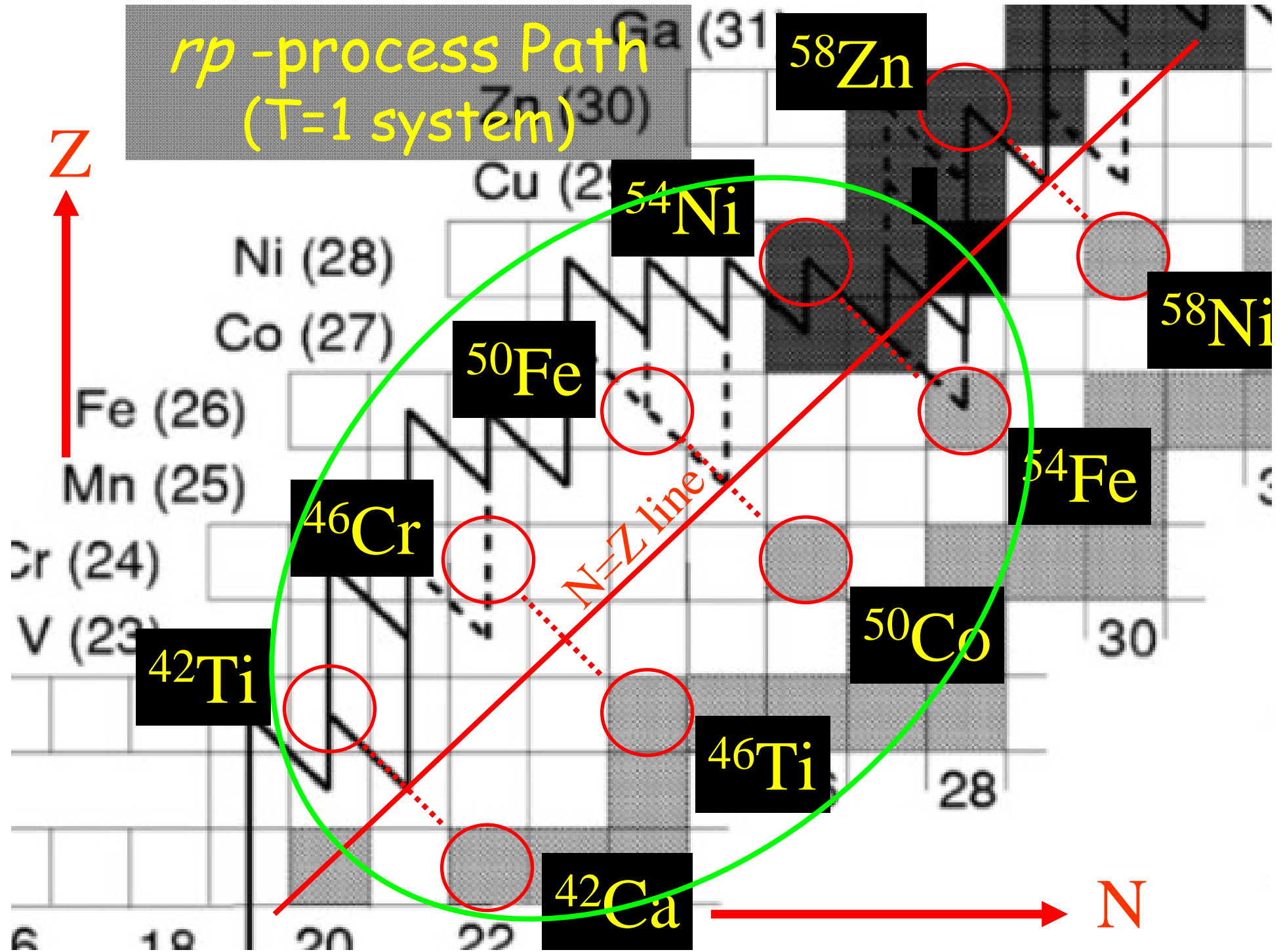


Comparison of (p, n) and (³He, t) 0° spectra

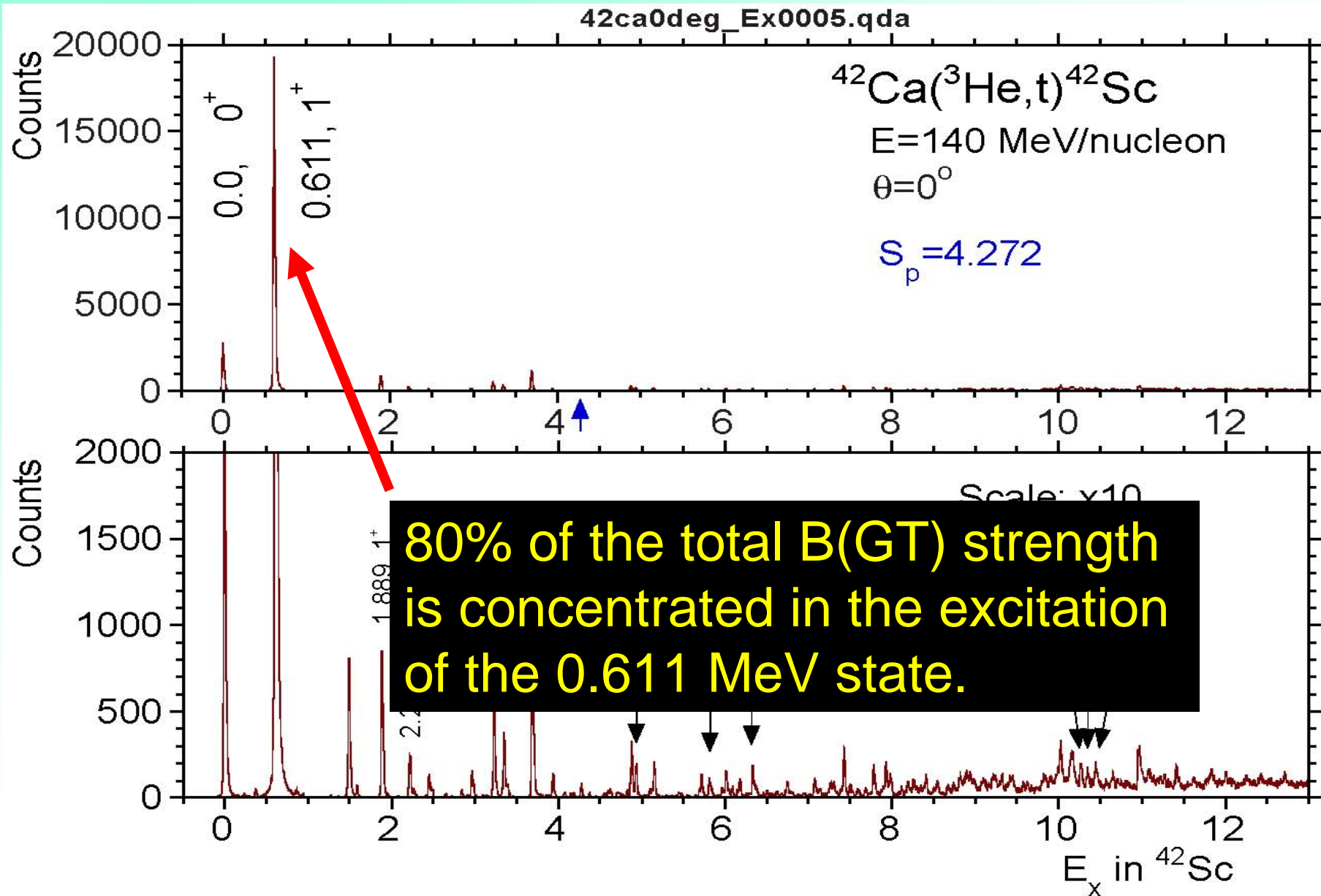


****GT transitions in each nucleus are
UNIQUE!**

- *pf*-shell nuclei -

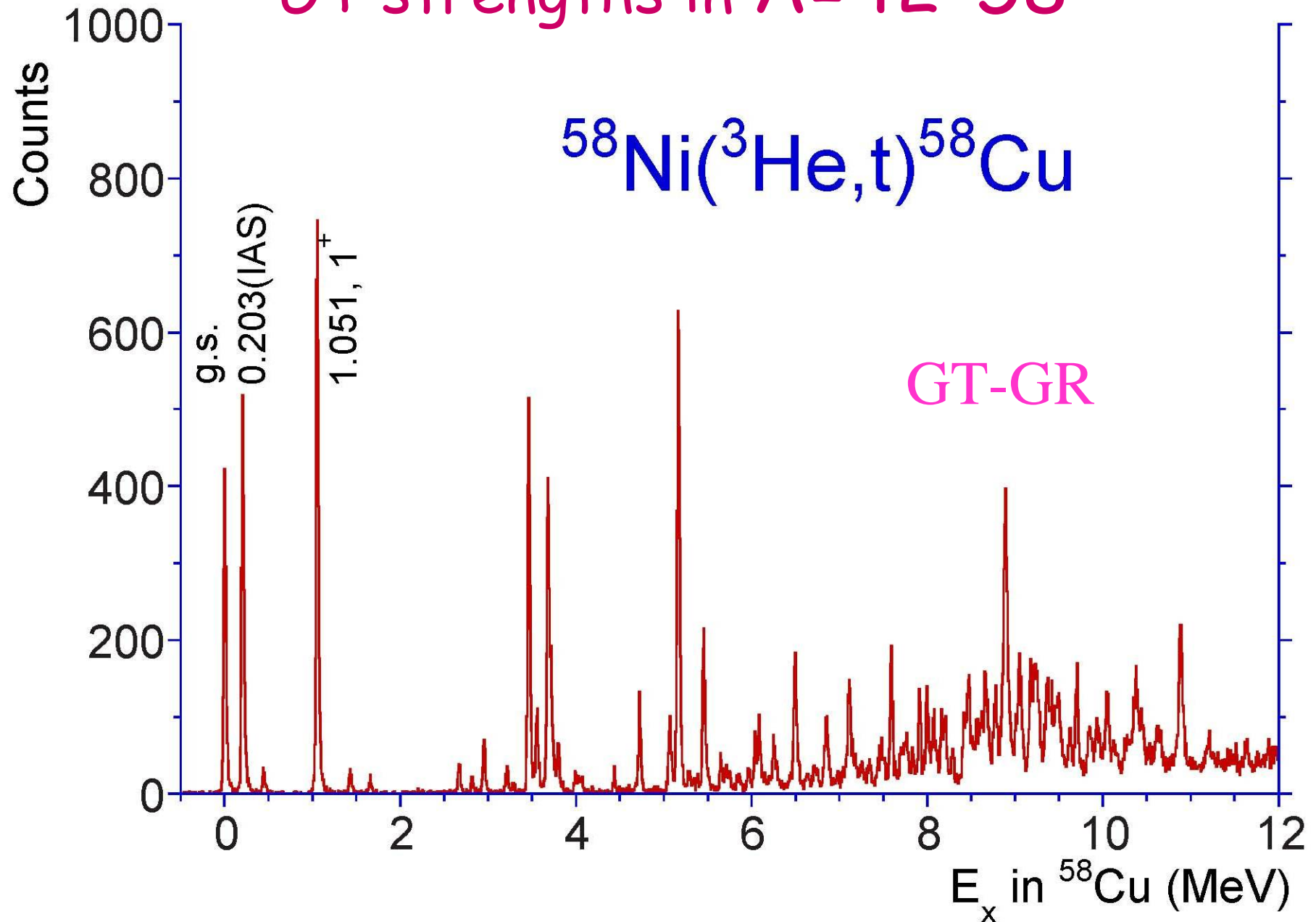


$^{42}\text{Ca}(^3\text{He},t)^{42}\text{Sc}$ in 2 scales



GT strengths in $A=42-58$

$^{58}\text{Ni}({}^3\text{He},t){}^{58}\text{Cu}$

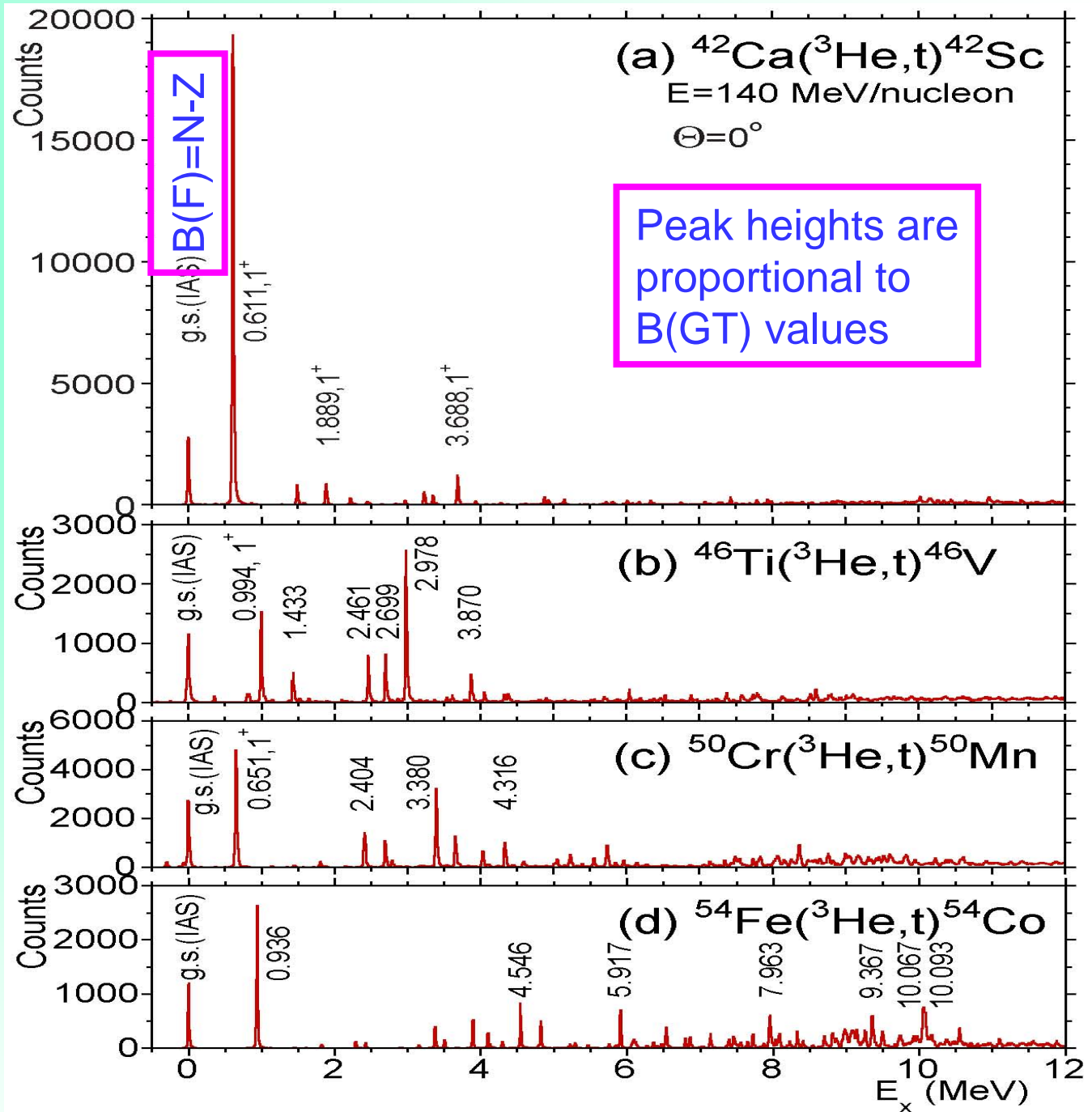


GT states
in
 $A=42-54$
 $T_z=0$ nuclei

T. Adachi et al.
PRC '06

Y. Fujita et al.
PRL '05

T. Adachi et al.
PRC '12

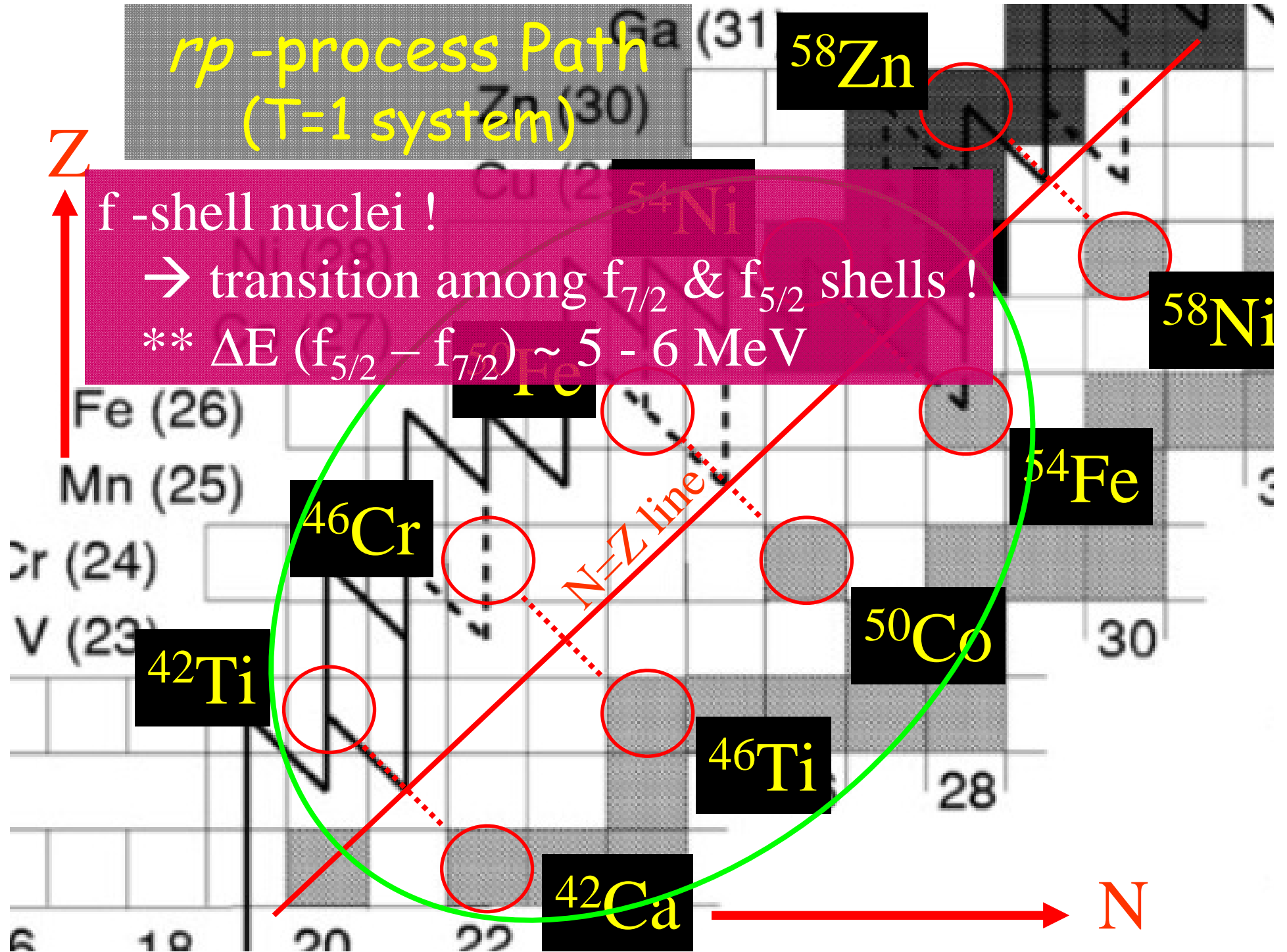


rp-process Path
(T=1 system)

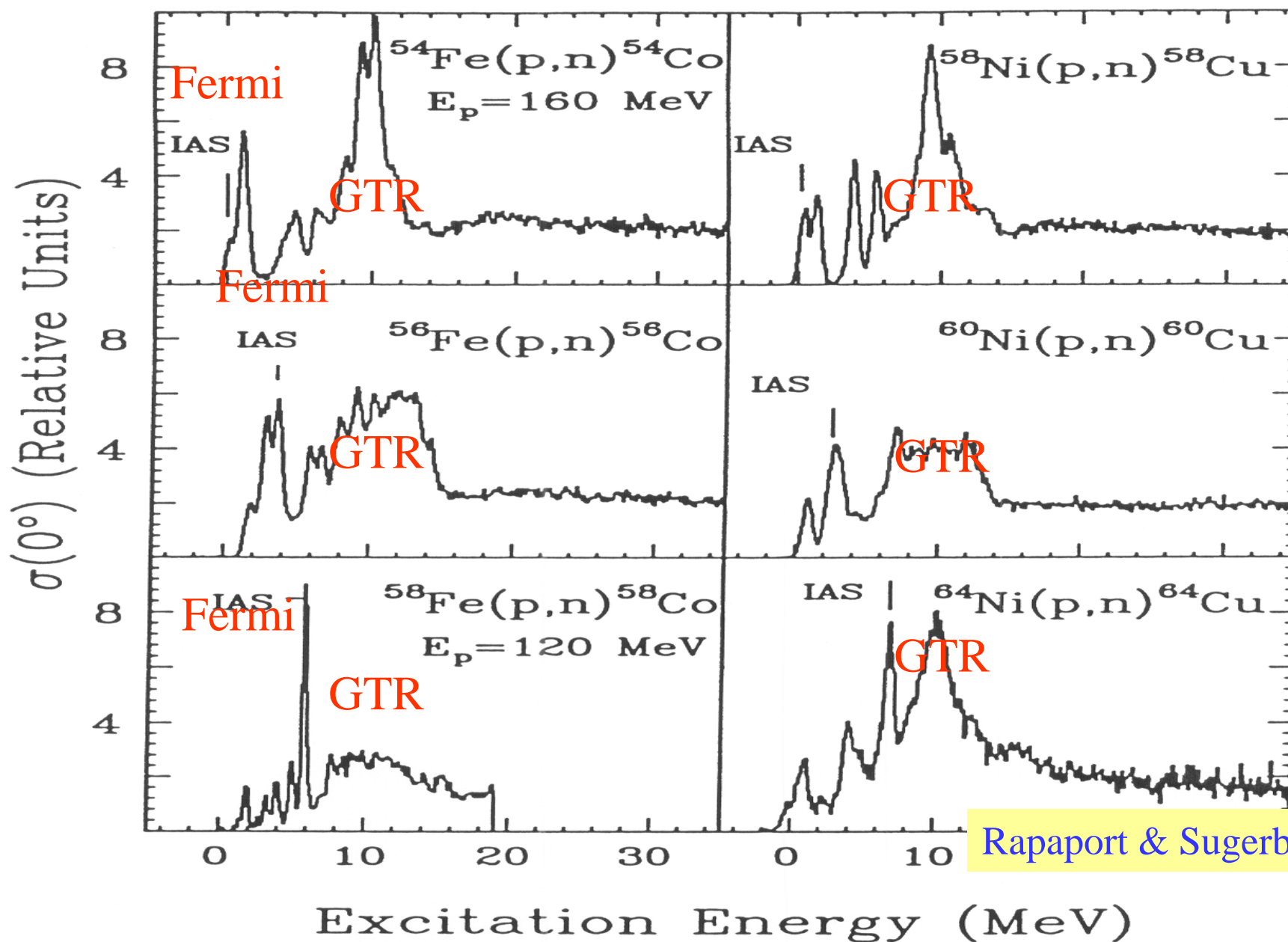
f-shell nuclei !

→ transition among $f_{7/2}$ & $f_{5/2}$ shells !

** $\Delta E (f_{5/2} - f_{7/2}) \sim 5 - 6 \text{ MeV}$

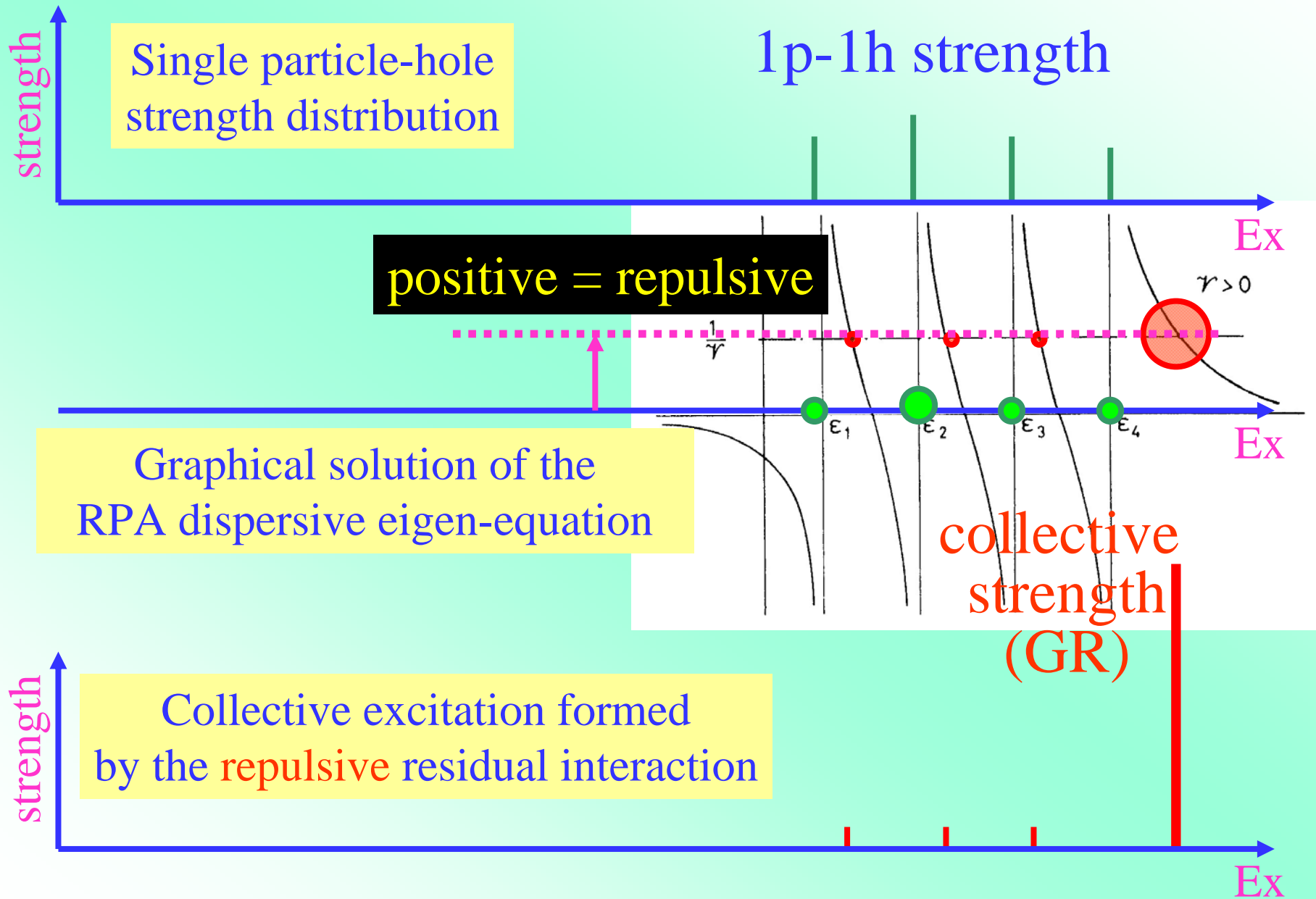


(p, n) spectra for Fe and Ni Isotopes

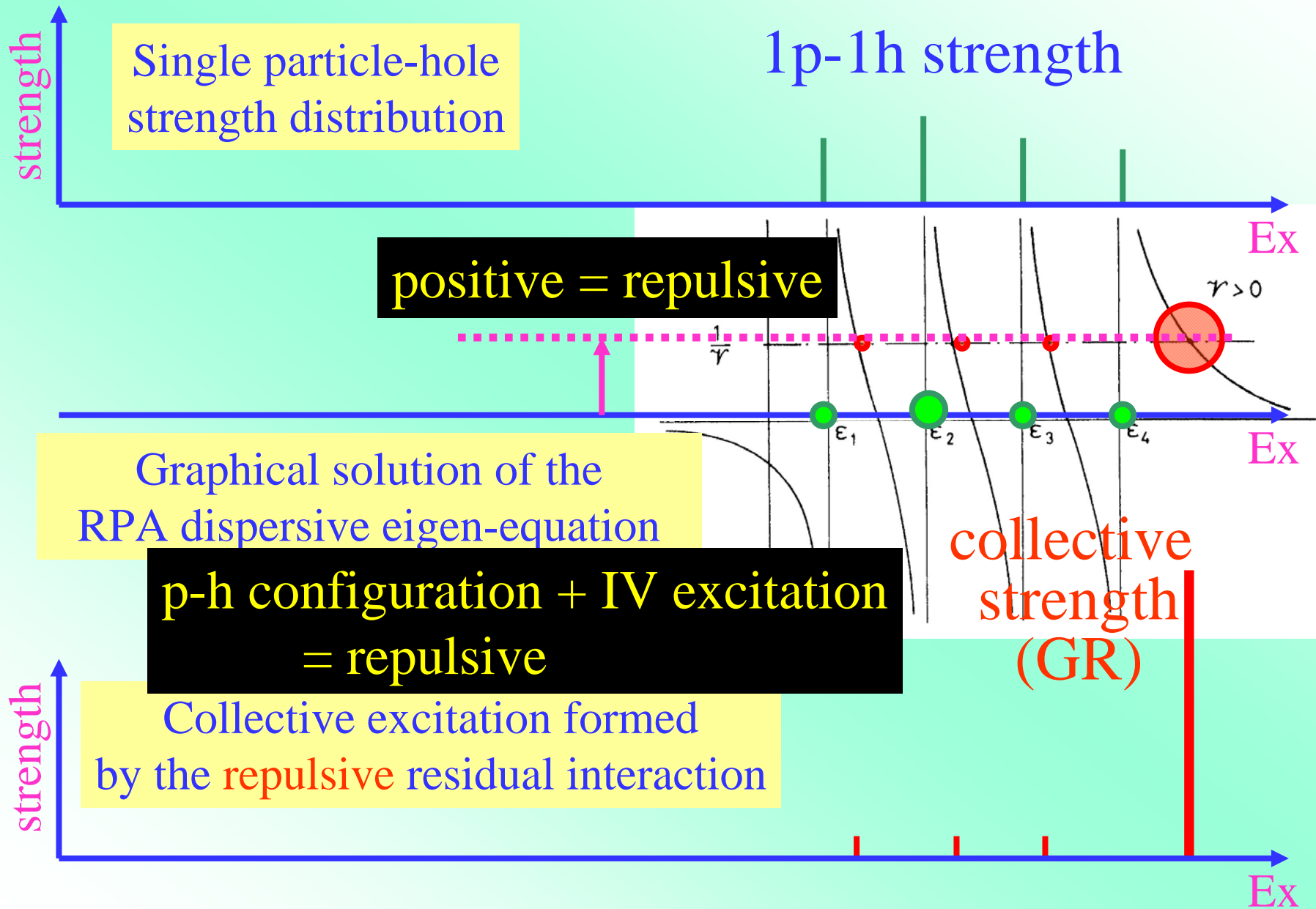


Rapaport & Sugerbaker

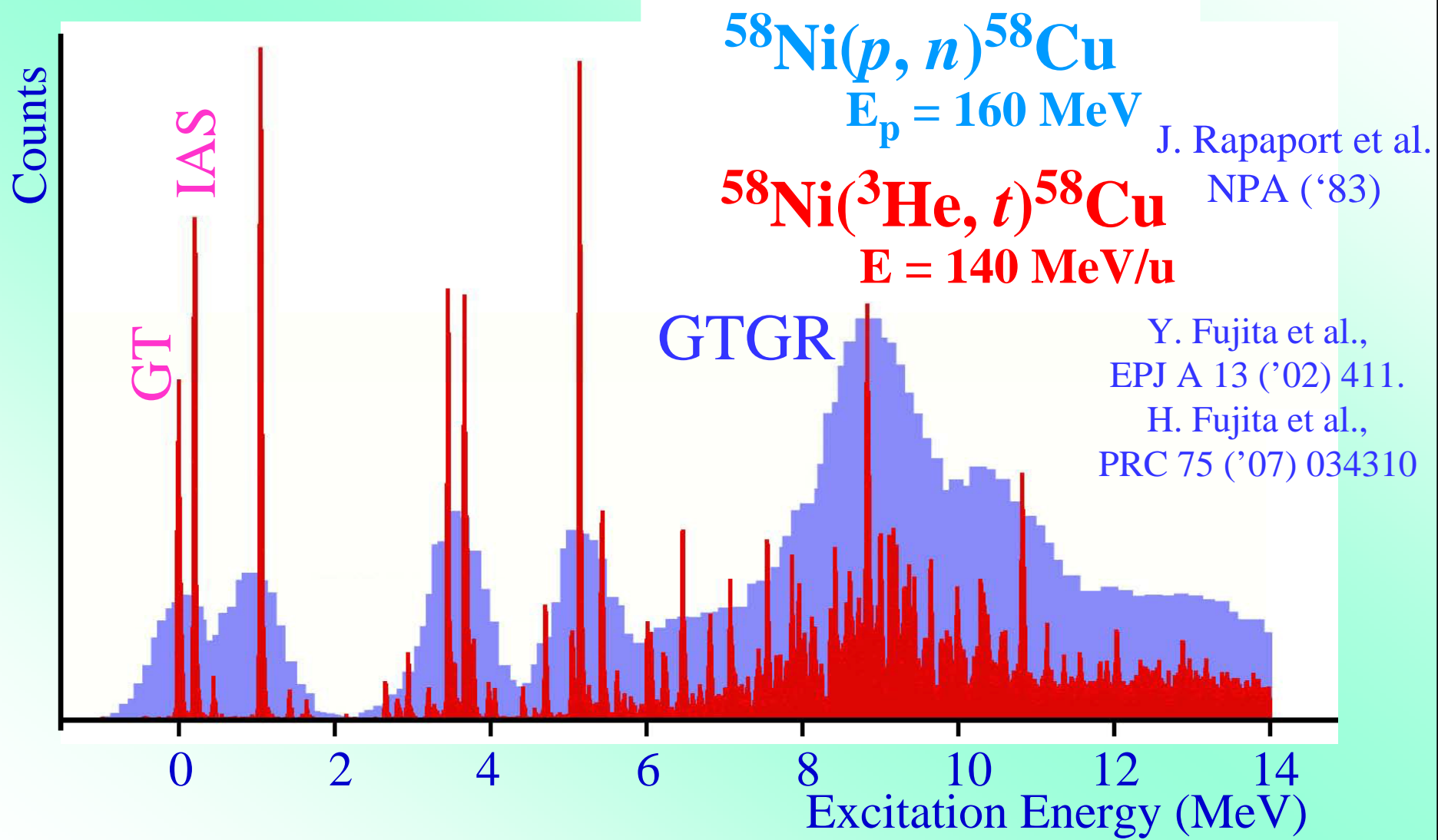
Role of Residual Int. (repulsive)



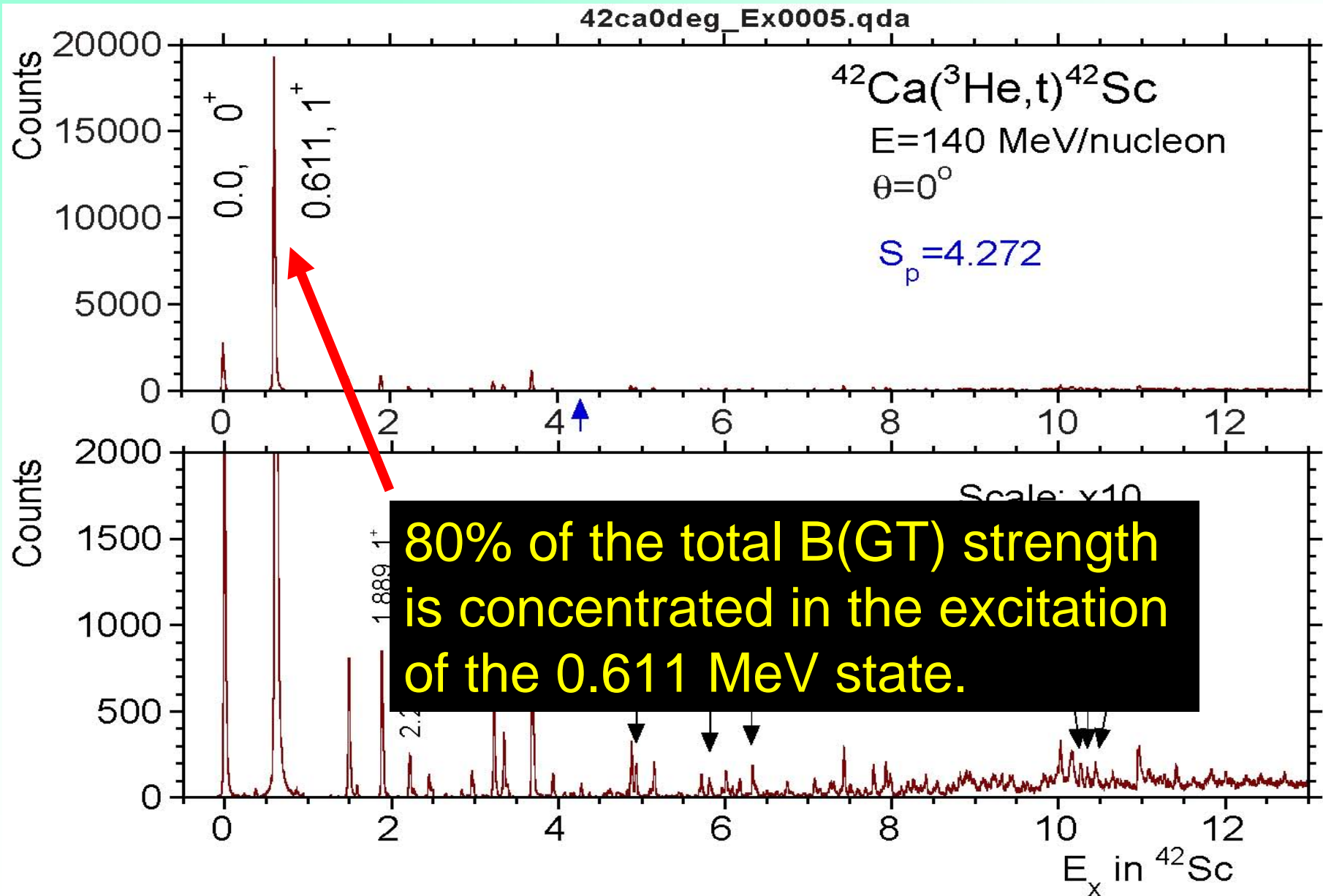
Role of Residual Int. (repulsive)



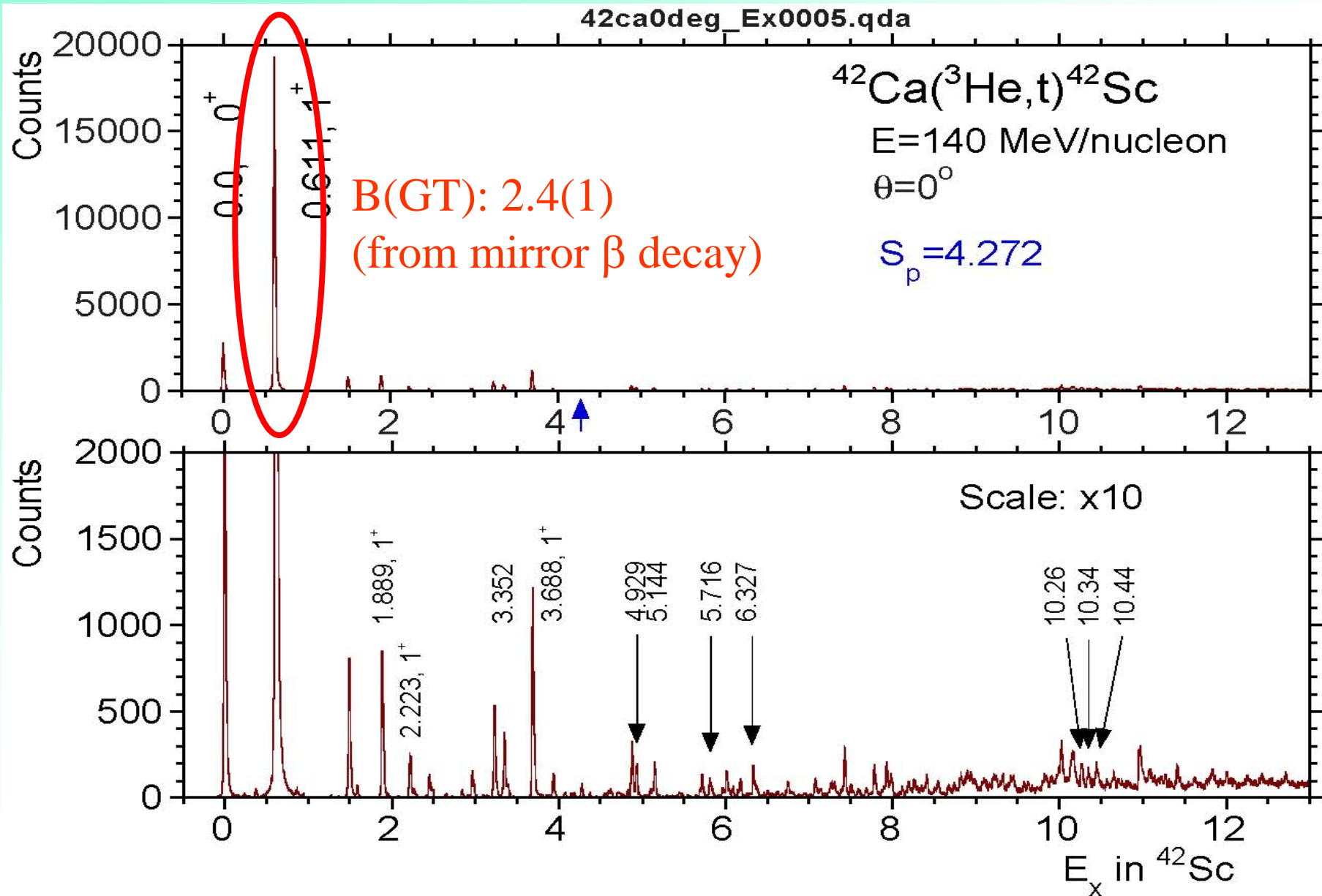
Comparison of (p, n) and (³He, t) 0° spectra



$^{42}\text{Ca}(^3\text{He},t)^{42}\text{Sc}$ in 2 scales



$^{42}\text{Ca}(^3\text{He},t)^{42}\text{Sc}$ in 2 scales

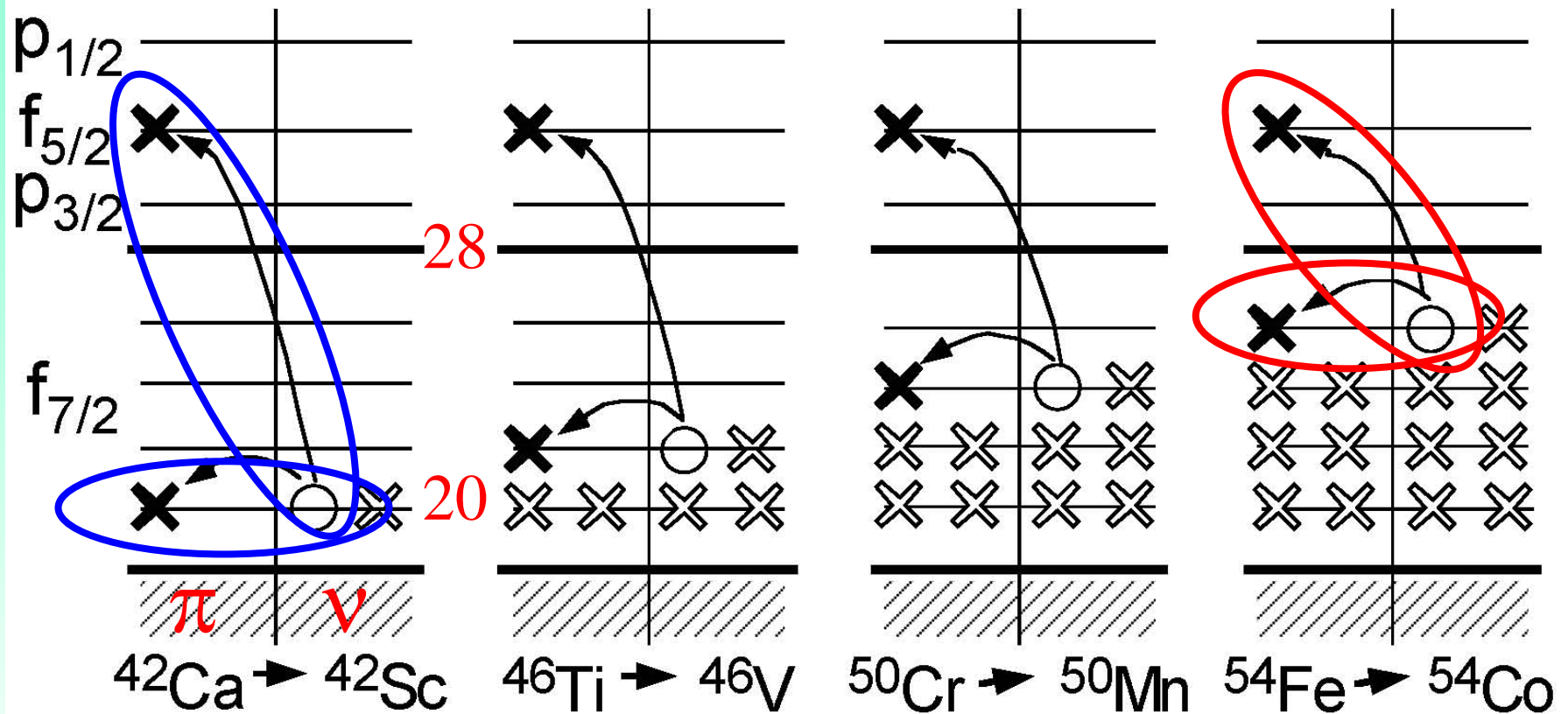


**Residual Interaction

- poor man's understanding! -

coupled pendulum

SM Configurations of GT transitions



Target nuclei: $N = Z + 2$ ($T_z = +1$)

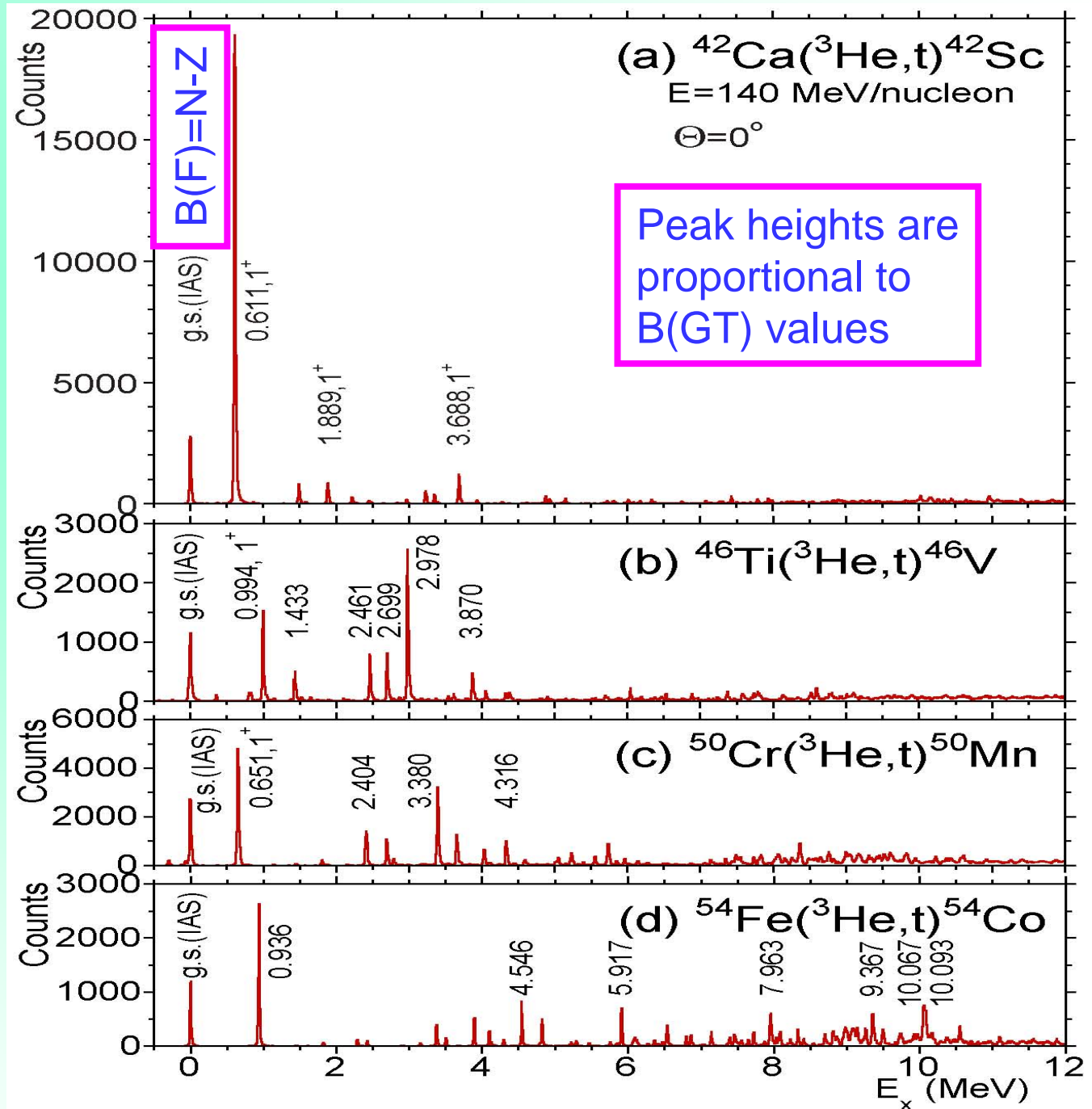
Final nuclei : $N = Z$ ($T_z = 0$)

GT states
in
 $A=42-54$
 $T_z=0$ nuclei

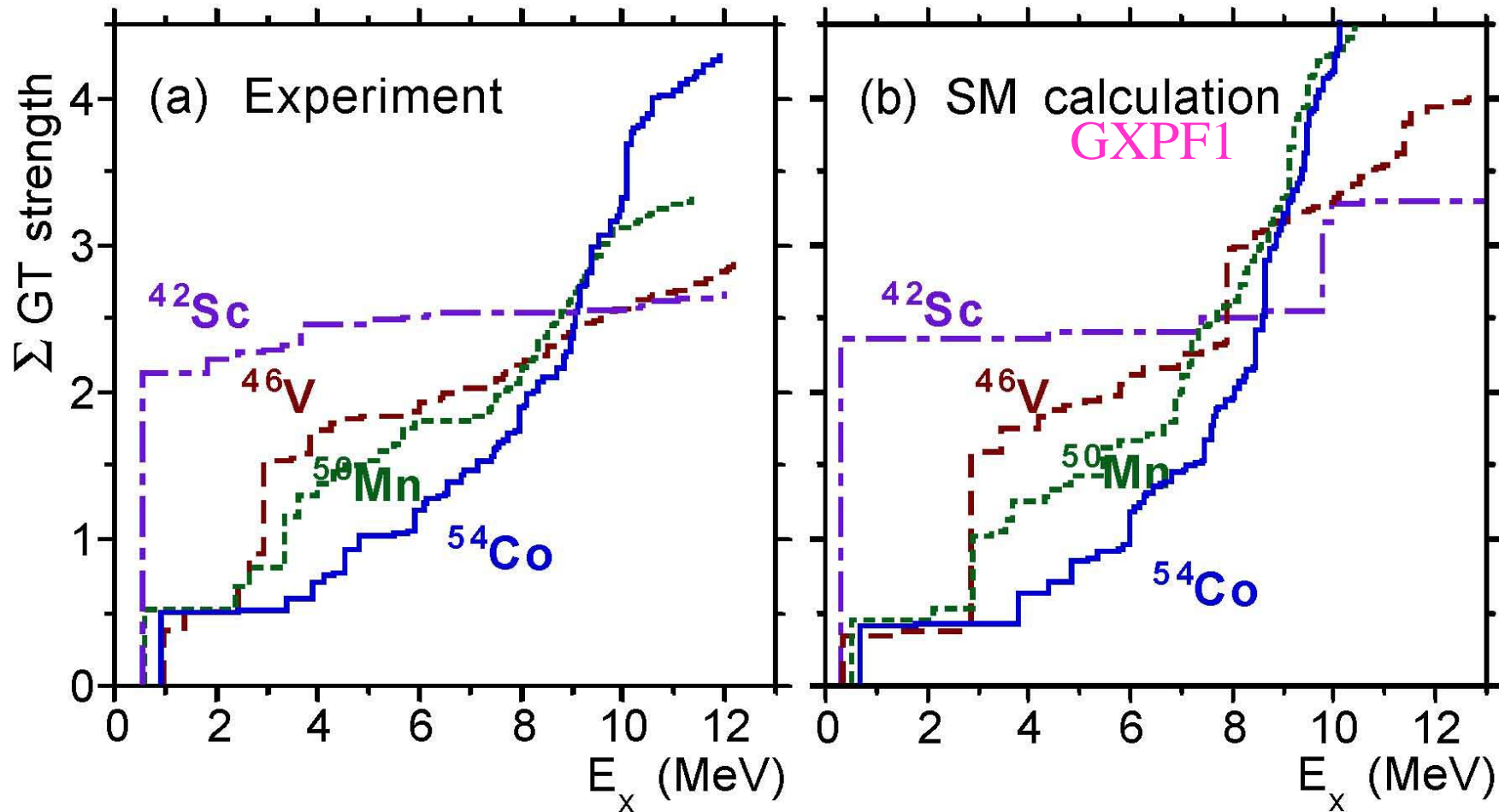
T. Adachi et al.
PRC '06

Y. Fujita et al.
PRL '05

T. Adachi et al.
PRC '12



GT-strength: Cumulative Sum

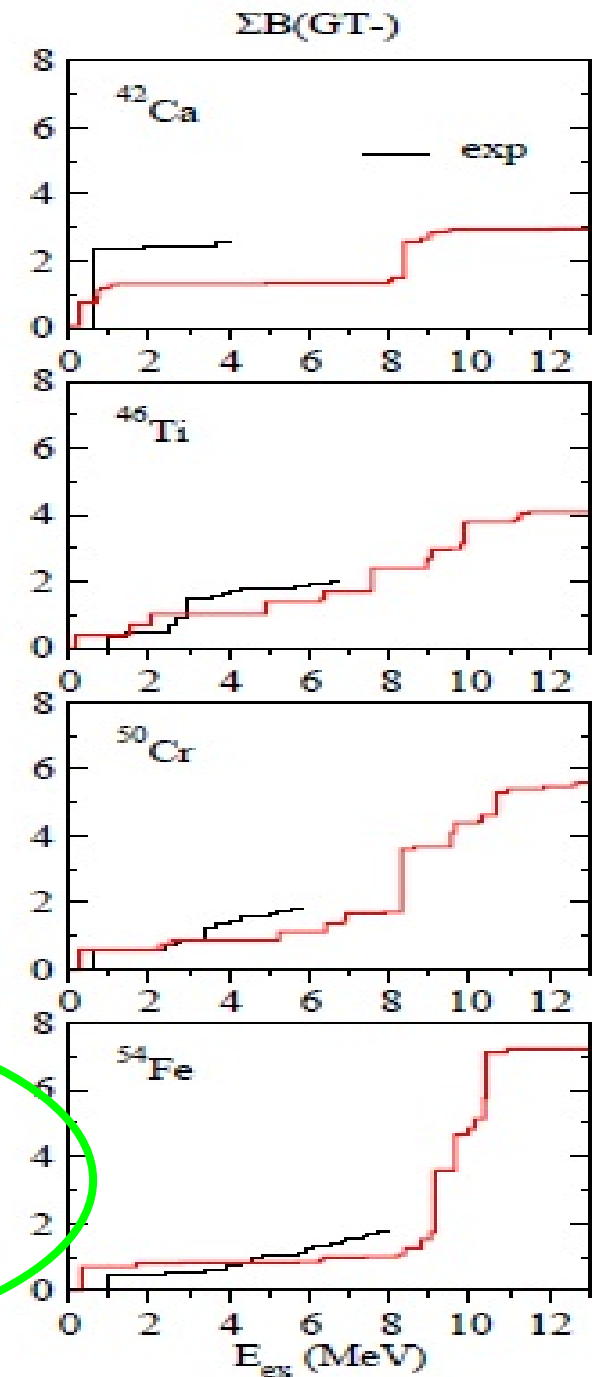
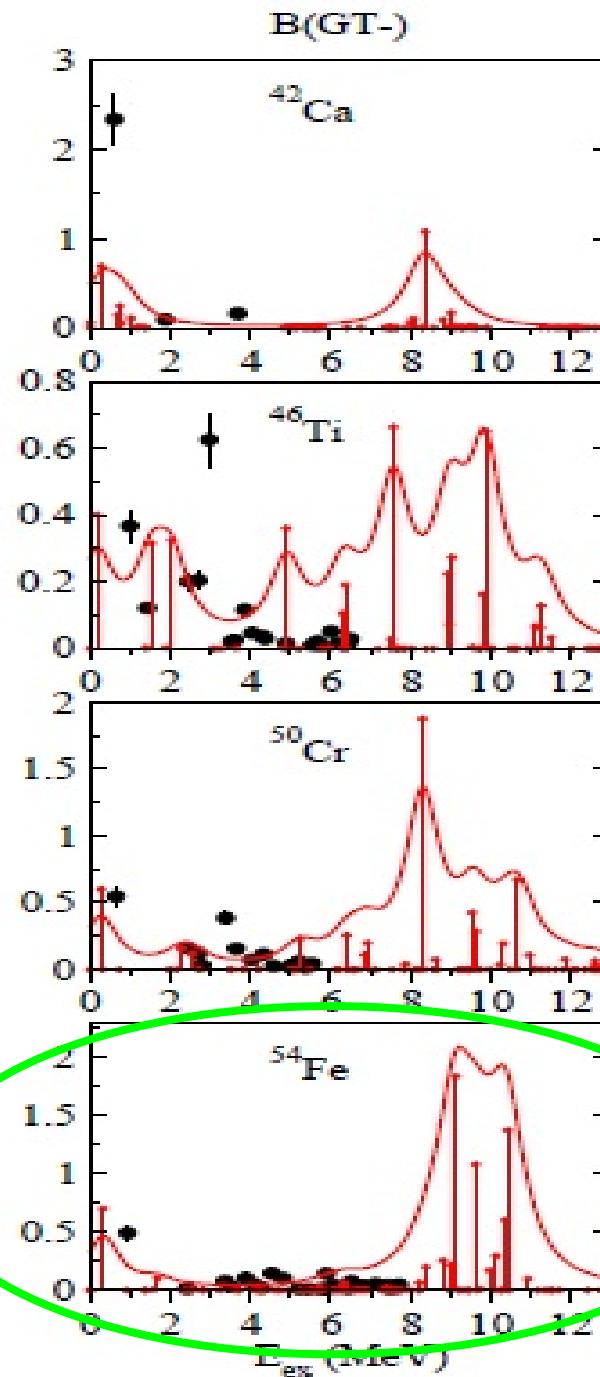


M. Homma et al.

QRPA calculations

using Skyrme int. SG1
(with IV pairing corr.)

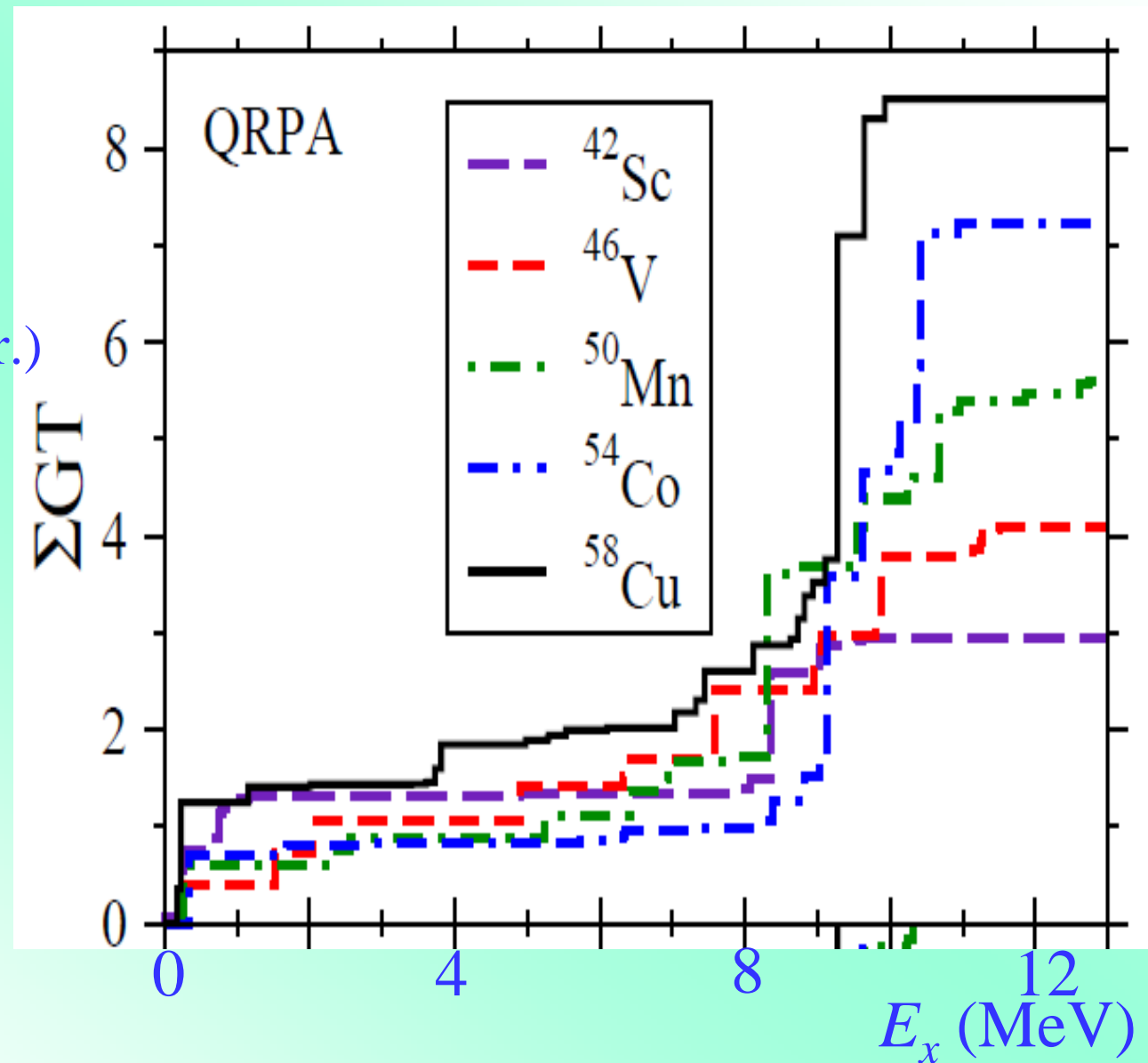
Calculation by
P. Sarrigren,
CSIC, Madrid



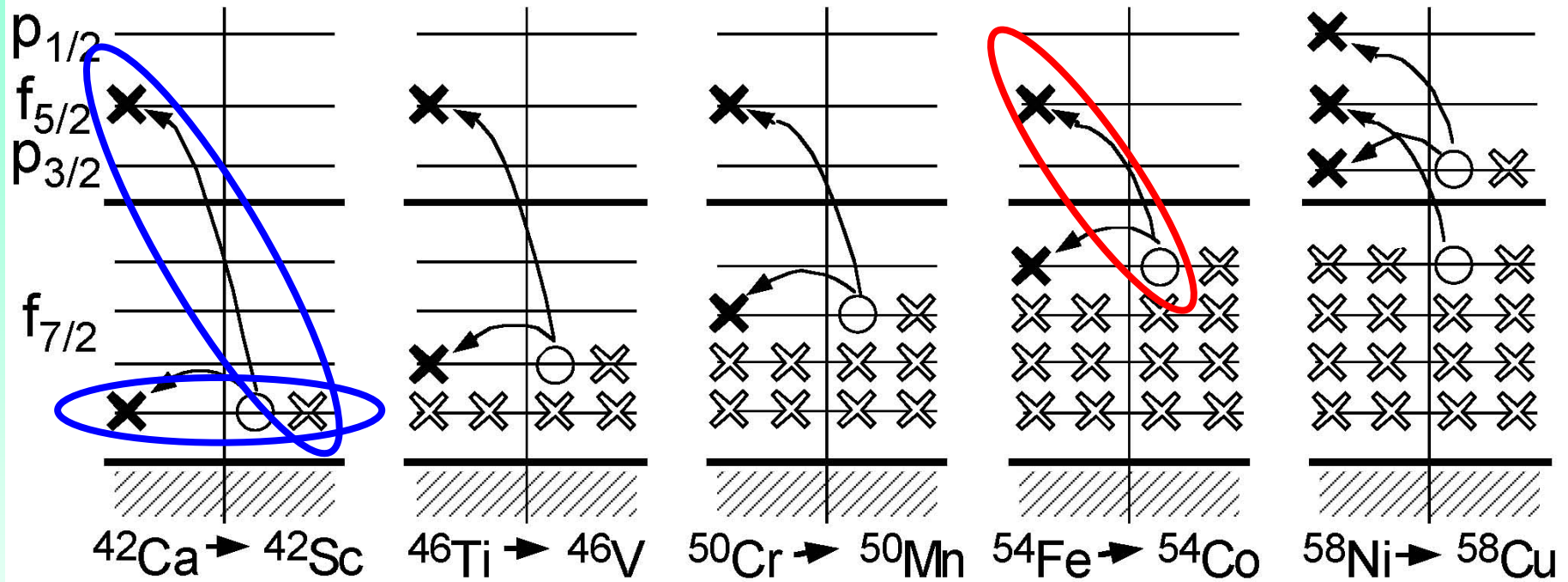
QRPA calculations

using Skyrme int.
(with IV pairing corr.)

Calculation by
P. Sarrigren,
CSIC, Madrid

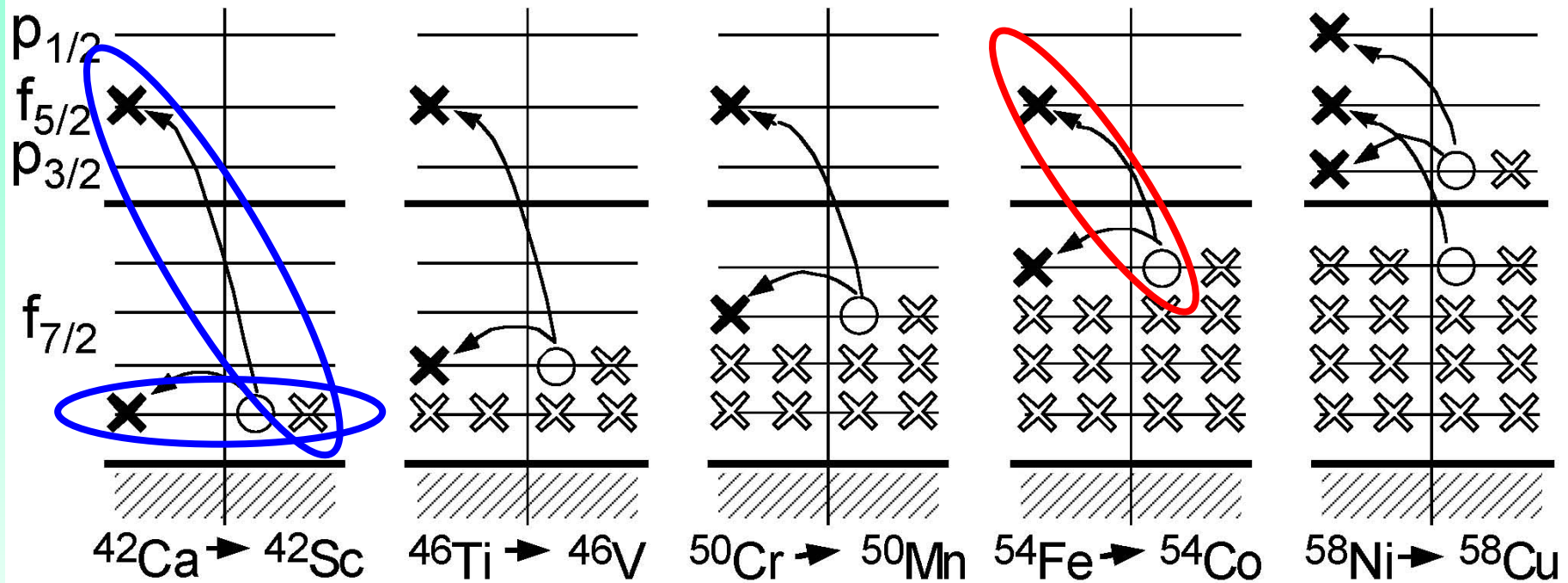


SM Configurations of GT transitions



particle-particle int. (attractive) → particle-hole int. (repulsive)

SM Configurations of GT transitions



π -p - ν -p configurations
sensitive to IS pairing int.

and

it is attractive

(spin-triplet, IS int. is stronger

than spin-singlet, IV int.)

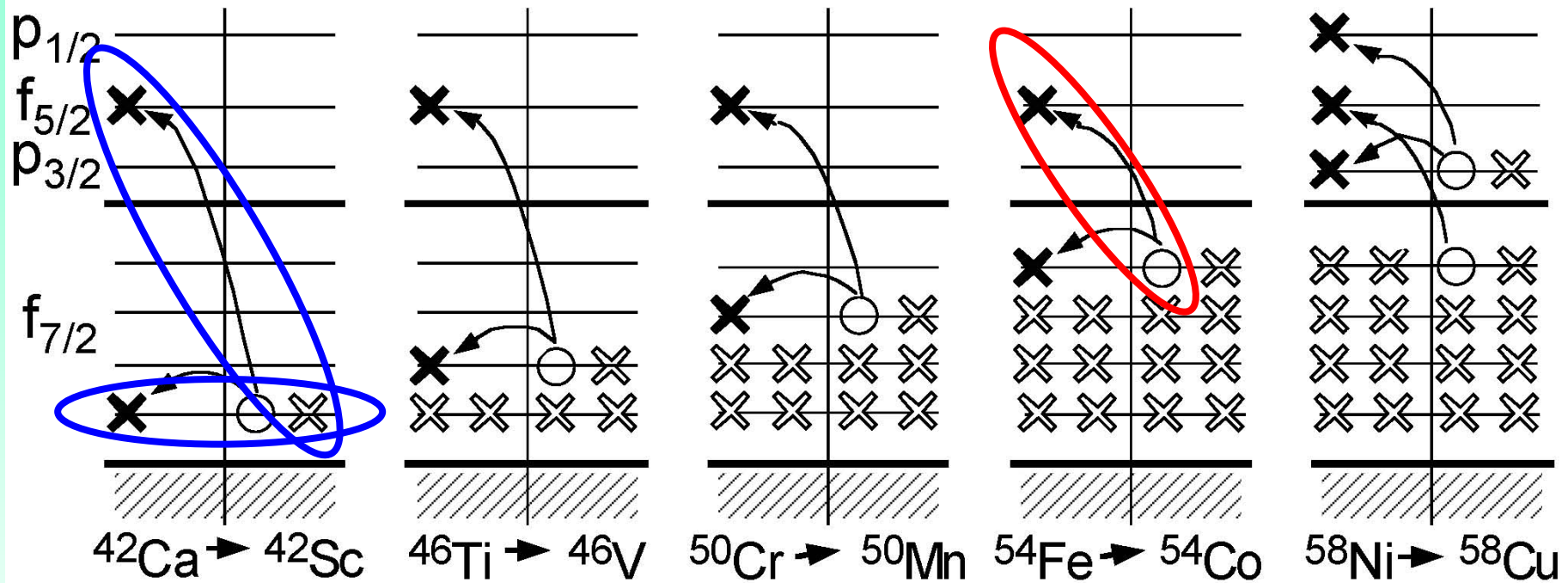
particle-hole configurations
+

IV-type excitation ($\sigma\tau$)

→ repulsive

by Engel, Bertsch, Macchiavelli

SM Configurations of GT transitions



particle-particle int. (attractive) →
 (T=0, IS p-n int. is attractive)

particle-hole int. (repulsive)

**Isoscalar interaction
 can play important roles !**

Cooperative with the repulsive
 nature of $\sigma\tau$ int. !

GT strength Calculations: HFB+QRPA + pairing int.

C.L. Bai, H. Sagawa et al., PL B 719 (2013) 116

The density dependent contact pairing interactions are adopted for both $T = 1$ and $T = 0$ channels,

$$\text{IV} \quad V_{T=1}(\mathbf{r}_1, \mathbf{r}_2) = V_0 \frac{1 - P_\sigma}{2} \left(1 - \frac{\rho(\mathbf{r})}{\rho_0} \right) \delta(\mathbf{r}_1 - \mathbf{r}_2), \quad (1)$$

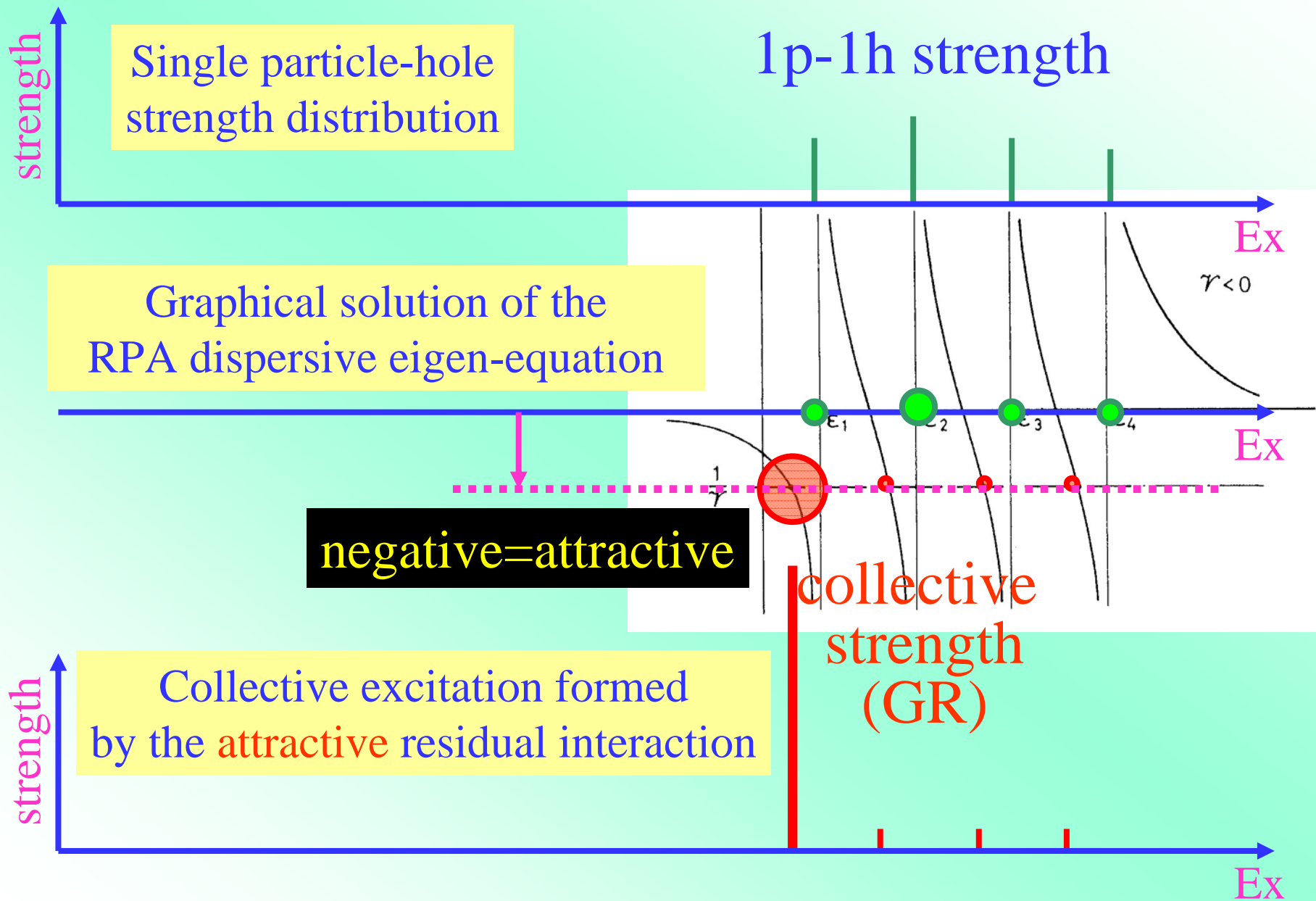
$$\text{IS} \quad V_{T=0}(\mathbf{r}_1, \mathbf{r}_2) = f V_0 \frac{1 + P_\sigma}{2} \left(1 - \frac{\rho(\mathbf{r})}{\rho_0} \right) \delta(\mathbf{r}_1 - \mathbf{r}_2), \quad (2)$$

Results (using Skyrme int. SGII)

at $f=0$: there is little strength in the lower energy part,

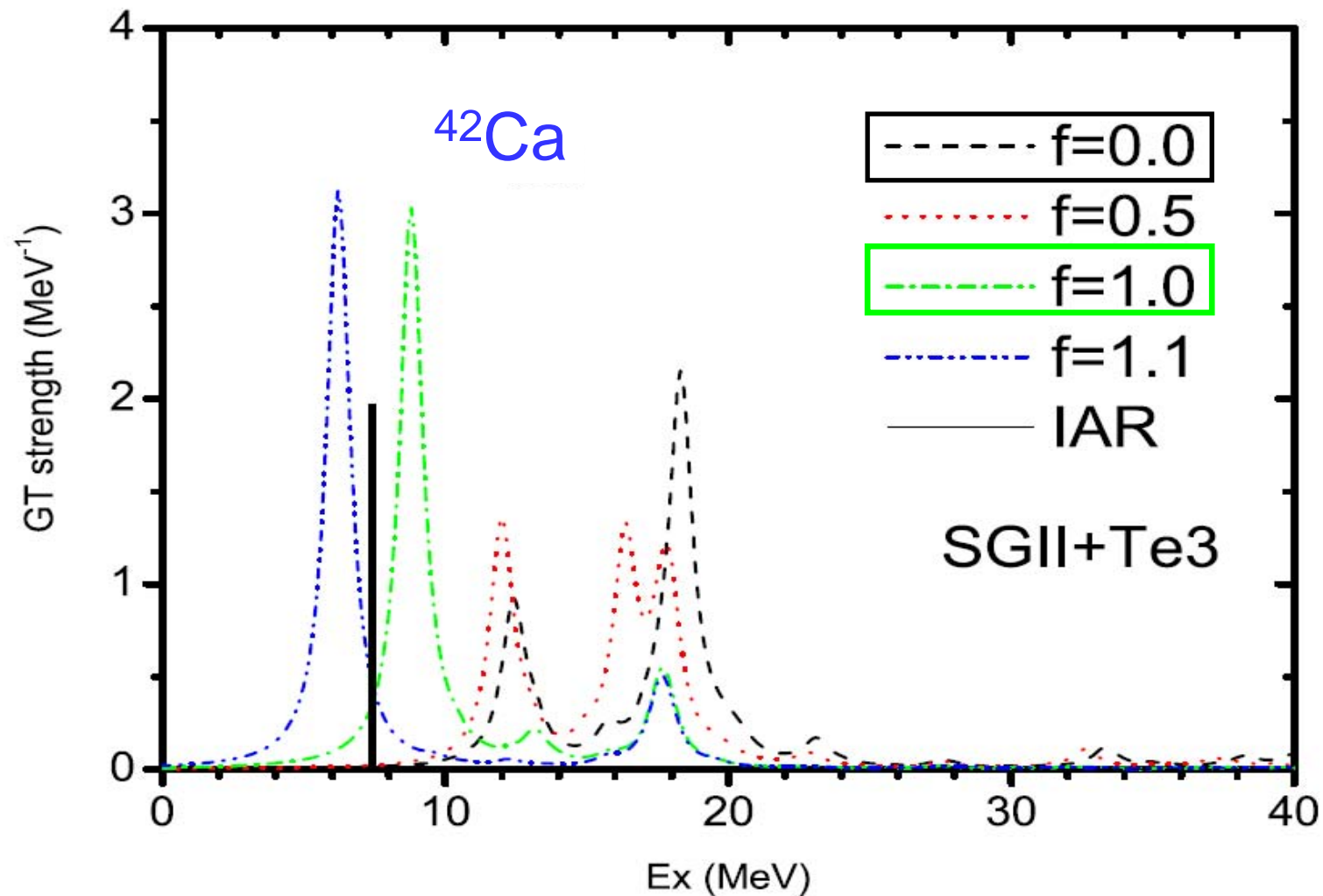
at $f=1.0 \sim 1.7$: coherent low-energy strength develops!

Role of Residual Int. (attractive)



QRPA-cal. GT-strength (with IS-int.)

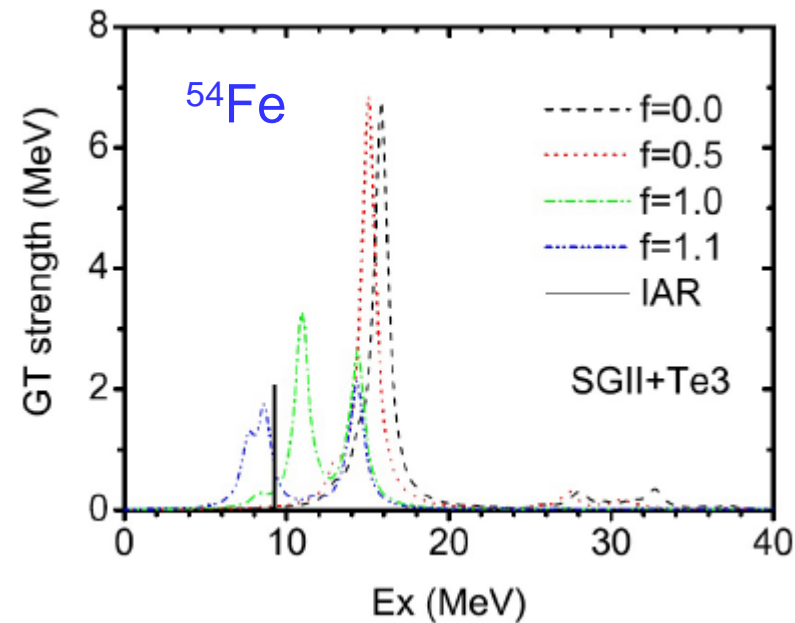
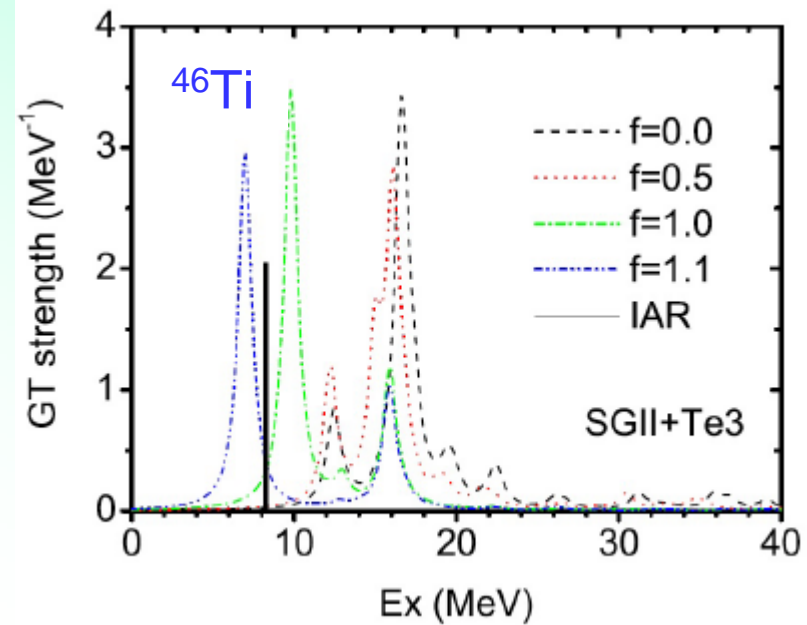
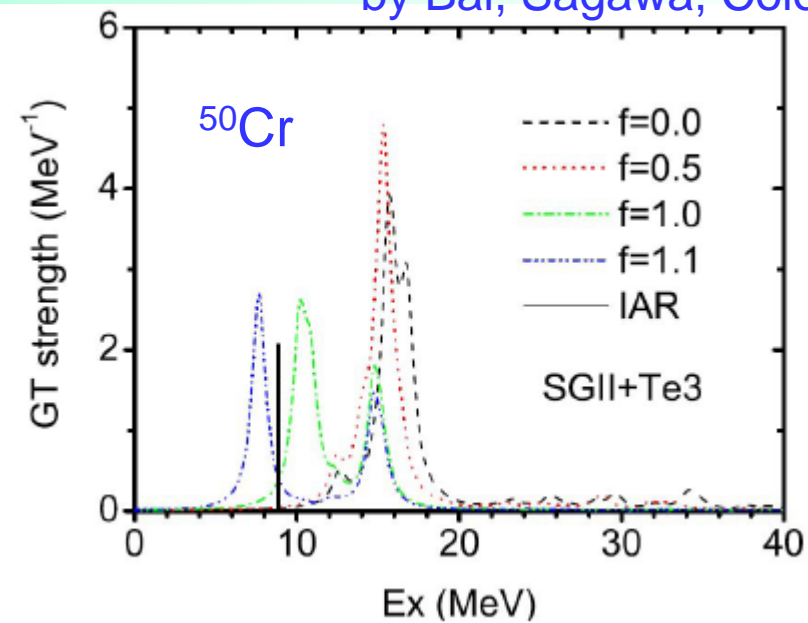
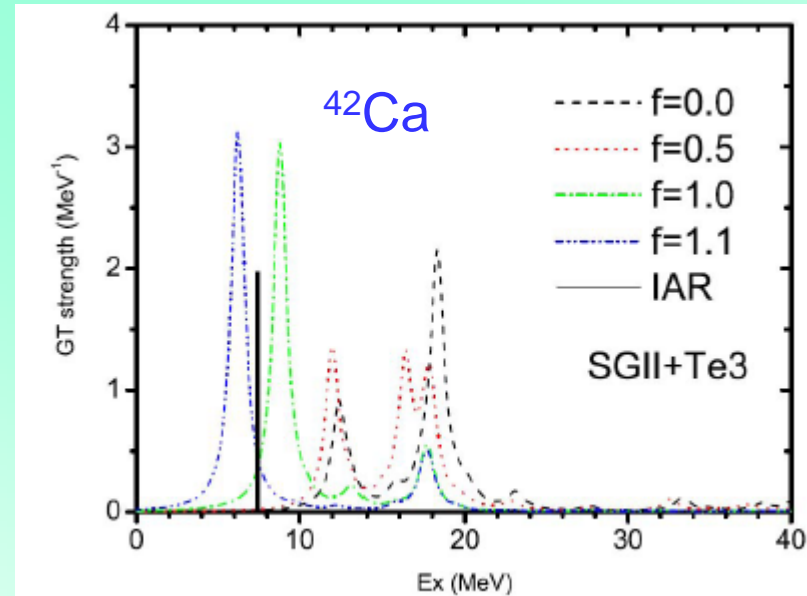
by Bai Sagawa Colo



$^{42}\text{Ca} \rightarrow ^{42}\text{Sc}$ (Q-value)

QRPA-cal. GT-strength (with IS-int.)

by Bai, Sagawa, Colo



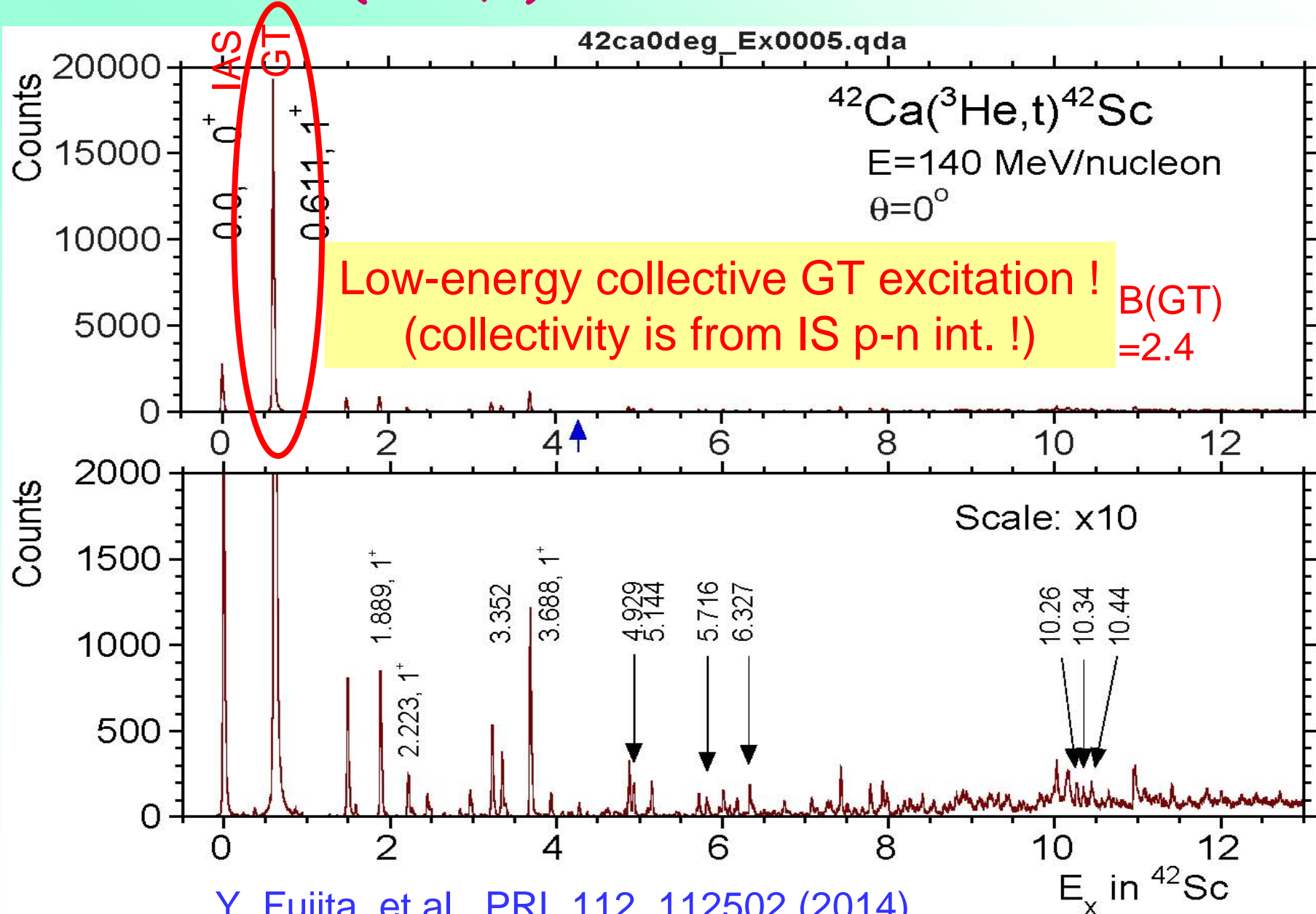
QRPA cal. including IS int.

C.L. Bai, H. Sagawa, G. Colo

f	Bnp			
0	1.34			
	neutron	proton	(Xupvn+Yunvp)	(Xupvn+Yunvp)*<p GT n>
	1f7/2	1f7/2	0.427	1.3689
0.5	2.051			
	1f7/2	1f7/2	0.432	1.384
1	4.75			
	1f5/2	1f7/2	0.053	0.2158
	1f7/2	1f5/2	0.129	0.474
	1f7/2	1f7/2	0.33	1.059

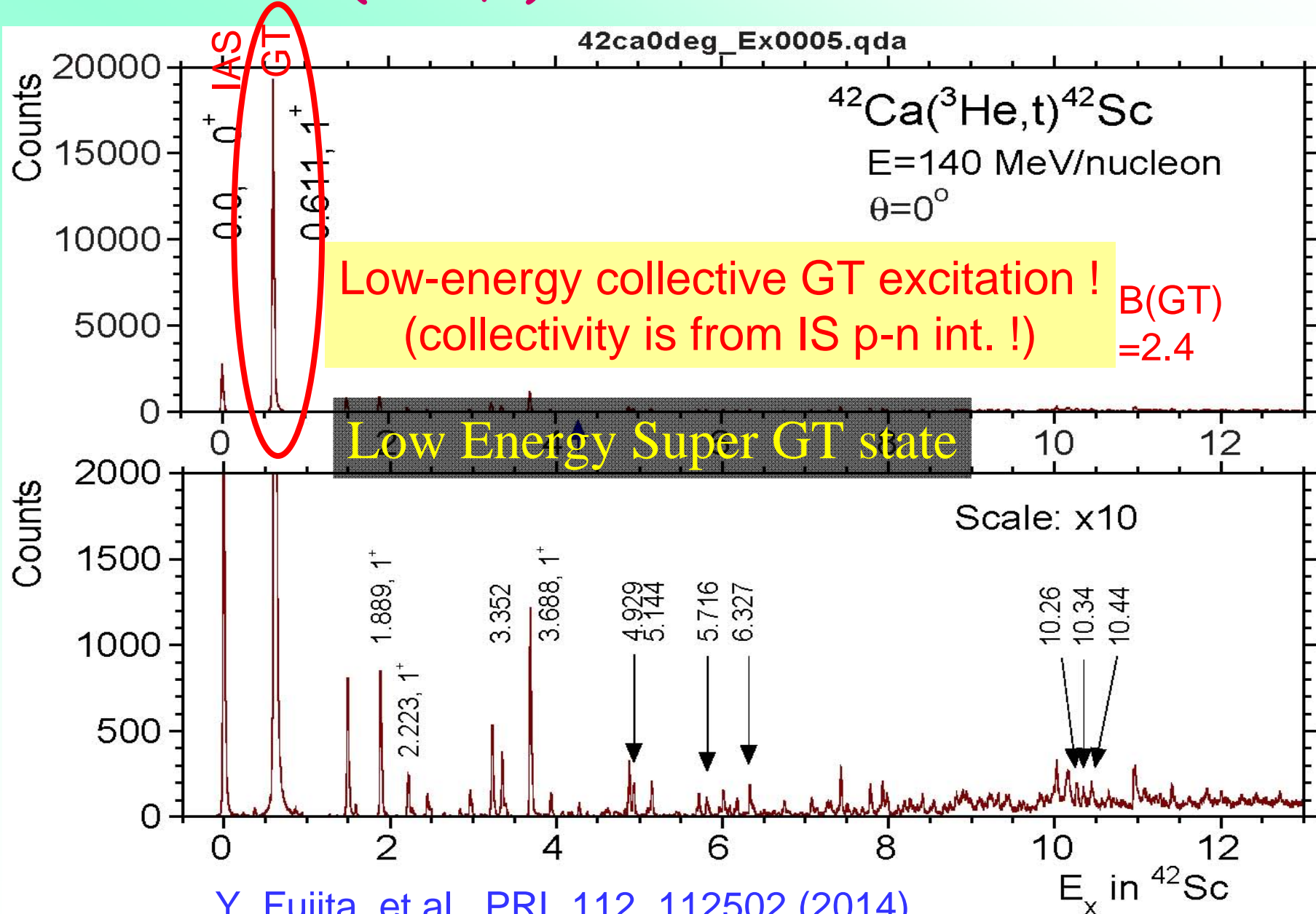
Configurations
are in phase!

$^{42}\text{Ca}(^3\text{He},t)^{42}\text{Sc}$ in 2 scales



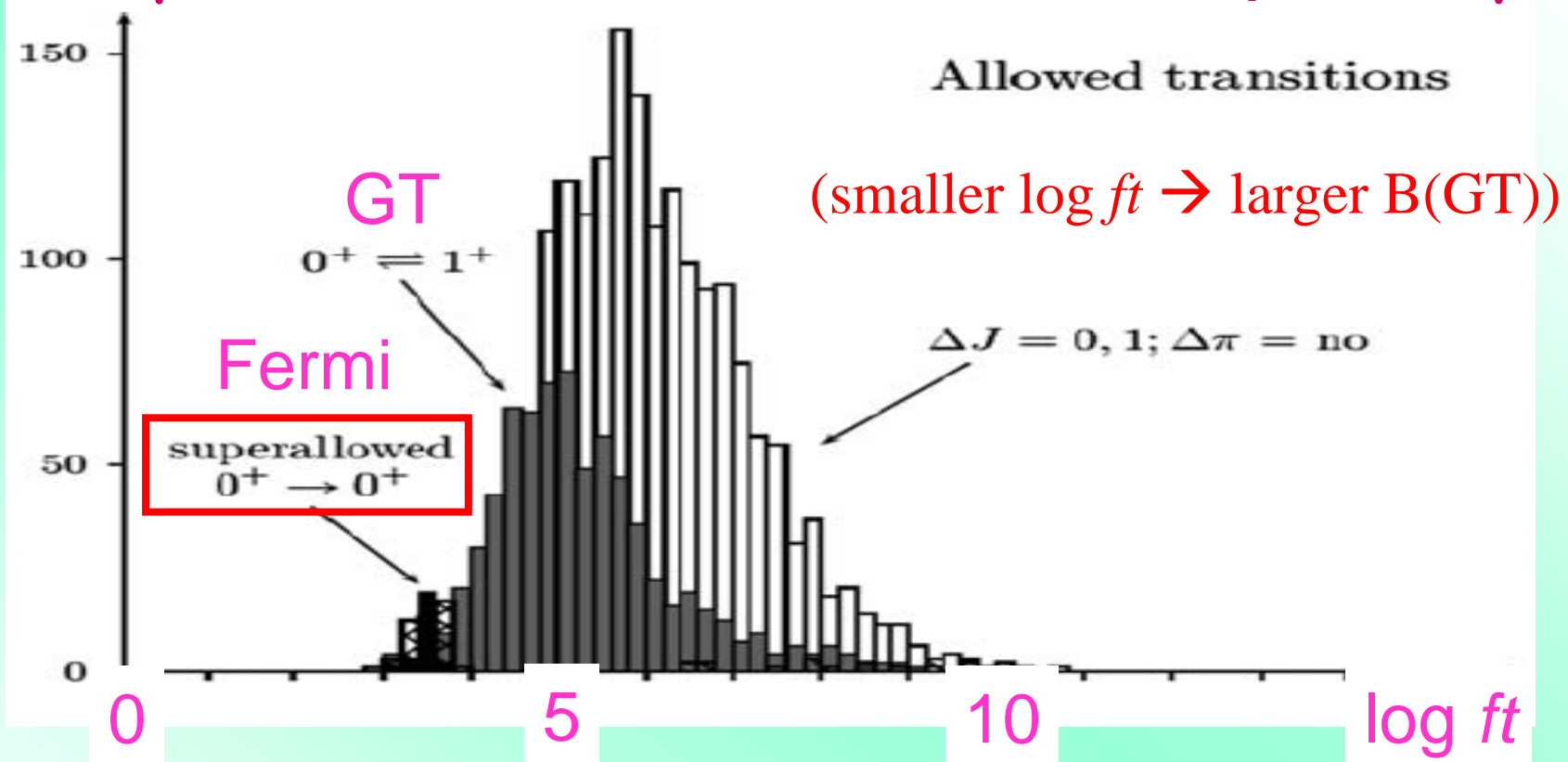
Y. Fujita, et al., PRL 112, 112502 (2014).

$^{42}\text{Ca}(^3\text{He},t)^{42}\text{Sc}$ in 2 scales



Y. Fujita, et al., PRL 112, 112502 (2014).

Super-allowed GT transitions in β decay



${}^6\text{He}, 0^+ \rightarrow {}^6\text{Li}, 1^+$	$\log ft = 2.9$
${}^{18}\text{Ne}, 0^+ \rightarrow {}^{18}\text{F}, 1^+$	$\log ft = 3.1$
${}^{42}\text{Ti}, 0^+ \rightarrow {}^{42}\text{Sc}, 1^+$	$\log ft = 3.2$

Super-allowed
GT transitions

Super-Multiplet State

*proposed by Wigner (1937)

In the limit of null $L \cdot S$ force, SU(4) symmetry exists.

We expect:

- a) GT excitation strength is concentrated in a low-energy GT state.
- b) excitation energies of both the IAS and the GT state are identical.

→ *Super-Multiplet State*

In ^{54}Co , we see a broken SU(4) symmetry.

In ^{42}Sc , we see a good SU(4) symmetry.

→ attractive IS residual int. restores the symmetry !

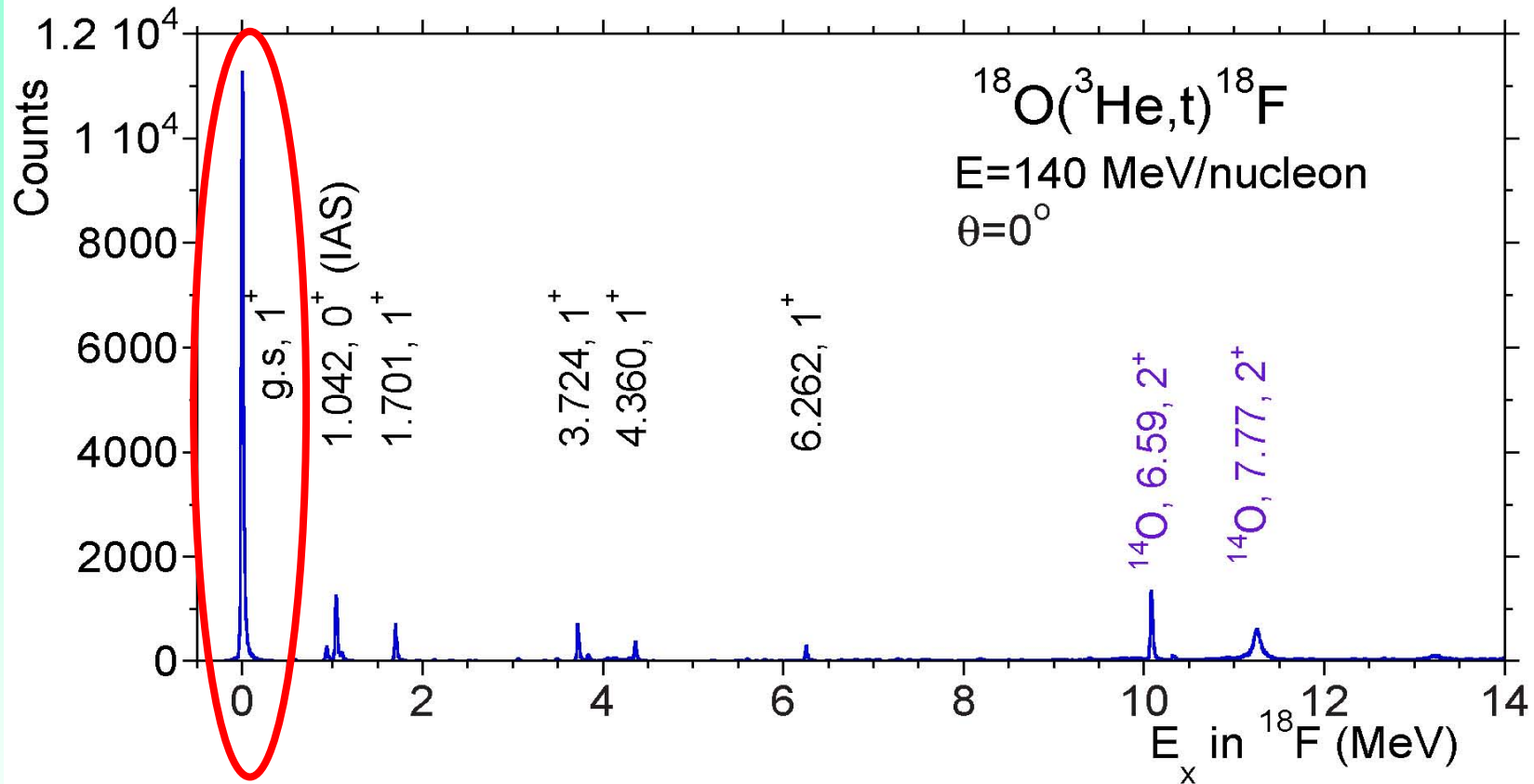
→ 0.611 MeV state in ^{42}Sc has a character close to

Super-Multiplet State !

We call this state the

Low-energy Super GT state !

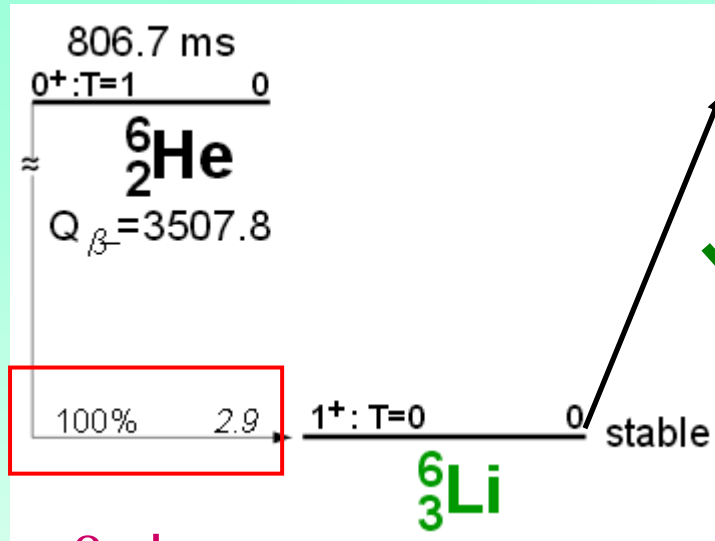
$^{18}\text{O}(^3\text{He},t)^{18}\text{F}$ at 0°



Low-energy collective GT excitation: $B(\text{GT})=3.1$

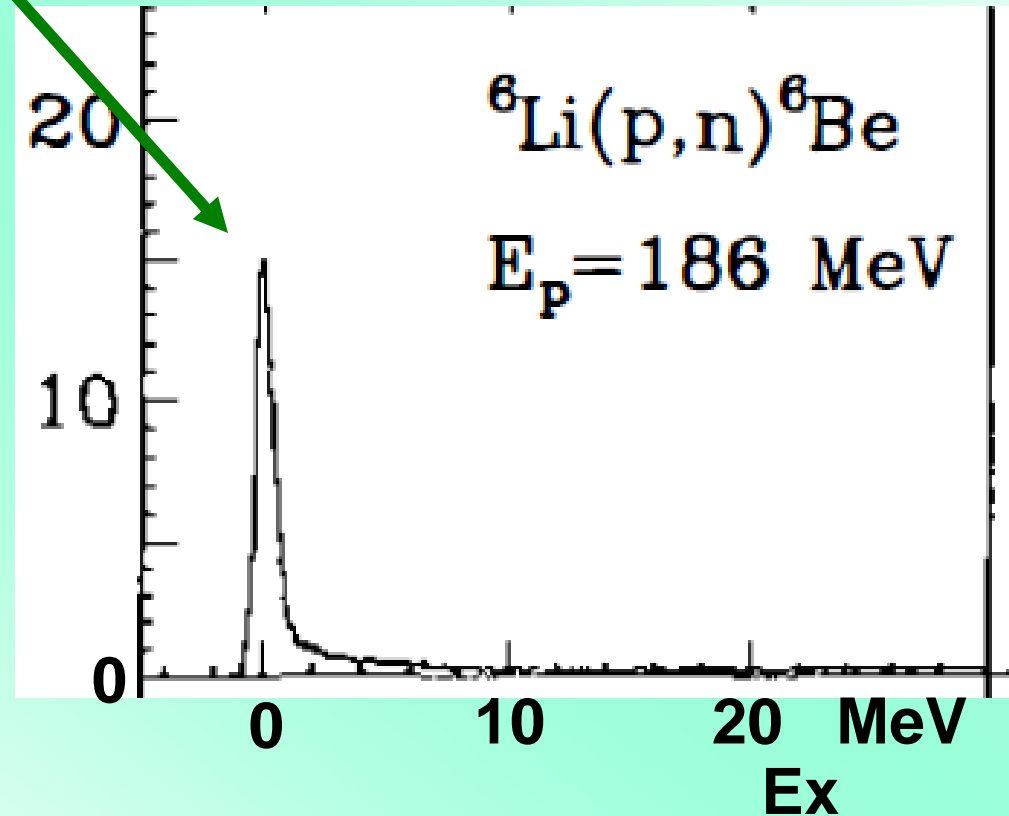
Low Energy Super GT state

${}^6\text{He}$ β -decay & ${}^6\text{Li}(pn){}^6\text{Be}$

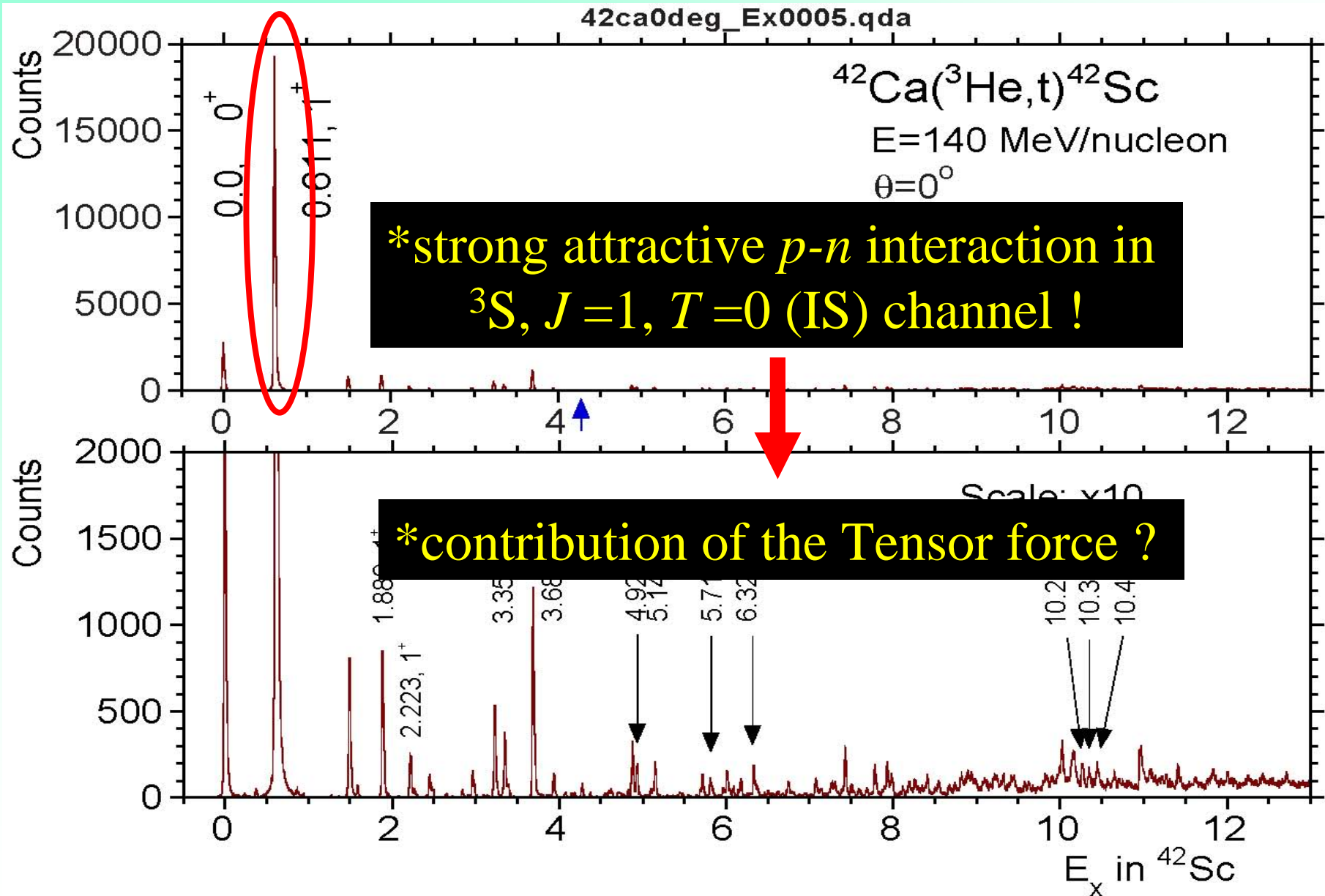


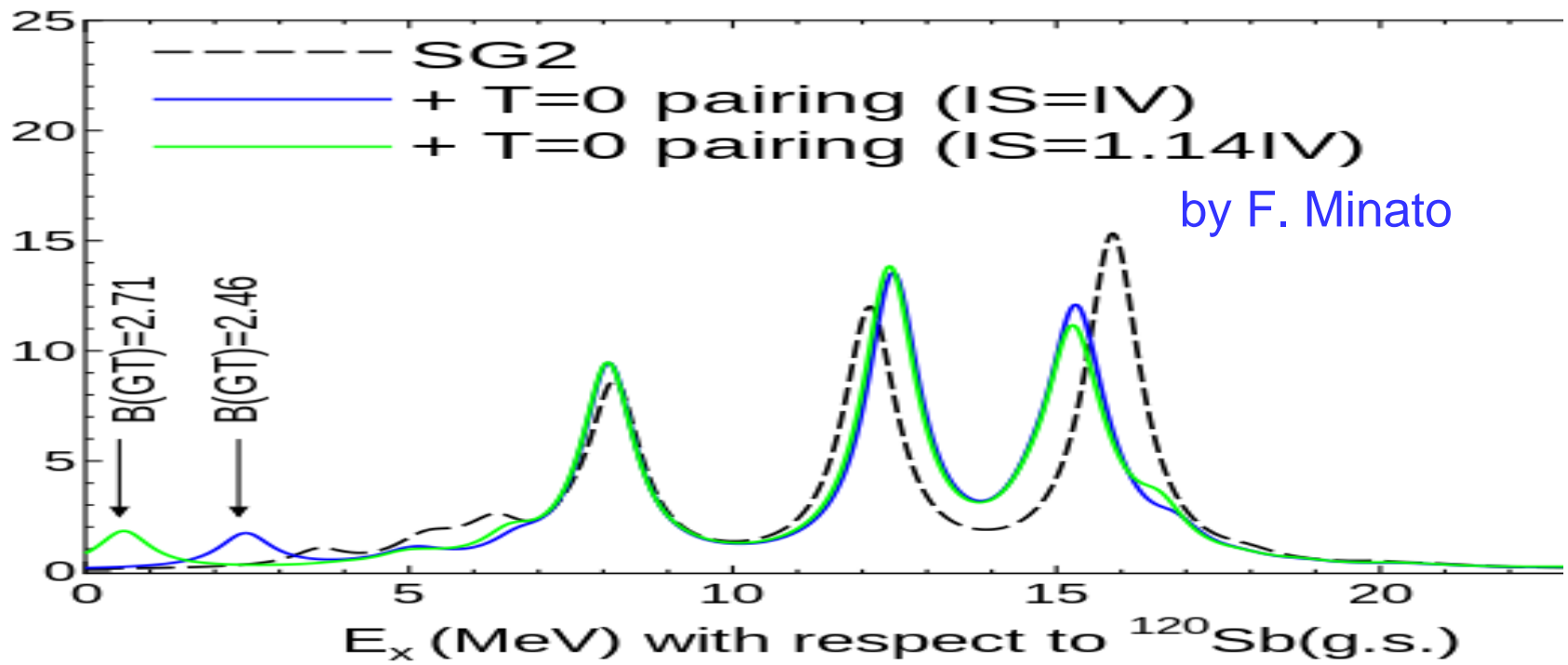
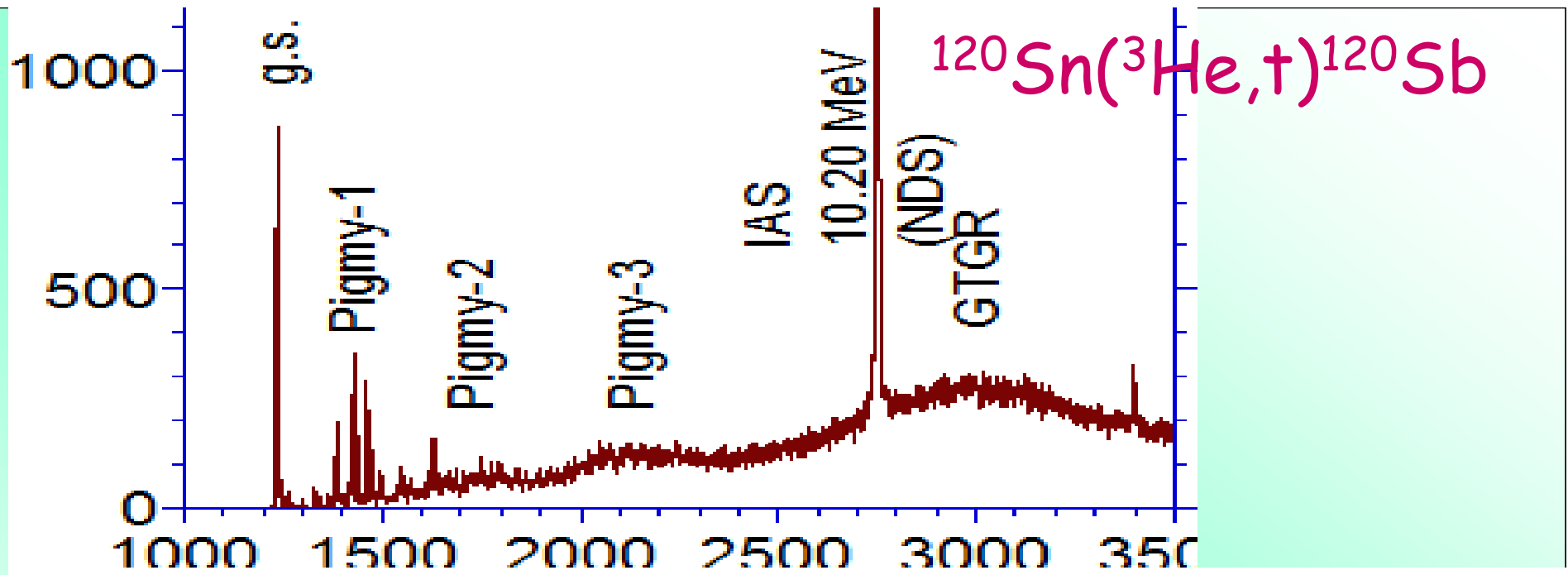
β -decay
 $\log ft = 2.9$
 very small !

${}^6\text{Be} \rightarrow 2p + \alpha$
 $\Gamma = 92 \text{ keV}$

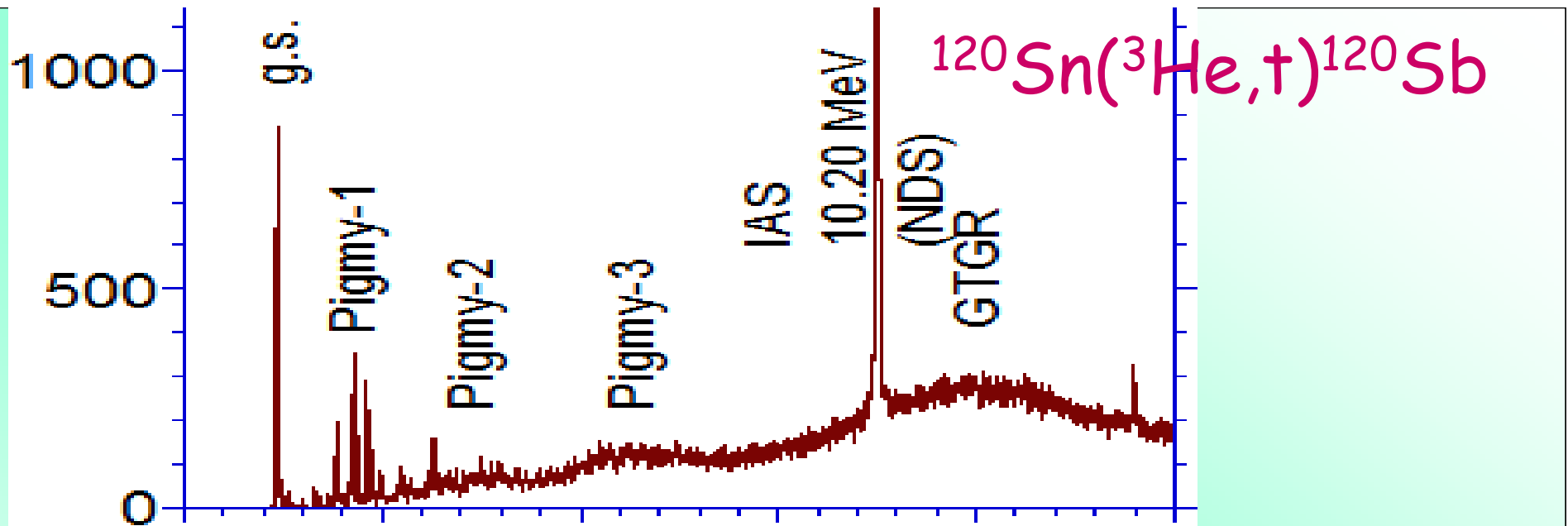


$^{42}\text{Ca}(^3\text{He},t)^{42}\text{Sc}$ in 2 scales

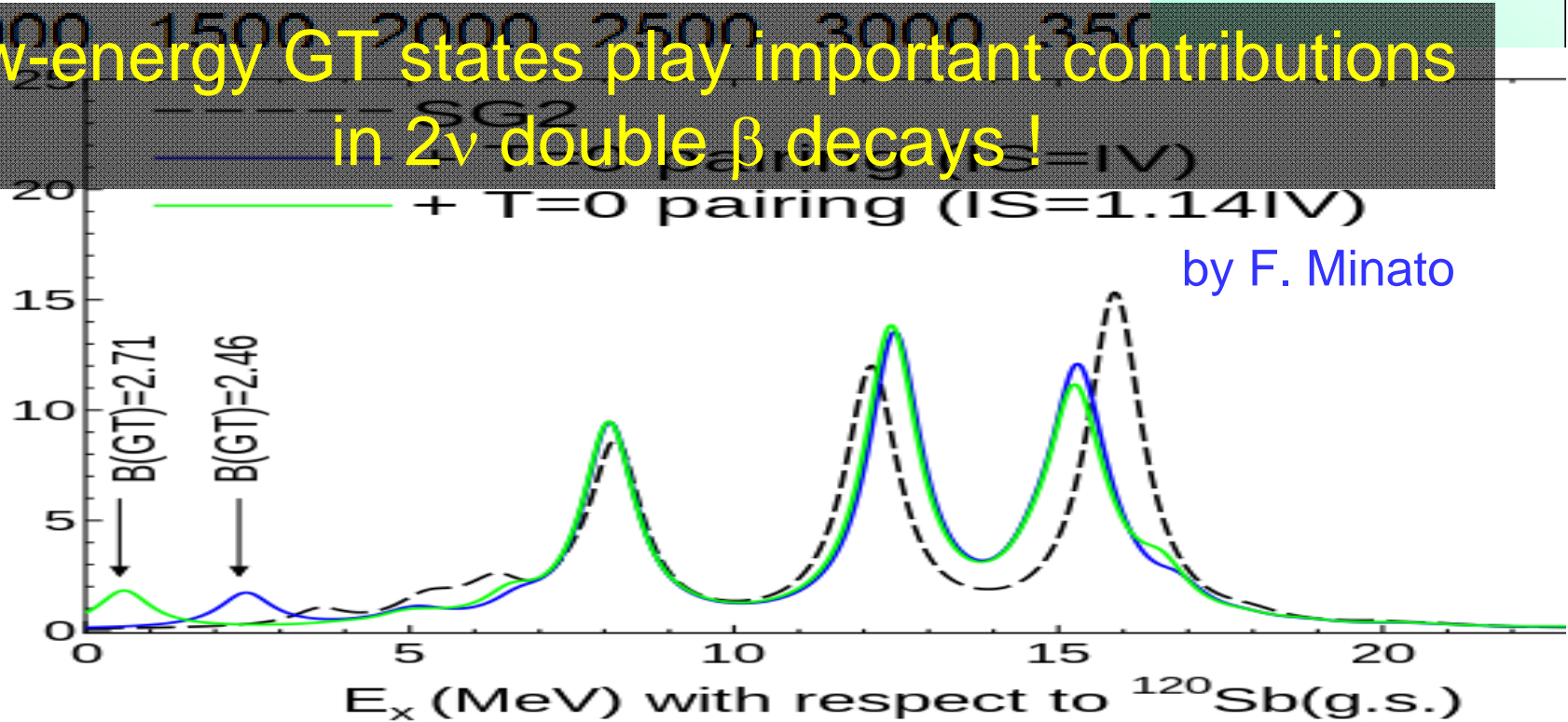




by F. Minato



Low-energy GT states play important contributions in 2ν double β decays!



List of double β -decay nuclei

^{48}Ca	CANDLES
^{64}Zn	COBRA
^{76}Ge	GERDA
^{82}Se	NEMO
^{96}Zr	NEMO
^{100}Mo	MOON/NEMO
^{116}Cd	COBRA
$^{128/130}\text{Te}$	CUORE
^{136}Xe	EXO, KamLAND-ZEN
^{150}Nd	SNO+

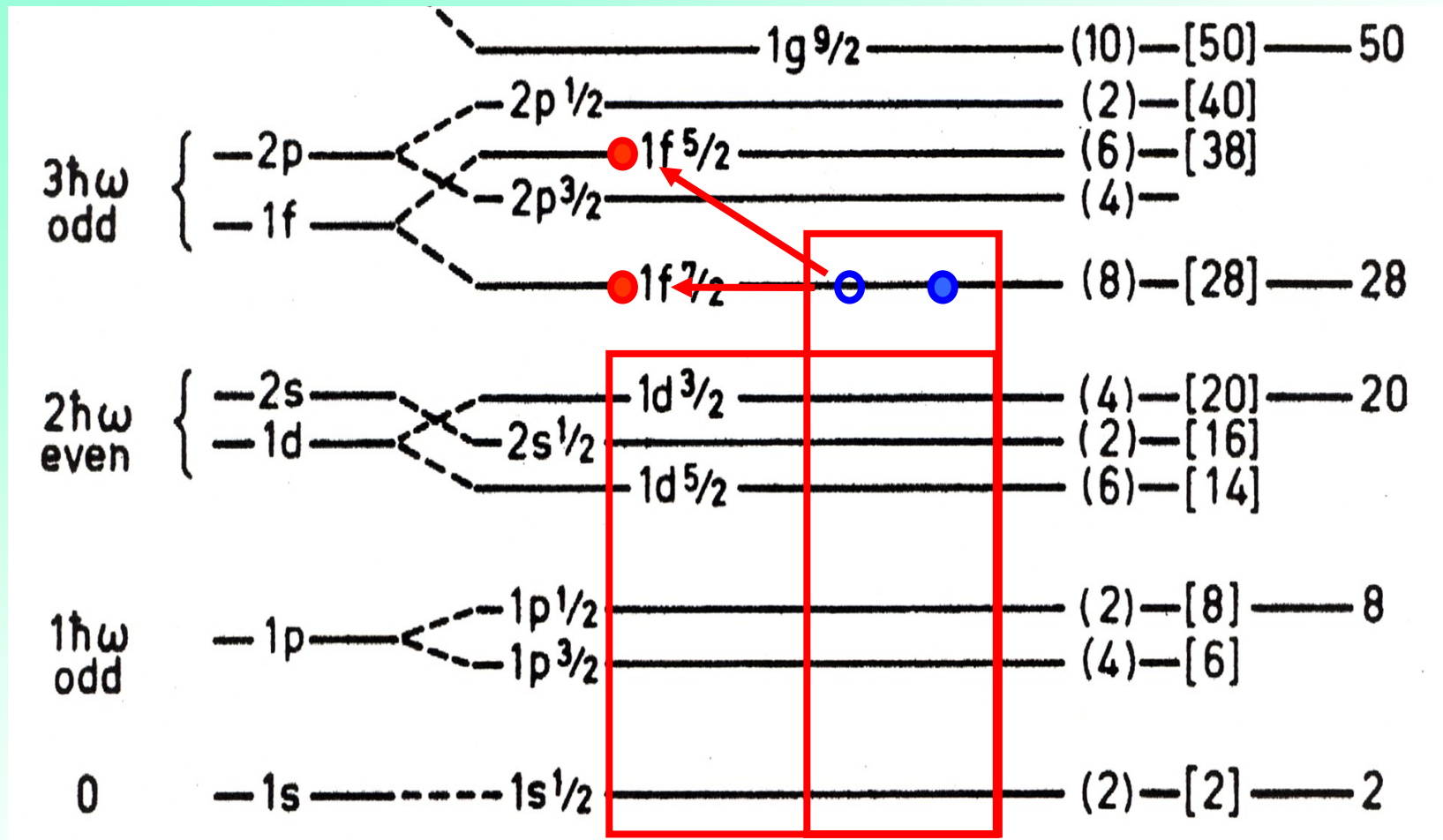
All of these nuclei have $N > Z$!

$B(GT)^-$ & $B(GT)^+$ strengths from Ca isotopes

Ikeda Sum Rule

$$\sum B(GT)\beta^- - \sum B(GT)\beta^+ = 3(N-Z)$$

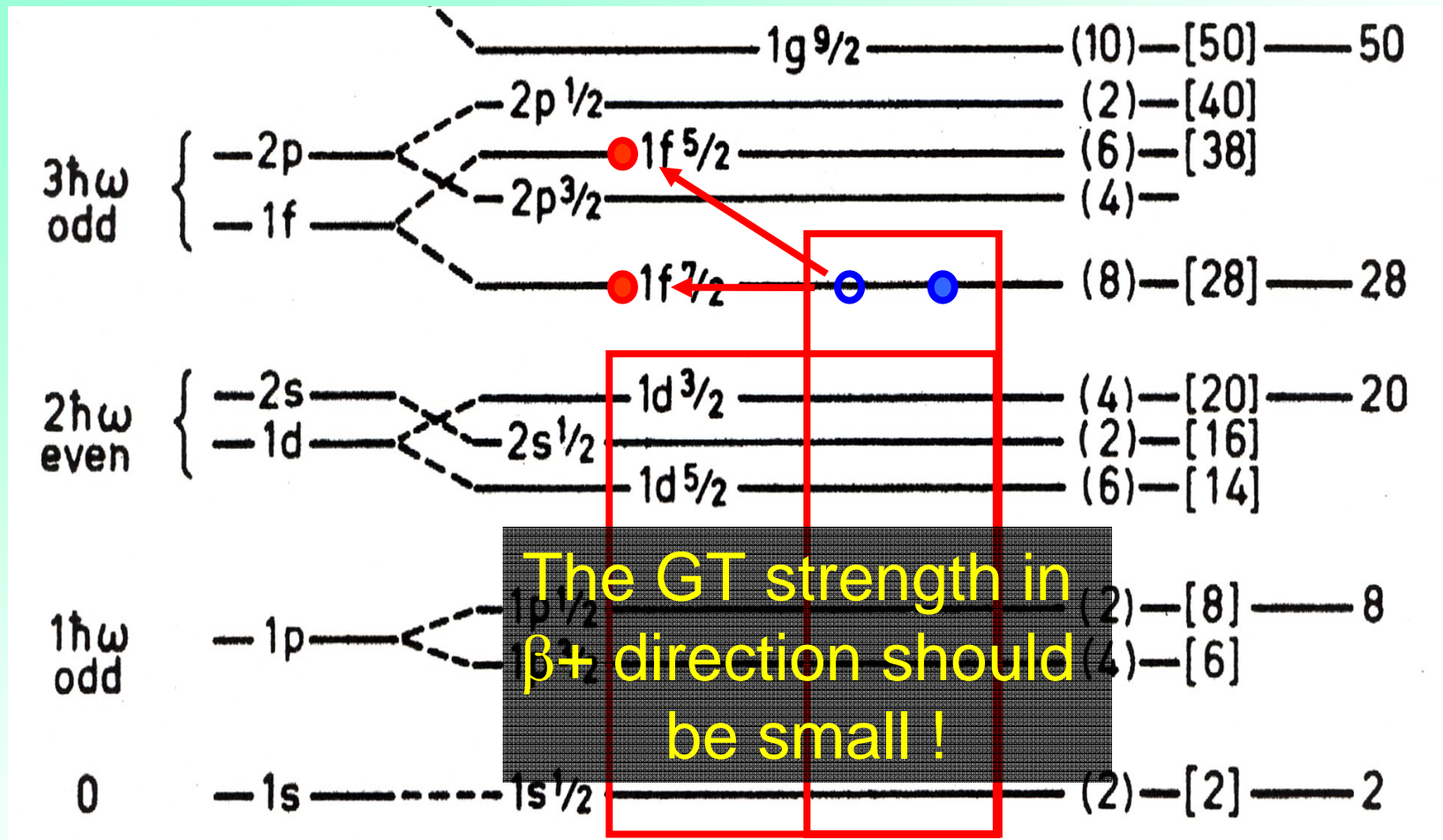
B(GT)⁻ & B(GT)⁺ strengths from Ca isotopes



neutron: $f_{7/2} \rightarrow$ proton $f_{7/2}$

neutron: $f_{7/2} \rightarrow$ proton $f_{5/2}$

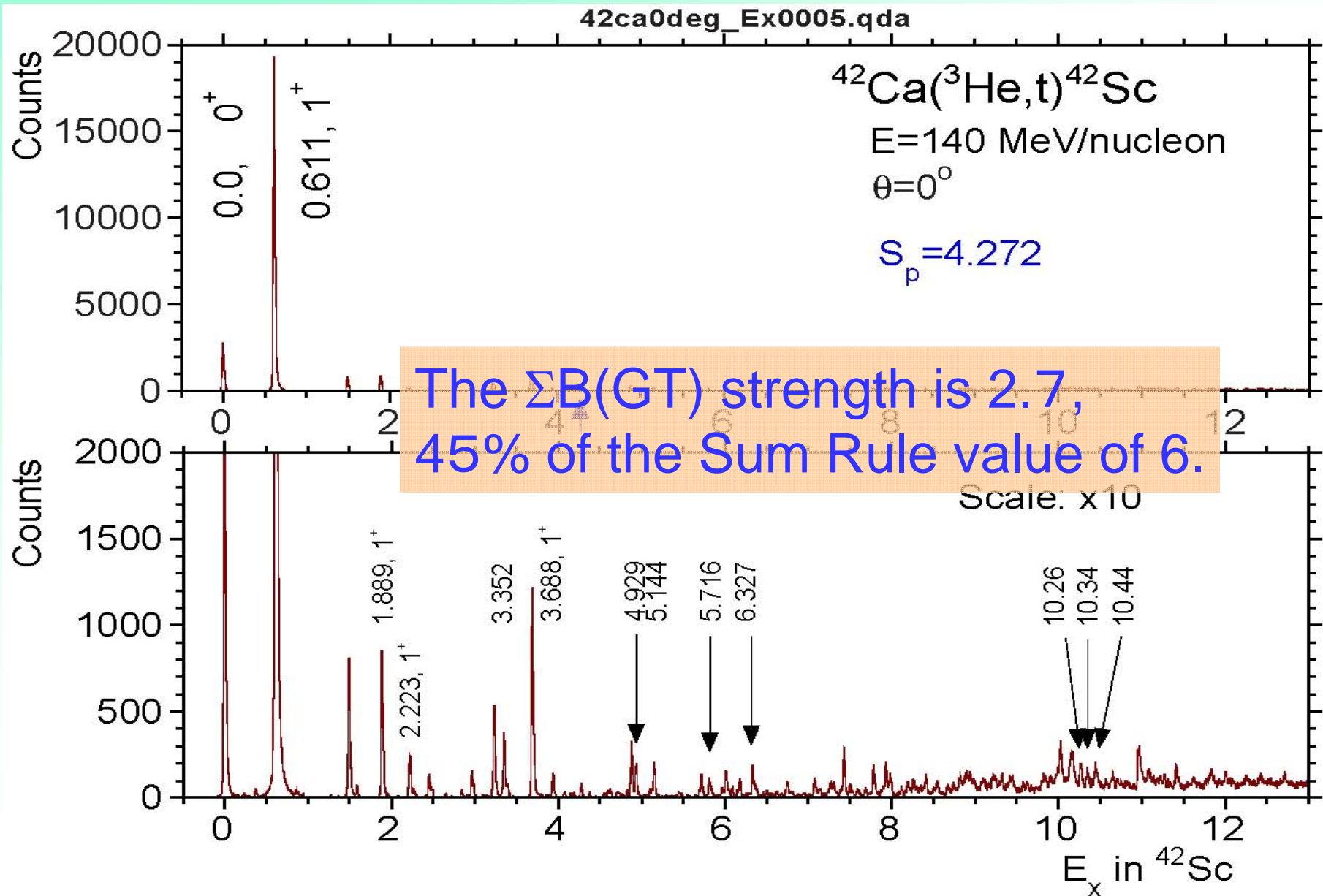
B(GT)⁻ & B(GT)⁺ strengths from Ca isotopes



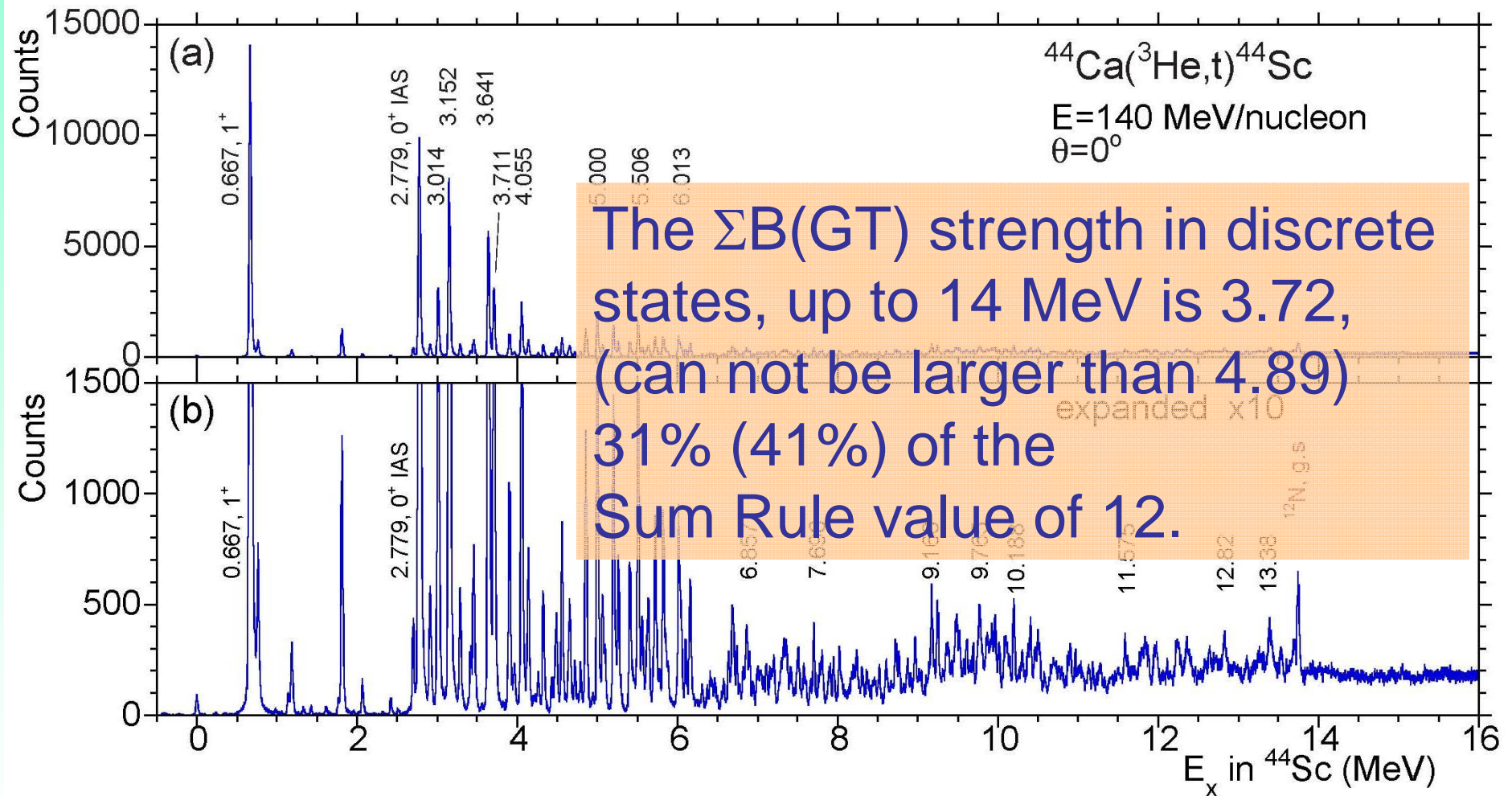
neutron: $f_{7/2} \rightarrow$ proton $f_{7/2}$

neutron: $f_{7/2} \rightarrow$ proton $f_{5/2}$

$^{42}\text{Ca}(^3\text{He},t)^{42}\text{Sc}$ in 2 scales



$^{44}\text{Ca}(^3\text{He},t)^{44}\text{Sc}$ in 2 scales



Y. Fujita et al., PRC in press

$\Sigma B(GT)$
in
 $^{48}\text{Ca}(p,n)^{48}\text{Sc}$

In $E_x < 30$ MeV,
 $\Sigma B(GT+IVSD = \Delta L=0)$
is 15.3,
which is 64(9)% of the
Ikeda Sum Rule value
of $3(N-Z) = 24$

K. Yako et al.,
PRL103 (2009)

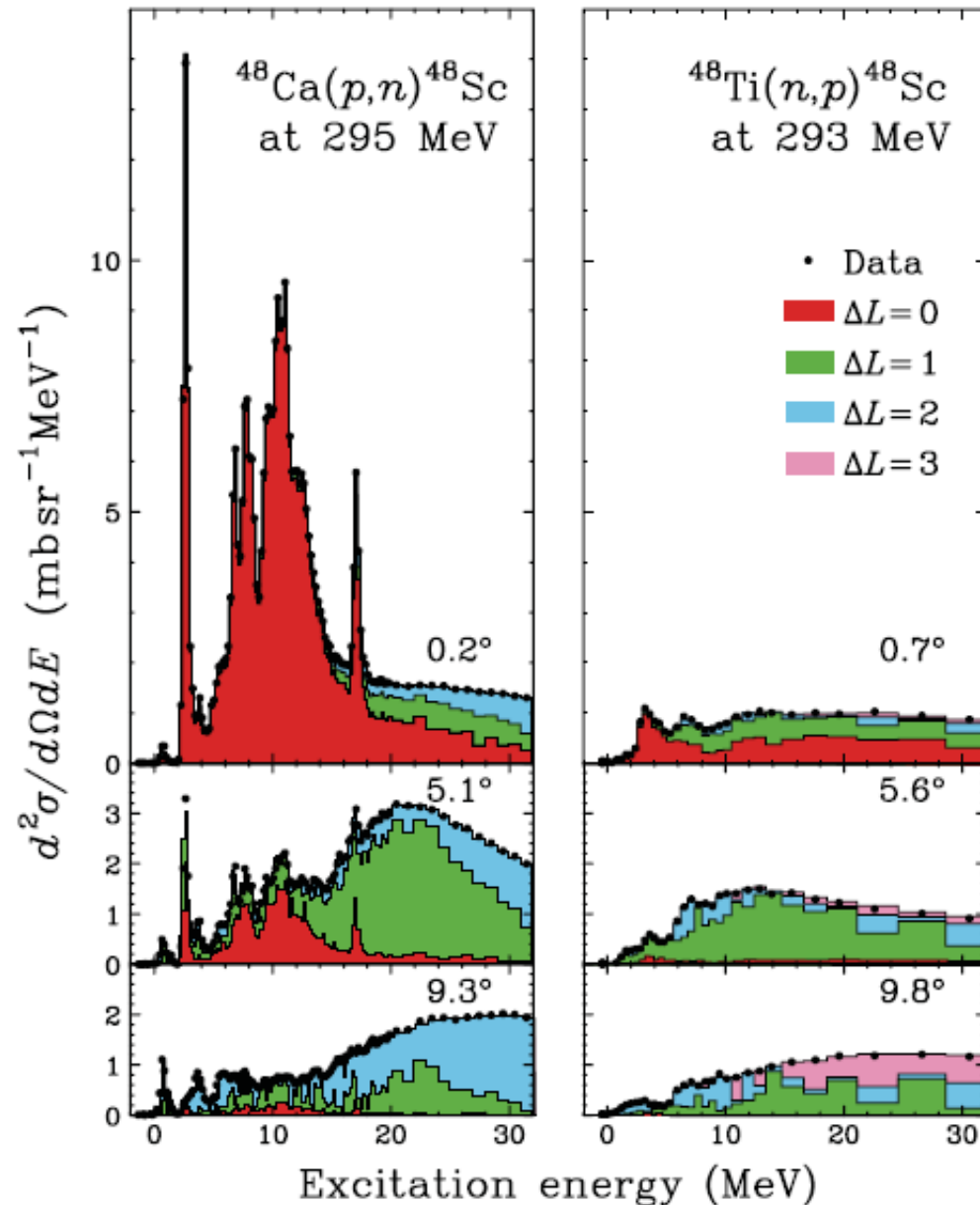


FIG. 1 (color). Double-differential cross sections for the $^{48}\text{Ca}(p,n)^{48}\text{Sc}$ (left-hand panel) and $^{48}\text{Ti}(n,p)^{48}\text{Sc}$ (right-hand panel) reactions. The histograms show the MD analysis results.

Summary

GT ($\sigma\tau$) operator : a simple operator !

* GT transitions: sensitive to the structure of $|i\rangle$ and $|f\rangle$

High resolution of the ($^3\text{He},t$) reaction

* Fine structures of GT transitions

(Precise comparison with mirror β -decay results)

→ Low-energy Super GT state (LESGT state)

→ Sum Rule values in Ca isotopes?

**We got a key to study the IS pn -interaction !
(May be connected to Tensor ?)**

GT-study Collaborations

Bordeaux (France) : β decay

GANIL (France) : β decay

Gent (Belgium) : (^3He , t), (d, ^2He), (γ , γ'), theory

GSI, Darmstadt (Germany) : β decay, theory

ISOLDE, CERN (Switzerland) : β decay

iThemba LABS. (South Africa) : (p, p'), (^3He , t)

Istanbul (Turkey): (^3He , t), β decay

Jyvaskyla (Finland) : β decay

Koeln (Germany) : γ decay, (^3He , t), theory

KVI, Groningen (The Netherlands) : (d, ^2He)

Leuven (Belgium) : β decay

LTH, Lund (Sweden) : theory

Osaka University (Japan) : (p, p'), (^3He , t), theory

Surrey (GB) : β decay

TU Darmstadt (Germany) : (e, e'), (^3He , t)

Valencia (Spain) : β decay

Michigan State University (USA) : theory, (t, ^3He)

Muenster (Germany) : (d, ^2He), (^3He ,t)

Univ. Tokyo and CNS (Japan) : theory, β decay

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Review

Spin–isospin excitations probed by strong, weak and electro-magnetic interactions

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