

International School of Nuclear Physics - 38th course
Nuclear matter under extreme conditions - relativistic heavy-ion collisions
September 2016



Quarkonium results in pA & AA: from RHIC to LHC

Roberta Araldi
INFN Torino

International School of Nuclear Physics - 38th course

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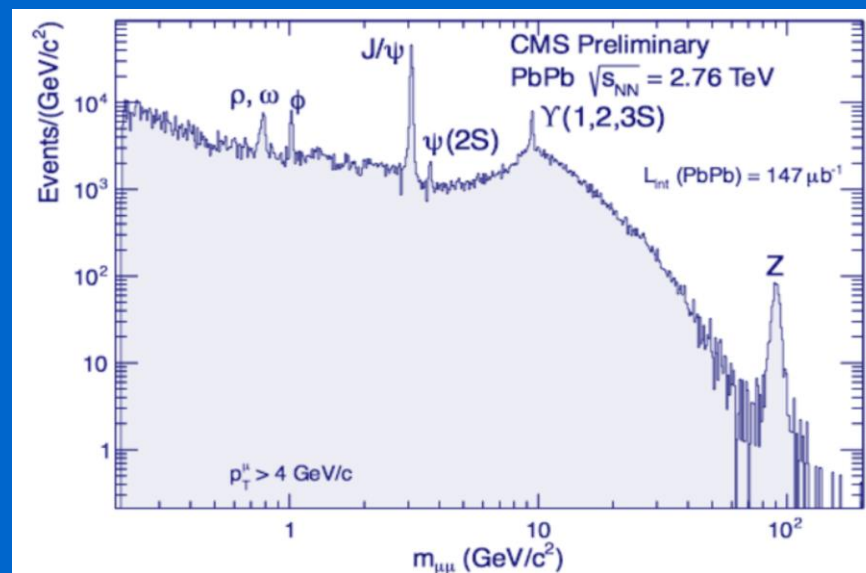


Outlook:

Selection on results on

- Charmonium: J/ψ and $\psi(2S)$
- Bottomonium: $\Upsilon(1S)$, $\Upsilon(2S)$, $\Upsilon(3S)$

in p-A, d-A and A-A collisions at RHIC and LHC energies



AA: hot matter effects

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→ the original idea

quarkonium production suppressed via color screening in the QGP

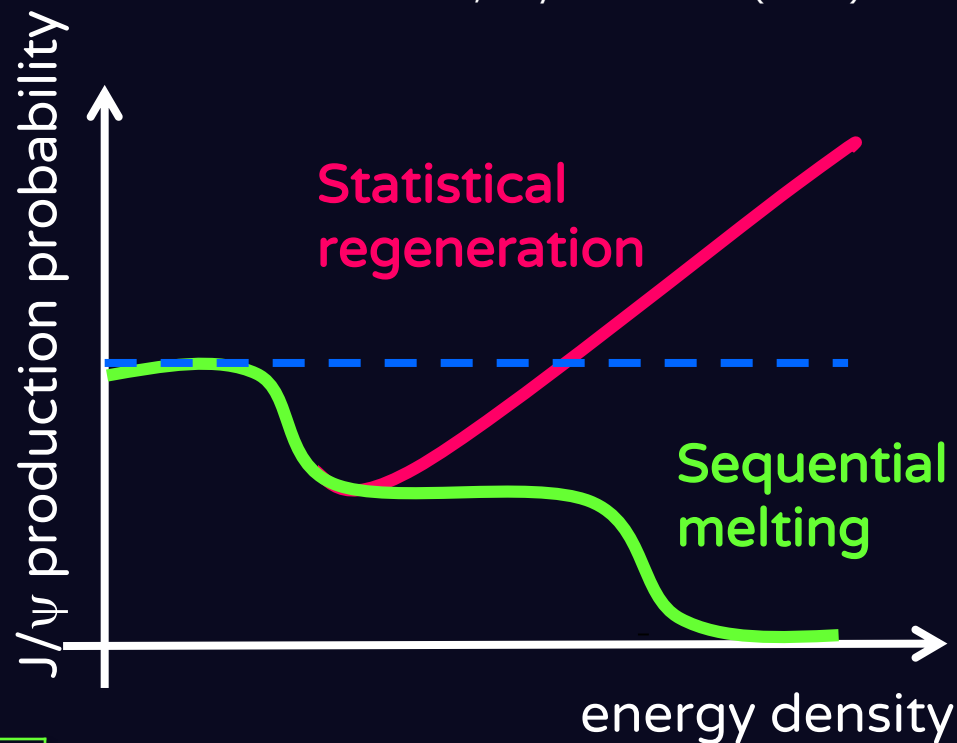
→ sequential melting

differences in quarkonium binding energies lead to a sequential melting with increasing temperature

→ (re)combination

enhanced quarkonium production through (re)combination during QGP phase or at hadronization

T.Matsui and H.Satz, Phys.Lett.B178 (1986) 416



Central AA collisions	SPS 20 GeV	RHIC 200 GeV	LHC 2.76TeV	LHC 5.02TeV
$N_{c\bar{c}}/\text{event}$	~0.2	~10	~85	~115

P. Braun-Muzinger, J. Stachel, PLB 490(2000) 196
R. Thews et al, Phys.Rev.C63:054905(2001)

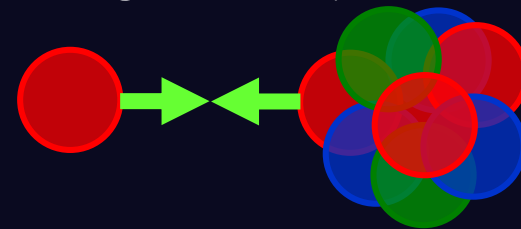
pA: CNM effects

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→ **Cold nuclear matter effects:** might affect quarkonium production on top of hot matter mechanisms

- nuclear parton shadowing/
color glass condensate
- energy loss
- $c\bar{c}$ in medium break-up

investigated in p-A collisions



→ the assessment of the size of these effects is fundamental to interpret quarkonium A-A results

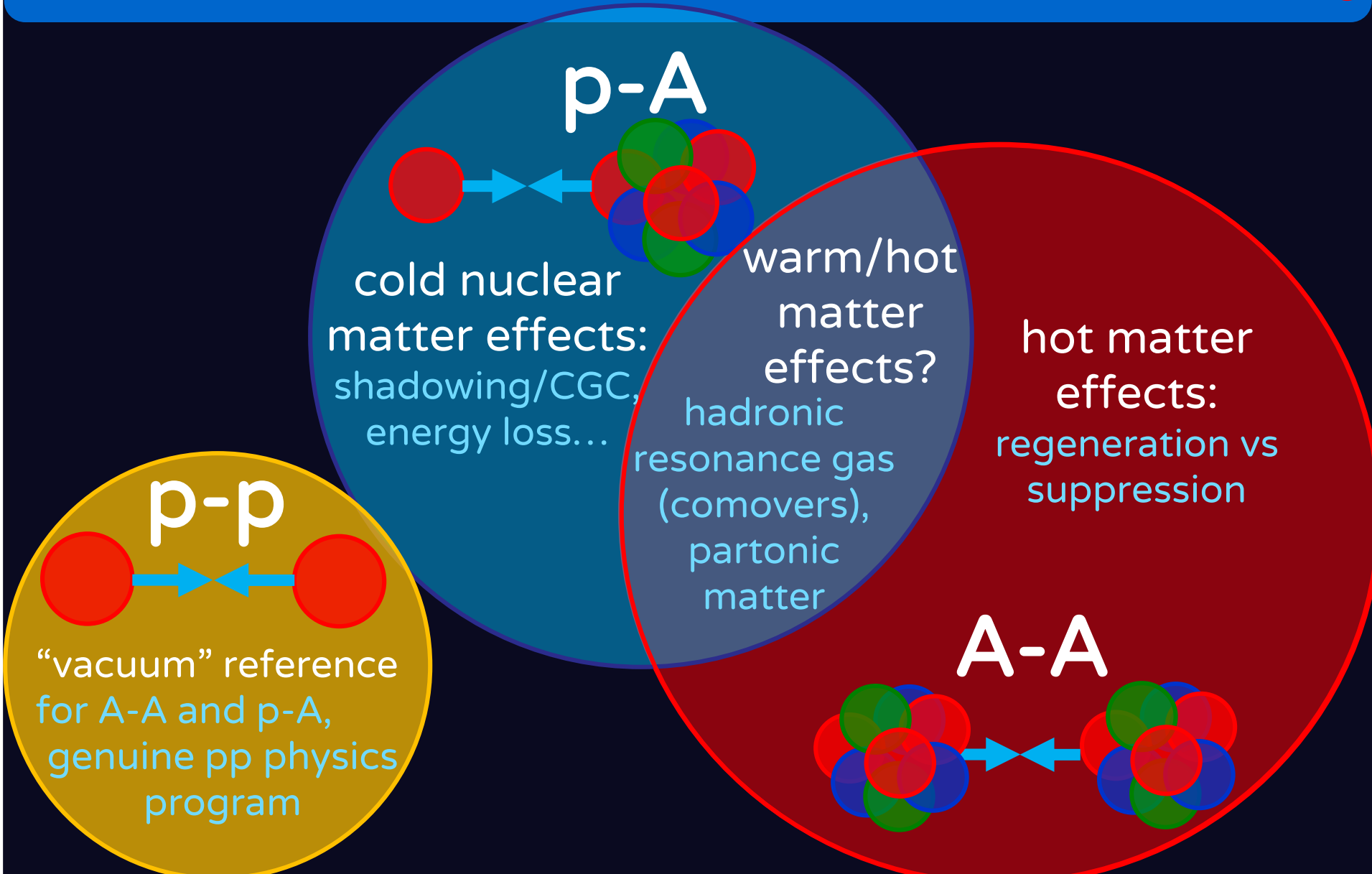
→ **Nuclear modification factor** Medium effects are quantified comparing the AA quarkonium yield with the pp one, scaled by a geometrical factor (from Glauber model)

$$R_{AA}^{J/\psi} = \frac{Y_{AA}^{J/\psi}}{\langle T_{AA} \rangle \sigma_{pp}^{J/\psi}}$$

- $R_{AA} = 1 \rightarrow$ no medium effects
- $R_{AA} \neq 1 \rightarrow$ hot/cold matter effects

Quarkonium in HI

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Quarkonium at RHIC & LHC

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Facility	Experiment	System	$\sqrt{s_{NN}}$ (GeV)	Data taking
RHIC	PHENIX STAR	Au-Au, Cu-Cu, Cu-Au, U-U	200, 193, 62, 39	2000-2016
		p-A, d-Au	200	
		pp	200-500	
LHC	ALICE ATLAS CMS LHCb	Pb-Pb	2760 5020	2010-2012 2015
		p-Pb	5020 (8000)	2013 (2016)
		pp	2760, 7000, 8000, 13000	2010-2016

Quarkonium at RHIC & LHC

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Facility	Experiment	System	$\sqrt{s_{NN}}$ (GeV)	Data taking
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		p-A, d-Au	200	
		pp	200-500	
LHC	ALICE ATLAS CMS LHCb	Pb-Pb	2760 5020	2010-2012 2015
		Quarkonium production investigated via collisions:		2013 (2016)
		<ul style="list-style-type: none"> with different beam species at various energies 		10-2016
				15000

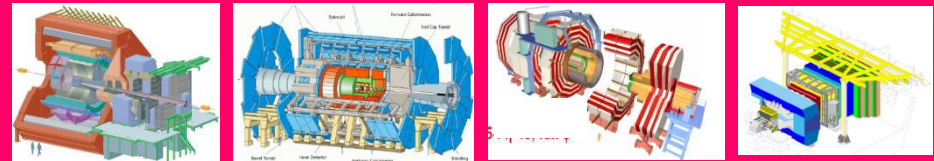
Quarkonium at RHIC & LHC

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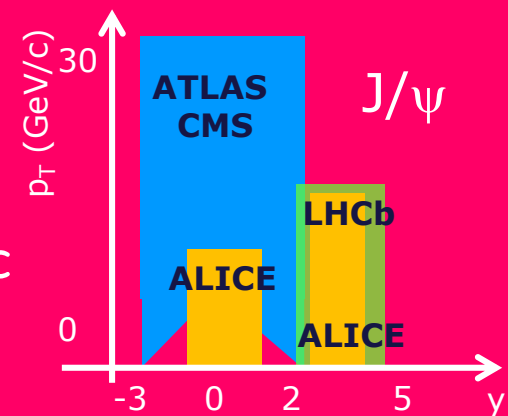
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		p-A, d-Au	200	

LHC	ALICE ATLAS CMS LHCb
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All LHC experiments investigate quarkonium production



complementary results due to different kinematic coverages



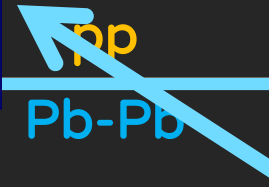
Quarkonium at RHIC & LHC

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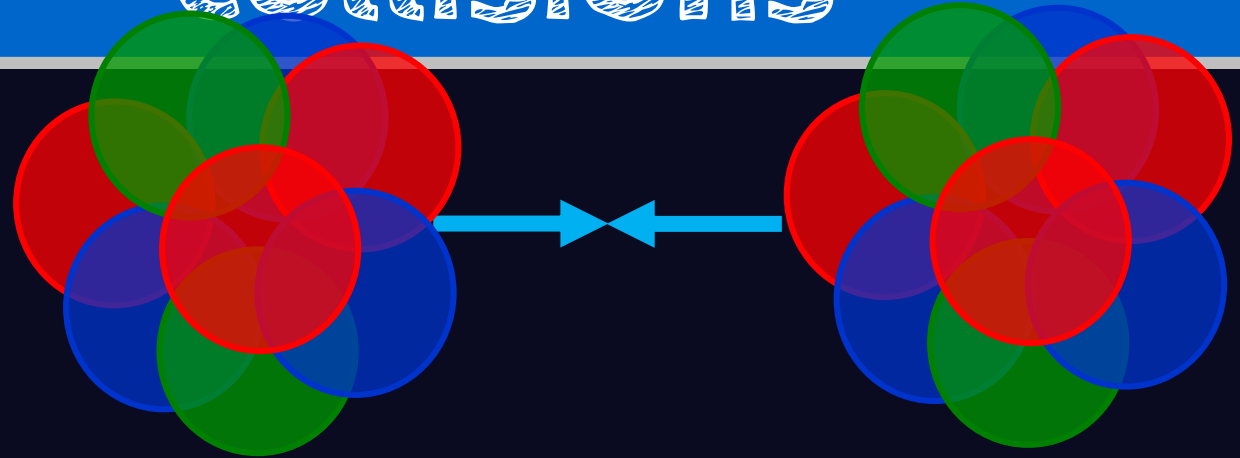
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RHIC	PHENIX	Au-Au, Cu-Cu,	200, 193,	2000-2016
		u-Au, U-U	62, 39	
		p-A, d-Au	200	
		pp	200-500	
LHC	ALICE ATLAS CMS LHCb	Pb-Pb	2760	2010-2012
			5020	2015
		p-Pb	5020	2013
			8000	(2016)
		pp	2760	2010-2016
			5020	
			7000	
		8000		
		13000		

LHC Run-1

LHC Run-2

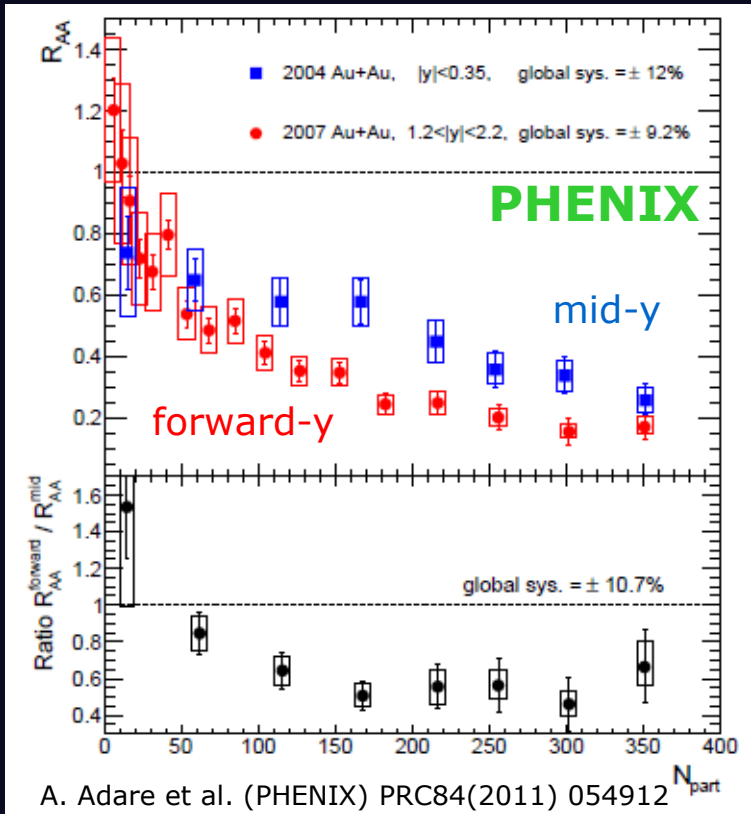


Quarkonium in AA collisions

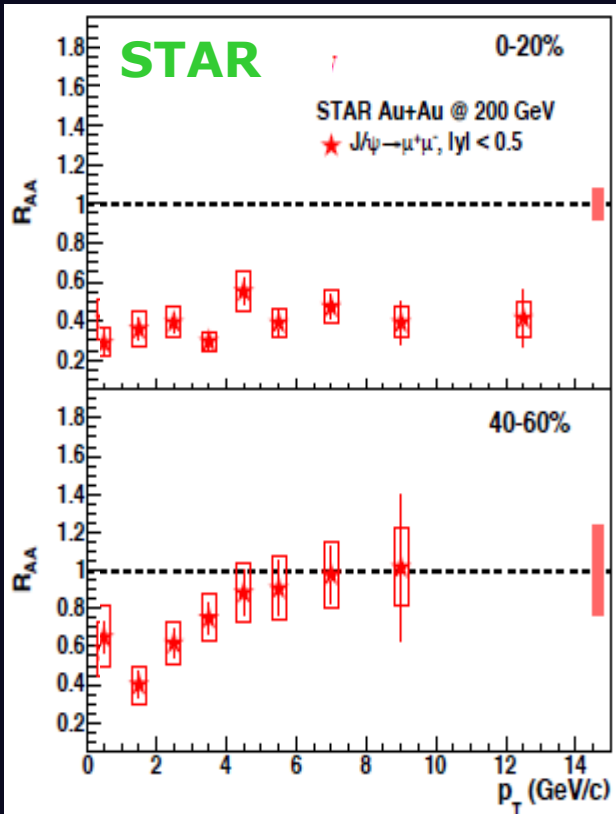
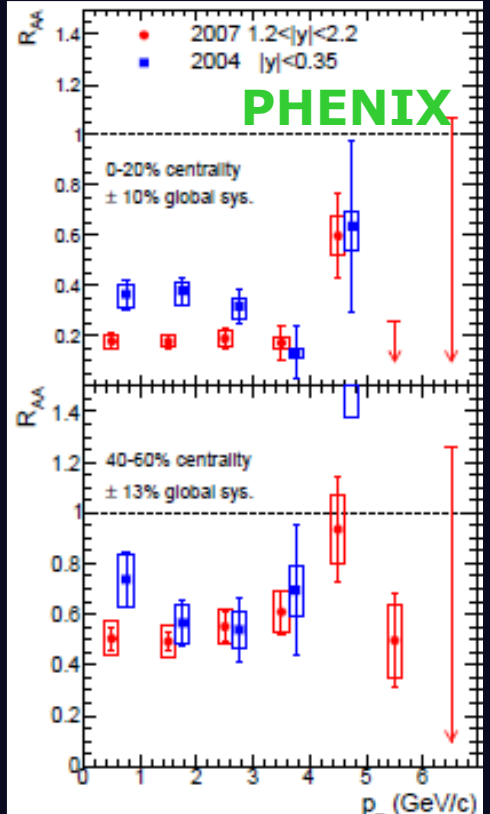


J/ψ suppression at RHIC

➔ Stronger J/ψ suppression at forward-y wrt mid-y in AuAu@200GeV



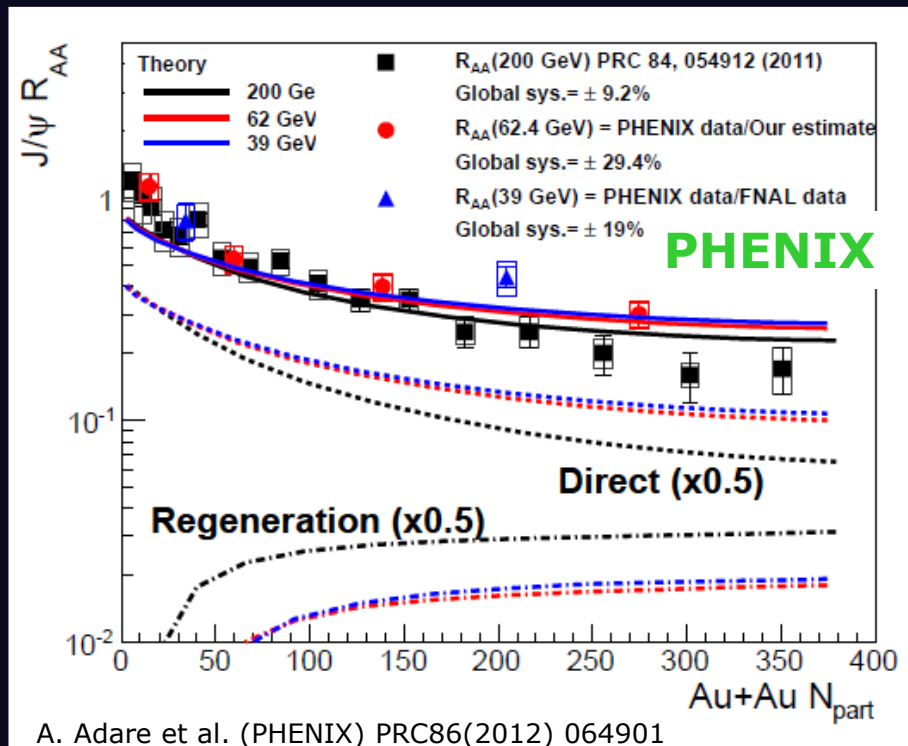
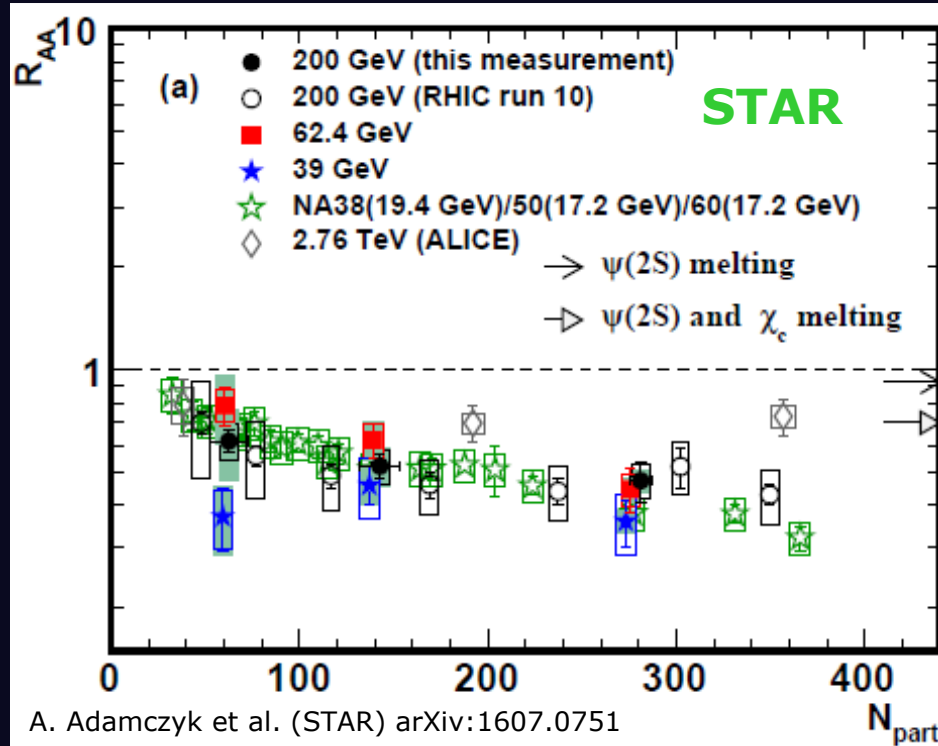
A. Adare et al. (PHENIX) PRC84(2011) 054912



- ➔ Strong centrality and low- p_T suppression
- ➔ Qualitative agreements with suppression + recombination models

Energy scan at RHIC

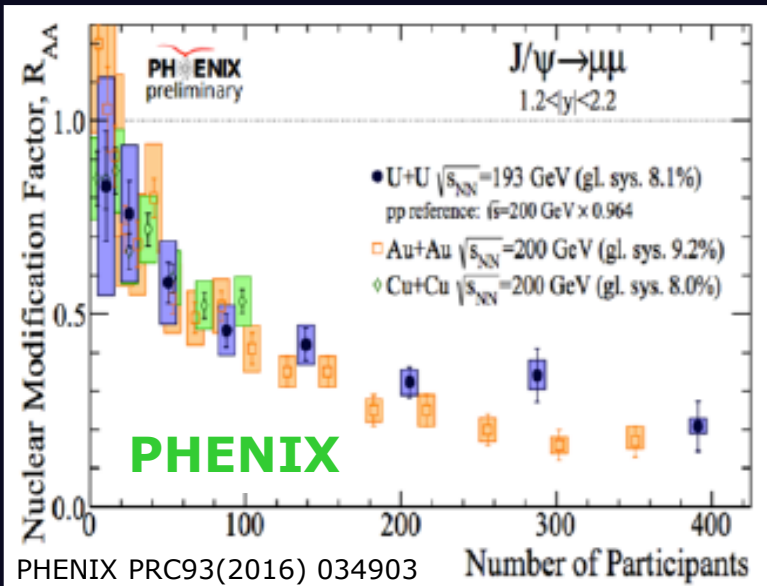
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- ➔ R_{AA} suppression visible at all energies
- ➔ No significant energy dependence, including SPS!
- ➔ Qualitative agreements with suppression + recombination models
→ pp reference at 39 & 62.4 GeV needed for quantitative comparison

System scan at RHIC

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- ➡ Various systems studied:
- rather similar suppression observed
 - hint for a weaker suppression in U-U

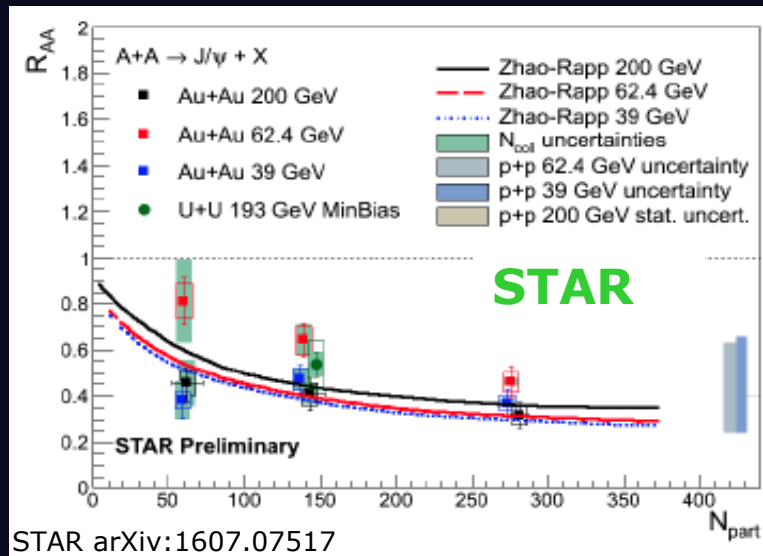
in central U-U collisions:

- 1) stronger color screening suppression

$$\varepsilon_{\text{AuAu}} \sim 80-85\% \varepsilon_{\text{UU}}$$

- 2) J/ψ recombination favoured by 25% larger N_{coll} in U-U

$$N_{J/\psi}^{\text{stat}} \sim N_C^2 \sim N_{\text{coll}}^2$$



- ➡ Dominant recombination in U-U over suppression?

- ➡ Quantitative conclusions depend on U Woods-Saxon description

J/ψ: LHC Run1 results

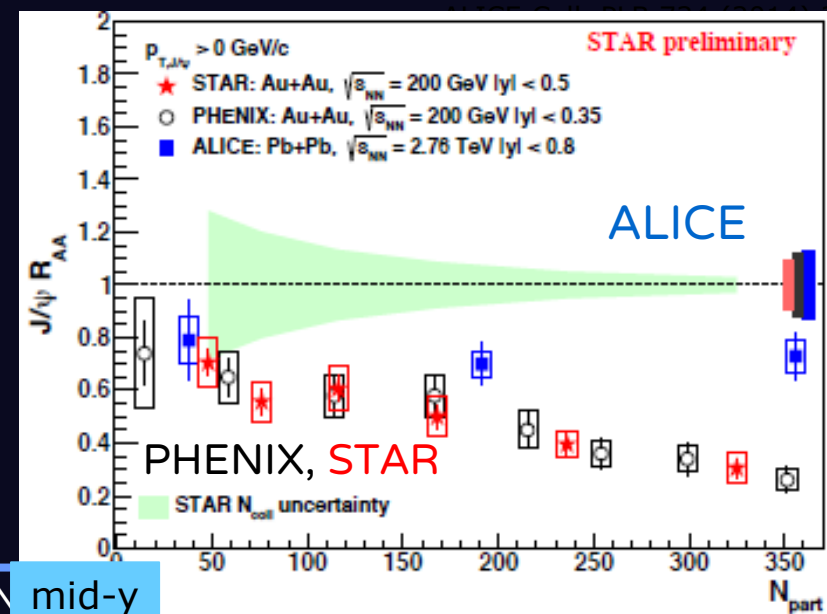
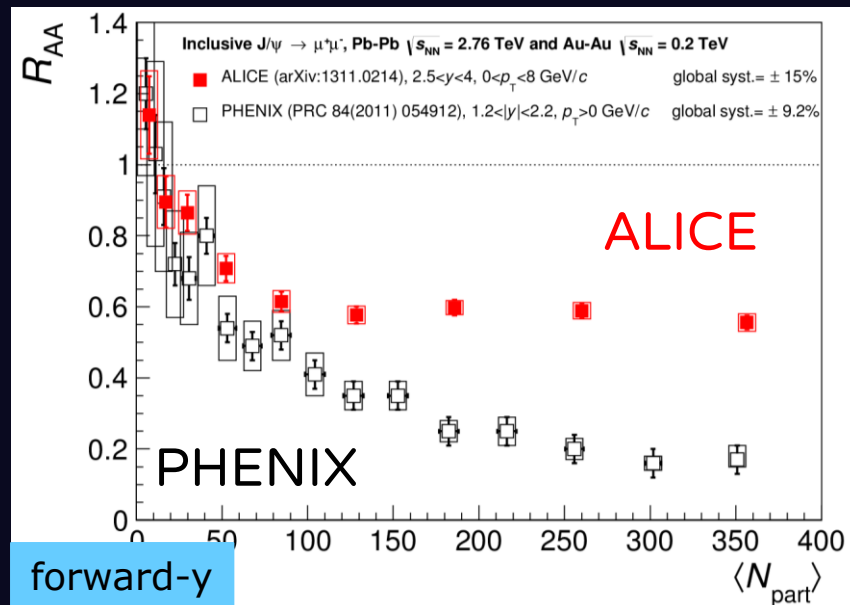
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➔ Evidence of recombination for low p_T J/ψ at LHC

Observation validated by the comparison of LHC results with

1) lower energy experiments

➔ J/ψ suppression vs centrality is stronger in PHENIX/STAR than in ALICE, in spite of the LHC larger energy densities



J/ψ: LHC Run1 results

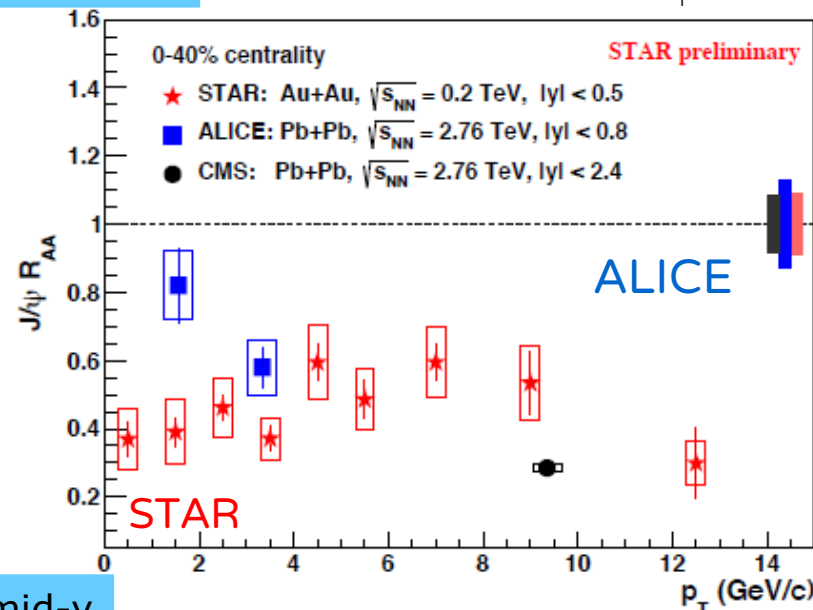
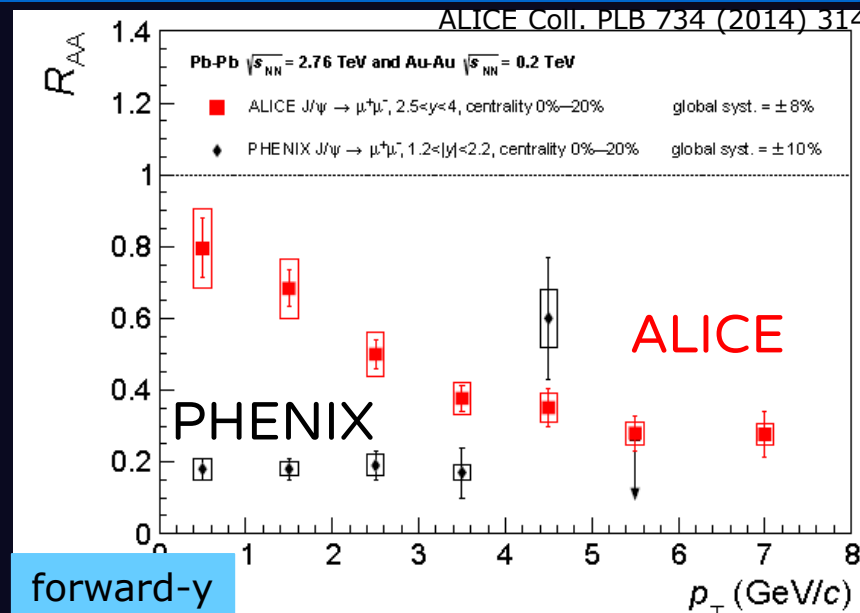
15

➔ Evidence of recombination for low p_T J/ψ at LHC

Observation validated by the comparison of LHC results with

1) lower energy experiments

➔ weaker suppression at low p_T observed by ALICE



J/ψ: LHC Run1 results

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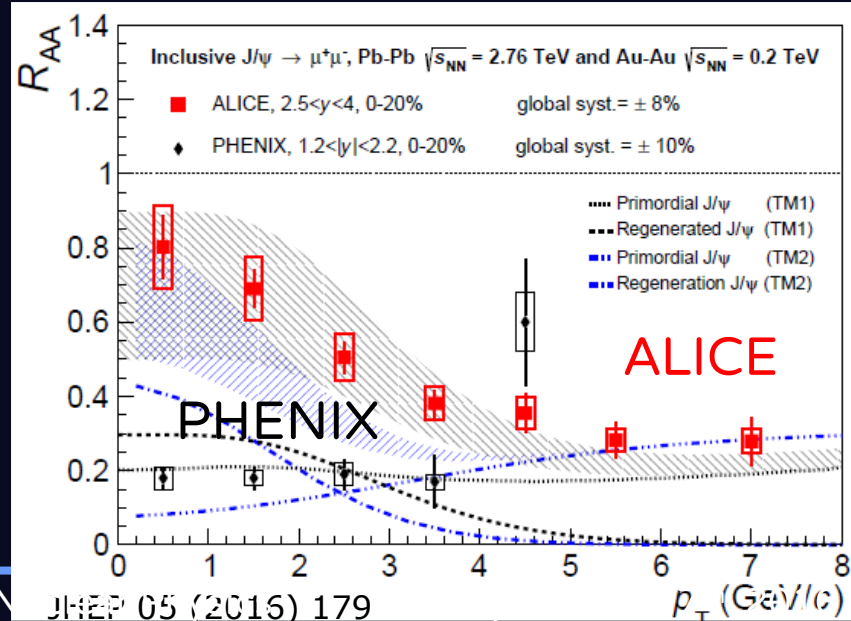
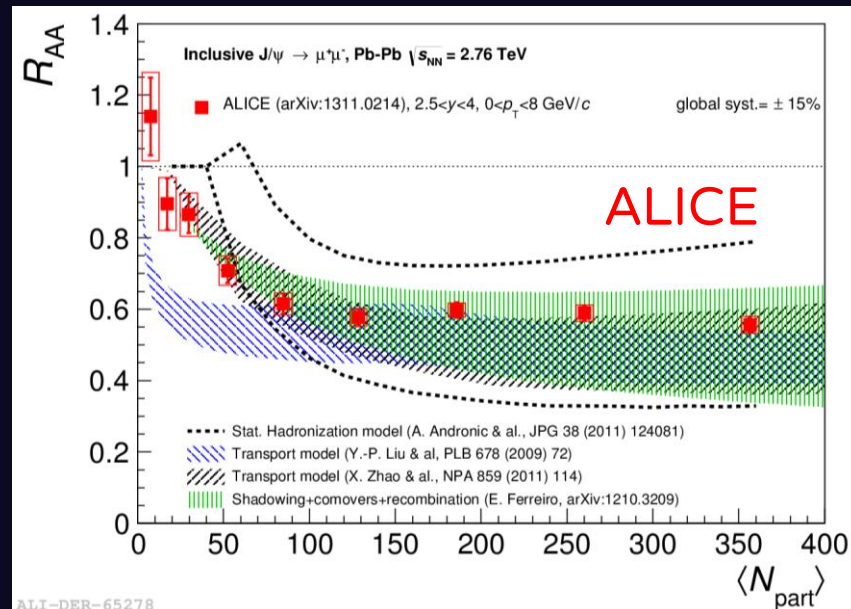
➔ Evidence of recombination for low p_T J/ψ at LHC

Observation validated by the comparison of LHC results with

- 1) lower energy experiments
- 2) theoretical models

➔ models including (re)combination of J/ψ in QGP or in the hadronic phase provide a reasonable description of ALICE results

➔ still rather large theory uncertainties: models will benefit from a precise measurement of σ_{cc} and CNM effects



J/ψ : LHC Run1 results

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➔ Evidence of recombination for low p_T J/ψ at LHC

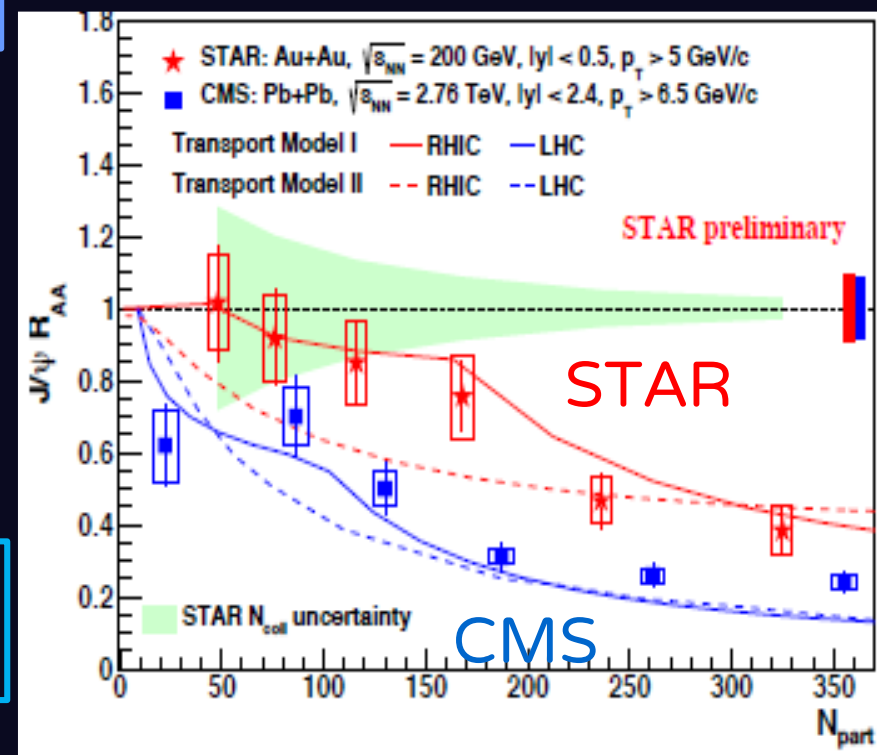
Observation validated by the comparison of LHC results with

- 1) lower energy experiments
- 2) theoretical models
- 3) high p_T J/ψ results

➔ suppression stronger at higher \sqrt{s} , as expected from QGP dissociation

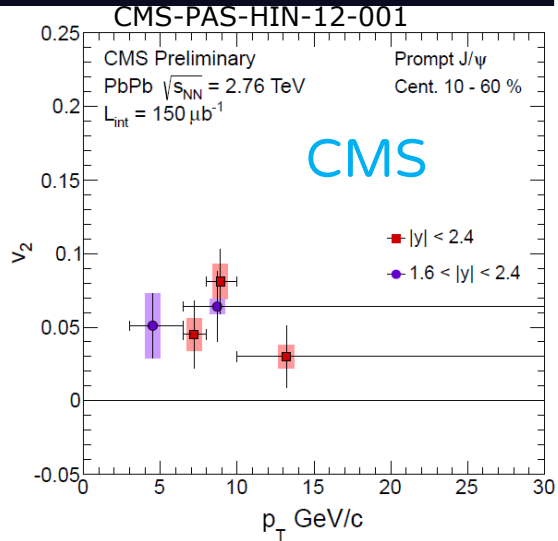
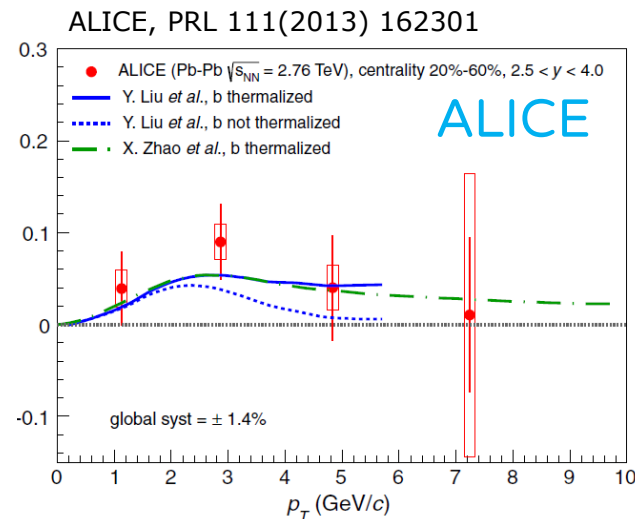
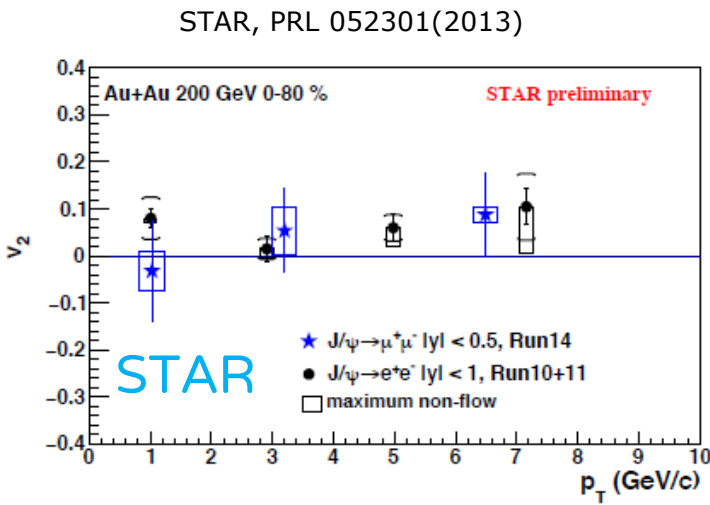
➔ opposite J/ψ behavior compared to low- p_T results

➔ negligible re(combination) effects expected at high p_T



Transport model:
Model I at RHIC: PLB 678 (2009) 72
Model I at LHC: PRC 89 (2014) 054911
Model II at RHIC: PRC 82 (2010) 064905
Model II at LHC: NPA 859 (2011) 114

➔ If c quarks participate to QGP collective motion, they should acquire elliptic flow
 ➔ J/ψ from (re)combination should inherit the flow of c quarks

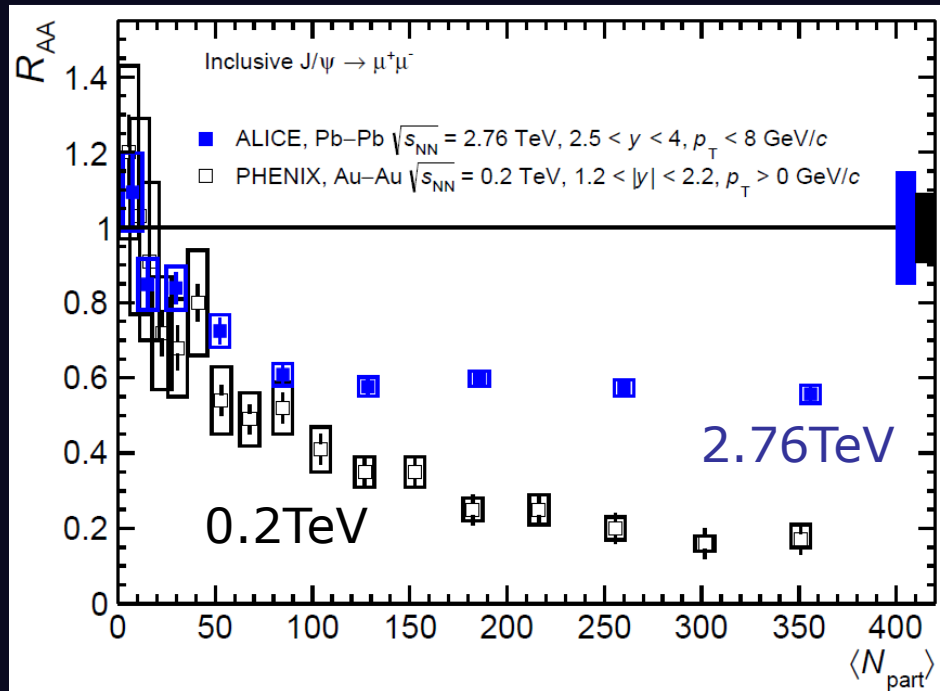


➔ Hint for J/ψ flow at LHC, contrary to $v_2 \sim 0$ observed at RHIC!

ALICE: qualitative agreement with transport models including regeneration
 CMS: path-length dependence of energy loss?

J/ψ results from LHC Run-2

LHC: Pb-Pb collisions @ $\sqrt{s_{NN}}=5.02\text{TeV}$

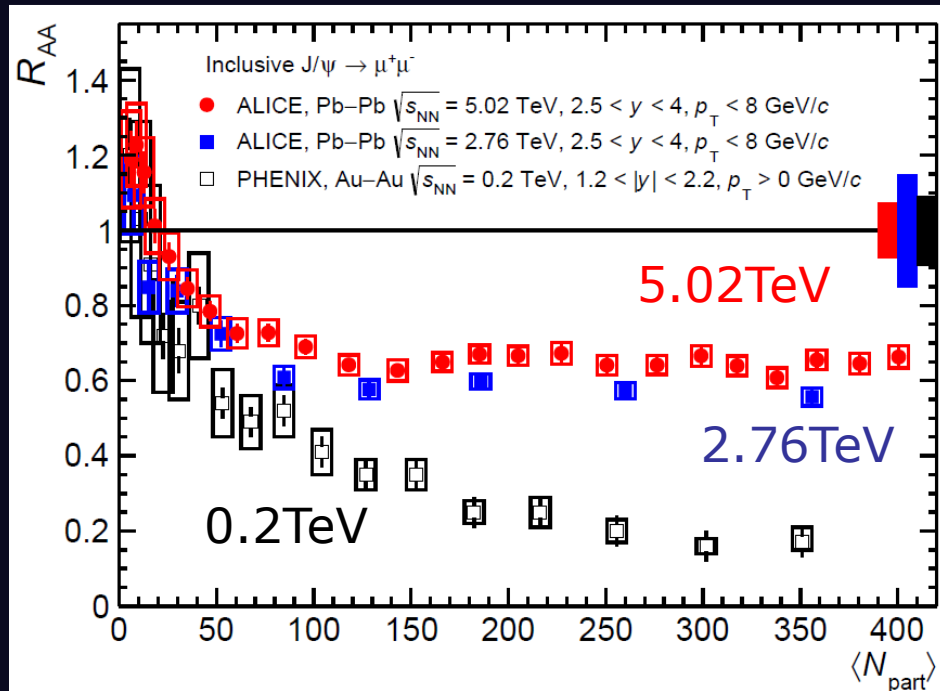


arXiv:1606.08197

J/ψ results from LHC Run-2²⁰

LHC: Pb-Pb collisions @ $\sqrt{s_{NN}}=5.02\text{TeV}$

High statistics Run-2 allows the R_{AA} evaluation in narrow centrality bins



Similar centrality dependence at the two energies, with an increasing suppression up to $N_{part} \sim 100$, followed by a plateau

R_{AA} @ 5.02TeV is ~15% higher than the one at 2.76TeV, even if within uncertainties

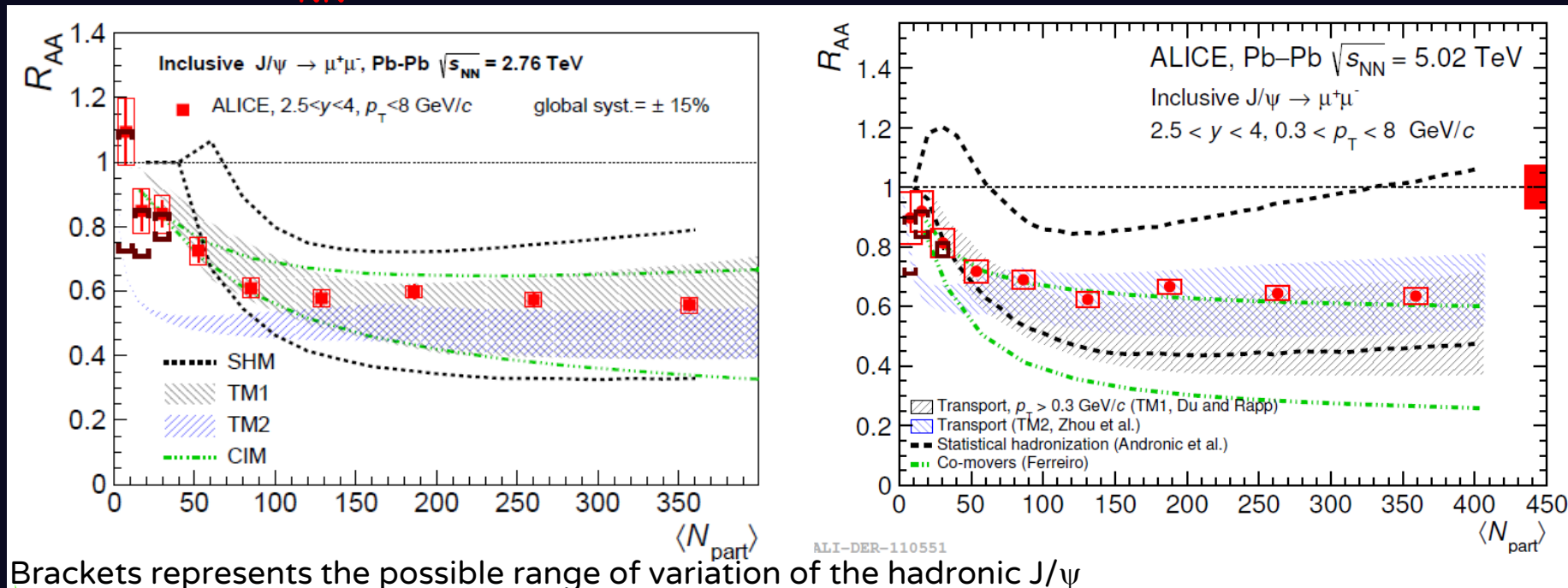
arXiv:1606.08197

J/ψ theory models

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$\sqrt{s_{NN}} = 2.76 \text{ TeV}$

$\sqrt{s_{NN}} = 5.02 \text{ TeV}$



Brackets represents the possible range of variation of the hadronic J/ψ

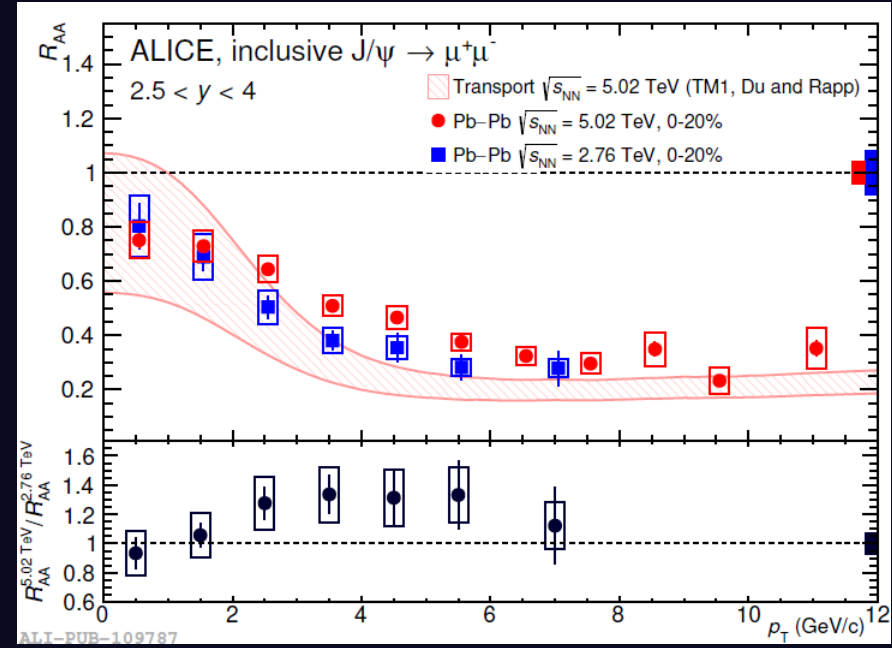
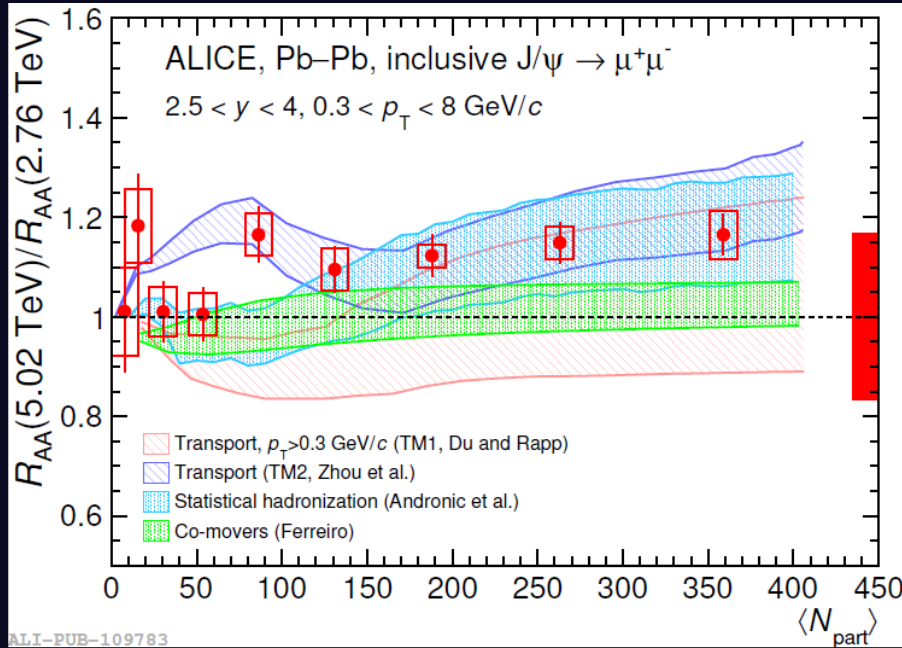
Comparison of same theory models at the two energies:

- TM1, TM2 (Du et al, Zhou et al): rate equation of suppression/regeneration in QGP
- SHM (Andronic et al): J/ψ produced by stat. hadronization at phase boundary
- CIM (Ferreiro): suppression by the comoving partonic medium and regeneration

- ➡ Data are compatible with theory models at both energies
- ➡ Still large uncertainties mainly due to the choice of σ_{cc}

Run-2 J/ ψ results

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➡ Theoretical and experimental uncertainties reduced in the R_{AA} double ratio

➡ Centrality dependence of the R_{AA} ratio is rather flat

➡ R_{AA} increases with p_T , at both energies, as expected in a regeneration scenario

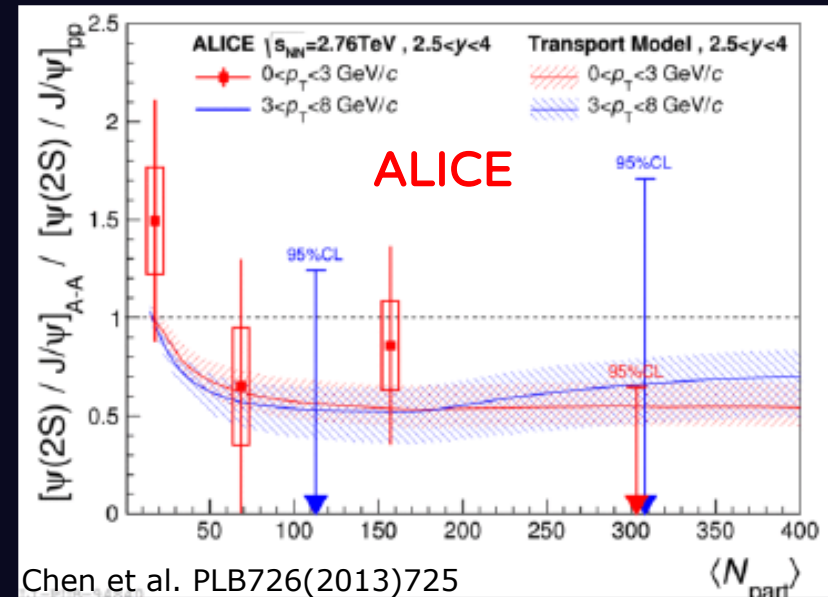
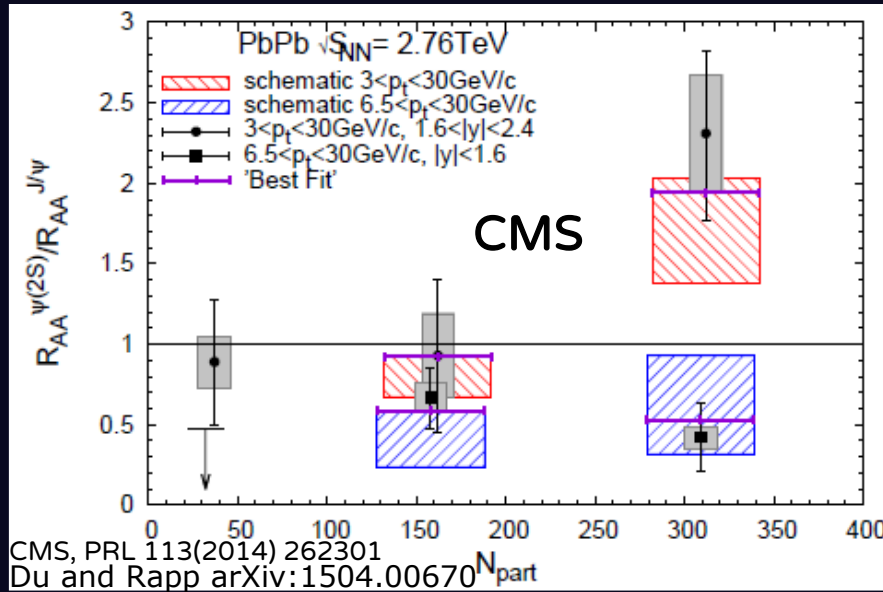
➡ Hint for an increase of R_{AA} , at 5.02 TeV, in $2 < p_T < 6 \text{ GeV}/c$

➡ Also $\sqrt{s_{NN}} = 5.02 \text{ TeV}$ results support a picture where a combination of J/ψ suppression and (re)combination occurs in the QGP

$\psi(2S)$ in AA collisions

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$\psi(2S)$ production modified in AA with a strong kinematic dependence



Fw-y, $3 < p_T < 30\text{GeV}/c \rightarrow R_{AA}^{J/\psi} < R_{AA}^{\psi(2S)}$

later $\psi(2S)$ regeneration, when radial flow is stronger, might explain the rise

Mid-y $6.5 < p_T < 30\text{GeV}/c \rightarrow R_{AA}^{J/\psi} > R_{AA}^{\psi(2S)}$

stronger suppression of $\psi(2S)$ wrt J/ψ

Fw-y, $0 < p_T < 3\text{GeV}/c \rightarrow R_{AA}^{J/\psi} > R_{AA}^{\psi(2S)}$

ALICE trend agrees with transport models and stat. hadronization approach

JHEP 05 (2016) 179

Run1 data not precise enough to conclude on $\psi(2S)$ behavior
Run2 results eagerly awaited!

Υ (ns) production in AA

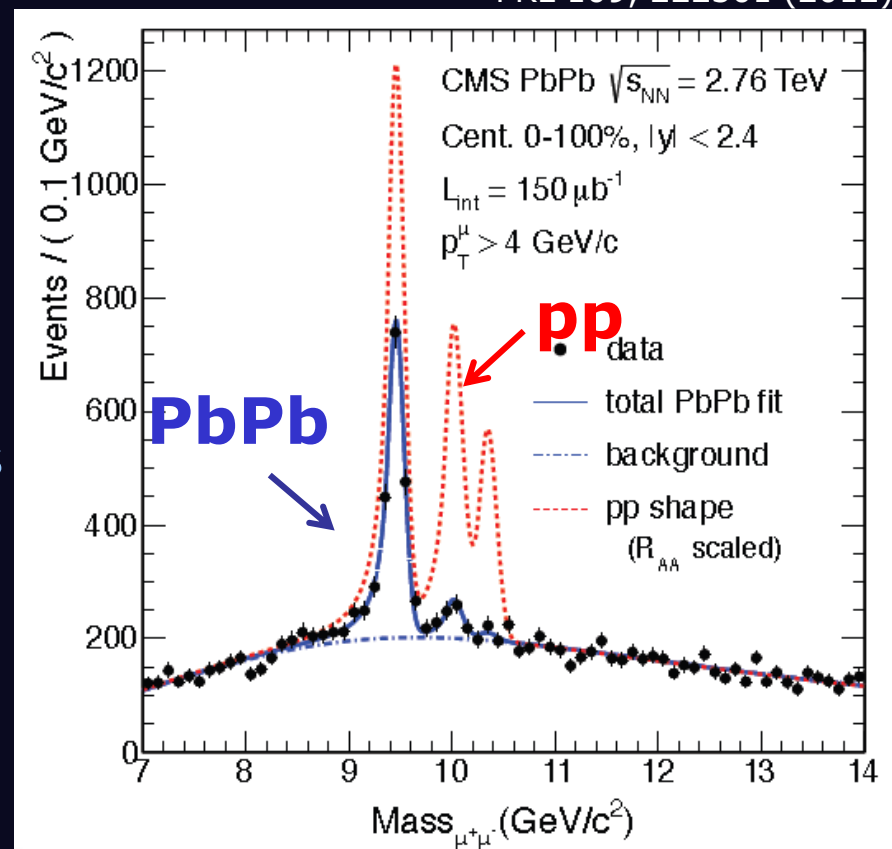
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PRL 109, 222301 (2012)

➔ Main features of bottomonium production wrt charmonium:

- no B hadron feed-down
- smaller gluon shadowing effects
- negligible (re)combination
- more robust theoretical predictions due to the higher b quark mass

with a drawback...smaller production cross-section



➔ Clear suppression of Υ states in PbPb at LHC energies with respect to pp collisions

Run-1 Υ (ns): where we stand? 25

➔ Sequential suppression observed at LHC in Run 1:

$$R_{AA}^{\Upsilon(3S)} < R_{AA}^{\Upsilon(2S)} < R_{AA}^{\Upsilon(1S)}$$

$$R_{AA}(\Upsilon(1S)) = 0.43 \pm 0.03 \pm 0.07$$

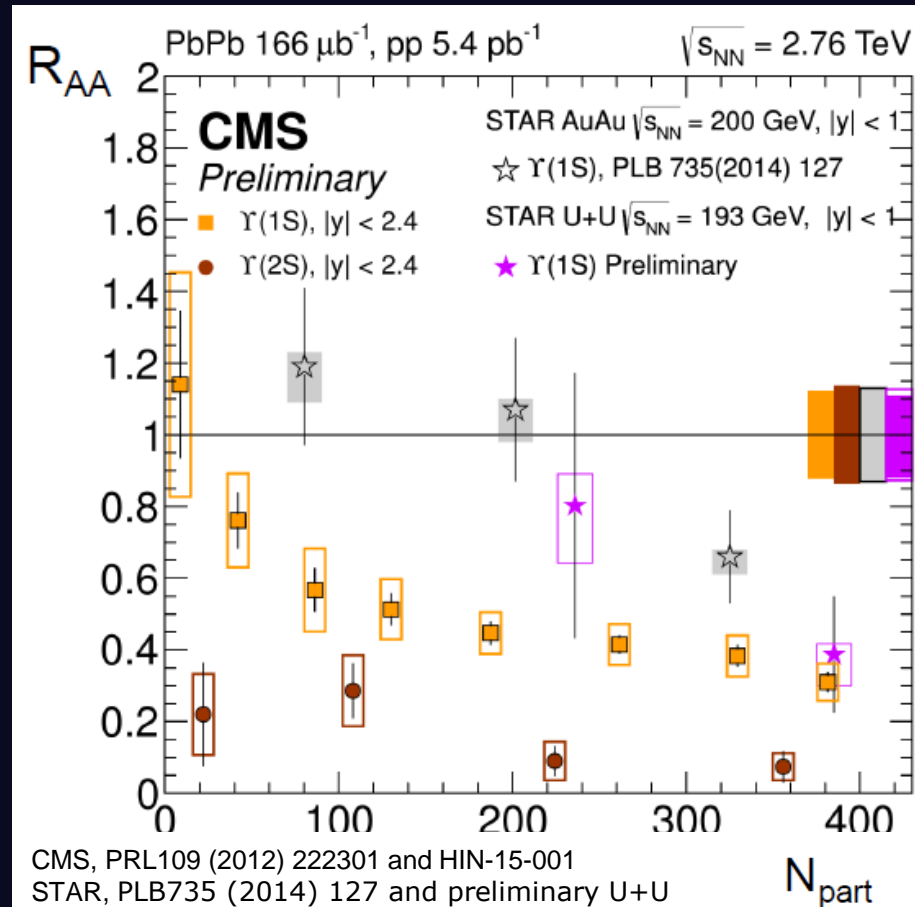
$$R_{AA}(\Upsilon(2S)) = 0.13 \pm 0.03 \pm 0.02$$

$$R_{AA}(\Upsilon(3S)) < 0.14 \text{ at } 95\% \text{ CL}$$

➔ centrality dependent suppression for $\Upsilon(1S)$ and $\Upsilon(2S)$

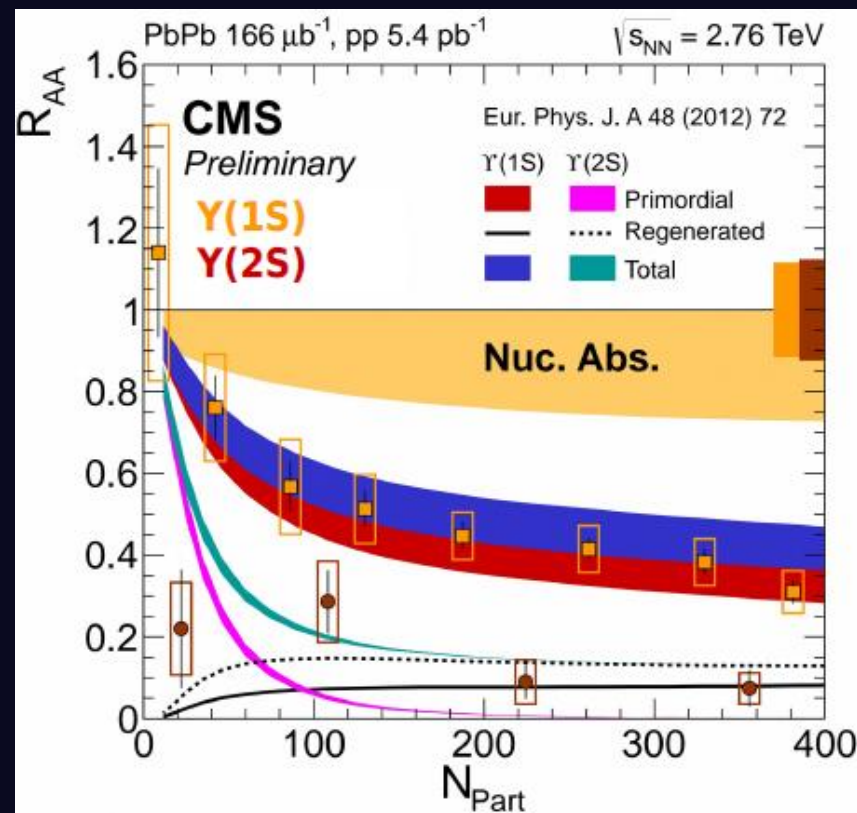
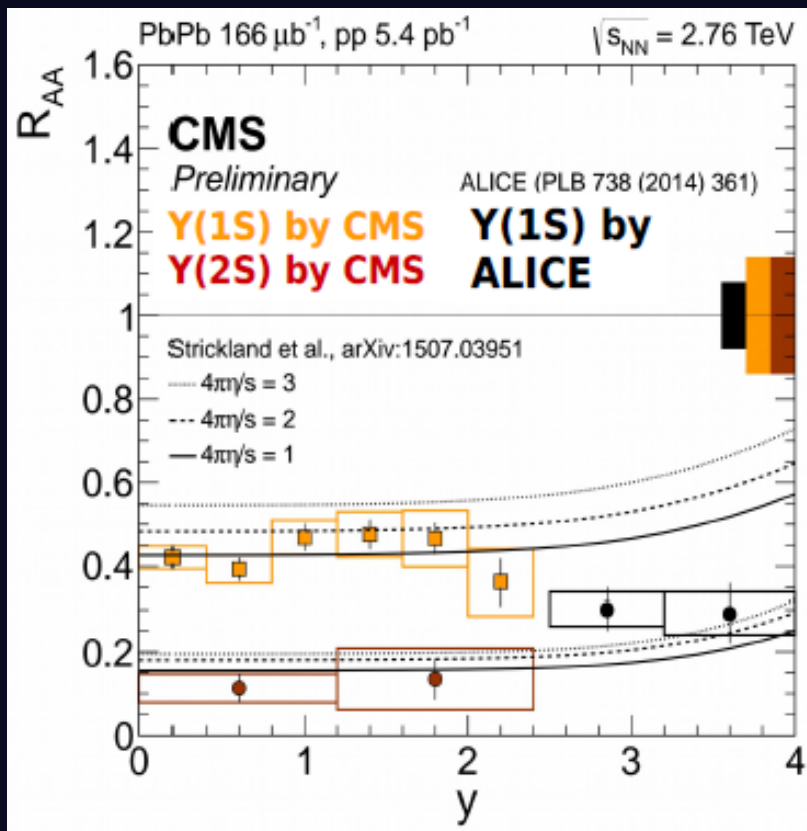
➔ at LHC $\Upsilon(1S)$ is already suppressed in semiperipheral collisions, while at RHIC only in the central ones

➔ feed-down from excited states + CNM are enough to explain the observed $\Upsilon(1S)$ suppression?



LHC Run-1 Υ (ns) results

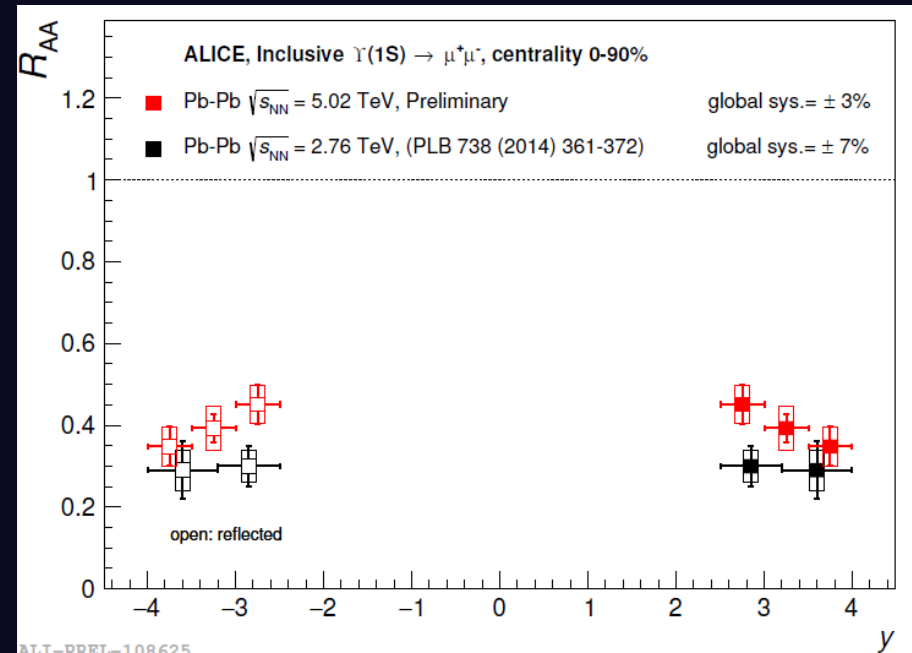
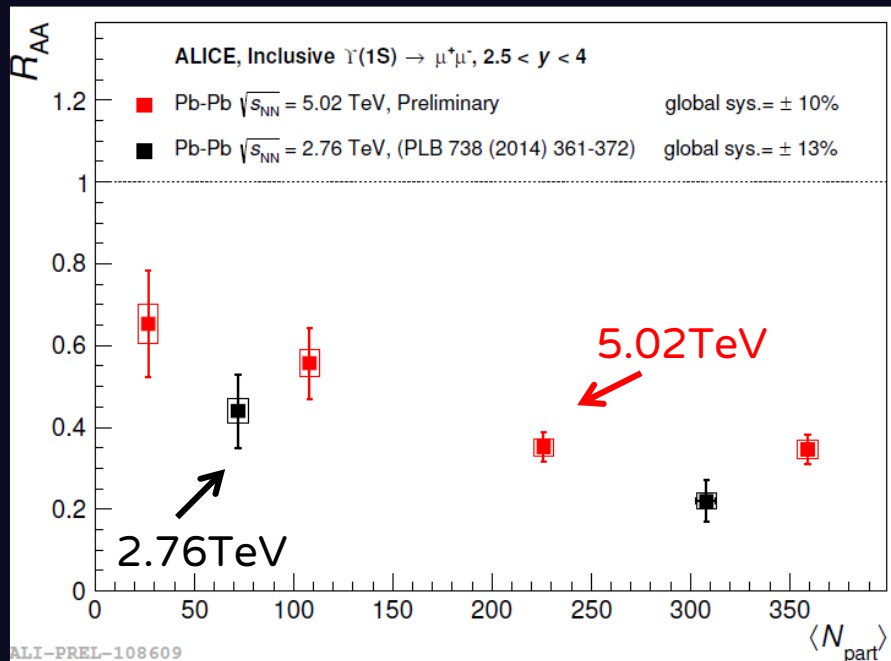
26



- ➔ no p_{T} or y dependence of the $\Upsilon(1S)$ and $\Upsilon(2S)$ suppressions
- ➔ models reproduce the p_{T} and centrality dependence
- ➔ rapidity description still needs tuning

LHC Run-2 $\Upsilon(1S)$ results

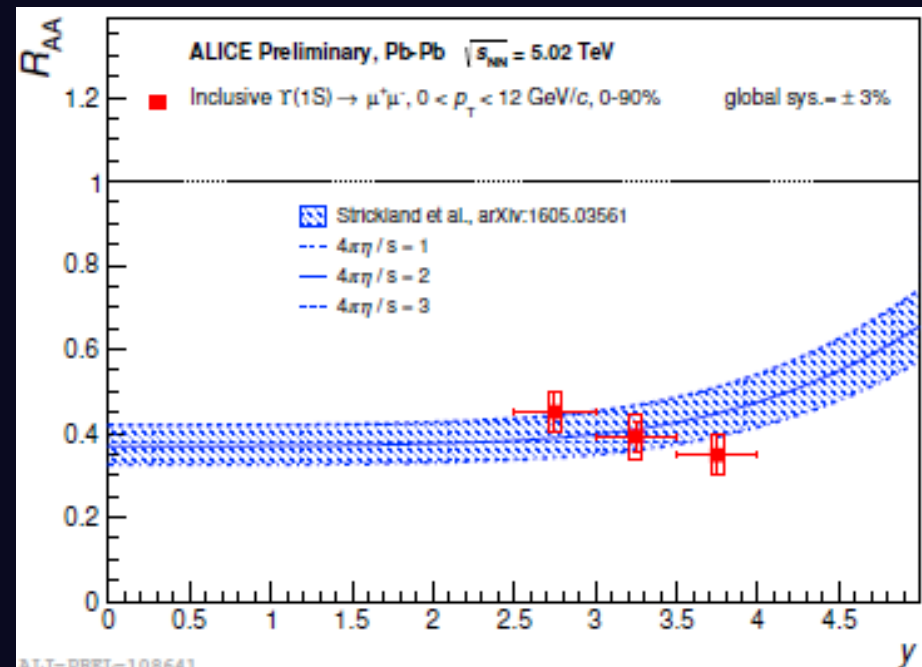
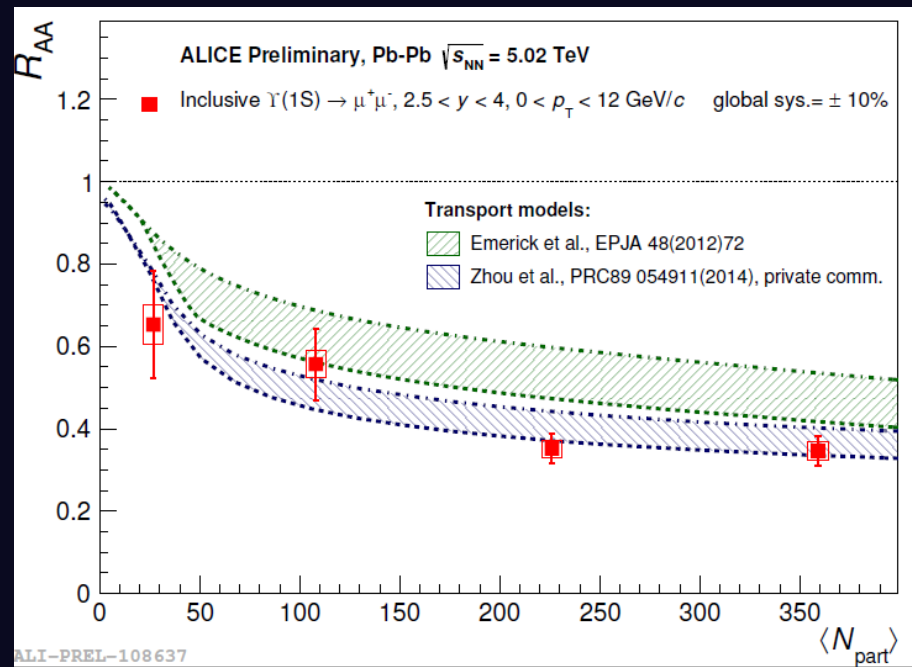
27



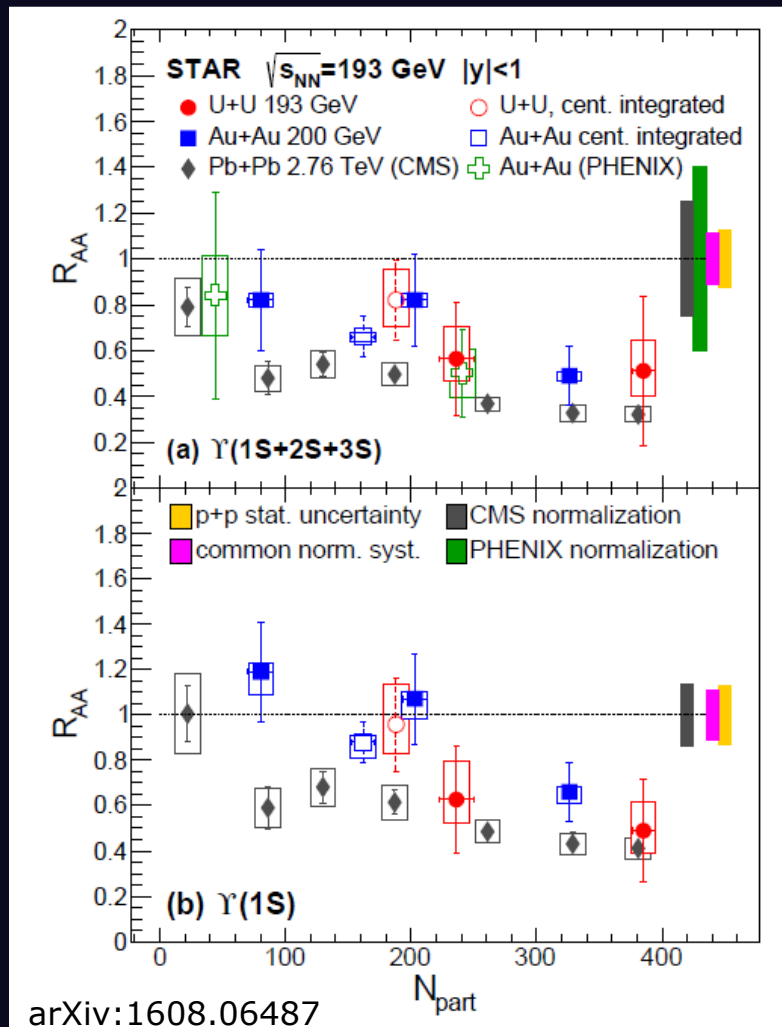
- ➔ Centrality dependent $\Upsilon(1S) R_{AA}$ suppression observed also at $\sqrt{s_{NN}} = 5.02$ TeV
- ➔ No firm conclusion on the R_{AA} energy dependence within the current uncertainties

Υ (1S) theory models

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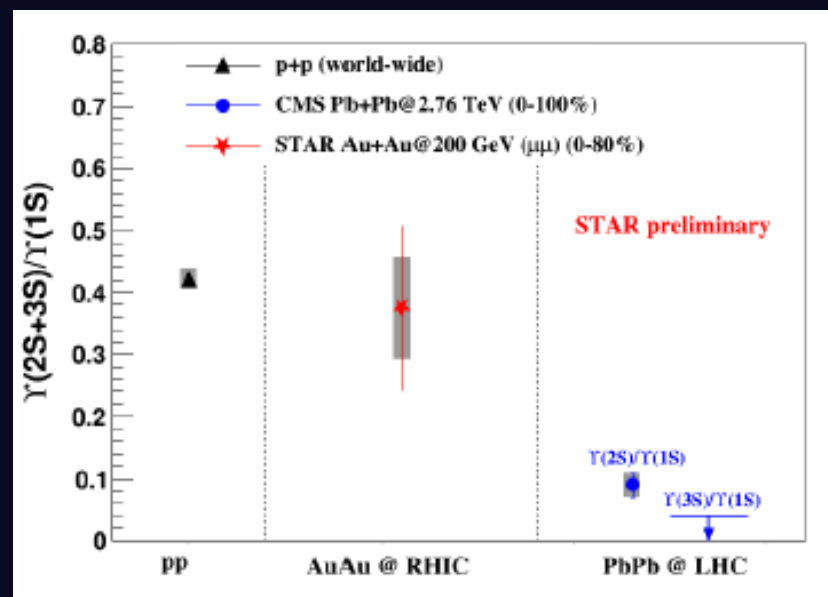


- ➔ Theory models, with (Emerick et al.) or without (Zhou et al.) regeneration component, qualitatively reproduce the data within uncertainties
- ➔ Different trend in data and theory for most forward- y ?



→ $\Upsilon(1S)$ is also suppressed at RHIC, in central collisions, even if less wrt LHC

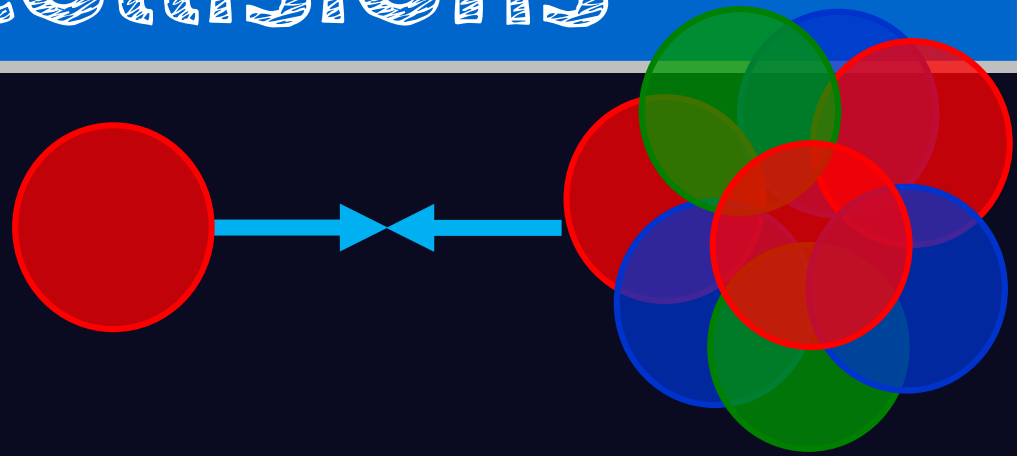
$$R_{AA}(\Upsilon(1S)) = 0.63 \pm 0.13 \pm 0.09 \text{ (AuAu+UU)}$$



→ STAR: excited states accessible in the muon channel

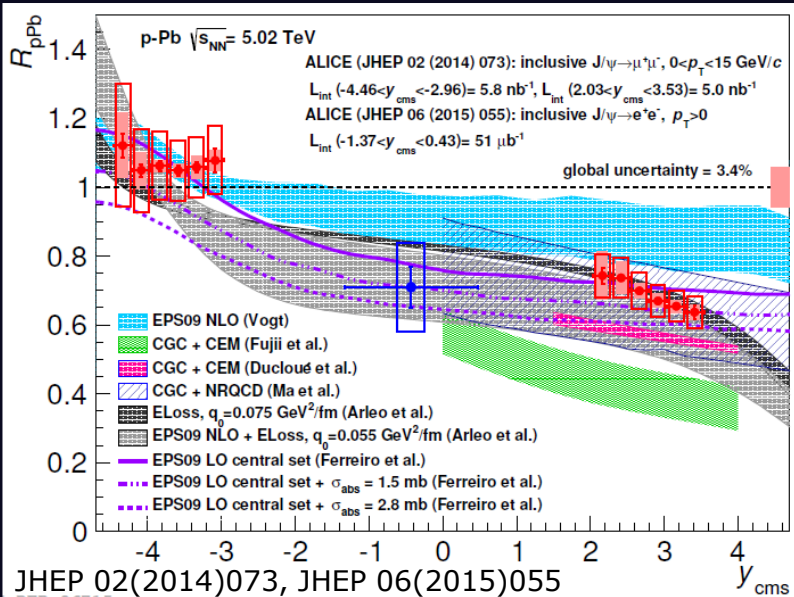
→ Hint of less suppression of excited states wrt LHC

Quarkonium in p-A collisions



R_{pA}/ψ results at LHC

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J/ψ affected by CNM effects, with a strong y and p_T dependence:

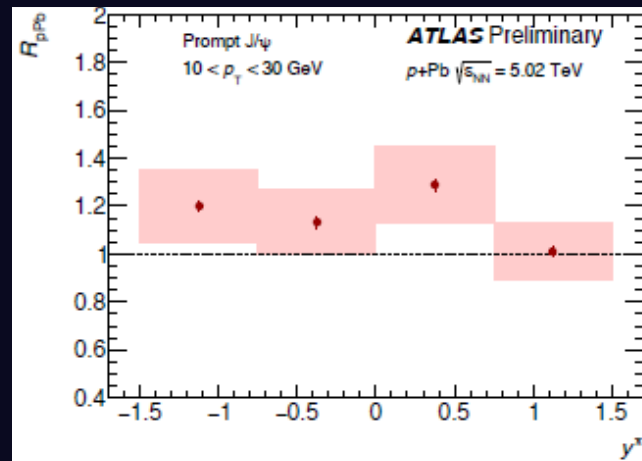
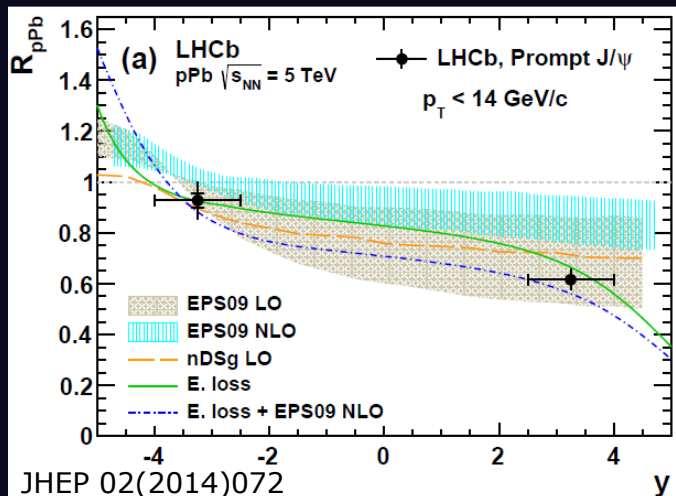
$\rightarrow R_{pA}$ decreases towards forward y

\rightarrow data consistent with shadowing and coherent parton energy loss models

\rightarrow agreement with CGC depends on implementation

\rightarrow good agreement between ALICE and LHCb (similar kinematic range)

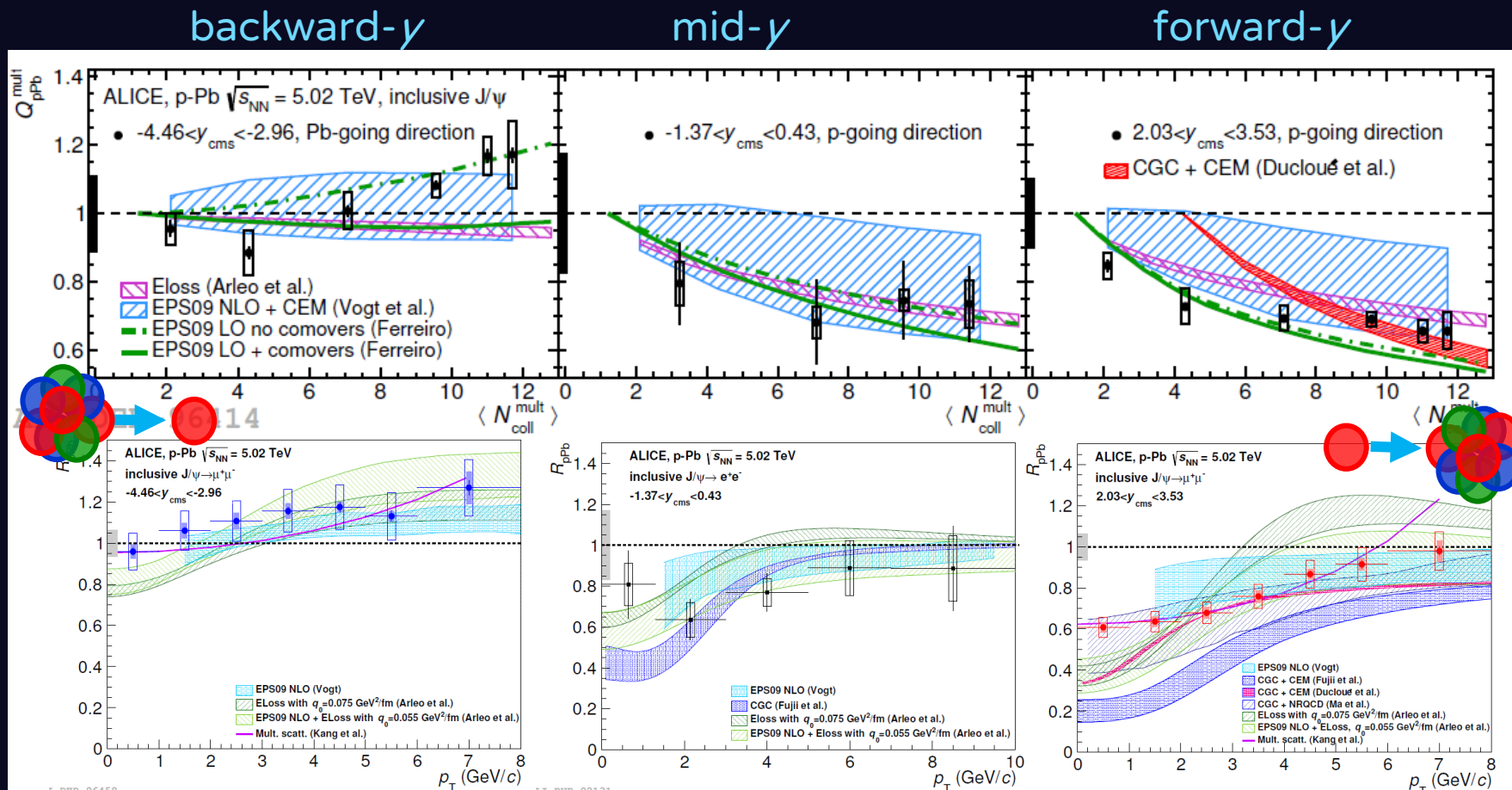
\rightarrow different behavior at mid- y for low and high p_T J/ψ



J/ψ vs p_T and centrality

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transv. momentum centrality



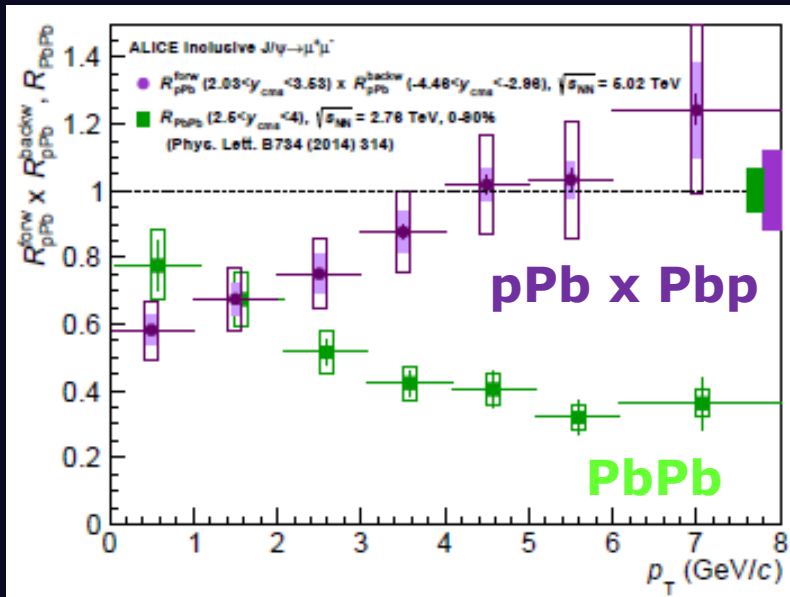
mid and fw-y: suppression increases vs centrality and is larger at low p_T
 backward-y: hint for increasing Q_{pA} vs centrality, with rather flat p_T trend

➔ Shadowing and coherent energy loss models in fair agreement with data

From pA to AA

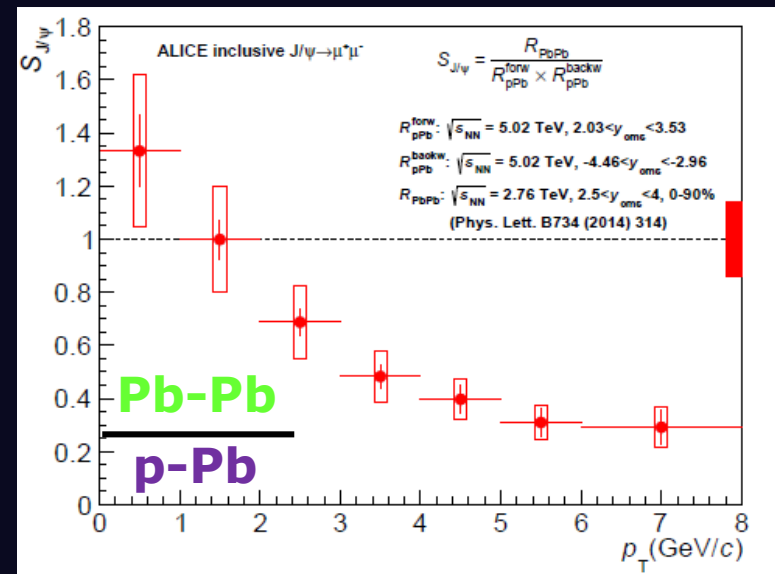
Once CNM effects are measured in pPb, what can we learn on J/ψ production in PbPb?

- Hypothesis:
- 2→1 kinematics for J/ψ production
 - CNM effects (dominated by shadowing) factorize in p-A
 - CNM obtained as $R_{pA} \times R_{Ap}$, similar x-coverage as PbPb



CNM effects not enough to explain PbPb data at high p_T

we get rid of CNM effects with **AA / pA x Ap**



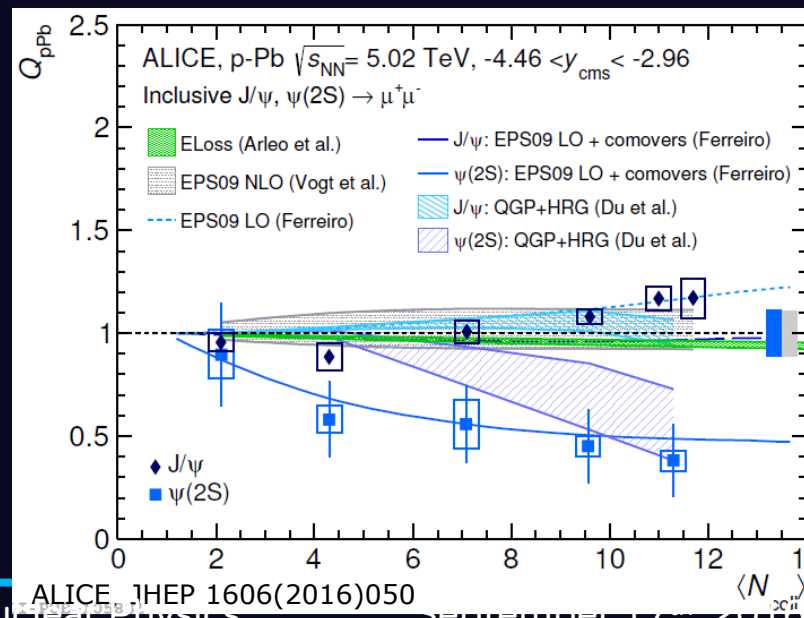
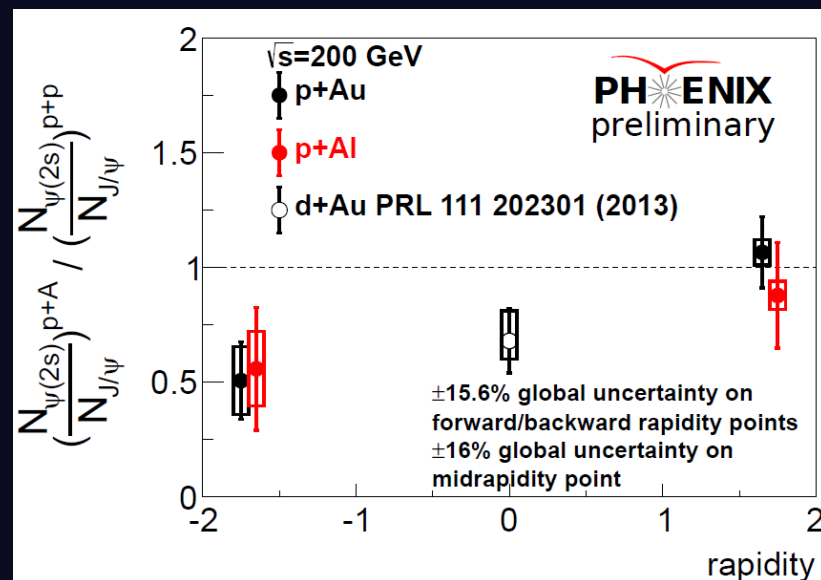
Evidence for hot matter effects in Pb-Pb!

$\psi(2S)$ production in pA, dA 34

➔ $\psi(2S)$ suppression is stronger than the J/ψ one, both at RHIC and LHC

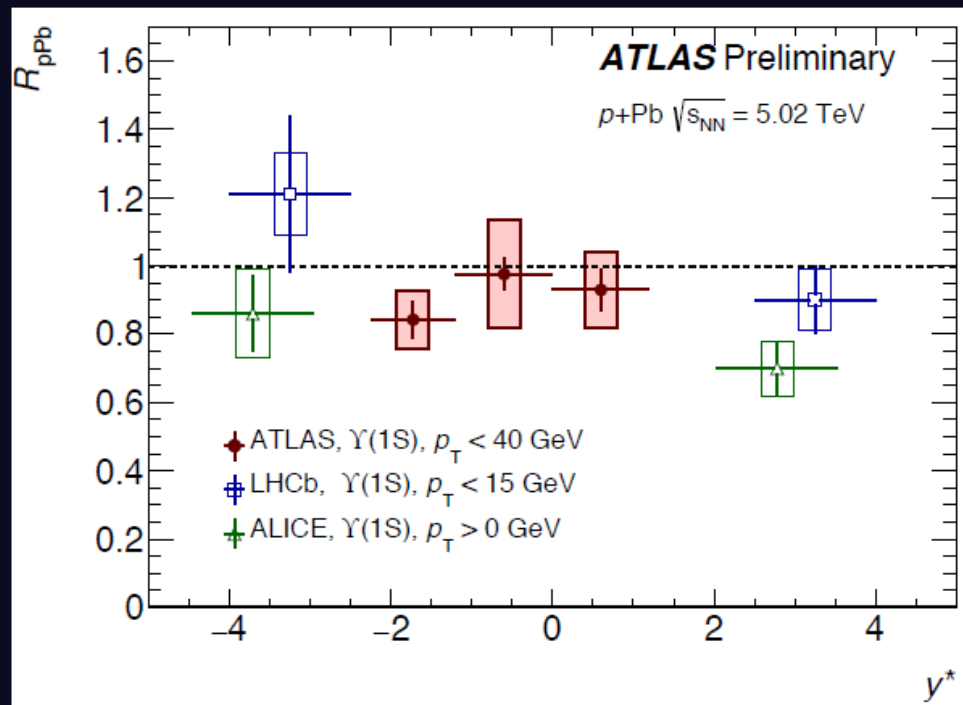
- ➔ unexpected since time spent by the $c\bar{c}$ in the nucleus (τ_c) is shorter than charmonium formation time (τ_f)
- ➔ shadowing and energy loss, almost identical for J/ψ and $\psi(2S)$, do not account for the different suppression

➔ Only models including QGP + hadron resonance gas or comovers describe the stronger $\psi(2S)$ suppression



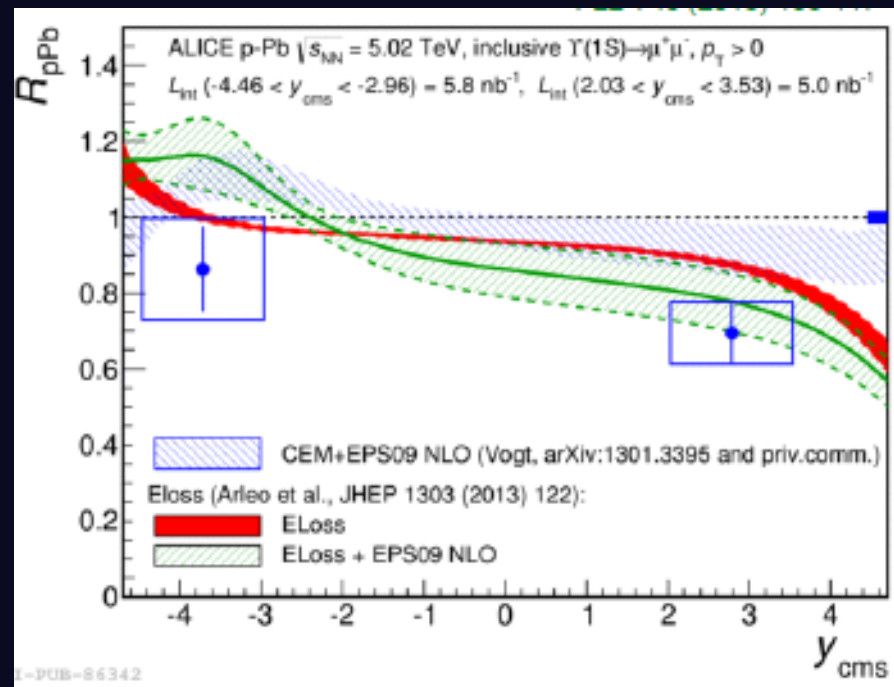
$\Upsilon(1S)$ in pA collisions

35



➔ No significant rapidity dependence of $\Upsilon(1S) R_{pA}$ (ALICE and LHCb agree within uncertainties)

➔ Shadowing and energy loss models are compatible at forward- y . At backward- y smaller anti-shadowing is suggested



ALICE, Phys. Lett. B 740 (2015) 105
 ATLAS-CONF-2015-050, LHCb, JHEP 07(2014)094

Υ excited states in pA

36

p-Pb vs pp @mid-y:

Stronger excited states suppression with respect to $\Upsilon(1S)$

Initial state effects similar for the three Υ states

→ Final state effects in p-Pb?

p-Pb vs PbPb @mid-y:

even stronger suppression of excited states in PbPb

→ ALICE (and LHCb) observes:

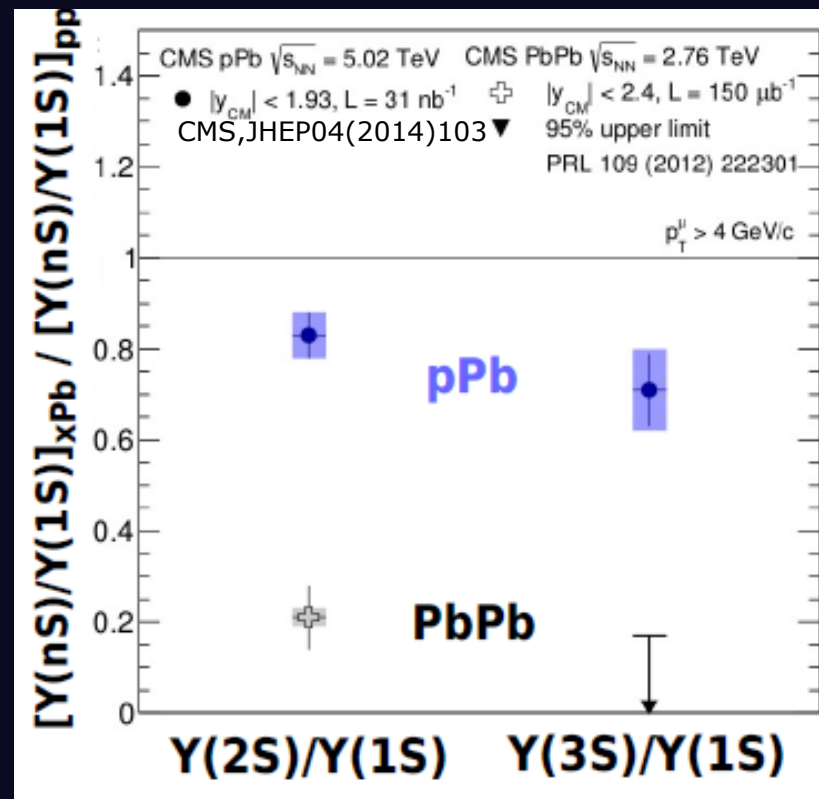
$\Upsilon(2S)/\Upsilon(1S)$ (ALICE)

$2.03 < y < 3.53$: $0.27 \pm 0.08 \pm 0.04$ (2012)

$-4.46 < y < -2.96$: $0.26 \pm 0.09 \pm 0.04$

compatible with pp results

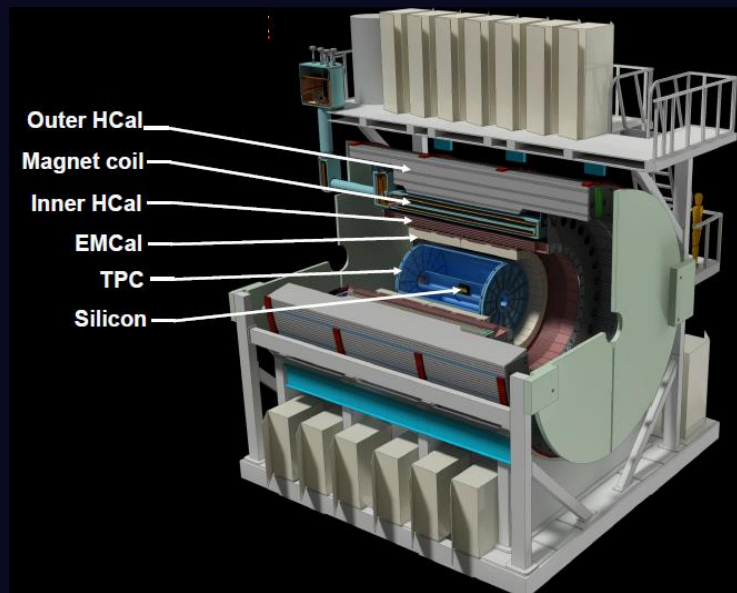
0.26 ± 0.08 (ALICE, pp@7TeV)



→ Rapidity dependent final state effects at play?

→ sPHENIX (>2020)

Precision Υ spectroscopy
(80 MeV resolution expected)



→ LHC heavy-ion program

2016: pA at $\sqrt{s_{NN}} = 5.02$ and 8 TeV

2018: PbPb

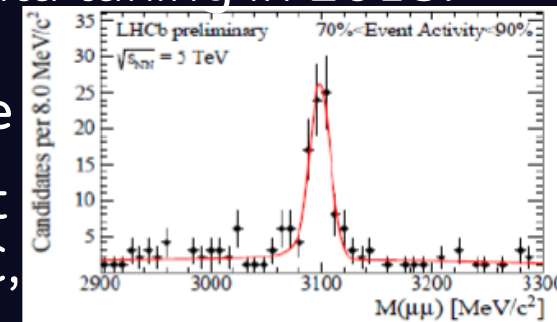
2021 - 2023: LHC Run3 – $L_{int} > 10 \text{nb}^{-1}$
for PbPb (is $\sim 1 \text{nb}^{-1}$ in Run2)

2026 – 2029 : LHC Run4

→ LHCb

Joined the PbPb data taking in 2015:
covers peripheral
semi-periph. Range

SMOG: fixed target
pA program at LHC,
up to $\sqrt{s} = 110$ GeV



Conclusions

38

A large sample of quarkonium results in various systems and at various energies is now available from both RHIC and LHC!

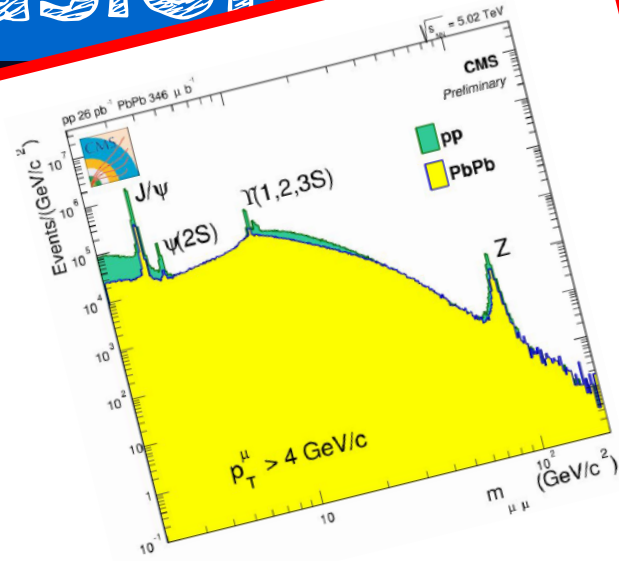
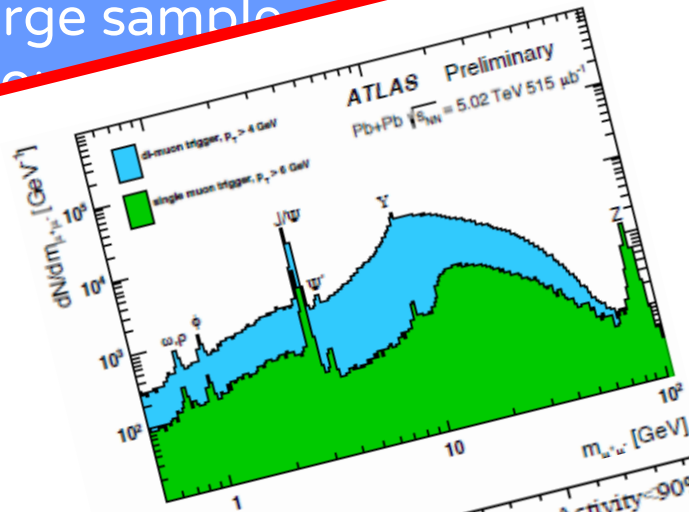
- ➔ A combination of suppression and regeneration mechanisms affects J/ψ production at RHIC and LHC
- ➔ Theory models qualitatively describe the data, but still large uncertainties (open charm cross section)
- ➔ CNM effects (mainly shadowing and energy loss) play an important role, as observed in pA collisions
- ➔ Bottomonium results might be compatible with sequential suppression in QGP

Thanks!

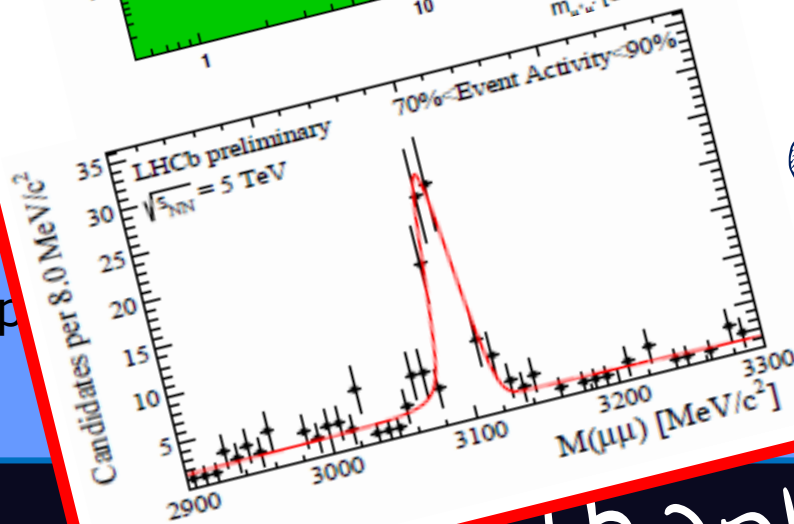
Conclusions

38

A large sample
various



Other new results
eagerly awaited!!!



Thanks!

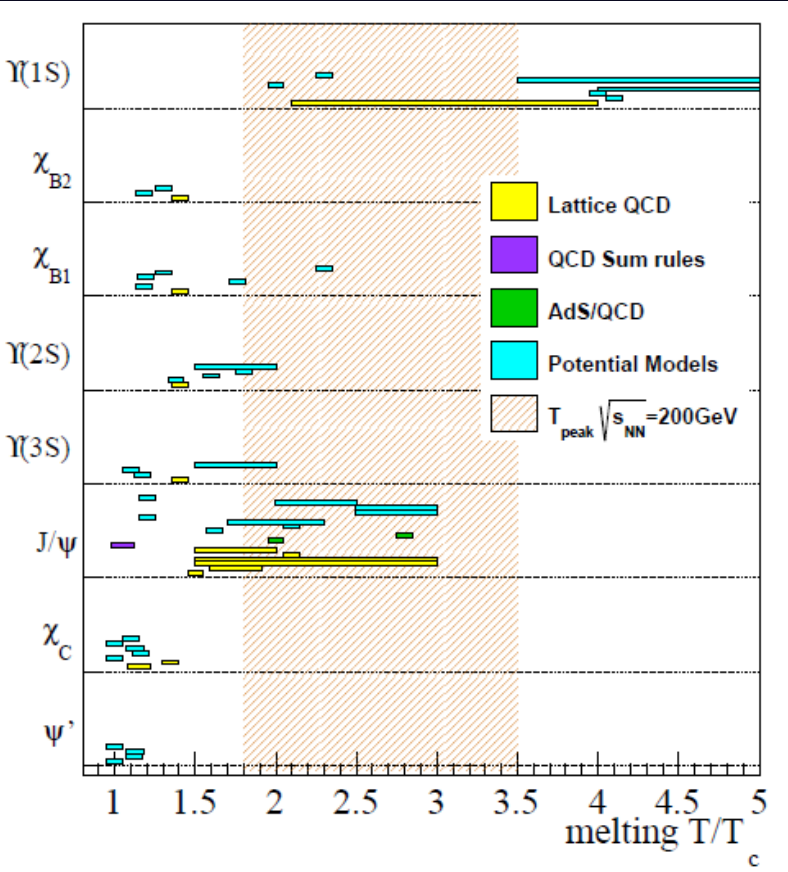
Backup slides

Backup slides

AA: from suppression...

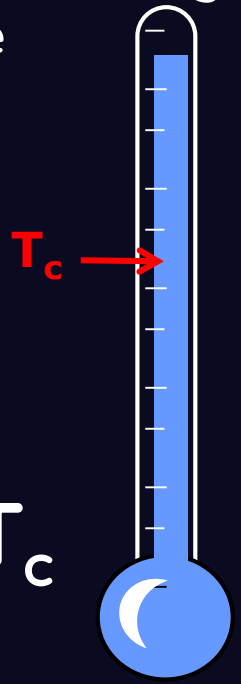
the original idea:
 quarkonium production suppressed via color screening in the QGP

sequential melting
 differences in the quarkonium binding energies lead to a sequential melting with increasing temperature



PHENIX, Phys.Rev C91, 024913

$\psi(2S)$ J/ψ $\Upsilon(1S)$



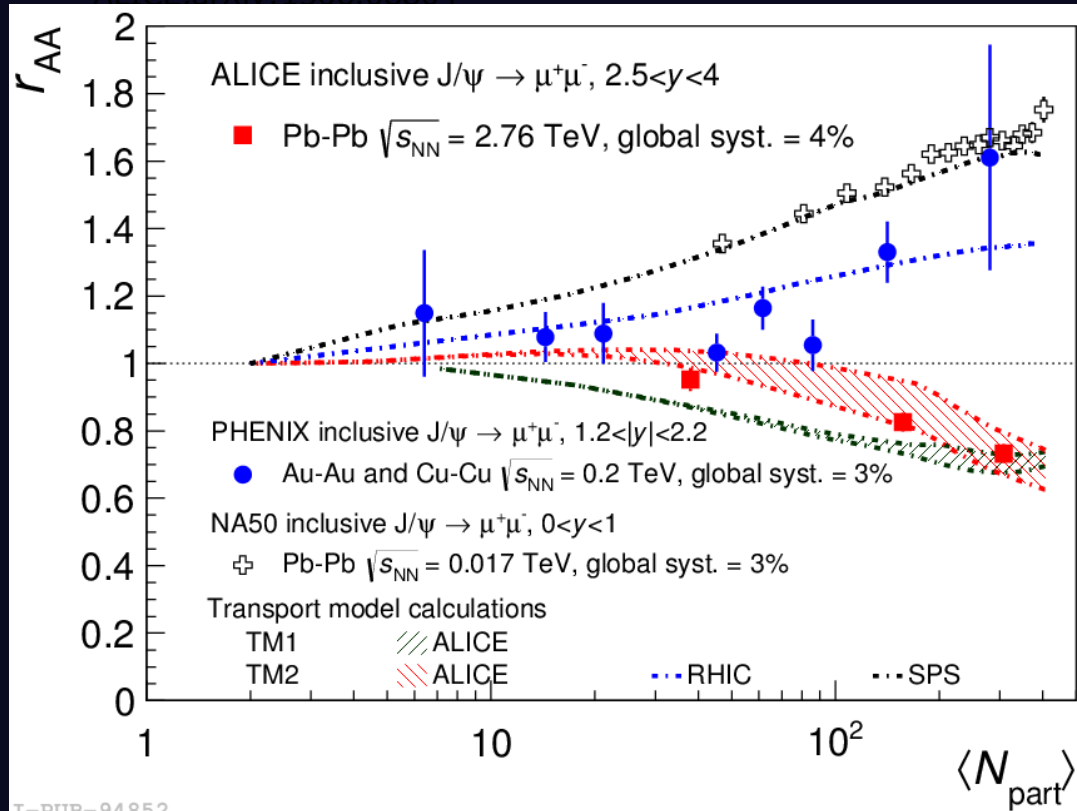
$T \gg T_c$

Quarkonium as QGP thermometer

Evolution of $J/\psi \langle p_T^2 \rangle$

42

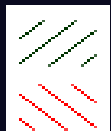
ALICE arXiv:1506.08804



$$r_{AA} = \frac{\langle p_T^2 \rangle_{AA}}{\langle p_T^2 \rangle_{pp}}$$

- r_{AA} centrality evolution strongly depends on \sqrt{s}
- decreasing r_{AA} trend, observed at LHC → due to (re)combination, which dominates J/ψ production at low p_T
- transport models, already describing $J/\psi R_{AA}$, also reproduce the r_{AA} evolution

I-PUB-94852

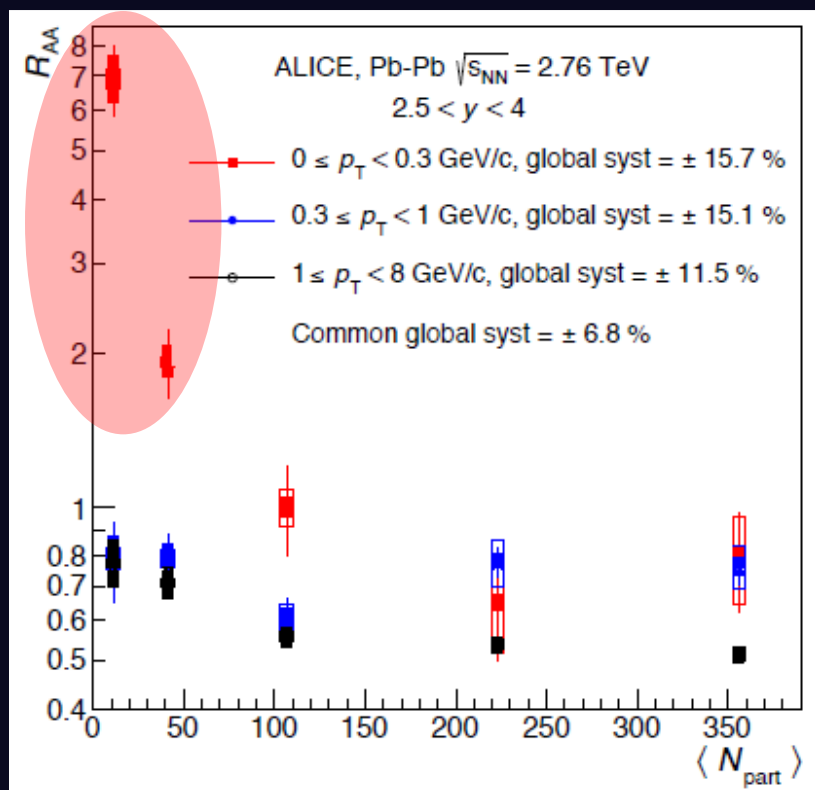


TM1: Zhao et al., Nucl.Phys.A859 (2011) 114
 TM2: Zhou et al. Phys.Rev.C89 (2014)054911

J/ψ at very low p_T

43

Strong R_{AA} enhancement in peripheral collisions for $0 < p_T < 0.3$ GeV/c

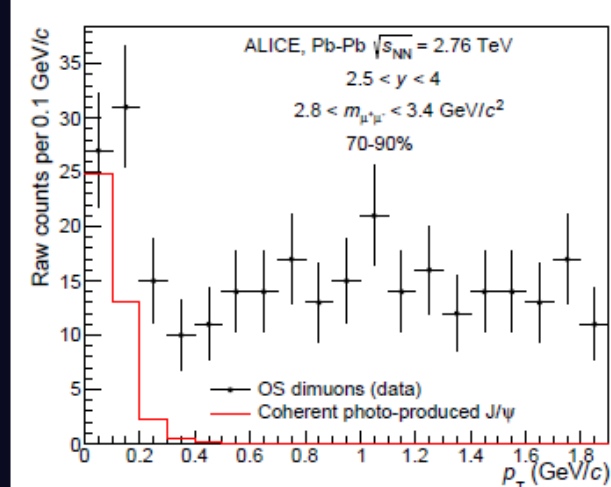


significance of the excess is 5.4 (3.4) σ in 70-90% (50-70%)

behaviour not predicted by transport models

excess might be due to coherent J/ψ photoproduction in PbPb (as measured also in UPC)

if excess is “removed” requiring $p_T^{J/\psi} > 0.3$ GeV/c \rightarrow ALICE R_{AA} lowers by 20% at maximum (in the most peripheral bin)

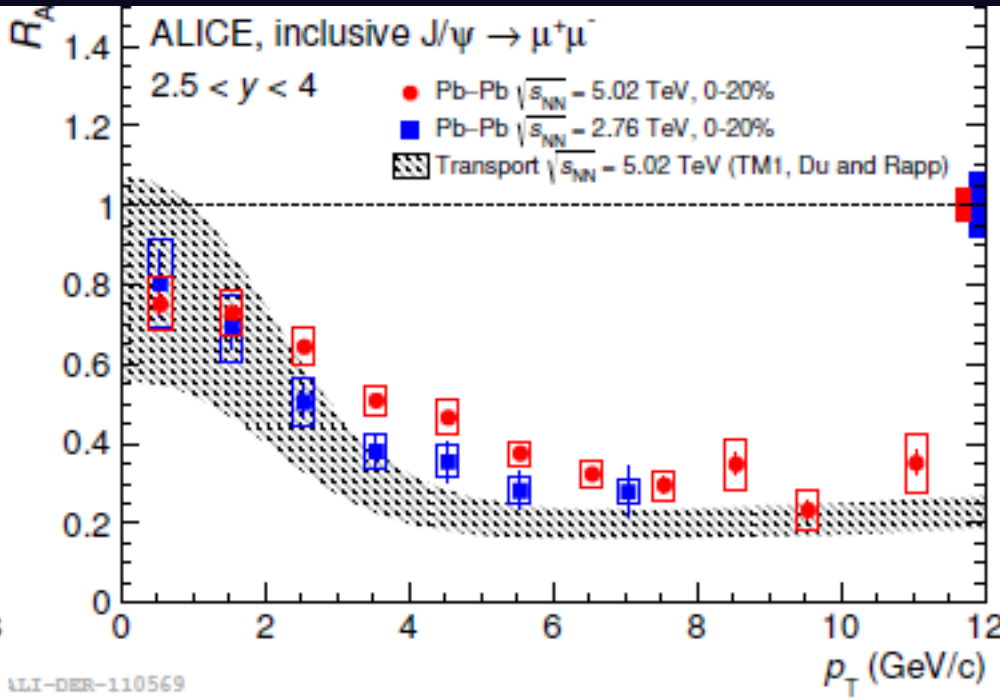
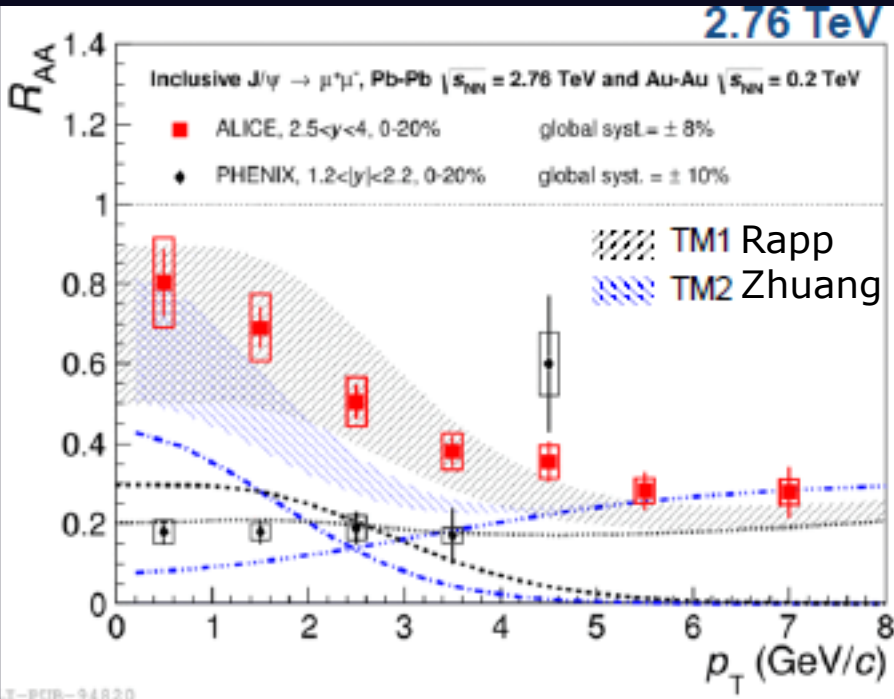


RAA vs pT

ALICE, PLB 734 (2014) 314, JHEP 07(2015)051, arXiv:1506.08804

2.76TeV

5.02TeV



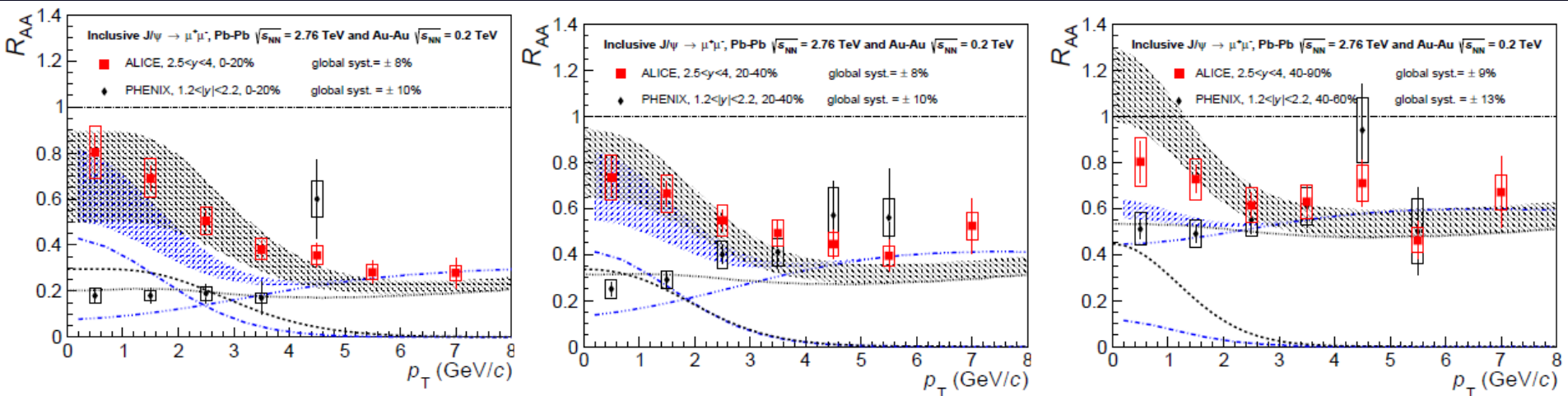
Multi-differential J/ψ studies 5

➔ p_T -centrality multi-differential studies allows detailed comparison with theory models

0-20%

20-40%

40-90%



- TM1 Zhao et al., Nucl.Phys.A859 (2011) 114
- TM2 Zhou et al. Phys.Rev.C89 (2014)054911

ALICE, arXiv:1506.08804

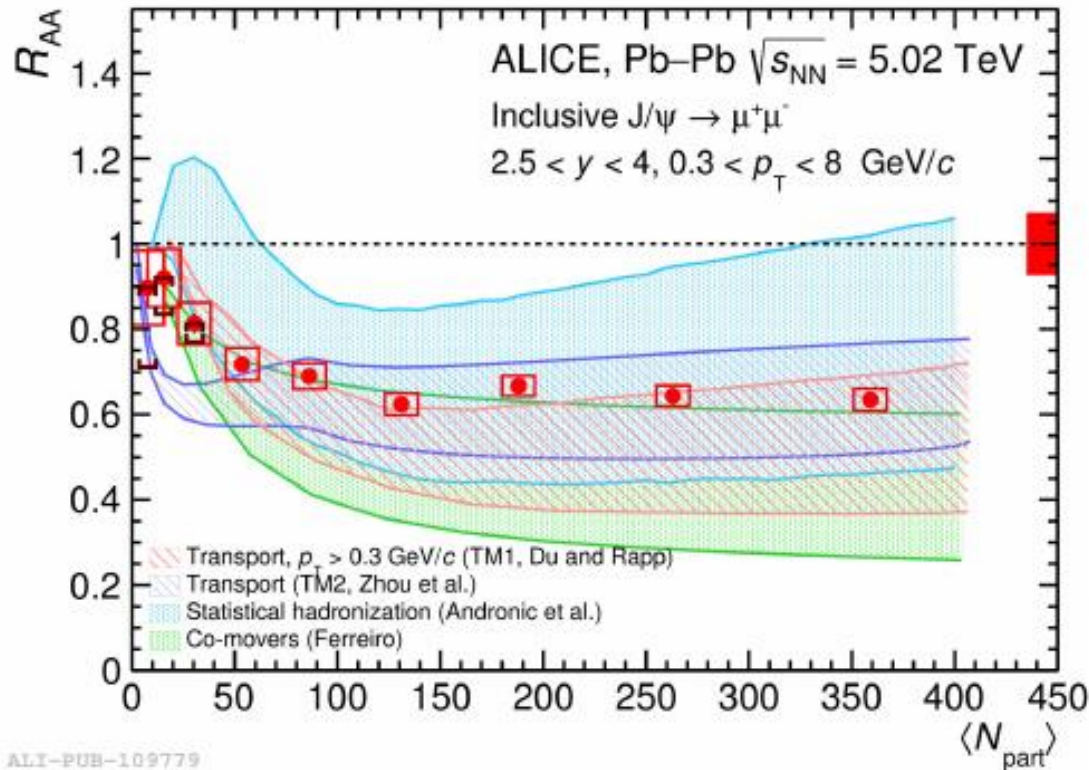
- Primordial J/ψ (TM1)
- Regenerated J/ψ (TM1)
- Primordial J/ψ (TM2)
- Regeneration J/ψ (TM2)

➔ Model provide a fair description of the data, even if with different balance of primordial/regeneration components

Still rather large theory uncertainties: models will benefit from precise measurement of σ_{cc} and CNM effects

LHC Run-2 J/ψ results

46



model	$\sigma_{c\bar{c}}$	N-N $\sigma_{J/\psi}$	comover $\sigma_{J/\psi}$	Shadowing
Transport(Rapp)	0.57 mb	3.14 μb	-	EPS09
Transport(Zhou)	0.82 mb	3.5 μb	-	EPS09
Stat. hadronization	0.45 mb	-	-	EPS09
Comovers	[0.45,0.7] mb	3.53 μb	0.65 mb	Glauber-Gribov theory

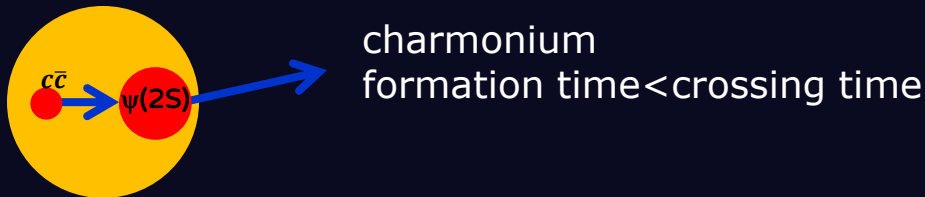
$\psi(2S)$ production in pA

47

- Being more weakly bound than the J/ψ , the $\psi(2S)$ is an interesting probe to have further insight on the charmonium behaviour in pA
- Low energy $\psi(2S)$ p-A results from NA50, E866 and HERA-B:

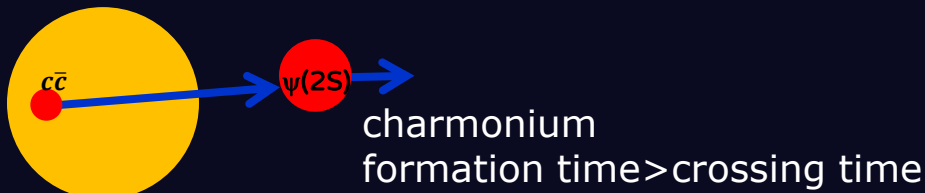
mid- y ($x_F \sim 0$):

$\psi(2S)$ suppression stronger than J/ψ one, interpreted via pair break-up
→ fully formed resonances traversing the nucleus

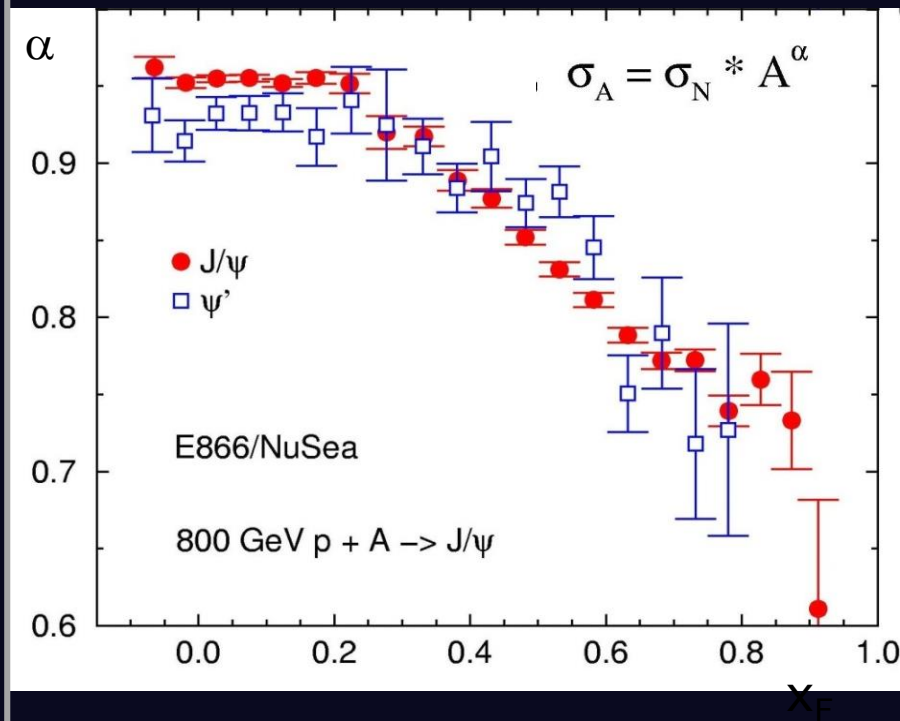


forward- y (high x_F):

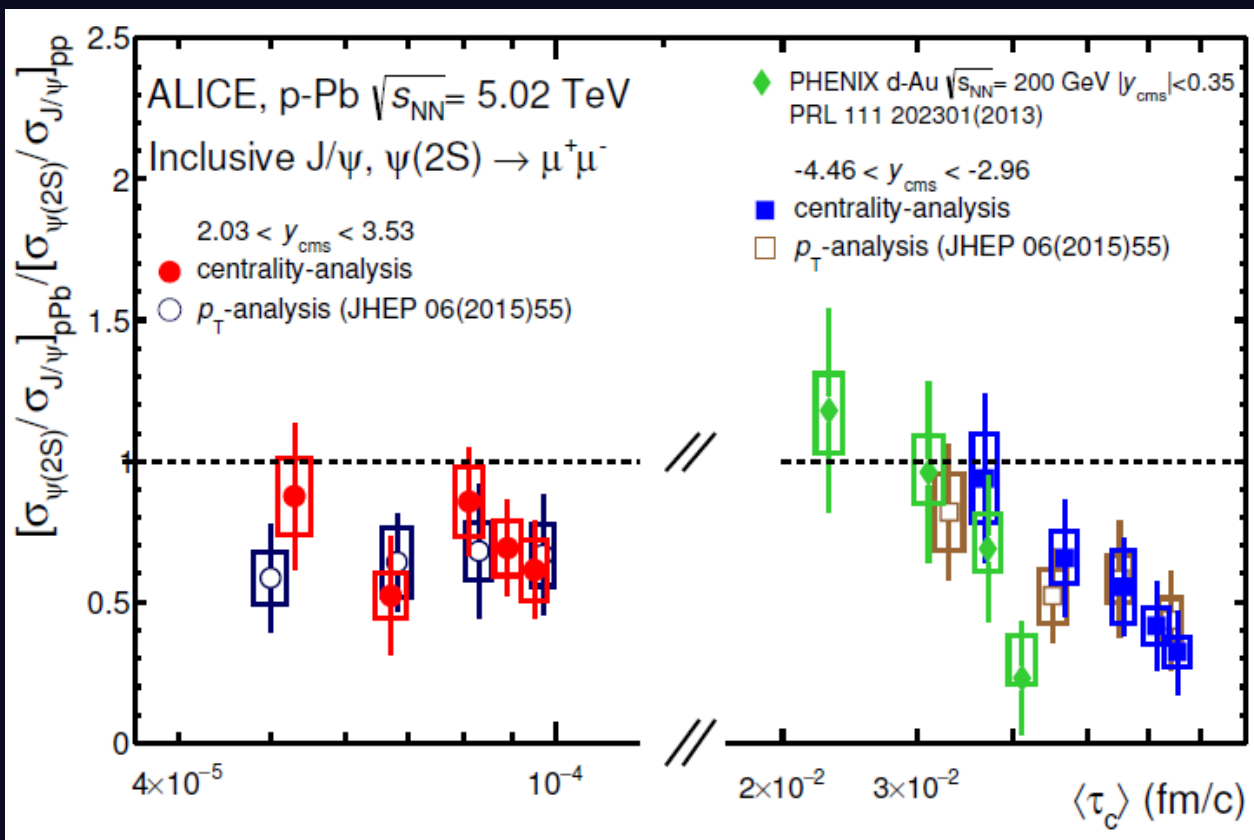
suppression becomes identical
→ dominated by energy loss



E866 Collab., PRL 84 (2000) 3256



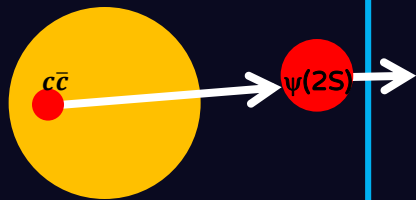
$\psi(2S)$ versus crossing time 48



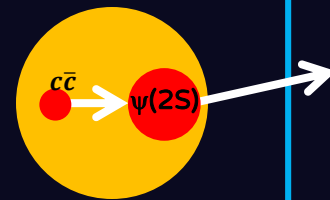
$$\tau_c = \frac{\langle L \rangle}{(\beta_z \gamma)}$$

D. McGlinchey, A. Frawley and R. Vogt, PRC 87,054910 (2013)

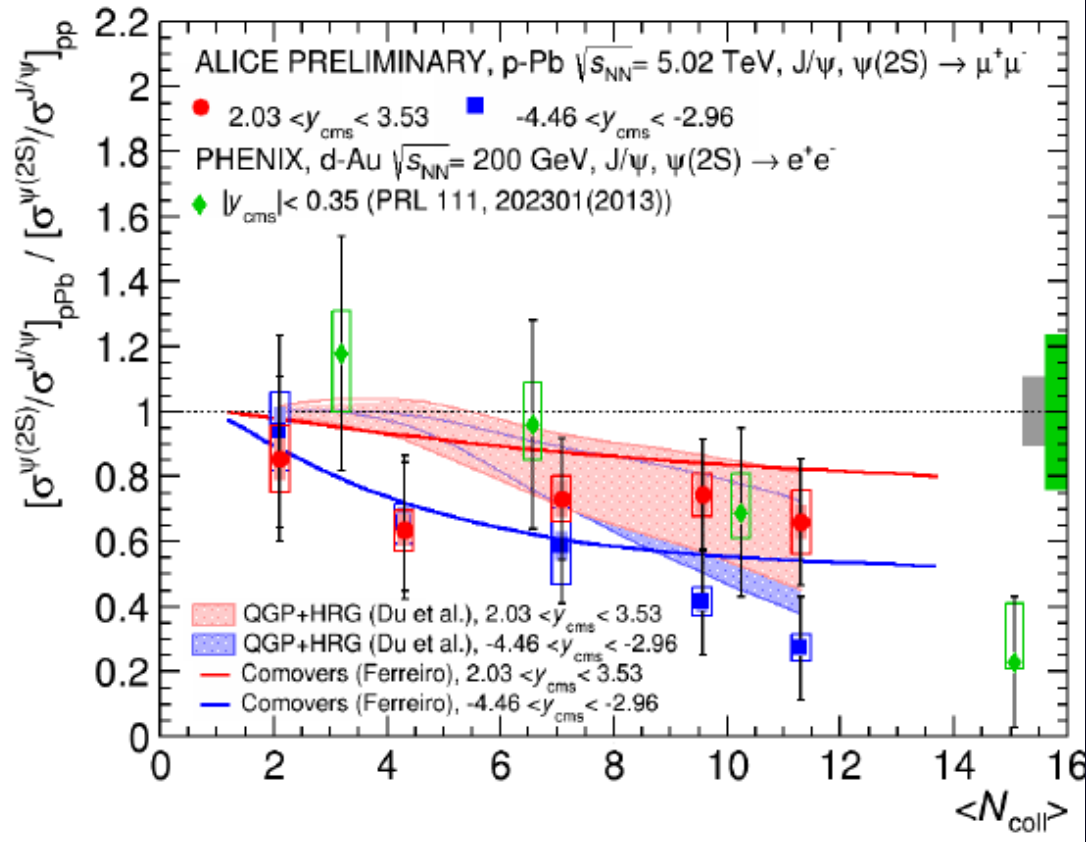
Forward- y : $\tau_c \ll \tau_f$
 interaction with
 nuclear matter
 cannot play a role



Backward- y : $\tau_c \simeq \tau_f$
 indication of effects
 related to break-up
 in the nucleus?



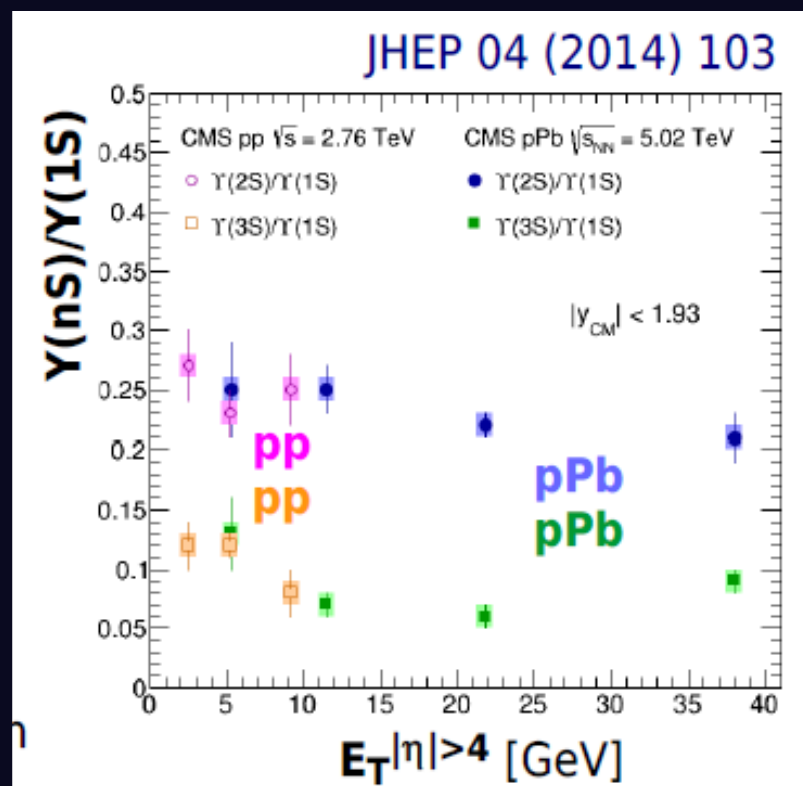
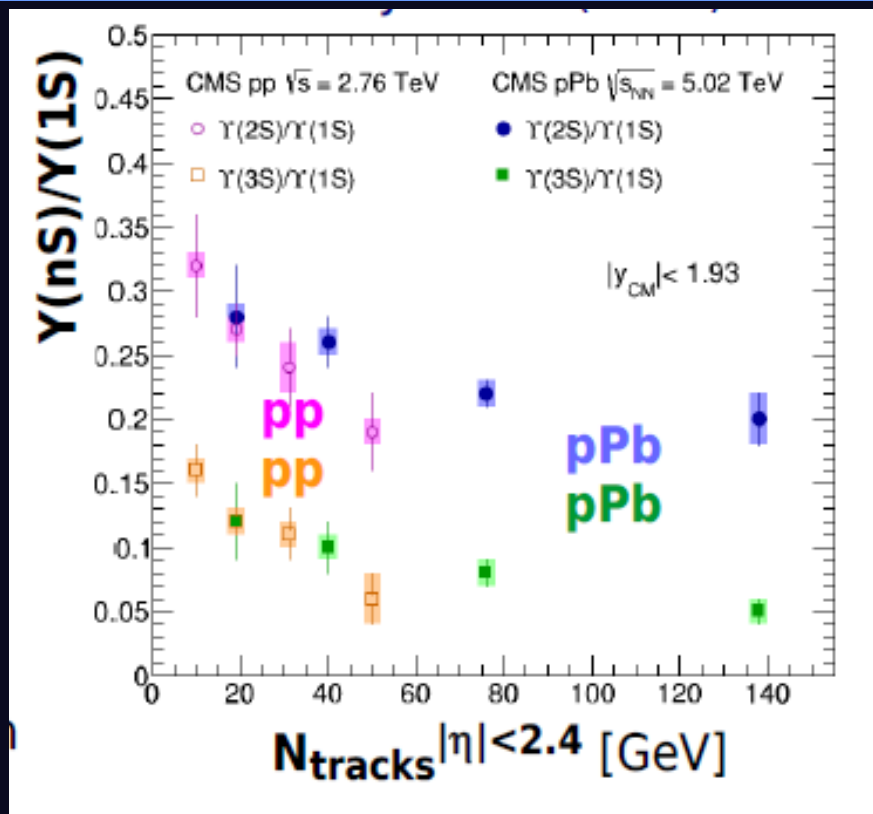
$\psi(2S) / J/\psi$ double ratio 49



Similar suppression trend observed versus centrality, by both ALICE and PHENIX

→ QGP+hadron resonance gas (Rapp) or comovers models (Ferreiro) describe the observed suppression

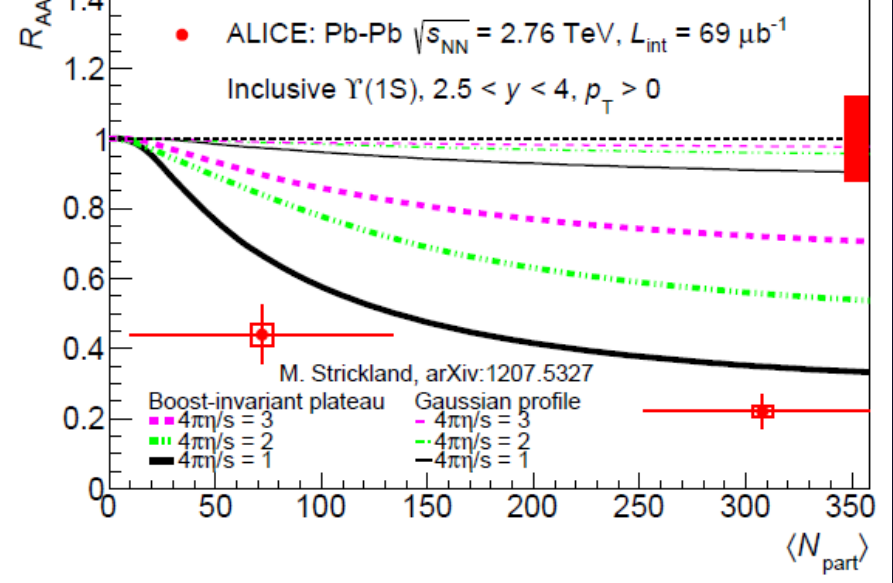
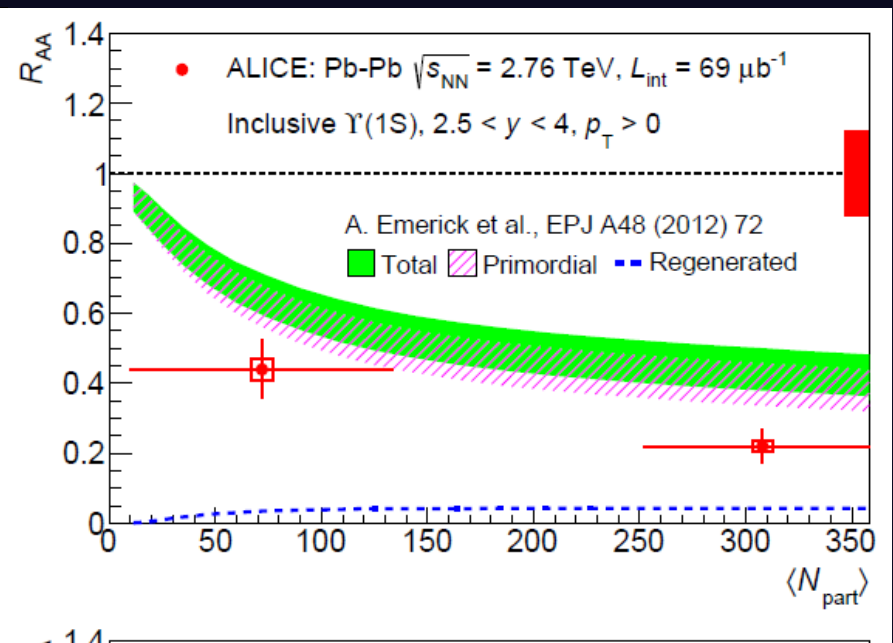
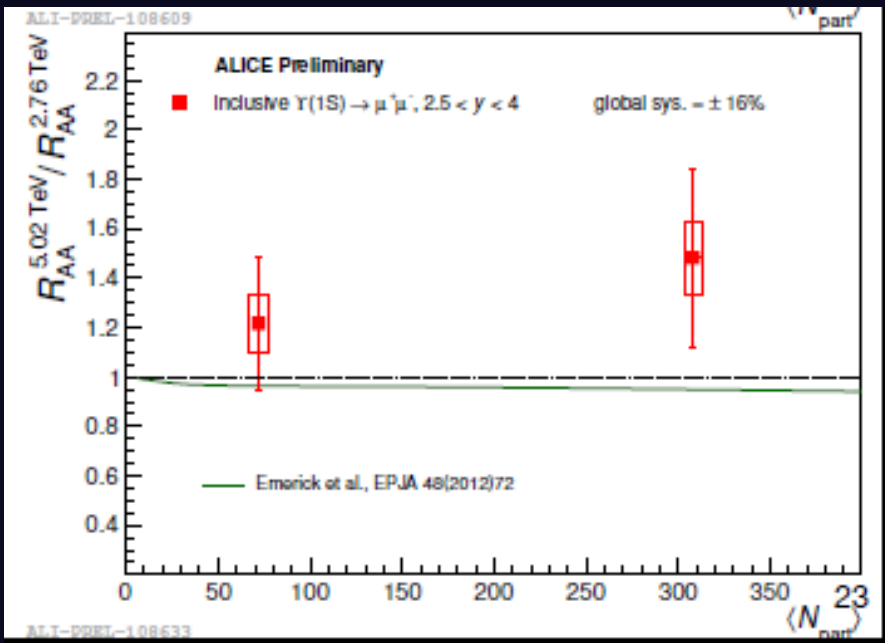
Υ vs ev. activity



- $Y(nS)/Y(1S)$ ratios fall with event-activity CMS
 - Is the multiplicity affecting the $Y(nS)$?
 - Are the $Y(nS)$ produced differently with multiplicity?

$-\eta$	Y	$+\eta$
HF [-5.2, -4]	Y [-1.93, 1.93]	HF [4, 5.2]
N_{tracks} [-2.4, 2.4]		

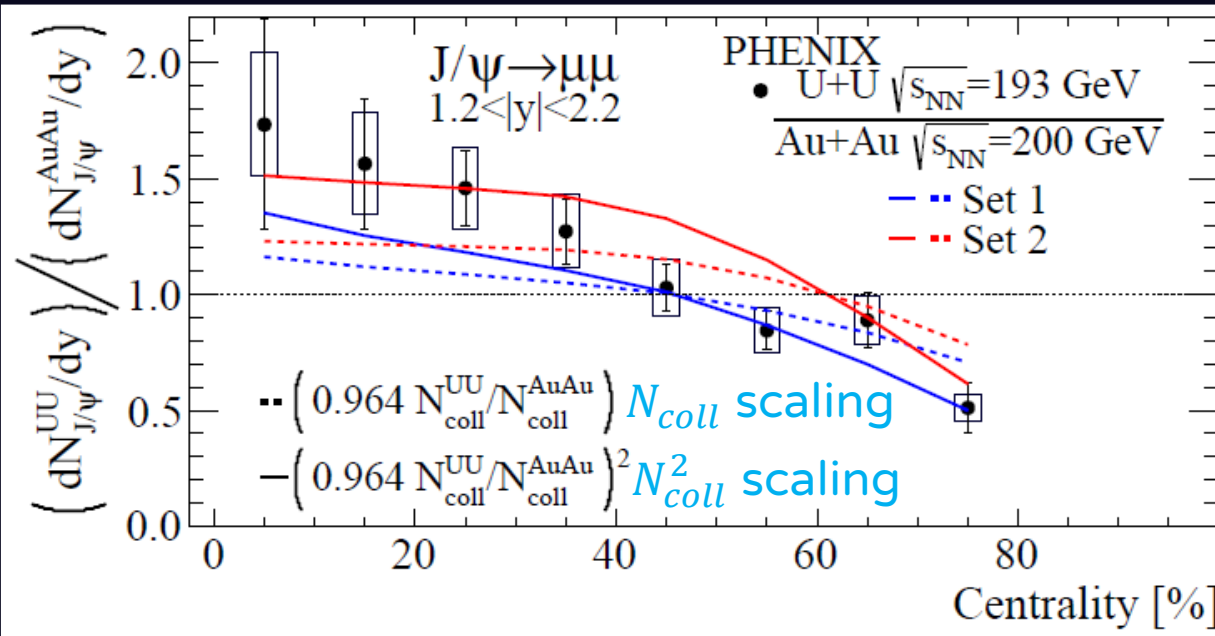
Y compared to theory



J/ψ production at RHIC

52

→ (re)combination/suppression role investigated comparing U-U and AuAu



in central U-U collisions:

1) stronger suppression due to color screening

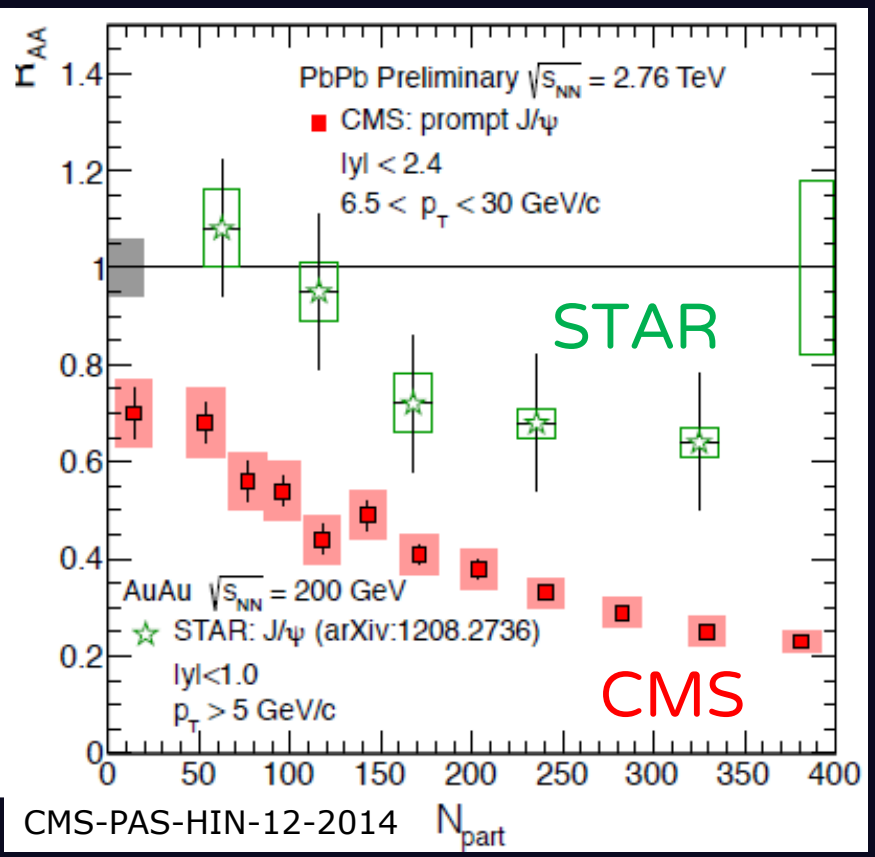
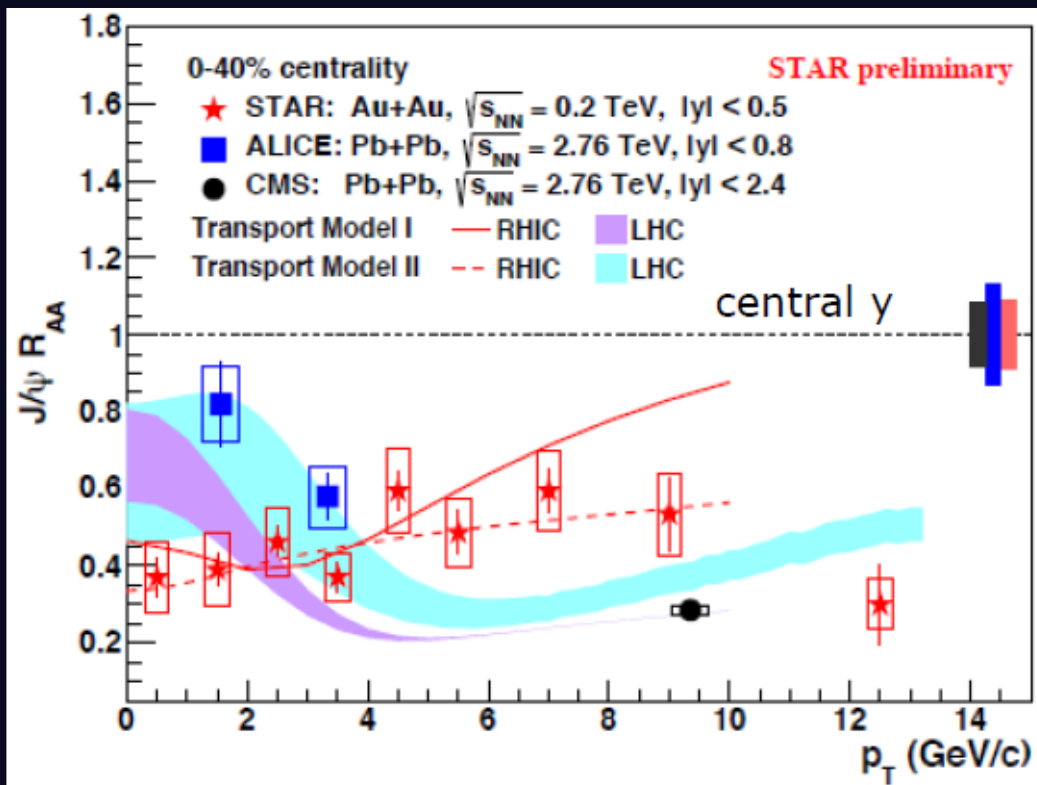
$$\varepsilon_{AuAu} \sim 80-85\% \varepsilon_{UU}$$

2) J/ψ recombination favoured by 25% larger N_{coll} in UU

$$N_{J/\psi}^{stat} \sim N_c^2 \sim N_{coll}^2$$

→ results slightly favour N_{coll}^2 scaling → dominant (re)combination over suppression when going from central U-U to Au-Au collisions

→ quantitative comparison depends on the choice of the uranium Woods-Saxon parametrizations



$\psi(2S)$ production in pA

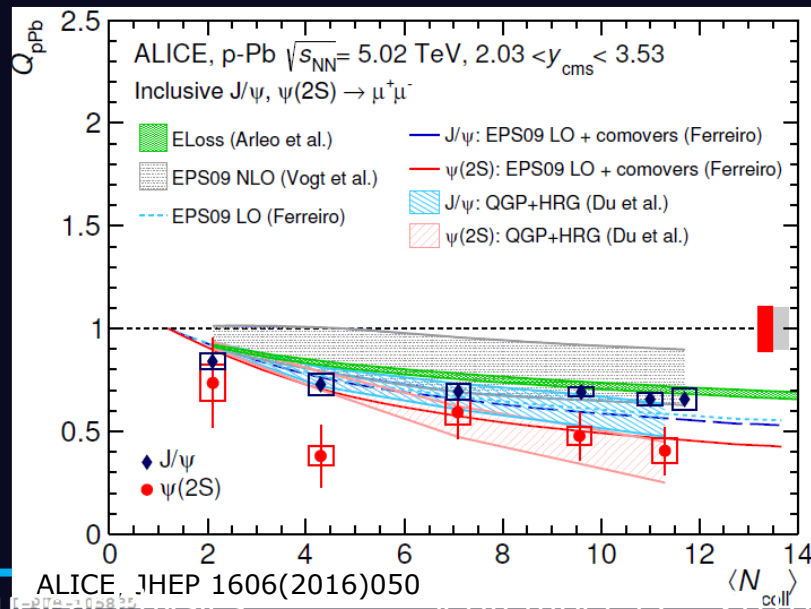
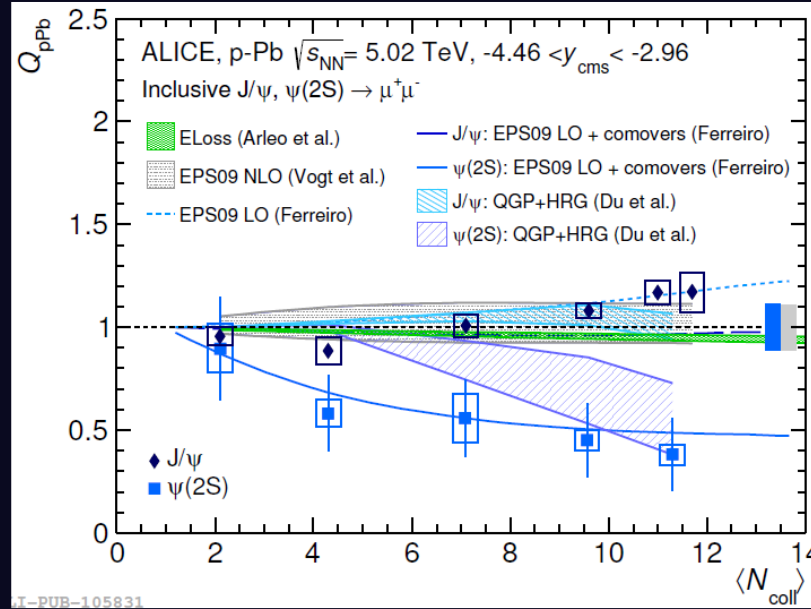
54

➔ $\psi(2S)$ suppression is stronger than the J/ψ one, both at RHIC and LHC

➔ unexpected since time spent by the $c\bar{c}$ in the nucleus (τ_c) is shorter than charmonium formation time (τ_f)

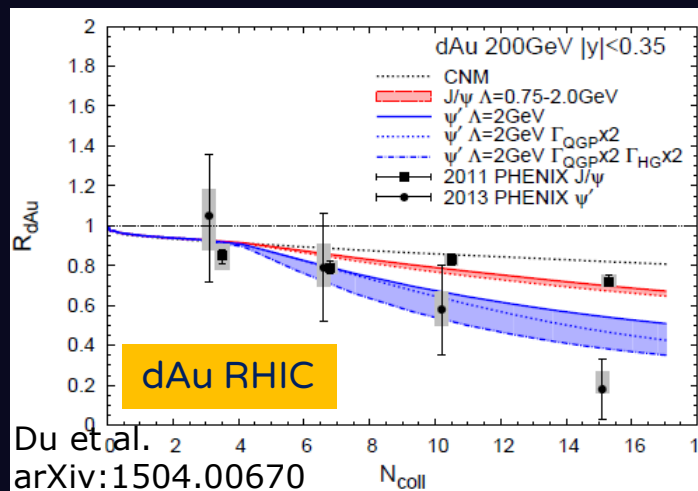
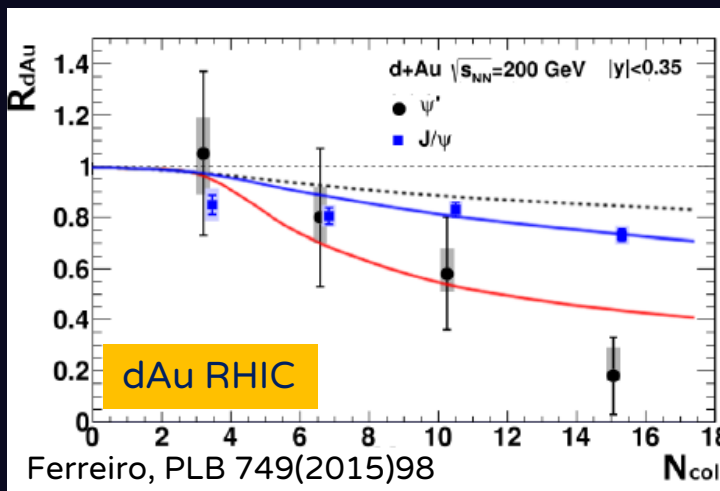
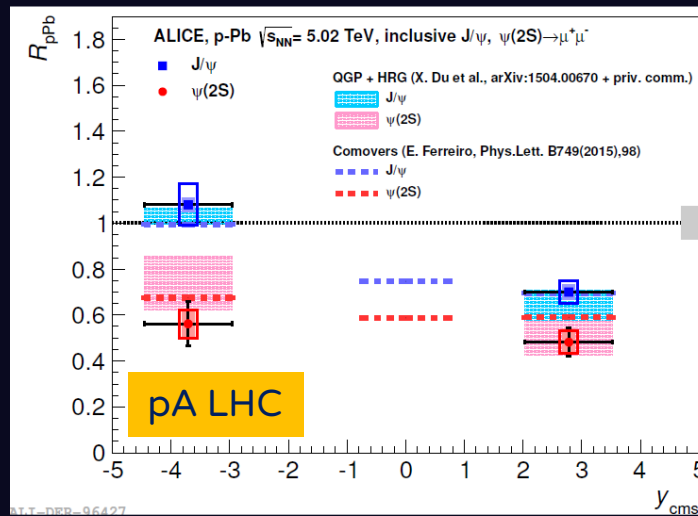
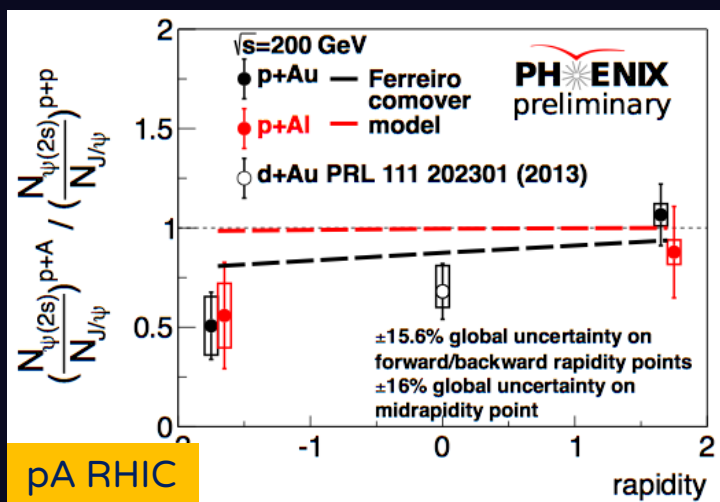
➔ shadowing and energy loss, almost identical for J/ψ and $\psi(2S)$, do not account for the different suppression

➔ Only models including QGP + hadron resonance gas or comovers describe the stronger $\psi(2S)$ suppression



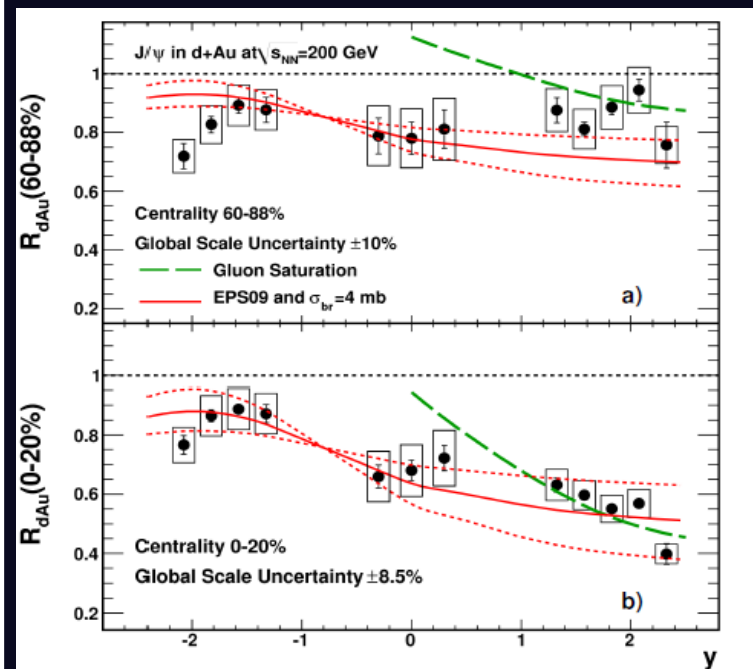
Comparison to theoretical models 55

➔ QGP+hadron resonance gas (Rapp) or comovers models (Ferreiro) reasonably describe both J/ψ and $\psi(2S)$ suppression at RHIC and LHC



CNM effects at RHIC

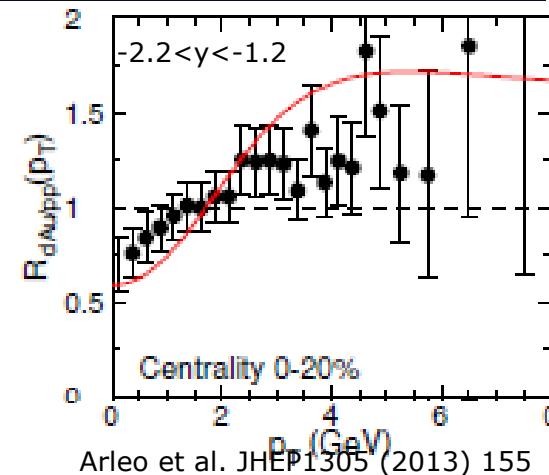
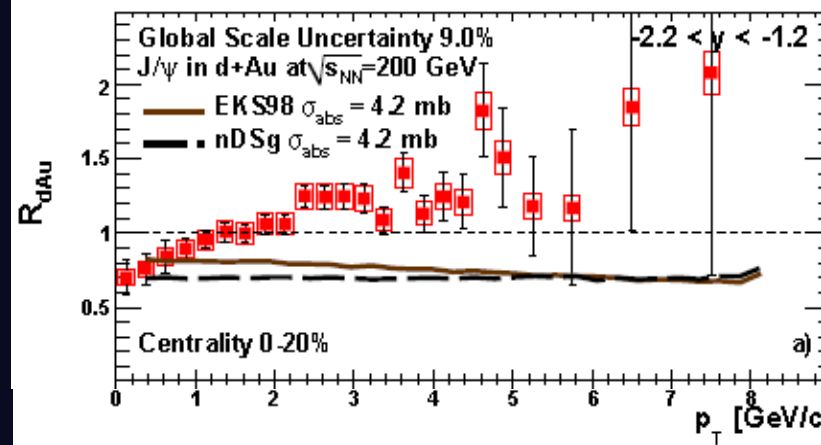
56

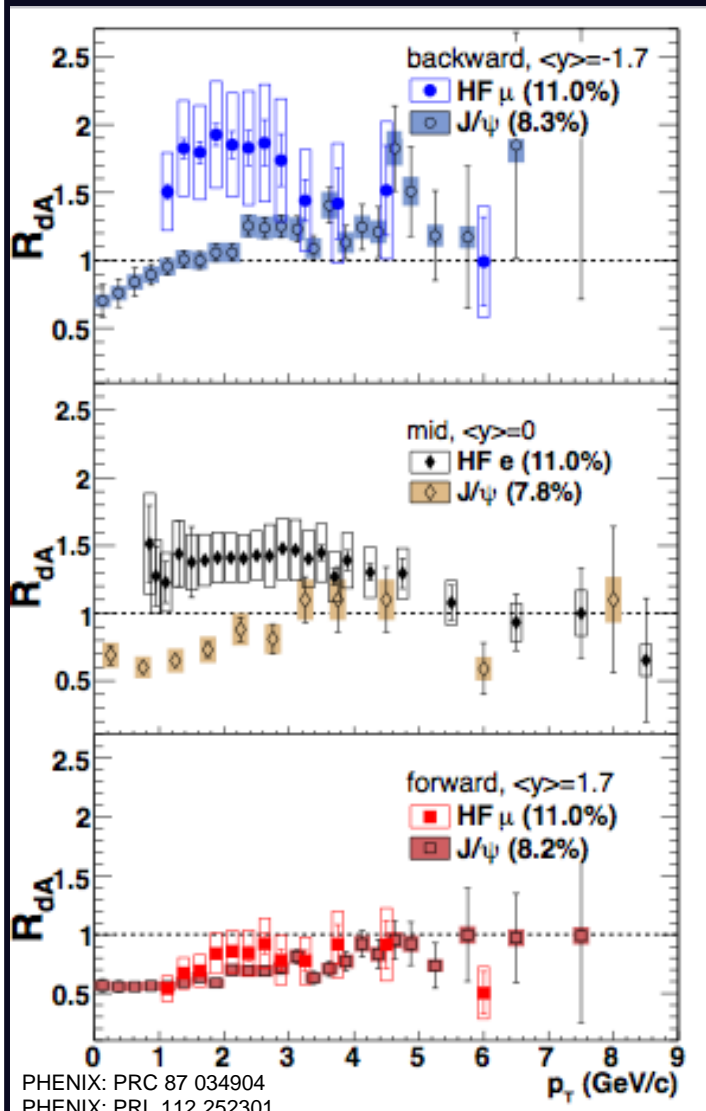


➔ Disentangling CNM mechanisms is challenging

➔ shadowing + cc break-up describe R_{dAu} vs y , but meets some difficulties for R_{dAu} vs p_T

➔ coherent energy loss contribution induces a less flat R_{dAu} dependence on p_T





➔ Comparison between heavy flavor and quarkonium:

R_{dA} of HF muon and J/ ψ are consistent at forward rapidity, but clearly different at backward rapidity

➔ charm production is enhanced but J/ ψ production is significantly suppressed due to nuclear breakup inside dense comovers at backward rapidity

➔ Contrarily to LHC, at RHIC energies a contribution from J/ ψ breakup in nuclear matter could be present ($\sigma_{J/\psi-N} \sim 4\text{mb}$)