

MASSIVE NEUTRINOS CIRCA 2017

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(ICREA U. Barcelona & YITP Stony Brook)

INTERNATIONAL SCHOOL OF NUCLEAR PHYSICS

Neutrinos in Cosmology, Astro-, Particle- and Nuclear Physics

Erice-Sicily: September 16-24, 2017

OUTLINE

The confirmed picture: 3ν Lepton Flavour Parameters

A partial list of Q&A

Neutrinos in the Standard Model

The SM is a gauge theory based on the symmetry group

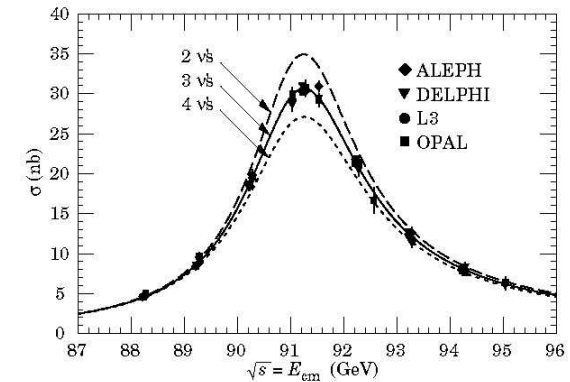
$$SU(3)_C \times SU(2)_L \times U(1)_Y \Rightarrow SU(3)_C \times U(1)_{EM}$$

With three generation of fermions

$(1, 2)_{-\frac{1}{2}}$	$(3, 2)_{\frac{1}{6}}$	$(1, 1)_{-1}$	$(3, 1)_{\frac{2}{3}}$	$(3, 1)_{-\frac{1}{3}}$
$\begin{pmatrix} \nu_e \\ e \end{pmatrix}_L$	$\begin{pmatrix} u^i \\ d^i \end{pmatrix}_L$	e_R	u^i_R	d^i_R
$\begin{pmatrix} \nu_\mu \\ \mu \end{pmatrix}_L$	$\begin{pmatrix} c^i \\ s^i \end{pmatrix}_L$	μ_R	c^i_R	s^i_R
$\begin{pmatrix} \nu_\tau \\ \tau \end{pmatrix}_L$	$\begin{pmatrix} t^i \\ b^i \end{pmatrix}_L$	τ_R	t^i_R	b^i_R

There is no ν_R

Three and only three



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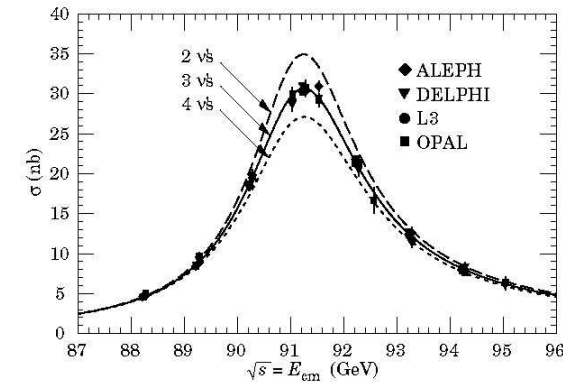


Accidental global symmetry: $B \times L_e \times L_\mu \times L_\tau$ (hence $L = L_e + L_\mu + L_\tau$)



ν strictly massless

Three and only three



- By 2017 we have observed with high (or good) precision:
 - * Atmospheric ν_μ & $\bar{\nu}_\mu$ disappear most likely to ν_τ (**SK, MINOS, ICECUBE**)
 - * Accel. ν_μ & $\bar{\nu}_\mu$ disappear at $L \sim 300/800$ Km (**K2K, T2K, MINOS, NO ν A**)
 - * Some accelerator ν_μ appear as ν_e at $L \sim 300/800$ Km (**T2K, MINOS, NO ν A**)
 - * Solar ν_e convert to ν_μ/ν_τ (**Cl, Ga, SK, SNO, Borexino**)
 - * Reactor $\bar{\nu}_e$ disappear at $L \sim 200$ Km (**KamLAND**)
 - * Reactor $\bar{\nu}_e$ disappear at $L \sim 1$ Km (**D-Chooz, Daya Bay, Reno**)

All this implies that L_α are violated

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- The *starting* path:

Precise determination of the low energy parametrization

The New Minimal Standard Model

- Minimal Extension to allow for LFV \Rightarrow give Mass to the Neutrino

* Introduce ν_R AND impose L conservation \Rightarrow Dirac $\nu \neq \nu^c$:

$$\mathcal{L} = \mathcal{L}_{SM} - M_\nu \overline{\nu}_L \nu_R + h.c.$$

* NOT impose L conservation \Rightarrow Majorana $\nu = \nu^c$

$$\mathcal{L} = \mathcal{L}_{SM} - \frac{1}{2} M_\nu \overline{\nu}_L \nu_L^C + h.c.$$

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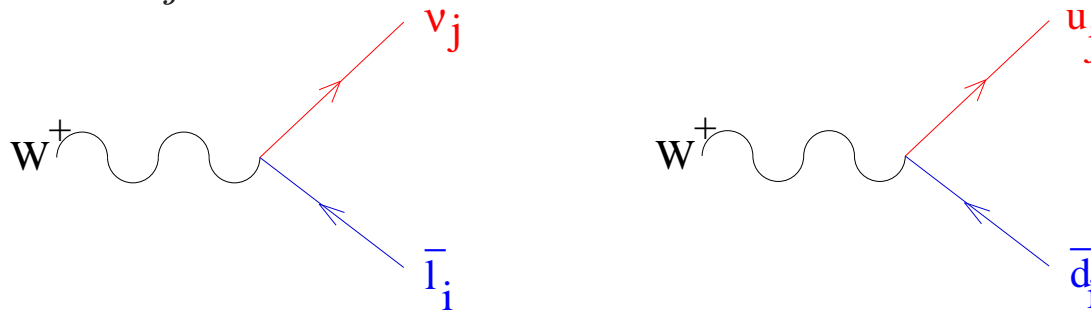
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- The charged current interactions of leptons are not diagonal (same as quarks)

$$\frac{g}{\sqrt{2}} W_\mu^+ \sum_{ij} (U_{LEP}^{ij} \bar{\ell}^i \gamma^\mu L \nu^j + U_{CKM}^{ij} \bar{U}^i \gamma^\mu L D^j) + h.c.$$



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- In general for $N = 3 + s$ massive neutrinos U_{LEP} is $3 \times N$ matrix

$$U_{\text{LEP}} U_{\text{LEP}}^\dagger = I_{3 \times 3} \quad \text{but in general} \quad U_{\text{LEP}}^\dagger U_{\text{LEP}} \neq I_{N \times N}$$

- U_{LEP} : $3 + 3s$ angles + $2s + 1$ Dirac phases + $s + 2$ Majorana phases

- If neutrinos have mass, a weak eigenstate $|\nu_\alpha\rangle$ produced in $l_\alpha + N \rightarrow \nu_\alpha + N'$

is a linear combination of the mass eigenstates ($|\nu_i\rangle$): $|\nu_\alpha\rangle = \sum_{i=1}^n U_{\alpha i} |\nu_i\rangle$

- After a distance L it can be detected with flavour β with probability

$$P_{\alpha\beta} = \delta_{\alpha\beta} - 4 \sum_{j \neq i}^n \text{Re}[U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*] \sin^2 \left(\frac{\Delta_{ij}}{2} \right) + 2 \sum_{j \neq i} \text{Im}[U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*] \sin(\Delta_{ij})$$

$$\frac{\Delta_{ij}}{2} = \frac{(E_i - E_j)L}{2} = 1.27 \frac{(m_i^2 - m_j^2)}{\text{eV}^2} \frac{L/E}{\text{Km/GeV}}$$

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- When osc between 2- ν dominates:

$$P_{\alpha\alpha} = 1 - P_{osc} \quad \text{Disappear}$$

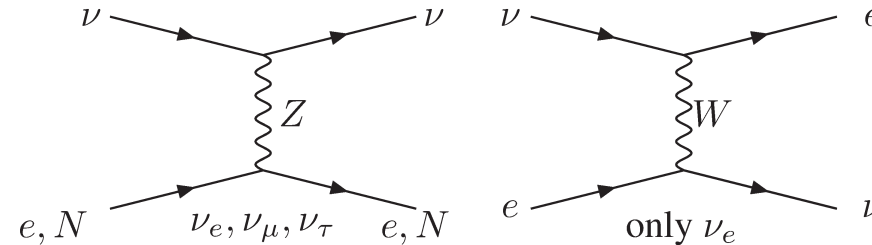
$$P_{osc} = \sin^2(2\theta) \sin^2 \left(1.27 \frac{\Delta m^2 L}{E} \right) \quad \text{Appear}$$

\Rightarrow No info on sign of Δm^2 and θ octant

Matter Effects

- If ν cross **matter** regions (Sun, Earth...) it interacts *coherently*

– But **Different flavours**
have **different interactions** :



\Rightarrow Effective potential in ν evolution : $V_e \neq V_{\mu,\tau} \Rightarrow \Delta V^\nu = -\Delta V^{\bar{\nu}} = \sqrt{2}G_F N_e$

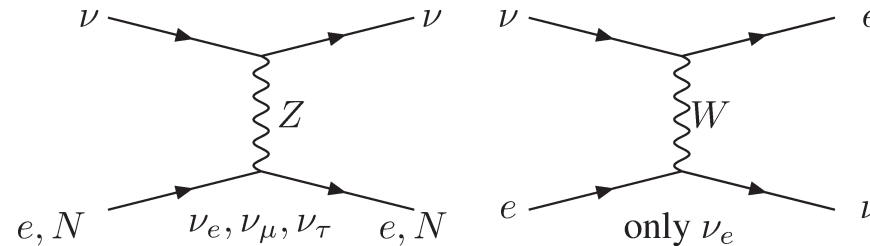
$$-i \frac{\partial}{\partial x} \begin{pmatrix} \nu_e \\ \nu_X \end{pmatrix} = \left[- \begin{pmatrix} V_e - V_X - \frac{\Delta m^2}{4E} \cos 2\theta & \frac{\Delta m^2}{4E} \sin 2\theta \\ \frac{\Delta m^2}{4E} \sin 2\theta & \frac{\Delta m^2}{4E} \cos 2\theta \end{pmatrix} \right] \begin{pmatrix} \nu_e \\ \nu_X \end{pmatrix}$$

\Rightarrow **Modification of mixing angle and oscillation wavelength (MSW)**

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⇒ **Modification of mixing angle and oscillation wavelength** (MSW)

- Mass difference and mixing in matter:

$$\Delta m_m^2 = \sqrt{(\Delta m^2 \cos 2\theta - 2E\Delta V)^2 + (\Delta m^2 \sin 2\theta)^2}$$

$$\sin(2\theta_m) = \frac{\Delta m^2 \sin(2\theta)}{\Delta m_{mat}^2}$$

⇒ For solar ν 's in adiabatic regime

$$P_{ee} = \frac{1}{2} [1 + \cos(2\theta_m) \cos(2\theta)]$$

Dependence on θ octant

⇒ In LBL terrestrial experiments

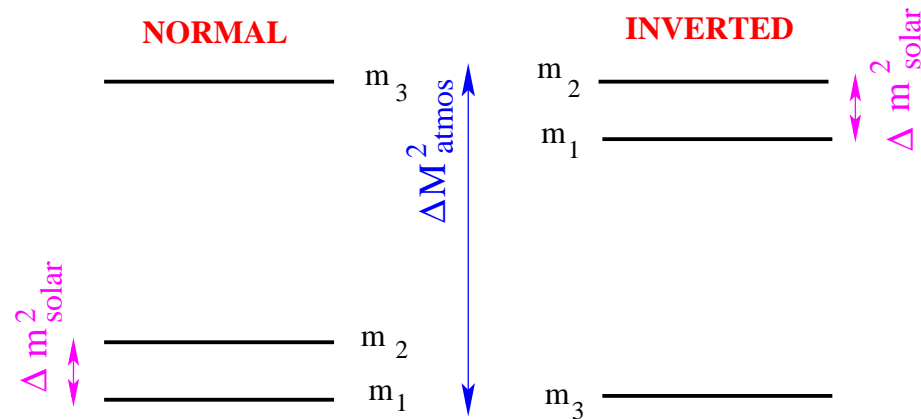
Dependence on **sign of Δm^2**

and θ octant

- For for 3 ν 's : 3 Mixing angles + 1 Dirac Phase + 2 Majorana Phases

$$U_{\text{LEP}} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{i\delta_{\text{CP}}} \\ 0 & 1 & 0 \\ -s_{13}e^{-i\delta_{\text{CP}}} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{21} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} e^{i\eta_1} & 0 & 0 \\ 0 & e^{i\eta_2} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

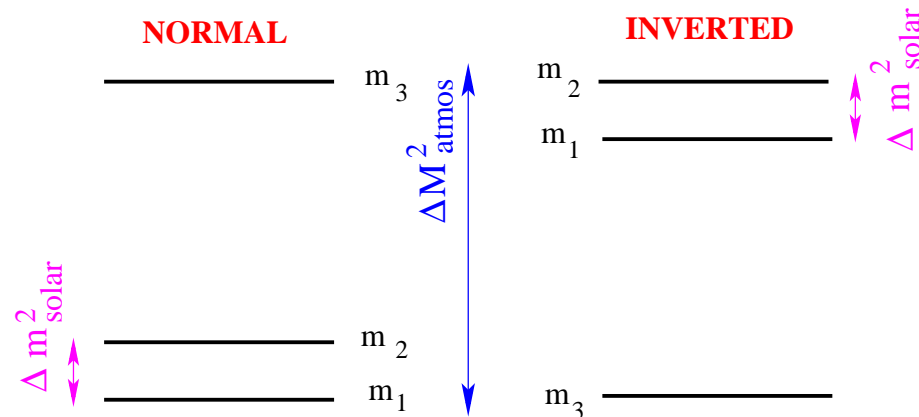
- Two Possible Orderings



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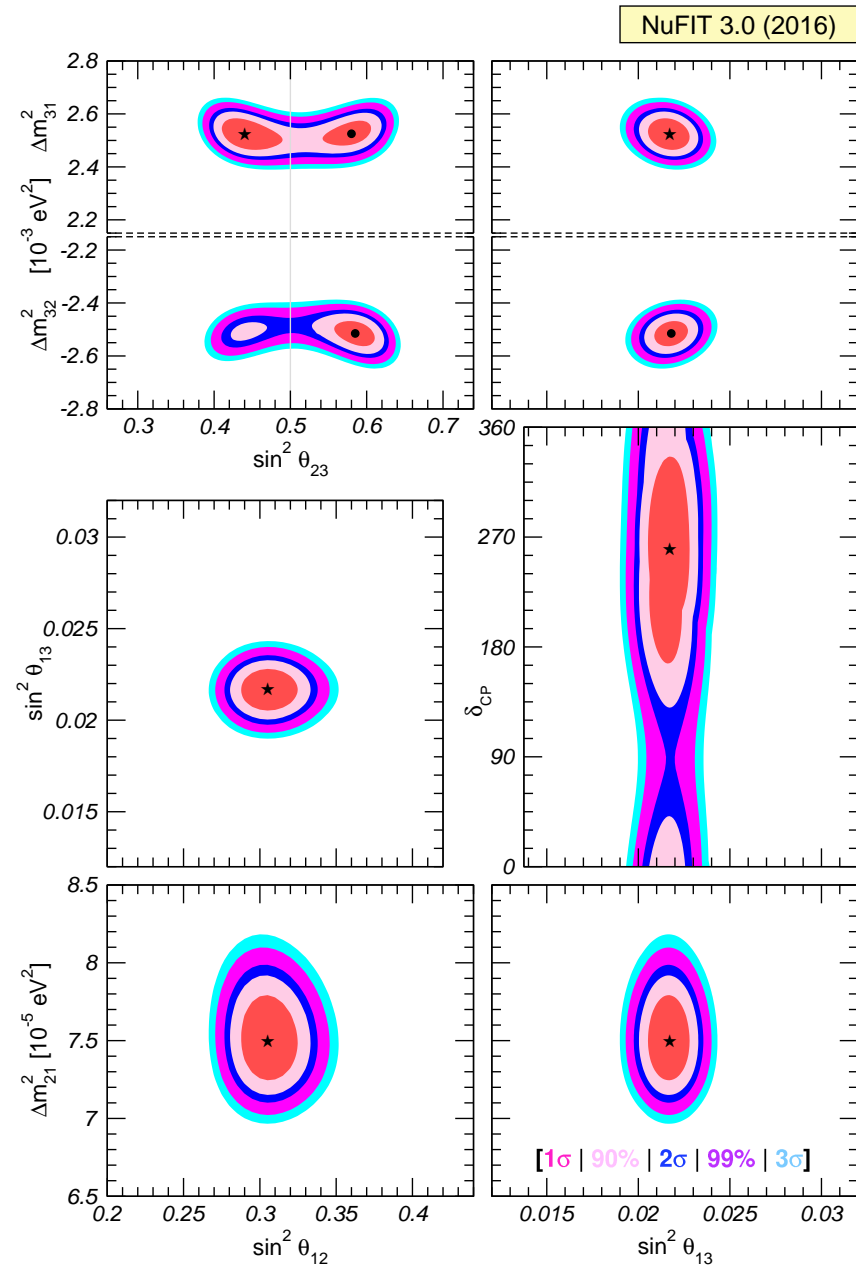
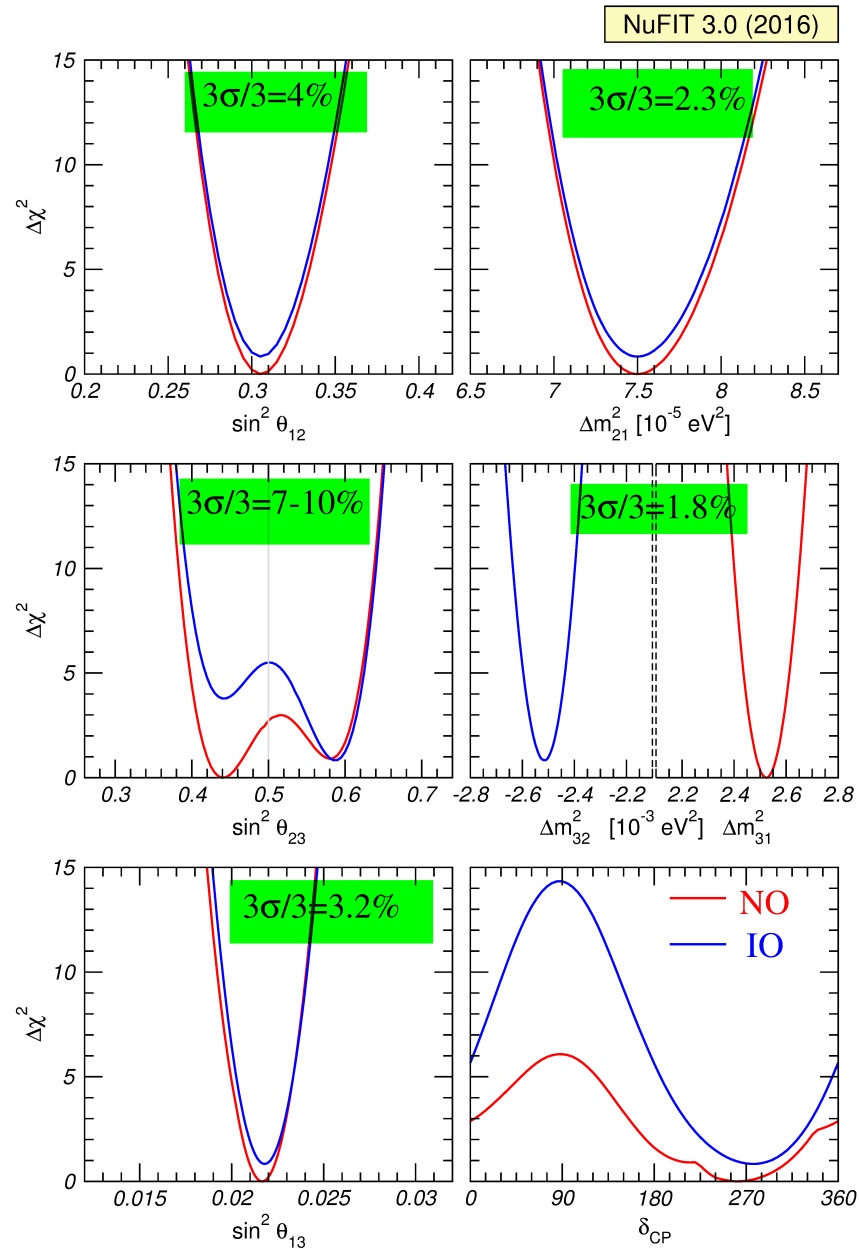
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Experiment	Dominant Dependence	Important Dependence
Solar Experiments	$\rightarrow \theta_{12}$	$\Delta m_{21}^2, \theta_{13}$
Reactor LBL (KamLAND)	$\rightarrow \Delta m_{21}^2$	θ_{12}, θ_{13}
Reactor MBL (Daya Bay, Reno, D-Chooz)	$\rightarrow \theta_{13}$	Δm_{atm}^2
Atmospheric Experiments	$\rightarrow \theta_{23}$	$\Delta m_{\text{atm}}^2, \theta_{13}, \delta_{\text{CP}}$
Acc LBL ν_μ Disapp (Minos, T2K, NOvA)	$\rightarrow \Delta m_{\text{atm}}^2$	θ_{23}
Acc LBL ν_e App (Minos, T2K, NOvA)	$\rightarrow \theta_{13}$	$\delta_{\text{CP}}, \theta_{23}$

Global 6-parameter fit <http://www.nu-fit.org>

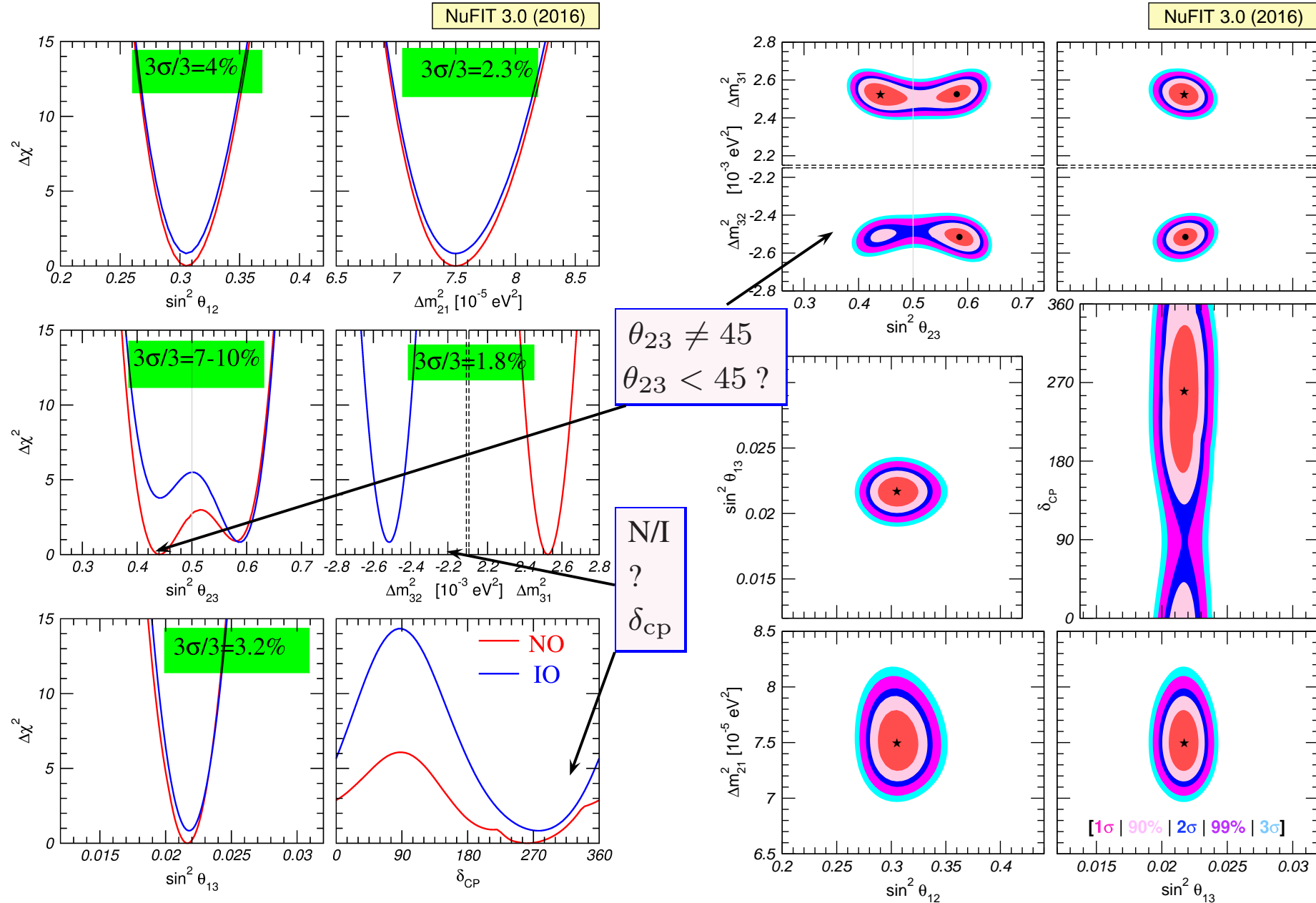
Esteban, Maltoni, Martinez-Soler, Schwetz, MCG-G ArXiv:1611:01514



3 ν Flavour Parameters: Status June 2017

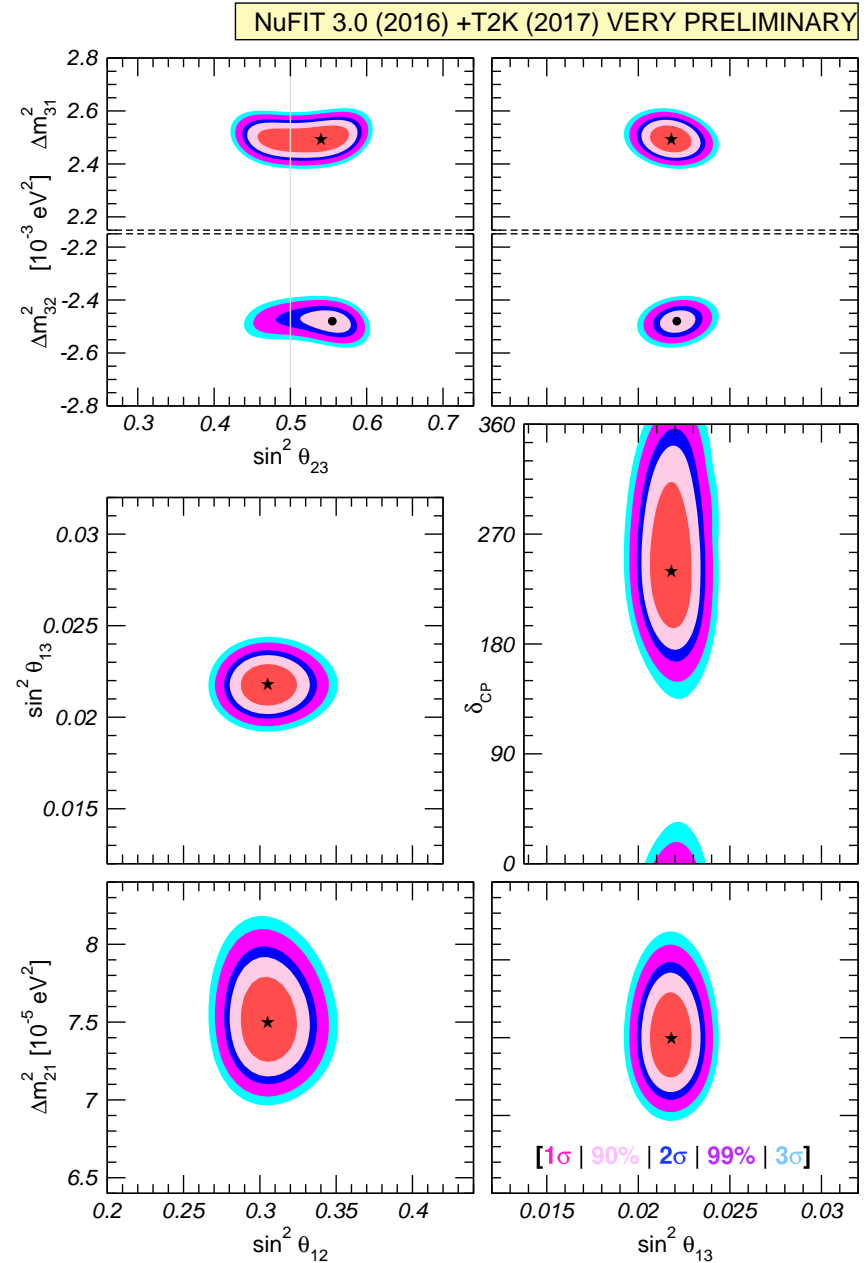
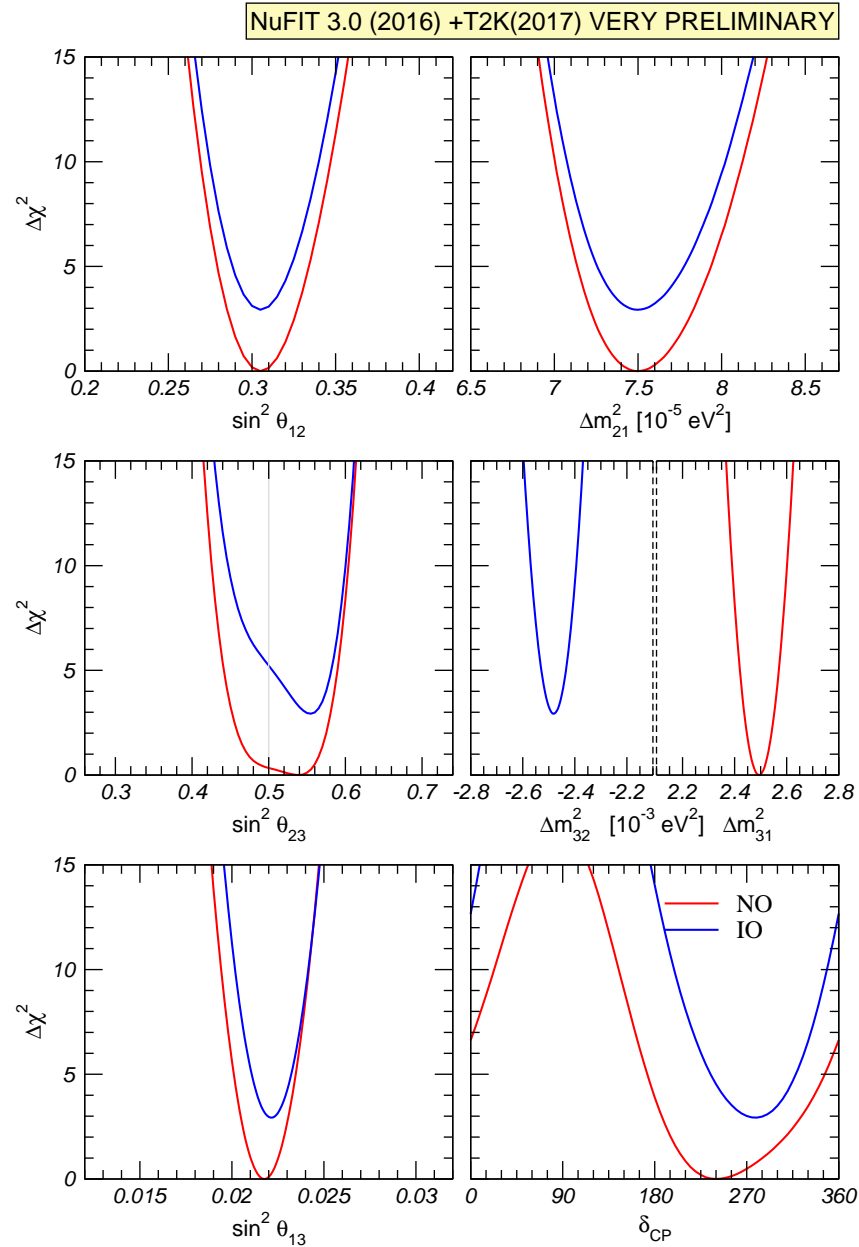
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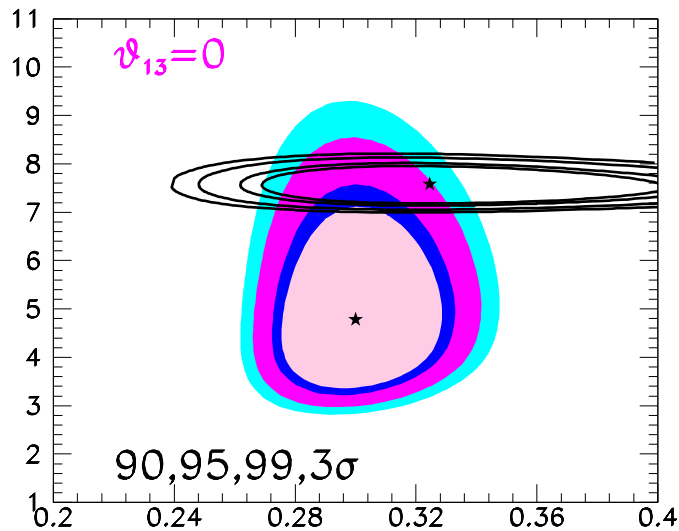
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Esteban, Maltoni, Martinez-Soler, Schwetz, MCG-G effect OF T2K (2017) VERY PRELIMINARY



3 ν Analysis: “12” Sector and θ_{13}

- For $\theta_{13} = 0$



- When θ_{13} increases

$$P_{ee} \simeq \begin{cases} \text{Solar High E} : c_{13}^4 \sin^2 2\theta_{12} \\ \text{Solar Low E} : c_{13}^4 (1 - \sin^2 2\theta_{12}/2) \\ \text{Kam} : c_{13}^4 \left(1 - \sin^2 2\theta_{12} \sin^2 \frac{\Delta m_{21}^2 L}{4E}\right) \end{cases}$$

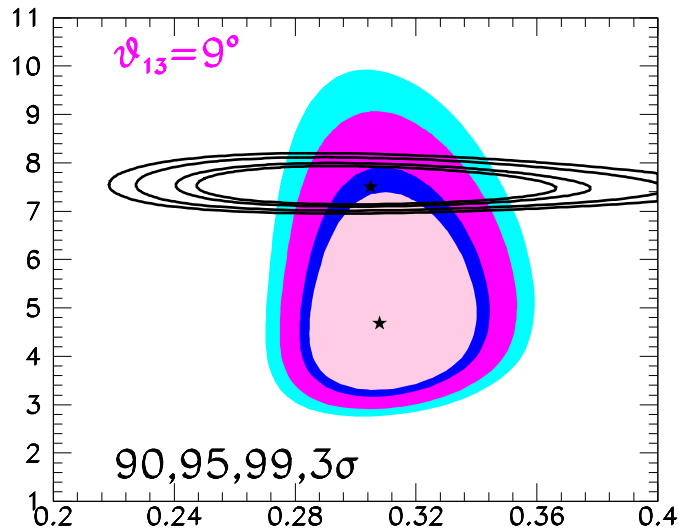
\Rightarrow KamLAND region shifts left

\Rightarrow Solar slight shifts right (due to High E)

$$\sin^2 \theta_{12} = \begin{cases} 0.3 \text{ From Solar} \\ 0.325 \text{ From KLAND} \end{cases}$$

3 ν Analysis: “12” Sector and θ_{13}

- For $\theta_{13} \simeq 9^\circ$



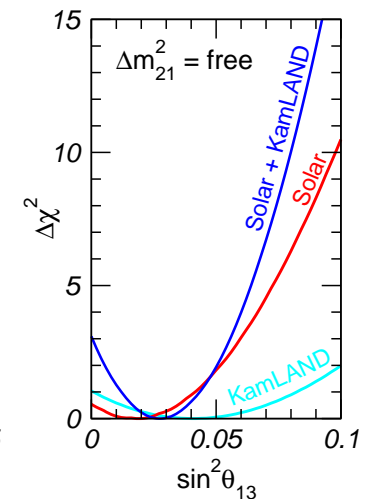
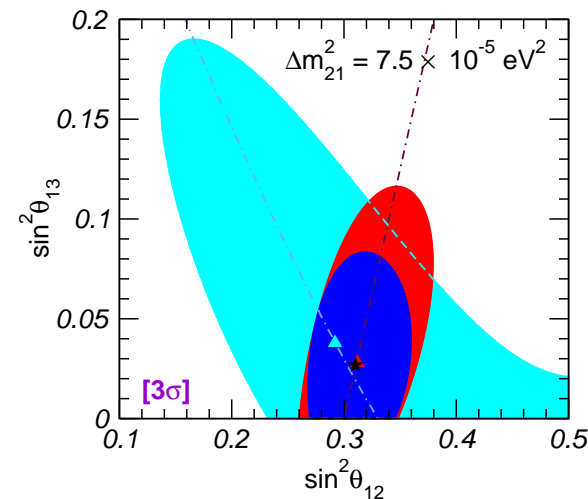
- \Rightarrow Good match of best fit θ_{12}
- \Rightarrow Residual tension on Δm_{21}^2

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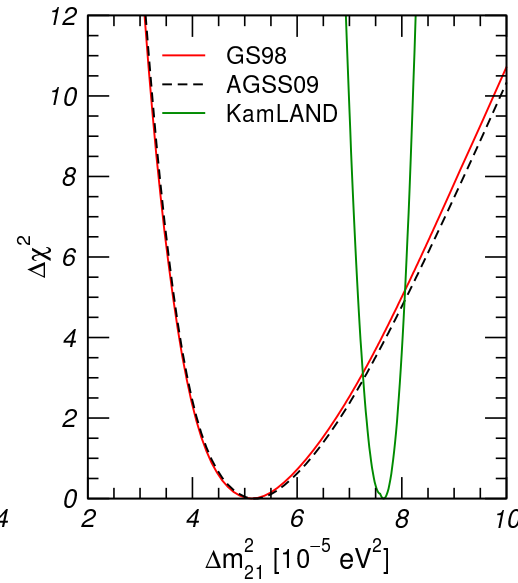
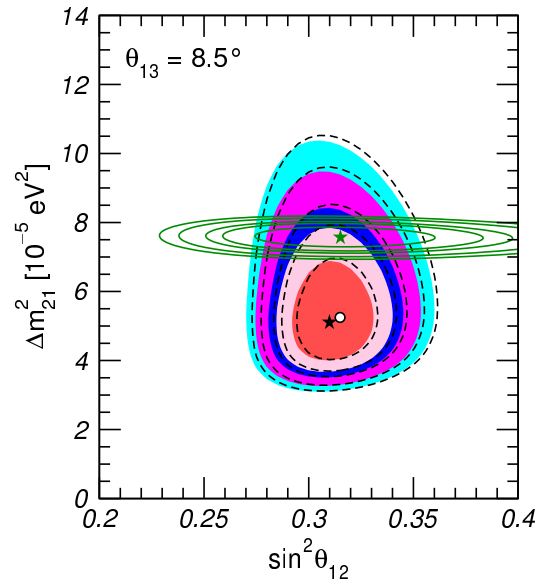
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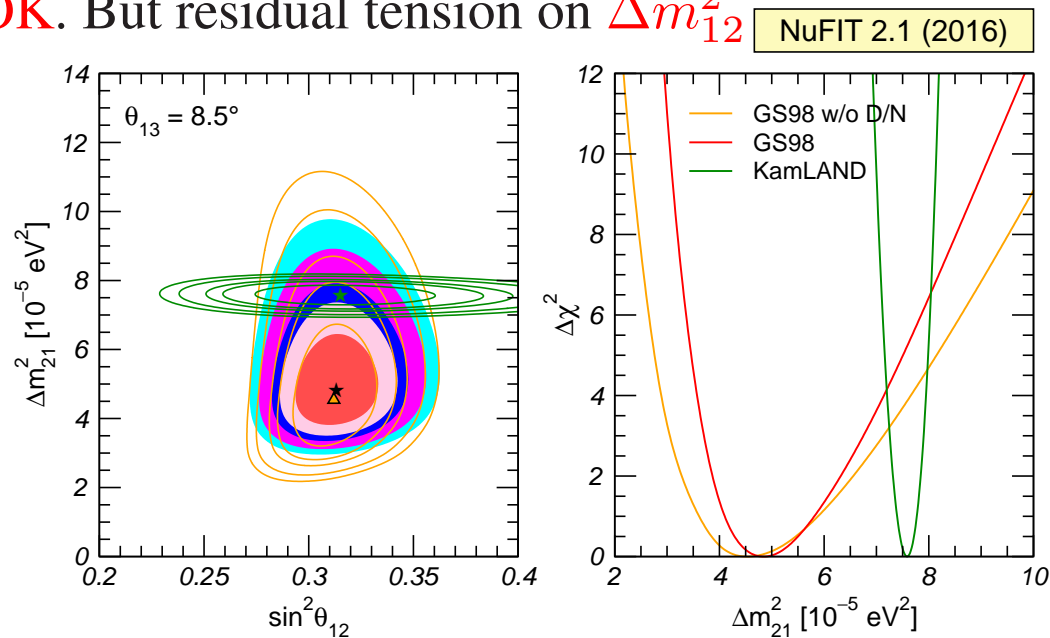
For $\theta_{13} \simeq 9^\circ$ θ_{12} OK. But residual tension on Δm_{12}^2

NuFIT 3.0 (2016)

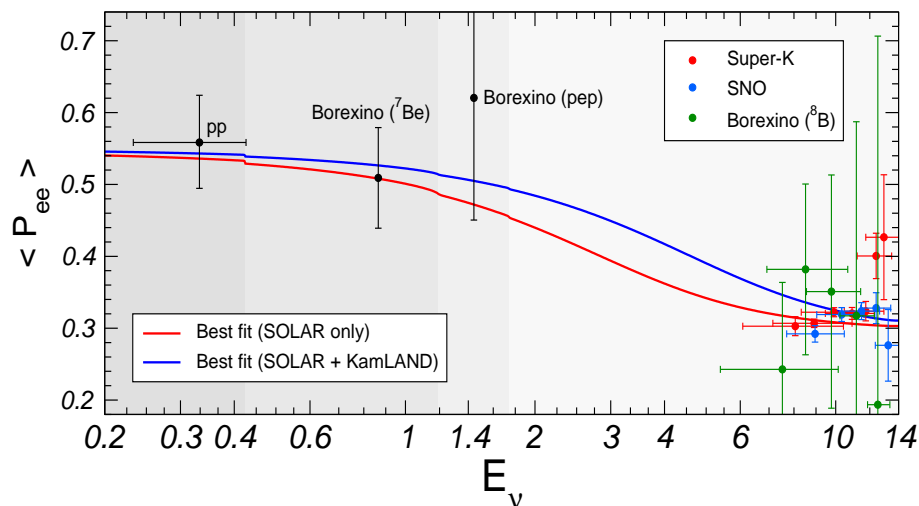


3 ν Analysis: Δm_{21}^2 KamLAND vs SOLAR

For $\theta_{13} \simeq 9^\circ$ θ_{12} OK. But residual tension on Δm_{12}^2



Tension related to: a) “too large” of Day/Night at SK



b) smaller-than-expected
low-E turn up from MSW
at best global fit

Modified matter potential? More latter ...

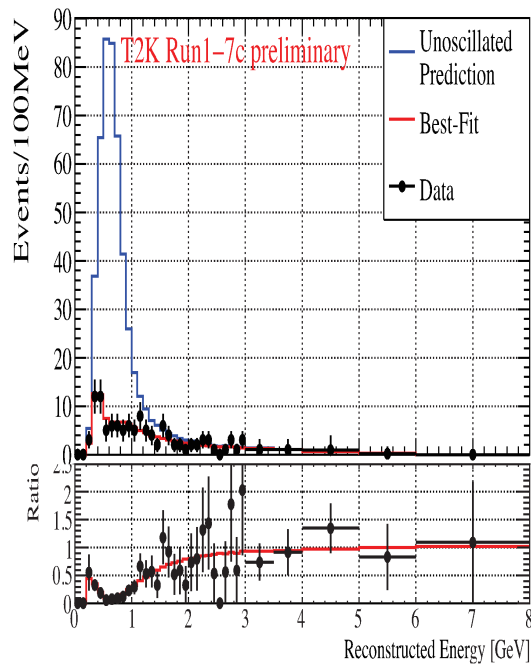
3 ν Analysis: θ_{23}

- Best determined in ν_μ and $\bar{\nu}_\mu$ disappearance in LBL

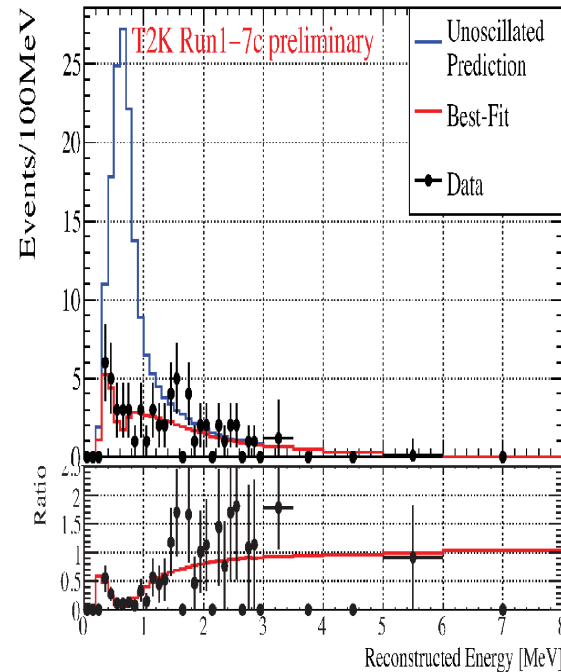
$$P_{\mu\mu} \simeq 1 - (c_{13}^4 \sin^2 2\theta_{23} + s_{23}^2 \sin^2 2\theta_{13}) \sin^2 \left(\frac{\Delta m_{31}^2 L}{4E} \right) + \mathcal{O}(\Delta m_{21}^2)$$

- At osc maximum $\sin^2 \left(\frac{\Delta m_{31}^2 L}{4E} \right) = 1 \Rightarrow P_{\mu\mu} \simeq 0$ for $\theta_{23} \simeq \frac{\pi}{4}$

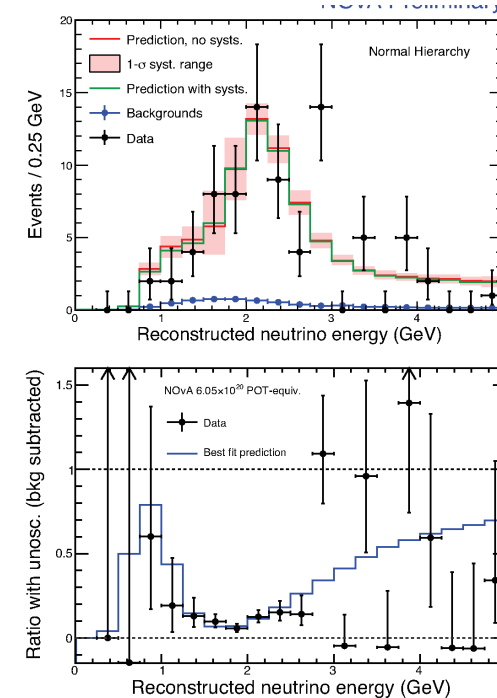
T2K $\nu_\mu \rightarrow \nu_\mu$



T2K $\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu$



NOvA $\nu_\mu \rightarrow \nu_\mu$

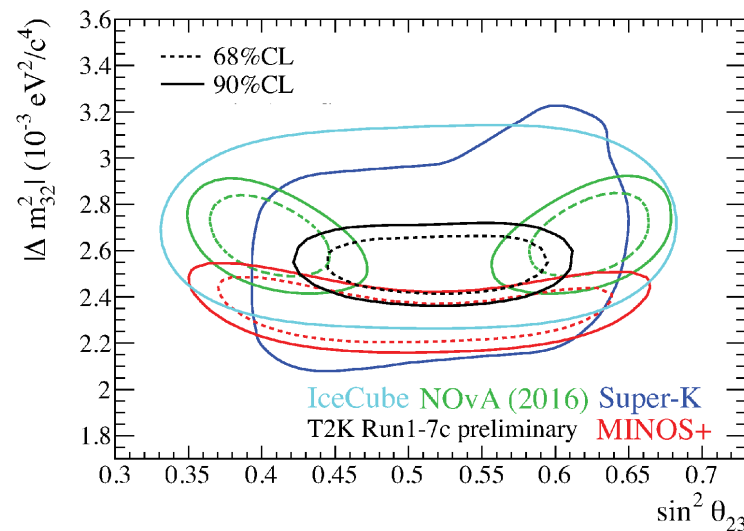


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- Allowed regions by the different experiments:



In making this figure θ_{13} is constrained by prior from reactor data

Caution: Not the same using θ_{13} reactor prior than combining with reactor results (because of Δm_{32}^2 in reactors)

Δm_{23}^2 in LBL vs Reactors

- At LBL determined in ν_μ and $\bar{\nu}_\mu$ disappearance spectrum

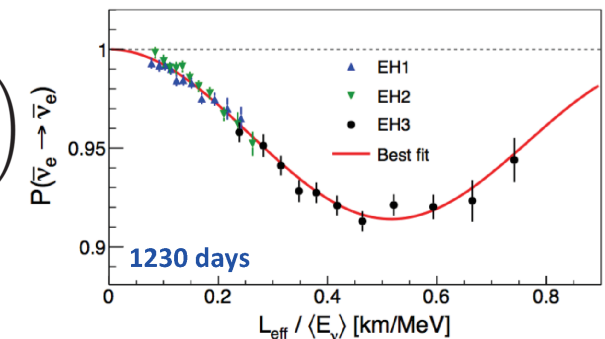
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- At MBL Reactors (Daya-Bay, Reno, D-Chooz) determined in $\bar{\nu}_e$ disapp spectrum

$$P_{ee} \simeq 1 - \sin^2 2\theta_{13} \sin^2 \left(\frac{\Delta m_{ee}^2 L}{4E} \right) - c_{13}^4 \sin^2 2\theta_{12} \sin^2 \left(\frac{\Delta m_{21}^2 L}{4E} \right)$$

$$\Delta m_{ee}^2 \simeq |\Delta m_{32}^2| \pm c_{12}^2 \Delta m_{21}^2 \simeq |\Delta m_{32}^2| \pm 0.05 \times 10^{-3} \text{ eV}^2$$

Nunokawa, Parke, Zukanovich (2005)

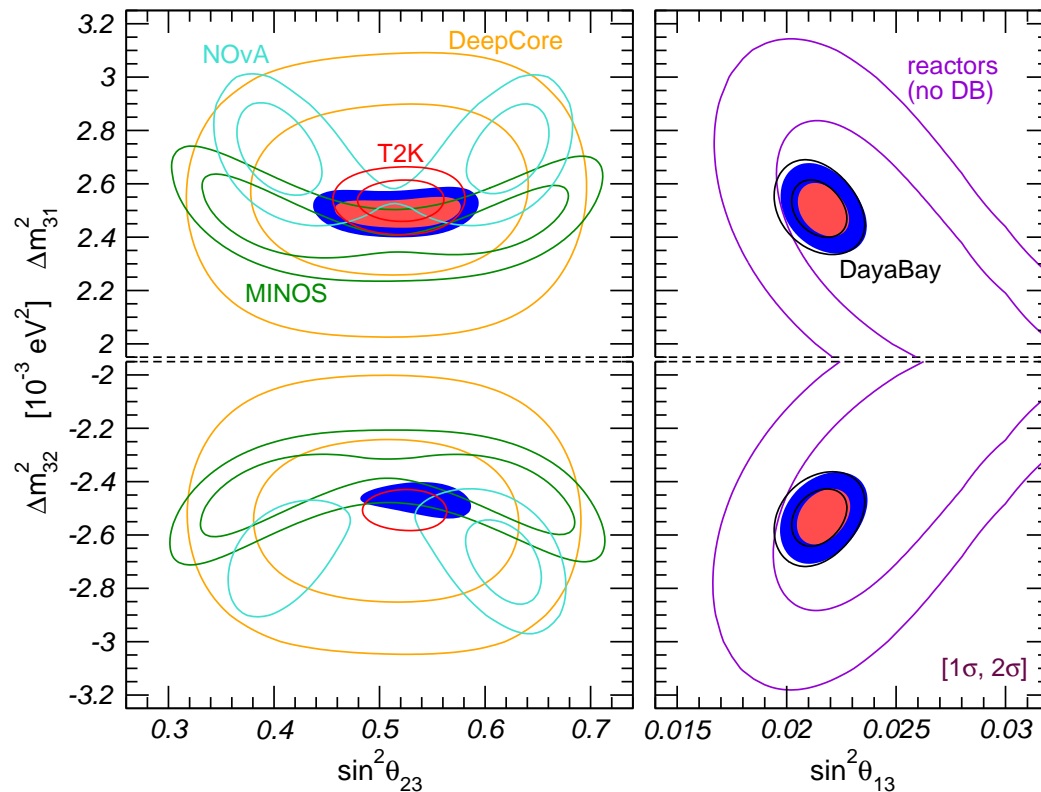


Δm_{23}^2 in LBL vs Reactors: Consistency

- At LBL determined in ν_μ and $\bar{\nu}_\mu$ disappearance spectrum
- At MBL Reactors (Daya-Bay, Reno, D-Chooz) determined in $\bar{\nu}_e$ disapp spectrum

LBL ν_μ disappearanceREAC ν_e disappearance

Sept (2017) PRELIM



- Consistent values of $|\Delta m_{32}^2|$
- Hint for non-maximal θ_{23} driven by NO ν A and MINOS
- T2K (2017) slight fav $\theta_{23} > 45$

Leptonic CP Violation

- Leptonic CP $\Rightarrow P_{\nu_\alpha \rightarrow \nu_\beta} \neq P_{\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta}$:

$$P_{\nu_\alpha \rightarrow \nu_\beta} - P_{\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta} \propto J \quad \text{with} \quad J = \text{Im}(U_{\alpha 1} U_{\alpha 2}^* U_{\beta 2} U_{\beta 1}^*) = J_{\text{LEP,CP}}^{\text{max}} \sin \delta_{\text{CP}}$$

$$J_{\text{LEP,CP}}^{\text{max}} = \frac{1}{8} c_{13} \sin^2 2\theta_{13} \sin^2 2\theta_{23} \sin^2 2\theta_{12}$$

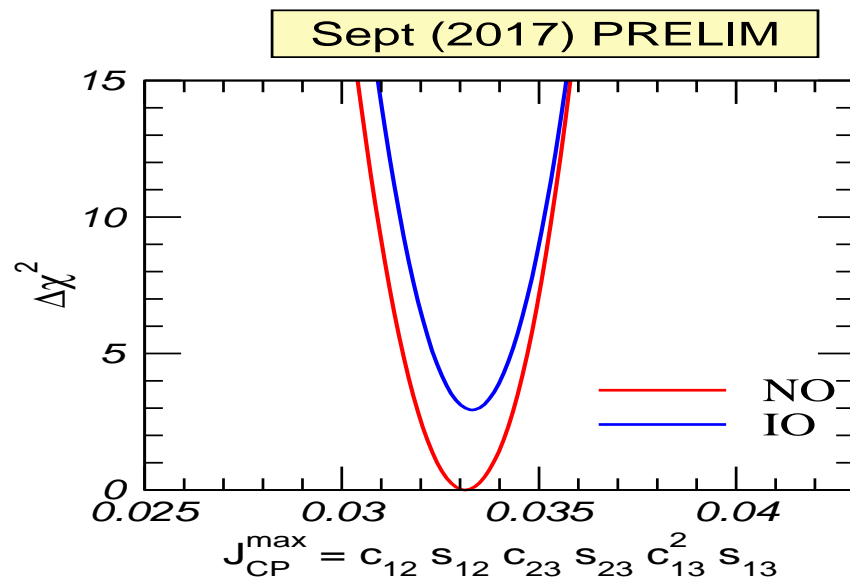
Leptonic CP Violation

- Leptonic CP $\Rightarrow P_{\nu_\alpha \rightarrow \nu_\beta} \neq P_{\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta}$:

$$P_{\nu_\alpha \rightarrow \nu_\beta} - P_{\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta} \propto J \quad \text{with} \quad J = \text{Im}(U_{\alpha 1} U_{\alpha 2}^* U_{\beta 2} U_{\beta 1}^*) = J_{\text{LEP,CP}}^{\text{max}} \sin \delta_{\text{CP}}$$

$$J_{\text{LEP,CP}}^{\text{max}} = \frac{1}{8} c_{13} \sin^2 2\theta_{13} \sin^2 2\theta_{23} \sin^2 2\theta_{12}$$

- Maximum Allowed Leptonic CPV:



$$J_{\text{LEP,CP}}^{\text{max}} = (3.29 \pm 0.07) \times 10^{-2}$$

to compare with

$$J_{\text{CKM,CP}} = (3.04 \pm 0.21) \times 10^{-5}$$

\Rightarrow Leptonic CPV may be largest CPV
in New Minimal SM

if $\sin \delta_{\text{CP}}$ not too small

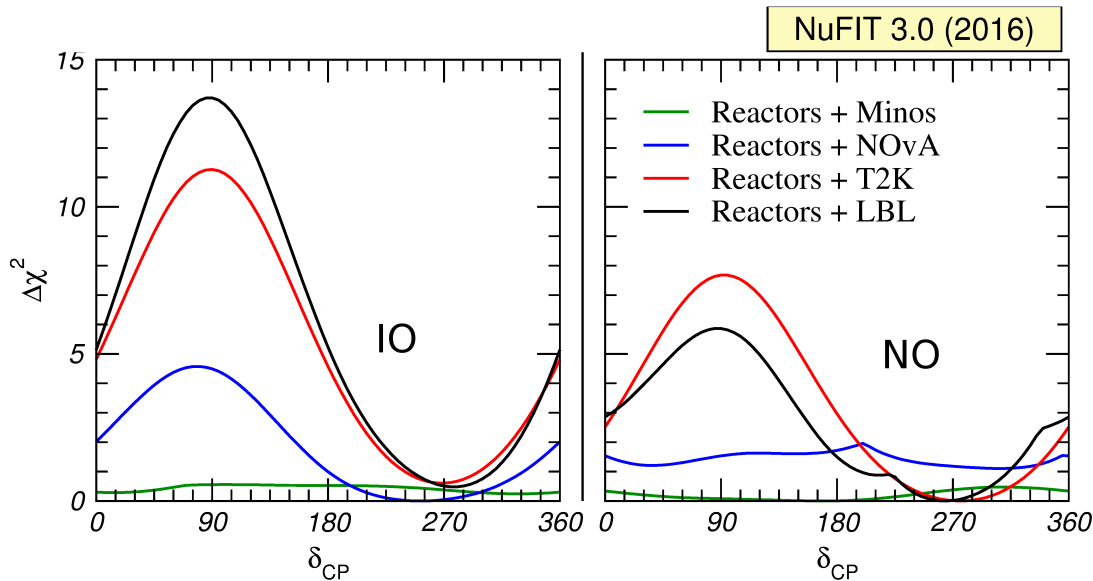
Leptonic CP Phase

- Leptonic CPV Phase: Mainly from $\nu_\mu \rightarrow \nu_e$ in LBL (complicated by **matter effects**)

$$P_{\mu e} \simeq s_{23}^2 \sin^2 2\theta_{13} \left(\frac{\Delta_{31}}{B_\mp} \right)^2 \sin^2 \left(\frac{B_\mp L}{2} \right) + 8 J_{\text{LEP,CP}}^{\text{max}} \frac{\Delta_{12}}{V_E} \frac{\Delta_{31}}{B_\mp} \sin \left(\frac{V_E L}{2} \right) \sin \left(\frac{B_\mp L}{2} \right) \cos \left(\frac{\Delta_{31} L}{2} \pm \delta_{CP} \right)$$

$$\Delta_{ij} = \frac{\Delta m_{ij}^2}{2E} \quad B_\pm = \Delta_{31} \pm V_E \quad J_{\text{LEP,CP}}^{\text{max}} = \frac{1}{8} c_{13} \sin^2 2\theta_{13} \sin^2 2\theta_{23} \sin^2 2\theta_{12}$$

Before T2K (2017)



- Best fit $\delta_{CP} \sim 270^\circ$
- CP conserv at 70% (NO), 97% (IO)
- Driven by “fluctuation” in T2K

Mass hierarchy	ν_e		$\bar{\nu}_e$	
	Normal	Inverted	Normal	Inverted
$\delta_{CP} = -\pi/2$	28.8	25.5	6.0	6.5
$\delta_{CP} = 0$	24.2	21.2	6.9	7.4
$\delta_{CP} = \pi/2$	19.7	17.2	7.7	8.4
$\delta_{CP} = \pm\pi$	24.2	21.6	6.8	7.4
Data	32		4	

\Rightarrow One concluded :
Significance may not grow soon

Leptonic CP Phase: T2K 2017

Accumulated 14.7×10^{20} protons-on-target (POT) in neutrino mode and 7.6×10^{20} POT in antineutrino mode - full data set presented here

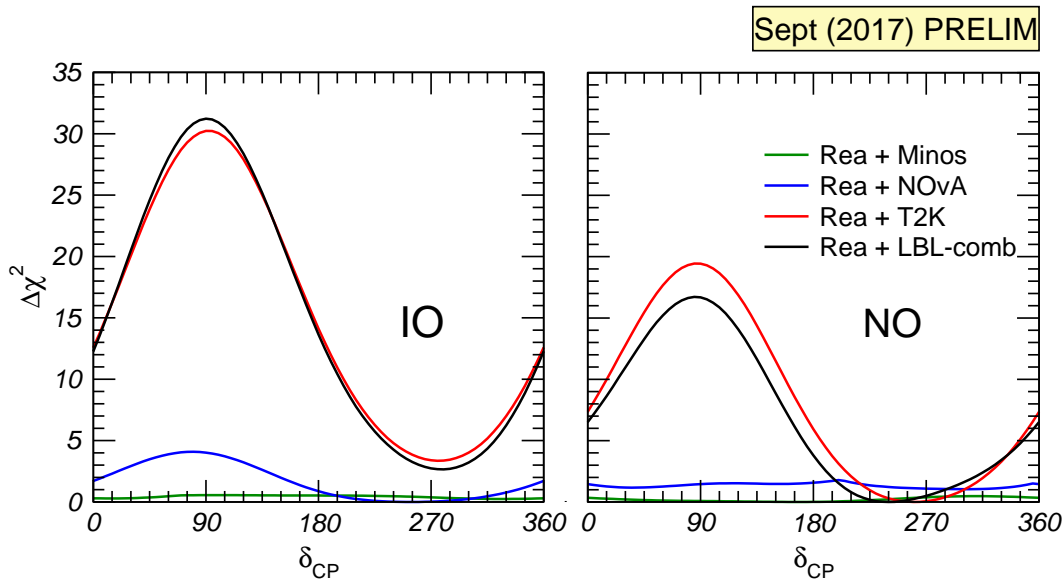
- 29% of the approved T2K POT

		Predicted Rates				Observed
Sample		$\delta_{cp}=-\pi/2$	$\delta_{cp}=0$	$\delta_{cp}=\pi/2$	$\delta_{cp}=\pi$	Rates
ν_e	CCQE 1-Ring e-like FHC	73.5	61.5	49.9	62.0	74
ν_e	CC1 π 1-Ring e-like FHC	6.92	6.01	4.87	5.78	15
$\bar{\nu}_e$	CCQE 1-Ring e-like RHC	7.93	9.04	10.04	8.93	7
ν_μ	CCQE 1-Ring μ -like FHC	267.8	267.4	267.7	268.2	240
$\bar{\nu}_\mu$	CCQE 1-Ring μ -like RHC	63.1	62.9	63.1	63.1	68

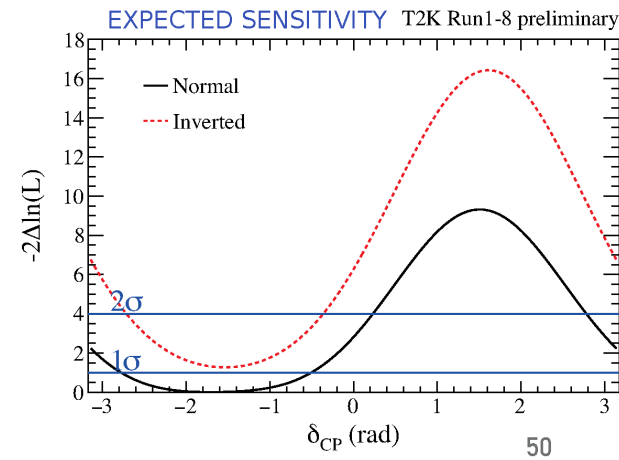
M. Hartz, KeK colloquim, August 2017

Leptonic CP Phase

Including T2K (2017) PRELIMINARY



- Best fit at $\delta_{CP} \sim 240^\circ$
- CP conserv at 95% (NO)
- $15^\circ - 160^\circ$ disfavoured at $\Delta\chi^2 > 9$
- Still more than expected sensitiv in T2K
- $\chi_{\min,IO}^2 - \chi_{\min,NO}^2 \simeq 3$



- Leptonic CP $\Rightarrow P_{\nu_\alpha \rightarrow \nu_\beta} \neq P_{\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta}$:

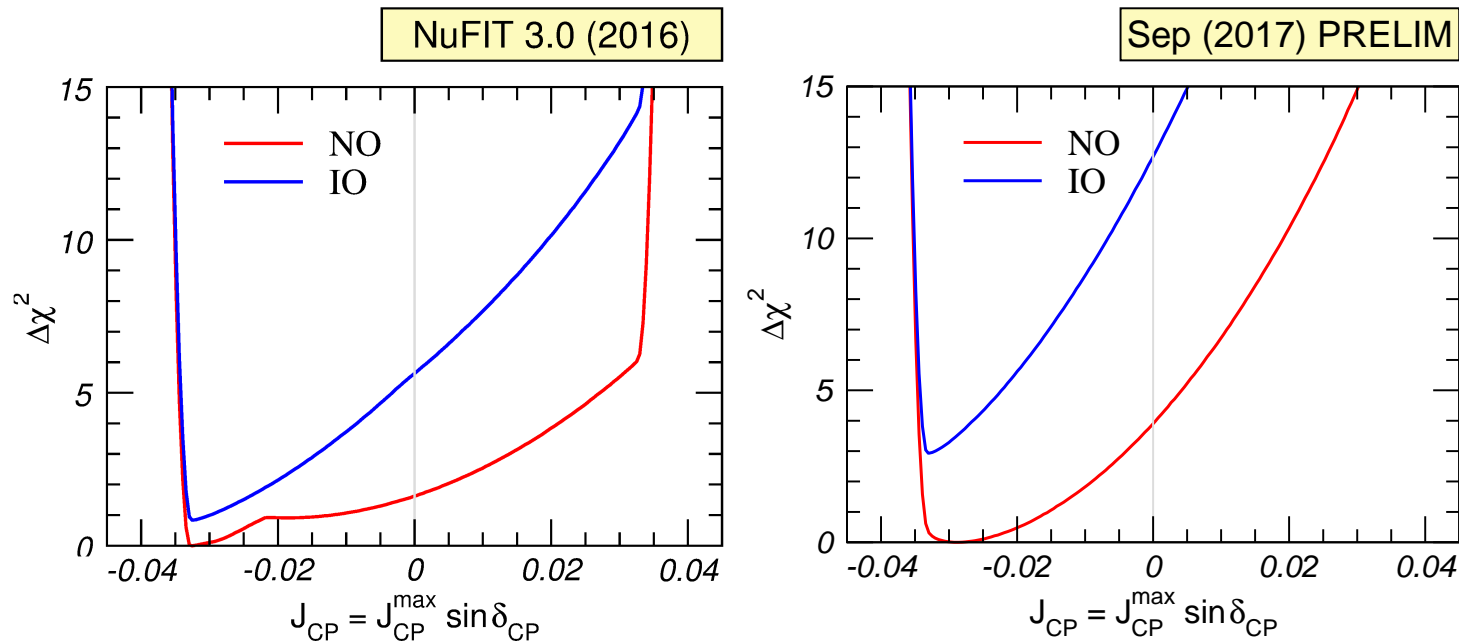
$$P_{\nu_\alpha \rightarrow \nu_\beta} - P_{\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta} \propto J \quad \text{with} \quad J = \text{Im}(U_{\alpha 1} U_{\alpha 2}^* U_{\beta 2} U_{\beta 1}^*) = J_{\text{LEP,CP}}^{\text{max}} \sin \delta_{\text{CP}}$$

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- Leptonic Jarlskog Invariant : Best fit $J_{\text{LEP,CP}} = -0.030$



Confirmed Low Energy Picture and MY List of Q&A

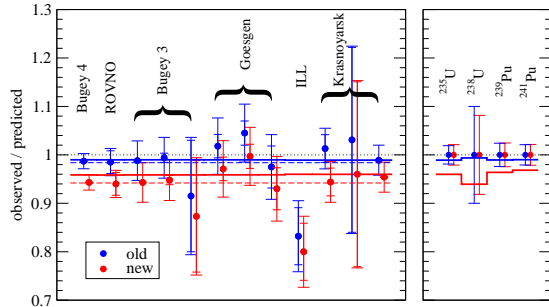
- At least **two** neutrinos **are massive** \Rightarrow **There is NP**
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- **Ordering**: No significant preference yet
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- **Oscillations DO NOT determine the** lightest mass
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Anxiously waiting for Katrin
- **Dirac or Majorana?**: We do not know, *anxiously* waiting for ν -less $\beta\beta$ decay
- **Only three light states?**

Beyond 3ν's: Light Sterile Neutrinos

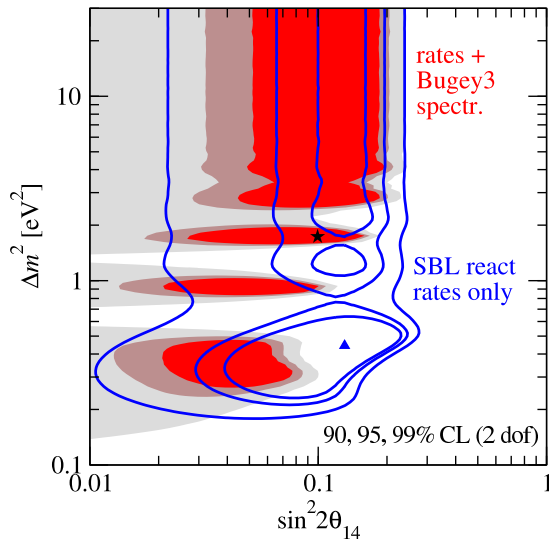
- Several Observations which can be Interpreted as Oscillations with $\Delta m^2 \sim eV^2$

Reactor Anomaly

New reactor flux calculation
 \Rightarrow Deficit in data at $L \lesssim 100$ m



Explained as ν_e disappearance

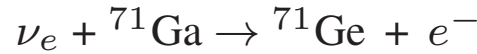


Kopp etal, ArXiv 1303.3011

Gallium Anomaly

Acero, Giunti, Laveder, 0711.4222
 Giunti, Laveder, 1006.3244

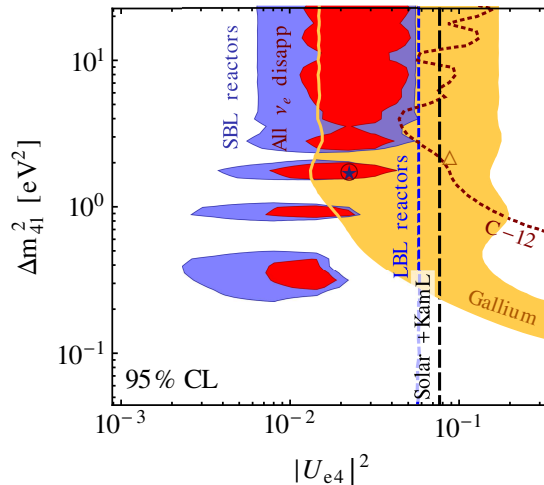
Radioactive Sources (^{51}Cr , ^{37}Ar)
 in calibration of Ga Solar Exp;



Give a rate lower than expected

$$R = \frac{N_{\text{obs}}}{N_{\text{Bahc}}^{\text{th}}} = 0.86 \pm 0.05 \quad (2.8\sigma)$$

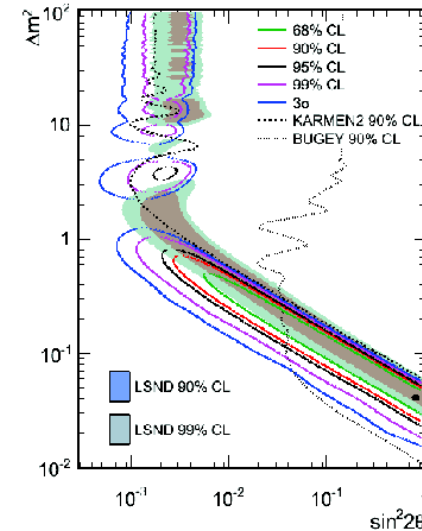
Explained as ν_e disappearance



Kopp etal, ArXiv 1303.3011

LSND, MiniBoone

$\nu_\mu \rightarrow \nu_e$ and $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$



- These explanations require $3+N_s$ mass eigenstates $\rightarrow N_s$ sterile neutrinos

$\nu_e \rightarrow \nu_e$ **disapp** (REACT, Gallium, Solar, LSND/KARMEN)

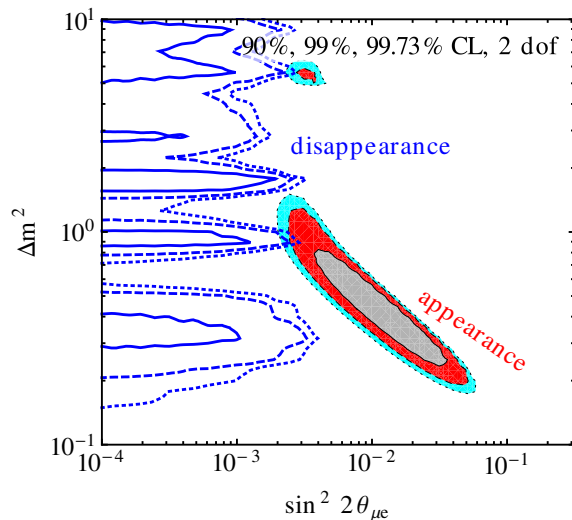
- Problem: fit together $\nu_\mu \rightarrow \nu_e$ **app** (LSND, KARMEN, NOMAD, MiniBooNE, E776, ICARUS)

$\nu_\mu \rightarrow \nu_\mu$ **disapp** (CDHS, ATM, MINOS, ICECUBE)

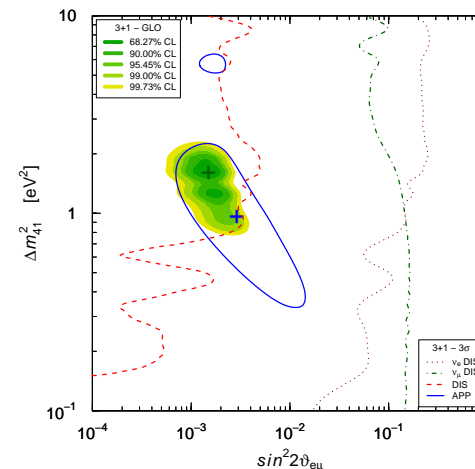
- Generically: $P(\nu_e \rightarrow \nu_\mu) \sim |U_{ei}^* U_{\mu i}|$ [i = heavier state(s)]

But $|U_{ei}|$ constrained by $P(\nu_e \rightarrow \nu_e)$ disappearance data
 And $|U_{\mu i}|$ constrained by $P(\nu_\mu \rightarrow \nu_\mu)$ disappearance data } \Rightarrow **Severe tension**

Kopp et al, ArXiv 1303.3011

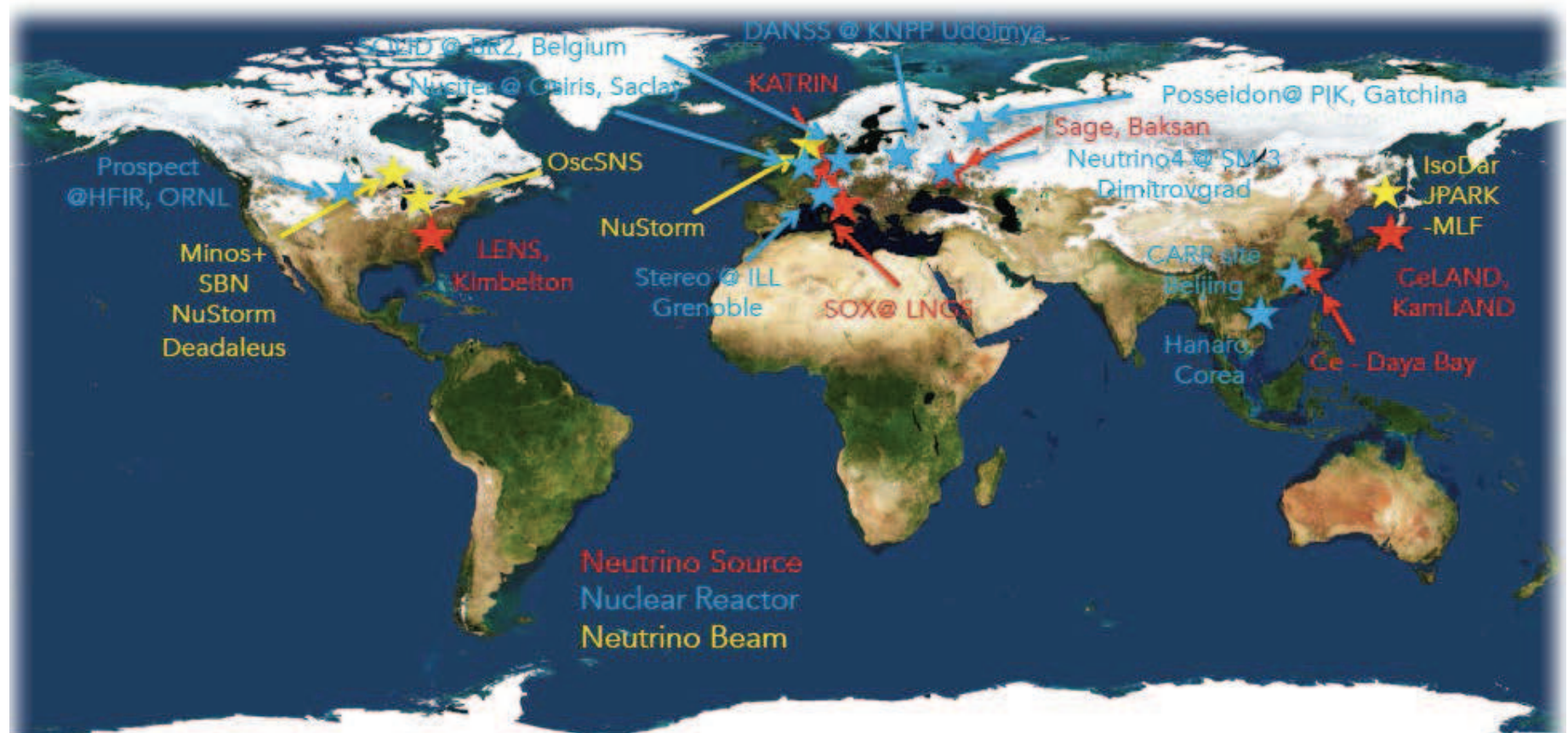


Giunti et al, ArXiv 1308.5288



- New generation of ν_e disappearance experiments \Rightarrow adding to the tension

Searches for eV sterile neutrinos



This talk: (anti-) ν_e disappearance only

$$P_{ee} = 1 - \sin^2 2\theta_{ee} \sin^2 \frac{\Delta m_{41}^2}{4E} \quad \& \quad \sin^2 2\theta_{ee} = |U_{e4}|^2 (1 - |U_{e4}|^2)$$

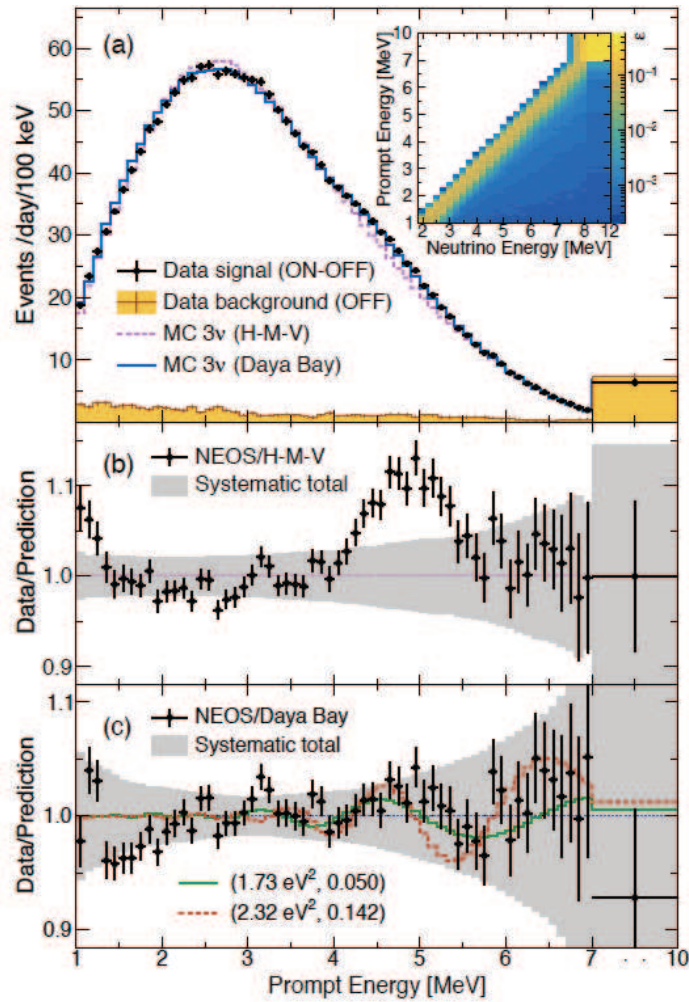
NEOS @ Hanbit (Korea)

arXiv:1610.05134v4 [hep-ex] 21 Mar 2017

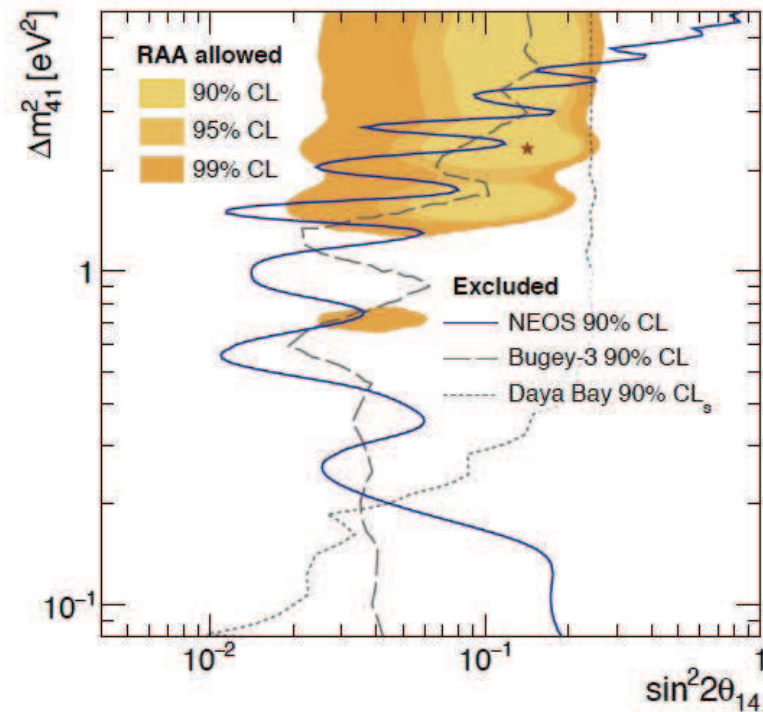
Single detector / single distance

Analysis done with Daya Bay spectrum as reference

Relies on 5 MeV bump!

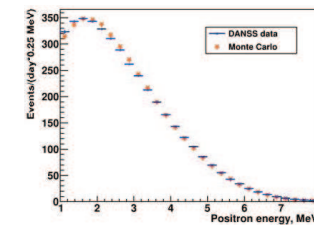
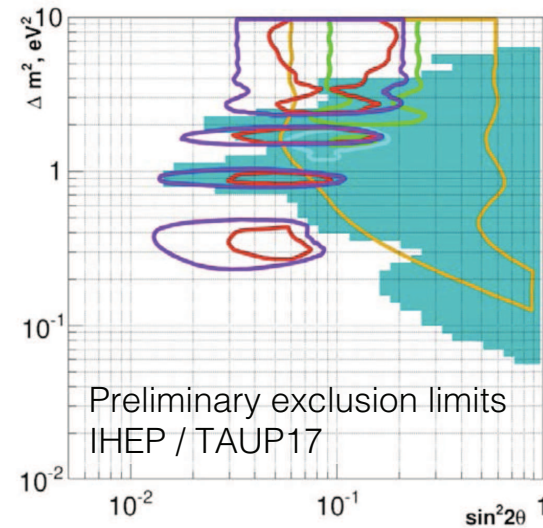
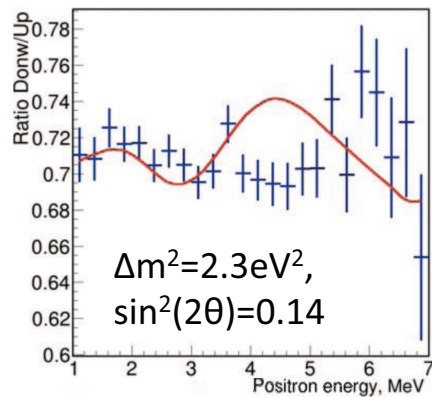
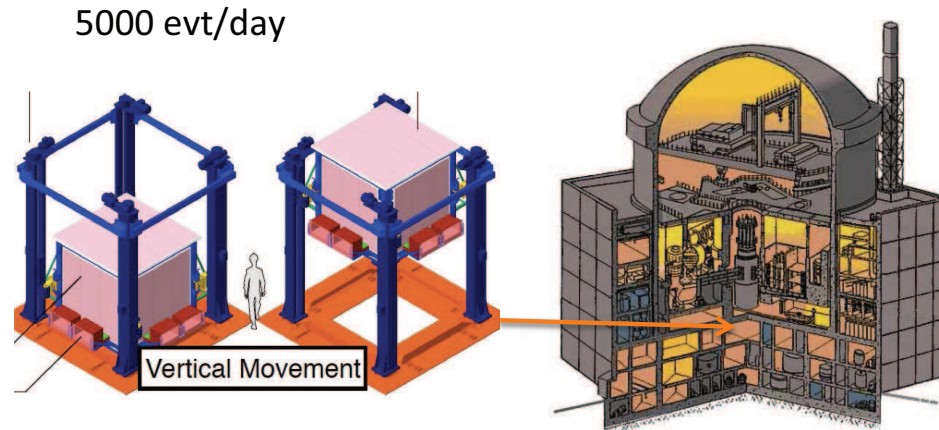


S. Schönert | TUM | Sterile neutrinos



DANSS @ Kalinin

- 3 GW extended core (5000 ev/day)
- Plastic strips with Gd-loaded interlayer, WLS fibers
- Vertical motion of the detector (9.7-12.2 m)
- Independent of burn-up or spectral feature



⇒ Decrease of the significance of the reactor anomaly

Dentler etal 1709.04294

⇒ Global fit with 3+N steriles severely disfavoured unless some data is dropped

Giunti etal 1703.00860

Confirmed Low Energy Picture and MY List of Q&A

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Only model independent probe of m_ν **β decay**: $\sum m_i^2 |U_{ei}|^2 \leq (2.2 \text{ eV})^2$
Anxiously waiting for Katrin
- **Dirac or Majorana?**: We do not know, *anxiously* waiting for ν -less $\beta\beta$ decay
- **Only three light states?** Tension between hints and bounds
New results from $\bar{\nu}_e$ disappearance further disfavour $\mathcal{O}(\text{eV})$ ν_s interpretation
- **Other NP at play?**

Alternative Oscillation Mechanisms

- Oscillations are due to:

- Misalignment between CC-int and propagation states: **Mixing** \Rightarrow **Amplitude**

- Difference phases of propagation states \Rightarrow **Wavelength**.

For Δm^2 -OSC $\lambda = \frac{4\pi E}{\Delta m^2}$

Alternative Oscillation Mechanisms

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- ν masses are not the only mechanism for oscillations

Violation of Equivalence Principle (VEP): Gasperini 88, Halprin, Leung 01

Non universal coupling of neutrinos $\gamma_1 \neq \gamma_2$ to gravitational potential ϕ

$$\lambda = \frac{\pi}{E|\phi|\delta\gamma}$$

Violation of Lorentz Invariance (VLI): Coleman, Glashow 97

Non universal asymptotic velocity of neutrinos $c_1 \neq c_2 \Rightarrow E_i = \frac{m_i^2}{2p} + c_i p$

$$\lambda = \frac{2\pi}{E\Delta c}$$

Interactions with space-time torsion: Sabbata, Gasperini 81

Non universal couplings of neutrinos $k_1 \neq k_2$ to torsion strength Q

$$\lambda = \frac{2\pi}{Q\Delta k}$$

Violation of Lorentz Invariance (VLI) Colladay, Kostelecky 97; Coleman, Glashow 99

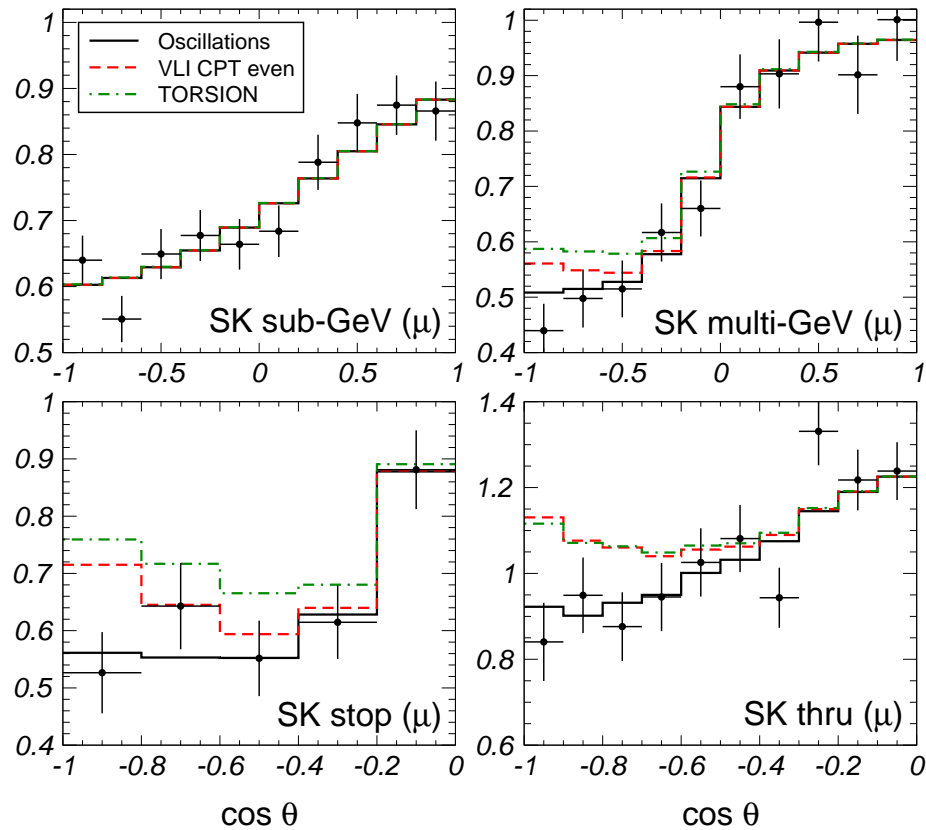
due to CPT violating terms: $\bar{\nu}_L^\alpha b_\mu^{\alpha\beta} \gamma_\mu \nu_L^\beta \Rightarrow E_i = \frac{m_i^2}{2p} \pm b_i$

$$\lambda = \pm \frac{2\pi}{\Delta b}$$

ATM ν 's: Subdominant NP Effects

- Using atmospheric neutrino data these effects can be constrained

MCG-G, M. Maltoni hepp-ph,0404085,0704.1800



At 90% CL:

$$\frac{|\Delta c|}{c} \leq 1.2 \times 10^{-24}$$

$$|\phi \Delta \gamma| \leq 5.9 \times 10^{-25}$$

$$|Q \Delta k| \leq 4.8 \times 10^{-23} \text{ GeV}$$

$$|\Delta b| \leq 3.0 \times 10^{-23} \text{ GeV}$$

- Including non-standard neutrino NC interactions with fermion f

$$\mathcal{L}_{\text{NSI}} = -2\sqrt{2}G_F \varepsilon_{\alpha\beta}^{fP} (\bar{\nu}_\alpha \gamma^\mu L \nu_\beta) (\bar{f} \gamma_\mu P f), \quad P = L, R$$

- In flavour basis $\vec{\nu} = (\nu_e, \nu_\mu, \nu_\tau)^T$ the neutrino evolution eq.:

$$i \frac{d}{dx} \vec{\nu} = H^\nu \vec{\nu} \quad \text{with} \quad H^\nu = H_{\text{vac}} + H_{\text{mat}} \quad \text{and} \quad H^{\bar{\nu}} = (H_{\text{vac}} - H_{\text{mat}})^*$$

$$H_{\text{mat}} = \sqrt{2}G_F N_e(r) \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} + \sqrt{2}G_F N_e(r) \begin{pmatrix} \varepsilon_{ee} & \varepsilon_{e\mu} & \varepsilon_{e\tau} \\ \varepsilon_{e\mu}^* & \varepsilon_{\mu\mu} & \varepsilon_{\mu\tau} \\ \varepsilon_{e\tau}^* & \varepsilon_{\mu\tau}^* & \varepsilon_{\tau\tau} \end{pmatrix}$$

$$\varepsilon_{\alpha\beta}(r) \equiv \sum_{f=ued} \frac{N_f(r)}{N_e(r)} \varepsilon_{\alpha\beta}^{fV} \Rightarrow 3\nu \text{ evolution depends on } 6 \text{ (vac)} + 8 \text{ per } f \text{ (mat)}$$

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\Rightarrow Parameters degeneracies (some well-known but being rediscovered lately ...)

In particular CPT \Rightarrow invariance under simultaneously:

$$\begin{aligned} \theta_{12} &\leftrightarrow \frac{\pi}{2} - \theta_{12}, & (\varepsilon_{ee} - \varepsilon_{\mu\mu}) &\rightarrow -(\varepsilon_{ee} - \varepsilon_{\mu\mu}) - 2, \\ \Delta m_{31}^2 &\rightarrow -\Delta m_{32}^2, & (\varepsilon_{\tau\tau} - \varepsilon_{\mu\mu}) &\rightarrow -(\varepsilon_{\tau\tau} - \varepsilon_{\mu\mu}), \\ \delta &\rightarrow \pi - \delta, & \varepsilon_{\alpha\beta} &\rightarrow -\varepsilon_{\alpha\beta}^* \quad (\alpha \neq \beta), \end{aligned}$$

NSI: Bounds/Degeneracies from/in Oscillation data

M.C G-G, M.Maltoni 1307.3092

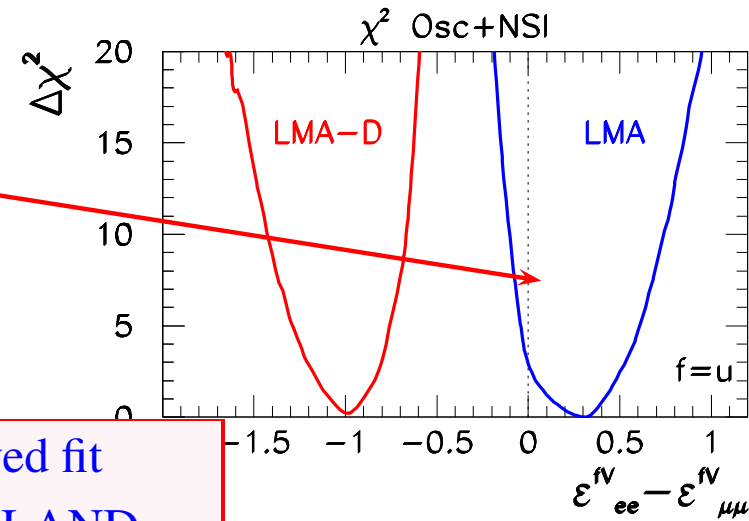
Param.	best-fit	90% CL	
		LMA	LMA-D
$\varepsilon_{ee}^u - \varepsilon_{\mu\mu}^u$	+0.298	[+0.00, +0.51]	\oplus [-1.19, -0.81]
$\varepsilon_{\tau\tau}^u - \varepsilon_{\mu\mu}^u$	+0.001	[-0.01, +0.03]	[-0.03, +0.03]
$\varepsilon_{e\mu}^u$	-0.021	[-0.09, +0.04]	[-0.09, +0.10]
$\varepsilon_{e\tau}^u$	+0.021	[-0.14, +0.14]	[-0.15, +0.14]
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- Bounds $\mathcal{O}(1 - 10\%)$
- Except $\varepsilon_{ee}^{q,V} - \varepsilon_{\mu\mu}^{q,V}$

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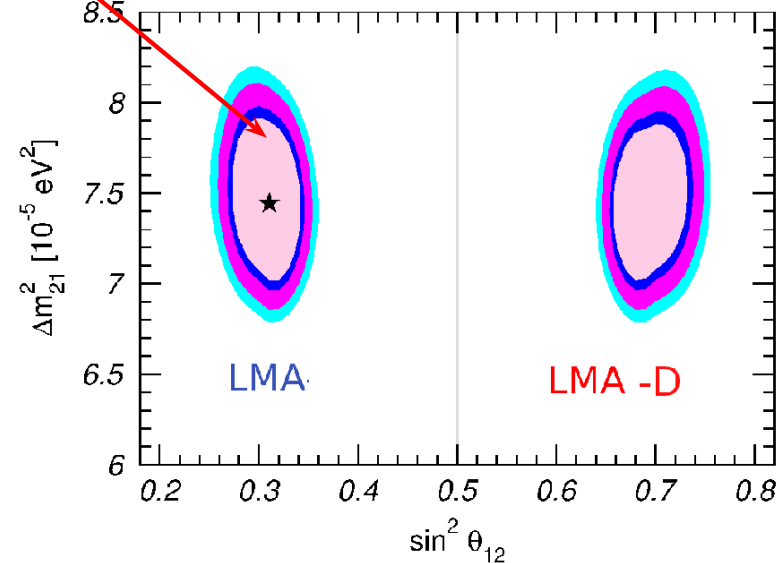
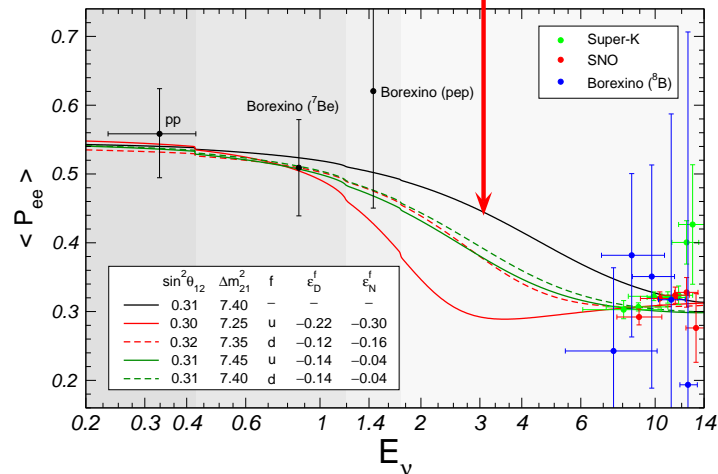
M.C G-G, M.Maltoni 1307.3092

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LMA: Improved fit to Solar+KamLAND



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M.C G-G, M.Maltoni 1307.3092

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– Bounds $\mathcal{O}(1 - 10\%)$

– Except $\varepsilon_{ee}^{q,V} - \varepsilon_{\mu\mu}^{q,V}$

Degenerate solution LMA-D ($\theta_{12} > 45^\circ$)

Miranda, Tortola, Valle, hep-ph/0406280

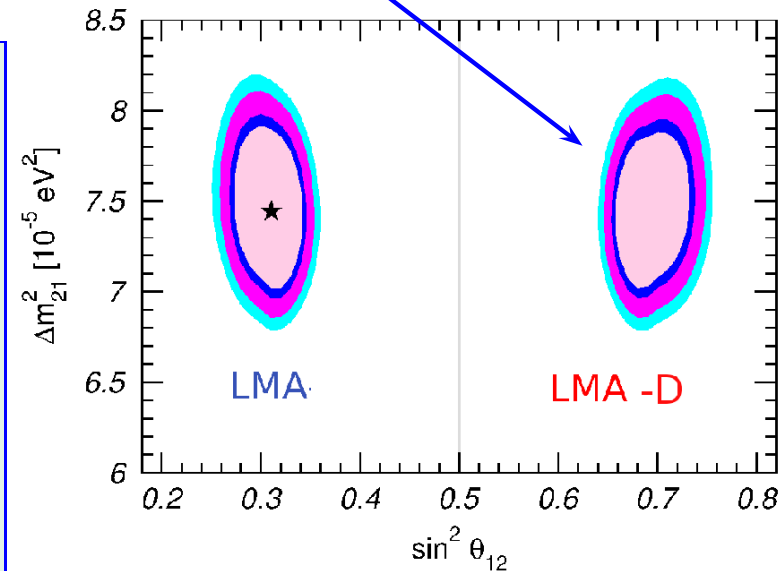
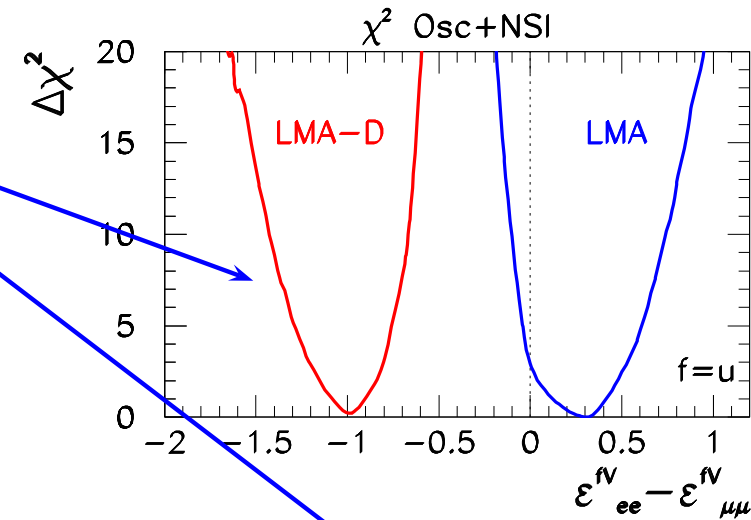
Cannot be resolved with osc-experiments

Requires NC scattering experiments

Coloma et al 1701.04828

Requires NSI $\sim G_F$ (light mediators?)

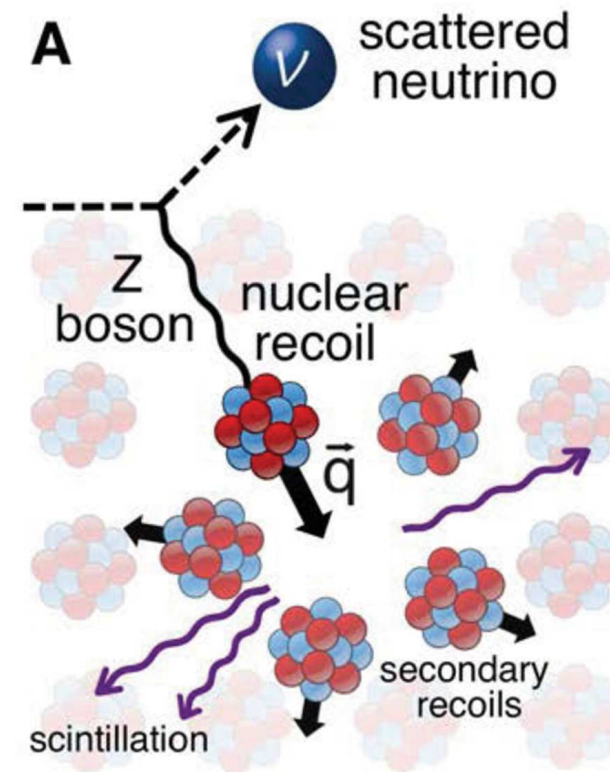
Farzan 1505.06906, and Shoemaker 1512.09147



COHERENT EXPERIMENT

Science 2017 [ArXiv:1708.01294]

- observation of coherent neutrino-nucleus scattering at 6.7σ at CsI[Na] detector
- neutrinos from stopped pion source at Oak Ridge NL
- 142 events observed, in agreement with Standard Model



NSI: Combination with COHERENT data

Coloma, MCGG, Maltoni, Schwetz ArXiv:1708.02899

- COHERENT has detected for first time Coherent νN scattering **1708.01294**:
 142($1 \pm 0.28(\text{sys})$) observed events over a steady bck of 405
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- In presence of NSI: $N_{\text{NSI}}(\varepsilon) = \gamma [f_{\nu_e} Q_{we}^2(\varepsilon) + (f_{\nu_\mu} + f_{\bar{\nu}_\mu}) Q_{w\mu}^2(\varepsilon)]$

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NSI: Combination with COHERENT data

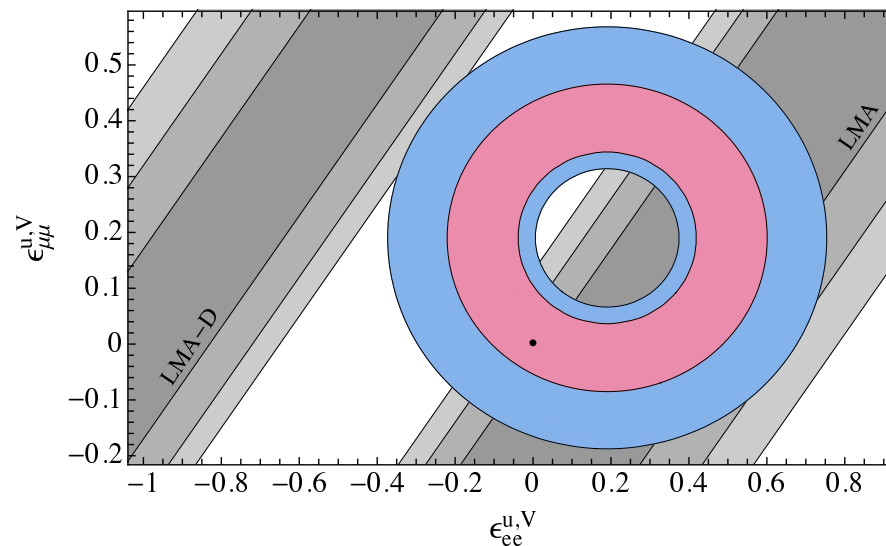
Coloma, MCGG, Maltoni, Schwetz ArXiv:1708.02899

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- Impact on LMA-D: Allowed COHERENT region vs LMA-D required range



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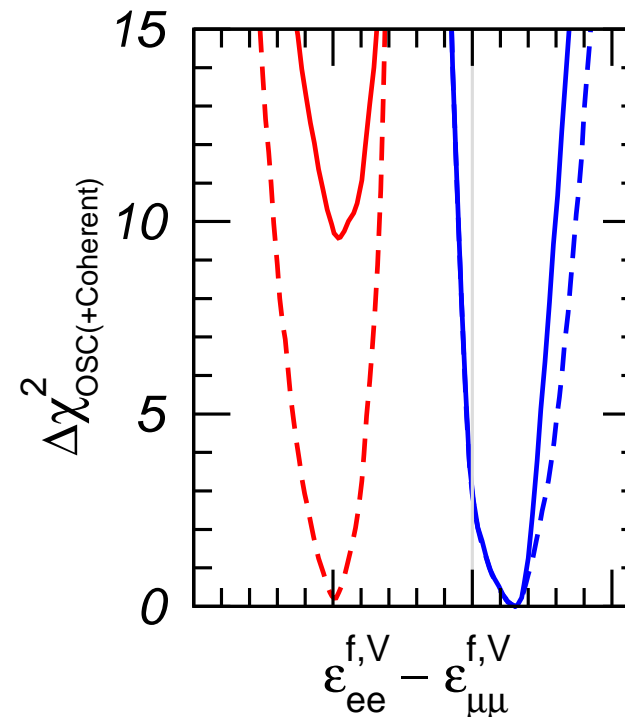
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- OSCILLATION + COHERENT \Rightarrow **LMA-D excluded at more than 3.1 σ**

All NSI's constrained

	$f = u$	$f = d$
$\varepsilon_{ee}^{f,V}$	[0.028, 0.60]	[0.030, 0.55]
$\varepsilon_{\mu\mu}^{f,V}$	[-0.088, 0.37]	[-0.075, 0.33]
$\varepsilon_{\tau\tau}^{f,V}$	[-0.090, 0.38]	[-0.075, 0.33]
$\varepsilon_{e\mu}^{f,V}$	[-0.073, 0.044]	[-0.07, 0.04]
$\varepsilon_{e\tau}^{f,V}$	[-0.15, 0.13]	[-0.13, 0.12]
$\varepsilon_{\mu\tau}^{f,V}$	[-0.01, 0.009]	[-0.009, 0.008]



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A Detour in the Sun

- Sun=Main sequence star
- Solar Models describes the Sun based on:
 - Mass: $M_{\odot} = 2 \times 10^{33}$ gr
 - Radius: $R_{\odot} = 7 \times 10^5$ km
 - Surf Lum: $L_{\odot} = 3.842 \times 10^{33} (1 \pm 0.004)$ erg/sec
 - Age: $\tau_{\odot} = 4.57 \times 10^9 (1 \pm 0.0044)$ yr
- Basic assumptions:
 - The Sun is spherically symmetric
 - Some Equation of State
- Incorporate:
 - Transport of Energy: Radiative and Convective
 - ⇒ Model of opacities
 - Chemical Evolution by Nuclear Reactions
 - ⇒ pp-chain and CNO cycles
 - Microscopic Diffusion

- Using inputs from:
 - Lab Measurements of Nuclear Rates
 - Element Abundance Determination By
 - ⇒ Spectroscopy of Photosphere: C, N, O
 - ⇒ Meteorites: Mg,Si,S,Fe
 - ⇒ Other methods: Ne, Ar
- They Predict Observables:
 - Neutrino Flux Spectrum
 - Relevant to Helioseismology :
 - ⇒ Surface He Abundance
 - ⇒ Inner Radius of Convective Zone
 - ⇒ Sound Speed Profile

The Solar Composition Problem

– Newer determination of abundances in solar surface give lower values

$$\log \epsilon_i \equiv \log N_i / N_H + 12$$

Element	GS98	AGSS09met
C	8.52 ± 0.06	8.43 ± 0.05
N	7.92 ± 0.06	7.83 ± 0.05
O	8.83 ± 0.06	8.69 ± 0.05
Mg	7.58 ± 0.01	7.53 ± 0.01
Si	7.56 ± 0.01	7.51 ± 0.01
S	7.20 ± 0.06	7.15 ± 0.02
Fe	7.50 ± 0.01	7.45 ± 0.01
Ar	6.40 ± 0.06	6.40 ± 0.13
Ne	8.08 ± 0.06	7.93 ± 0.10

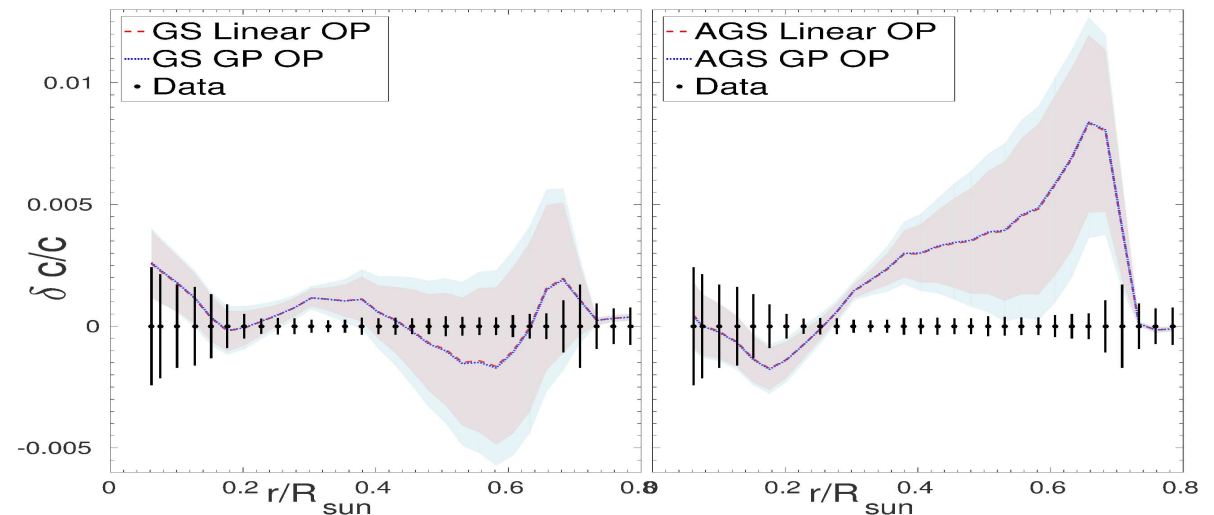
⇒ Two sets of SSM:

Starting from Bahcall *etal* 05, Serenelli *etal* 2016

B16-GS98 with old abund

B16-AGSS09met with new abund

– Solar Models with lower metallicities fail in reproducing helioseismology data

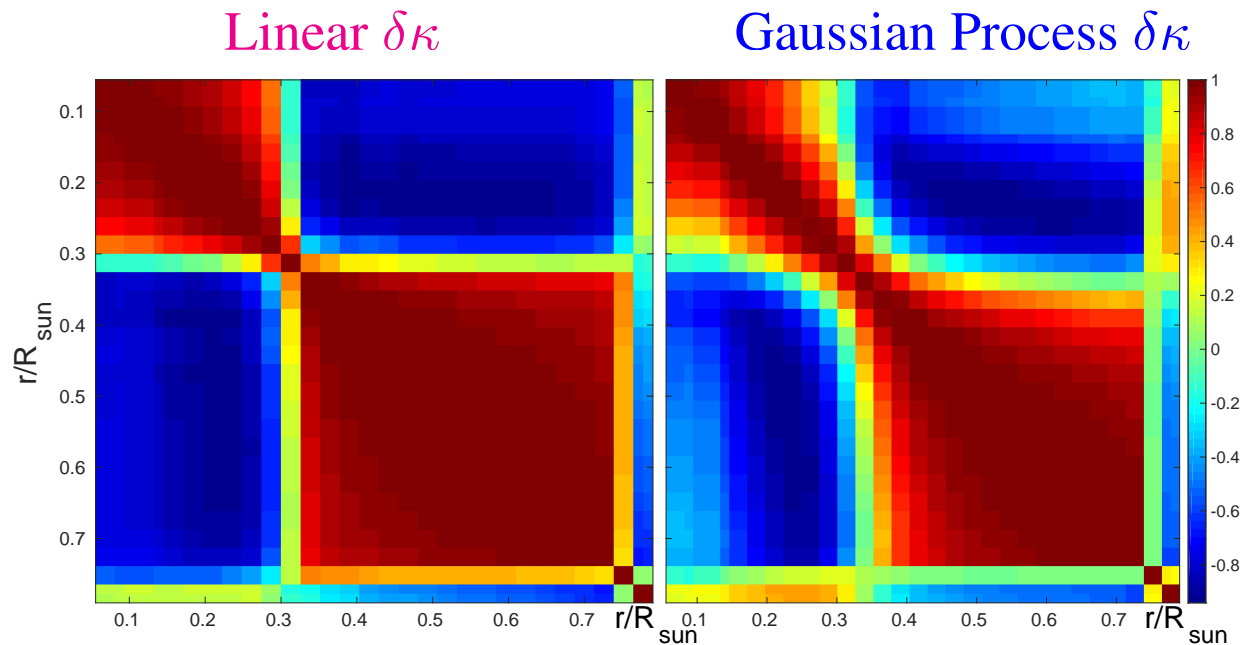


Predictions very strongly correlated

- B16-GS98 (dis)agreement at 2.5σ
- B16-AGSS09 disagreement 4.7σ
- Bayes factor B16-AGSS09/B16-GS98 < -13
(very strong disfavouring)

Modeling the uncertainty in the opacity profile

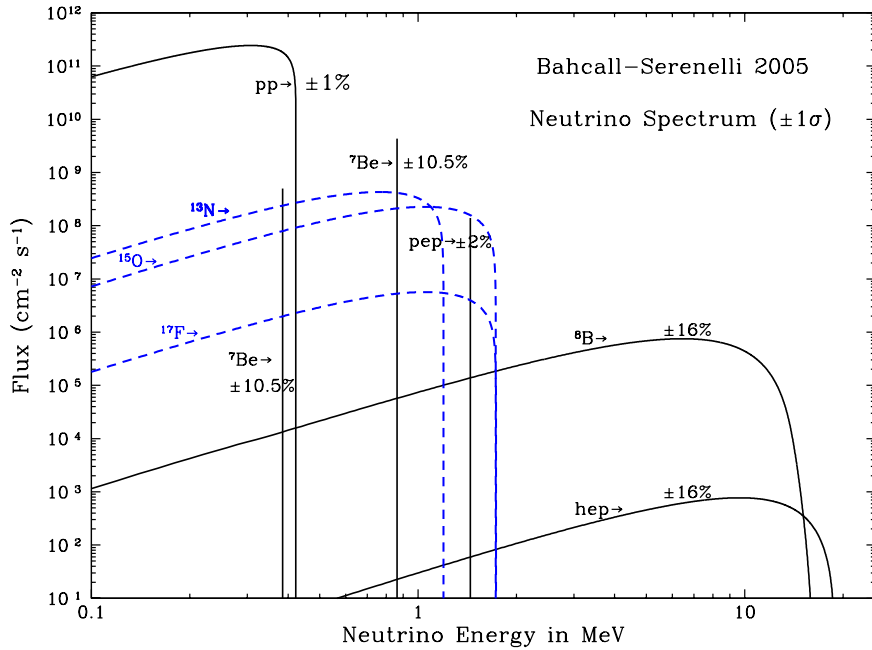
- Opacity is a function $\kappa(T, \rho, X_i = N_i/N_H)$. How to parametrize its uncertainty?
- Generically $(1 + \delta\kappa(T))\langle\kappa(T, \rho, X_i)\rangle$
 - \Rightarrow Most studies $\delta\kappa(T) = C$ or $\delta\kappa(T) = a + b \log T$ with prior for σ_C (or σ_a, σ_b)
 - \Rightarrow only very rigid variations allowed
- Alternative: **Gaussian Process** ansatz with same $\sigma(T)$ but correlation length $L < 1$



Song, MCG-G, Serenelli, Villante (17)

Still, even with GP opacity uncertainty Bayes factor B16-AGSS09/B16-GS98=-4.1
(Moderate to strong disfavour)

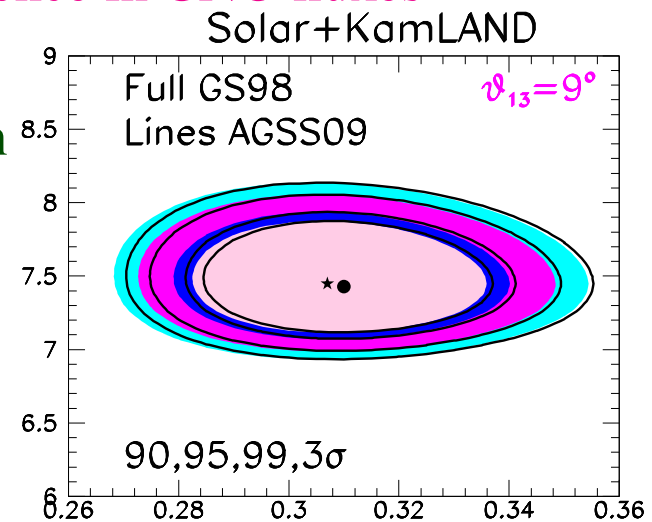
The Neutrino Fluxes



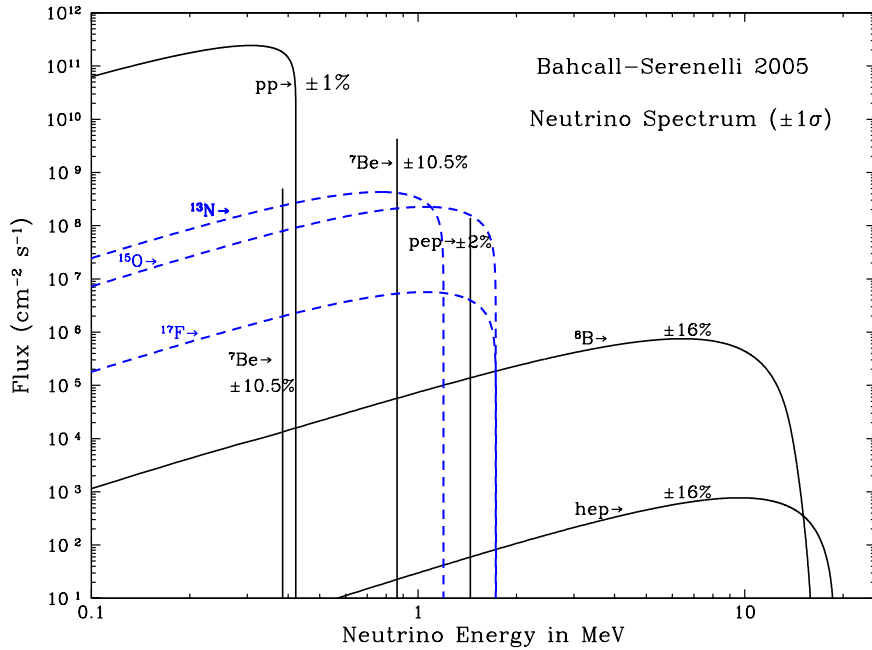
Flux $\text{cm}^{-2} \text{s}^{-1}$	B16GS98	B16-AGSS09met	Diff (%)
pp/ 10^{10}	5.98	6.03 (1 ± 0.005)	0.8
pep/ 10^8	1.44	1.46 (1 ± 0.01)	2.1
hep/ 10^3	7.98	8.25 (1 ± 0.30)	3.4
$^{7}\text{Be}/10^9$	4.93	4.40 (1 ± 0.06)	8.8
$^{8}\text{B}/10^6$	5.46	4.50 (1 ± 0.12)	17.7
$^{13}\text{N}/10^8$	2.78	2.04 (1 ± 0.14)	26.7
$^{15}\text{O}/10^8$	2.05	1.44 (1 ± 0.16)	30.0
$^{17}\text{F}/10^{16}$	5.29	3.26 (1 ± 0.18)	38.4

Most difference in CNO fluxes

– Negleageable Impact in Osc Parameter Determination



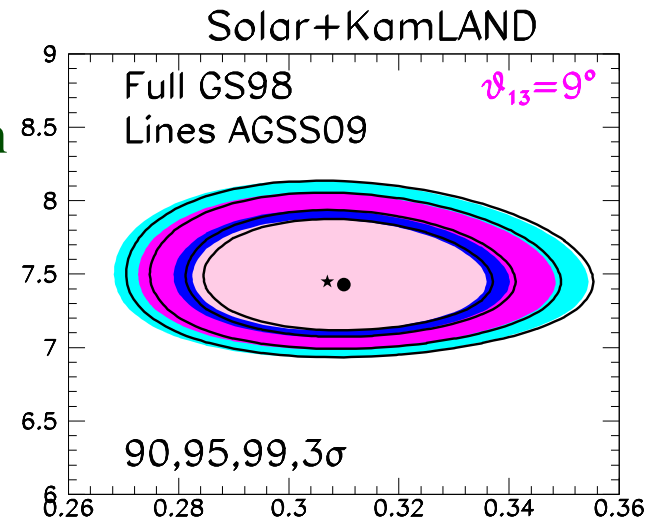
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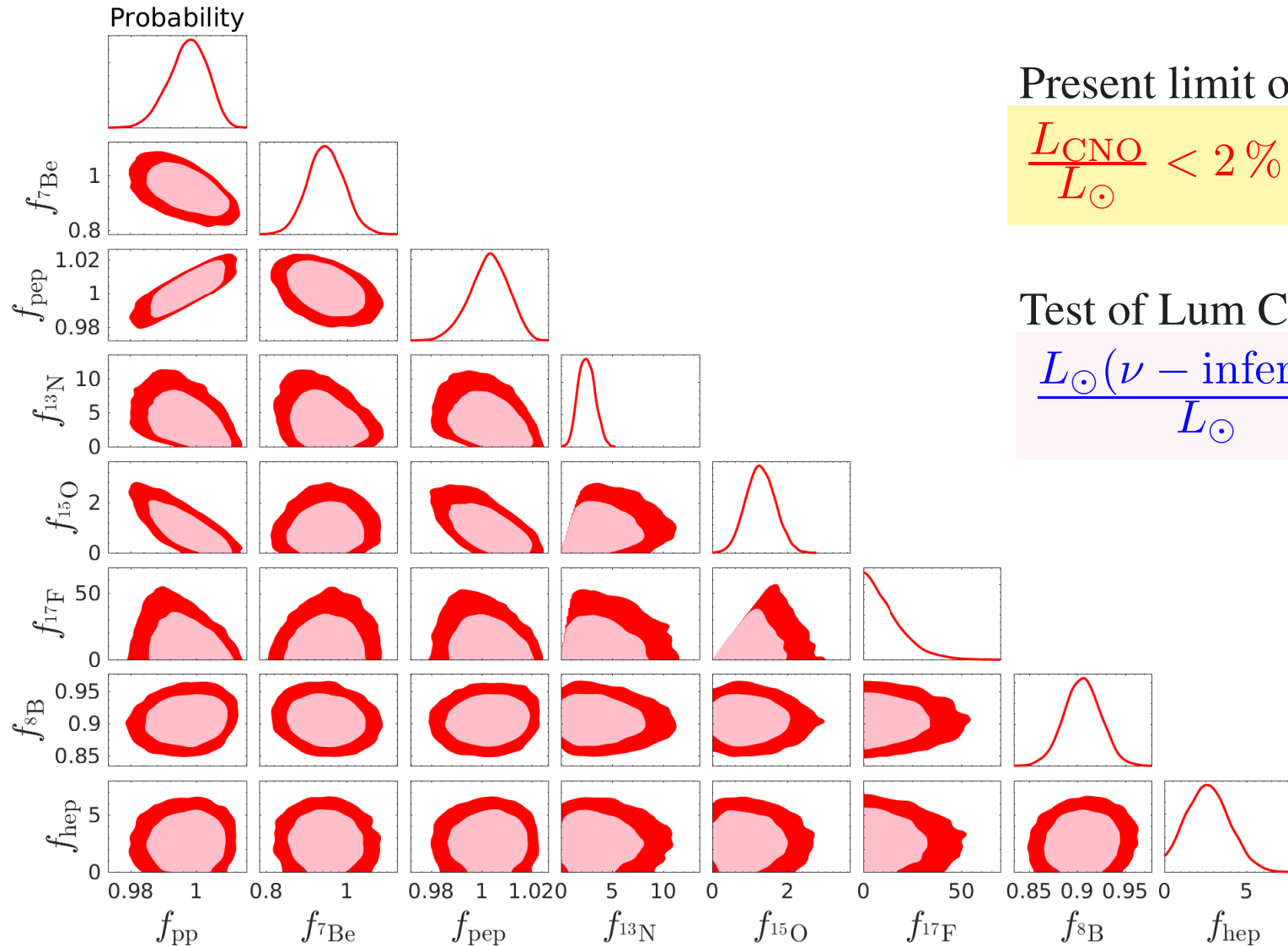
Most difference in CNO fluxes

– Negleageable Impact in Osc Parameter Determination
 \Rightarrow Possible to extract fluxes for data



Testing How the Sun Shines with ν 's

Results of Oscillation analysis with solar flux normalizations free: $f_i = \frac{\Phi_i}{\Phi_{GS98}}$



Present limit on CNO:

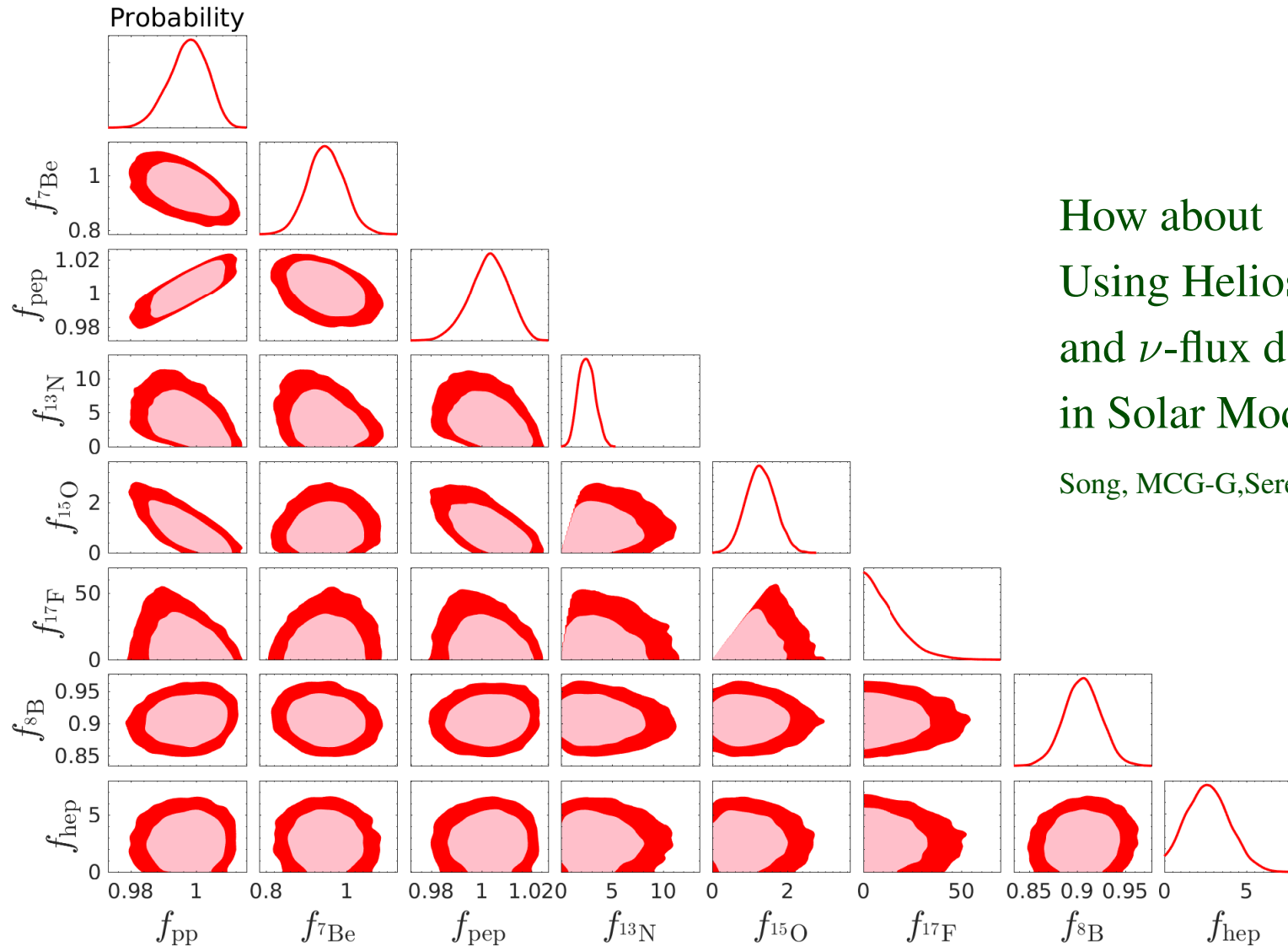
$$\frac{L_{\text{CNO}}}{L_{\odot}} < 2\% (3\sigma)$$

Test of Lum Constraint:

$$\frac{L_{\odot}(\nu - \text{inferred})}{L_{\odot}} = 1.04 \pm 0.07$$

Testing How the Sun Shines with ν 's

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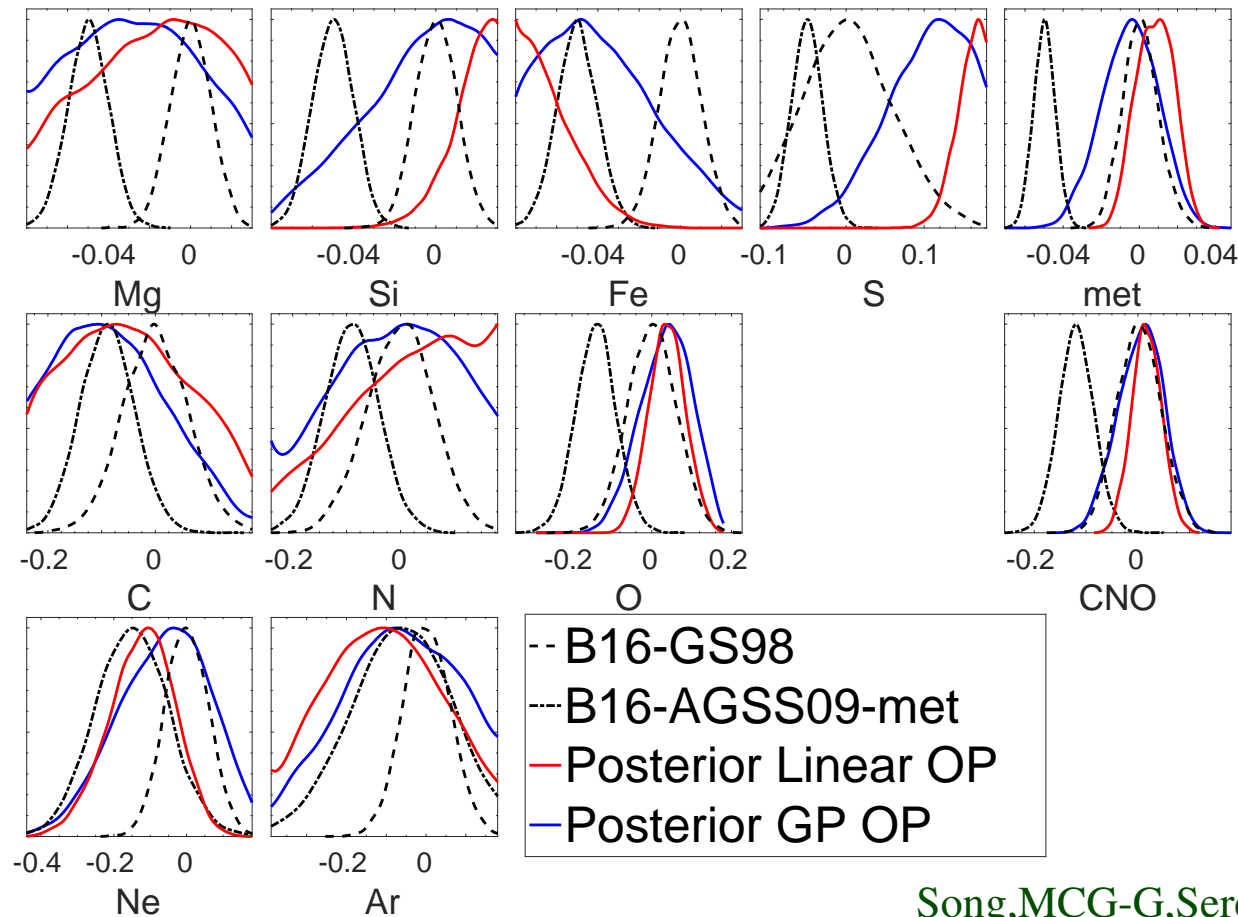
How about
Using Helioseismic
and ν -flux data
in Solar Modeling?

Song, MCG-G, Serenelli, Villante (17)

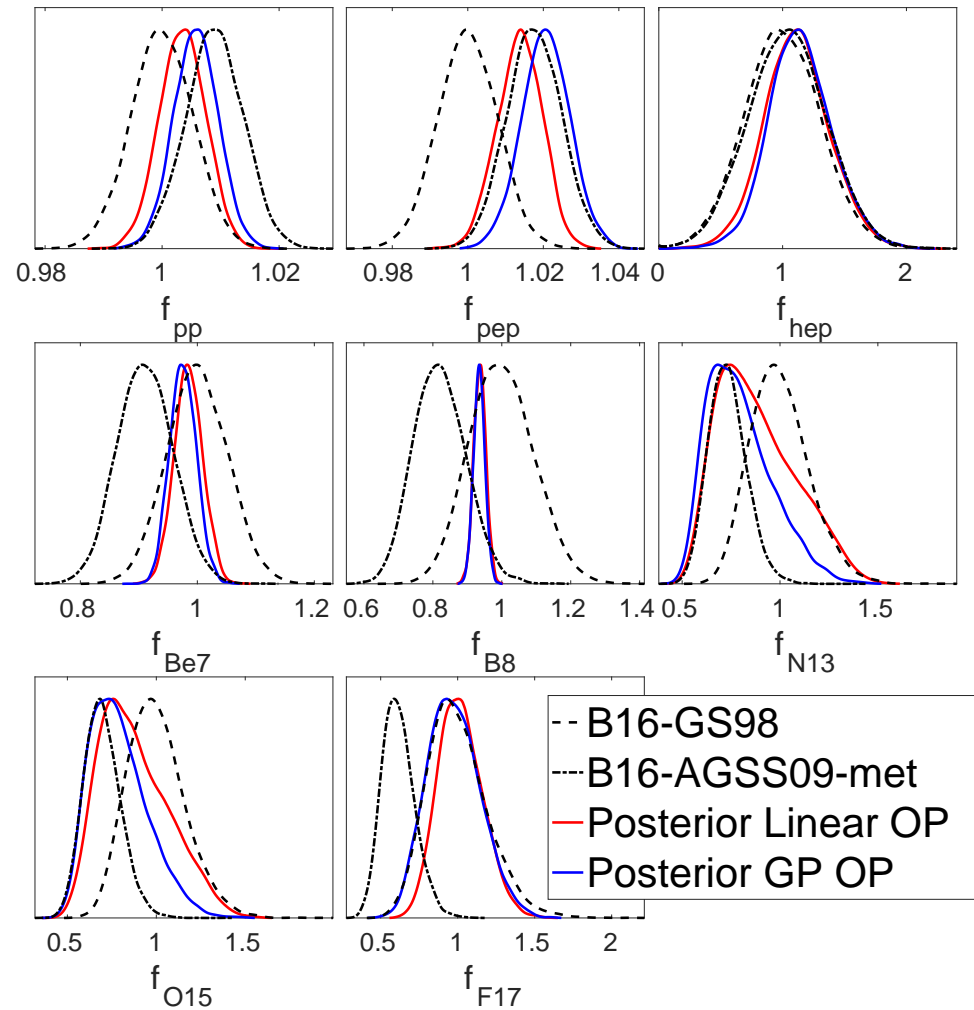
Using ν and Helioseismic Data in Sun Modeling

- Proposal: Invert approach and use the ν and helioseismic data in construction of SSM
- Method: Bayesian Inference of Abundance Posterior Distrib (from Uniform Priors)
- Test effects of effects of other modeling aspects (f.e. opacity uncertainty profiles)

$$x = \ln \frac{N_i}{N_H} - \left\langle \ln \frac{N_i}{N_H} \right\rangle_{GS98}$$



ν fluxes from ν +helioseismic data



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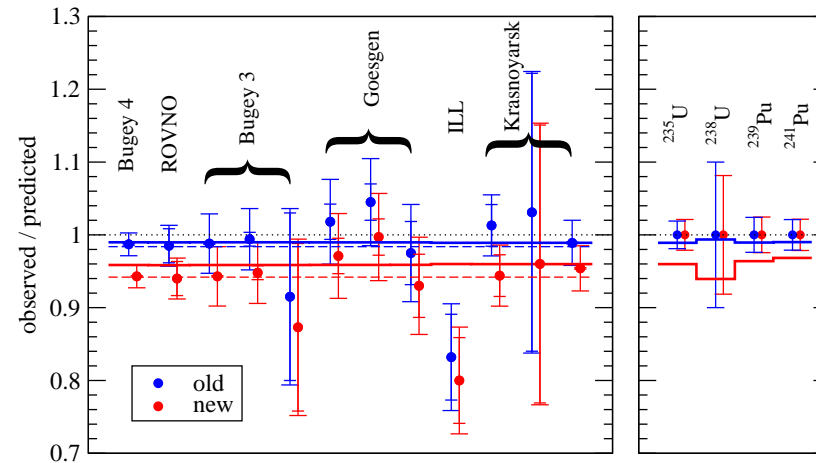
Thank You

3 ν Analysis: Reactor Flux anomaly and θ_{13}

- The reactor $\bar{\nu}_e$ fluxes was recalculated about 6 yrs ago

T.A. Mueller et al., [arXiv:1101.2663]; P. Huber, [arXiv:1106.0687].

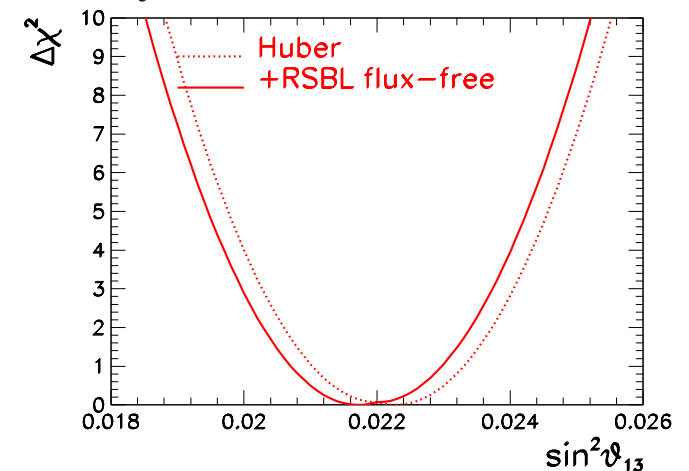
- Both found higher fluxes $\sim 3.5\%$
- \Rightarrow *negative* reactor experiments
at short baselines (RSBL) indeed
observed a deficit



- For 3ν analysis a consistent approach (T. Schwetz et. al. [arXiv:1103.0734]):
 - Fit oscillation parameters and reactor fluxes simultaneously
 - Use calculated fluxes (a) or RSBL data (b) as priors

Difference at $\lesssim 0.3\sigma$ level

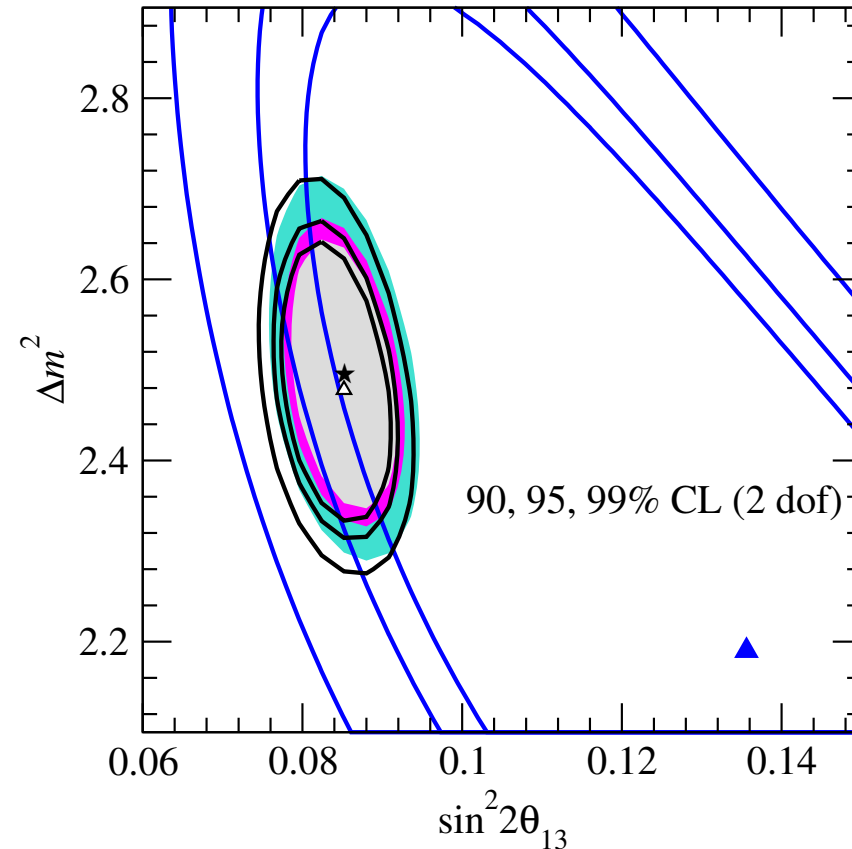
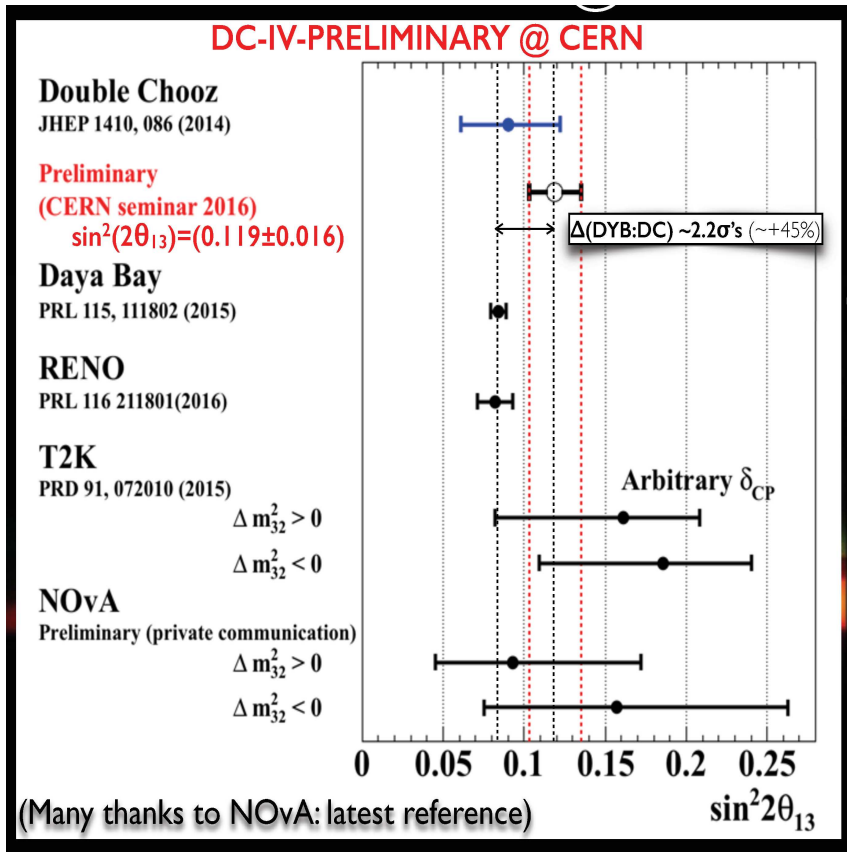
$$\chi_{min,a}^2 - \chi_{min,b}^2 \sim 7$$



Issues in 3 ν Analysis: Consistency of θ_{13}

Daya Bay vs Double Chooz?

Allowed regions of DC vs Daya Bay



From DC (Anatael Cabrera) Talk CERN Sep 16

Fig. Courtesy of T. Schwetz

No significant discrepancy

Lepton Mixing Unitarity

- Previous results assume U_{LEP} to be unitary
- If ν_L mixed with m extra states $U_{\text{LEP}} = (K_l, K_h)$ Schechter, Valle (1980)
And $U_{\text{LEP}} U_{\text{LEP}}^\dagger = I_{3 \times 3}$ but in general $U_{\text{LEP}}^\dagger U_{\text{LEP}} \neq I_{(3+m) \times (3+m)}$
- If m states are heavy ($M \gg E_\nu$) oscillations measure K_L (not unitary)

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Flavour Changing Neutral Currents

- But this **unitarity violation** \Rightarrow Flavour Violation in Charged Lepton Processes
Universality Violation of Charge Current ...

- Constraints on these processes limit leptonic unitarity violation to

$$|K_l K_l^\dagger| = \begin{pmatrix} 0.9979 - 0.9998 & < 10^{-5} & < 0.0021 \\ < 10^{-5} & 0.9996 - 1.0 & < 0.0008 \\ < 0.0021 & < 0.0008 & 0.9947 - 1.0 \end{pmatrix}$$

Antusch *et al* ArXiv:1407.6607

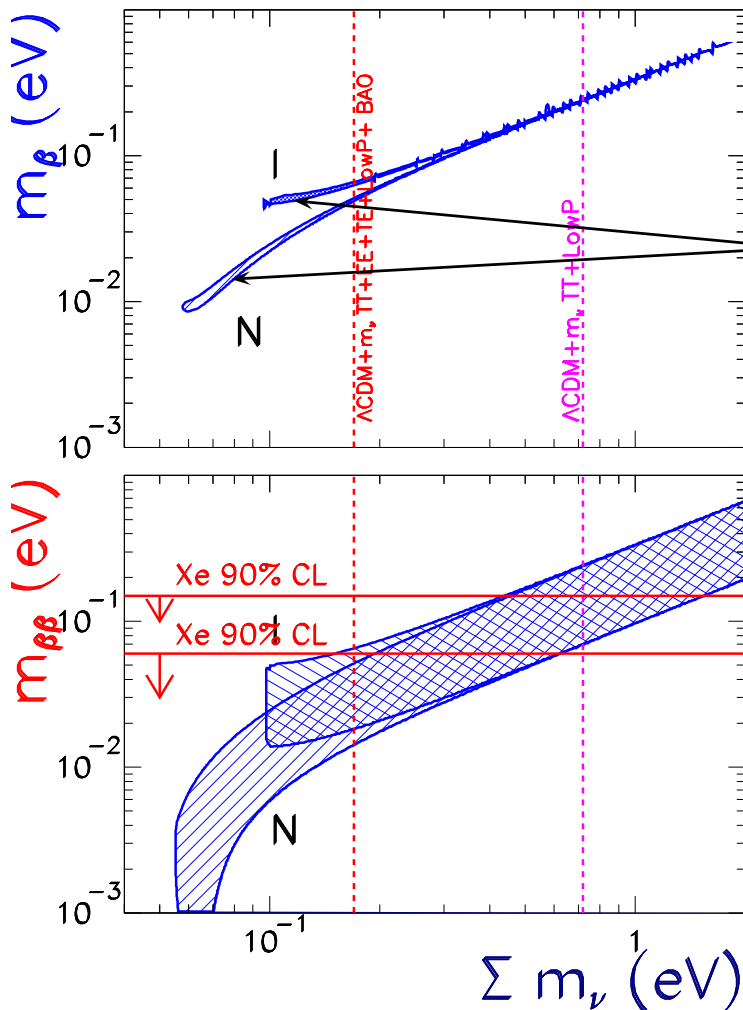
or equivalently $K_l \simeq (I + \epsilon)U(\theta_{ij}, \delta, \eta_i)$ with $|\epsilon_{\alpha j}| \leq \text{few} \times 10^{-3}$ while $K_h \sim \mathcal{O}(\epsilon)$

Neutrino Mass Scale: The Cosmo-Lab Connection

Global oscillation analysis

⇒ Correlations m_{ν_e} , m_{ee} and $\sum m_i$
(Fogli *et al* (04))

Nufit (95%)



Lower bound on $\sum m_i$ depends on ordering

Precision determination/bound of $\sum m_i$ can give information on ordering ?

Hannestad, Schwetz 1606.04691, Simpson *et al* 1703.03425, Capozzi *et al* 1703.04471 ...

Or much ado about nothing?

Cosmo data will only add to N/I likelihood when accuracy on $\sum m_{\nu}$ better than 0.02 eV (to see a 2σ N/I difference between 0.06 and 0.1)

Hannestad, Schwetz 1606.04691