

Testing the low scale seesaw and leptogenesis

Juraj Klarić (TU München)

based on 1606.6690 and 1609.09069 with Marco Drewes, Björn Garbrecht and Dario Gueter

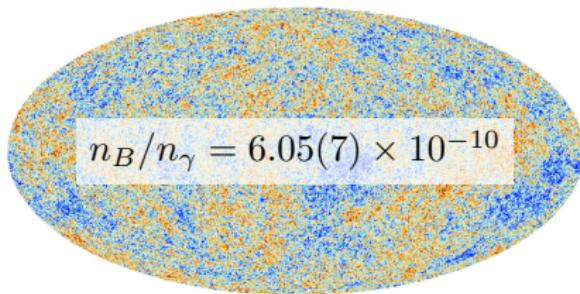
Neutrinos in Cosmology, in Astro- Particle- and Nuclear physics

EMFCSC, Erice, Sicily, 17. September 2017

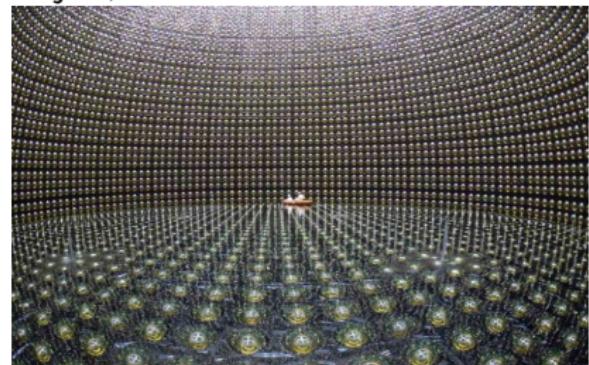


Some of the missing pieces of the standard model:

BAU baryon asymmetry of the universe
WMAP, Planck and Big bang nucleosynthesis:

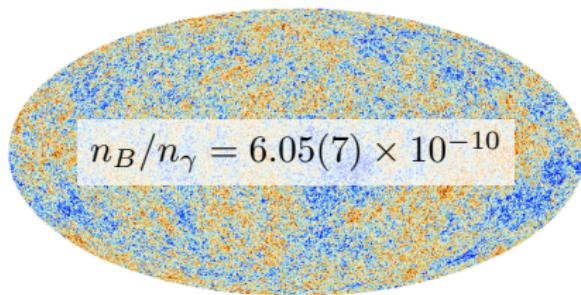


Neutrino masses
Nobel prize 2015
Kajita, McDonald

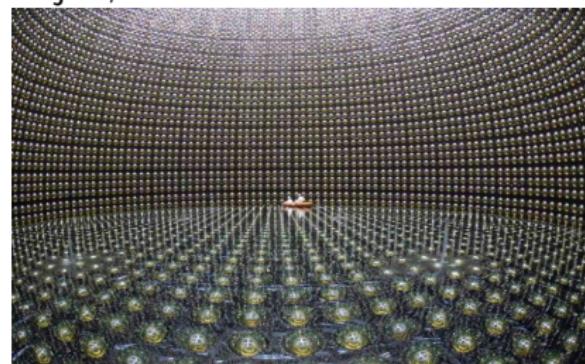


Some of the missing pieces of the standard model:

BAU baryon asymmetry of the universe
WMAP, Planck and Big bang nucleosynthesis:

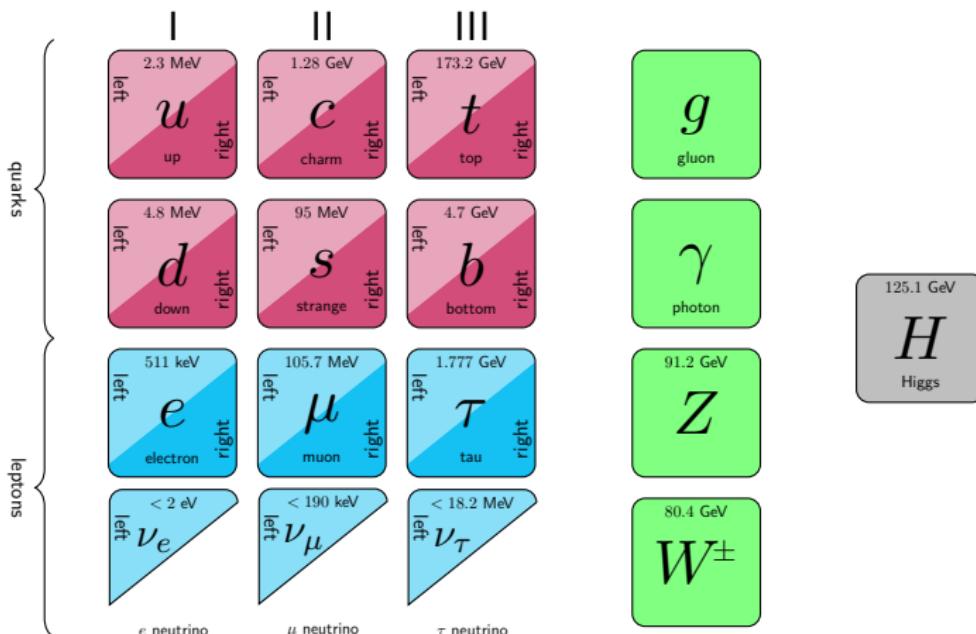


Neutrino masses
Nobel prize 2015
Kajita, McDonald

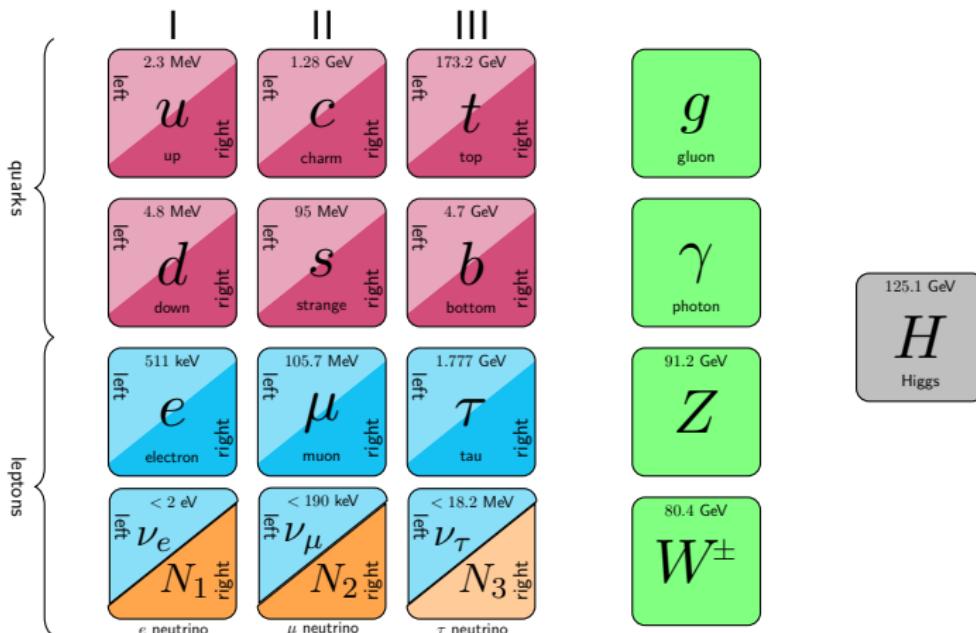


Is there a way to explain both?

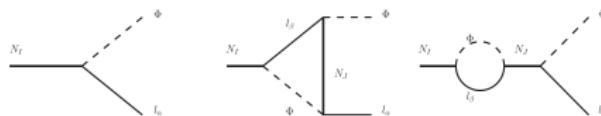
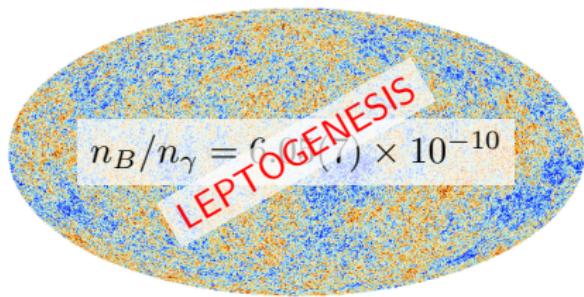
Standard Model



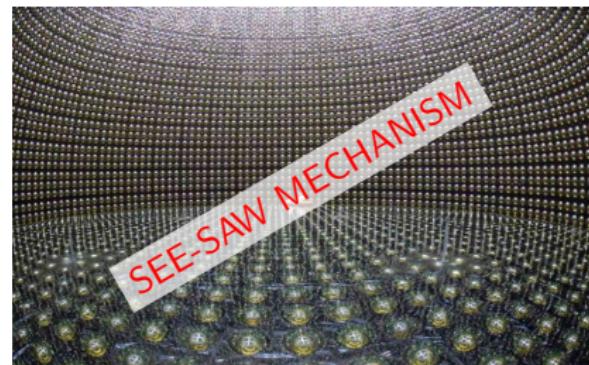
Standard Model



BAU baryon asymmetry of the universe
 WMAP, Planck and Big bang nucleosynthesis:



Neutrino masses
 Nobel prize 2015
 Kajita, McDonald



$$m_\nu = -v^2 Y^\dagger M_M^{-1} Y^*$$

Seesaw Mechanism

- Dirac Mass $m_D = vY^\dagger$
- Right handed neutrino (RHN) Majorana mass M_M

$$\mathcal{L} \supset \frac{1}{2} \begin{pmatrix} \overline{\nu_L} \\ \bar{N} \end{pmatrix} \begin{pmatrix} 0 & m_D \\ m_D^T & M_M \end{pmatrix} \begin{pmatrix} \nu_L & N \end{pmatrix}$$

Active neutrino masses

$$m_\nu = -m_D M_M^{-1} m_D^T$$

Mixing with RHN

$$|U_{ai}|^2 = \left| \left(m_D M_M^{-1} \right)_{ai} \right|^2$$

Constraints on RHN parameters

- Direct constraints - past experiments
- Seesaw constraints - neutrino oscillation data
radiatively corrected Casas-Ibarra parametrization

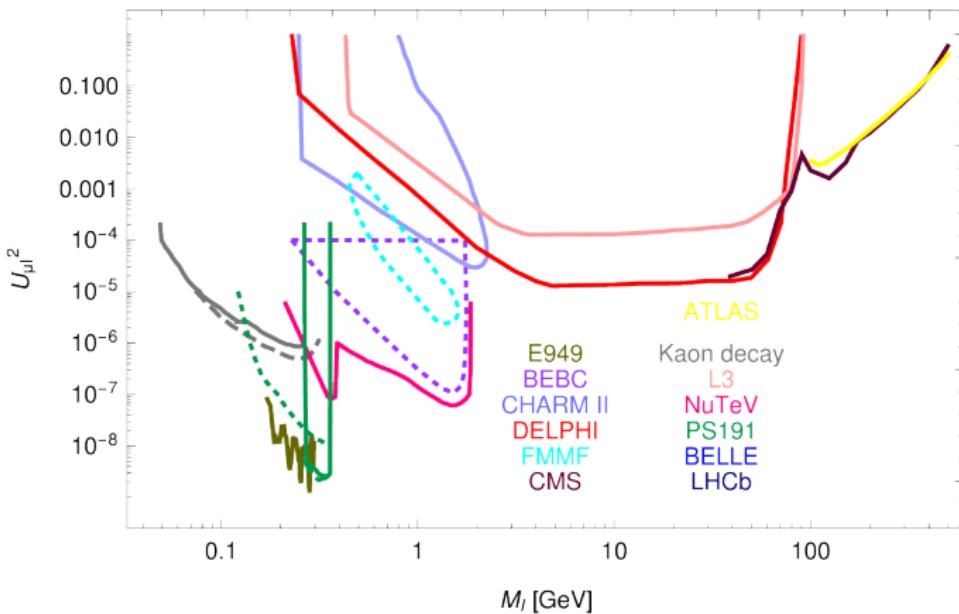
[Lopez-Pavon/Molinaro/Petcov 1506.05296]

- Cosmological constraints - BBN $\tau_N < 0.1\text{s}$

- Indirect constraints

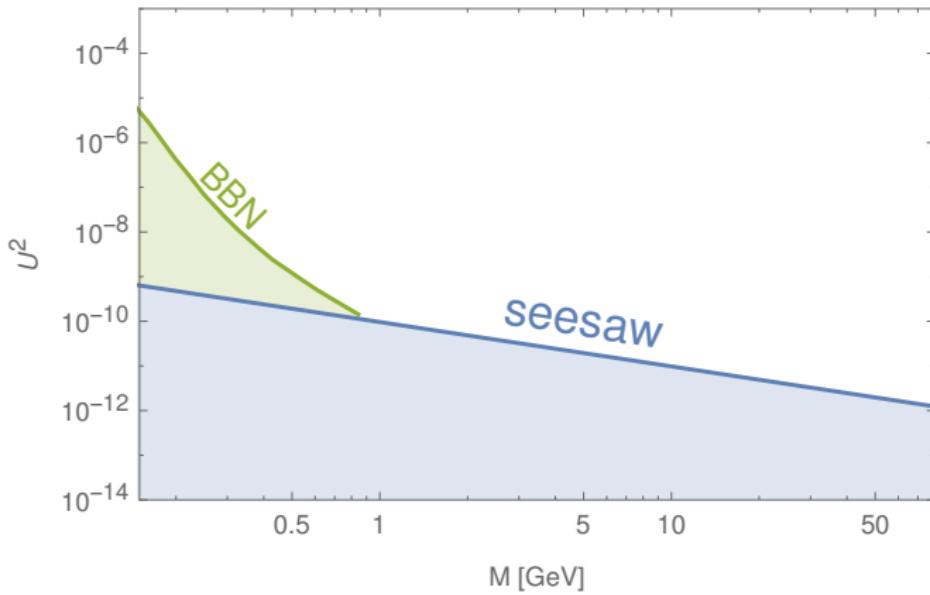
- neutrinoless double β decay
- lepton universality
- CKM universality
- electroweak precision data
- LFV in rare lepton decays:
 - $\mu \rightarrow e\gamma$
 - $\tau \rightarrow e\gamma$
 - $\tau \rightarrow \mu\gamma$

Direct constraints

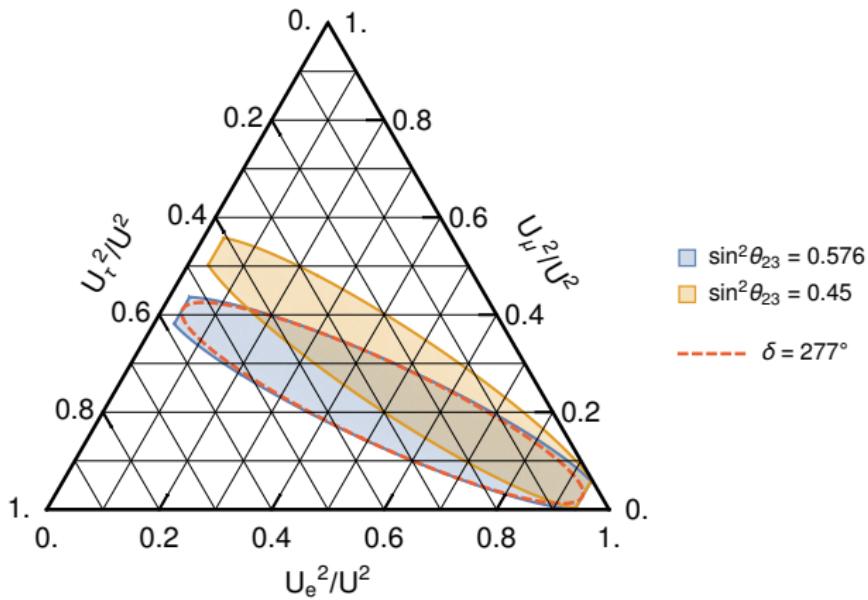


[Plot from arXiv:1502.00477]

Seesaw and BBN constraints

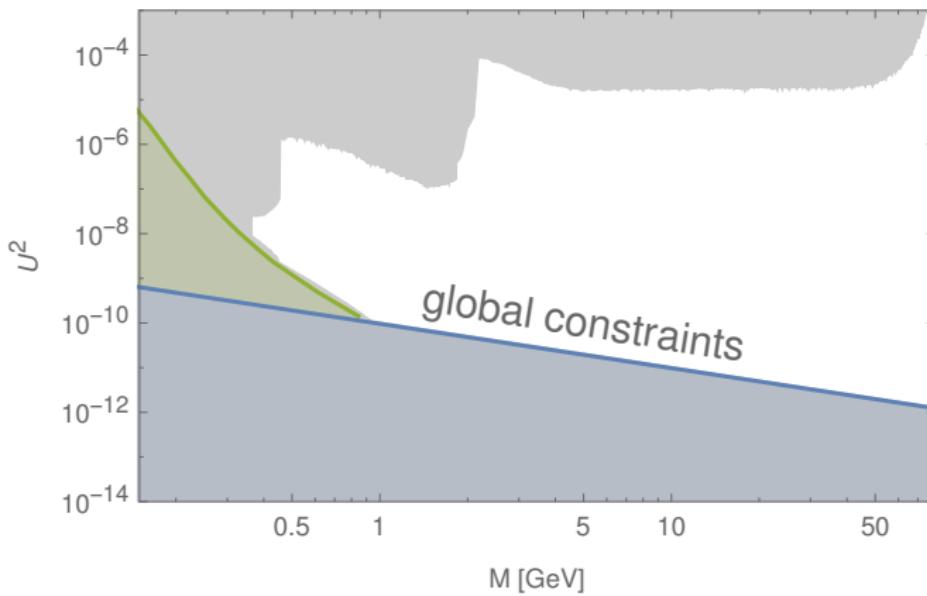


Constraints on flavour patterns: Inverted hierarchy

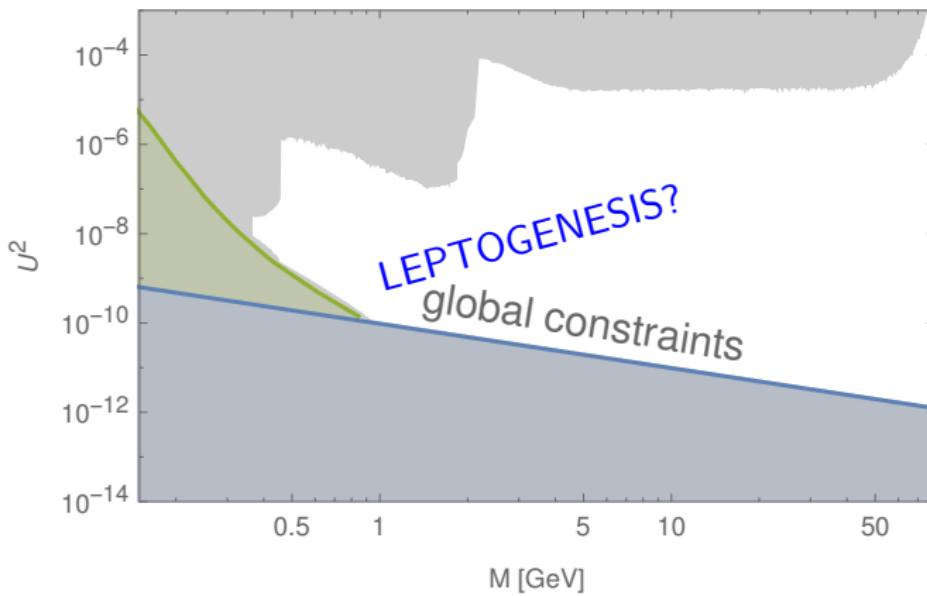


[Drewes/Garbrecht/Gueter/JK 1609.09069]

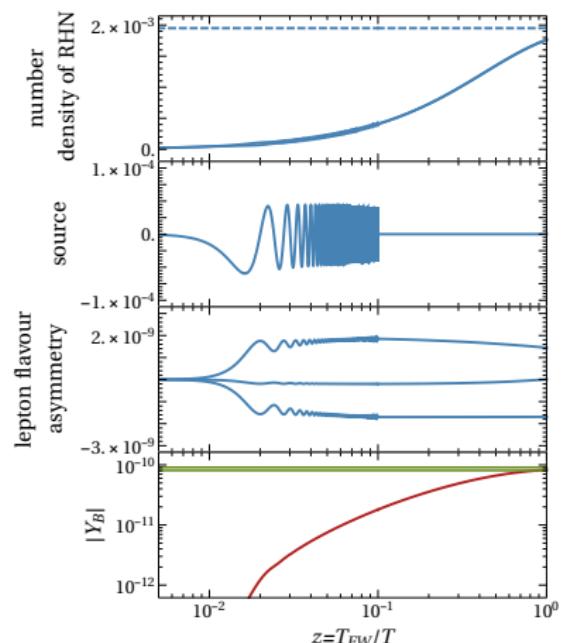
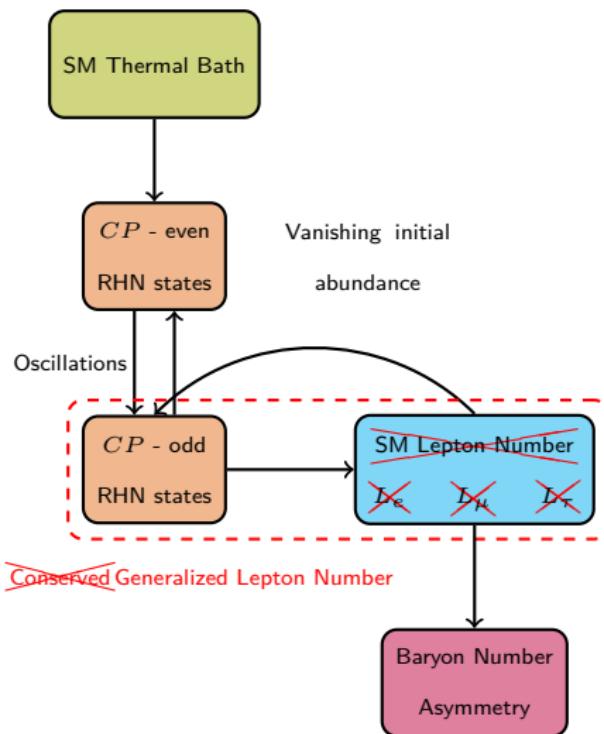
Global constraints: Inverted hierarchy



Global constraints: Inverted hierarchy

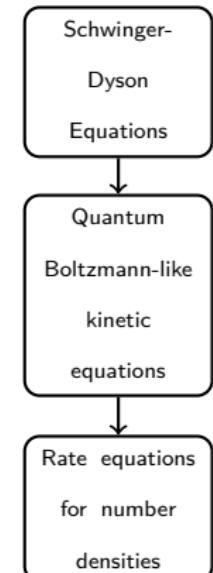


Leptogenesis via Neutrino Oscillations



Goals of this work:

- derivation of the **density matrix** equations from **first principles** [Drewes/Garbrecht/Gueter/JK 1606.06690]
- use the more recent **calculations of rates**
[Anisimov/Besak/Bödeker 1012.3784] [Garbrecht/Glowna/Schwaller 1303.5498]
- inclusion of **spectator effects** [Barbieri/Creminelli/Strumia/Tetradis
hep-ph/9911315] [Garbrecht/Schwaller 1404.2915]
- **analytical approximations** for different regimes
[Drewes/Garbrecht/Gueter/JK 1606.06690]
- explore parameter space/ **phenomenological implications** [Hernández/Kekic/López-Pavón/Racker/Salvado 1606.06719]
[Drewes/Garbrecht/Gueter/JK 1609.09069]
- *violation of generalized lepton number* [Hambye/Teresi
1606.00017] [Eijima/Shaposhnikov 1703.06085] [Laine/Ghiglieri 1703.06087]



Evolution Equations

RHN density matrix

$$\frac{dn}{dz} = -\frac{i}{2} [H, n] - \frac{1}{2} \{ \Gamma, n - n^{\text{eq}} \} - \tilde{\Gamma} q_\ell$$

Active lepton equations

$$\frac{dq_\ell}{dz} = \frac{S_\ell(n)}{T} - W q_\ell + \tilde{W} q_N$$

- Density matrix of the RHN

$$n = \begin{pmatrix} n_{11} & n_{12} \\ n_{21} & n_{22} \end{pmatrix}$$
- Effective Hamiltonian H of the RHN $\sim M^2$
- Production rate $\Gamma \sim Y^2 T$
- Source term S_ℓ of the active neutrinos
- Washout term W

Evolution Equations

RHN density matrix

$$\frac{d\mathbf{n}}{dz} = -\frac{i}{2} [\mathbf{H}, \mathbf{n}] - \frac{1}{2} \{\boldsymbol{\Gamma}, \mathbf{n} - n^{\text{eq}}\} - \tilde{\Gamma} q_\ell$$

Active lepton equations

$$\frac{dq_\ell}{dz} = \frac{S_\ell(n)}{T} - \mathbf{W} q_\ell + \tilde{\mathbf{W}} q_N$$

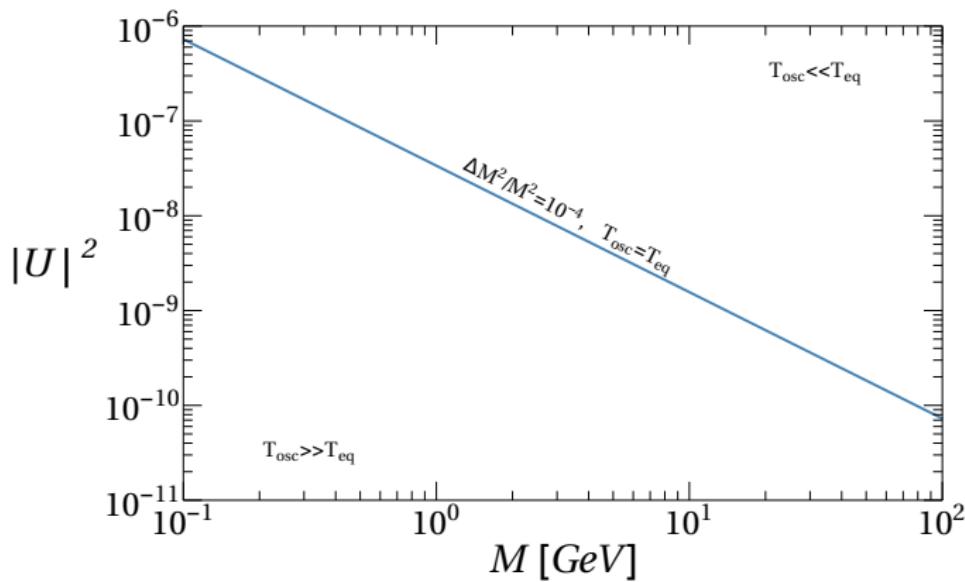
Temperature (time) scales

$$T_{\text{osc}} = \sqrt[3]{T_{\text{com}} (M_{11}^2 - M_{22}^2)}$$

$$T_{\text{eq}} = T_{\text{com}} \gamma_{\text{av}} \text{Tr} (YY^\dagger)$$

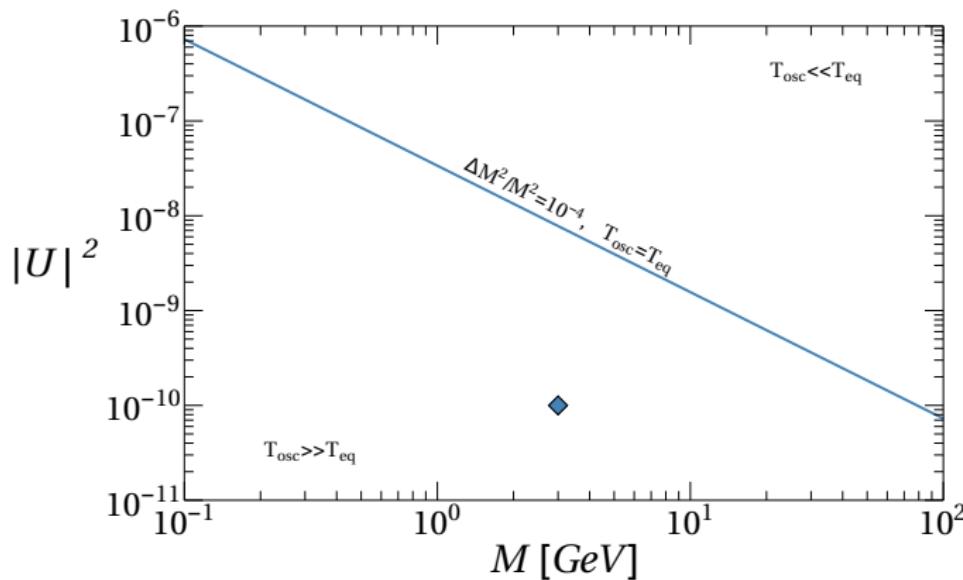
- Possible to solve numerically
- Approximations needed for parameter scans

Temperature scales and observables



[Drewes/Garbrecht/Gueter/JK 1606.06690]

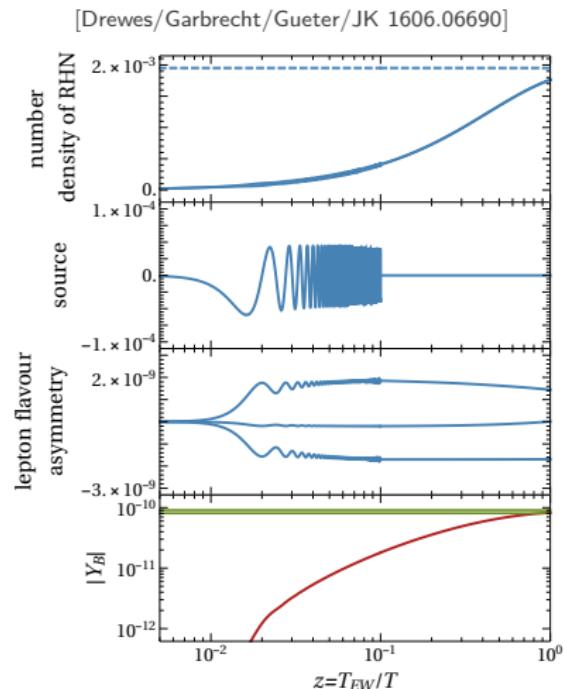
Temperature scales and observables



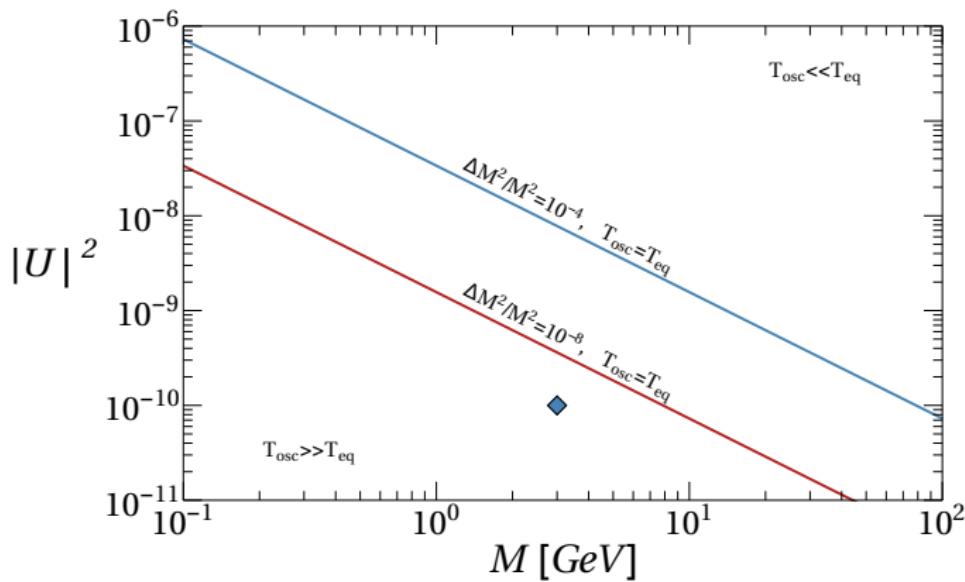
[Drewes/Garbrecht/Gueter/JK 1606.06690]

Oscillatory regime: $T_{\text{osc}} \gg T_{\text{eq}}$ (small mixing angles)

- **oscillations begin long before relaxation to equilibrium**
- almost all lepton flavour asymmetry produced during first few oscillations
- lepton number asymmetry produced only through **flavour asymmetric washout**

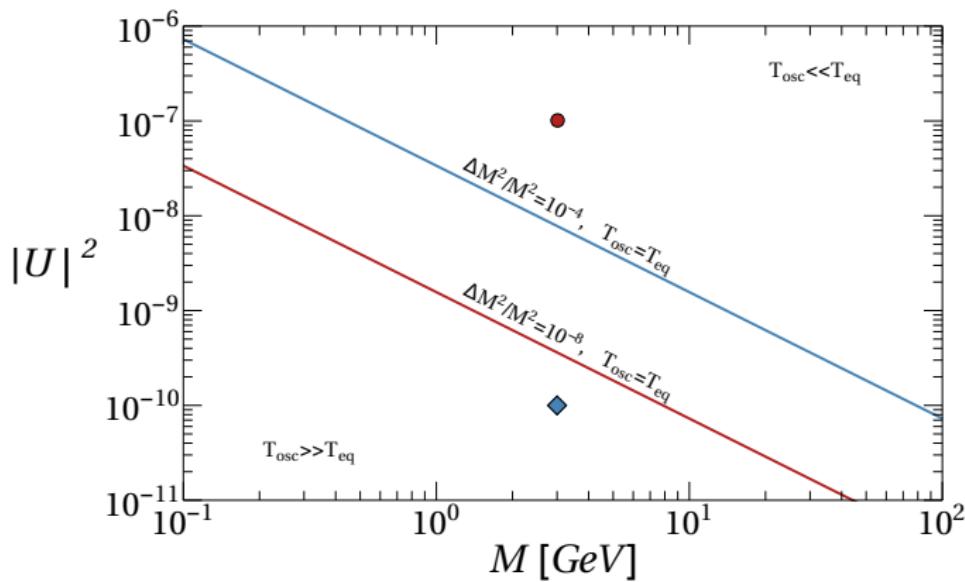


Temperature scales and observables



[Drewes/Garbrecht/Gueter/JK 1606.06690]

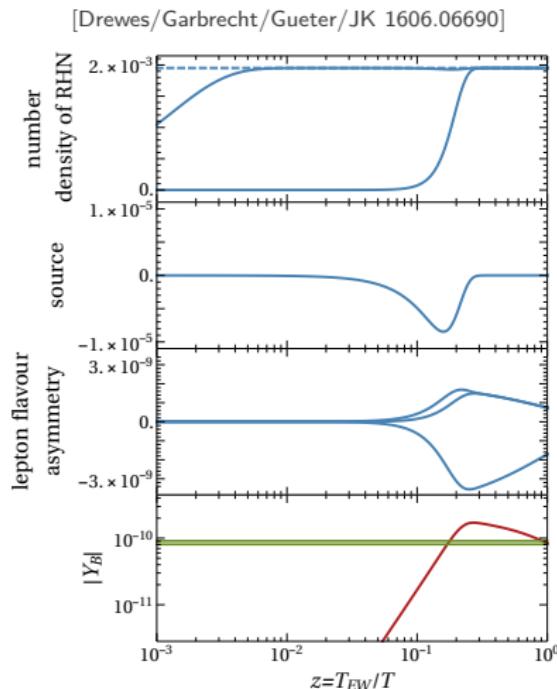
Temperature scales and observables



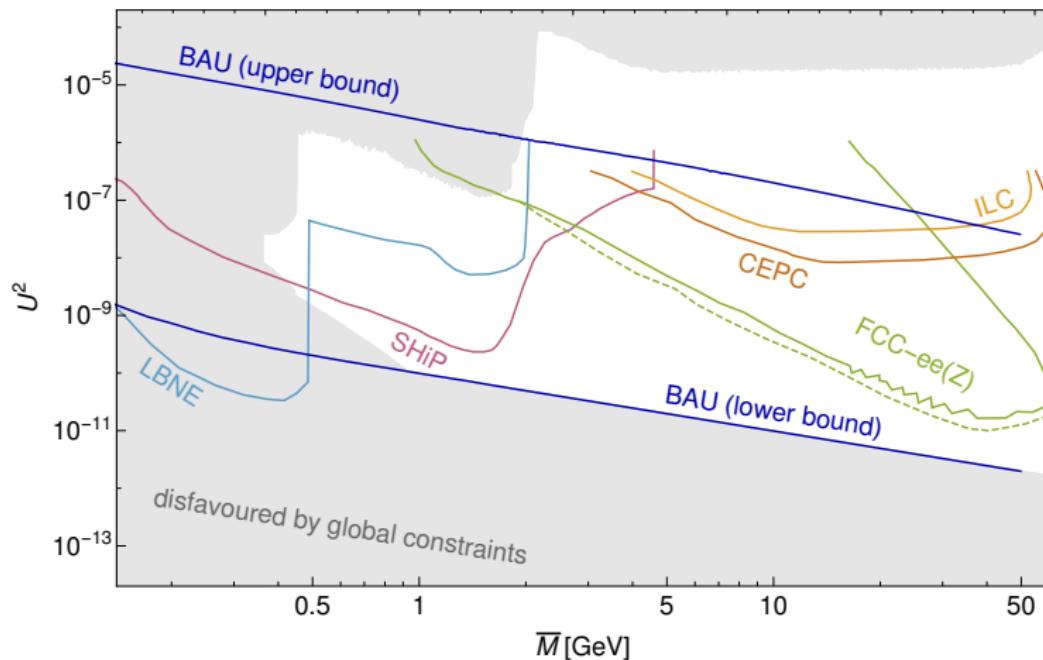
[Drewes/Garbrecht/Gueter/JK 1606.06690]

Overdamped regime: $T_{\text{osc}} \ll T_{\text{eq}}$ (large mixing angles)

- naively for $T_{\text{osc}} < T_{\text{eq}}$, already in equilibrium
- requirement of reproducing the neutrino masses only allows **one interaction eigenstate** to equilibrate
- mixing between interaction eigenstates \rightarrow equilibration
- approximate $B - L$ can postpone the production of BAU, **preventing** too much **washout**

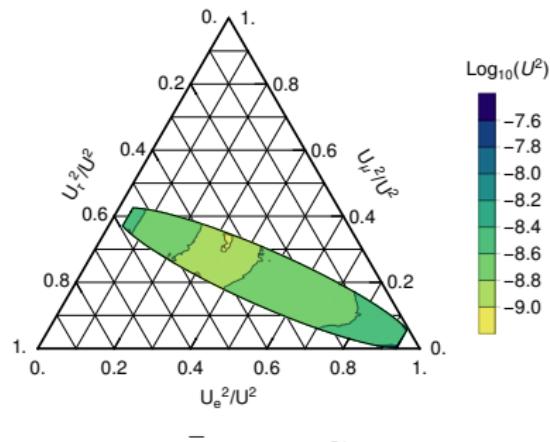


Results: Inverted hierarchy



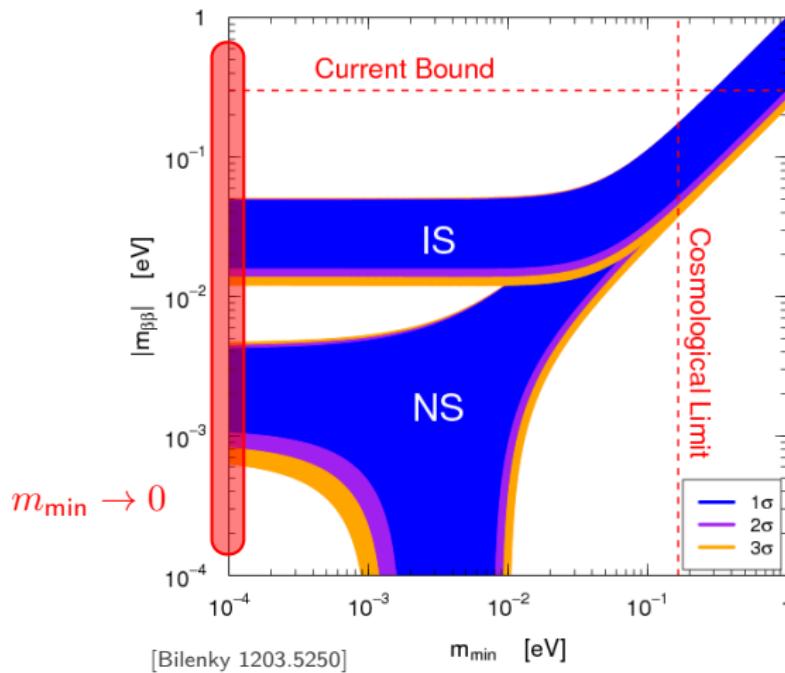
Flavour patterns from leptogenesis: Inverted hierarchy

- large mixing angles require a **flavour asymmetric washout**, which corresponds to a flavour asymmetric mixing
- together with **seesaw constraints** this imposes constraints on the mixing patterns for **large mixing angles**

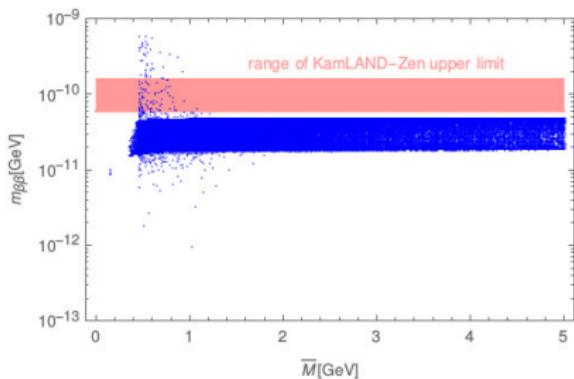


[Drewes/Garbrecht/Gueter/JK 1609.09069]

Leptogenesis and neutrinoless double β decay



Leptogenesis and neutrinoless double β decay



[Eijima/Drewes 1606.06221,
Hernández/Kekic/López-Pavón/Salvado 1606.06719]

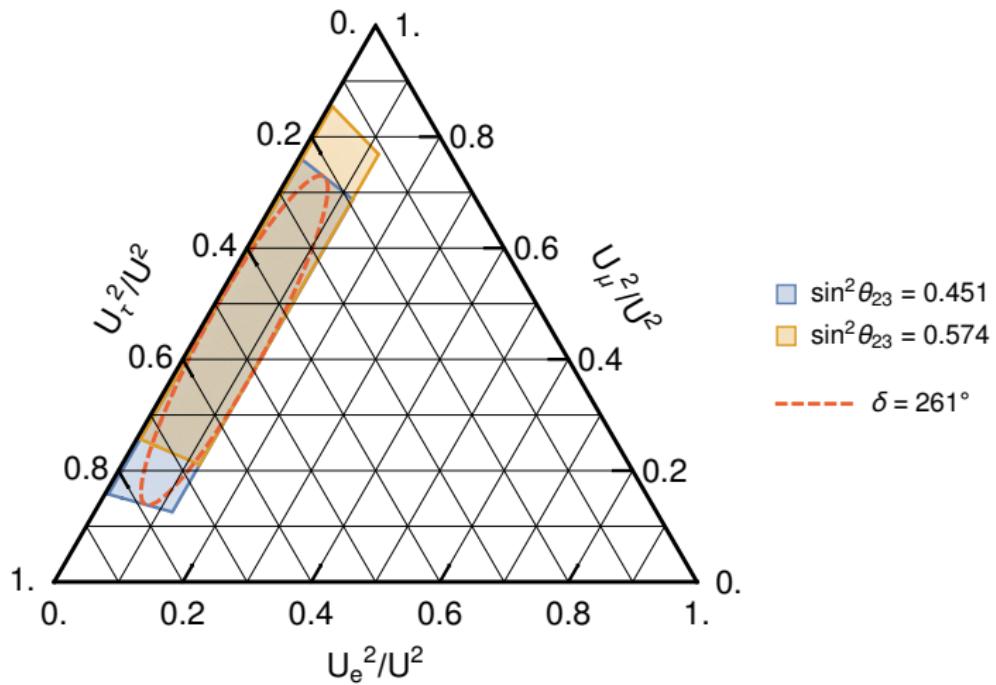
- RHN can contribute to $m_{\beta\beta}$
- large mass splitting is required to have an observable effect (not always compatible with leptogenesis)
- some leptogenesis scenarios can already be excluded by current results

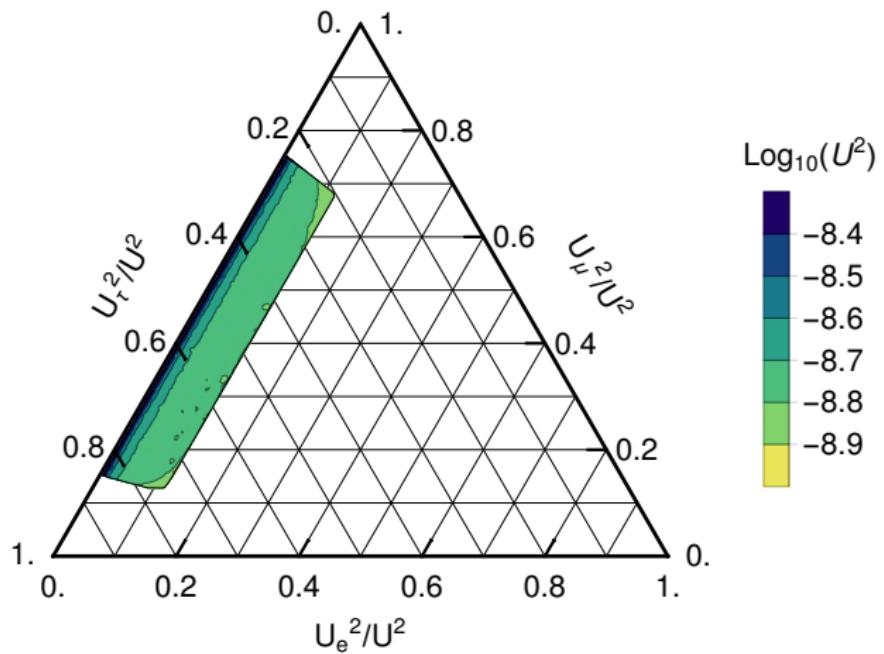
Full Testability?

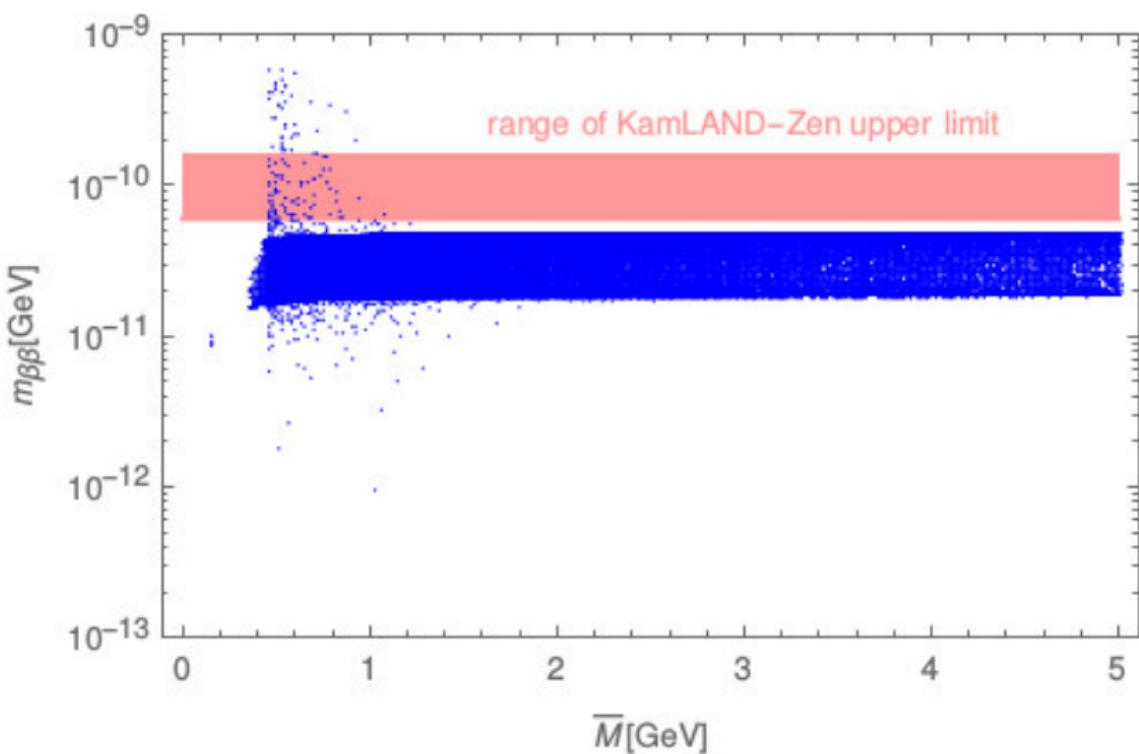
- full testability requires a complete determination of the RHN parameters
- in principle possible from a measurement of all mixing angles and masses
- leptogenesis requires degenerate masses - hard to resolve in experiments
- remaining parameters could be probed by:
 - neutrinoless double β decay requires large ΔM and U^2
[Eijima/Drewes 1606.06221, Hernández/Kekic/López-Pavón/Salvado 1606.06719]
 - CP violation requires ΔM comparable to the decay width
 - lepton number violation requires ΔM comparable to the decay width

Conclusions

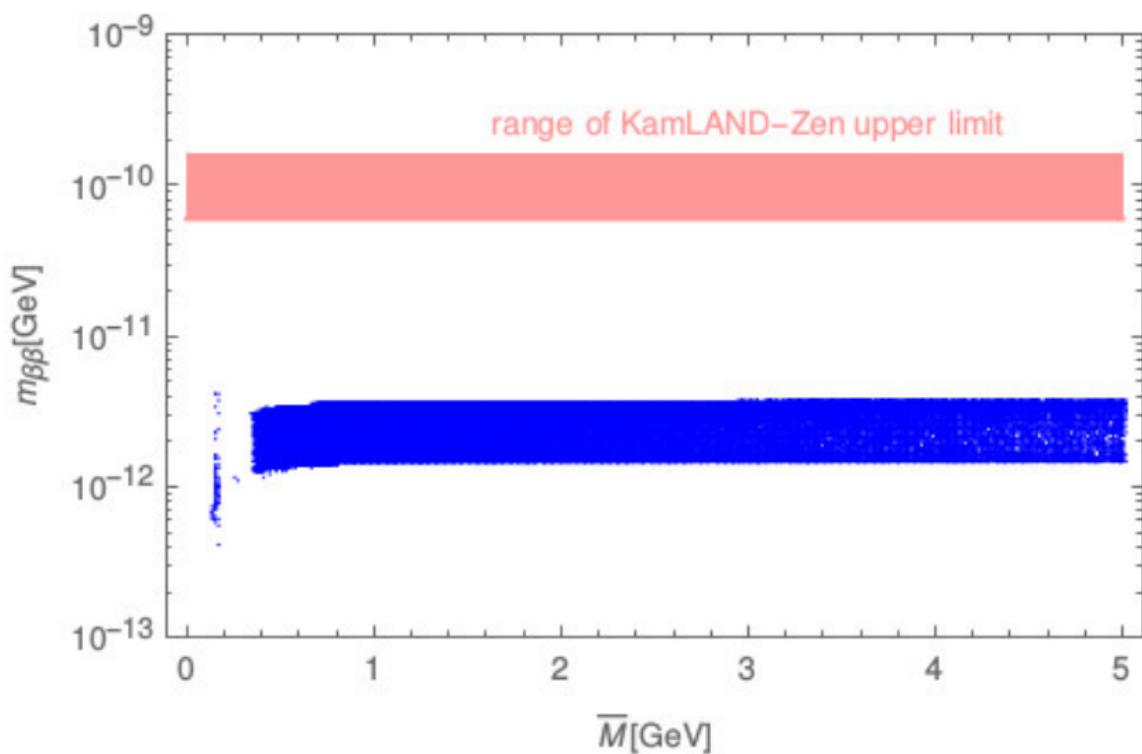
- adding GeV-scale RHNs to the standard model can explain both the observed **neutrino masses** and the **BAU**
- the seesaw mechanism gives constraints on **RHN mixing patterns** (stronger if δ is measured!)
- testable leptogenesis within reach of future experiments (SHiP, LBNE , FCC, ILC, CEPC)
- large mixing angles + leptogenesis → even stronger predictions on the flavour patterns
- while complete determination of the RHN parameters is possible in principle, it requires extreme experimental sensitivity







[Eijima/Drewes 1606.06221]



[Eijima/Drewes 1606.06221]