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Precision mass measurements in the  
context of neutrino-nuclear physics

-----  $0\nu\beta\beta$  decay and  $g_A$  -----

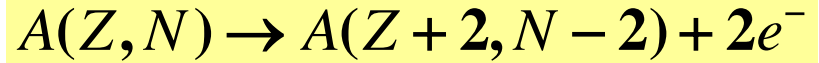
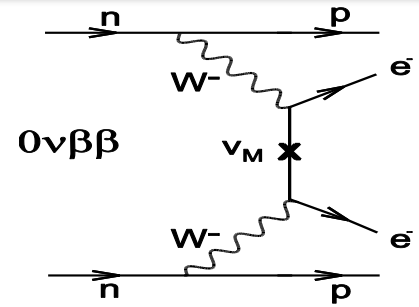
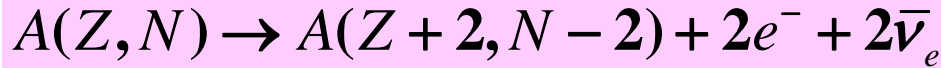
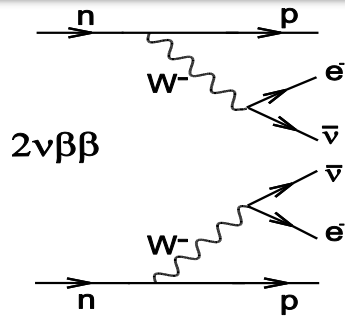
Erice, September, 2017,



# Outline

- Double beta ( $\beta\beta$ ) decay
  - neutrinoless double beta ( $0\nu\beta\beta$ ) decay NME
  - the specialties of  $^{96}\text{Zr}/^{96}\text{Nb}$  for  $\beta$  and  $\beta\beta$  decay
- Mass measurements using the JYFLTRAP ion trap
- Results and the issue of the axial vector coupling  $g_A$
- Proposal for a  $\mu$ -capture experiment at RCNP to “measure”  $g_A$

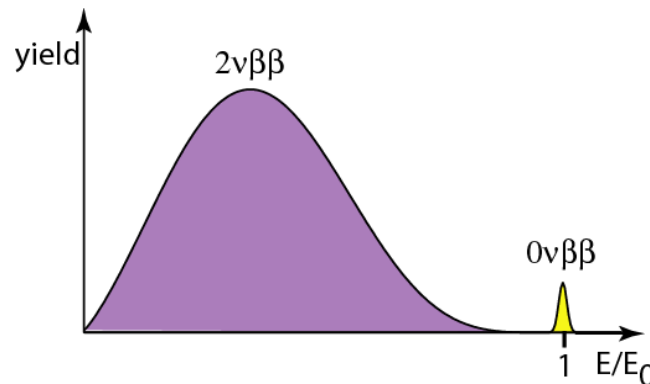
# Double beta decay



- allowed in Standard Model ( $\Delta L=0$ )
- observed experimentally
- NME can be measured (charge-exchange)
- no dependence on  $\nu$  mass
- low- $q$  phenomenon ( $q_{tr} \sim 0.01 \text{ fm}^{-1}$ )

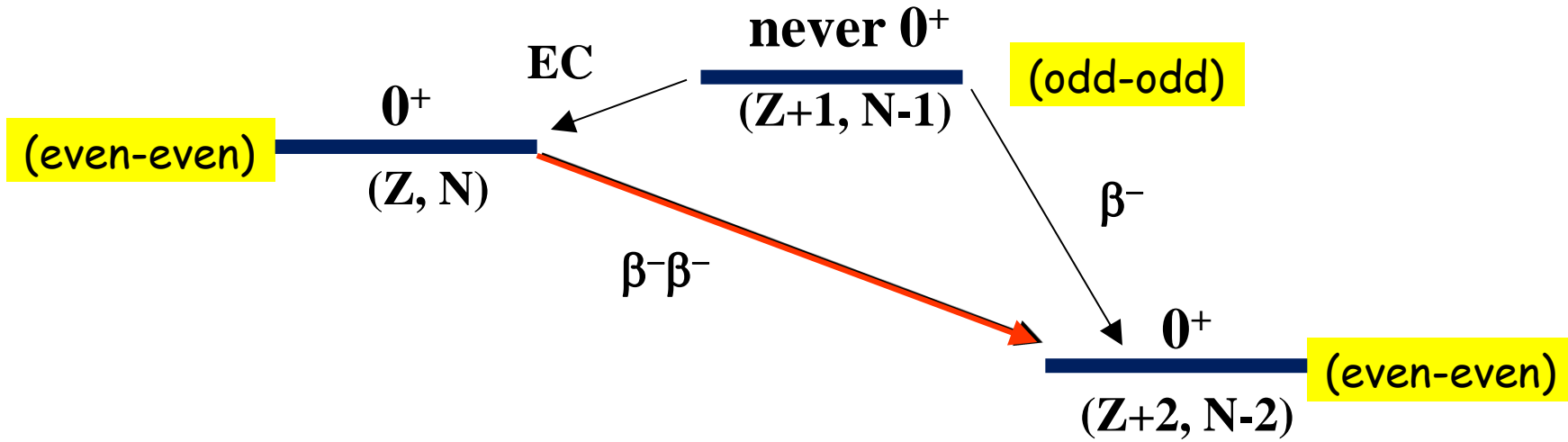
- forbidden in Standard Model ( $\Delta L=2$ )
- not observed yet
- NME only calculated
- $\nu$  has Majorana mass
- high- $q$  phenomenon ( $q_{tr} \sim 0.5 \text{ fm}^{-1}$ )

**$2\nu\beta\beta$ -decay**



**$0\nu\beta\beta$ -decay**

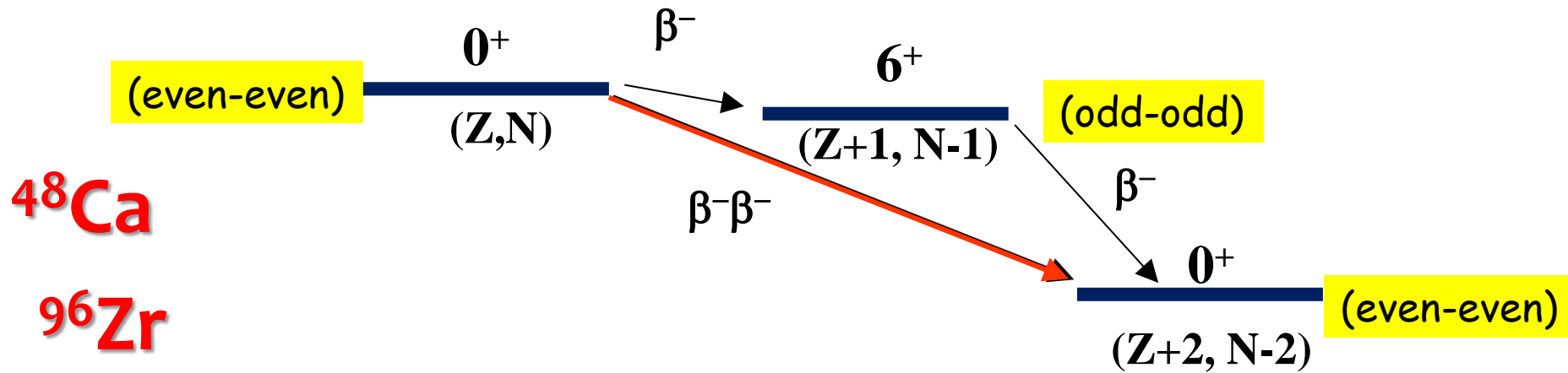
# $\beta\text{-}\beta\text{-}$ decay



1.  $^{48}\text{Ca}$
2.  $^{150}\text{Nd}$
3.  $^{96}\text{Zr}$
4.  $^{100}\text{Mo}$
5.  $^{82}\text{Se}$
6.  $^{116}\text{Cd}$
7.  $^{130}\text{Te}$
8.  $^{136}\text{Xe}$
9.  $^{124}\text{Sn}$
10.  $^{76}\text{Ge}$
11.  $^{110}\text{Pd}$

the  $\beta\text{-}\beta\text{-}$  decay candidates with  
 $Q\text{-value} > 2 \text{ MeV}$

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the  $\beta\text{-}\beta\text{-}$  decay candidates with  
 $Q\text{-value} > 2 \text{ MeV}$

# $\beta\text{-}\beta\text{-}$ decay rate

$2\nu\beta\text{-}\beta\text{-}$  decay:

$$T_{1/2} \approx 10^{19-21} \text{ y}$$

$$\Gamma = G(Q, Z) \times g_A^4 \times \left| NME_{\text{allowed (GT)}} \right|^2$$

5-body

$$\propto Q^{11}$$

from charge  
exchange reactions

$0\nu\beta\text{-}\beta\text{-}$  decay:

$$T_{1/2} > 10^{24} \text{ y}$$

$$\Gamma = G(Q, Z) \times g_A^4 \times \left| NME \right|^2 \times \left| \sum_{i=1}^3 U_{ei}^2 m_i \right|^2$$

3-body

$$\propto Q^5$$

calculated within  
models

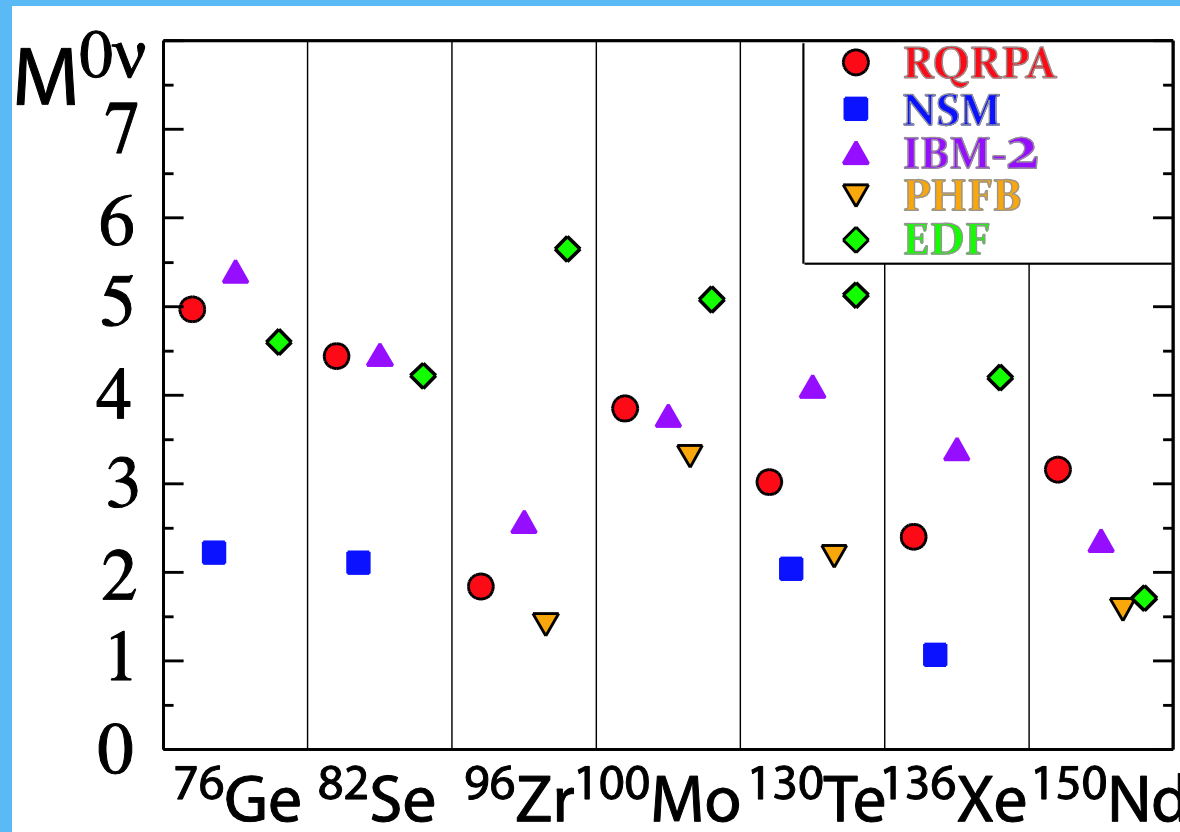
effective Majorana  
 $\nu$  mass,  $m_{\beta\beta}$

**favorable:**

1. high Q-value
2. large Z

# $0\nu\beta\beta^-$ $N_{\text{ucl.}}$ $M_{\text{atrix}}$ $E_{\text{lement}}$

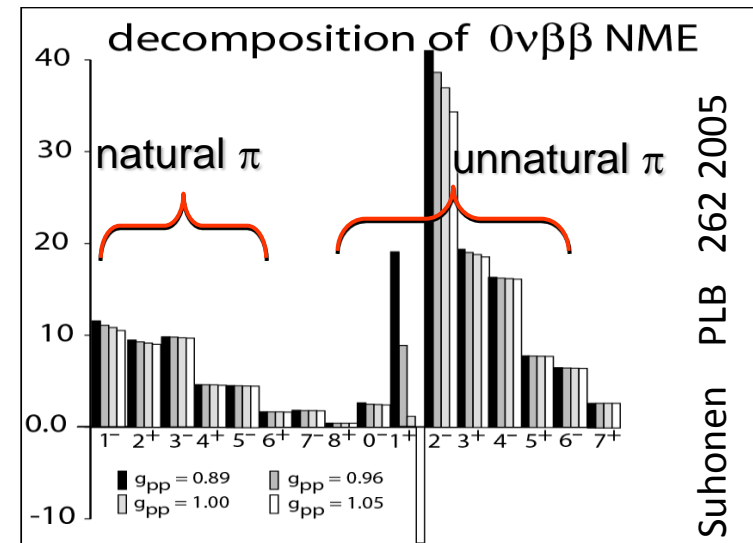
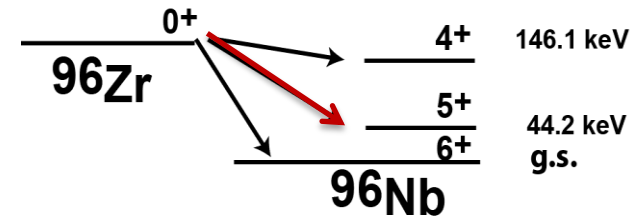
P.Vogel, J. Phys. G, NPP39, 2012



The calculated  $0\nu\beta\beta$  decay NME via different models differ by more than a factor of 2-3 ( i.e. half-life 4-9)

# Idea for $^{96}\text{Zr}$

- (i) measure **Q-value** for  $^{96}\text{Zr} \rightarrow ^{96}\text{Nb}$  **single  $\beta$ -decay** by precision mass measurement and
- (ii) measure the **single  $\beta$ -decay** rate
- (iii)  $\rightarrow$  ft-value
- determine the  $^{96}\text{Zr}$  **4-fold forbidden  $\beta$ -decay NME** and confront with theory
- confront with same theories aimed at calculating  **$0\nu\beta\beta$ -decay NME for the same nucleus!!**





# Competition between $\beta$ & $\beta\beta$ decay of $^{96}\text{Zr}$

two conflicting half-lives:

NEMO-3:  $T_{1/2}^{2\nu\beta\beta} = (2.3 \pm 0.2) \times 10^{19} \text{ y}$   
 geo-chem:  $T_{1/2}^{\beta} = (0.94 \pm 0.32) \times 10^{19} \text{ y}$  1

can this difference be reconciled ?  
 yes, if single  $\beta$  competes with  $\beta\beta$  decay

$$(T_{1/2})^{-1} = (T_{1/2}^{2\nu\beta\beta})^{-1} + (T_{1/2}^{\beta})^{-1}$$

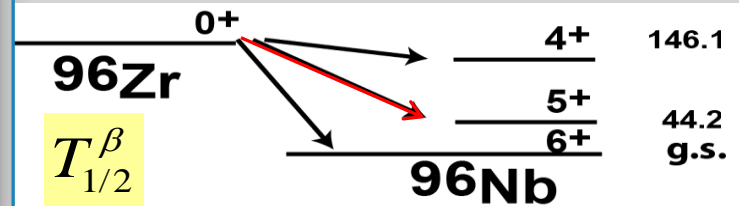
expected  $T_{1/2}^{\beta} = (1.6 \pm 0.9) \times 10^{19} \text{ y}$

experiment  $T_{1/2}^{\beta} > 2.6 \times 10^{19} \text{ y}$  2

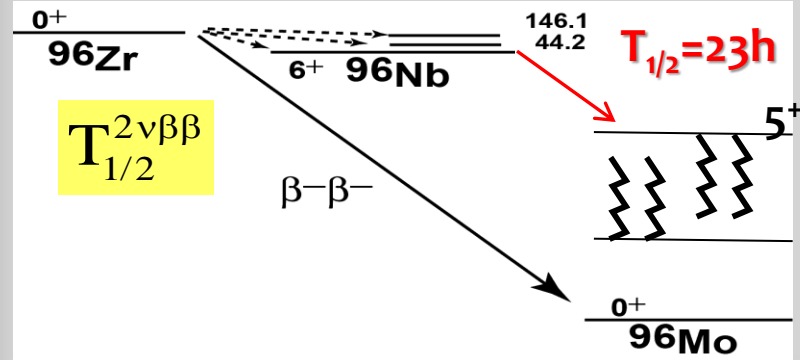
pred. (QRPA)  $T_{1/2}^{\beta} = 24 \times 10^{19} \text{ y}$  3

**BUT**

$$(T_{1/2}^{\beta})^{-1} \propto o(Q^{13}) g_A^2 \langle M_{\beta}^{4u} \rangle^2$$



$0^+ \rightarrow 6^+$  6-fold non-unique (unobservably long)  
 $0^+ \rightarrow 5^+$  4-fold unique (possible)  
 $0^+ \rightarrow 4^+$  4-fold non-unique (no phase space)

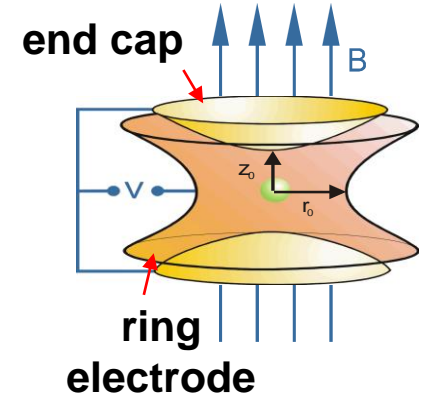
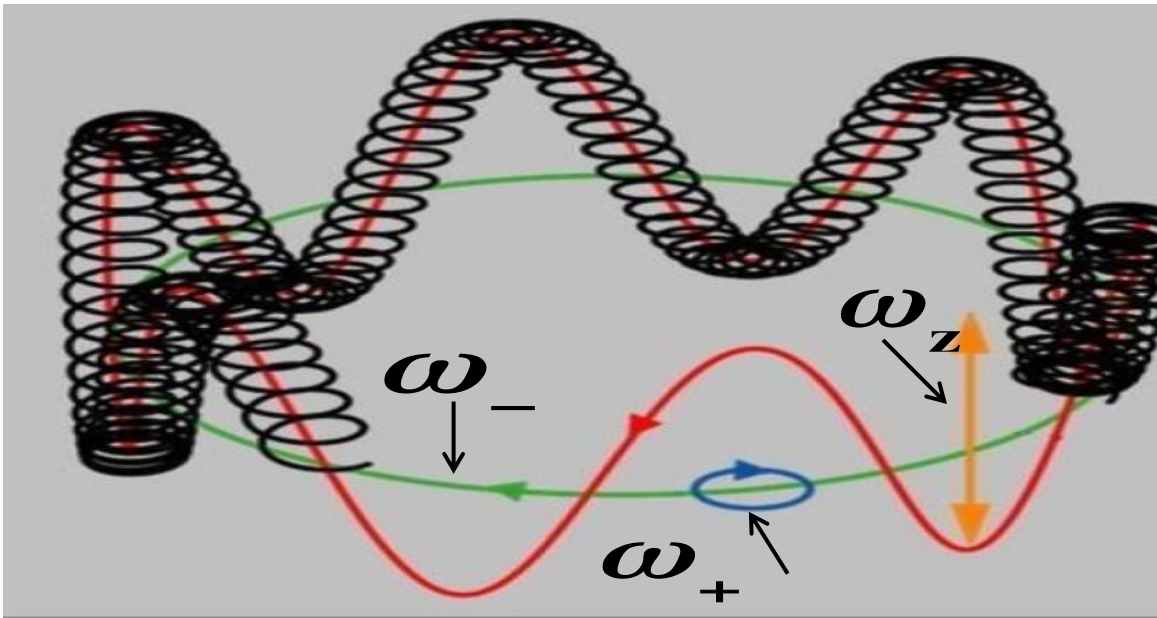


**Q-value**  $\rightarrow M_{\beta}^{4u} \rightarrow (T_{1/2}^{0\nu\beta\beta})^{-1} \propto Q^5 |M_{\beta\beta}^{0\nu}|^2 \langle m_{\beta\beta} \rangle^2$

# Ion motion in a Penning trap

Homogeneous magnetic field + static electric field provides 3D confinement  
results in three eigenmotions:

1. Magnetron motion  $\omega_-$
2. Reduced cyclotron motion  $\omega_+$
3. Axial motion  $\omega_z$



$$\omega_c = \omega_- + \omega_+$$

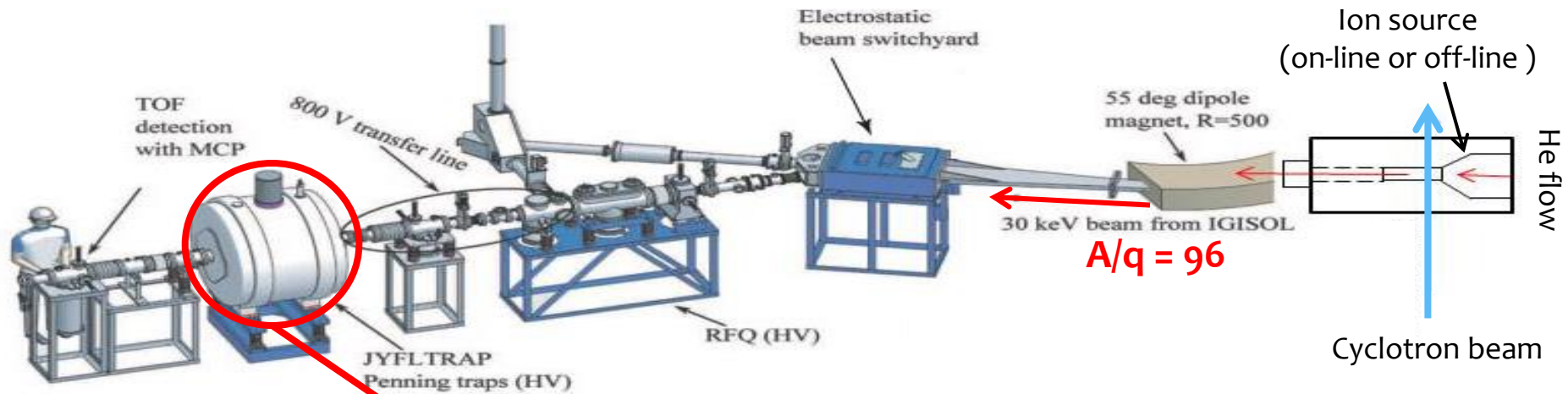
$$\omega_- \approx 1 \text{ kHz},$$

$$\omega_+ \approx 1 \text{ MHz}$$

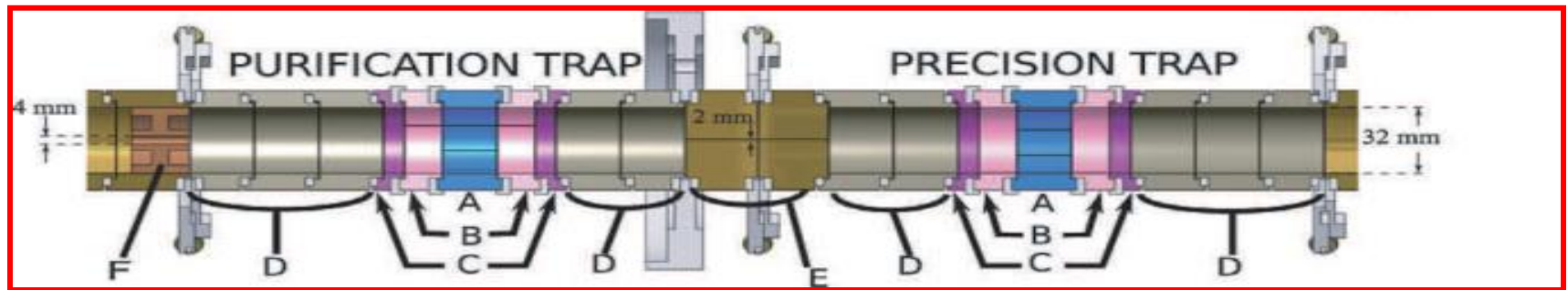
$$\omega_c = \frac{q}{m} \cdot B$$

# IGISOL / JYFLTRAP setup

Figures from Eronen EPJA 48-2012



$$\Delta B/B = 8.18 \times 10^{-12} / \text{min}$$



purification & isobar separation  
by buffer-gas cooling technique

mass measurement via cyclotron  
frequency,  $p < 10^{-7}$  mbar.

# Beam production at IGISOL

Off-line measurements

$^{96}\text{Zr}$  and  $^{96}\text{Mo}$



off-line ion source

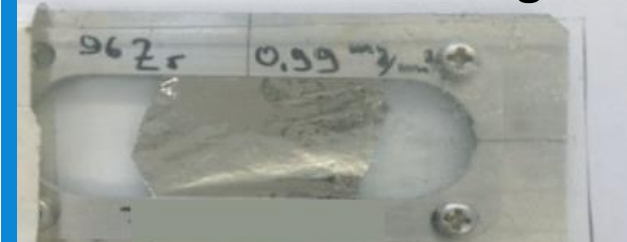
On-line measurements

$^{96}\text{Zr}$  (p, n) $^{96}\text{Nb}$  reaction  
for production of  $^{96}\text{Nb}$



10 MeV proton  
beam

57% enriched  $^{96}\text{Zr}$  target



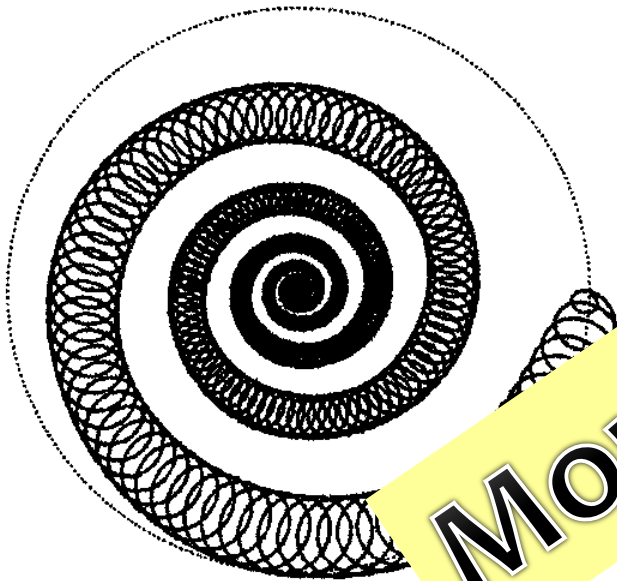
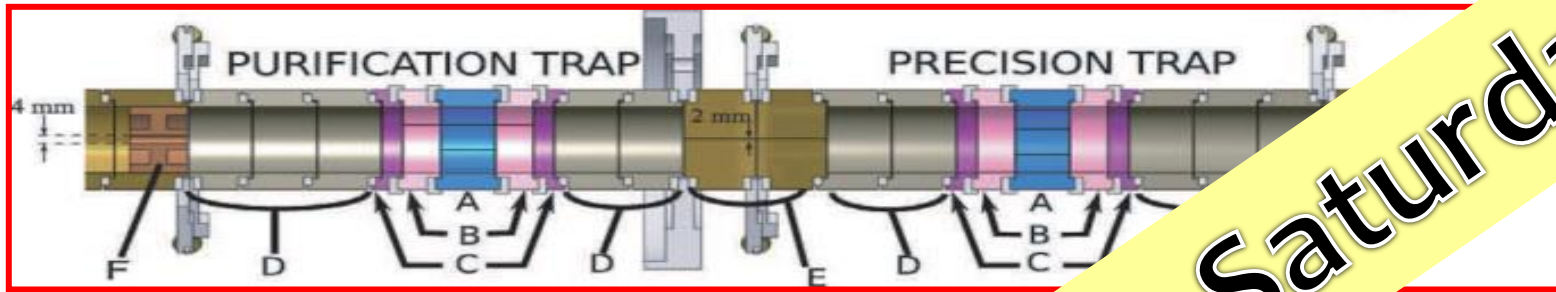
target (on-line) ion source

**Note** IGISOL produces on-line :  
 $^{96}\text{Zr}$  ----- from target material  
 $^{96}\text{Nb}$  ----- (p,n) charge-ex reaction  
 $^{96}\text{Mo}$ ----- from havar<sup>®</sup> separator foil

mother & daughter ions  
are produced at  
IGISOL at the same  
time

# Principle of isobar separation in purification trap

Excite ion eigenmotion by a dipolar or a quadrupolar electric field with a corresponding frequency ( $\nu_-$  or  $\nu_c$ )



More T. Eronen, Saturday

excite ion eigenmotion with  $\nu_c$  causes  
excitation of this species in the  
trap, the buffer gas  
trap:  $R > 500/1$  )

# Mass measurements in a Penning trap

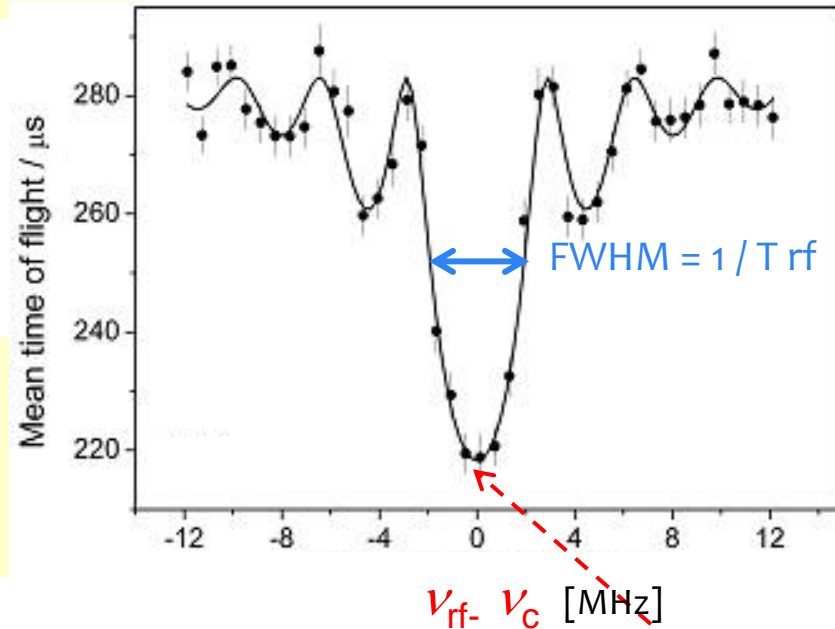
- Performing precision mass determination via **cyclotron frequency**  $\nu_c$  measurement

$$\nu_c = \frac{1}{2\pi} \frac{q}{m} \cdot B$$

$$\omega_c = 2\pi \nu_c$$

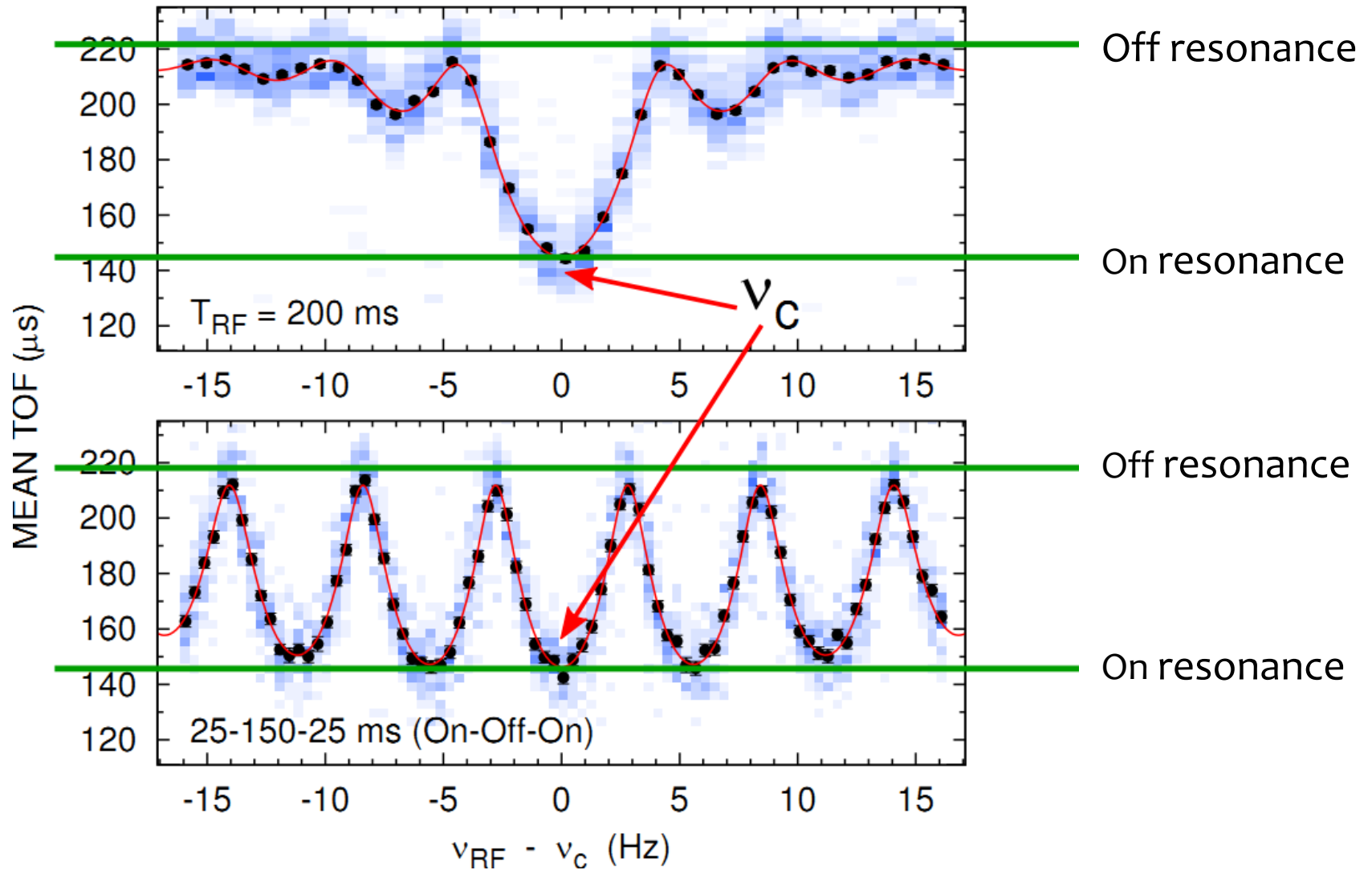
- Cyclotron frequency  $\nu_c$  determination done by **TOF-ICR** technique

- Frequency ratio  $r = \frac{n_{c \text{ daughter}}}{n_{c \text{ mother}}}$

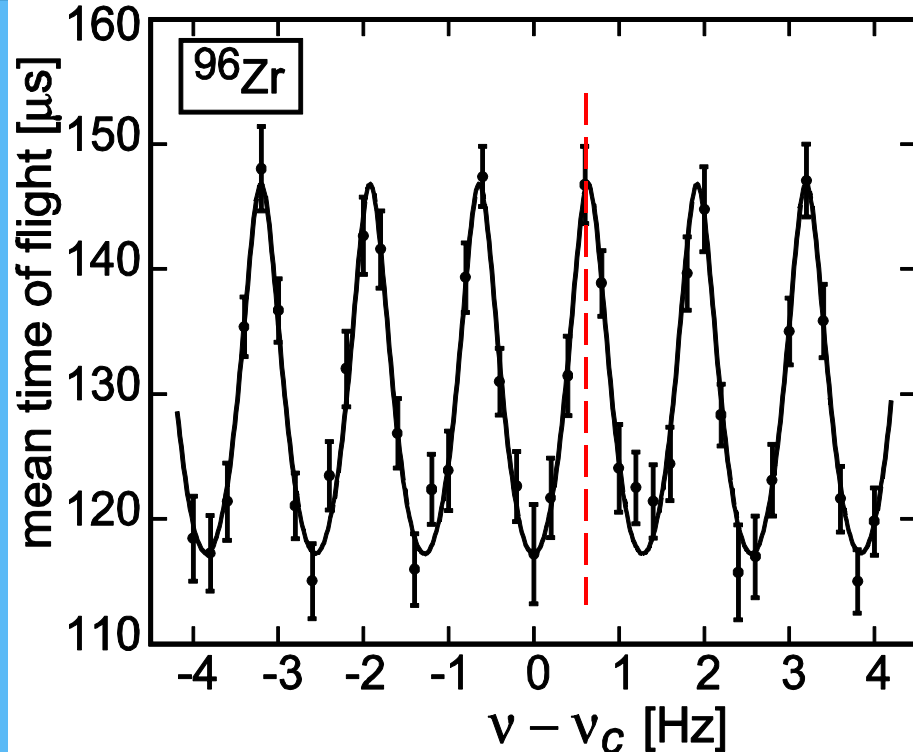


at  $\nu_c$  ion acquires extra energy  
→ shorter time of flight

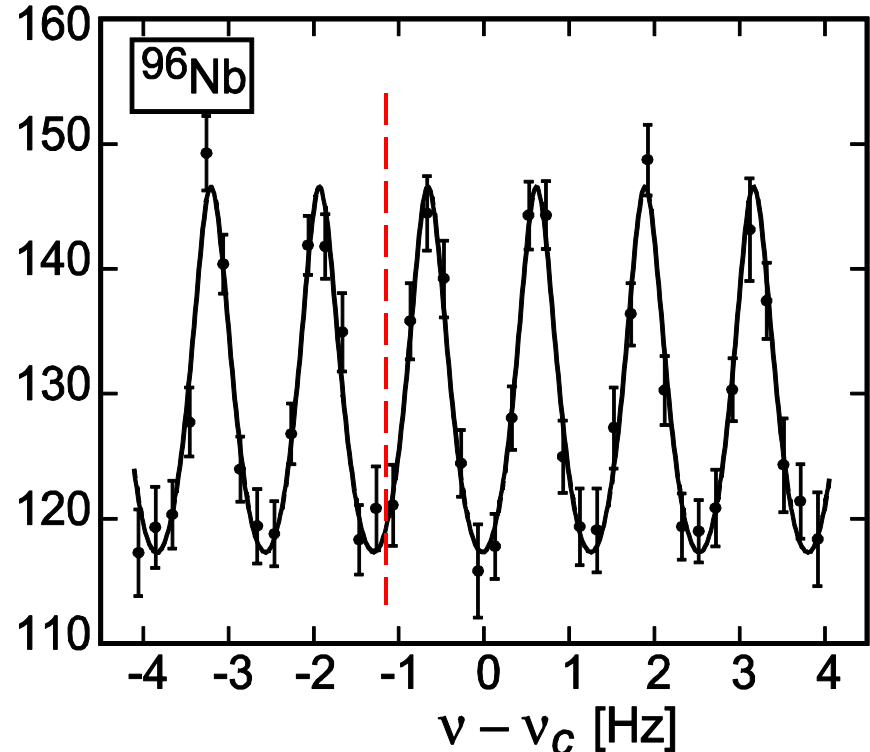
# Example resonances



# Cyclotron frequency results



$\nu_C = 1120477.608$  Hz



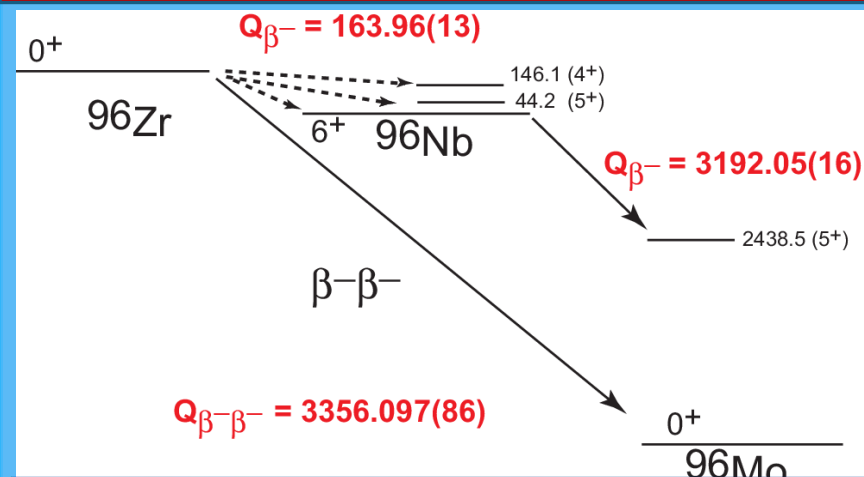
$\nu_C = 1120479.664$  Hz

The measured cyclotron frequency of  $^{96}\text{Zr}$  and  $^{96}\text{Nb}$  by

Ramsey excitation pattern 25-750-25 ms (on-off-on)



# Q-value results



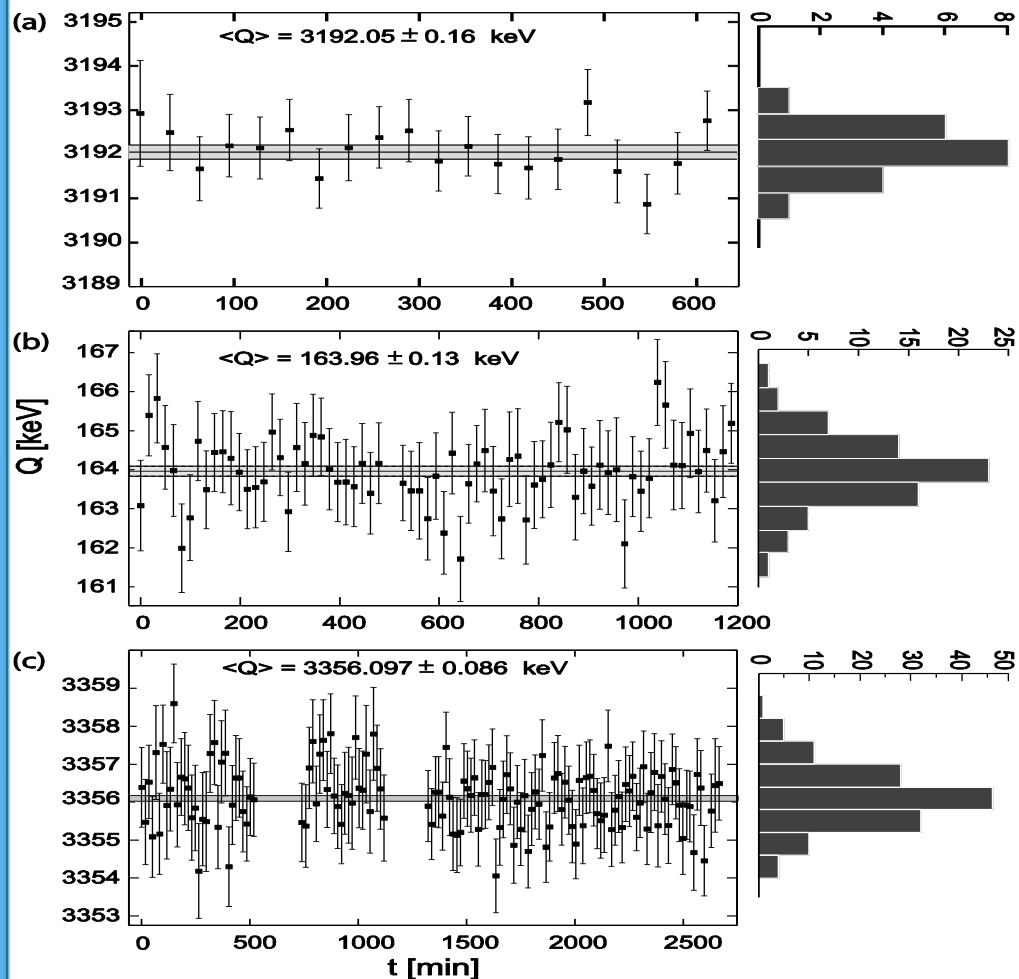
## 96Zr

$$Q_{\beta\beta} = 3356.097 \pm 0.086 \text{ keV}$$

(7.1 keV higher than AME2012)

$$Q_{\beta} = 163.96 \pm 0.13 \text{ keV}$$

$$Q_{\beta} = 3192.05 \pm 0.16 \text{ keV}$$



The measured single  $\beta^-$  and  $\beta\beta^-$  decay Q-values of the A=96 triplet

# Next: need $T_{1/2}$ of single $\beta$ decay

$$T_{1/2}(\text{QRPA}) = \frac{24}{g_A^2} \times 10^{19} \text{ yr}$$

$$T_{1/2}(\text{SM}) = \frac{11}{g_A^2} \times 10^{19} \text{ yr}$$

$$T_{1/2}(\text{exp}) > 2.3 \times 10^{19} \text{ yr}$$

## Important side effect:

single  $\beta$  decay depends on  $g_A^2$

$2\nu/0\nu\beta\beta$  decay depends on  $g_A^4$

A measurements of single  $\beta$  decay gives experimental handle on the quenching of  $g_A$

# The muon capture and $g_A$ in weak decays

## Title

Exclusive  $\mu$ -capture on  $^{24}\text{Mg}$ ,  $^{32}\text{S}$  and  $^{56}\text{Fe}$  populating low-lying  $1^+$  states to probe the weak axial current at high momentum transfer

M. Alanssari,<sup>1</sup> I. H. Hashim,<sup>2</sup> L. Jokiniemi,<sup>3</sup> H. Ejiri,<sup>4</sup>  
E. Ideguchi,<sup>4</sup> A. Sato,<sup>4</sup> J. Suhonen,<sup>3</sup> and D. Frekers<sup>1</sup>

<sup>1</sup>Institut für Kernphysik, Westfälische Wilhelms-Universität, D-48149 Münster, Germany

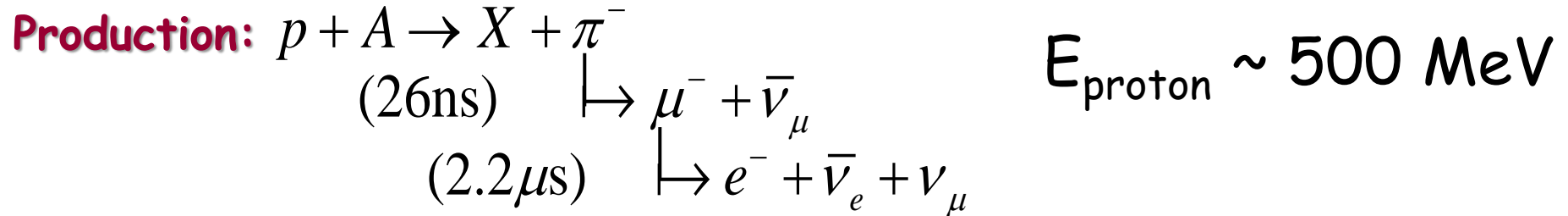
<sup>2</sup>Universiti Teknologi Malaysia, 81310 Johor Bahru, Malaysia

<sup>3</sup>University of Jyväskylä, Department of Physics, FI-40014, Finland

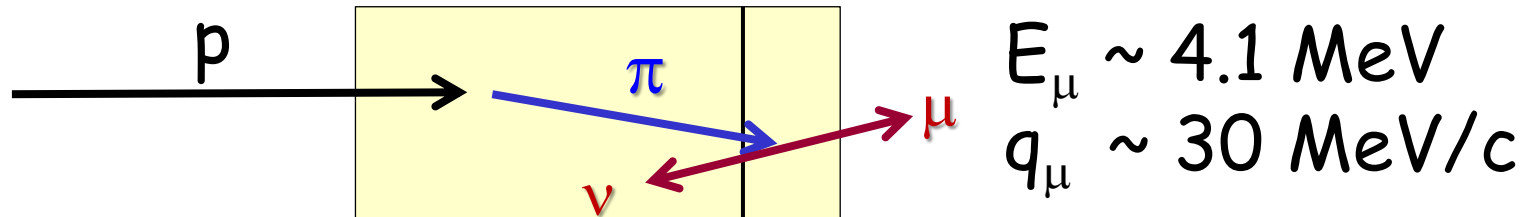
<sup>4</sup>Research Center for Nuclear Physics, Osaka University, Ibaraki, Osaka 567-0047, Japan

# The amazing muon

There has not been any other elementary particle so „successful“ in advancing our knowledge in so many different areas of physics.



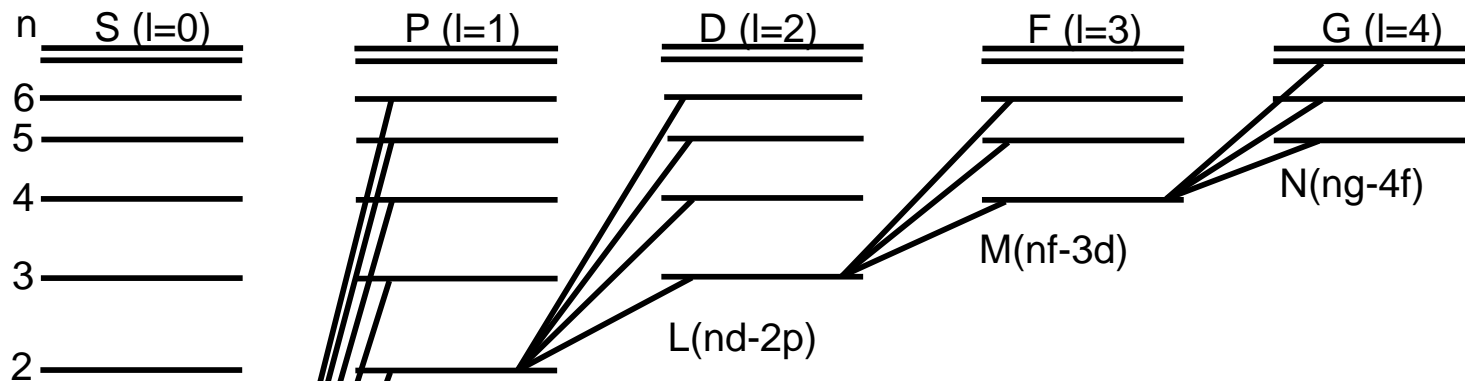
**Surface muons:** produced from stopped (usually negative) pions at end of target



**Life-time:**  $\Gamma_\mu = \frac{1}{\tau_\mu} = \frac{G_F^2 m_\mu^5}{192\pi^3} (1 - \varepsilon) = (2,196981(2) \mu\text{sec})^{-1} \quad (\varepsilon \approx 10^{-3})$

$G_F = 1,16637(2) \cdot 10^{-5} \text{ GeV}^{-2} \quad m_\mu = 105.6583745(24) \text{ MeV}$

**Muonium:**  $\mu^+ + e^-$  an exotic hydrogen  $I \sim 13.6 \text{ eV}$



neutrino

K( $np-1s$ )

prompt Lyman  $\alpha$ -series of atomic  $\mu$ -capture followed by delayed nucl. capture

$$\lambda_{cap} = \lambda_{total} - Q\lambda_{decay}$$

Huff-factor

$$\lambda_{decay} = (2.2\mu s)^{-1} = 4.54 \cdot 10^5 s^{-1}$$

$$\lambda_{total} \sim (10 - 1000 ns)^{-1} \sim 10^8 - 10^6 s^{-1}$$

$$Q \sim 0.9 - 1.0$$

- the neutrino takes most of the energy
- $E_x(\text{nucl}) < 10 - 20 \text{ MeV}$

captured by the nucleus

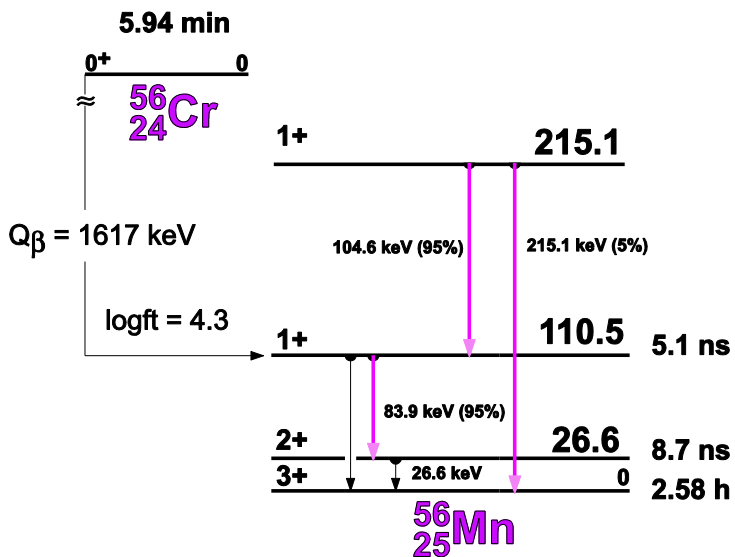
# Motivations

- $\mu$ -cap features momentum transfers similar to  $0\nu\beta\beta$  decay ( $q_{tr} \sim 0.5\text{fm}^{-1} \sim 100\text{MeV}/c$ )
- $\mu$ -cap processes to  $1^+$  states in  $A(\mu^-, \nu)B$  may be compared with charge-ex reactions of (n,p) type.
- $\mu$ -cap may give access to  $g_A$  quenching issue

## However

- only the  $0\nu$ -channel ( $\sim 5\%$ ) is most relevant for  $0\nu\beta\beta$  decay
- level scheme of final odd-odd nucleus is extremely!! poorly known

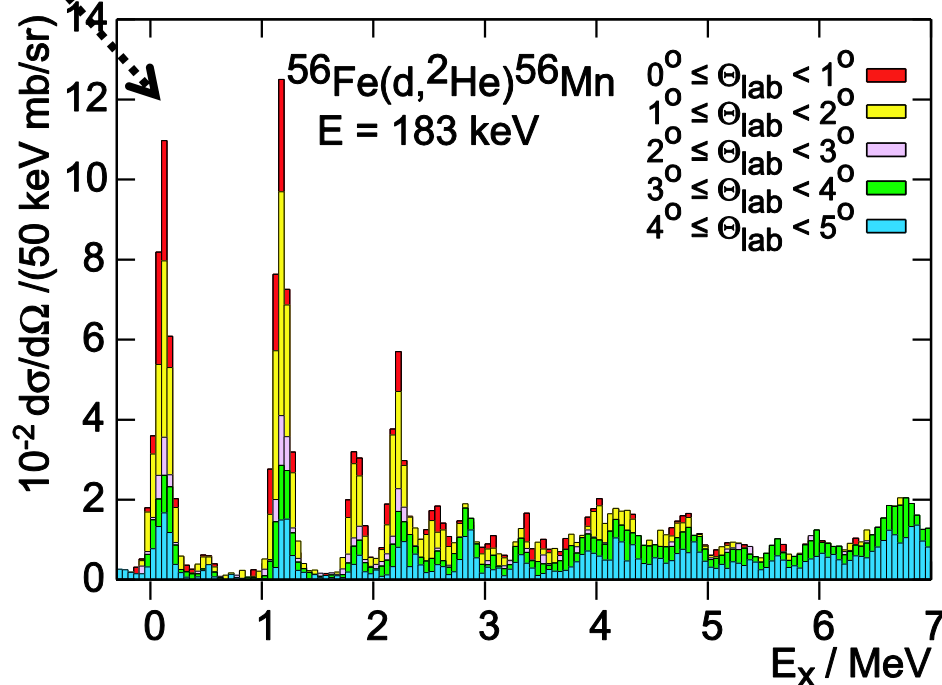
# The issue of $g_A$ queching



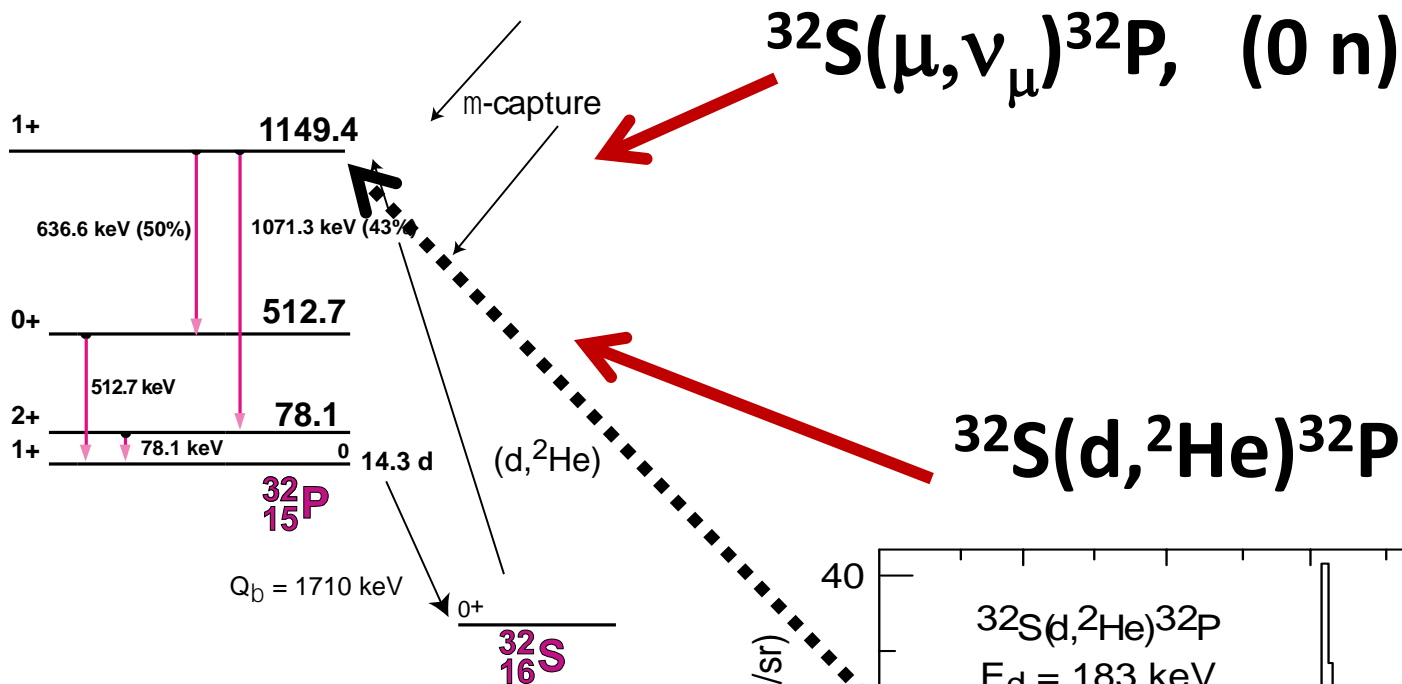
$^{56}\text{Fe}(\mu, \nu_\mu)^{56}\text{Mn}, (0 \text{ n})$

$^{56}\text{Fe}(d, ^2\text{He})^{56}\text{Mn}$

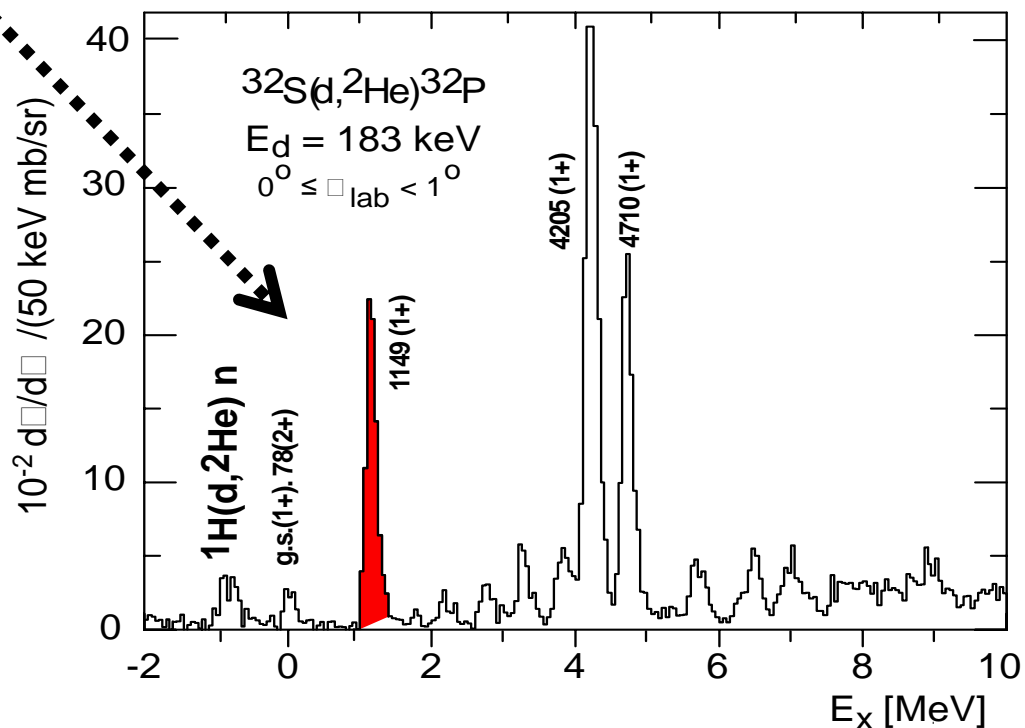
**Example:**  
 Compare transition strength in  $\mu$ -cap and  $(d, ^2\text{He})$  charge-ex



# The issue of $g_A$ queching

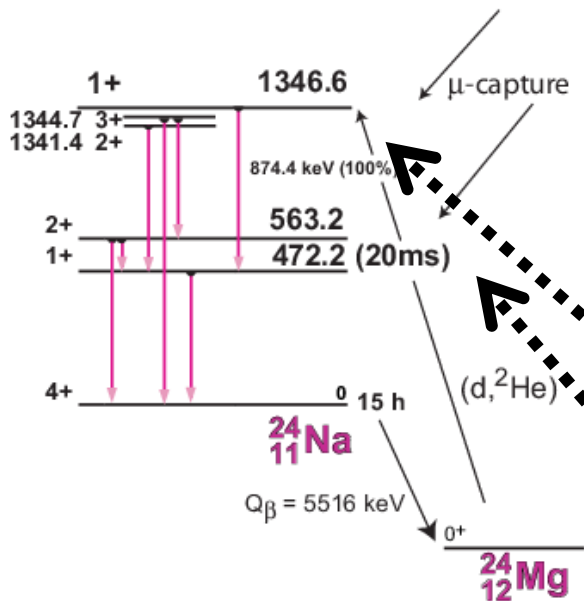


**Example:**  
 Compare transition strength in  $\mu$ -cap and  $(d, ^2\text{He})$  charge-ex





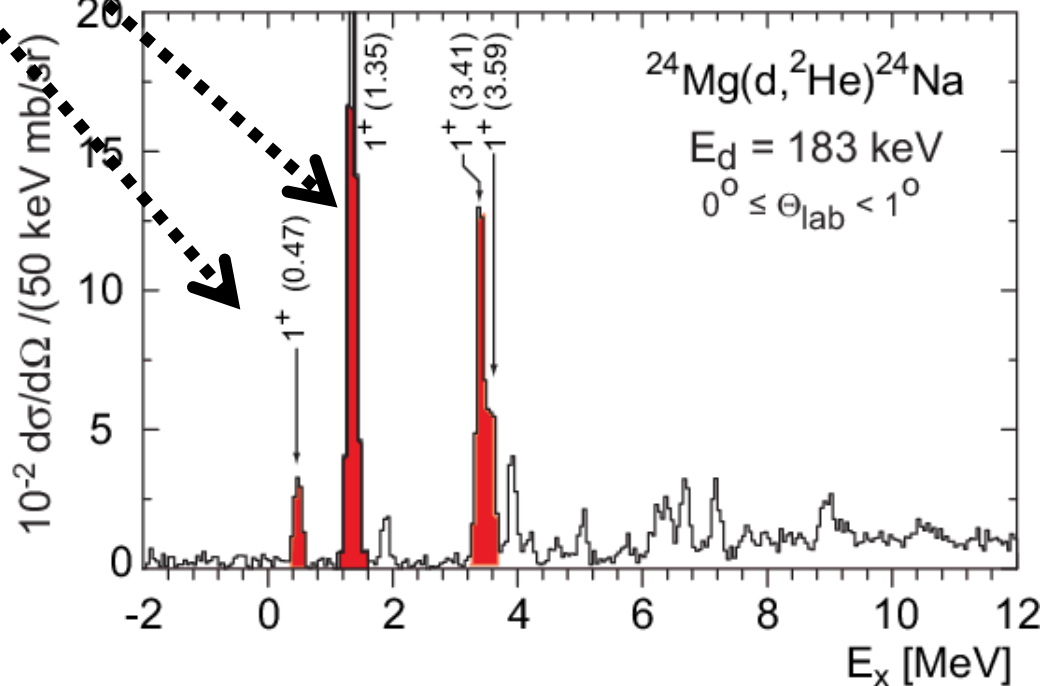
# The issue of $g_A$ queching



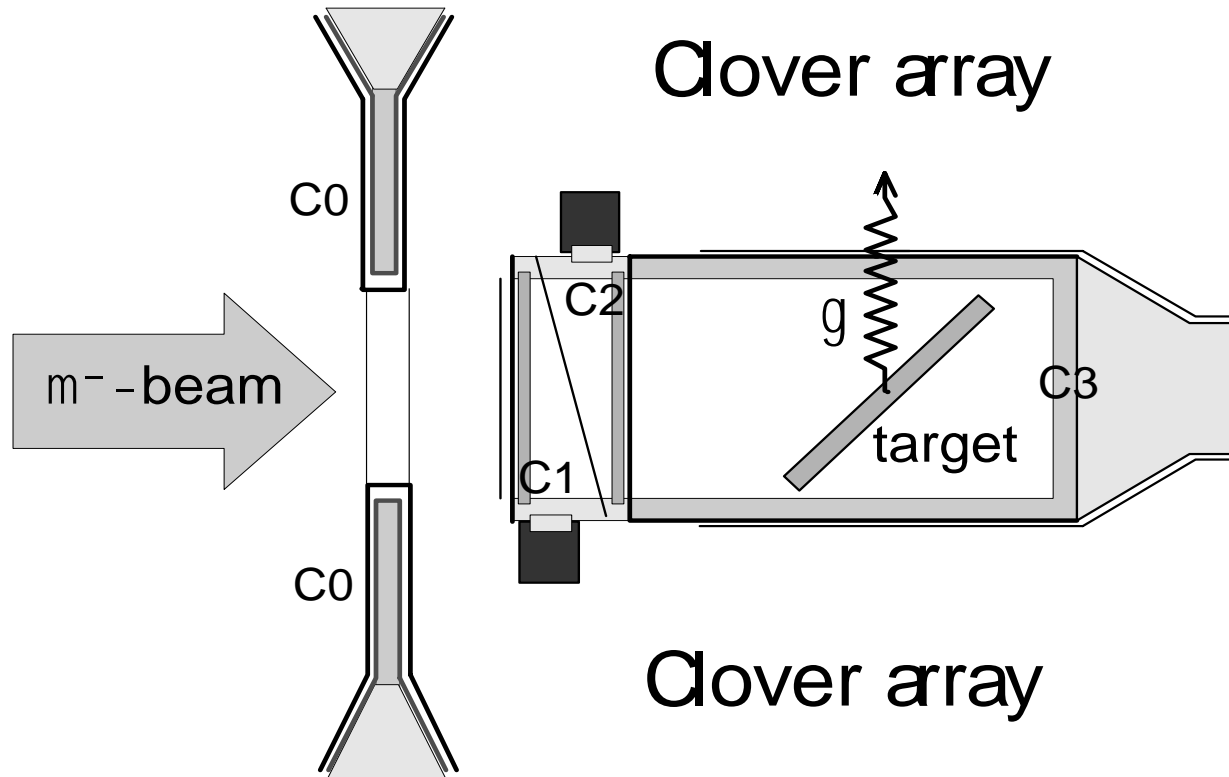
$^{24}\text{Mg}(\mu, \nu_\mu)^{24}\text{Na}$ , (0 n)

$^{24}\text{Mg}(d, ^2\text{He})^{24}\text{Na}$

**Example:**  
Compare transition strength in  $\mu$ -cap and  $(d, ^2\text{He})$  charge-ex



# Schematics of set-up

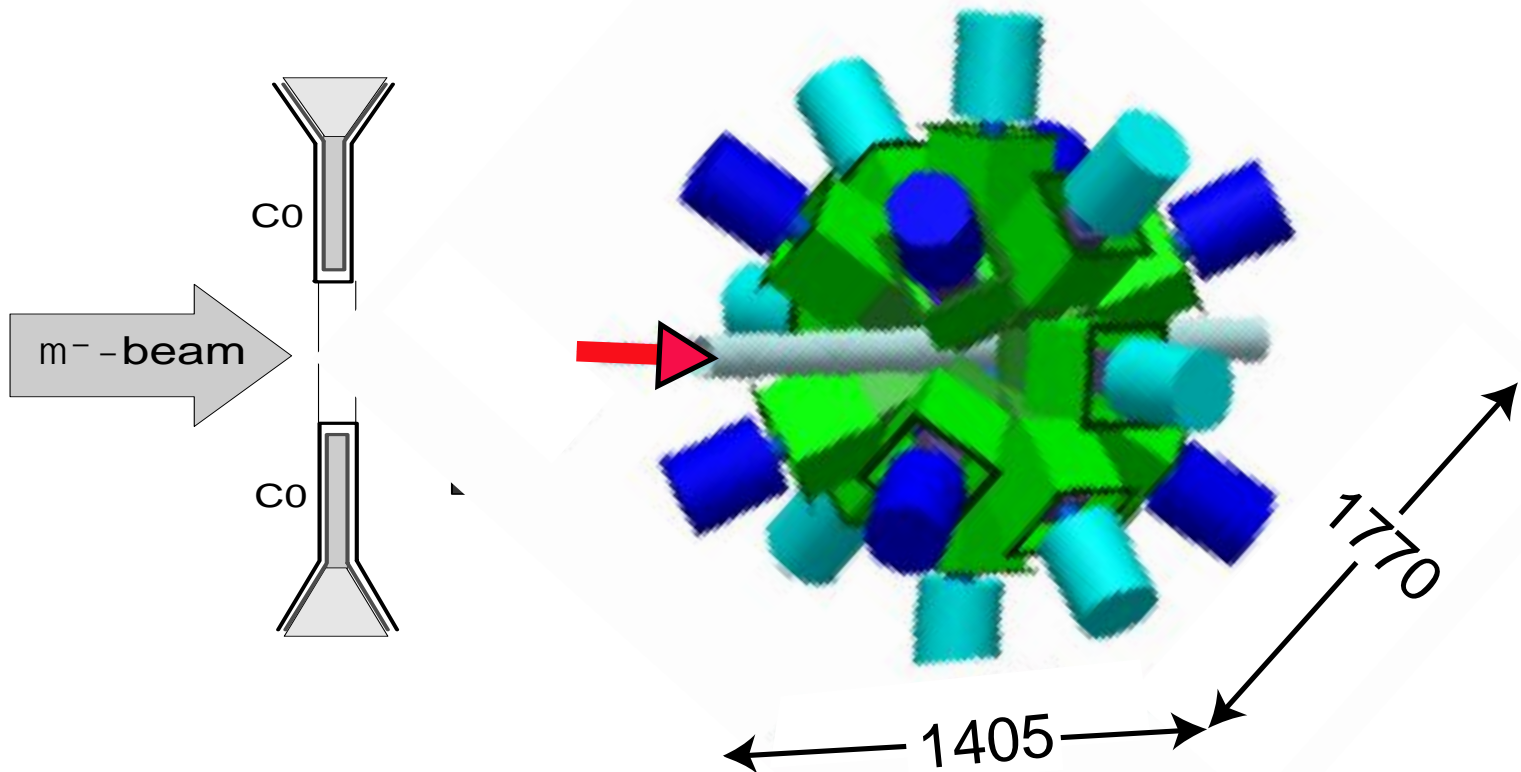


$$\mu_{stop} = \overline{C0} \wedge C1 \wedge C2 \wedge \overline{C3}$$

# of  $\mu$ -stop =  $8 - 25 \times 10^3$  with  $20 - 30$  MeV/c

target can be used for solids and gas

# Schematics of set-up muon beam line facility “MuSIC” + “CAGRA” at RCNP Osaka



CAGRA = Clover Array Gamma RAY spectrometer

MuSIC = MUon Science Innovative muon beam Channel

# The issue of $g_A$ quenching

**But: things are a bit more complicated**

there is a pseudo-scalar coupling effective in  $\mu$ -capture with a constant  $g_p$  ( also badly known !!)

**What is this ????**

Inside the nucleus the muon can decay back into a virtual pion (lots of energy available!!), and the pion generates a final state imprinting it with the parity of the pion. ( $P(\pi) = -1$ )

**The effect depends on how many protons there are.**

$^{24}\text{Mg} \rightarrow 12$  protons

$^{32}\text{S} \rightarrow 16$  protons

$^{56}\text{Fe} \rightarrow 26$  protons

# Conclusion

- Precision mass measurement:

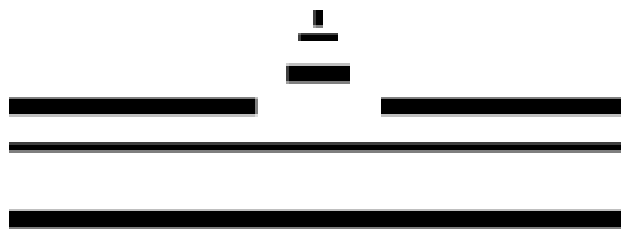
$^{96}\text{Zr}$  → golden case for testing NME's of  $0\nu\beta\beta$  decay  
gives handle on  $g_A$  for a unique forbidden decay

- $\mu$ -capture

maybe the only viable tool to study weak response at high momentum transfer ( $0\nu\beta\beta$  decay) and to fix the  $g_A$  problem by comparing with ( $d, ^2\text{He}$ )

# Thank you

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WILHELMS-UNIVERSITÄT  
MÜNSTER**