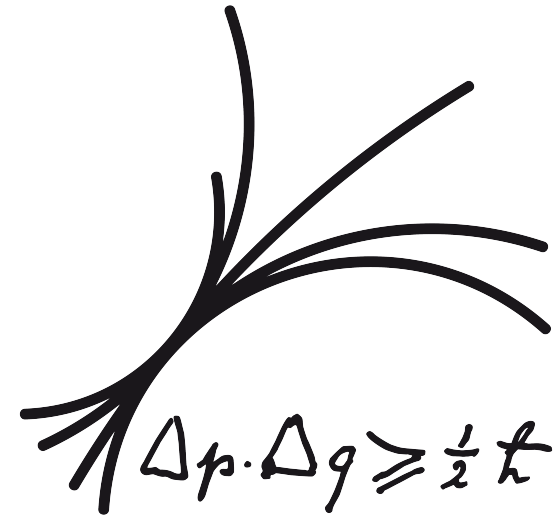


# Sterile Neutrinos as Dark Matter: **NOT** warm **BUT** non-thermal



*Alexander Merle*  
MPP Munich  
Germany



Based on:

MAX-PLANCK-GESELLSCHAFT

JCAP **1506** (2015) 011, Phys. Lett. **B749** (2015) 283, Int. J. Mod. Phys. **D22** (2013) 1330020, JCAP **1403** (2014) 028, JCAP **1107** (2011) 023, Phys. Rev. **D88** (2013) 113004, JCAP **1604** (2016) 003, JCAP **1701** (2017) 025, JCAP **1611** (2016) 038, 1609.04671 [hep-ex], Astrophys. J. **836** (2017) 61 , 1704.07838 [astro-ph], ...

Erice 2017 Workshop, 19 September 2017



## Contents:

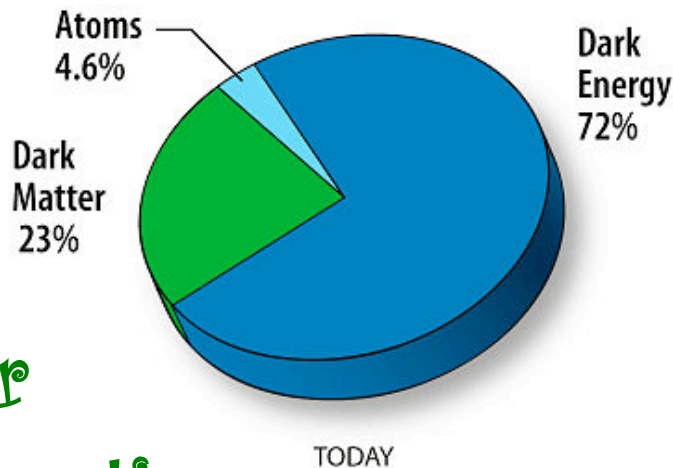
1. To WIMP or not to WIMP
2. Sterile neutrinos
3. Non-thermal DM production
4. Decay production of sterile neutrinos
5. Structure Formation
6. Where to go from here?!?

# 1. To WIMP or not to WIMP

*What do we actually know about Dark Matter?!?*

**KNOWN** 😊

- abundance
- rough distribution
- important for structure formation



**UNKNOWN** ☹️

- identity
- production mechanism
- exact velocity spectrum

**THUS:** We should be careful not to overlook possibilities just because they are called "non-standard"!!!

# 1. To WIMP or not to WIMP

*Maybe our most natural guess for the identity of DM is a yet unknown elementary particle:*

- historically most natural possibility: **WIMP**
  - ↳ because:
    - weak interaction known → “NATURAL”
    - stable WIMPs are predicted in particular by SUSY and by extra dimensions → “THEORY MOTIVATION”
    - comparatively good detection prospects → “EXPERIMENTAL INTEREST”
- **BUT**: unfortunately no clear detection so far...

**→ let us think the unthinkable...**

**WHAT IF DARK MATTER IS NOT A WIMP?!?**

## 2. Sterile neutrinos

What is a **sterile neutrino** and why could it be a good Dark Matter candidate?!?

- ordinary (“active”) neutrino  $\nu_a$ : known elementary particle with very small mass and only weak interactions
- sterile neutrino  $\nu_s$ : may have a larger mass (value theoretically not predicted) and does not at all participate in standard interactions (**BUT**: small mixing with  $\nu_a$ )
- **thus**: if produced in the right amounts and with a suitable velocity spectrum,  $\nu_s$  could act as DM if they are sufficiently stable

*NB: NO constraint from oscillations!!!*

## 2. Sterile neutrinos

Indeed, a sterile neutrino with a (typical) mass of a few keV may act as DM, but...

- needs **non-standard production mechanism** (ordinary thermal freeze-out does not work due to tiny coupling)  
→ **warm/cold/non-thermal?** (interesting for structure and/or galaxy formation)
- typically, this is decaying Dark Matter:  $N_1 \rightarrow \nu + \gamma$   
→ **monoenergetic X-ray signal** e.g. from galaxies
- strong connection to ordinary neutrinos  
→ **concrete models can be tested using light neutrinos**

# 3. Non-thermal production mechanisms



# 3. Non-thermal production mechanisms

**Non-thermal...**

*What does this mean?!?*

**THERMAL**

$$f(p) = \frac{1}{\exp\left(\frac{\sqrt{p^2+m^2}}{T}\right) \pm 1}$$

**HOT**

$$T \gg m$$

**WARM/COOL**

$$T \sim m$$

**COLD**

$$T \ll m$$

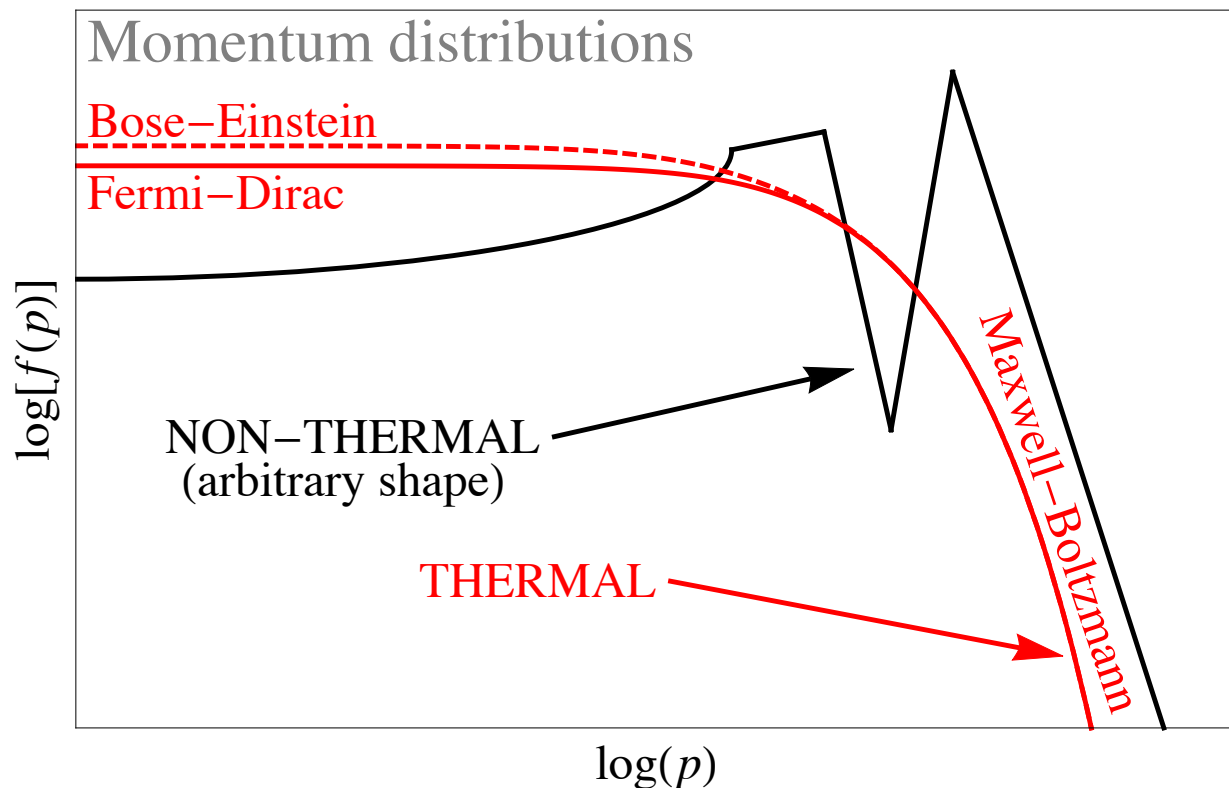
**NON-THERMAL**

$f(p)$  arbitrary

**with:**

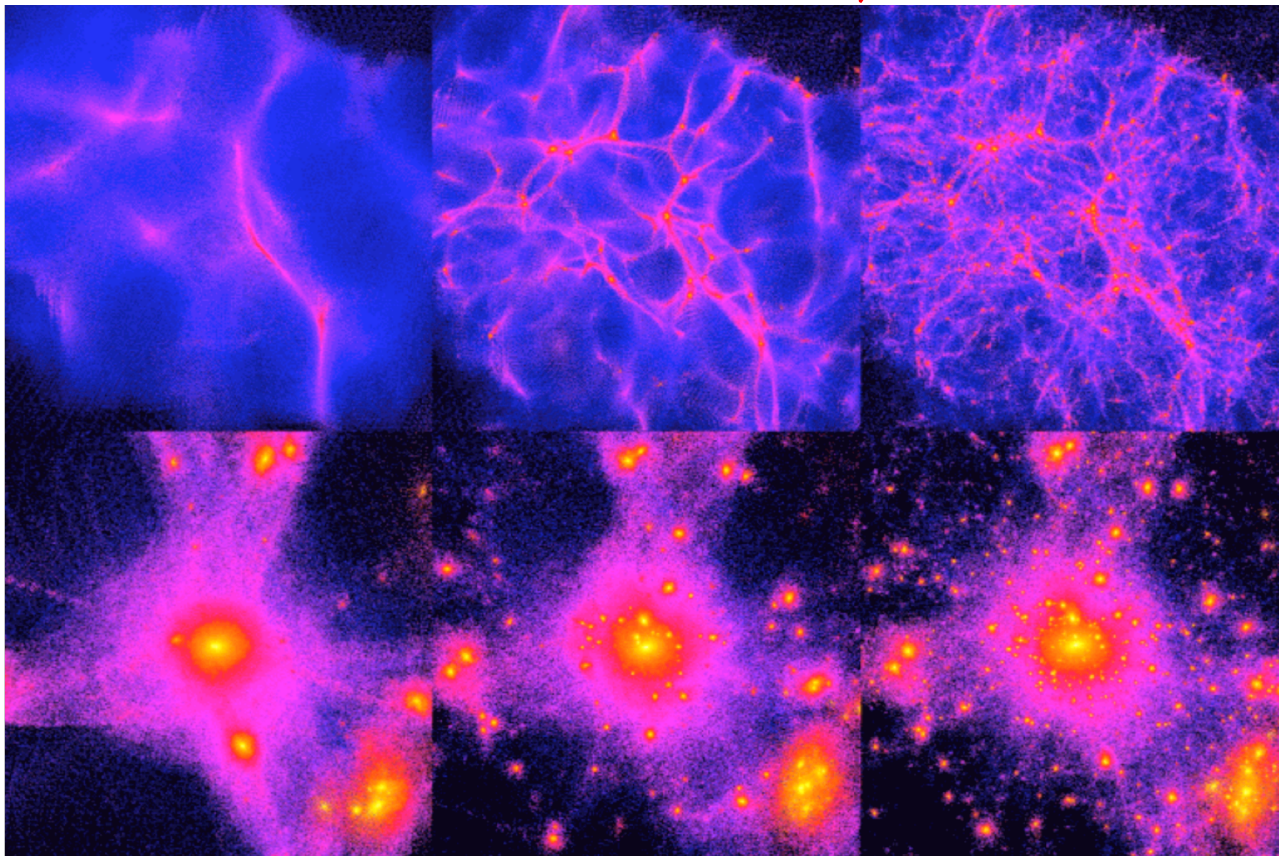
$$\int_{p=0}^{\infty} p^2 f(p) dp < \infty$$

$T$  not defined!!!



### 3. Non-thermal production mechanisms

But what about structure formation?!



[<http://www.ctac.uzh.ch/gallery/>]

**HOT**

**WARM**

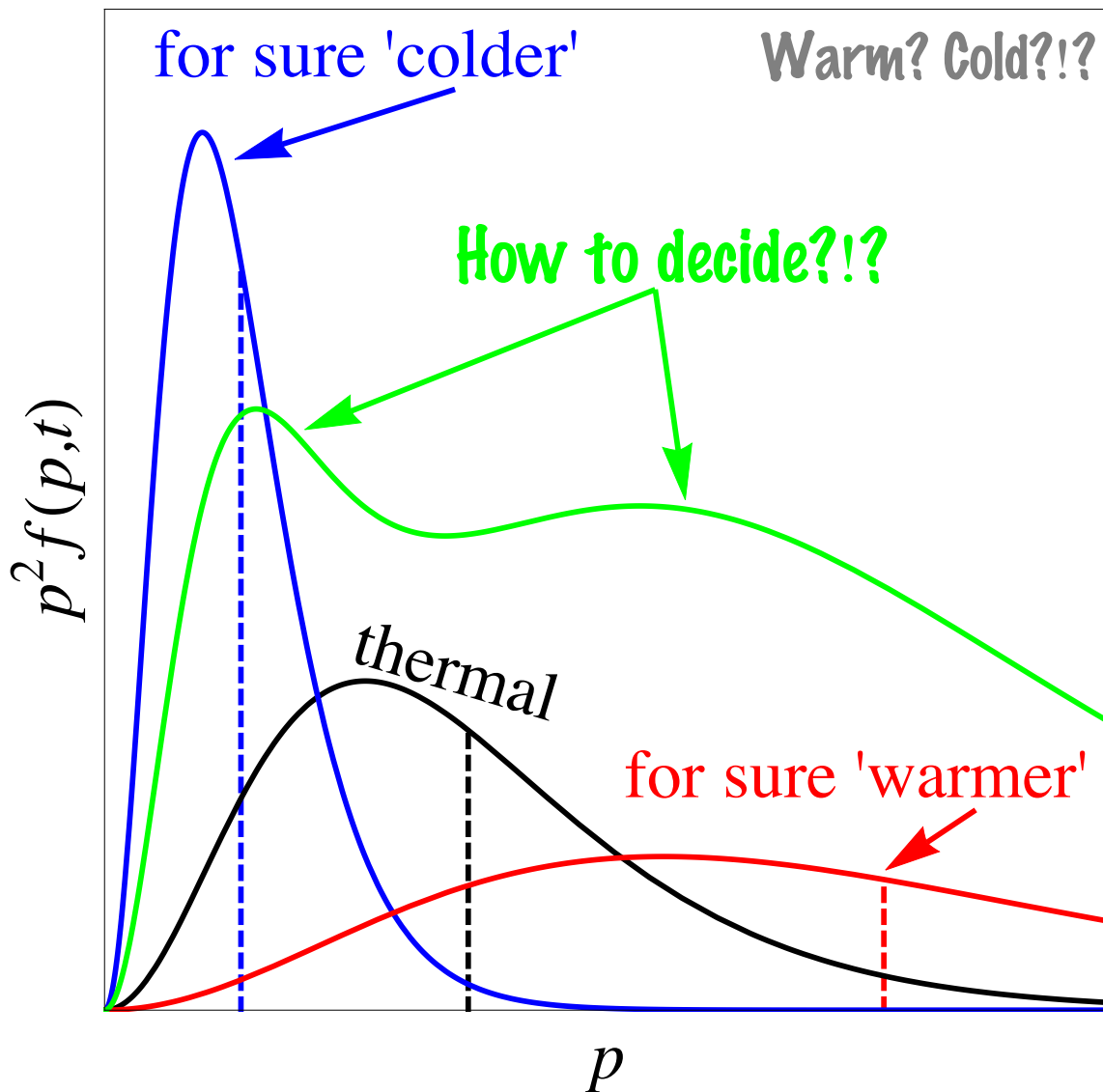
**COLD**

More involved for non-thermal spectra!!!

# 3. Non-thermal production mechanisms

**PROBLEM:**

**HOW TO COMPARE TO THE THERMAL CASE?**



→ Warm Dark Matter is an INSUFFICIENT DESCRIPTION!

→ non-thermal spectra are more complicated!

→ will become clear by the end of my talk

# 4. Production of Sterile Neutrinos

# 4. Production of Sterile Neutrinos

## 4 main mechanisms for sterile neutrinos:

### Non-resonant transitions

First: Langacker

DM: **Dodelson & Widrow**

Idea: active-sterile mixings  
gradually produce  $\nu_s$

### Resonant transitions

First: Enquist, Kainulainen

DM: **Shi & Fuller, Abazajian,  
Shaposhnikov, Laine, ...**

Idea: “MSW”-enhancement  
of active-sterile transitions

### Decay production

$\nu_s$ -DM: **Kusenko, AM,  
Boyanovsky, Shaposhnikov, ...**

Idea: other new particle is  
produced and decays into  $\nu_s$

### Diluted thermal production

$\nu_s$ -DM: **Bezrukov, Lindner,  
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Idea: produce entropy to  
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DM: Shi & Fullinwider, Senjanovic, ...

resonance enhancement

active-sterile transitions

**BUT: When confronted with data...  
WHICH MECHANISMS SURVIVE?!?**

## $\nu_s$ -DM production

$\nu_s$ -DM: Kusenko, AM, Boyanovsky, Shaposhnikov, ...

Idea: other new particle is produced and decays into  $\nu_s$

## Diluted thermal production

$\nu_s$ -DM: Bezrukov, Lindner, Nemevsek, Senjanovic, AM, ...

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**EXCLUDED BY  
STRUCTURE  
FORMATION  
(TOO "HOT" DM)**

## Resonant transitions

Minulainen

Abazajian,

$\nu$ , Laine, ...

" $\nu$ - $\nu$ "-enhancement  
or active-sterile transitions

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STRONGLY  
PUSHED BY  
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(TOO "WARM")

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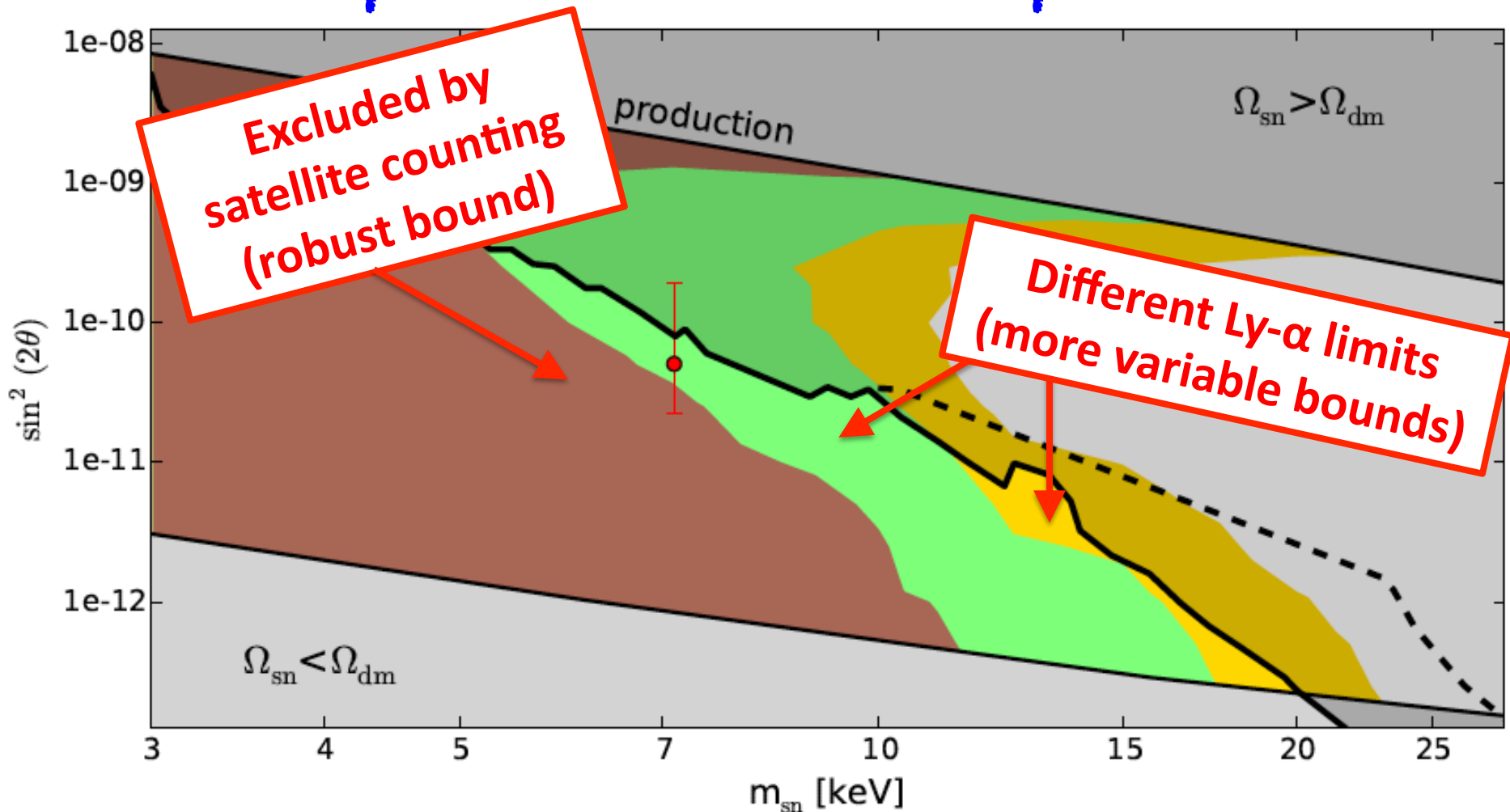
$\nu_s$ -DM: Bezrukov, Lindner, Nemevsek, Senjanovic, AM, ...

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# Resonant production ("standard")

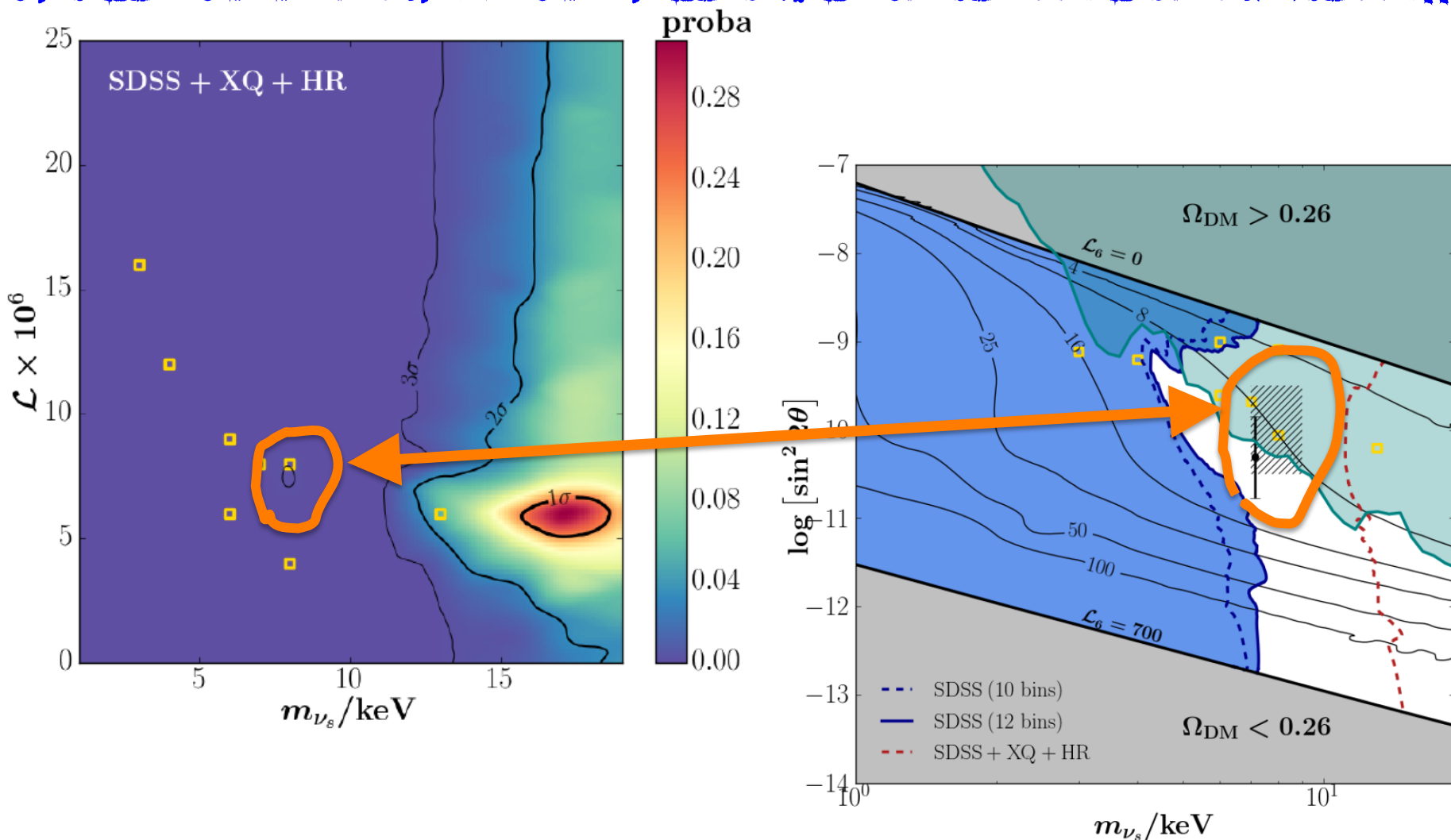
Resonant production under pressure!!



[e.g. Schneider: JCAP 1604 (2016) 059]

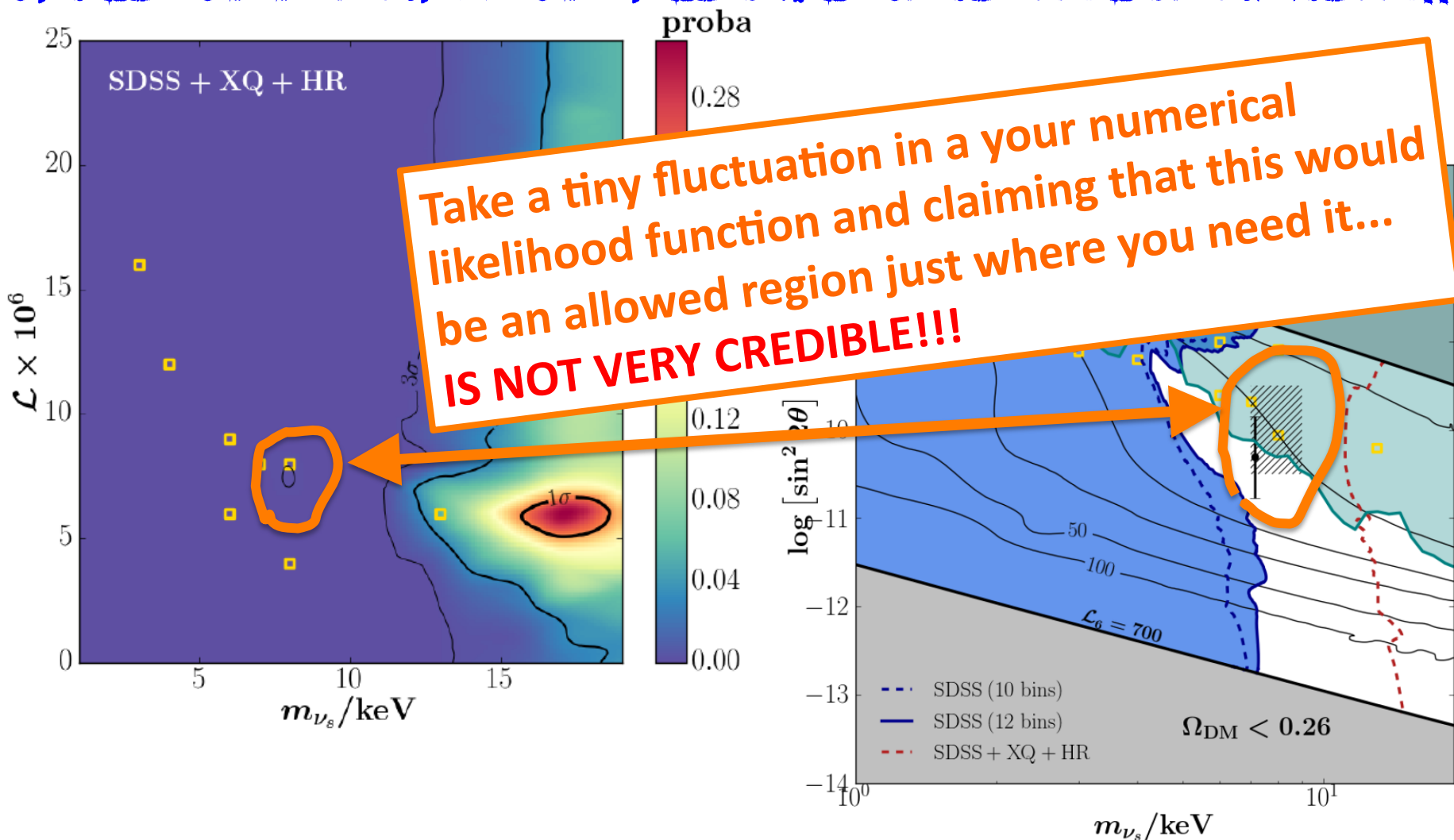
# Resonant production (“standard”)

Thus, better be skeptical if somebody tries too hard to rescue this mechanism....



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### Resonant transitions

First: Enquist, Kainulainen

DM: Shi & Fuller, Abazajian, Shaposhnikov, Laine, ...

Idea: "MSW"-enhancement of active-sterile transitions

$\nu_s$ -DM Possible, but BBN bound is non-trivial to fulfill!!

Idea: other new particle is produced and decays into  $\nu_s$

### Diluted thermal production

$\nu_s$ -DM: Bezrukov, Lindner, Nemevsek, Senjanovic, AM, ...

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OKAY FOR A SIGNIFICANT PART OF  
THE PARAMETER SPACE  
(BUT: partially "boring", i.e. like CDM)

### Decay production

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# 4. Production of Sterile Neutrinos

4 main mechanisms for sterile neutrinos:

## Non-resonant transitions

... **LET US DISCUSS A WORKING EXAMPLE**

DM: **Douglas**  
Idea: active-sterile mixing  
gradually

## Resonant transitions

First: **Enquist, Kainulainen**

**Chen & Fuller, Ab**

**OKAY FOR A SIGNIFICANT FRACTION OF THE PARAMETER SPACE**

**(BUT: partially "boring", i.e. like CDM)**

## Decay prod

$\nu_s$ -DM: **Kusenko, AM, Boyanovsky, Shaposhnikov, ...**

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## 4. *The one that works*: Decay production

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Scalar (or other) decays: e.g.  $S \rightarrow N_1 N_1$

- **decaying inflaton**

[Asaka *et al.*: Phys. Lett. **B638** (2006) 401]

[Anisimov *et al.*: Phys. Lett. **B671** (2009) 211]

[Bezrukov, Gorbunov: JHEP **1005** (2010) 010]

- **singlet scalar that freezes out**

[Kusenko: Phys. Rev. Lett. **97** (2006) 241301; Kusenko, Petraki: Phys. Rev. **D77** (2008) 065014]

[Frigerio, Yaguna: Eur. Phys. J. **C75** (2015), 1]

[AM, Schneider: Phys. Lett. **B749** (2015) 283; AM, Totzauer: JCAP **1506** (2015) 011]

- **singlet scalar that freezes in**

[AM, Niro, Schmidt: JCAP **1403** (2013) 028]

[Adulpravitchai, Schmidt: JHEP **1501** (2015) 006]

[AM, Schneider: Phys. Lett. **B749** (2015) 283; AM, Totzauer: JCAP **1506** (2015) 011]

[König, AM, Totzauer: JCAP **1611** (2016) 038]

[Klasen, Yaguna: JCAP **1311** (2013) 039]

- **other particle that decays**

[Lello, Boyanovsky: Phys. Rev. **D91** (2015) 063502]

[Lello, Boyanovsky: JCAP **1606** (2016) 011]



## 4. *The one that works*: Decay production

- **SINGLET SCALAR “S” FREEZES IN OR OUT BEFORE DECAY**

[Kusenko: Phys. Rev. Lett. **97** (2006) 241301; Kusenko, Petraki: Phys. Rev. **D77** (2008) 065014]

[AM, Niro, Schmidt: JCAP **1403** (2013) 028; AM, Totzauer: JCAP **1506** (2015) 011]

Two-step process:

scalar  $S$  must be produced before it can decay

# 4. The one that works: Decay production

## • SINGLET SCALAR “S” FREEZES IN OR OUT BEFORE DECAY

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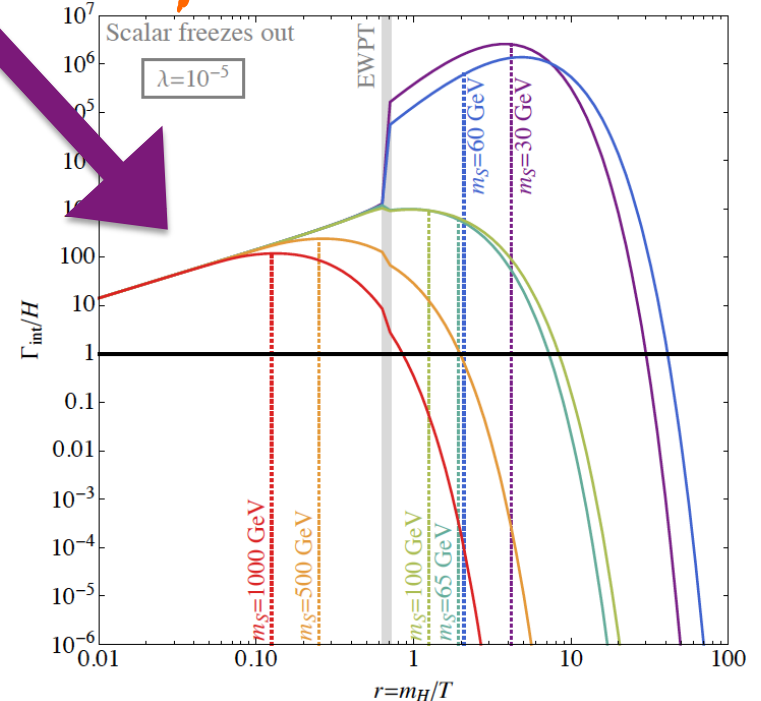
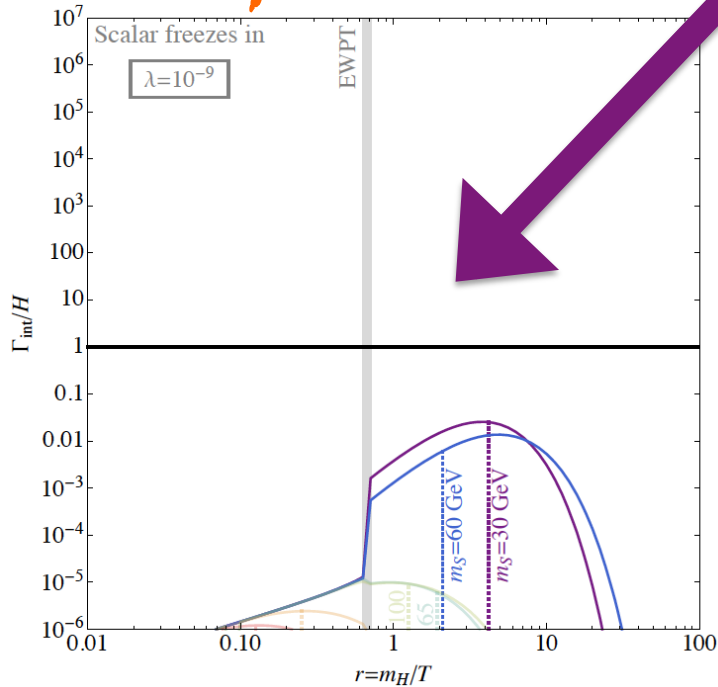
[AM, Niro, Schmidt: JCAP **1403** (2013) 028; AM, Totzauer: JCAP **1506** (2015) 011]

Two-step process:

scalar S must be produced before it can decay

small coupling to SM  
→ freeze-in!

Large coupling to SM  
→ freeze-out!



# 4. *The one that works*: Decay production

## • SINGLET SCALAR “S” FREEZES IN OR OUT BEFORE DECAY

[Kusenko: Phys. Rev. Lett. **97** (2006) 241301; Kusenko, Petraki: Phys. Rev. **D77** (2008) 065014]

[AM, Niro, Schmidt: JCAP **1403** (2013) 028; AM, Totzauer: JCAP **1506** (2015) 011]

### Two-step process:

scalar S must be produced before it can decay

→ solution: distribution of scalar translates directly into distribution of sterile neutrinos

$$f_N(x, r) = \int_0^r dr' 2C_\Gamma \frac{r'^2}{x^2} \int_{\hat{x}_{\min}}^{\infty} d\hat{x} \frac{\hat{x}}{\sqrt{\hat{x}^2 + r'^2}} f_S(\hat{x}, r')$$

$$\hat{x}_{\min} = \left\| x - r^2 / (4x) \right\|.$$

$$\begin{array}{l} x=p/T \\ r=m_s/T \end{array}$$

**direct translation**

# 4. *The one that works*: Decay production

- **SINGLET SCALAR “S” FREEZES IN OR OUT BEFORE DECAY**

[Kusenko: Phys. Rev. Lett. 97 (2006) 241301; Kusenko, Pospelov, 2007] [4]

[AM, Niro, Schmidt: JCAP 1403 (2013) 038]

Two  
sc

→

**We set the bar high in JCAP 1611 (2016) 038 (König, AM, Totzauer):**

- **FULL PARAMETER SPACE**
- **NO APPROXIMATIONS**
- **STEP TO ASTROPHYSICS MADE**

$f_N$

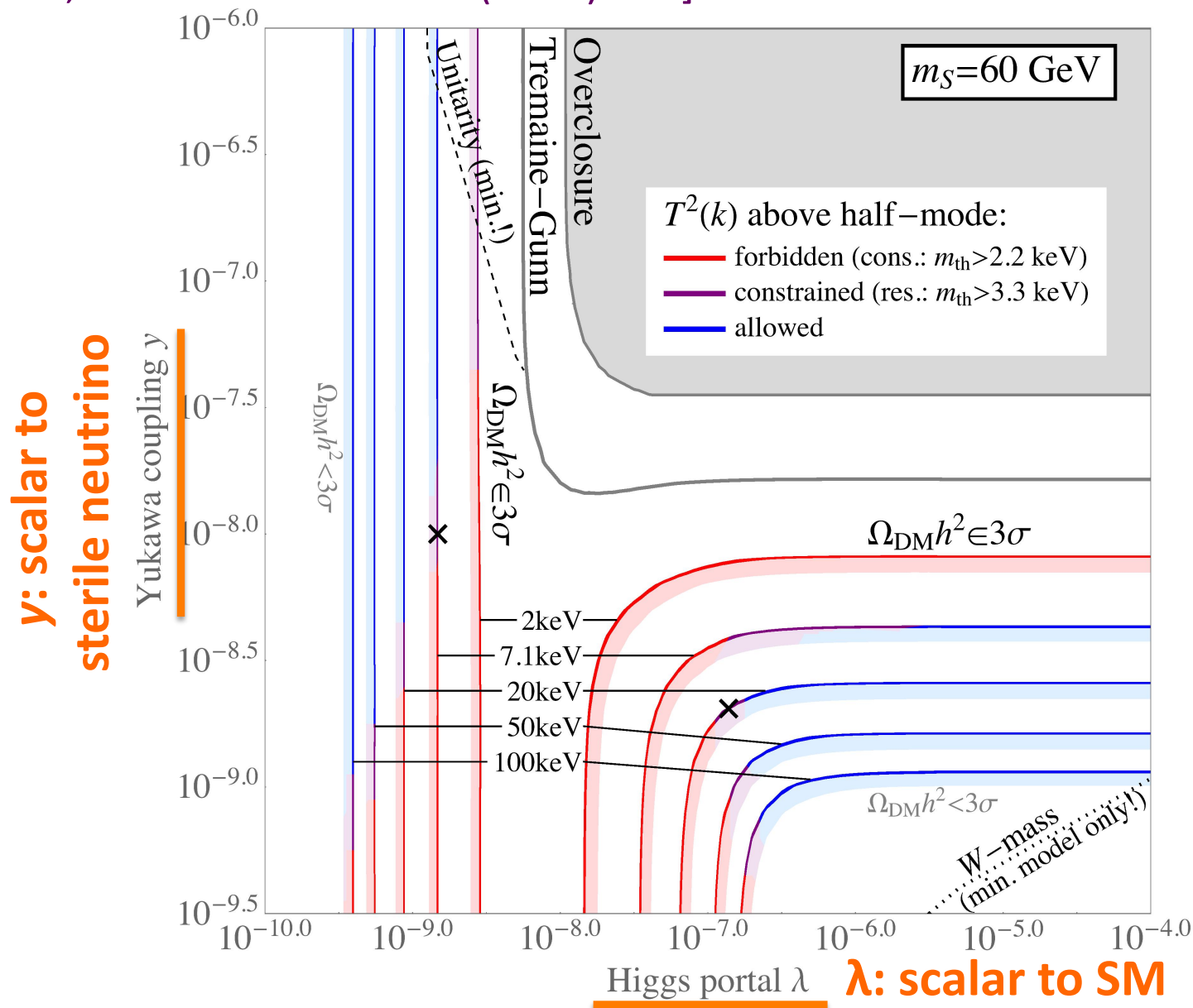
$$x_{\min} = \left\| x - r^2 / (4x) \right\|.$$

$x = p/T$
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**direct translation**

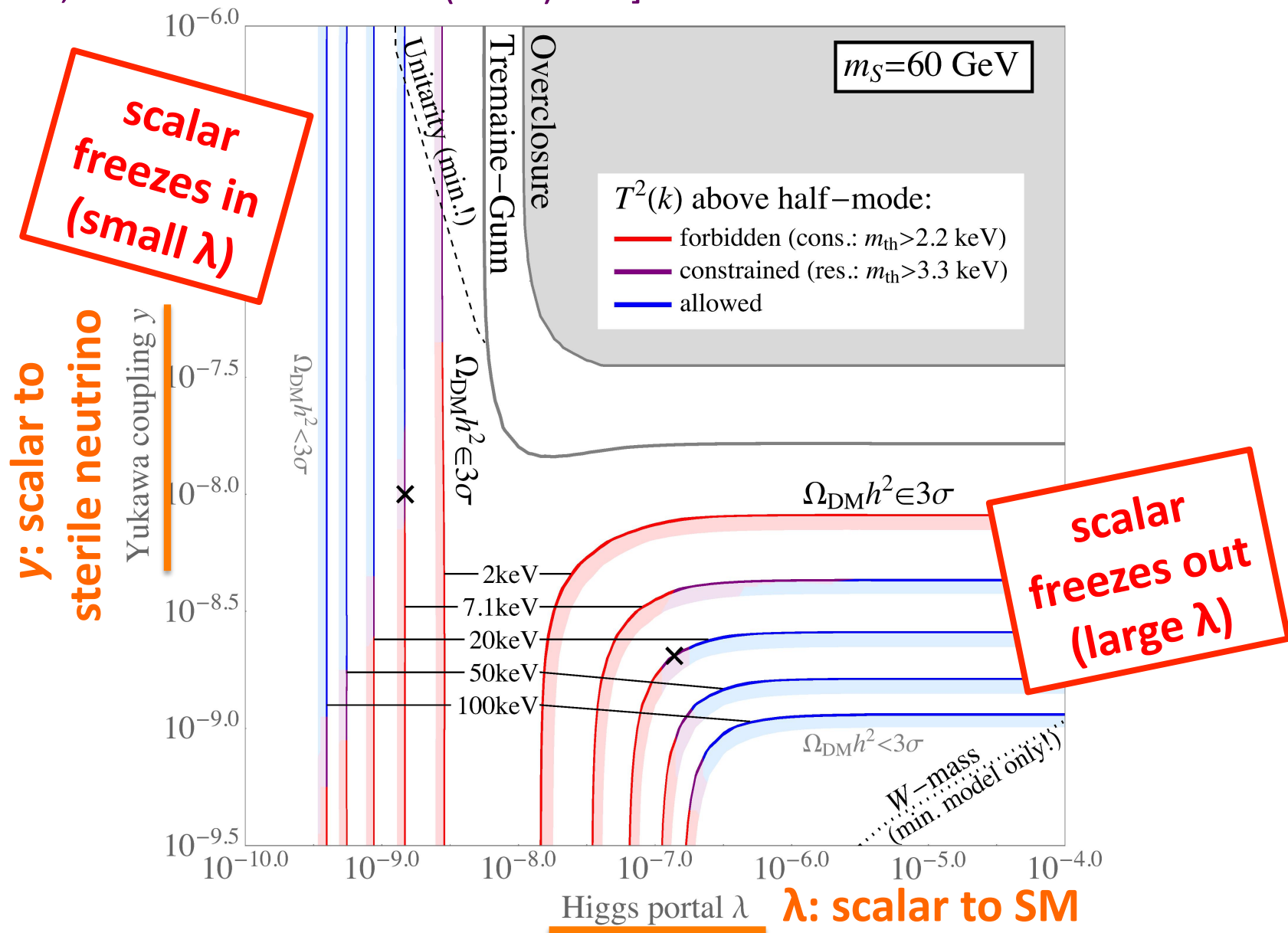
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[König, AM, Totzauer: JCAP **1611** (2016) 038]



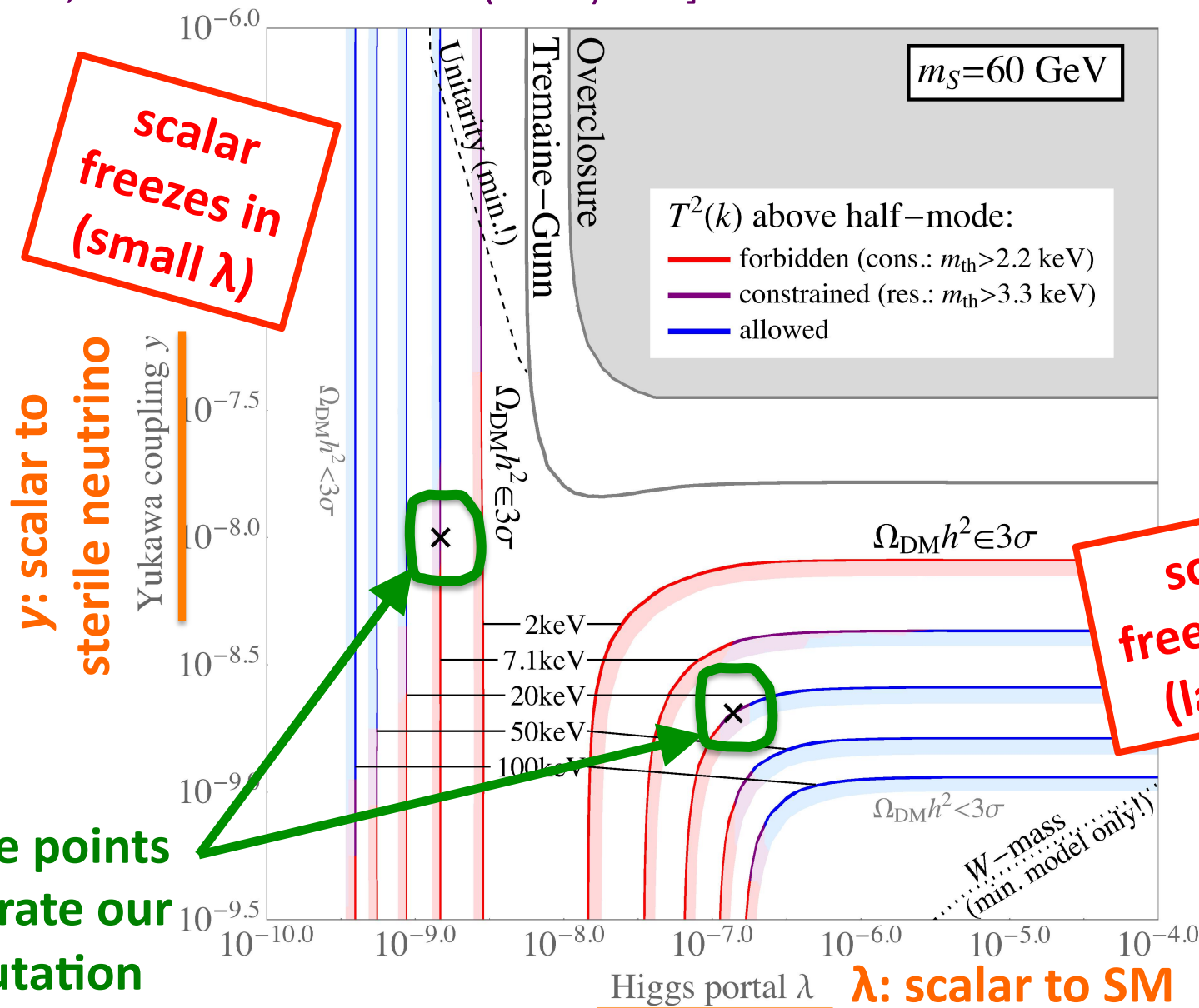
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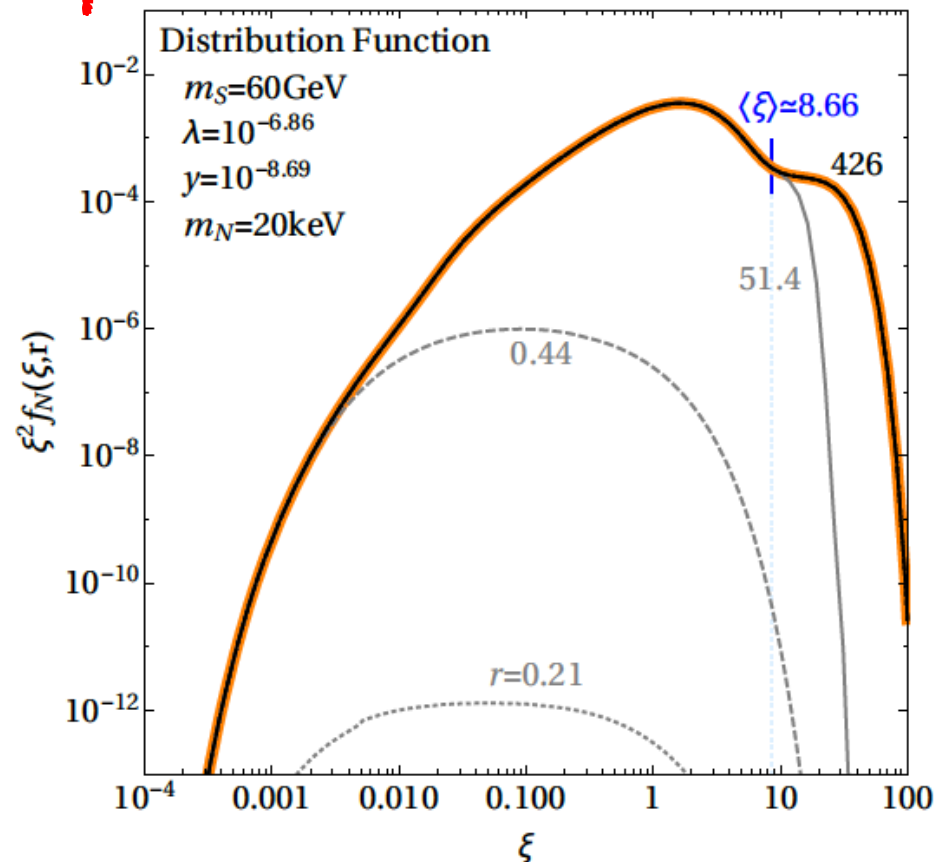
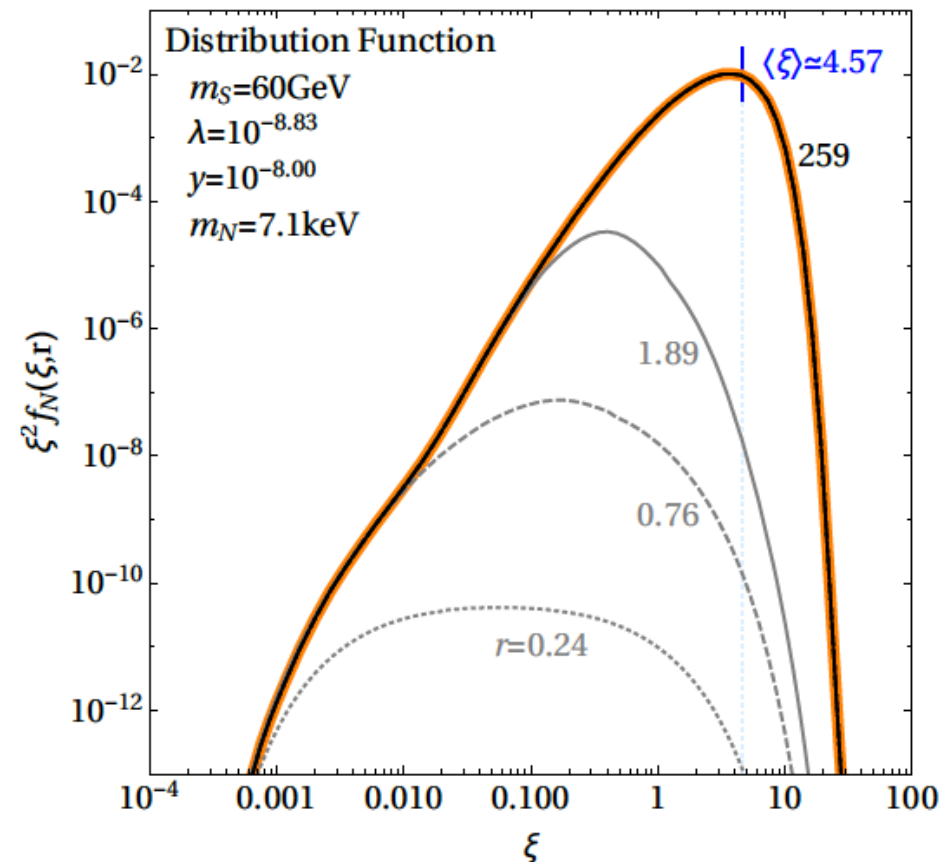
[König, AM, Totzauer: JCAP **1611** (2016) 038]



# 4. *The one that works*: Decay production

[König, AM, Totzauer: JCAP 1611 (2016) 038]

**Distribution functions computed for each case:**



$$\xi = \frac{1}{T_0} \frac{a(t)}{a(t(T_0))} p = \left( \frac{g_s(T_0)}{g_s(T)} \right)^{1/3} \frac{p}{T}$$

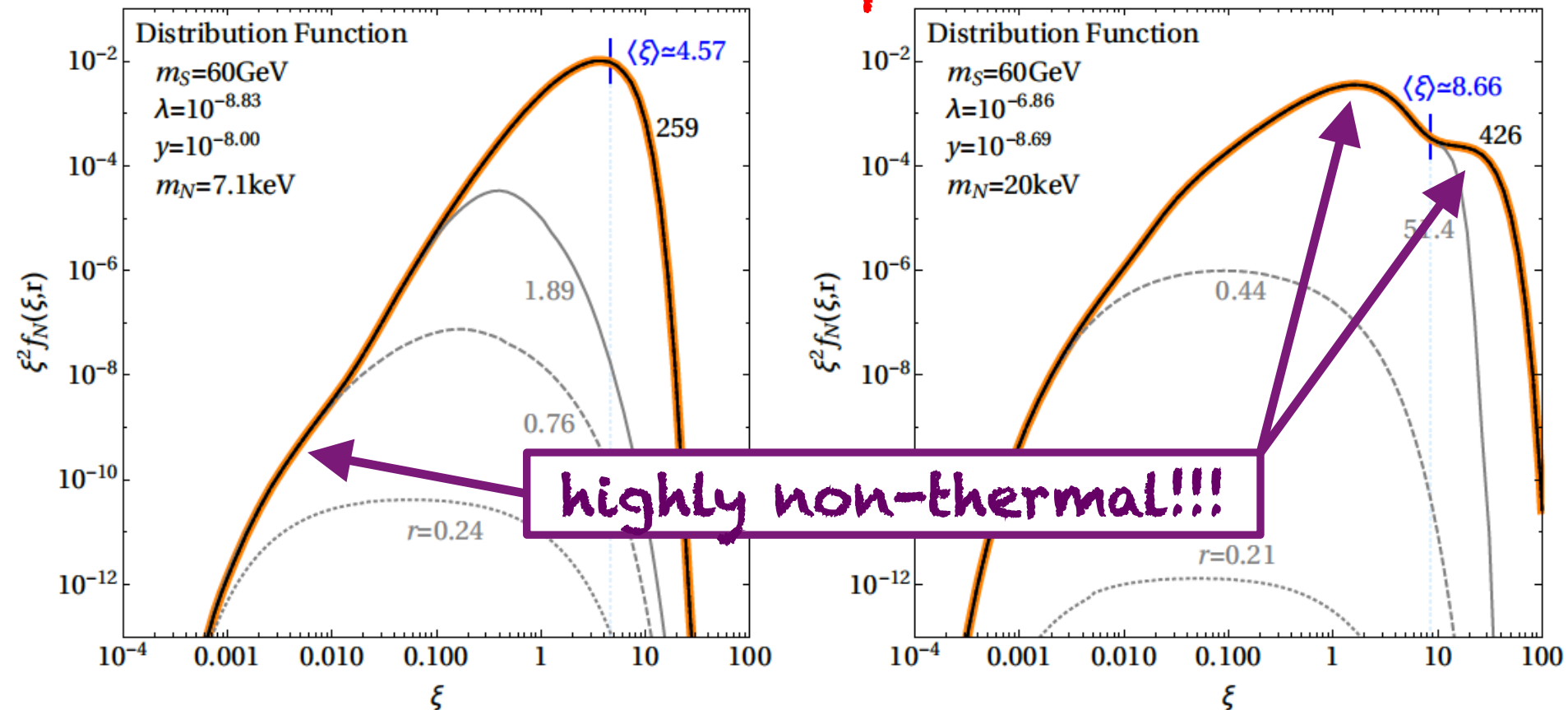
$\xi$  = momentum variable that includes redshift



# 4. The one that works: Decay production

[König, AM, Totzauer: JCAP 1611 (2016) 038]

**Distribution functions computed for each case:**



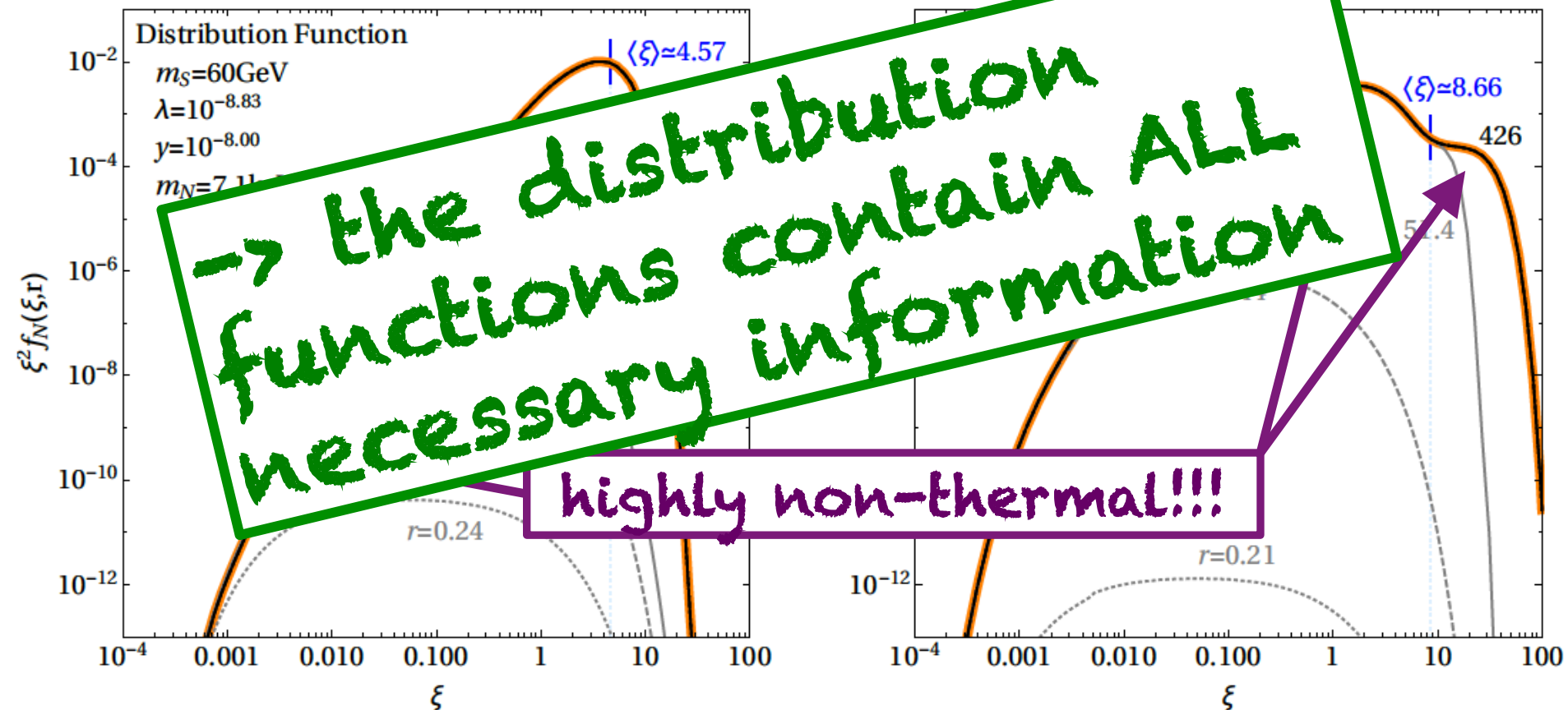
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# 4. *The one that works*: Decay production

[König, AM, Totzauer: JCAP 1611 (2016) 038]

Distribution functions computed for *case:*



→ the distribution functions contain ALL necessary information

highly non-thermal!!!

$$\xi = \frac{1}{T_0} \frac{a(t)}{a(t(T_0))} p = \left( \frac{g_s(T_0)}{g_s(T)} \right)^{1/3} \frac{p}{T}$$

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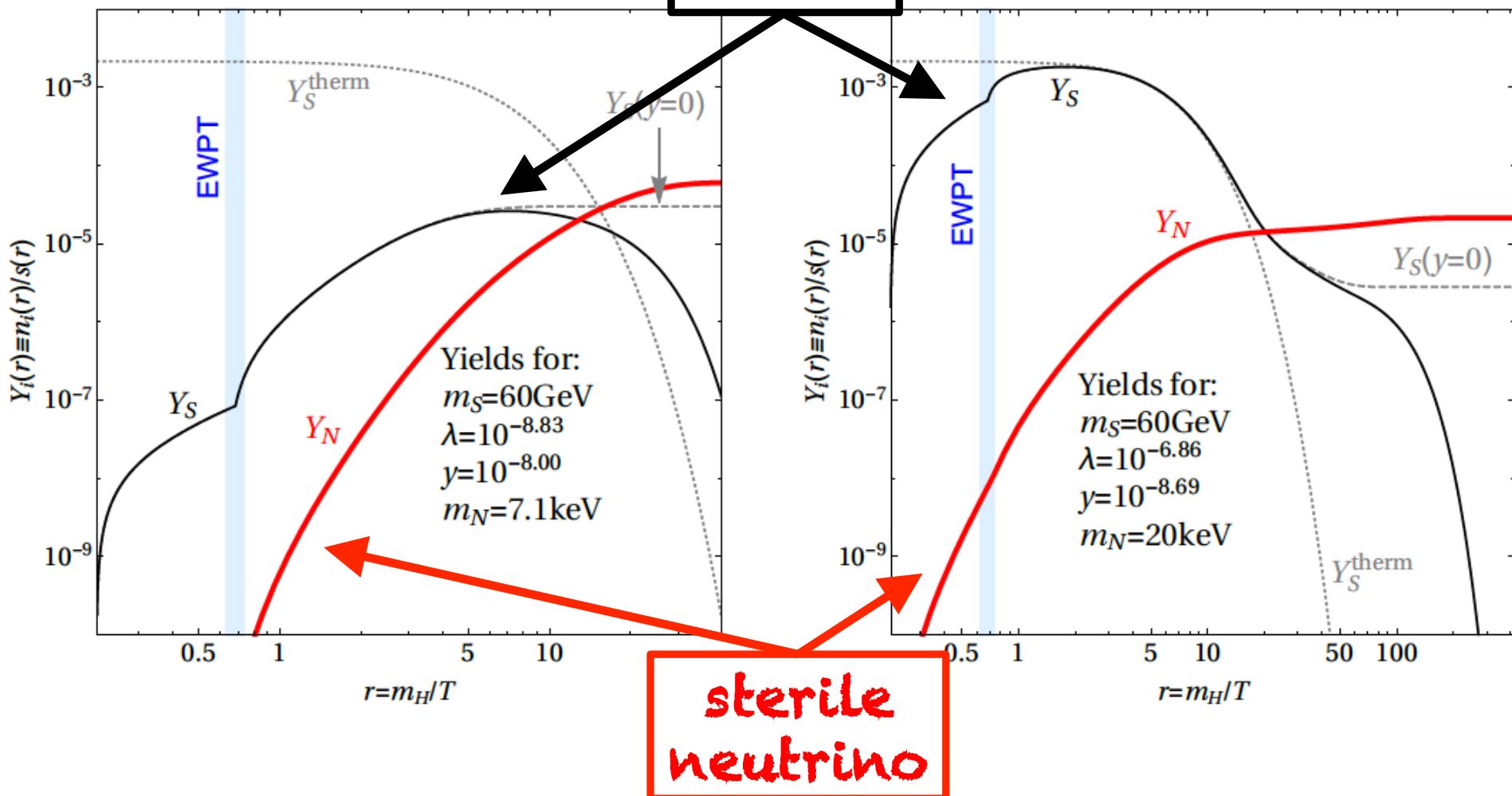
[König, AM, Totzauer: JCAP 1611 (2016) 038]

Example: Abundance/Yield by 2-step process

FIMP

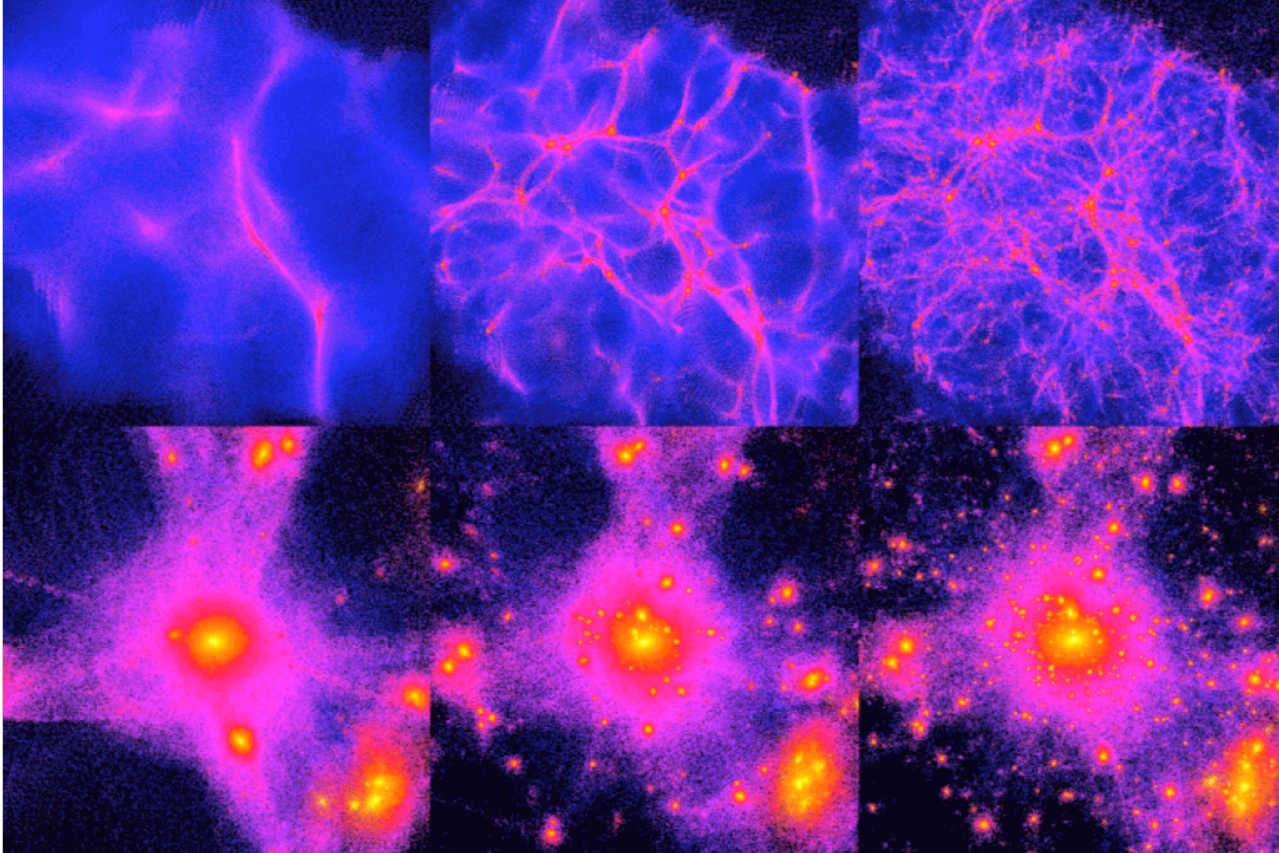
scalar

WIMP



# 5. Structure Formation

Sooo, what about structure formation...?!?



[<http://www.ctac.uzh.ch/gallery/>]

HOT

WARM

COLD

... with non-thermal spectra!!!  
-> We will see that now!!

# 5. Structure Formation

## • SINGLET SCALAR “S” FREEZES IN OR OUT BEFORE DECAY

[Kusenko: Phys. Rev. Lett. **97** (2006) 241301; Kusenko, Petraki: Phys. Rev. **D77** (2008) 065014]  
[AM, Niro, Schmidt: JCAP **1403** (2013) 028; AM, Totzauer: JCAP **1506** (2015) 011]

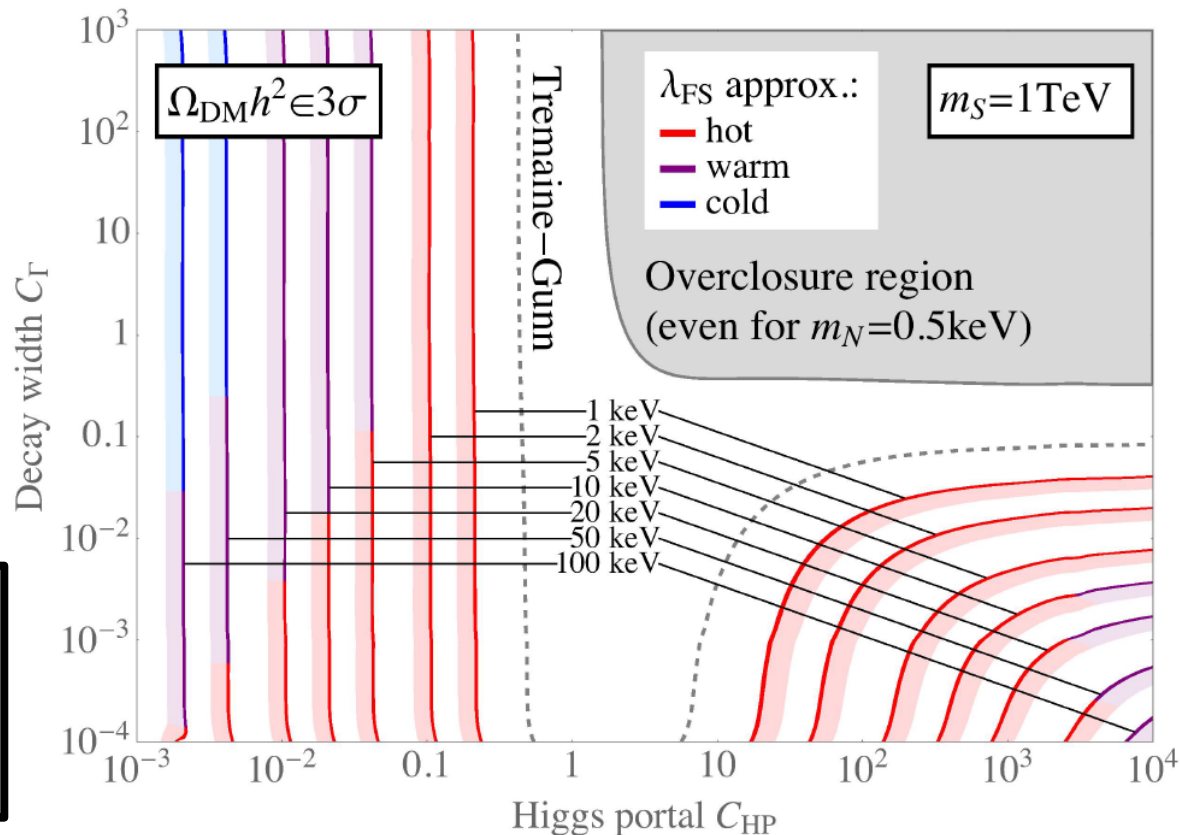
→ “Standard treatment”: **FREE STREAMING**

Free-streaming horizon

$$r_{\text{FS}} = \int_{t_{\text{in}}}^{t_0} \frac{\langle v(t) \rangle}{a(t)} dt$$



Decides about whether the keV sterile neutrinos are **HOT**, **WARM**, or **COLD**



# 5. Structure Formation

## • SINGLET SCALAR “S” FREEZES IN OR OUT BEFORE DECAY

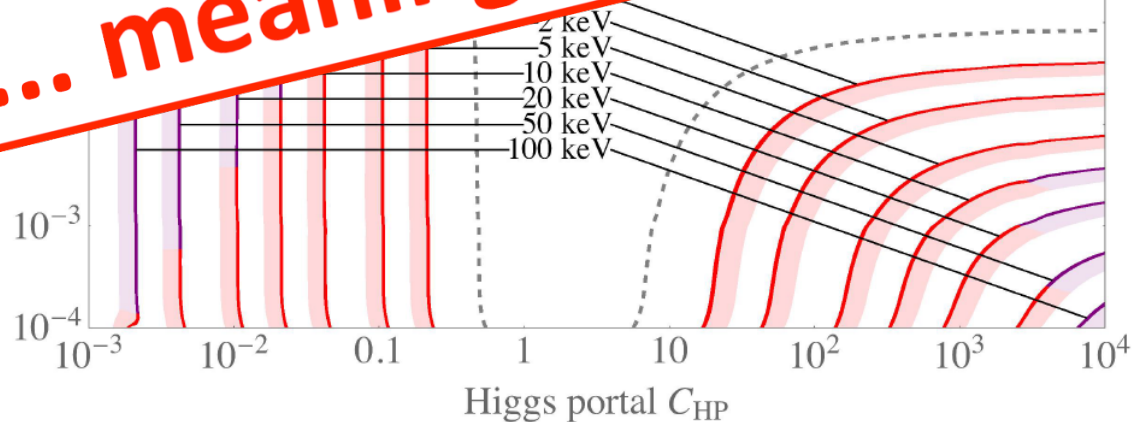
[Kusenko: Phys. Rev. Lett. 97 (2006) 241301; Kusenko, Petrucci, 065014]  
[AM, Niro, Schmidt: JCAP 1403 (2013) 028: AM]

→ “Standard”

Problem: **INSUFFICIENT!!!!**

- free-streaming is inspired from thermal distributions
- too crude! (based on average velocity... meaningless?!?)

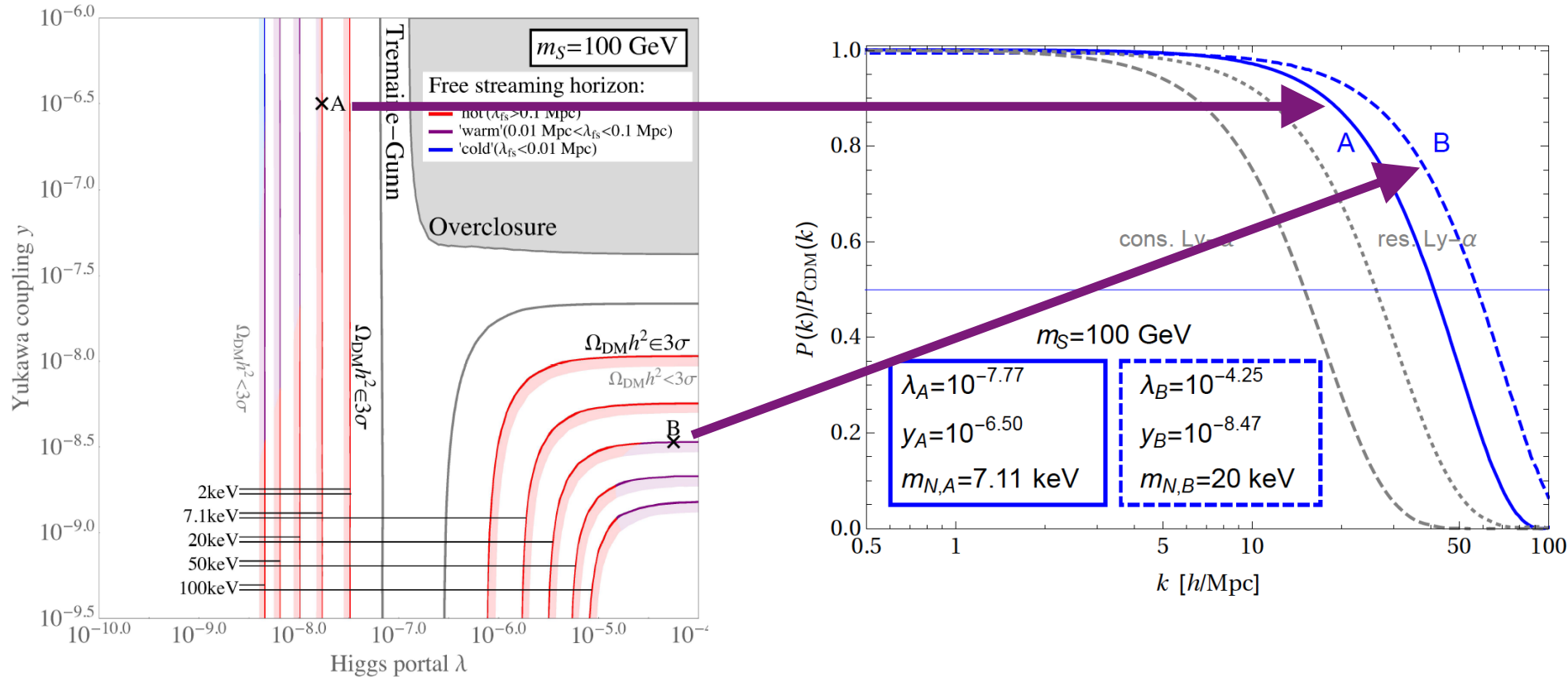
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# 5. Structure Formation

## • HOW FREE STREAMING CAN FAIL:

[König, AM, Totzauer: JCAP 1611 (2016) 038]



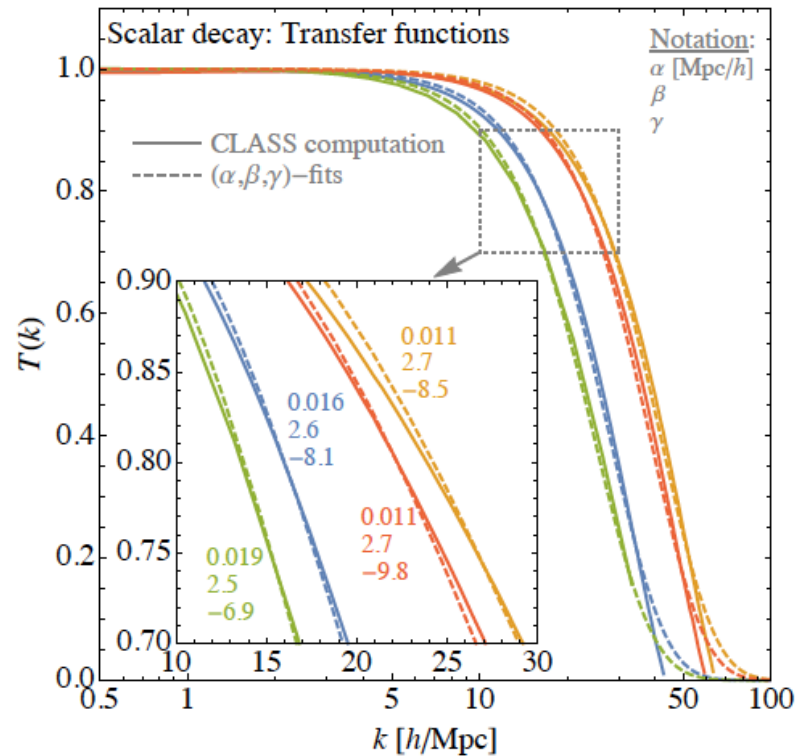
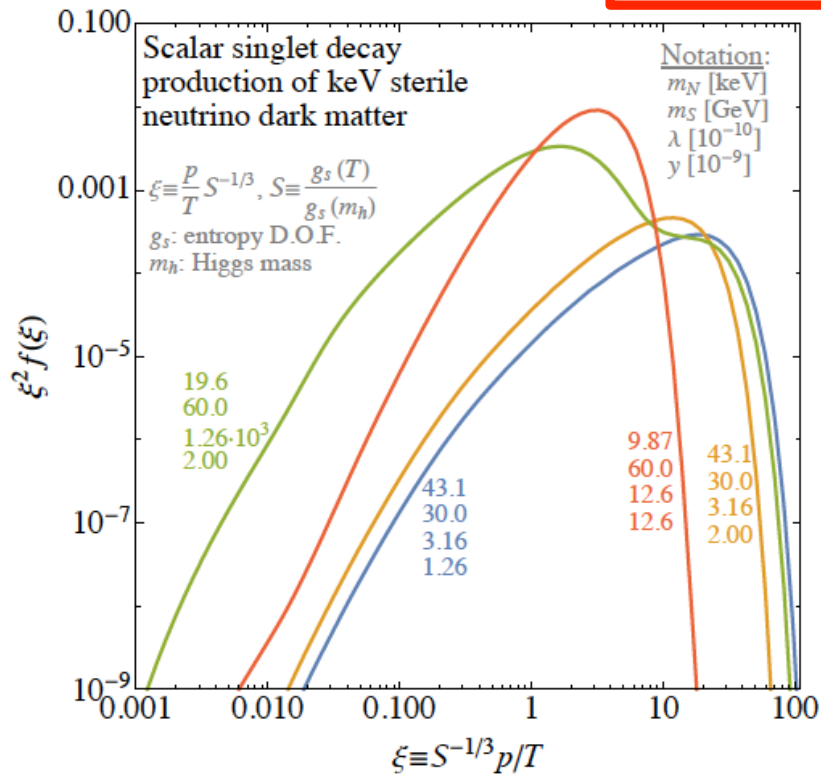
Points that fail according to a free-streaming analysis can actually be in full agreement with data....

# 5. Structure Formation

- **WAY OUT** (VALID FOR ANY NON-THERMAL MODEL ON THE MARKET, AT LEAST SO FAR):

**Generalised TF**

$$T(k) = [1 + (\alpha k)^\beta]^\gamma$$



[Murgia, AM, Viel, Totzauer, Schneider: 1704.07838 [astro-ph]]

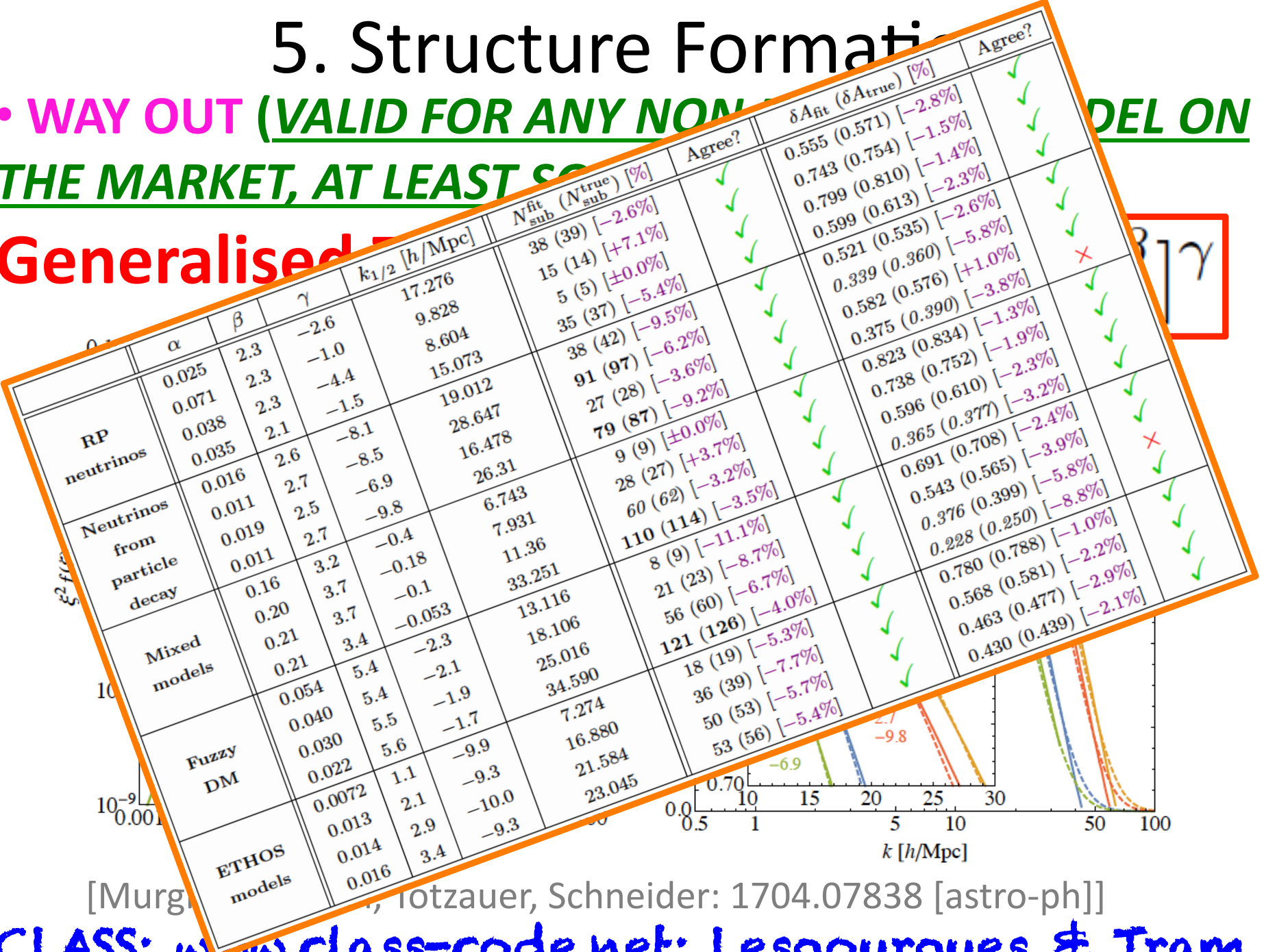
**CLASS: [www.class-code.net](http://www.class-code.net): Lesgourgues & Tram**



# 5. Structure Formation

- WAY OUT (VALID FOR ANY NON-RELATIVISTIC MODEL ON  
THE MARKET, AT LEAST SO FAR)

## Generalised



[Murgia, Totzauer, Schneider: 1704.07838 [astro-ph]]

CLASS: [www.class-code.net](http://www.class-code.net): Lesgourgues & Tram

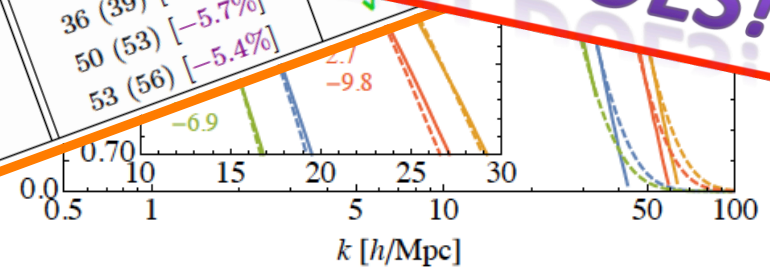
# 5. Structure Formati

• WAY OUT (VALID FOR ANY NON-THERMAL MODEL ON THE MARKET AT LEAST SO FAR)

**Can handle practically any non-thermal spectrum to a sufficient accuracy... NO EXCUSE LEFT TO DODGE SUCH CASES - DON'T BELIEVE ANY PAPER THAT DOES!**

$N_{\text{fit sub}} (N_{\text{true sub}})$ [%]	Agree?	$\delta A_{\text{fit}} (\delta A_{\text{true}})$ [%]	Agree?
(39) [-2.6%]	✓	0.555 (0.571) [-2.8%]	✓
[+7.1%]	✓	0.743 (0.754) [-1.5%]	✓
	✓	0.799 (0.810) [-1.4%]	✓
	✓	0.599 (0.613) [-2.3%]	✓
	✓	0.521 (0.535) [-2.6%]	✓
	✓	0.339 (0.360) [-5.8%]	✗
	✓	0.582 (0.576) [+1.0%]	✓
	✓	0.390 [-3.8%]	✓
	✓	[-1.3%]	✓
	✓	0%	✓

Neutrino from particle decay	0.011	3.2	-0.1	13.1
Mixed models	0.16	3.7	-0.053	18.106
	0.20	3.7	-2.3	25.016
	0.21	3.4	-2.1	34.590
Fuzzy DM	0.054	5.4	-1.9	7.274
	0.040	5.4	-1.7	16.880
	0.030	5.6	-9.9	21.584
ETHOS models	0.0072	1.1	-9.3	23.045
	0.013	2.1	-10.0	
	0.014	2.9	-9.3	
	0.016	3.4		



[Murgia, Totzauer, Schneider: 1704.07838 [astro-ph]]

CLASS: [www.class-code.net](http://www.class-code.net): Lesgourgues & Tram

# ... and it confirms our half-mode analysis:

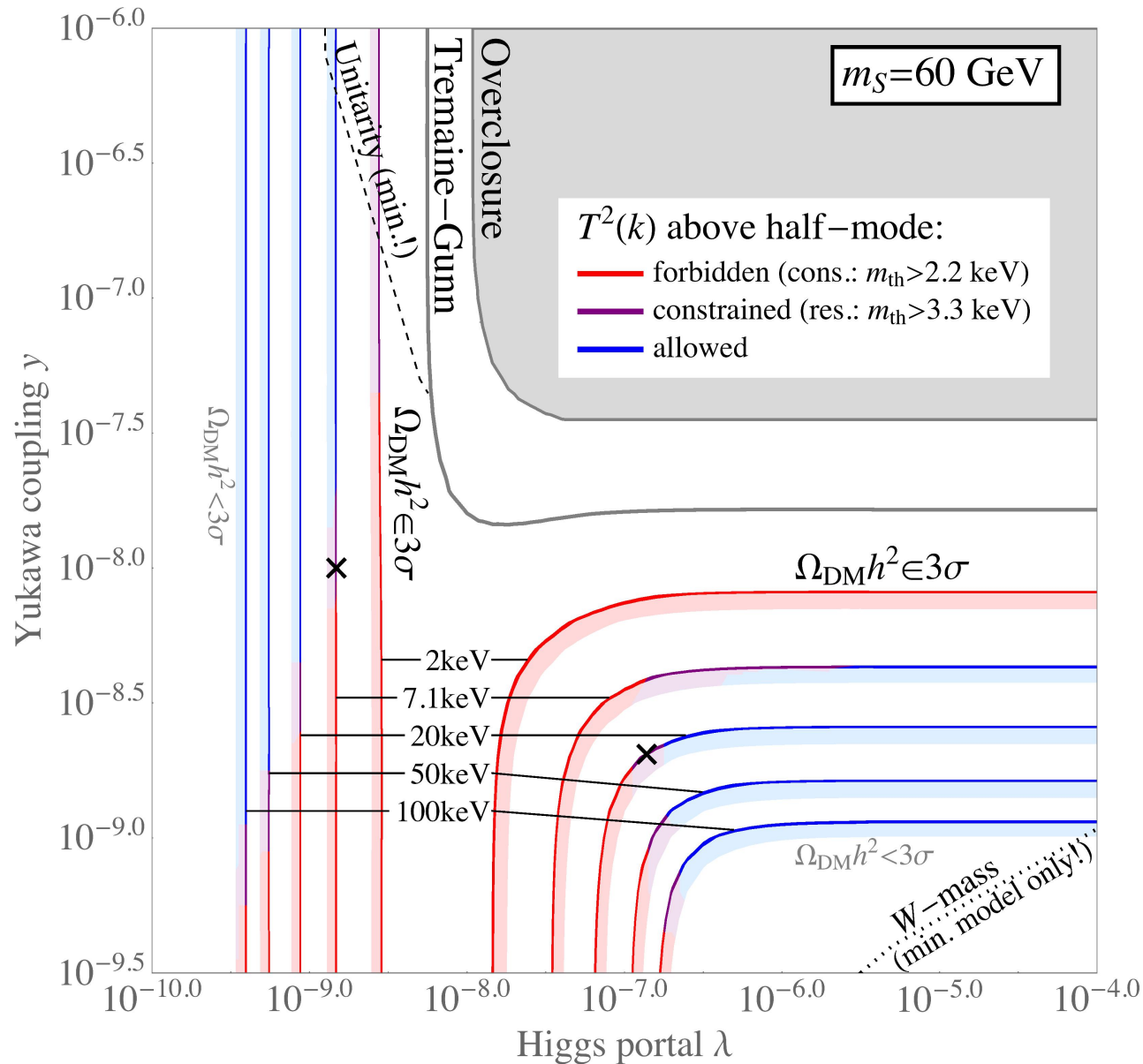
[König, AM, Totzauer: JCAP **1611** (2016) 038]

Allowed regions  
from Ly- $\alpha$  data:

Coservative ✘  
Restrictive ✘

Coservative ✔  
Restrictive ✘

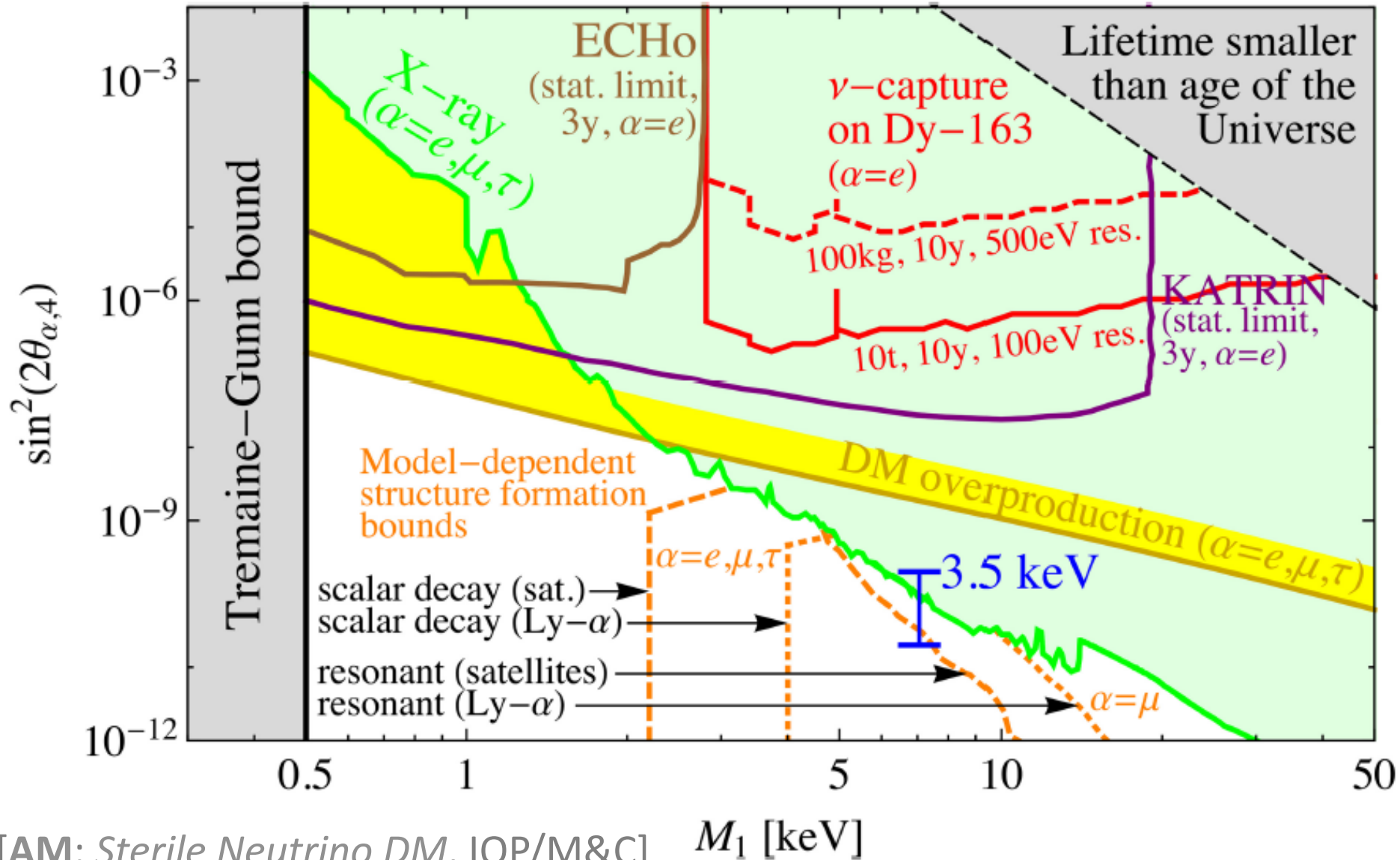
Coservative ✔  
Restrictive ✔



6. Where to go from here?!?

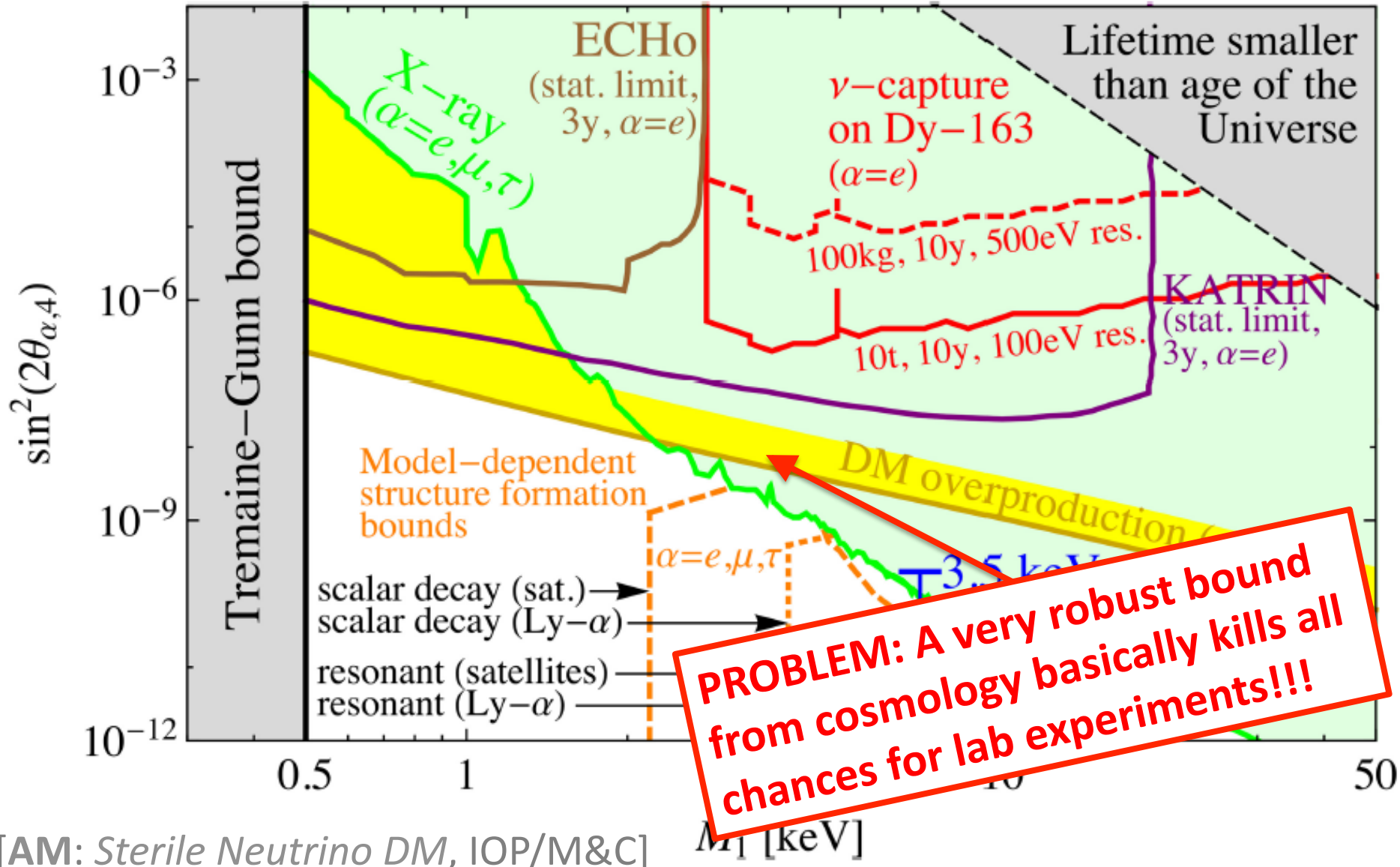
# 6. Where to go from here?!?

## Summary of all constraints:



# 6. Where to go from here?!?

## Summary of all constraints:



# This elephant can't be hidden:

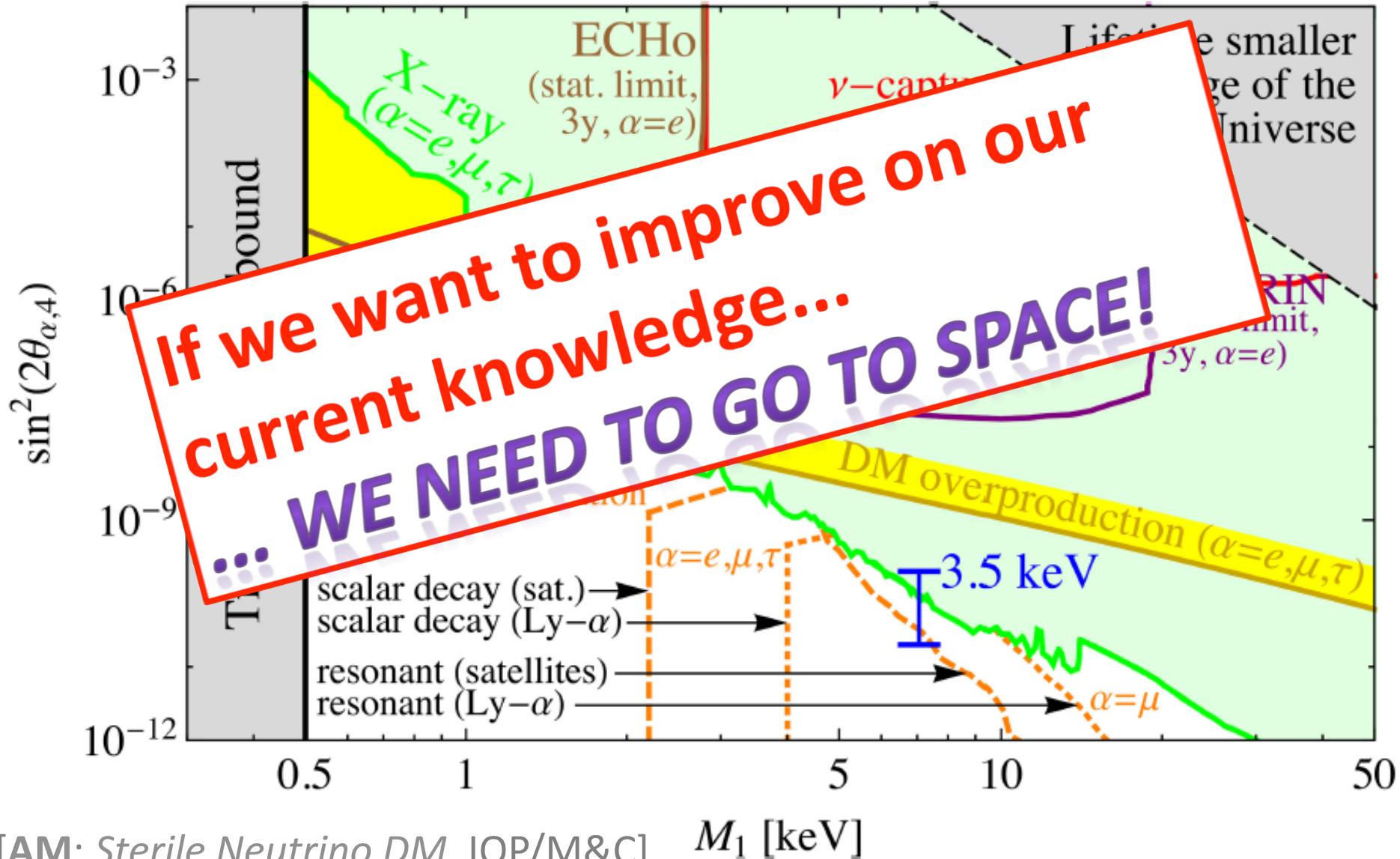
[JCAP 1706 (2017) 051] claims that keV sterile neutrino interactions can be made  $T$ -dependent, so that one can actually detect them in kinematical experiments without spoiling DM production and early Universe cosmology.

**BUT:** Their arguments rely, to quite some extent, on incorrect assumptions and/or oversimplifications...

- DW-distribution of thermal shape:  
**disproved** [AM, Schneider, Totzauer: JCAP 1604 (2016) 003]
- abundance-calculation unreliable:  
**successive error due to incorrect shape**
- hidden scalar sector very sketchy:  
**no actual computation of the decisive scalar production**
- crucial phase transition not actually computed:  
**again, only very sketchy estimates are presented**

# 6. Where to go from here?!?

## Summary of all constraints:





# 6. Where to go from here?!?

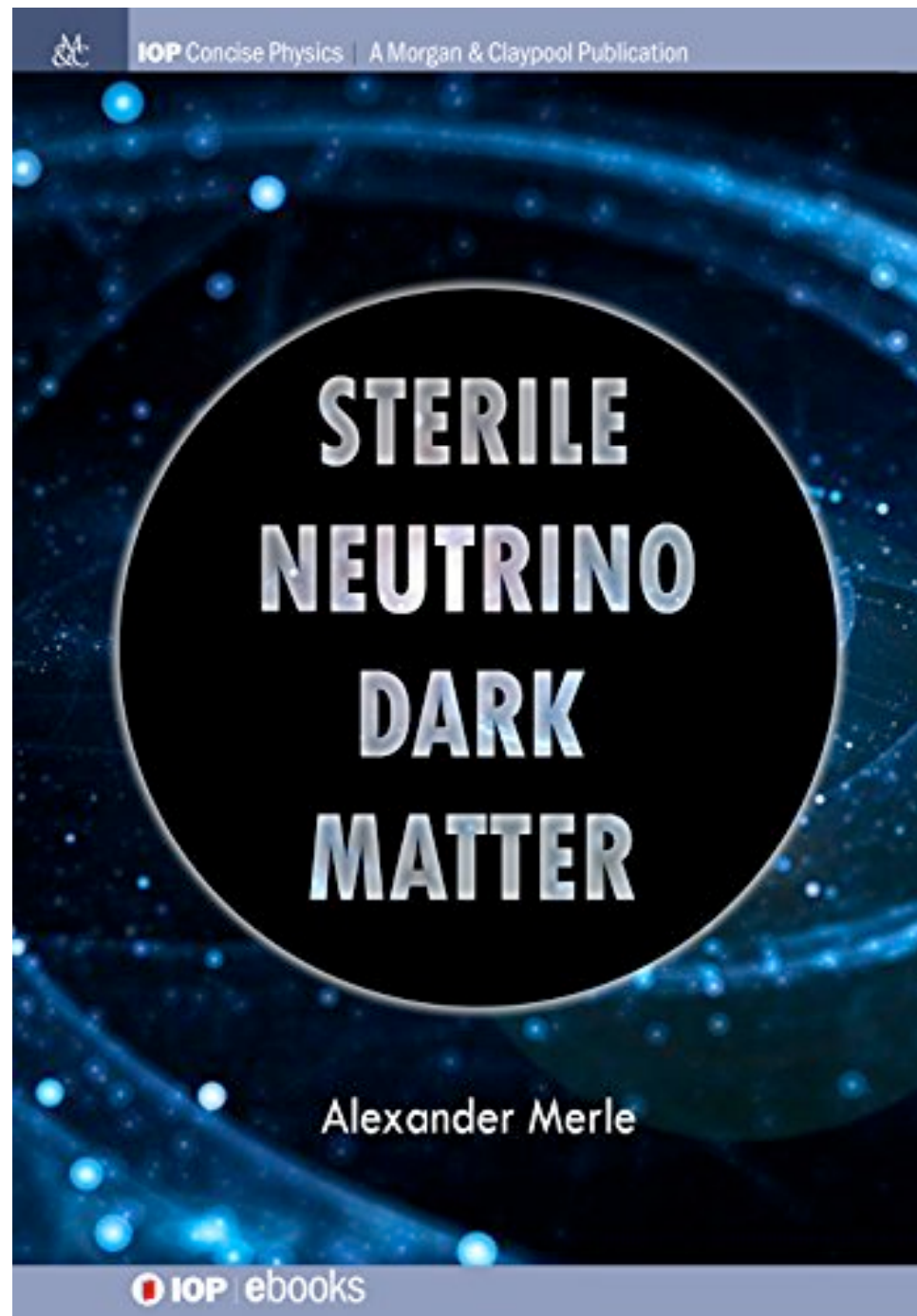
## THREE ROADS LOOK PROMISING:

- **X-ray observations from  $N_1 \rightarrow \nu + \gamma$  decay**  
→ would be good but will take time given that Hitomi (Astro-H) went to the happy hunting ground
- **new astrophysical signatures to constrain structure formation at smaller scales**  
→ e.g. high-redshift galaxies provide strong support for Lyman- $\alpha$  limits
- **seriously including non-thermal DM into high-resolution  $N$ -body simulations**  
→ we have shown that this is possible with less effort than anticipated



And if you want  
to know more...

... **YOU MAY WANNA  
CHECK OUT MY  
NEW BOOK!!**





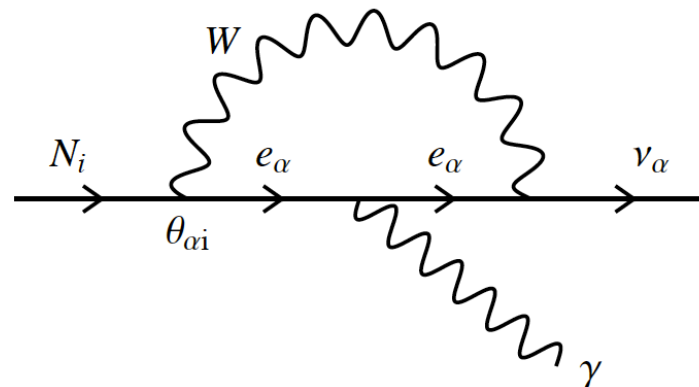
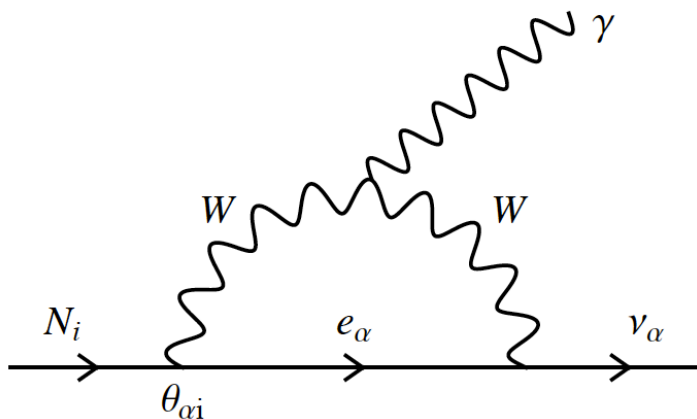
# BACK-UP SLIDES

# 4. Example A: Active-Sterile Transitions

It could all be sooo simple...

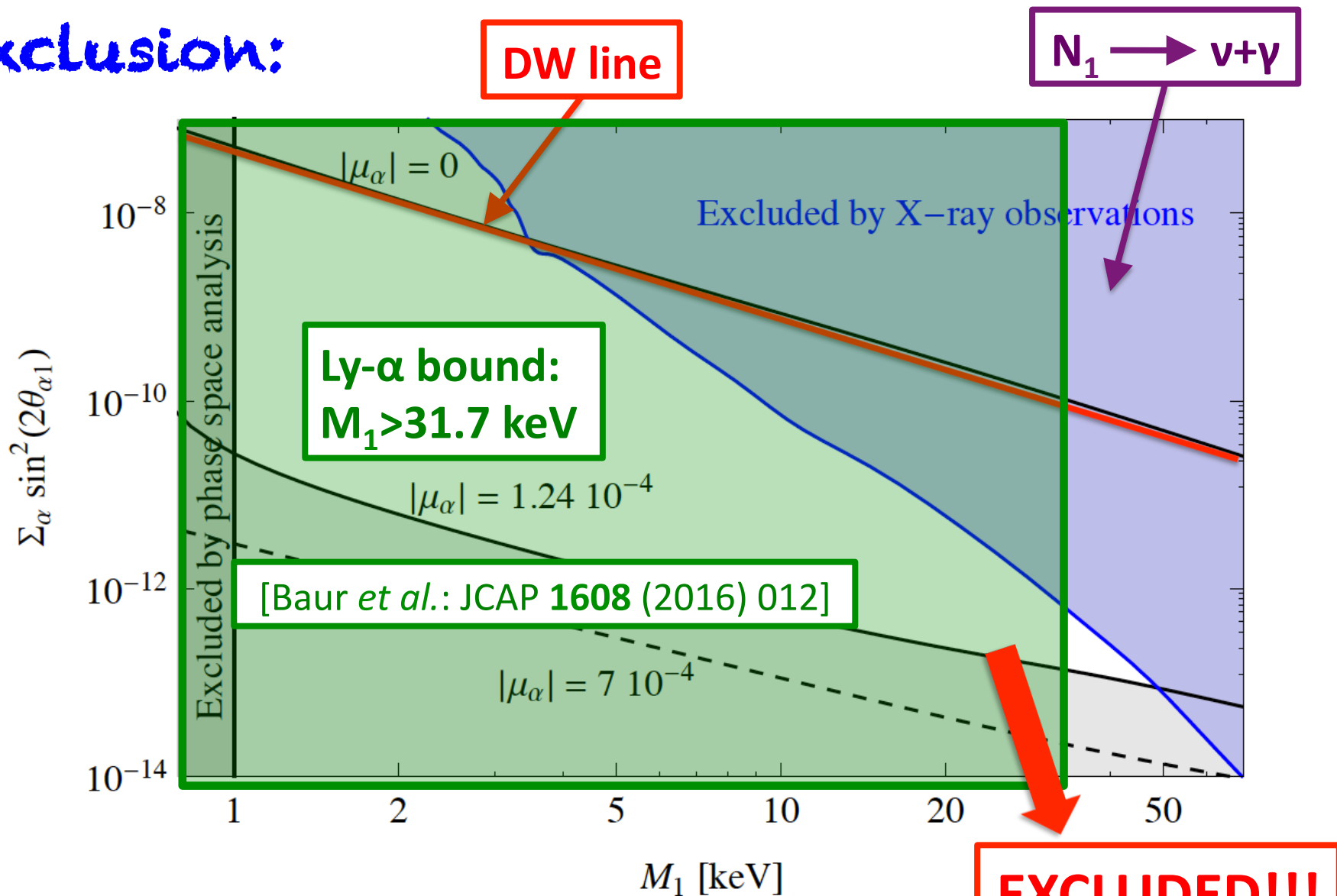
[Dodelson, Widrow: Phys. Rev. Lett. 72 (1994) 17]

- slow non-resonant “oscillations” of active into sterile neutrinos can gradually produce the DM from the thermal plasma (just like “freeze-in”)  $\rightarrow$  nice & simple
- this mechanism produces relatively hot DM  $\rightarrow$  large mass  $M_1$  needed, BUT decay into X-rays scales like  $M_1^5$ :



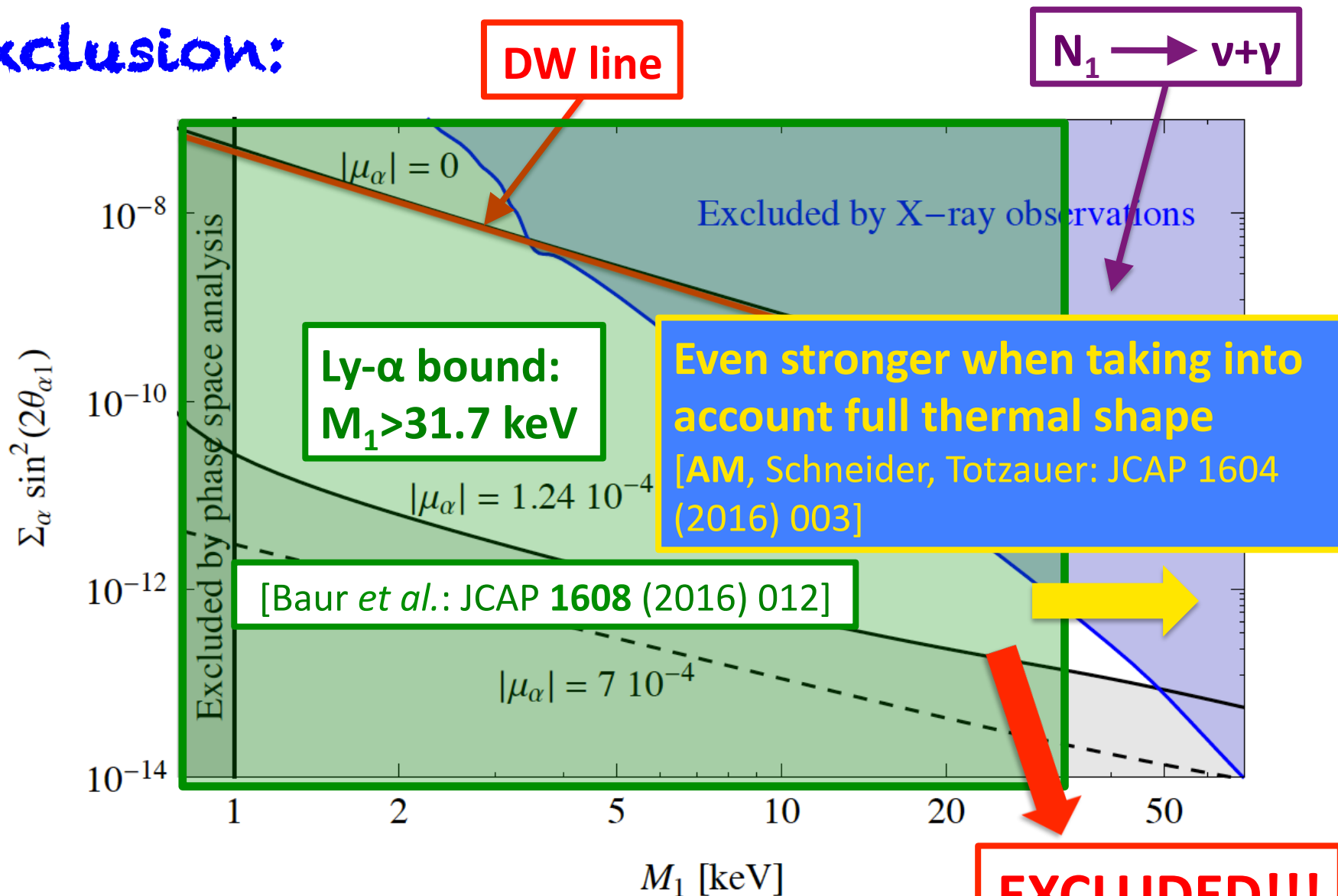
# The Dodelson-Widrow mechanism

Exclusion:



# The Dodelson-Widrow mechanism

Exclusion:



# The Shi-Fuller mechanism

Is there a good way out?!?

[Shi, Fuller: Phys. Rev. Lett. 82 (1999) 2832]

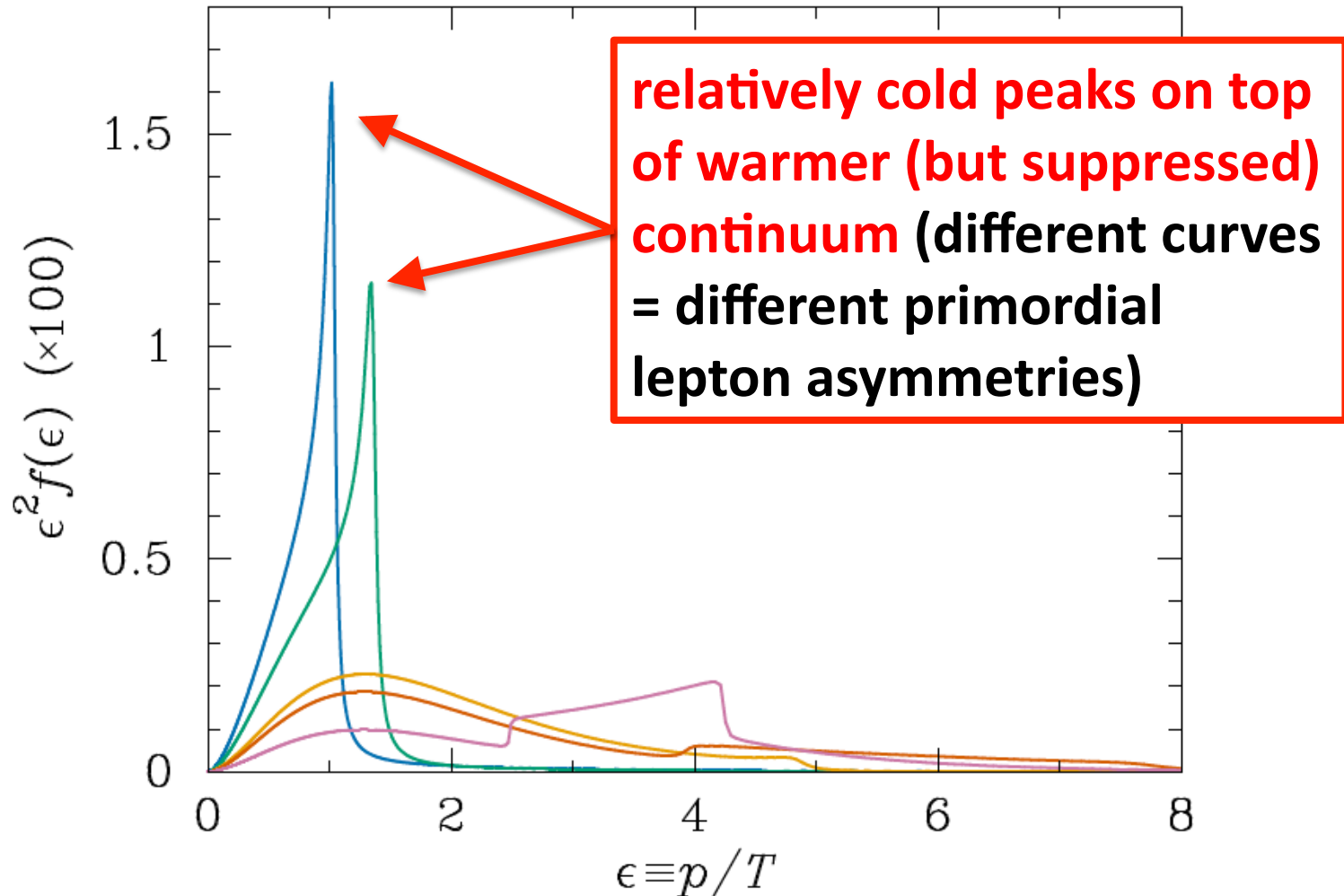
- just like for ordinary neutrinos in the Sun, active-sterile neutrino **transitions could be resonantly enhanced by a sizeable lepton number asymmetry  $|\mu_\alpha|$**  present in the early Universe
- this would produce a large amount of  $\nu_s$  at a specific (momentum-dependent) resonance temperature  
→ *cooler spectrum*
- **BUT: the origin of such a primordial lepton number asymmetry is unclear... XXXX**



# The Shi-Fuller mechanism

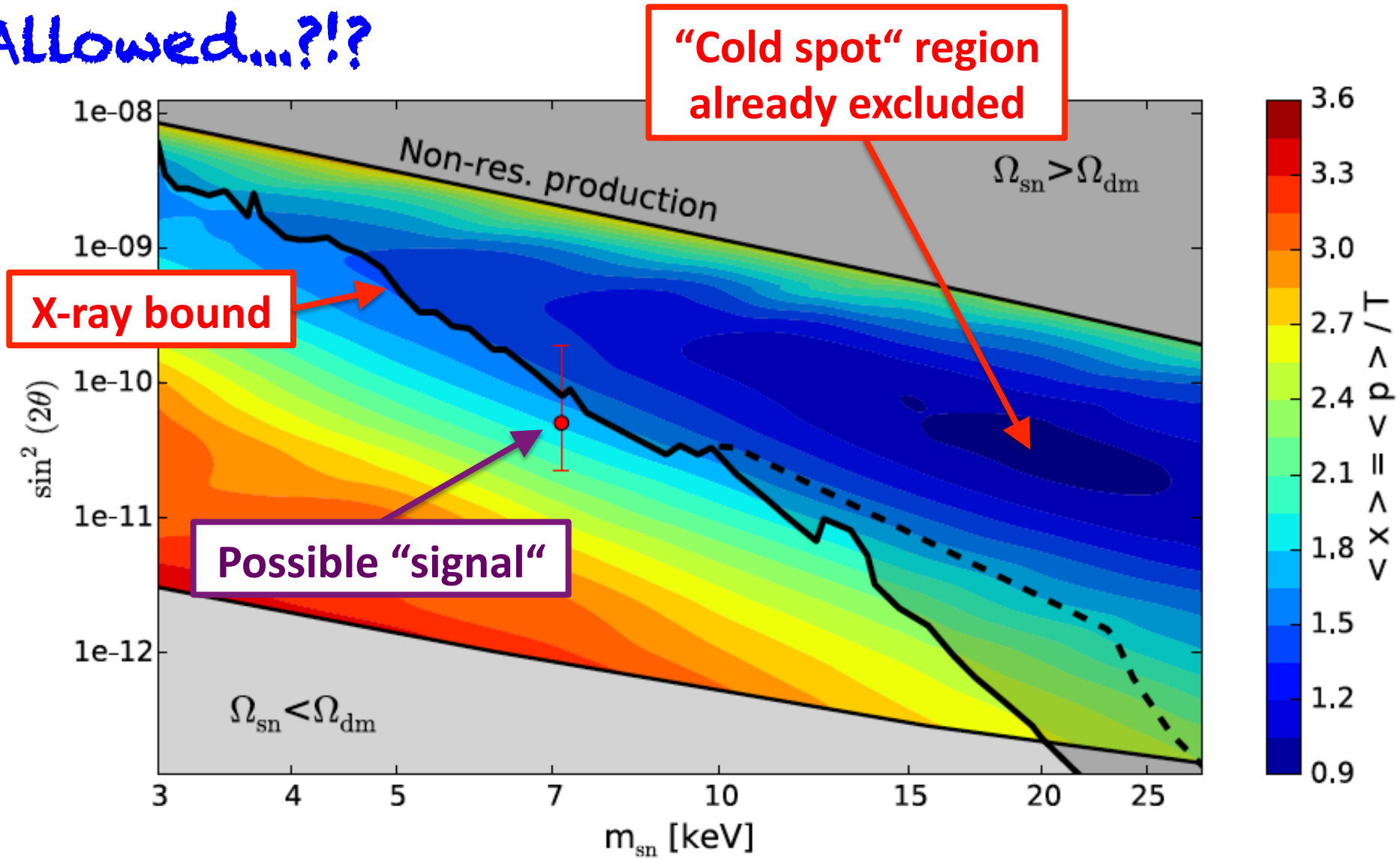
This produces to non-trivial spectra:

[Abazajian: Phys. Rev. Lett. 112 (2014) 161303]



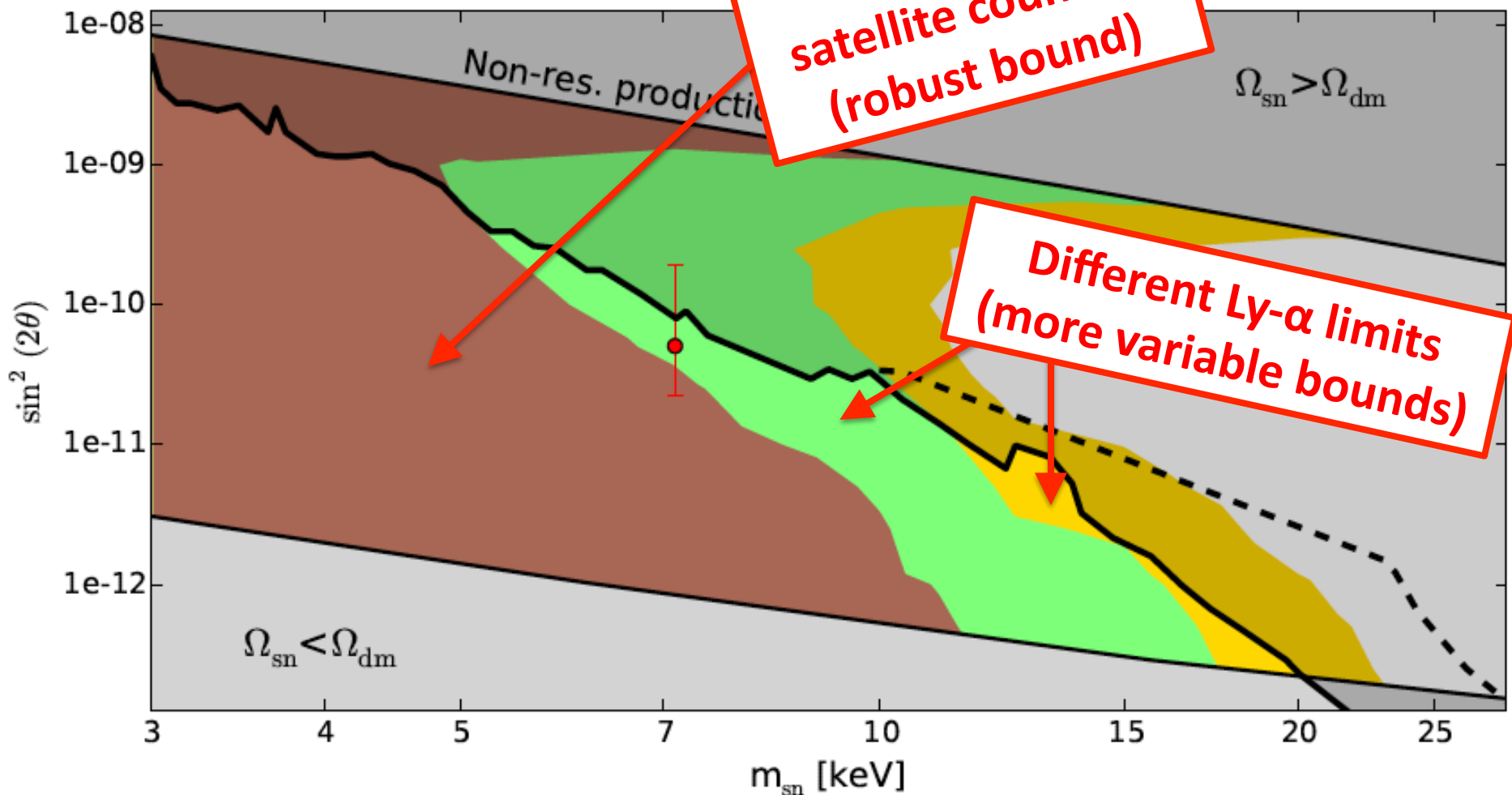
# The Shi-Fuller mechanism

Allowed...?!?



# The Shi-Fuller mechanism

... or threatened?!?



## 4. *Example B*: DW with initial abundance

## 4. Example B: DW with initial abundance

**BUT: DW is an unavoidable correction!!!**

*We have looked into this and, in passing, disproved two common prejudices about DW-production...*

[AM, Schneider, Totzauer: JCAP **1604** (2016) 003]

### **1st prejudice:**

**DW produces spectrum with thermal shape**

**- appears as approximation in the original papers:**

[Dodelson, Widrow: Phys. Rev. Lett. **72** (1994) 17; DW & Colombi: Astrophys. J. **458** (1996) 1]

$$f(q) = \frac{\beta_{\text{DW}}}{\exp(q/T_{\text{DW}}) + 1}$$

**- used in structure/galaxy formation computations:**

[Viel *et al.*: Phys. Rev. **D71** (2005) 063534; Herpich *et al.*: Mon. Not. Roy. Astron. Soc. **442** (2014) 176; Menci, Fiore, Lamastra: Mon. Not. Roy. Astron. Soc. **421** (2012) 2384; Lovell *et al.*: Mon. Not. Roy. Astron. Soc. **439** (2014) 300,...]

# 4. Example B: DW with initial abundance

We have found the **FULL SOLUTION** (semi-analytically)

[AM, Schneider, Totzauer: JCAP 1604 (2016) 003]

$$f_N(T_f, p) = \mathcal{S}(T_f, T_{\text{ini}}, T_f, p) \left[ f_{\text{ini}} \left( \frac{T_{\text{ini}}}{T_f} \left( \frac{g_S(T_{\text{ini}})}{g_S(T_f)} \right)^{1/3} p \right) + f_{\text{DW}}(T_f, T_{\text{ini}}, p) \right]$$

With:

$$\mathcal{S}(T_a, T_b, T_c, p) \equiv \exp \left[ \int_{T_a}^{T_b} dT_2 h \left( T_2, \frac{T_2}{T_c} \left( \frac{g_S(T_2)}{g_S(T_c)} \right)^{1/3} p \right) \right]$$

$$f_{\text{DW}}(T_a, T_b, p) \equiv - \int_{T_a}^{T_b} dT_2 \mathcal{S}^{-1}(T_2, T_b, T_a, p) (h f_{\text{th}}) \left( T_2, \frac{T_2}{T_f} \left( \frac{T_2 g_S(T_2)}{T_f g_S(T_f)} \right)^{1/3} p \right)$$

$$\mathcal{S}(T_f, T_{\text{ini}}, T_f, p) f_{\text{DW}}(T_f, T_{\text{ini}}, p) \stackrel{!}{=}$$

**"Stop the clock"**

$$\mathcal{S}(T_f, T_3, T_f, p) \left[ \mathcal{S} \left( T_3, T_{\text{ini}}, T_3, \frac{T_3}{T_f} \left( \frac{g_S(T_3)}{g_S(T_f)} \right)^{1/3} p \right) f_{\text{DW}} \left( T_3, T_{\text{ini}}, \frac{T_3}{T_f} \left( \frac{g_S(T_3)}{g_S(T_f)} \right)^{1/3} p \right) + \right.$$

$$\left. f_{\text{DW}}(T_f, T_3, p) \right],$$

**CONSERVATION LAW**

# 4. Example B: DW with initial abundance

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[AM, Schneider, Totzauer: JCAP 1604 (2016) 003]

$$f_N(T_f, p) = \mathcal{S}(T_f, T_{\text{ini}}, T_f, p) \left[ f_{\text{ini}} \left( \frac{T_{\text{ini}}}{T_f} \left( \frac{g_S(T_{\text{ini}})}{g_S(T_f)} \right)^{1/3} p \right) + f_{\text{DW}}(T_f, T_{\text{ini}}, p) \right]$$

coupling between both parts ( $\mathcal{S} \approx 1$ )

Redshifted initial distribution

DW-part

$$f_{\text{DW}}(T_a, T_{\text{ini}}, p) \equiv \exp \left[ \int_{T_a}^{T_b} dT_c \ln \left( \frac{T_2}{T_c} \left( \frac{g_S(T_2)}{g_S(T_c)} \right)^{1/3} \right) \right] f_{\text{DW}} \left( T_2, \frac{T_2}{T_f} \left( \frac{T_2 g_S(T_2)}{T_f g_S(T_f)} \right)^{1/3} p \right)$$

$$\mathcal{S}(T_f, T_{\text{ini}}, T_f, p) f_{\text{DW}}(T_f, T_{\text{ini}}, p) \stackrel{!}{=}$$

"Stop the clock"

$$\mathcal{S}(T_f, T_3, T_f, p) \left[ \mathcal{S} \left( T_3, T_{\text{ini}}, T_3, \frac{T_3}{T_f} \left( \frac{g_S(T_3)}{g_S(T_f)} \right)^{1/3} p \right) f_{\text{DW}} \left( T_3, T_{\text{ini}}, \frac{T_3}{T_f} \left( \frac{g_S(T_3)}{g_S(T_f)} \right)^{1/3} p \right) + f_{\text{DW}}(T_f, T_3, p) \right],$$

CONSERVATION LAW

## 4. Example B: DW with initial abundance

**BUT: DW is an unavoidable correction!!!**

*We have looked into this and, in passing, disproved two common prejudices about DW-production...*

[AM, Schneider, Totzauer: JCAP 1604 (2016) 003]

**1st prejudice:**

**DW produces spectrum with thermal shape**

**- in reality, even if  $g_S$  is taken constant, which it is NOT during DW production:**

$$f_N^{\text{DW}}(T_f, p) \approx \frac{1}{\exp\left(\frac{p}{T_f} \left(\frac{\langle g_S \rangle}{g_S(T_f)}\right)^{1/3}\right) + 1} \int_{T_{\text{ini}}}^{T_f} dT_2 h\left(T_2, \frac{T_2}{T_f} \left(\frac{g_S(T_2)}{g_S(T_f)}\right)^{1/3}, p\right)$$

This function  $h$  (contains active-sterile mixing, mass difference, etc.) and the whole integral needs to vary **SLOWLY** with the momentum  $p$ !

**→ not the case in reality... → NO THERMAL SHAPE!!!**



# 4. Example B: DW with initial abundance

**BUT: DW is an unavoidable correction!!!**

We have looked into this and, in passing, have uncovered two common prejudices about DW.

1st prejudice

DW

- it

depends

$f_N^{\text{DW}}$

$(\dots)$

$(\dots)$

$(\dots)$

$(\dots)$

**→ thus, all bounds based on the assumption of a thermal shape are, strictly speaking, wrong (how far from reality depends on the details...)**

This function  $h$  (contains active-sterile mixing, mass difference, etc.) and the whole integral needs to vary **SLOWLY** with the momentum  $p$ !

**→ not the case in reality... → NO THERMAL SHAPE!!!**

## 4. Example B: DW with initial abundance

**BUT: DW is an unavoidable correction!!!**

*We have looked into this and, in passing, disproved two common prejudices about DW-production...*

[AM, Schneider, Totzauer: JCAP 1604 (2016) 003]

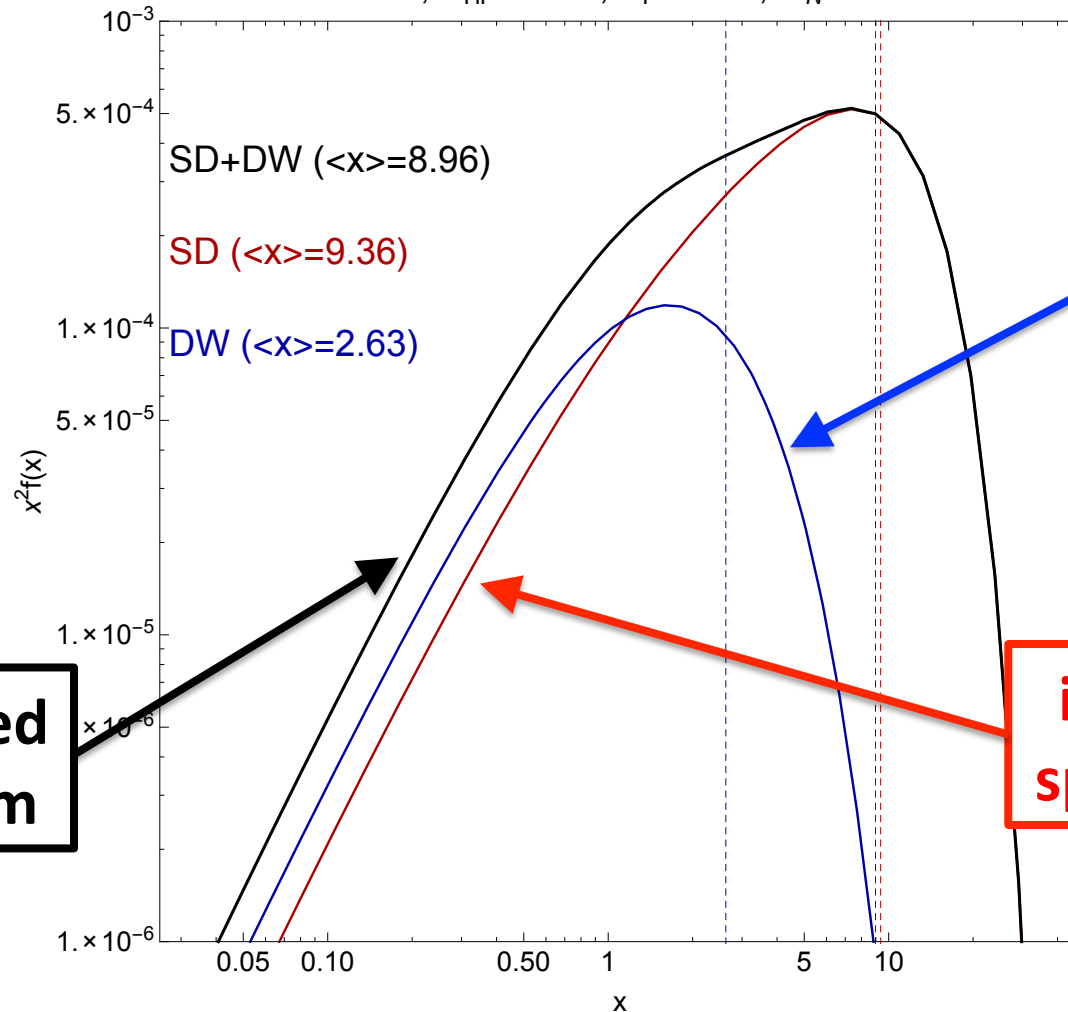
**2nd prejudice:**

**DW always produces a too hot spectrum**

- *different for **non-zero initial abundance**: if DM is already present before DW sets in, the spectrum may experience non-trivial modifications*
- *illustration: scalar decay production + **SUBSEQUENT** modification by the Dodelson-Widrow mechanism*

# 4. Example B: DW with initial abundance

$T=10$  MeV,  $C_{HP}=10^{-1.5}$ ,  $C_{\Gamma}=10^{-3.0}$ ,  $m_N=6.35$  keV



DW-modification

Combined spectrum

initial scalar decay spectrum (see later)

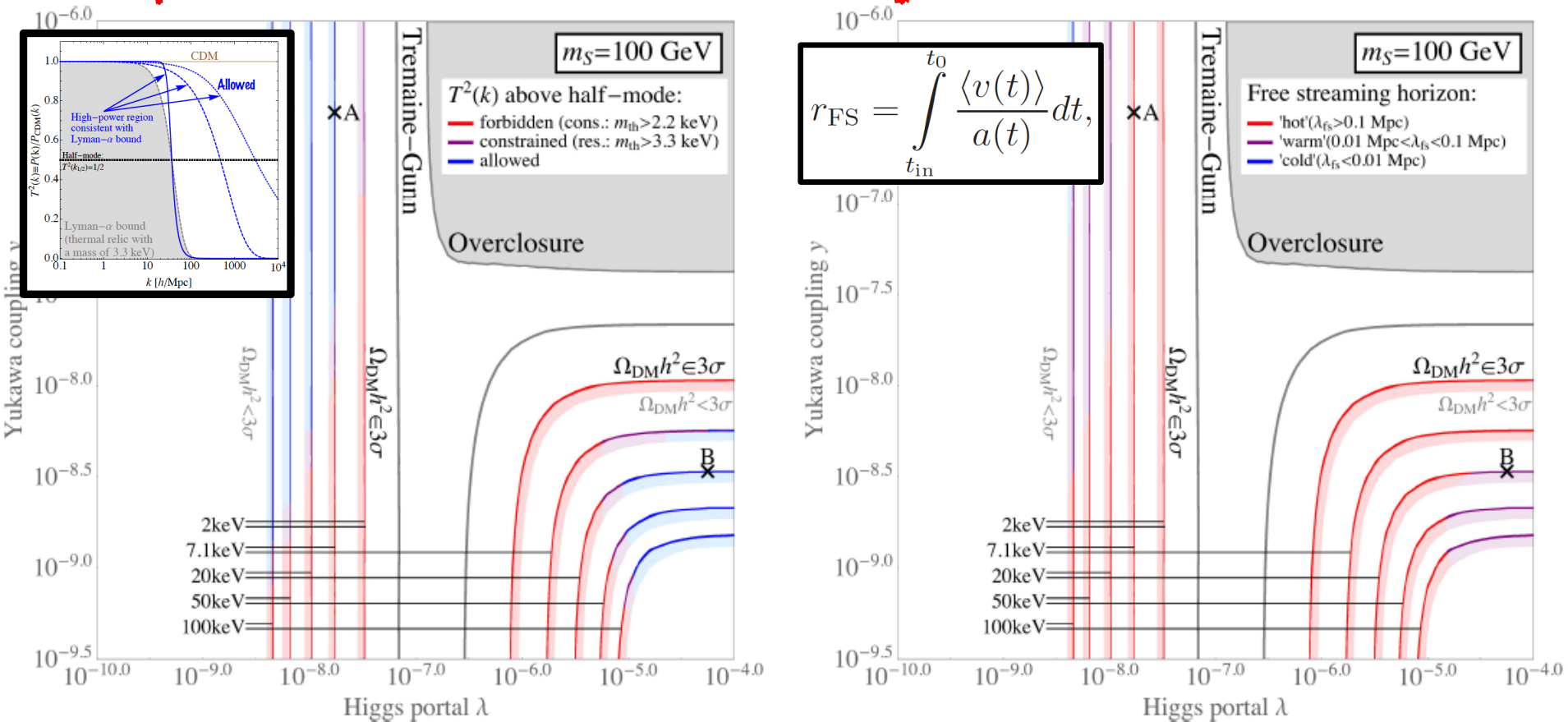
may shift the average momentum to LOWER values (“DW-cooling”)

**BUT: SMALL EFFECT!!!** (DW-part can at most be a  $\sim 25\%$  modification)

# 4. Example 2: Decay production

[König, AM, Totzauer: 1609.01289]

## Comparison: Half-mode analysis vs. free streaming

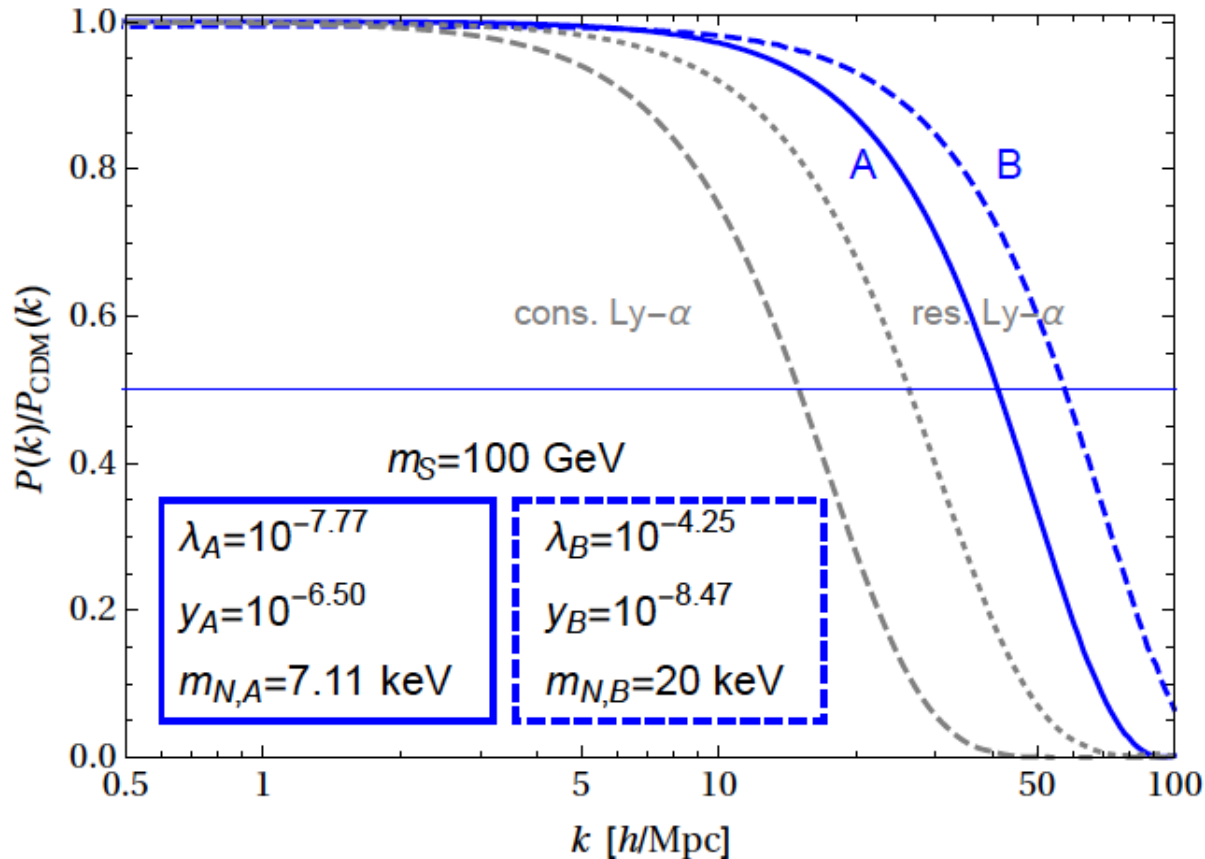


-> free-streaming would classify points A/B as "hot"/"warm" and discard at least one...

# 4. Example 2: Decay production

[König, AM, Totzauer: 1609.01289]

... while in reality both are not even touched by structure formation constraints!!!!

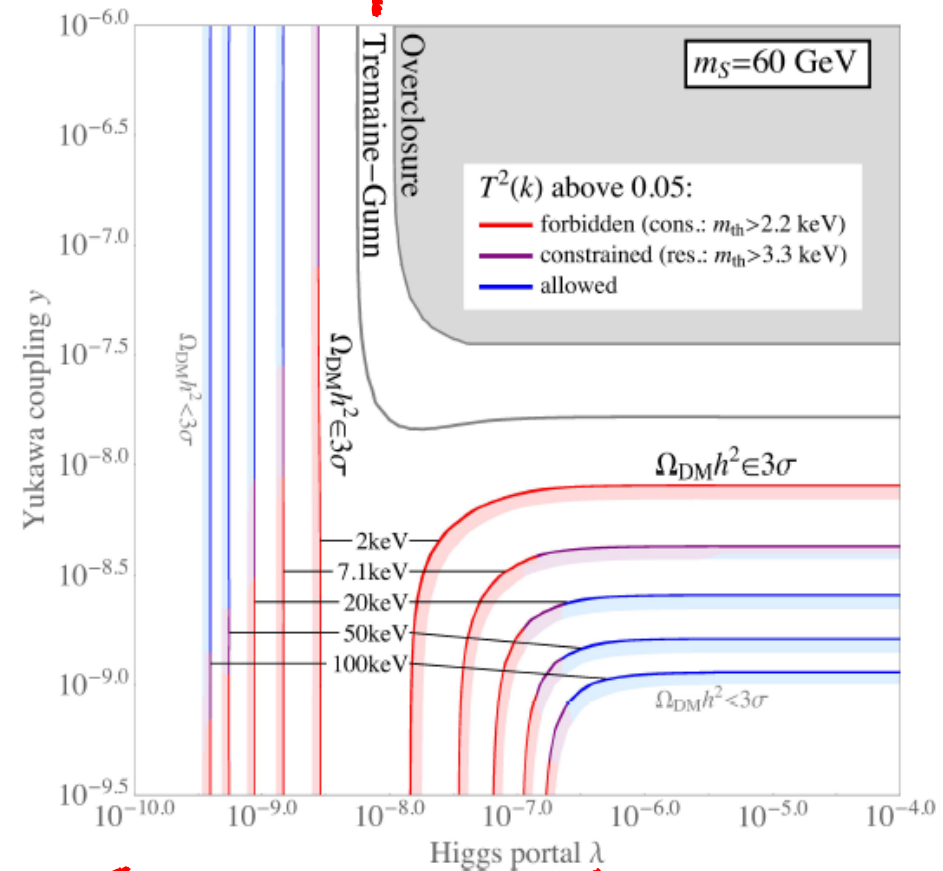
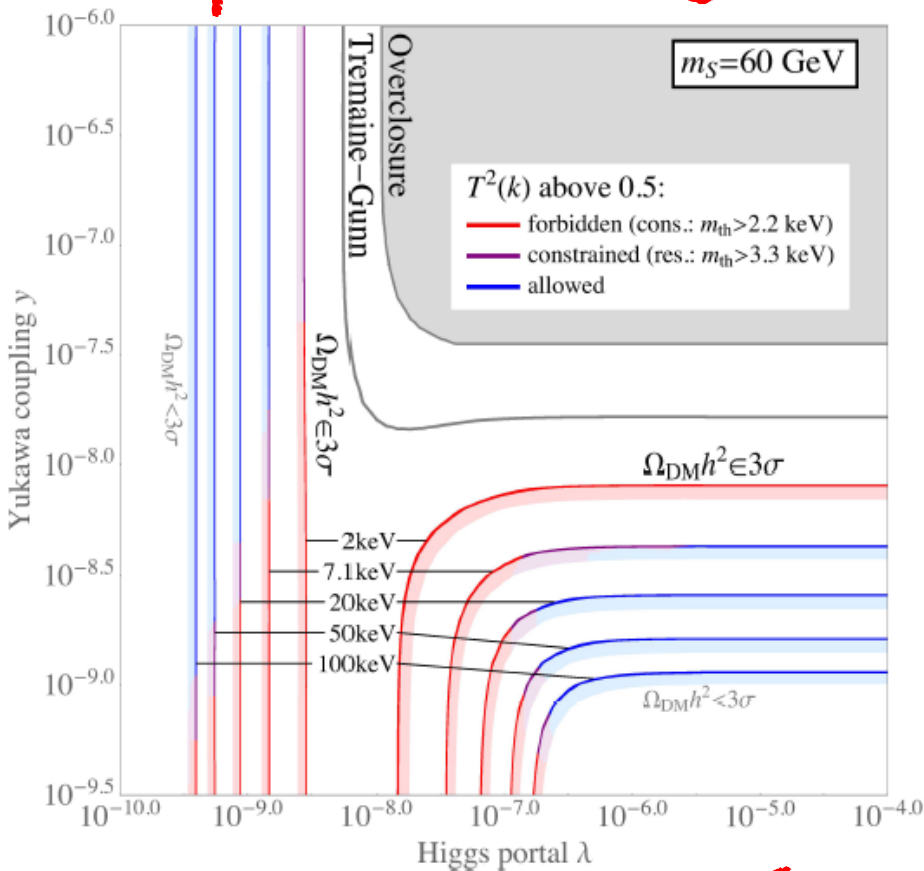


→ free-streaming cannot describe non-thermal distributions correctly!!!!

# 4. Example 2: Decay production

[König, AM, Totzauer: 1609.01289]

... whereas our analysis is very robust with respect to changes in the assumption:

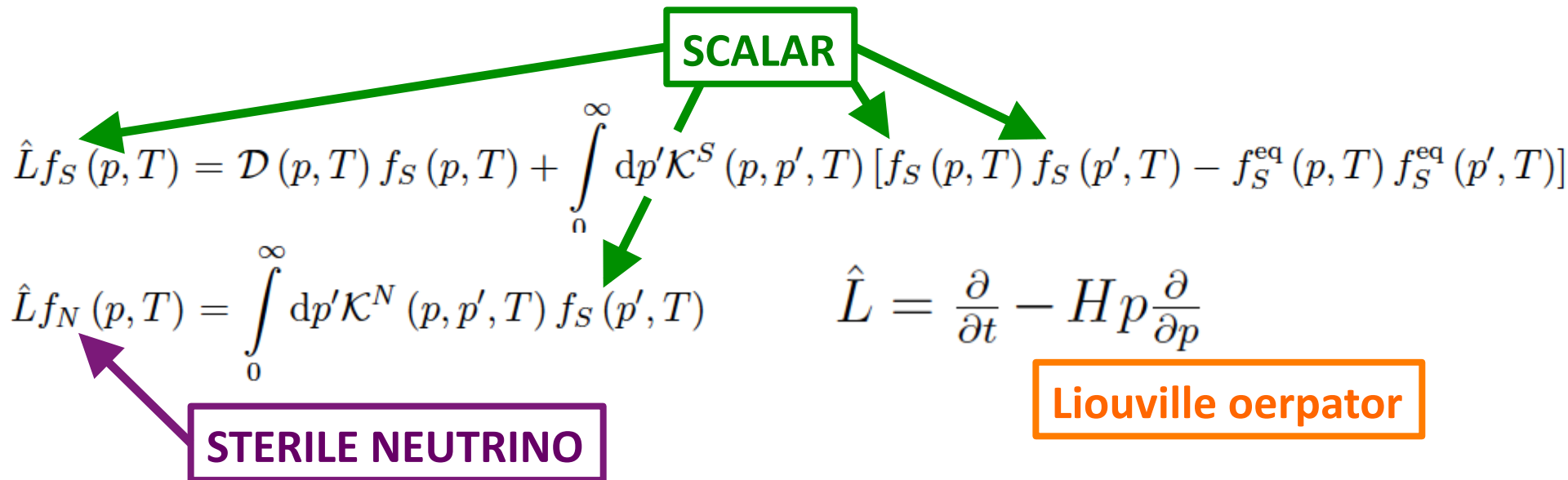


→ even a change of the reference point by one order of magnitude hardly affects our results!

# 4. *The one that works*: Decay production

[König, AM, Totzauer: JCAP 1611 (2016) 038]

- **TECHNICALITY:** How to crack a coupled system of non-linear partial integro-differential equations in two variables?



→ extremely tricky if not approached in the correct way (no known type of integro-differential equation, like e.g. Volterra-type, because of integrals extending over the whole domain in the momentum  $p$ )

# 4. *The one that works*: Decay production

[König, AM, Totzauer: JCAP **1611** (2016) 038]

- **Our strategy:**

- **family of new variables:**

$$\left. \begin{matrix} T \\ p \end{matrix} \right\} \rightarrow \left\{ \begin{matrix} r(T, p) = r(T) \\ \xi(T, p) = f \left( \frac{a[t(T)]}{a[t(T_0)]} p \right) \end{matrix} \right\} \rightarrow \left\{ \begin{matrix} r = m_0/T \\ \xi = \frac{p}{T} \left( \frac{g_s(T_0)}{g_s(T)} \right)^{1/3} \end{matrix} \right.$$

-> **equations decoupled, numerics stabilised ( $m_0$  &  $T_0$ !!)**

- **discretise the “momentum”  $\xi$  into fine slices, to turn the PDE into a system of ODEs:**

$$\frac{d}{dr} f_S^i(r) = \tilde{\mathcal{D}}^i(T) f_S^i(r) + \sum_{j=1, \dots, M} \tilde{\mathcal{K}}^{ij}(r) [f_S^i(r) f_S^j(r) - f_S^{i, \text{eq}}(r) f_S^{j, \text{eq}}(r)]$$

- **use a solver that is NOT list-based (like Mathematica), but matrix-based (like *ode15s* in Matlab) to exploit semi-analytical methods**



# 4. *The one that works*: Decay production

[König, AM, Totzauer: JCAP 1611 (2016) 038]

- **Our strategy:**

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$$\left. \begin{matrix} T \\ p \end{matrix} \right\} \rightarrow \left\{ \begin{matrix} r(T, p) = r(T) \\ \xi(T, p) = f\left(\frac{a[t(T)]}{a[t(T_0)]} p\right) \end{matrix} \right\} \rightarrow \left\{ r = m_0/T \right.$$

-> equations de

**THIS WAY, WE COULD SOLVE THE FULL SET OF EQUATIONS GOVERNING DARK MATTER PRODUCTION IN THE WHOLE PARAMETER SPACE**

return the PDE

$$\frac{d}{dt} \dots + \sum_{j=1, \dots, M} \tilde{\mathcal{K}}^{ij}(r) [f_S^i(r) f_S^j(r) - f_S^{i, \text{eq}}(r) f_S^{j, \text{eq}}(r)]$$

- **use a solver that is NOT list-based (like Mathematica), but matrix-based (like *ode15s* in Matlab) to exploit semi-analytical methods**

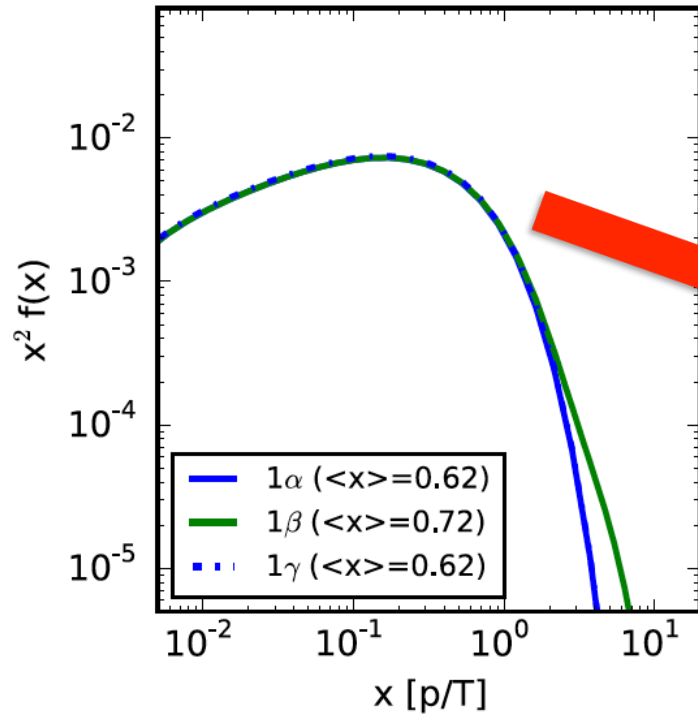
# 5. Structure Formation

- **OUR PROPOSAL** (for stupid particle guys only):

**Half-mode analysis** [König, AM, Totzauer: JCAP 1611 (2016) 038]

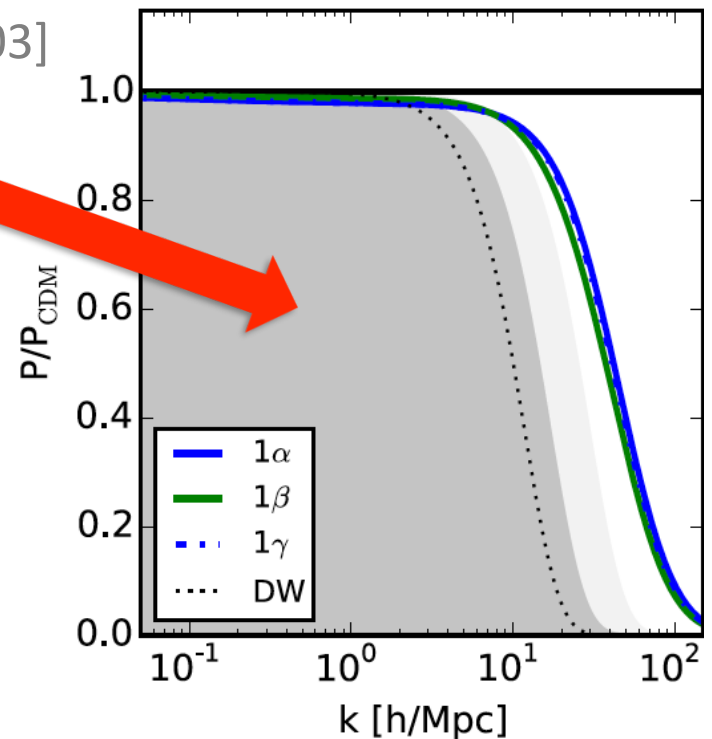
- **Linear power spectrum with CLASS**

→ [www.class-code.net](http://www.class-code.net): Lesgourgues & Tram



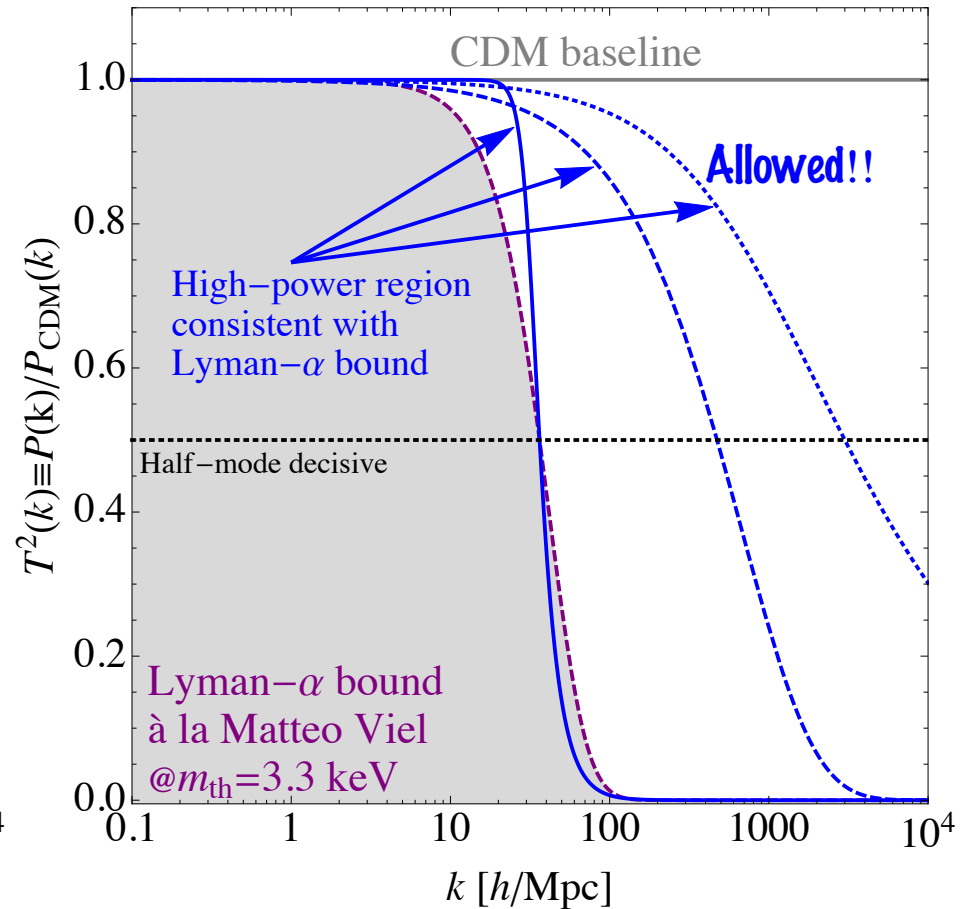
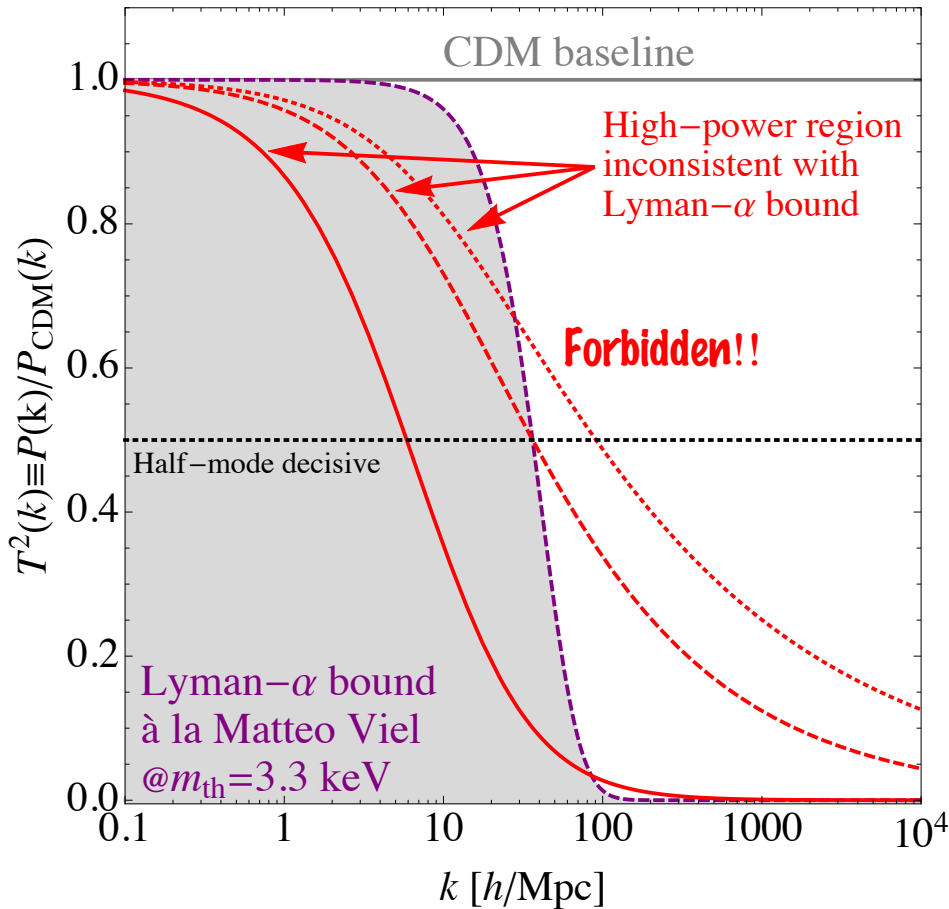
**INPUT: momentum distribution  $f(p,T)$**

**OUTPUT: (squared) transfer function**



# 5. Structure Formation

- use half-mode crossing, instead of free streaming: WORKS FOR ANY DISTRIBUTION, NO MATTER HOW NON-THERMAL!!



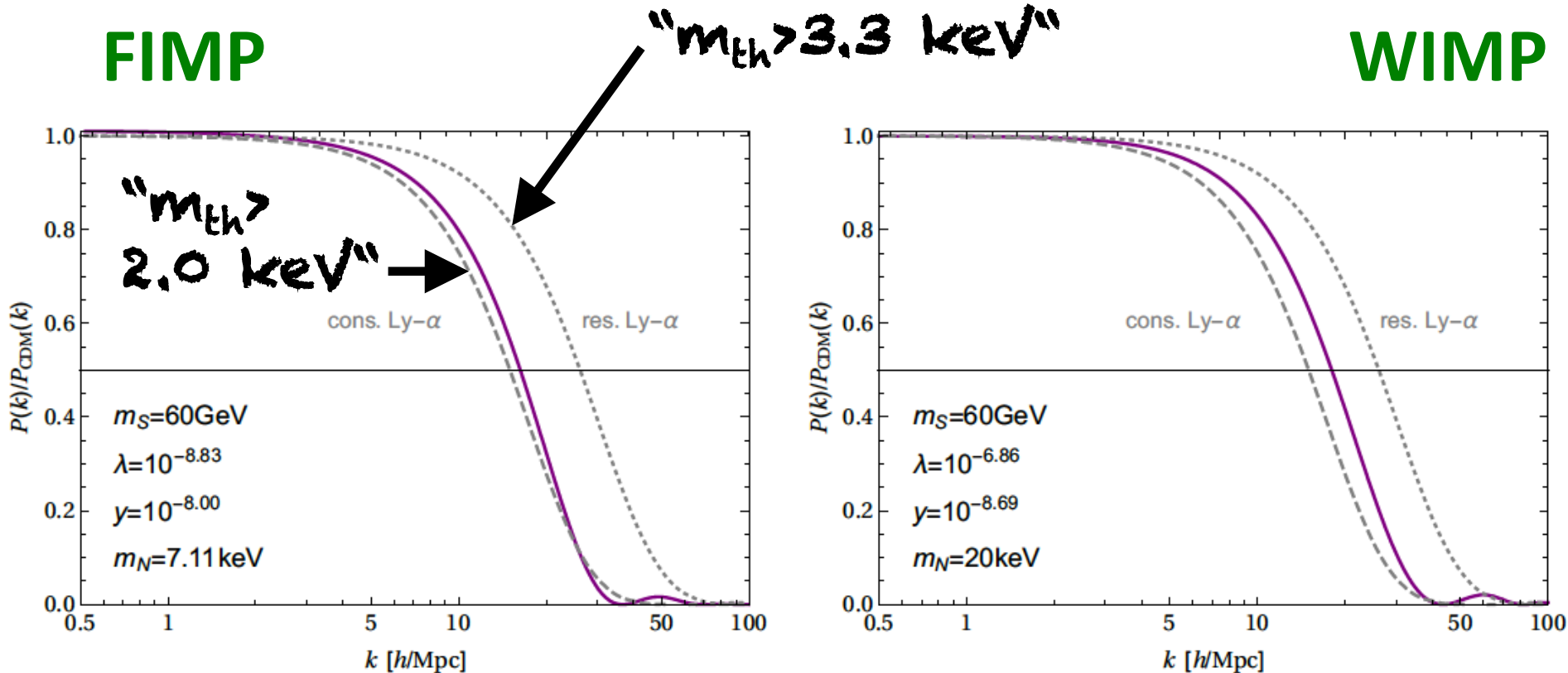
# Applied to Decay production

[König, AM, Totzauer: JCAP 1611 (2016) 038]

Ly-alpha Limits on the two example points:

FIMP

WIMP



→ both agree with conservative Ly-alpha bound but violate restrictive one  
(NB: thermal relic mass  $\neq$  physical mass!!!)

# Applied to Decay production

[König, AM, Totzauer: JCAP **1611** (2016) 038]

Allowed regions  
from Ly- $\alpha$  data:

Coservative ✘  
Restrictive ✘

Coservative ✔  
Restrictive ✘

Coservative ✔  
Restrictive ✔

