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Influence of the electromagnetic fields on hadronic observables in proton-induced collisions

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QCD PHASE DIAGRAM



High energy heavy ion collisions

- ✓ allow to experimentally investigate the QCD phase diagram
- recreate the extreme condition of temperature and density required to form the QUARK-GLUON PLASMA

Large Hadron Collider (LHC)



Relativistic Heavy Ion Collider (RHIC)



Facility for Antiproton and Ion Research (FAIR)



Nuclotron-based Ion Collider fAcility (NICA)



EXPANDING FIREBALL

the evolution lasts about $t \sim 10-20 \text{ fm/c} \sim 10^{-23} \text{ s}$

initial temperature is about $T \sim 300-600 \text{ MeV} \sim 10^{12} \text{ K}$

Quark-Gluon Plasma (QGP)

an "almost perfect fluid" with very low viscosity and the formation of collective flows

Anisotropic radial flow is described by the Fourier coefficients of the azimuthal particle distributions w.r.t. the reaction plane p_y

$$\frac{\mathrm{d}n}{\mathrm{d}\phi} \propto 1 + \sum_{n} 2v_n(p_T) \cos[n(\phi - \Psi_n)]$$



 p_{x}

QGP initially expected only in high energy collisions of two heavy ions Small colliding systems initially regarded as control measurements



 \mathcal{X}

Signatures of collective flow found in small systems p+Pb collisions at LHC, p/d/³He+Au at RHIC

COLLECTIVITY IN SMALL SYSTEMS AS SIGN OF QGP DROPLETS?





collision overlap zone

PHENIX Collaboration, 1805.02973

 \mathcal{X}



Intense magnetic field $eB_v \sim 5-50 \ \underline{m_\pi^2} \sim 10^{18} \cdot 10^{19} \ G$

Kharzeev, McLerran and Warringa, NPA 803 (2008) 227 Skokov, Illarionov and Toneev, IJMPA 24 (2009) 5925



Earth's magnetic field ~ 1 G



laboratory ~ 10⁶ G

magnetar ~ 10¹⁴-10¹⁵ G

A consistent non-equilibrium transport approach to study heavy ion collisions (HICs) on a miscoscopic level

Cassing and Bratkovskaya, PRC 78 (2008) 034919; NPA831 (2009) 215 Cassing, EPJ ST 168 (2009) 3; NPA856 (2011) 162





more details by Wolfgang Cassing

GOAL

study the phase transition from hadronic to partonic matter and the properties of the quark gluon plasma from a microscopic origin

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 string formation in primary nucleon-nucleon collisions

 string decay to pre-hadrons (baryons and mesons)



INITIAL A+A COLLISIONS

nucleon-nucleon collisions between the two incoming nuclei lead to the formation of strings that decay to pre-hadrons

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- ➤ the Dynamical Quasi-Particle Model (DQPM) defines parton spectral functions, i.e. masses $M_{q,g}(\varepsilon)$ and widths $\Gamma_{q,g}(\varepsilon)$
- $\succ \text{ mean-field potential } U_q \text{ at given } \varepsilon$ related by 1QCD EoS to the local temperature



FORMATION OF QUARK-GLUON PLASMA if the energy density is above the critical value pre-hadrons dissolve in massive quarks and gluons



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PARTONIC STAGE

evolution based on off-shell transport equations and the Dynamical Quasi-Particle Model (DQPM)



- quarks and gluons as 'dynamical quasiparticles' with off-shell spectral functions
- self-generated mean-field potential
- Equation of state from lattice QCD
- (quasi-)elastic and inelastic partonparton interactions



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- massive off-shell quarks and antiquarks with broad spectral functions hadronize to off-shell mesons and baryons or strings
- ➢ local covariant off-shell transition rate for $q + \bar{q}$ fusion which lead to meson formation

HADRONIZATION

massive off-shell quarks with broad spectral functions hadronize to off-shell mesons and baryons



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PHST

- off-shell propagation
- elastic and inelastic hadron-hadron interactions



HADRONIC PHASE

evolution based on off-shell transport equations with hadron-hadron interactions

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PHSI

FINAL OBSERVABLES

good description of bulk observables (rapidity and transverse momentum distributions, flow coeficients, ...) for A+A collisions from SPS to LHC energies

PHSD + electromagnetic fields

PHSD has been extended including the dynamical formation and evolution of the retarded electomagnetic field (EMF) and its influence on the quasi-particle (QP) dynamics

Voronyuk *et al.*, PRC 83 (2011) 054911 Toneev *et al.*, PRC 85 (2012) 034910; PRC 86 (2012) 064907; PRC 95 (2017) 034911

TRANSPORT EQUATION

$$\left\{\frac{\partial}{\partial t} + \left(\frac{\mathbf{p}}{p_0} + \nabla_{\mathbf{p}} U\right) \nabla_{\mathbf{r}} + (-\nabla_{\mathbf{r}} U + e\mathbf{E} + e\mathbf{v} \times \mathbf{B}) \nabla_{\mathbf{p}} \right\} f = C_{\text{coll}}(f, f_1, \dots, f_N)$$

Lorentz force

MAXWELL EQUATIONS

$$\nabla \cdot \mathbf{B} = 0 \qquad \nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} \qquad \nabla \cdot \mathbf{E} = 4\pi\rho \qquad \nabla \times \mathbf{B} = \frac{\partial \mathbf{E}}{\partial t} + \frac{4\pi}{c}\mathbf{j}$$

charge distribution

electric current

consistent solution of particle and field evolution equations



Retarded electromagnetic fields

$$\mathbf{B} = \nabla \times \mathbf{A}, \quad \mathbf{E} = -\nabla \Phi - \frac{\partial \mathbf{A}}{\partial t}$$

General solution of the wave equation for the electromagnetic potentials

$$\begin{aligned} \mathbf{A}(\mathbf{r},t) &= \frac{1}{4\pi} \int \frac{\mathbf{j}(\mathbf{r}',t') \ \delta(t-t'-|\mathbf{r}-\mathbf{r}'|/c)}{|\mathbf{r}-\mathbf{r}'|} \ d^3r' dt' \\ \Phi(\mathbf{r},t) &= \frac{1}{4\pi} \int \frac{\rho(\mathbf{r}',t') \ \delta(t-t'-|\mathbf{r}-\mathbf{r}'|/c)}{|\mathbf{r}-\mathbf{r}'|} \ d^3r' dt' \end{aligned}$$

$$\mathbf{r}' \equiv \mathbf{r}(t')$$
$$t' = t - \frac{\mathbf{r} - \mathbf{r}'}{c}$$

Liénard-Wiechert potentials for a moving point-like charge

$$\Phi(\mathbf{r},t) = \frac{e}{4\pi} \left[\frac{1}{R(1-\mathbf{n}\cdot\boldsymbol{\beta})} \right]_{\text{ret}} \qquad \mathbf{A}(\mathbf{r},t) = \frac{e}{4\pi} \left[\frac{\boldsymbol{\beta}}{R(1-\mathbf{n}\cdot\boldsymbol{\beta})} \right]_{\text{ret}}$$

 $\mathbf{R} = \mathbf{r} - \mathbf{r}'$ $\mathbf{n} = \frac{\mathbf{R}}{R}$ $\boldsymbol{\beta} = \frac{\mathbf{v}}{c}$

ret: evaluated at the times t'

Voronyuk et al., PRC 83 (2011) 054911

Retarded electromagnetic fields

Retarded electric and magnetic fields for a moving point-like charge

$$\mathbf{E}(\mathbf{r},t) = \frac{e}{4\pi} \left[\frac{\mathbf{n} - \mathbf{\beta}}{(1 - \mathbf{n} \cdot \mathbf{\beta})^3 \gamma^2 R^2} + \frac{\mathbf{n} \times \left((\mathbf{n} - \mathbf{\beta}) \times \dot{\mathbf{\beta}} \right)}{(1 - \mathbf{n} \cdot \mathbf{\beta})^3 c R} \right]_{\text{ret}} \quad \mathbf{B}(\mathbf{r},t) = [\mathbf{n} \times \mathbf{E}(\mathbf{r},t)]_{\text{ret}}$$
elastic Coulomb inelastic bremsstrahlung

Neglecting the acceleration

$$e\mathbf{E}(t, \mathbf{r}) = \alpha_{em} \frac{1 - \beta^2}{\left[(\mathbf{R} \cdot \boldsymbol{\beta})^2 + R^2 (1 - \beta^2) \right]^{3/2}} \mathbf{R}$$
$$e\mathbf{B}(t, \mathbf{r}) = \alpha_{em} \frac{1 - \beta^2}{\left[(\mathbf{R} \cdot \boldsymbol{\beta})^2 + R^2 (1 - \beta^2) \right]^{3/2}} \boldsymbol{\beta} \times \mathbf{R}$$

Voronyuk et al., PRC 83 (2011) 054911

magnetic field created by a single freely moving charge



EM fields in nuclear collisions

in a nuclear collision the magnetic field is a superposition of solenoidal fields from different moving charges

Voronyuk et al. (PHSD team), PRC 83 (2011) 054911



Au+Au @RHIC 200 GeV - b = 10 fm



EM fields in nuclear collisions

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Voronyuk et al. (PHSD team), PRC 83 (2011) 054911



Au + Au @RHIC 200 GeV - b = 10 fm



 $RHIC 200 \ GeV - b = 7 \ fm$





SYMMETRIC SYSTEMS (Au+Au, Pb+Pb)

partial compensation of electric and magnetic forces

ASYMMETRIC SYSTEMS (e.g. Cu+Au)

electric field strongly asymmetric inside the overlap region

Voronyuk et al. (PHSD team), PRC 90, 064903 (2014)

EM fields in proton-induced collisions



EM fields in proton-induced collisions



Centrality in heavy ion collisions

Centrality characterizes the amount of overlap or size of the fireball in the collision region

e.g. (MC-)Glauber model

 $\begin{array}{l} \text{INITIAL STATE QUANTITIES} \\ \text{b, } N_{\text{part}}, \{N_{\text{part}}, N_{\text{coll}}\}, N_{\text{qp}} \end{array}$

FINAL STATE OBSERVABLES N_{ch}, E_T, N_{neutron}



from talk of Jiangyong Jia at MIAPP (2018)

CENTRALITY FLUCTUATION

- main uncertainty for many measurements
- Iarge in peripheral collisions or small collision systems

PROBABILITY DISTRIBUTION IN THE NUMBER OF PARTICIPANTS AND CHARGED HADRON MULTIPLICITY AT MIDRAPIDITY





➤ correlation between $N_{ch}(|\eta| < 0.5)$ and N_{part}

 large dispersion respect to AA collisions

1

PROBABILITY DISTRIBUTION IN THE NUMBER OF PARTICIPANTS AND CHARGED HADRON MULTIPLICITY AT MIDRAPIDITY p+Pb collisions @ LHC 5.02 TeV



Konchakovski, Cassing and Toneev JPG 41, 105004 (2014)



CENTRALITY SELECTION FROM MINIMUM BIAS EVENTS



η <2	$< dN_{ch}/d\eta>$
0-5%	11.9
5-10%	8.9
10-20%	7.2
20-40%	5.1
40-60%	3.5
60-80%	2.2
80-100%	0.7

$$\eta = \frac{1}{2} \log \frac{1 + \cos \theta}{1 - \cos \theta}$$



Adare et al. (PHENIX Collaboration), 1807.11928



- enhanced particle production in the Au-going directions
- asymmetry increases
 with centrality of
 the collision





Experimental data: Aidala et al. (PHENIX Collaboration), PRC 95 (2017) 034910



Experimental data: Adare et al. (PHENIX Collaboration), PRC 97 (2018) 064904



CONCLUDING....

- □ The Parton-Hadron-String-Dynamics (PHSD) describes the entire dynamical evolution of heavy ion collisions within one single theoretical framework
- PHSD has been extended to include in a consistent way the intense electromagnetic fields produced in the very early stage of the collision
- □ Preliminary study of p+Au collisions at top RHIC energy:
 - ✓ the electric field is strongly asymmetric inside the overlap region
 - ✓ asymmetry of charged-particle rapidity distributions increasing with centrality

...LOOKING FORWARD



- □ Evolution dynamics and properties of the matter created in small and asymmetric systems (e.g. p+Au, d+Au, ³He+Au @ RHIC, p+Pb @ LHC)
- Influence of the intense electric field created in asymmetric collisions on the formation of the quark-gluon plasma

Thank you for your attention!