

In-Medium Dynamics of ω -Mesons

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ω -Mesons in Free Space

$$J^{PC} (I^G) = 1^{--} (0^-)$$

$$m_\omega = (781.94 \pm 0.12) \text{ MeV} \quad \Gamma_0 = (8.41 \pm 0.09) \text{ MeV}$$

$$\omega \rightarrow \pi^0 \pi^+ \pi^- \text{ and } \omega \rightarrow \rho \pi$$

$\omega \rightarrow$	Γ_i/Γ_0 in %	Γ_i in MeV
$\pi^0 \pi^+ \pi^-$	88.8 ± 0.7	7.486
$\pi^0 \gamma$	8.5 ± 0.5	0.717
$\pi^+ \pi^-$	2.21 ± 0.3	0.186
$\eta \gamma$	$(6.5 \pm 1.0) \cdot 10^{-4}$	$5.47 \cdot 10^{-3}$
$\pi^0 e^+ e^-$	$(5.9 \pm 1.9) \cdot 10^{-4}$	$5.0 \cdot 10^{-3}$
$\pi^0 \mu^+ \mu^-$	$(9.6 \pm 2.3) \cdot 10^{-5}$	$8.0 \cdot 10^{-4}$
$e^+ e^-$	$(7.07 \pm 0.19) \cdot 10^{-5}$	$5.9 \cdot 10^{-4}$

G-parity forbidden,
 $\omega \leftrightarrow \rho$ mixing

Why exposing Mesons to Nuclear Matter?

Nuclear physics – study of structure of matter in all its forms

- Most of the mass and energy in the universe around us comes from nuclei and nuclear reactions.
- The nucleus is a unique form of matter in that all the forces of nature are present : (strong, electromagnetic, weak, gravity).

History of the Universe

About 1 second after the Big Bang, protons and neutrons were formed

In today's universe, 99% visible matter are atomic nuclei (protons and neutrons).

Component	Percentage
Dark Energy	71.4%
Dark Matter	24%
Atoms	4.6%

atom $\sim 10^{-9}$ cm

nucleus $\sim 10^{-12}$ cm

proton (neutron) $\sim 10^{-13}$ cm

quark $< 10^{-16}$ cm

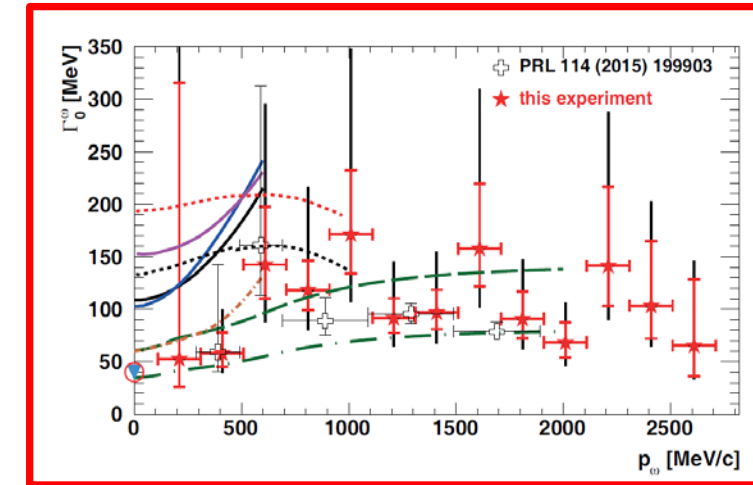
electron $< 10^{-19}$ cm

Brookhaven National Laboratory

c/o Haiyan Gao

Agenda

- Mesons in a Nuclear Environment
- Lagrangian for ω NN and ω NN* Interactions
- N^*N^{-1} particle-hole Polarization Modes
- ω -Mesons in Infinite Nuclear Matter
- ω -Mesons in a ^{93}Nb Nucleus
- $\omega + ^{93}\text{Nb}$ Bound States
- Outlook



CBELSA/TAPS Data:
Friedrich et al., Eur. Phys. J. A (2016) 52: 297

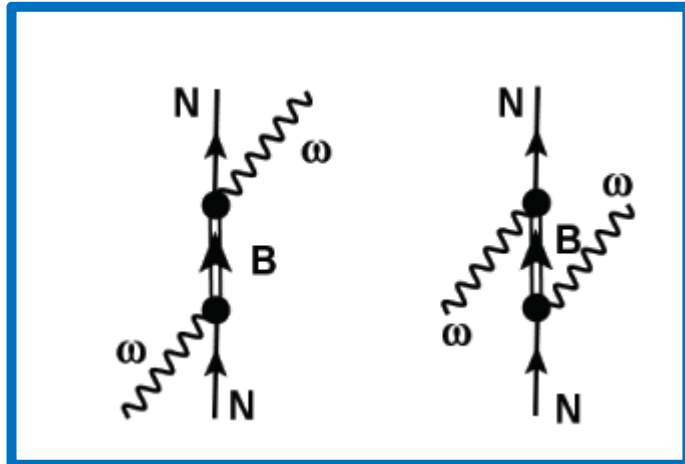
W. Peters, H.L. et al.
The spectral function of the **rho meson** in nuclear matter
Nucl. Phys. A 632 (1998) 109

H.L.
Interactions of **ω Mesons** in Nuclear Matter and with Nuclei
Eur. Phys. J A (in print)

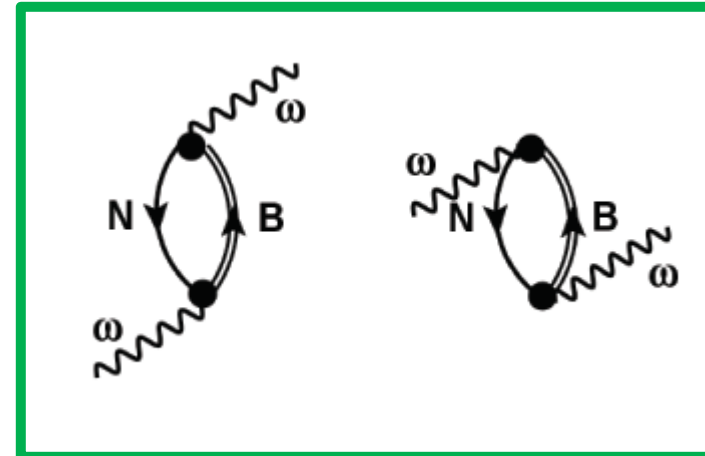
Polarization Tensors and Self-Energies

ω -Mesons interacting with Nuclear Matter: The Leading Order Processes

$$\approx O(g_{\omega NB}^2)$$



$$T_{\omega N} = V_{\omega N} + V_{\omega N} G_{\omega N} T_{\omega N}$$



$$\Pi_{\alpha\beta}^{\mu\nu} = -tr_A \langle A | \hat{\Gamma}_{\beta}^{\mu} \mathcal{G}_A \hat{\Gamma}_{\alpha}^{\nu} | A \rangle$$

The ω -nucleon scattering amplitudes (left) and the corresponding particle-hole type diagrams (right). Nucleons and hole states are denoted by N. The intermediate $B = N, N^*$ states (left), which become particle states (right), are indicated by double line

NN⁻¹ and N*N⁻¹ Polarization Tensors and Self-Energies in Infinite Matter (B=N,N*)

$$\Pi_{NB}^{\mu\nu}(w, \mathbf{q}) = \int \frac{d^3k}{(2\pi)^3} \frac{\theta(k_{FN}^2 - \mathbf{k}^2)\theta((\mathbf{k} + \mathbf{q})^2 - k_{FB}^2)}{2E_N^*(\mathbf{k})} \text{Tr}_s(\Gamma^\mu(\not{k}^* + M_N^*)\Gamma^\nu G_B(\mathbf{k} + \mathbf{q}|k_F))$$

Interaction Vertices

$$\begin{aligned} \Gamma_{NB}^\mu &= \gamma^\mu && \text{if B=N,N* is a positive parity state,} \\ \Gamma_{NB}^\mu &= \gamma_5 \gamma^\mu && \text{if B=N* is a negative parity state,} \end{aligned}$$

Omega In-Medium Self-Energy Tensor

$$\mathcal{S}_{\omega A}^{\mu\nu}(w, \mathbf{q}) = \sum_{N;B} g_{\omega NB}^2 \Pi_{NB}^{\mu\nu}(w, \mathbf{q}),$$

Longitudinal and Transversal In-Medium Self-Energies

$$\mathcal{S}^{\mu\nu}(w, \mathbf{q}) = (P_L^{\mu\nu} + P_T^{\mu\nu})^2 \mathcal{S}^{\mu\nu}(w, \mathbf{q}) = P_L^{\mu\nu} \Sigma_L(w, \mathbf{q}) + P_T^{\mu\nu} \Sigma_T(w, \mathbf{q})$$

$$\Sigma_{L/T}(w, \mathbf{q}) = P_{L/T}^{\mu\nu} \mathcal{S}_{\mu\nu}(w, \mathbf{q}) = \sum_{N=p,n;B} g_{BN\omega}^2 P_{L/T}^{(NB)}(w, \mathbf{q}).$$

$$P_{L/T}^{(NB)}(w, \mathbf{q}) = P_{L/T}^{\mu\nu} \Pi_{\mu\nu, NB}(w, \mathbf{q}).$$

ω -Meson Interactions on ^{93}Nb

Choice of States and Methods

- NN^{-1} and N^*N^{-1} with P-wave (positive parity) and S-wave (negative parity) resonances, $M_{N^*} < 2.5$ GeV
- Only isospin $I=\frac{1}{2}$ resonances are allowed
- $g_{N^*N\omega}$ coupling constants are determined by fit to data

Particle	Mass/MeV	total Width/MeV	$g_{N^*N\omega}$	M^*/M
ω	782.7	8.68	–	1
$P_{11}(940)$	938.92	–	2.076	0.60
$P_{11}(1440)$	1386	350	0.218	0.72
$P_{11}(1710)$	1670	140	2.163	0.78
$P_{11}(1880)$	1880	230	$< 10^{-3}$	0.80
$P_{11}(2100)$	2100	260	$< 10^{-3}$	0.82
$P_{11}(2300)$	2300	340	1.714	0.84
$P_{13}(1720)$	1720	225	0.022	0.78
$P_{13}(1920)$	1920	215	0.017	0.80
$S_{11}(1535)$	1535	150	$< 10^{-3}$	0.76
$S_{11}(1650)$	1650	125	0.008	0.77
$S_{11}(1895)$	1880	120	0.006	0.80
$S_{13}(1520)$	1520	110	0.003	0.75
$S_{13}(1700)$	1700	200	0.011	0.78
$S_{13}(1875)$	1875	200	1.554	0.70
$S_{13}(2120)$	2120	300	0.008	0.82

In-Medium Dynamics, Self-Energies, and Widths

- N, N* mean-field dynamics \rightarrow static (time-like) vector fields (v) and scalar fields (s), i.e. $M_B \rightarrow M_B^*(\rho)$
- Mean-field couplings $f_{N^*Na} = f_{NNa}$ for $a=s,v$
- P-wave N* lead to purely transversal self-energies
- S-wave N* lead to longitudinal and transversal self-energies

$$\Sigma_L(\omega, \mathbf{q}) \equiv \Sigma_L^{(S)}(\omega, \mathbf{q})$$

$$\Sigma_T(\omega, \mathbf{q}) = \Sigma_T^{(S)}(\omega, \mathbf{q}) + \Sigma_T^{(P)}(\omega, \mathbf{q})$$

$$\Gamma_{L/T}(\mathbf{q}) = -\frac{1}{m_\omega} \text{Im} (\Sigma_{L/T}(E(\mathbf{q}), \mathbf{q}))$$

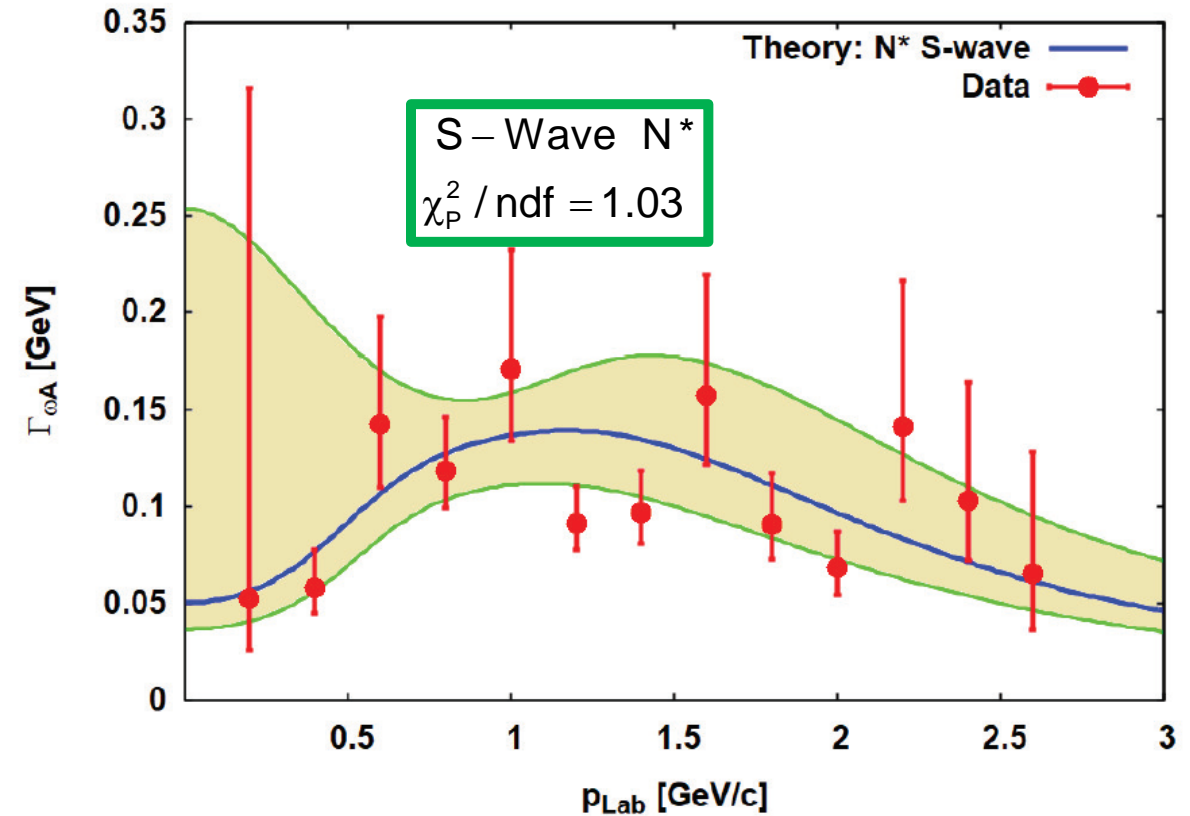
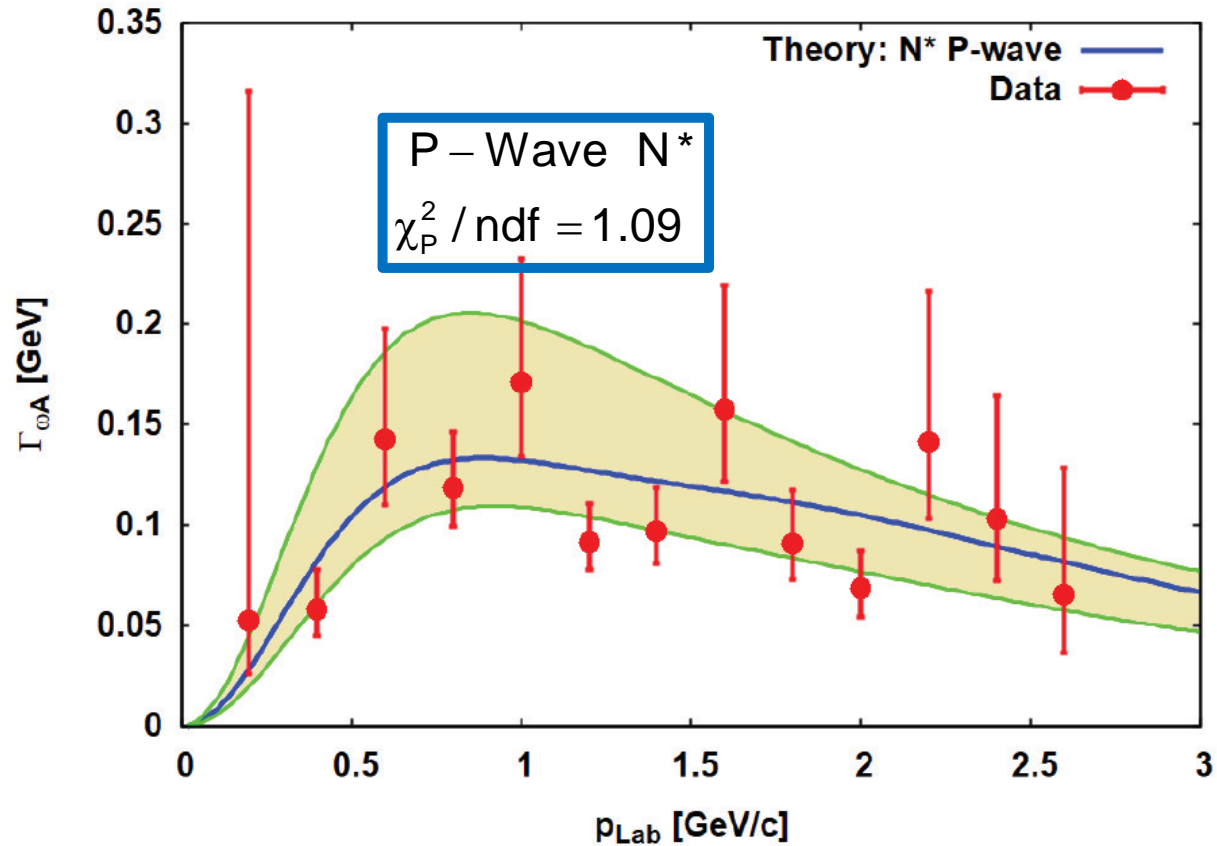
$$\omega = \sqrt{q^2 + m_\omega^2} \quad ; \quad q \rightarrow 0:$$

$$\Sigma_L^{(S)}(\omega, \mathbf{q}) \approx \text{const.}$$

$$\Sigma_T^{(S)}(\omega, \mathbf{q}) \approx \left(1 + \frac{q^2}{m_\omega^2}\right) \text{const.} \quad ; \quad \Sigma_T^{(P)}(\omega, \mathbf{q}) \approx \frac{q^2}{m_\omega^2} \text{const}$$

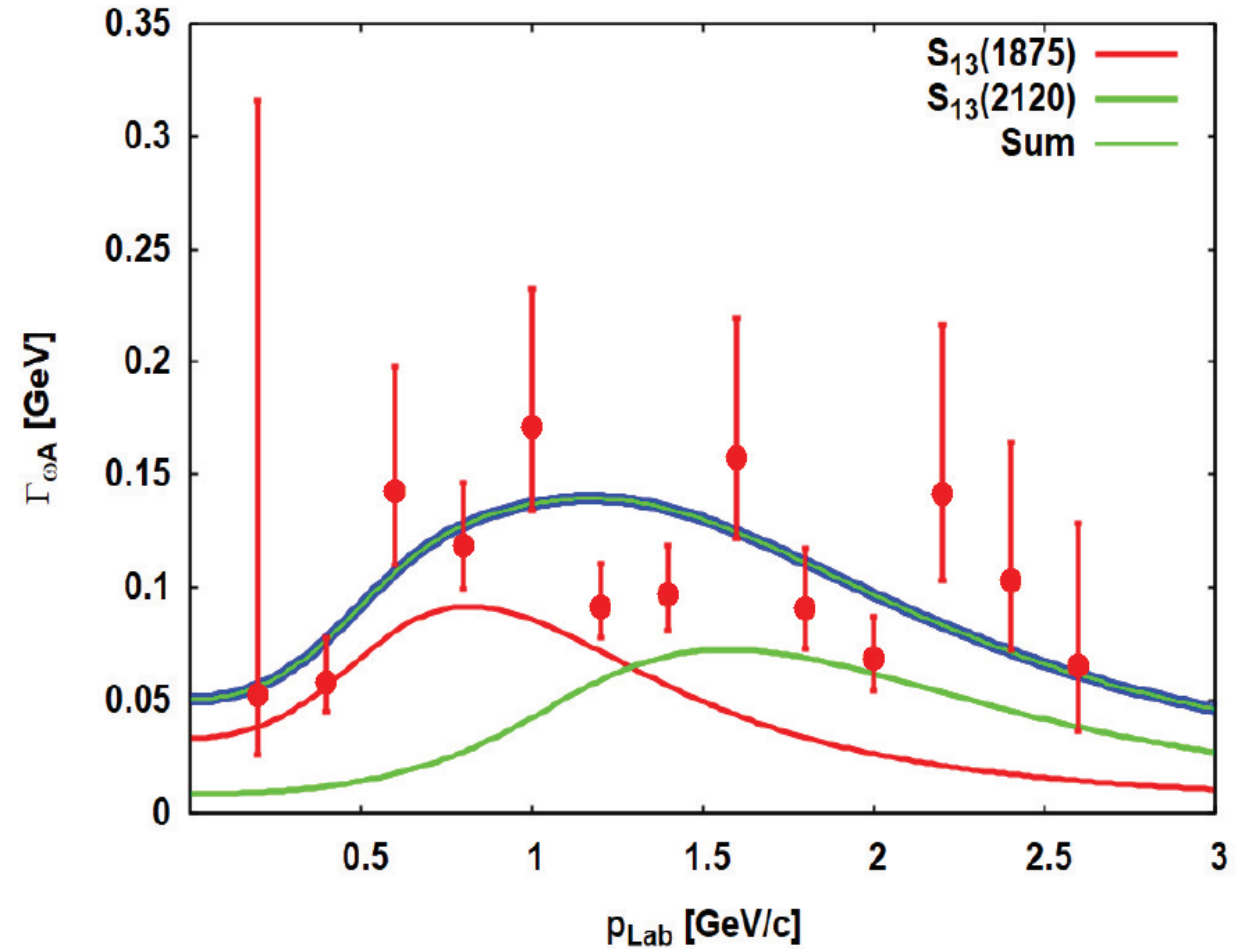
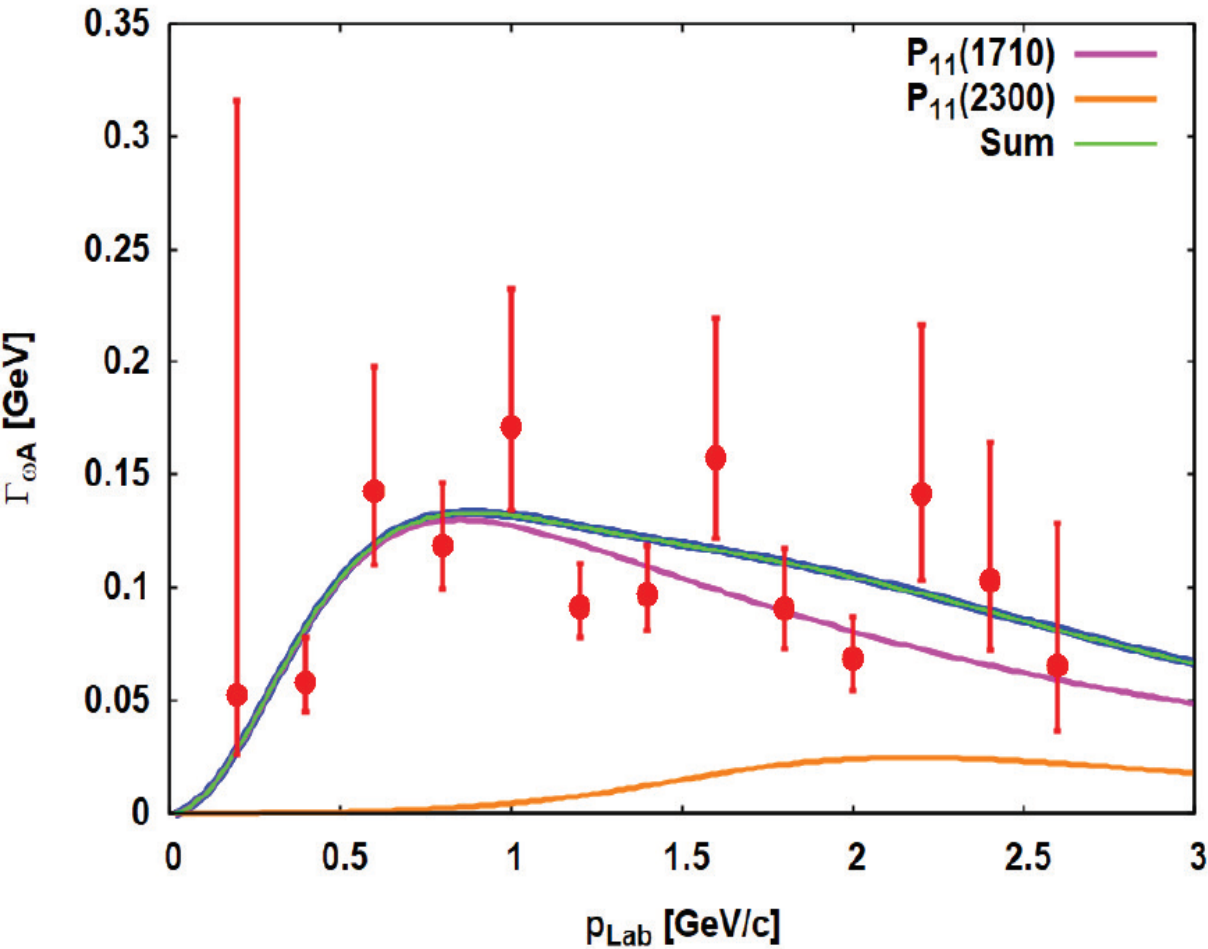
Exploratory **Independent** P-wave and S-Wave Studies

Independent P-wave and S-wave Fits to the observed Widths (at central density of ^{93}Nb : $\rho=0.140 \text{ fm}^{-3}$)



**CBELSA/TAPS Data:
Friedrich et al., Eur. Phys. J. A (2016) 52: 297**

Dominant P-wave and S-wave Components

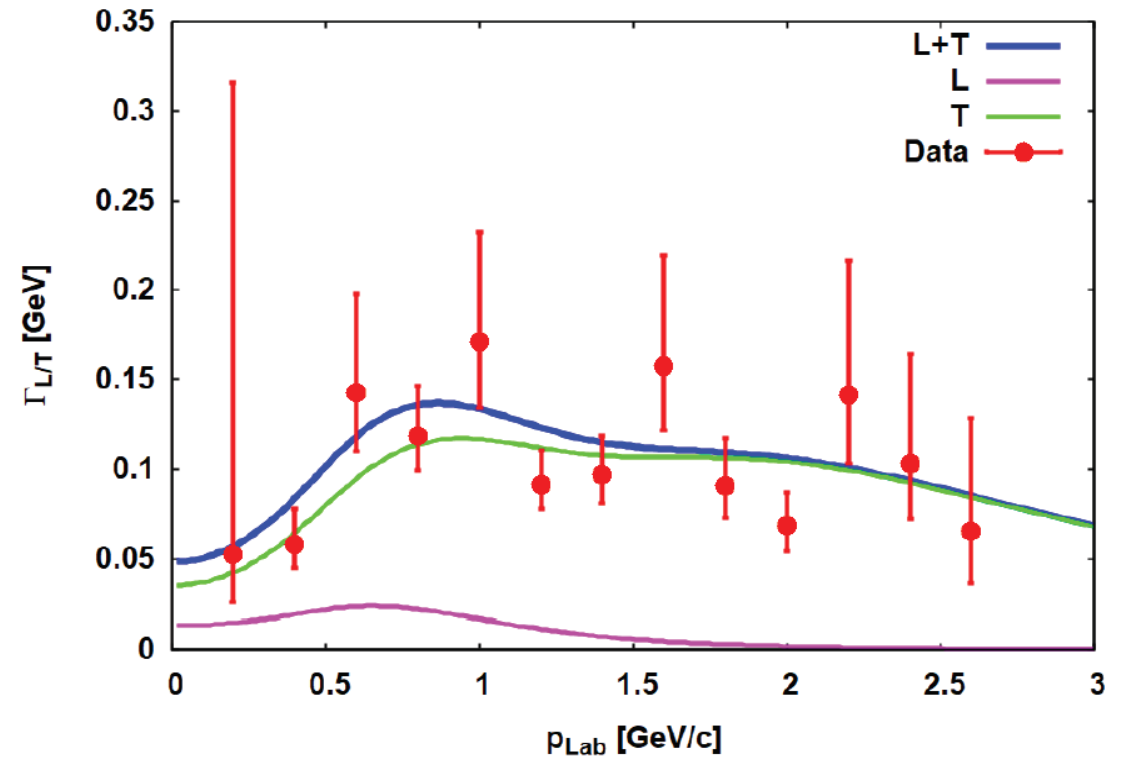
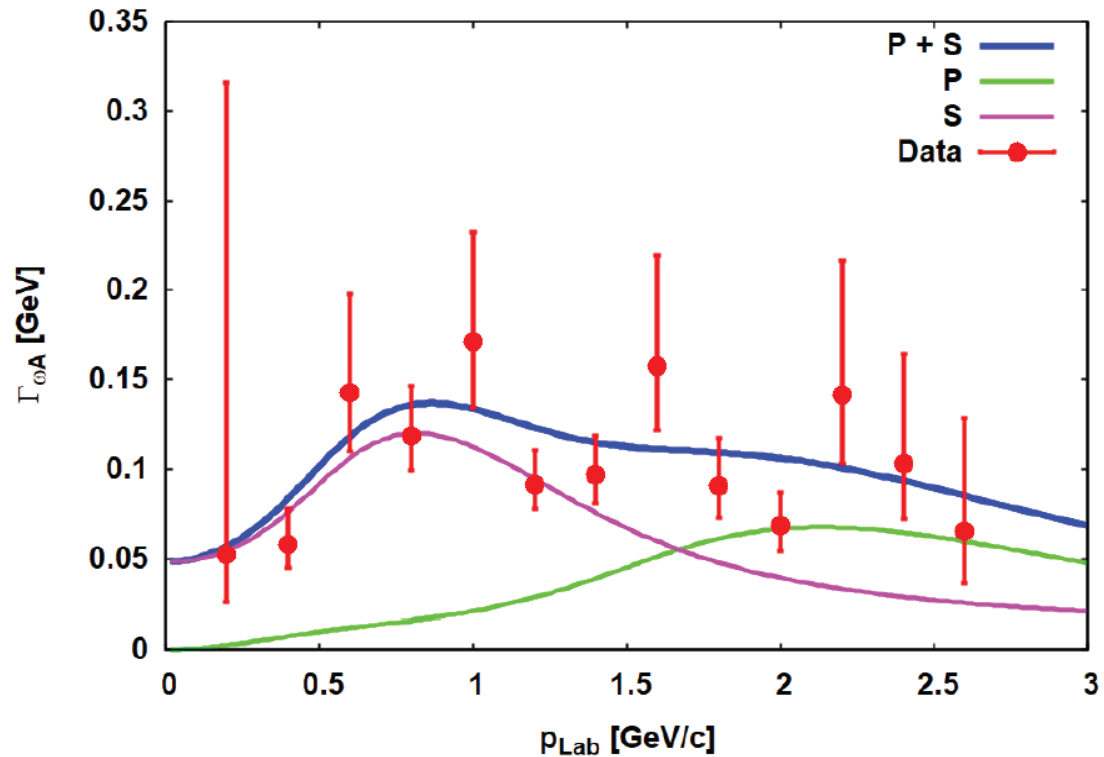


- P-wave and S-wave polarization self-energies are compatible with the data
- The complete description must include both combined
- Width at threshold is decisive for and selective on the N^* states

Combined P-wave and S-wave Description

Results of Combined P- and S-wave Approach

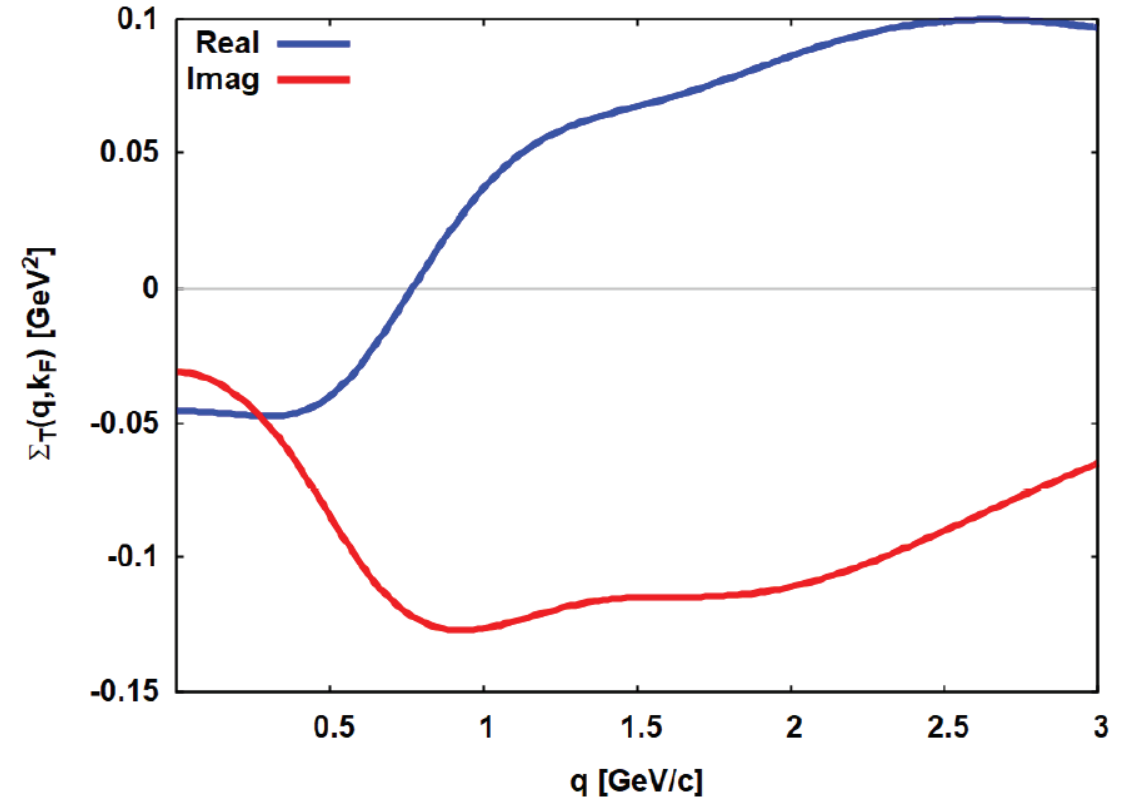
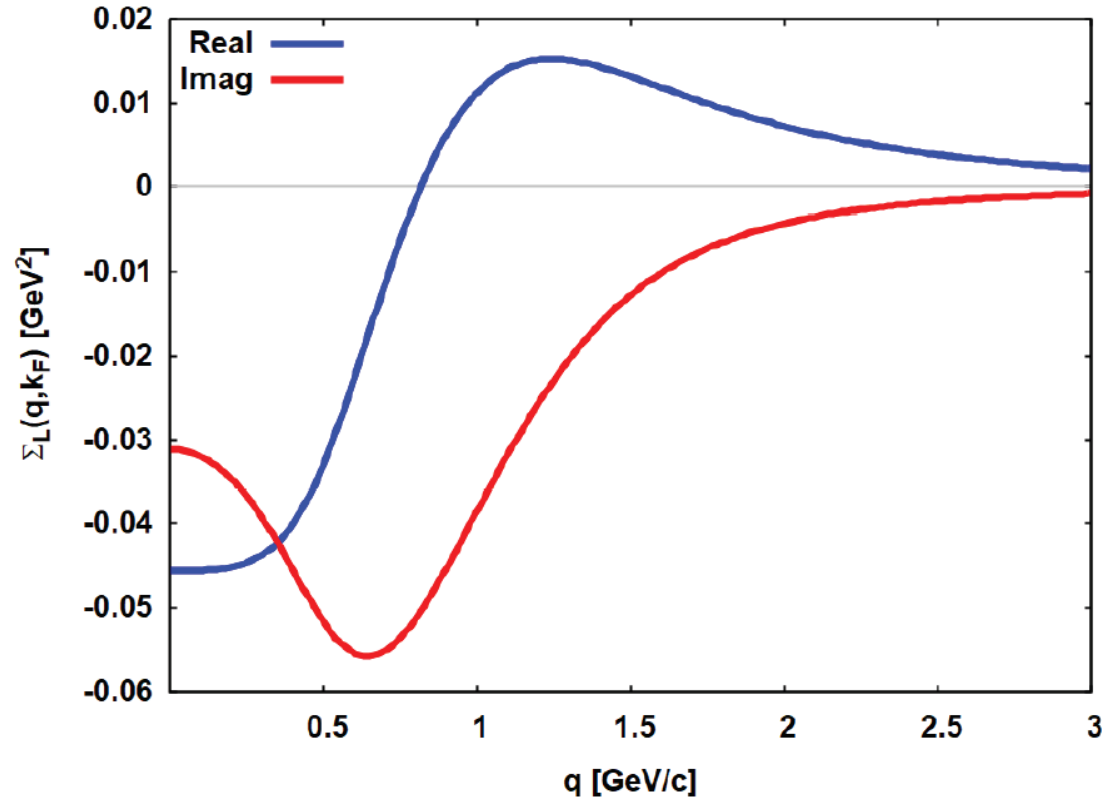
Unconstrained Fit at $\rho=0.140 \text{ fm}^{-3}$ ($\chi^2=1.014$)



$$\Gamma_{\text{thr}} = \Gamma_{\text{free}} + \Gamma_A = 48.41 \text{ MeV}$$
$$\Rightarrow \Gamma_A = 39.73 \text{ MeV}$$

Self-Energies and Spectral Distributions

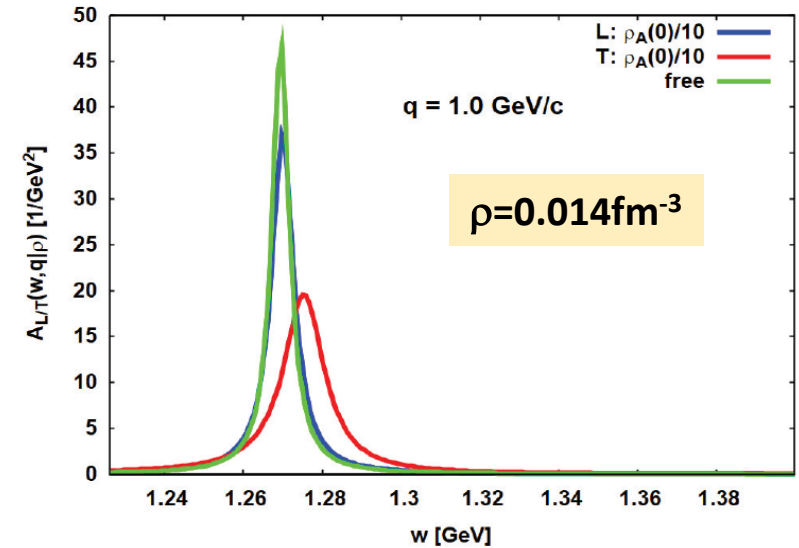
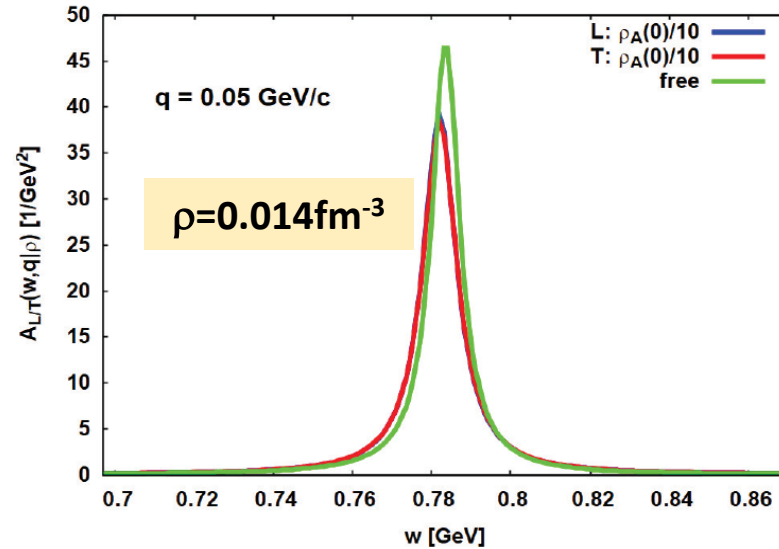
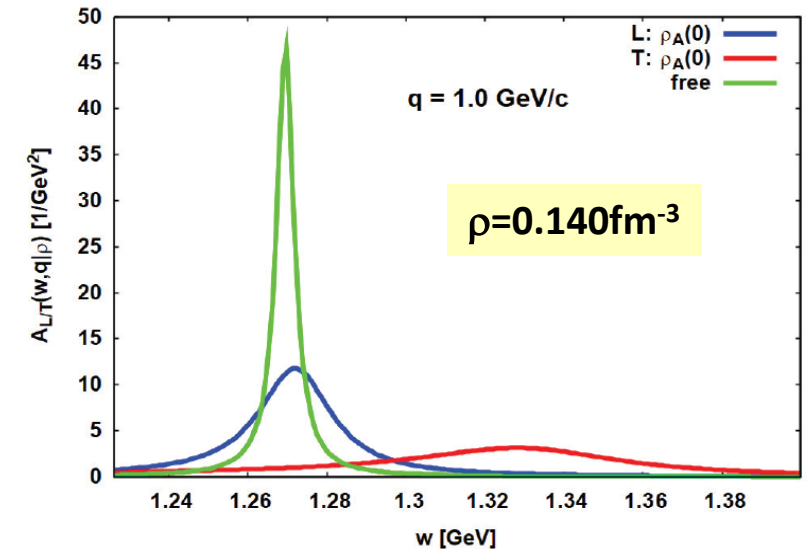
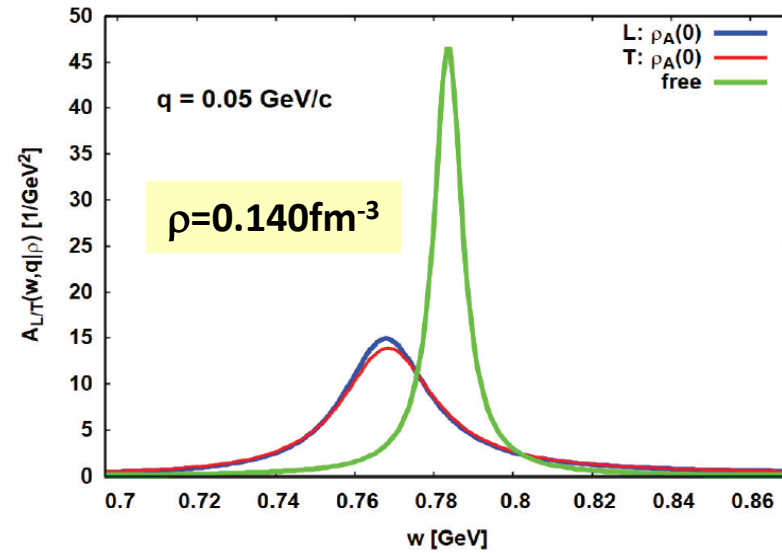
Energy Dependence of ω +A Polarization Self-Energies in the Center of ^{93}Nb ($\rho=0.14 \text{ fm}^{-3}$)



- **Re($\Sigma_{\omega A}$)** change from attractive to repulsive
- **Downward mass shifts** close to threshold
- **Upward mass shifts** at higher energies

ω In-Medium Spectral Distributions

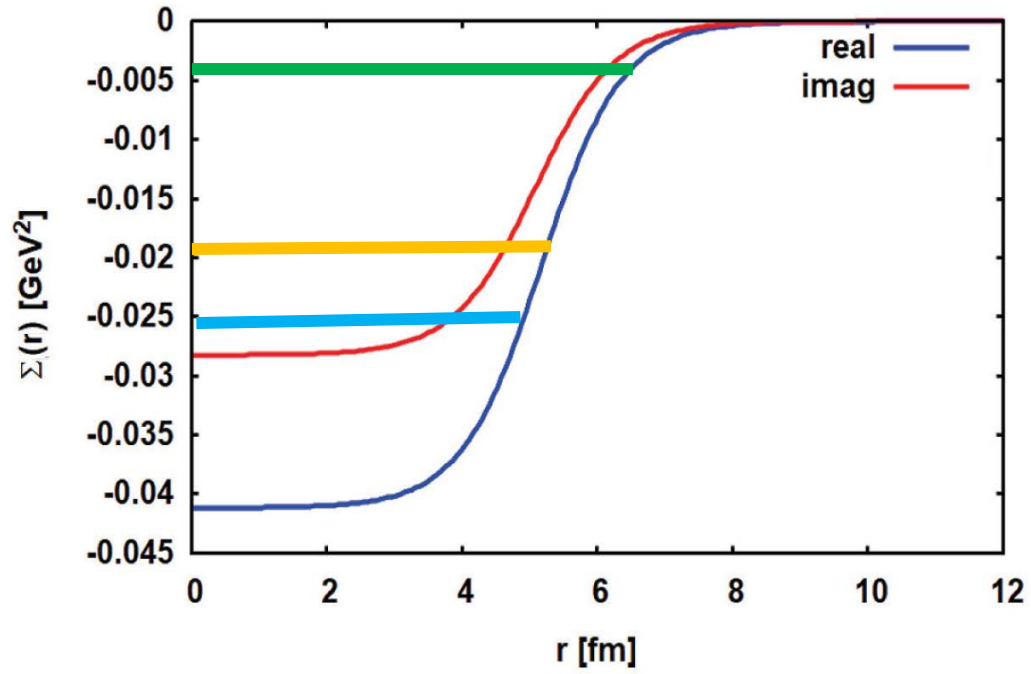
- At $q = 0.05$ GeV/c and central density (upper left) the distributions are moved to energies $\omega < m_w$
- At $q = 1$ GeV/c (upper right), A_L is centered close to A_{free} , A_T is shifted to $\omega > m_w$ with $\Gamma_T \gg \Gamma_L$.
- At low densities (lower row), these differences prevail on a less pronounced level.



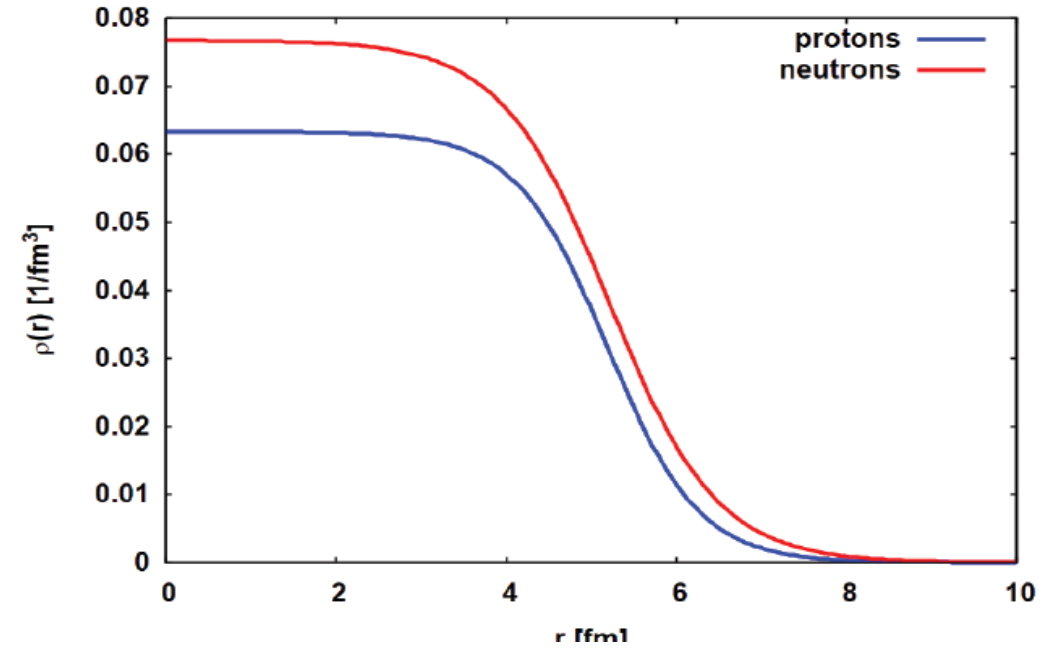
$$A_{L/T}(w, \mathbf{q}|\rho) = -\frac{1}{\pi} \text{Im} \left(\frac{1}{w^2 - \mathbf{q}^2 - m_w^2 - \Sigma_{L/T}(w, \mathbf{q}, k_F) - \Sigma_{free}(w, \mathbf{q})} \right)$$

$\omega+^{93}\text{Nb}$ Bound States?

$\omega+^{93}\text{Nb}$ Self-Energy at Threshold and Low-Energy Parameters



$\omega+^{93}\text{Nb}$ g.s. self-energies at threshold

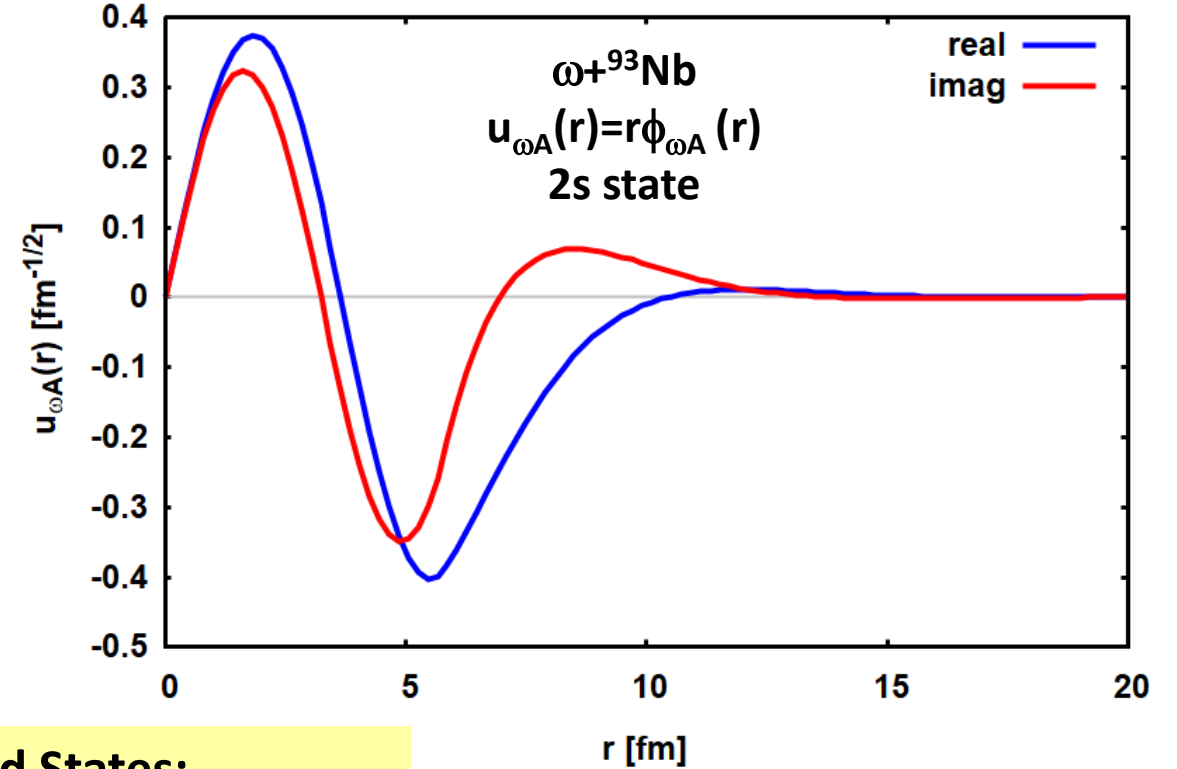
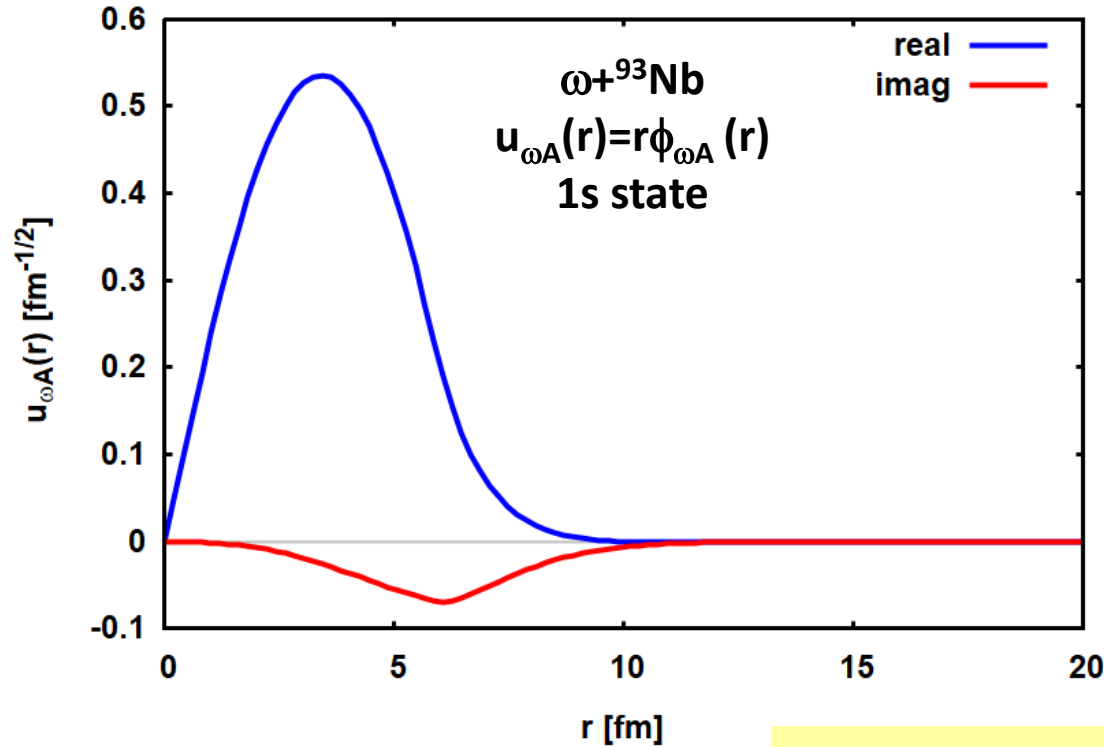


^{93}Nb RMF proton and neutron g.s densities

Low-energy Constants:
 $a_s = 5.6542 - i0.9041 \text{ fm}$ and $r_s = 3.7372 - i0.5183 \text{ fm}$

$\rightarrow \varepsilon_B = -0.488 \text{ MeV} ; \Gamma_B = 13.125 \text{ MeV}$

$\omega+^{93}\text{Nb}$ Bound States by Solution of the Wave Equation



Bound States:

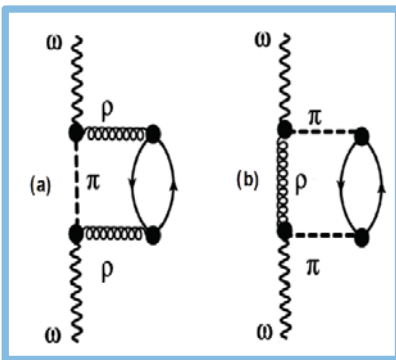
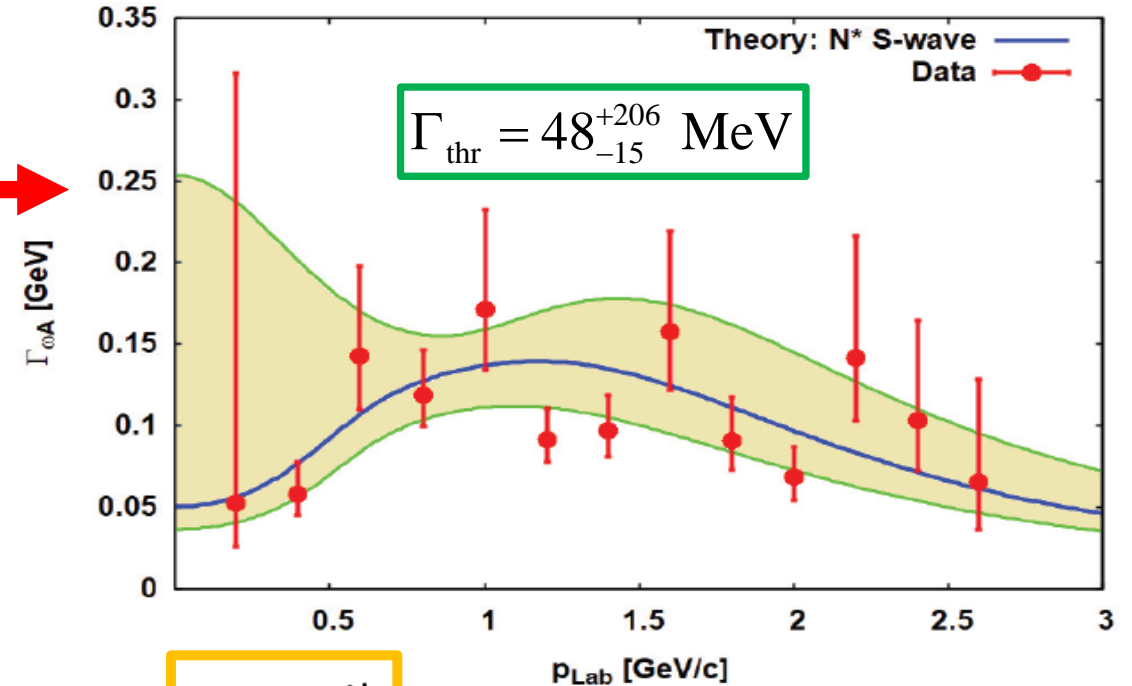
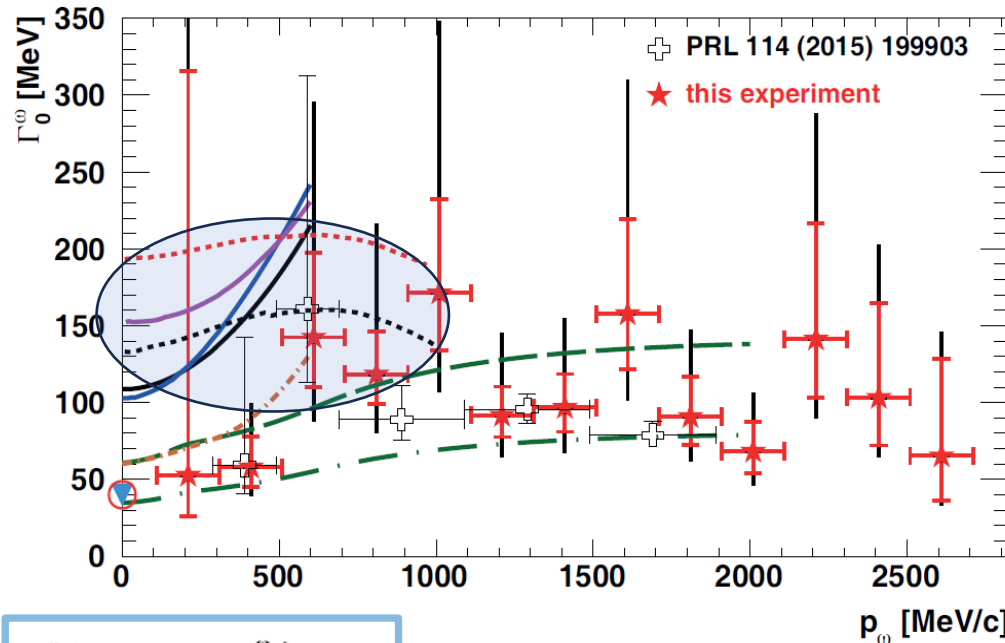
$$\varepsilon_B(2s) = -1.910 - i12.027 \text{ MeV}$$

$$\varepsilon_B(1p) = -11.804 - i13.575 \text{ MeV}$$

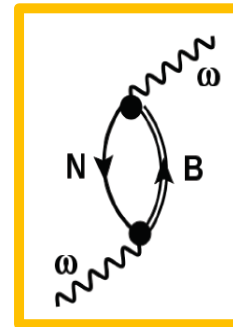
$$\varepsilon_B(1s) = -17.266 - i14.136 \text{ MeV}$$

$$c\tau_{\text{tot}} \sim 5 \dots 6 \text{ fm} \ll c\tau_{\text{free}} \sim 22 \text{ fm}$$

Meson Cloud vs. Polarization Self-Energies



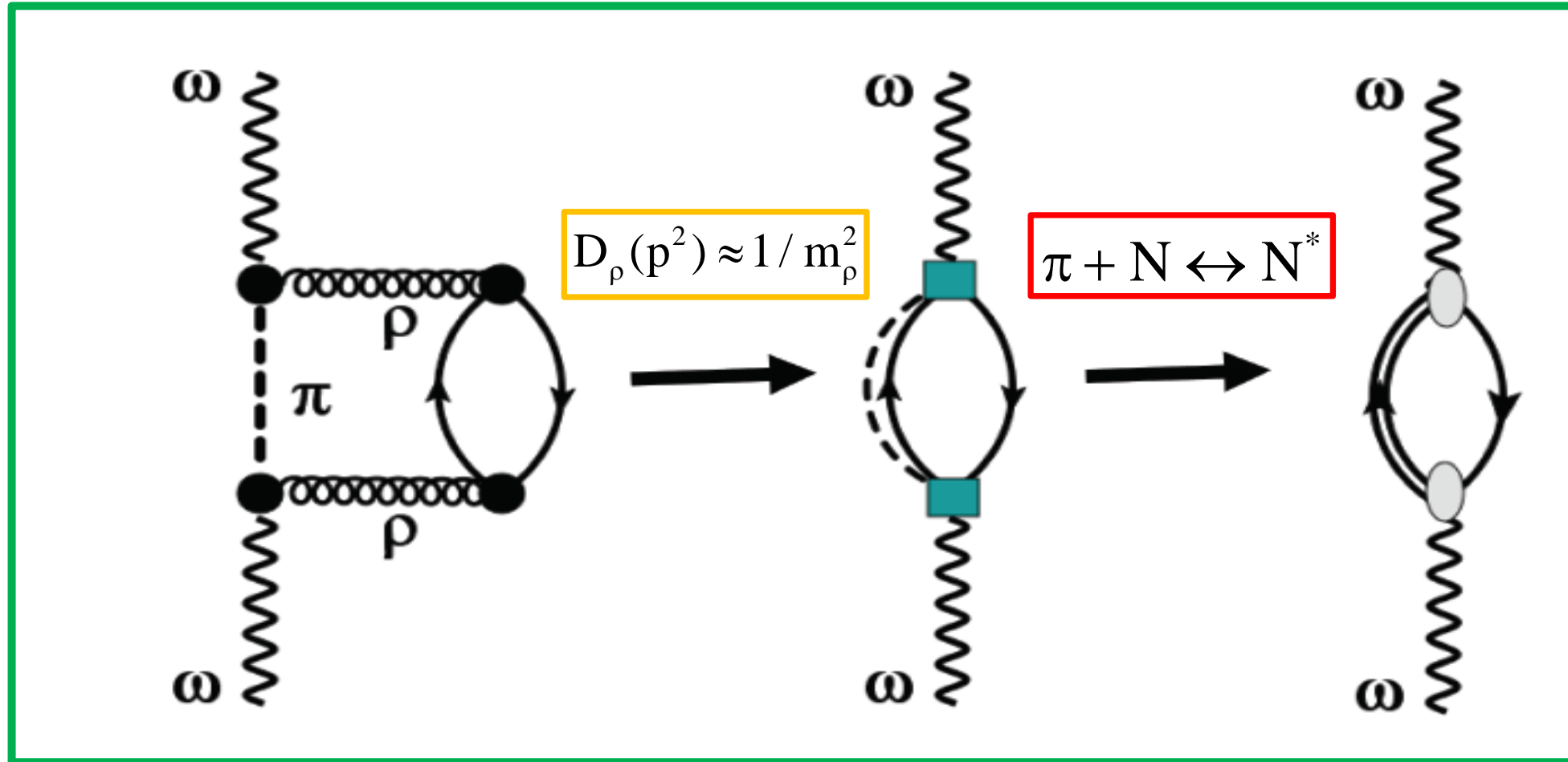
- **Meson Cloud Self-energies**
- **Solid:** Ramos et al., EPJ A 49, 148 (2013)
- **Short dashed:** Cabrera & Rapp, PLB 729, 67 (2014)



- **Polarization Modes**
- $\omega \rightarrow N^* N^{-1} \rightarrow \omega$
- **see also left figure, Long dashed:** Mühlich et al., NPA 780, 187 ('06)

from Friedrich et al., Eur. Phys. J. A (2016) 52: 297

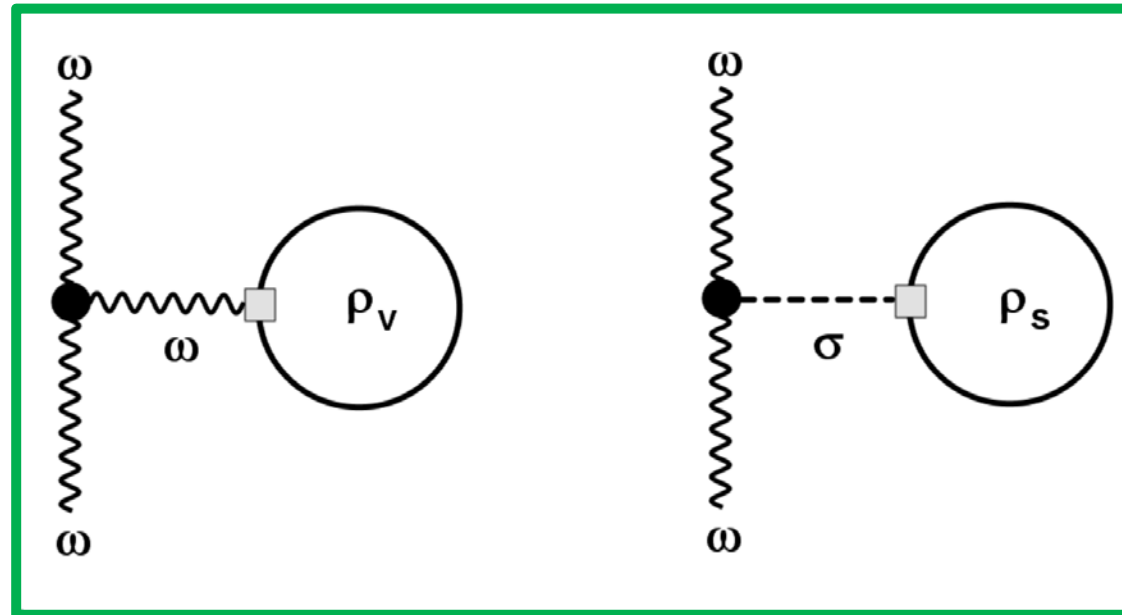
Overlap of N^*N^{-1} and Meson Cloud Self-Energies



Illustrating the relation of in-medium meson loops to baryon particle-hole loops for the $\omega \rightarrow \pi + \rho$ decay for the case where the rho-meson couples to an (isovector) excitation of the nuclear background medium.

...and what is missing?

...e.g. mean-field dynamics!



$\omega\omega\omega$ -Coupling unknown
but from
SU(3) symmetry

$\omega\omega\sigma$ -Coupling unknown
but weak
if $\sigma \sim [\pi\pi]$

Summary

- ω -mesons in nuclear matter and a finite nucleus
- NN^{-1} and N^*N^{-1} polarization self-energies
- Spectral distributions in nuclear matter
- ω self-energies in ^{93}Nb
- ω - ^{93}Nb bound states
- Meson cloud and medium polarization

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Interactions of ω Mesons in Nuclear Matter and with Nuclei
Eur. Phys. J A (in print)

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