



# Spontaneous Fission from Self-Consistent Calculations

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“Nuclei in the laboratory and in stars”

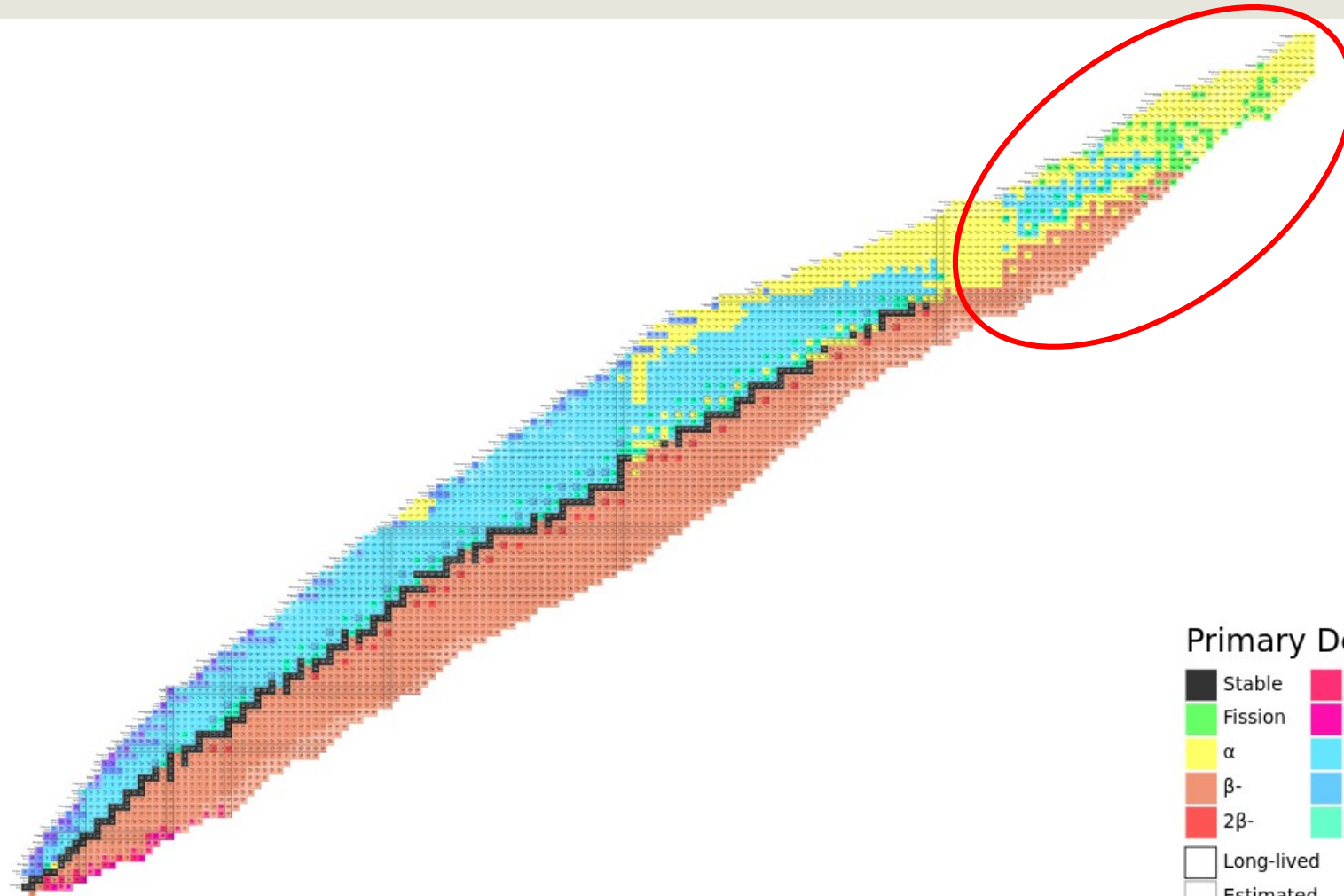


U.S. DEPARTMENT OF  
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Office of  
Science

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# The Landscape

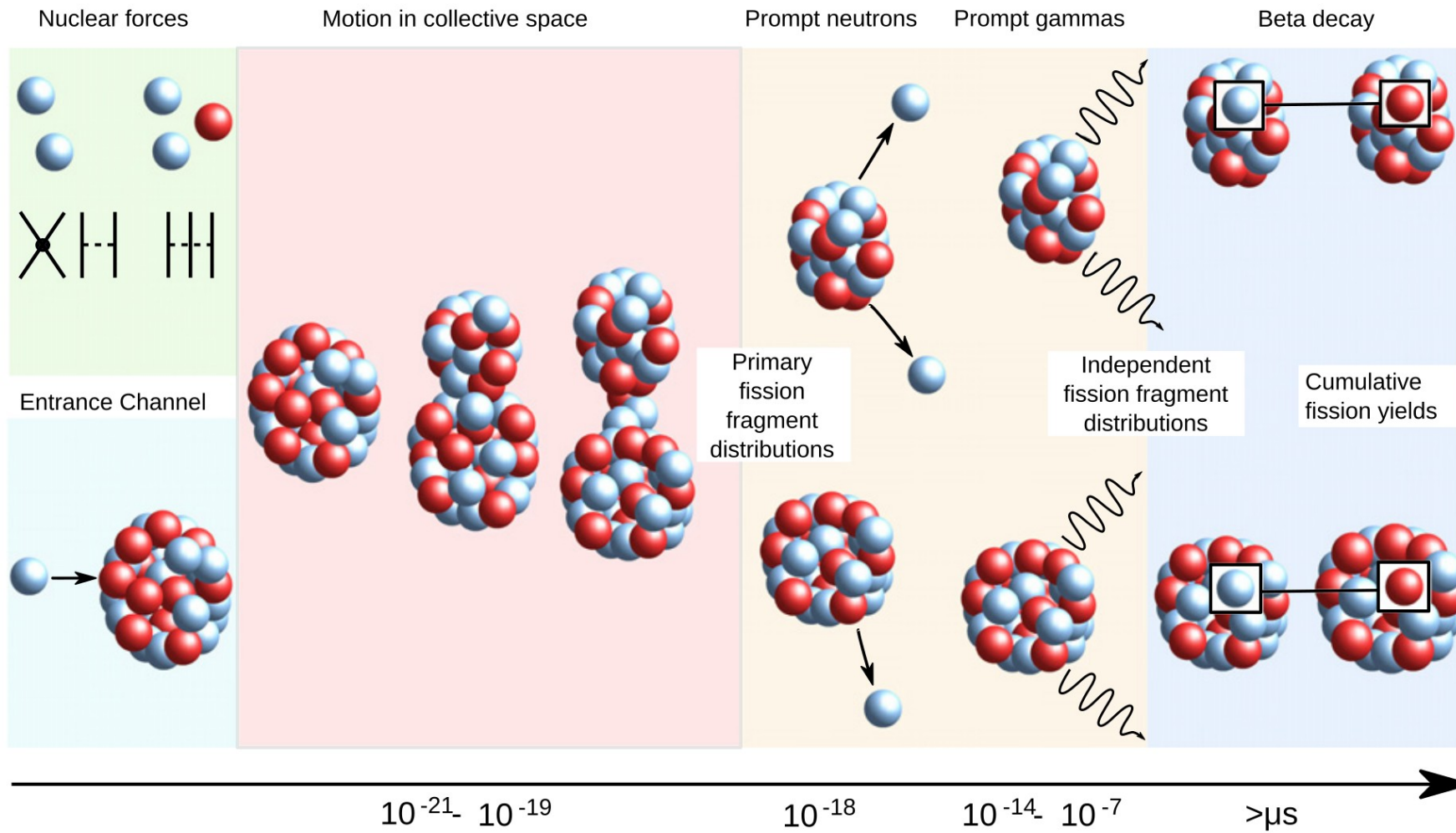


From Colorful  
Nuclear Chart

## Primary Decay Mode

|                 |                 |                        |
|-----------------|-----------------|------------------------|
| Stable          | n               | e <sup>-</sup> capture |
| Fission         | 2n              | p                      |
| α               | β <sup>+</sup>  | 2p                     |
| β <sup>-</sup>  | 2β <sup>+</sup> | 3p                     |
| 2β <sup>-</sup> | e <sup>+</sup>  |                        |
| Long-lived      |                 |                        |
| Estimated       |                 |                        |
| Unknown         |                 |                        |

# Time Scales



Bender et. al. , J. Phys. G: Nucl. Part. Phys. 47 (2020) 113002

# Stationary Density Functional Theory Approach to Fission

- We use static, constrained HFB density functional theory (DFT).
- The nuclear shape is parameterized by multipole moments in the intrinsic frame.

$$\hat{Q}_{\mu\nu} = \hat{r}^\mu Y_{\mu\nu}(\theta, \phi)$$

- We calculate the energy of the nucleus by constraining the energy at given multipole moments.

$$E[\rho, \kappa, \kappa^*, Q_{\mu\nu}] = \langle \psi | H - \sum_{\mu\nu} \lambda_{\mu\nu} \hat{Q}_{\mu\nu} | \psi \rangle$$

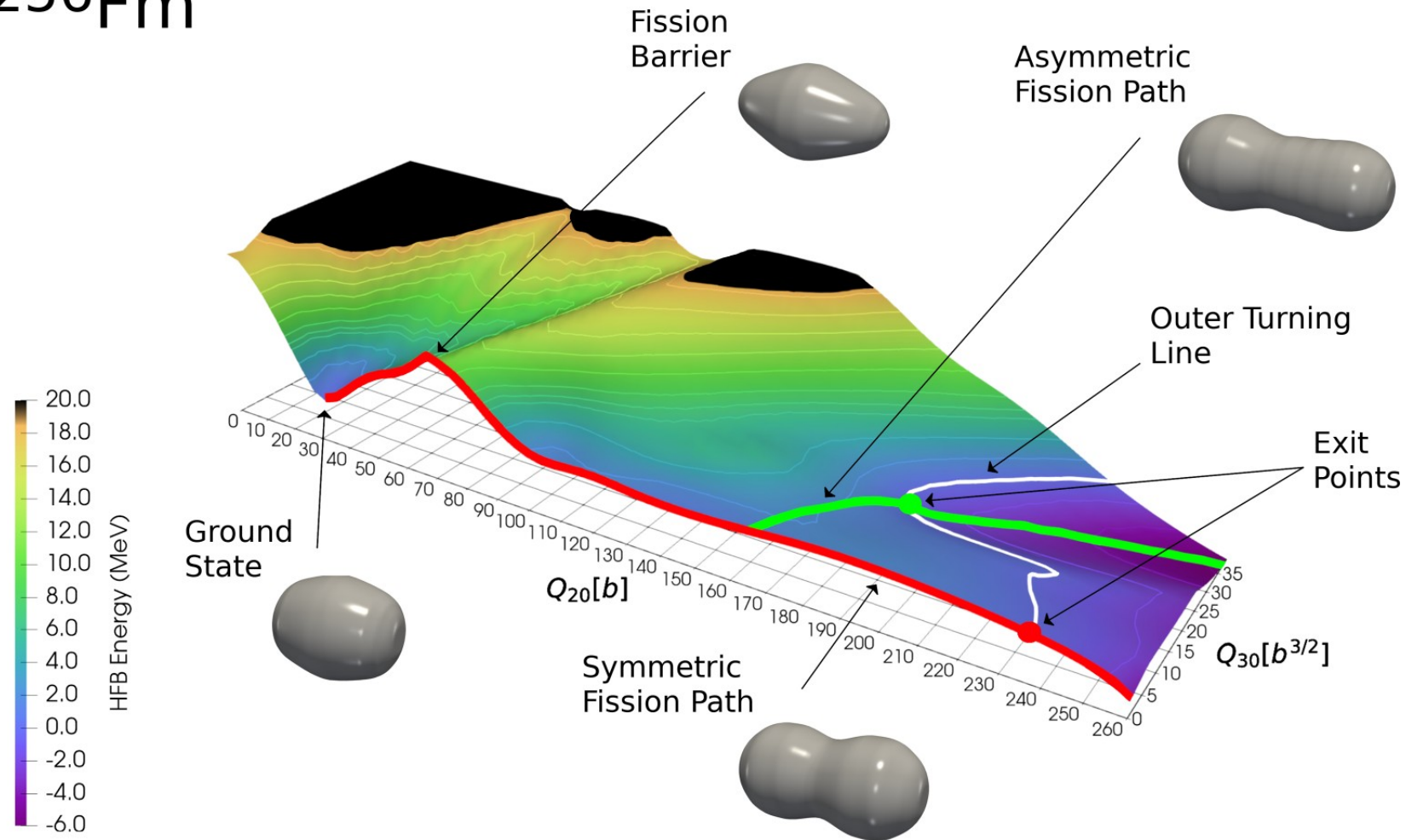
- Potential energy surfaces (PES) in this work were computed using HFBTHOv3.00 [1] and HFBaxial [2]

[1] Perez, R. Navarro et. al *Computer Physics Communications* 220 (2017): 363-375.

[2] Robledo, Luis M., et. al. *Physical Review C* 84, no. 1 (2011): 014312.

# Summary of DFT approach to Fission

$^{256}\text{Fm}$



Flynn et. al.  
Phys. Rev. C  
105, 054302

# Fission as a Finite Dimensional Tunneling Problem

- **Main Problem:** Find the most probable fission pathways in the collective space
- We can view spontaneous fission as a tunneling problem in a collective space.
- Given a PES, we can solve the collective Schrodinger equation with multidimensional WKB theory:

$$\psi(\vec{q}) \sim e^{-S_0[L]} \quad \text{as } \hbar \longrightarrow 0$$

$$S_0[L] \geq \min \frac{1}{\hbar} \int_{s_{\text{in}}}^{s_{\text{out}}} \sqrt{2\mathcal{M}_{\text{eff}}(s) \left( V_{\text{eff}}(s) - \Delta E \right)} ds$$

$\vec{q}(s_{\text{in}}) = \vec{q}_{\text{in}}$   
 $\vec{q}(s_{\text{out}}) = \vec{q}_{\text{out}}$

Effective Inertia                      Effective potential                      Zero-point correction to GS

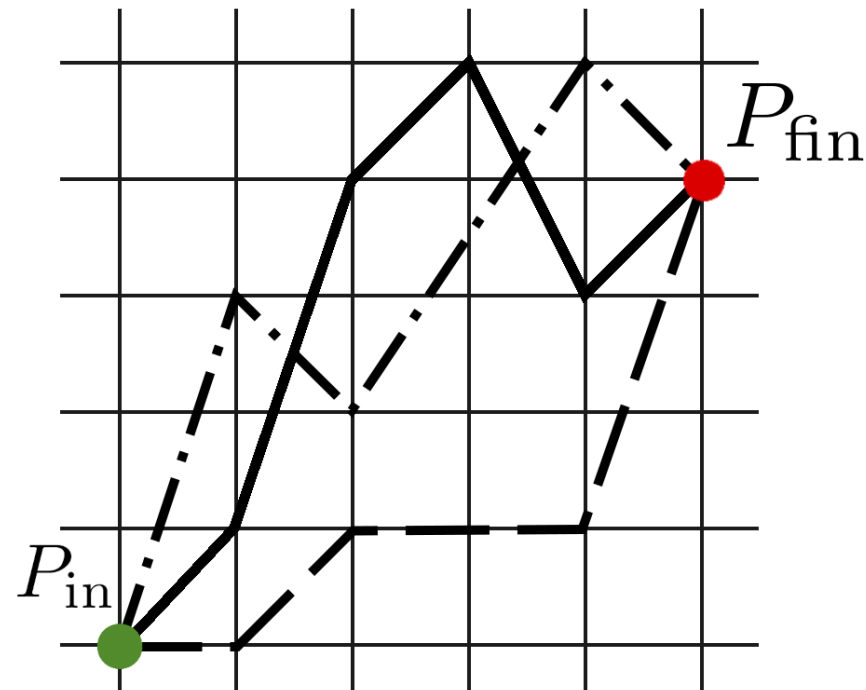
A. Schmid, Ann. Phys. (NY) 170, 333–369 (1986).

# Previous Approaches

- **Main Problem:** Find the most probable fission pathways in the collective space
- Previous approaches have used grid-based methods to find least action paths (LAPs)
- Baran et. al. (1981) [1] used Dynamic Programming Method (DPM) to search for the LAP for fixed initial and final positions.

## Drawbacks of grid-based methods

- Path is restricted to a grid.
- Can only find global minimum.
- Doesn't scale well with respect to dimension and grid size.

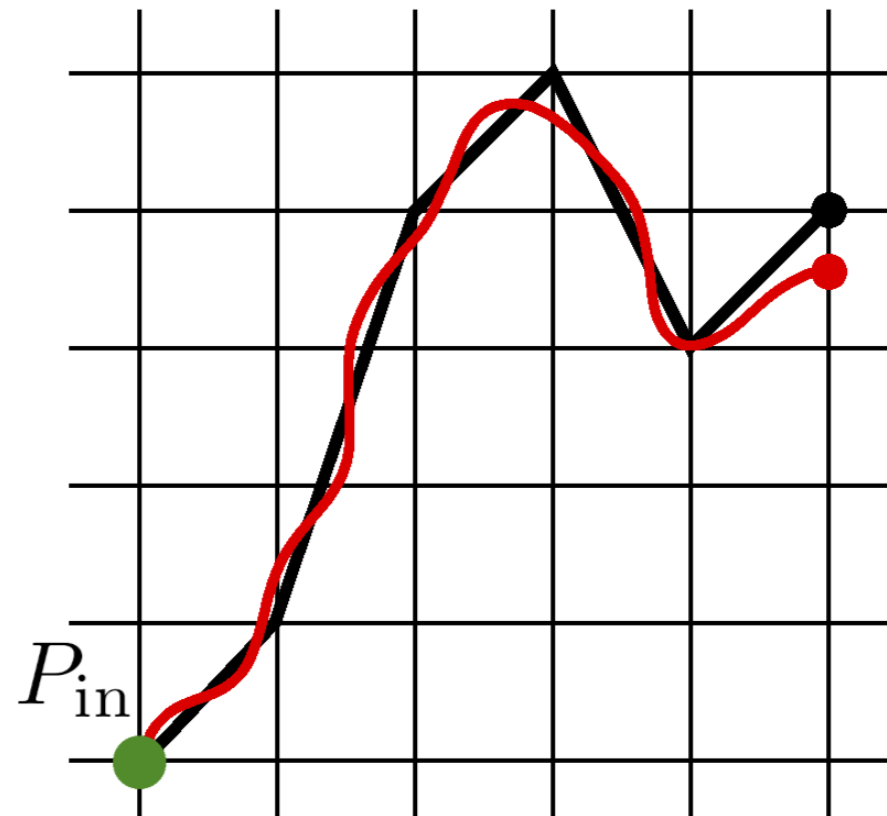


Flynn et. al.  
Phys. Rev. C  
105, 054302

[1] Baran, A. et. al., Nucl. Phys. A 361, 83 (1981).

# The Nudged Elastic Band (NEB) Method

- The nudged elastic band method (NEB) was introduced in molecular chemistry by Henkelman et. al [1,2]
  - Originally used to find minimum energy paths (MEPs)
  - Later extended to LAPs.
- Requires local evaluations of an interpolation function
- Scales well with grid size and dimension compared to grid-based methods



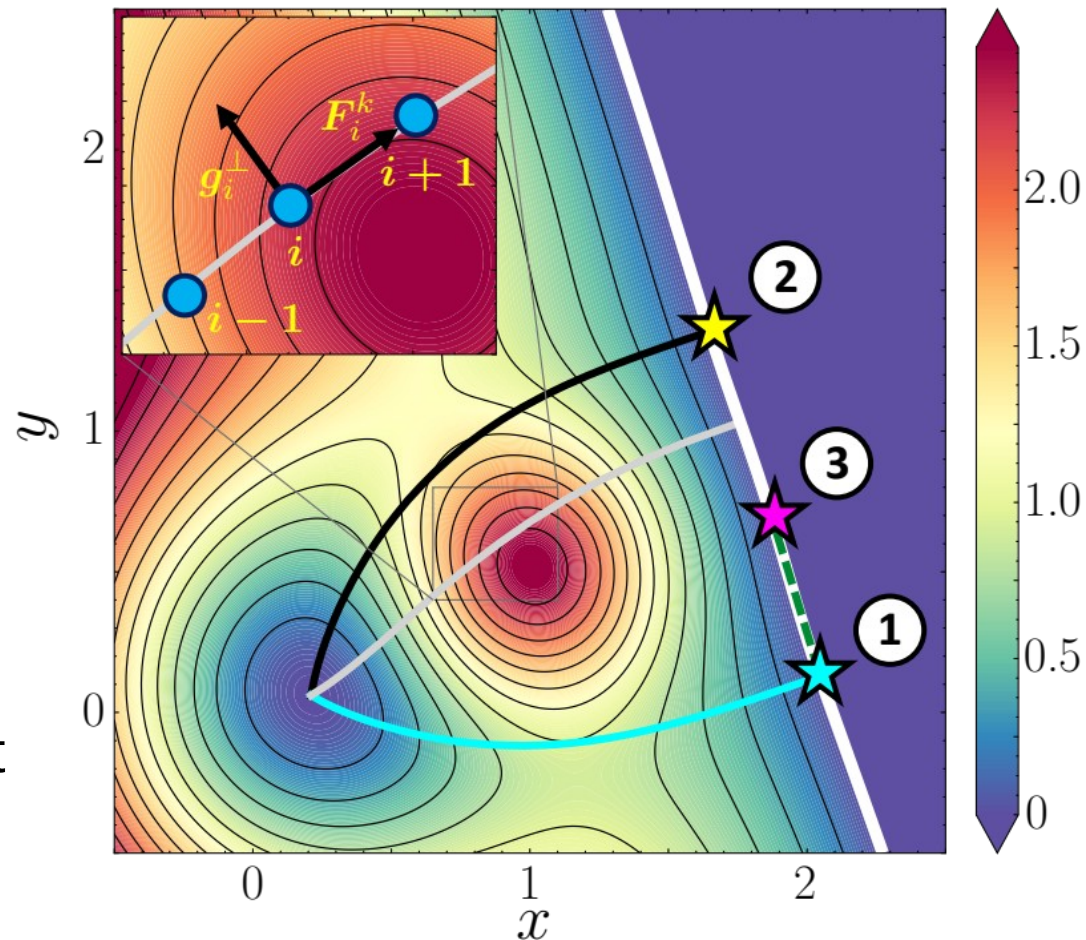
[1] Henkelman, Graeme, and Hannes Jónsson. The Journal of chemical physics 113, no. 22 (2000): 9978-9985.

[2] Henkelman, Graeme, Blas P. Uberuaga, and Hannes Jónsson. The Journal of chemical physics 113, no. 22 (2000): 9901-9904.



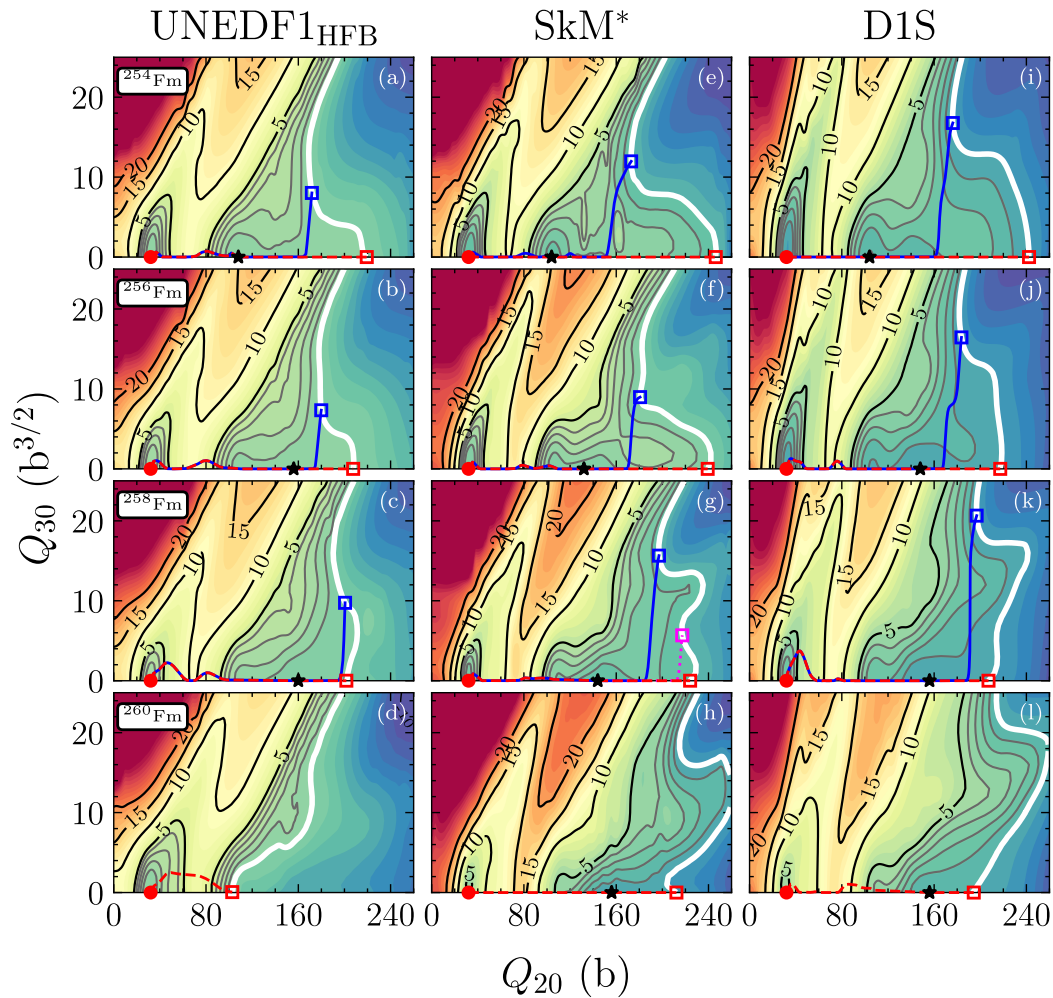
# The Nudged Elastic Band (NEB) Method

- Uses a string of “particles” with mass in a potential (the PES), connected by springs.
- Given an initial guess pathway, the particles adjust their positions according to the gradient a net force.
  - Net force is the sum of the spring force  $\vec{F}_i^k$  and the gradient of the action  $\vec{\nabla} S_i$ .
- The particles are updated locally
  - Can find local minima easily so it can be applied to multi-modal fission.
- Harmonic force is applied to constrain the end point to an energy contour.
- **Demo:** [www.pyneb.dev/](http://www.pyneb.dev/)



Flynn, Eric, Daniel Lay, at. el. Physical Review C 105, no. 5 (2022): 054302.

# Results: Fermium Isotope Chain Pathways

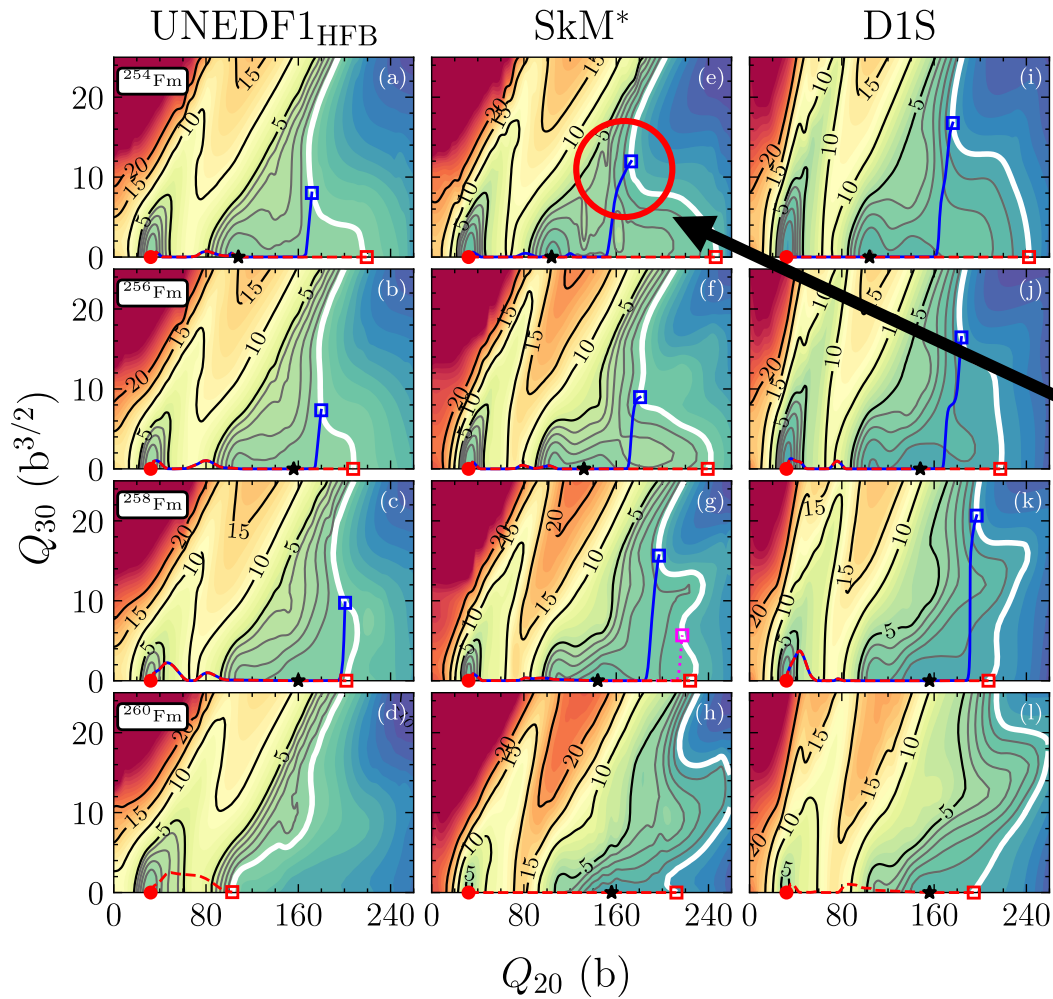


- At the exit points (square points on left), we identify a pre-scission configuration using a nucleon localization function.
- We use the hybrid approach developed by Sadhukhan et al. Phys. Rev. C 101, 065803 2020.

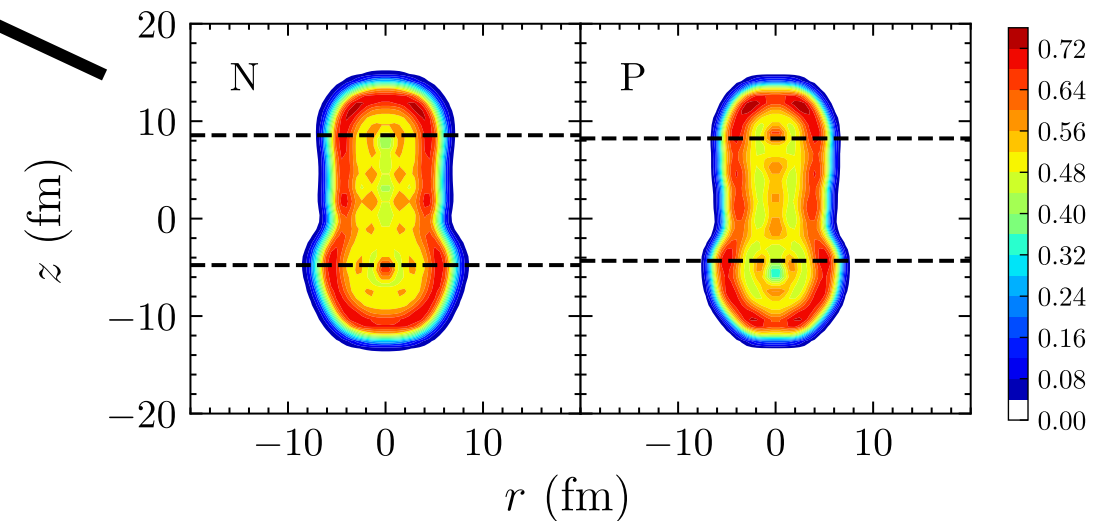
Lay, Daniel, Eric Flynn, et. al. Physical Review C 109, no. 4 (2024): 044306.



# Results: Fermium Isotope Chain Pathways



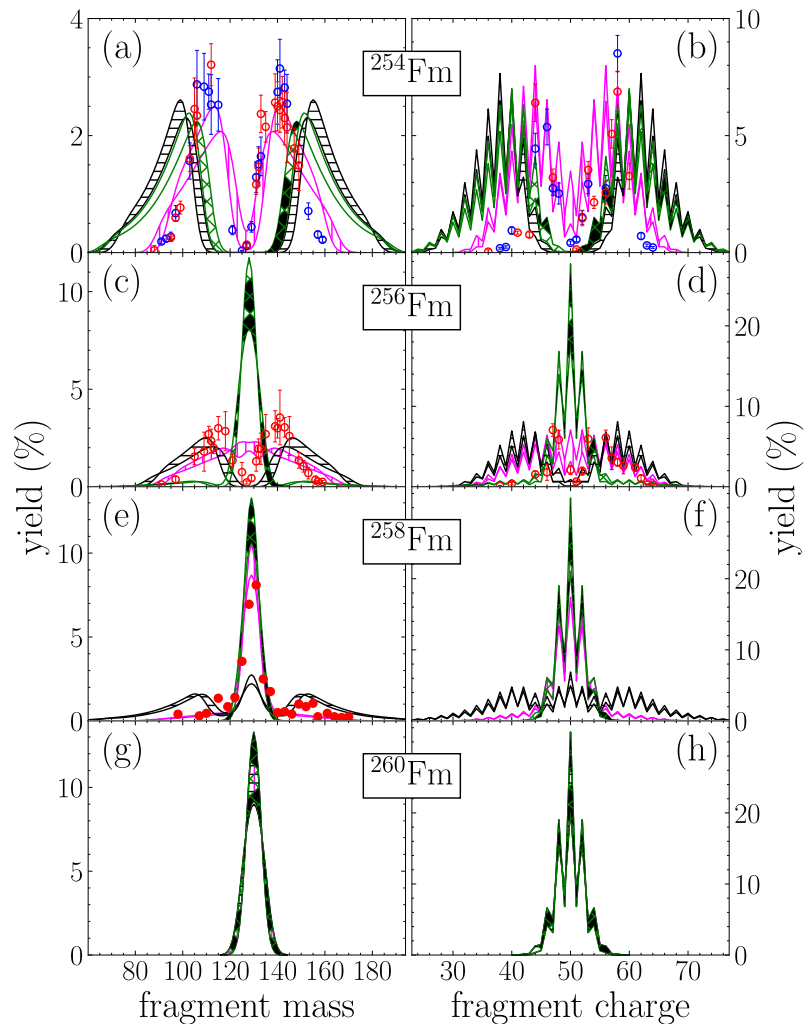
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Lay, Daniel, Eric Flynn, et. al. Physical Review C 109, no. 4 (2024): 044306.



# Results: Fermium Isotope Chain Yields

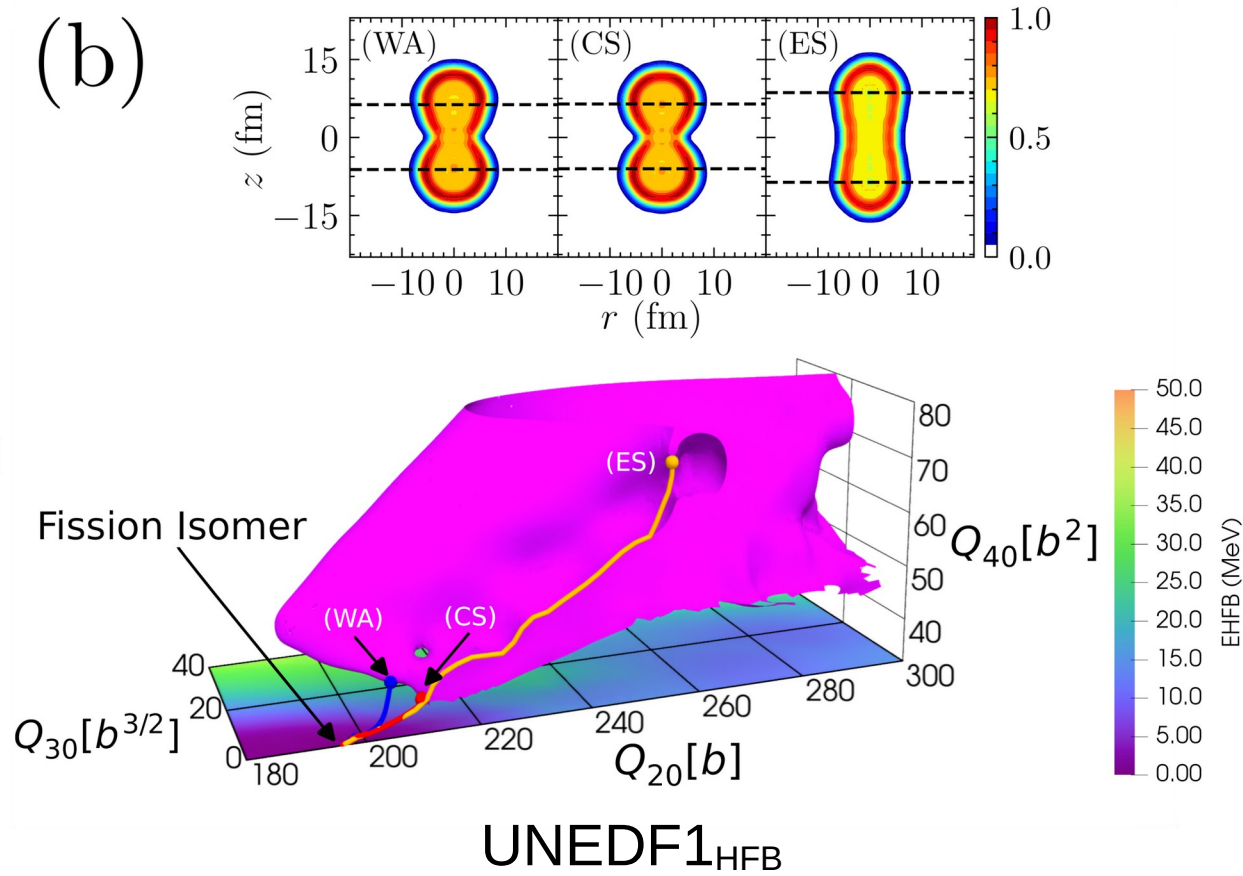
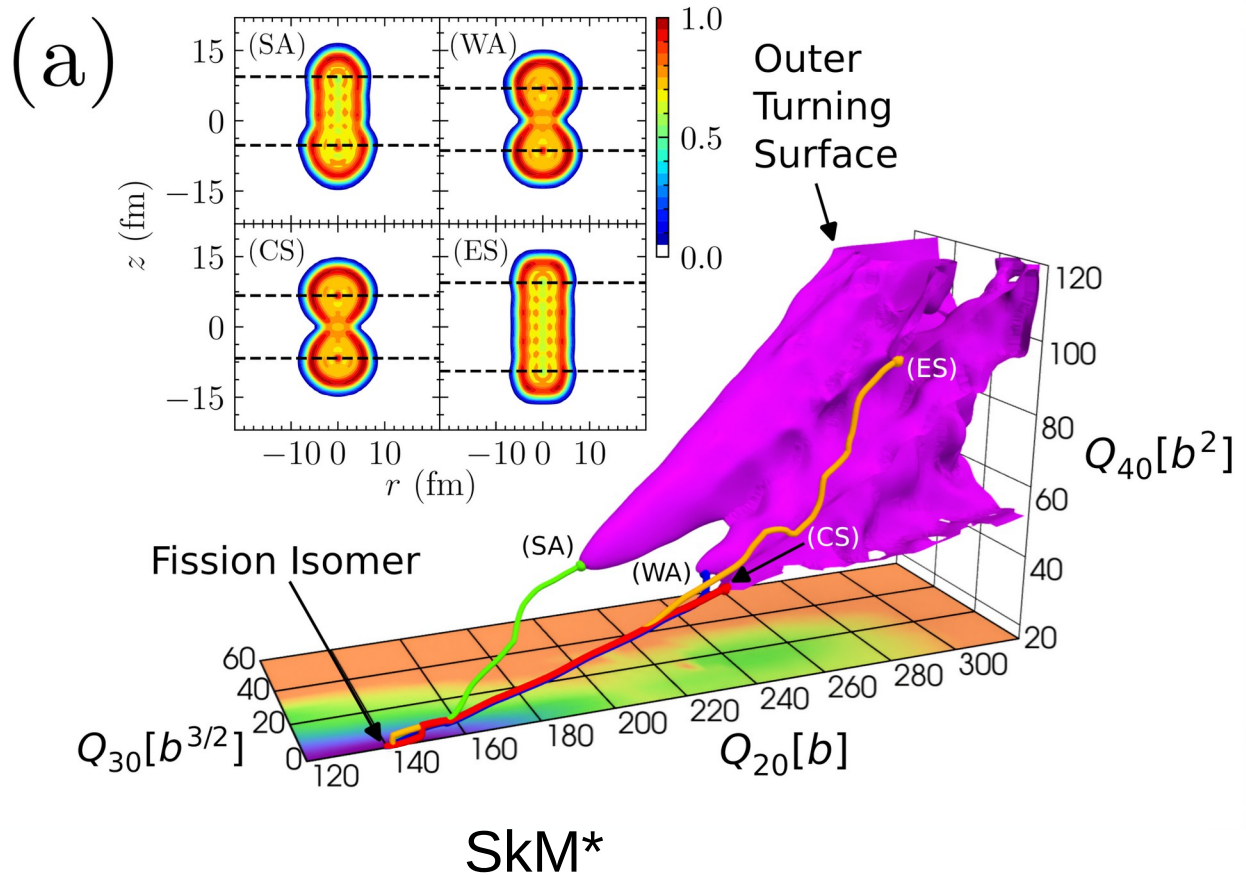


Black = SkM\*, Magenta = UNEDF1<sub>HFB</sub>, Green = D1S

- Weighting each mode by its respective action, we can estimate the total yield.
- The asymmetric mode becomes less significant.
- Observed shift from asymmetric dominant fission to symmetric dominant in the Fermium chain.
  - This is due to the shell closure at  $^{132}\text{Sn}$
- PyNEB can identify multiple modes of fission and classify them based on the action.

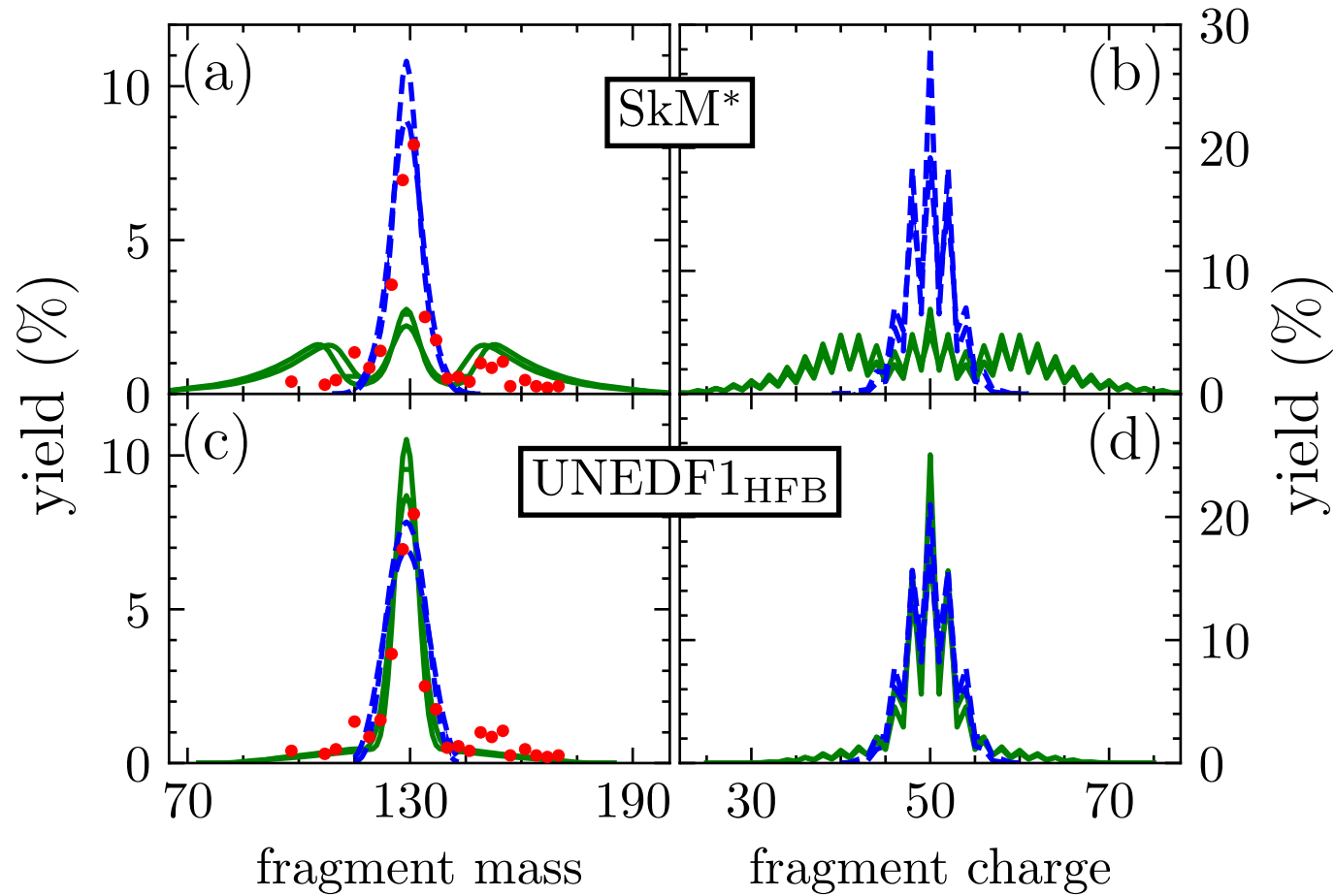
Lay, Daniel, Eric Flynn, et. al. Physical Review C 109, no. 4 (2024): 044306.

# Results: $^{258}\text{Fm}$ in 3-D



Lay, Daniel, Eric Flynn, et. al. Physical Review C 109, no. 4 (2024): 044306.

# Results: $^{258}\text{Fm}$ 2D vs 3D



Lay, Daniel, Eric Flynn, et. al. Physical Review C 109, no. 4 (2024): 044306.



# Future Directions: Instantons

- PES and inertia calculations become difficult problem as dimension increases.
- We can instead look for mean-field instantons using framework developed by Levit, Negele, et. al Phys. Rev. C 22, 1979 and H. Reinhardt Nucl. Phys. A, Volume 367, Issue 2, 1981
- By conservation of effort, we exchange the curse of dimensionality for a more tricky numerical problem.

$$\mathrm{Tr}_N \left( \frac{1}{\hat{H} - E} \right) = i \int_{\gamma} e^{-E\beta} \mathrm{Tr}_N \left( e^{-\beta \hat{H}} \right) d\beta$$

$$Z = \mathrm{Tr} \left( e^{-\beta \hat{H}} \right) = \int \prod_{rs} D\sigma_{rs} e^{-S_{eff}[\sigma]}$$

# Future Directions: Instantons

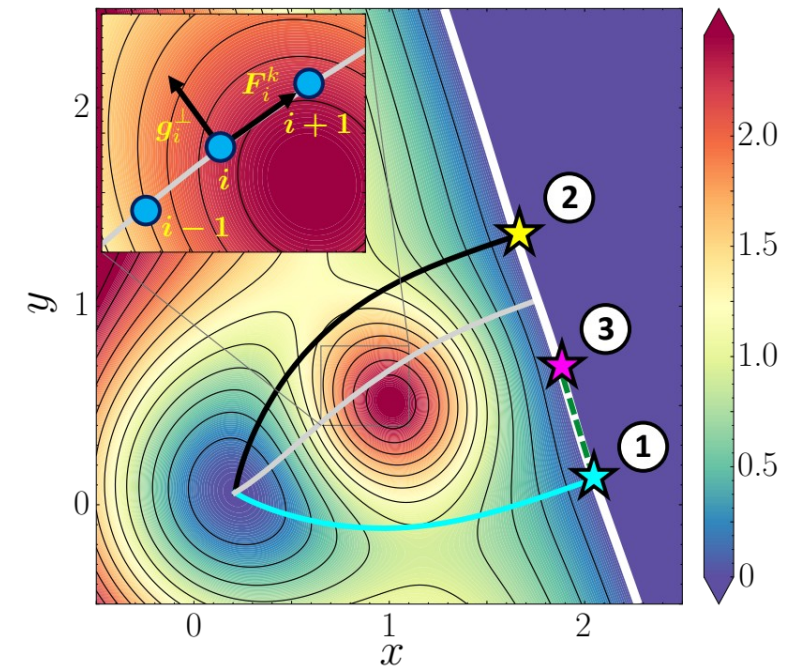
$$\delta S[\sigma] = 0$$

- Analytic expressions for the saddle points are two coupled partial differential equations: forward propagating and backward propagating imaginary time equations.
- Solutions must be periodic in imaginary time.
- Numerically unstable and resistant to standard self-consistent mean-field theory methods.
- Saddle points are generically complex.
- **Idea:** Use the holomorphic gradient flow technique used by Scott Lawrence and Yukari Yamauchi Phys. Rev. D 103, 114509 2021 to find saddle points of the functional integral representation of the partition function.



# Conclusions

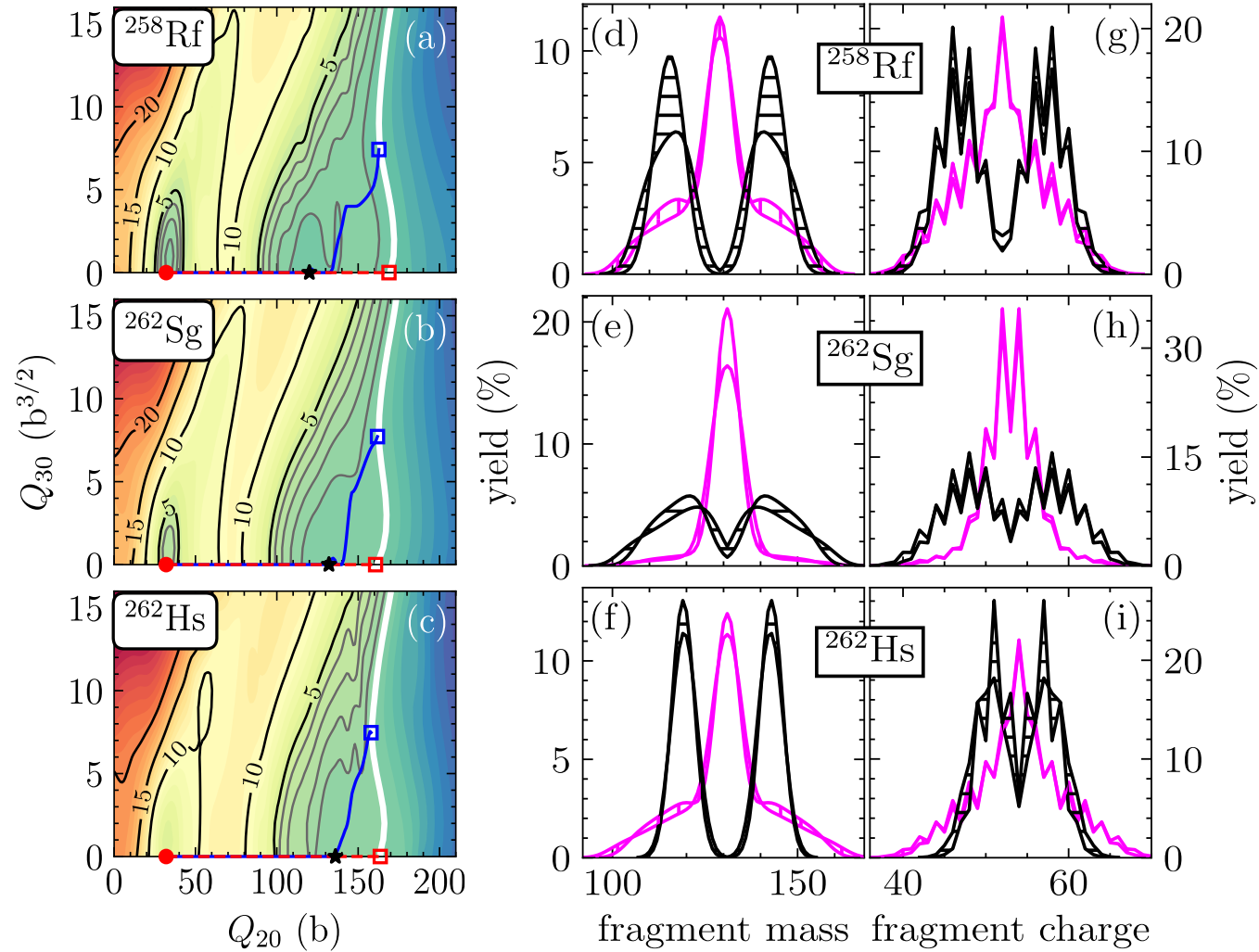
- NEB can efficiently classify different modes of fission in 2 and 3 dimensions.
- NEB overcomes limitations of grid-based algorithms.
- Combined with the micro-canonical hybrid model, NEB can efficiently predict fission yields and potentially predict realistic fission lifetimes.
- Mean-field instantons?
- Thanks to all collaborators and group members: Sylvester Agbemava, Samuel Giuliani, Kyle Godbey, Daniel Lay, Witek Nazarewicz, Jhilam Sadhukhan
- Thanks to Yukari Yamauchi and Scott Lawrence for helpful discussions



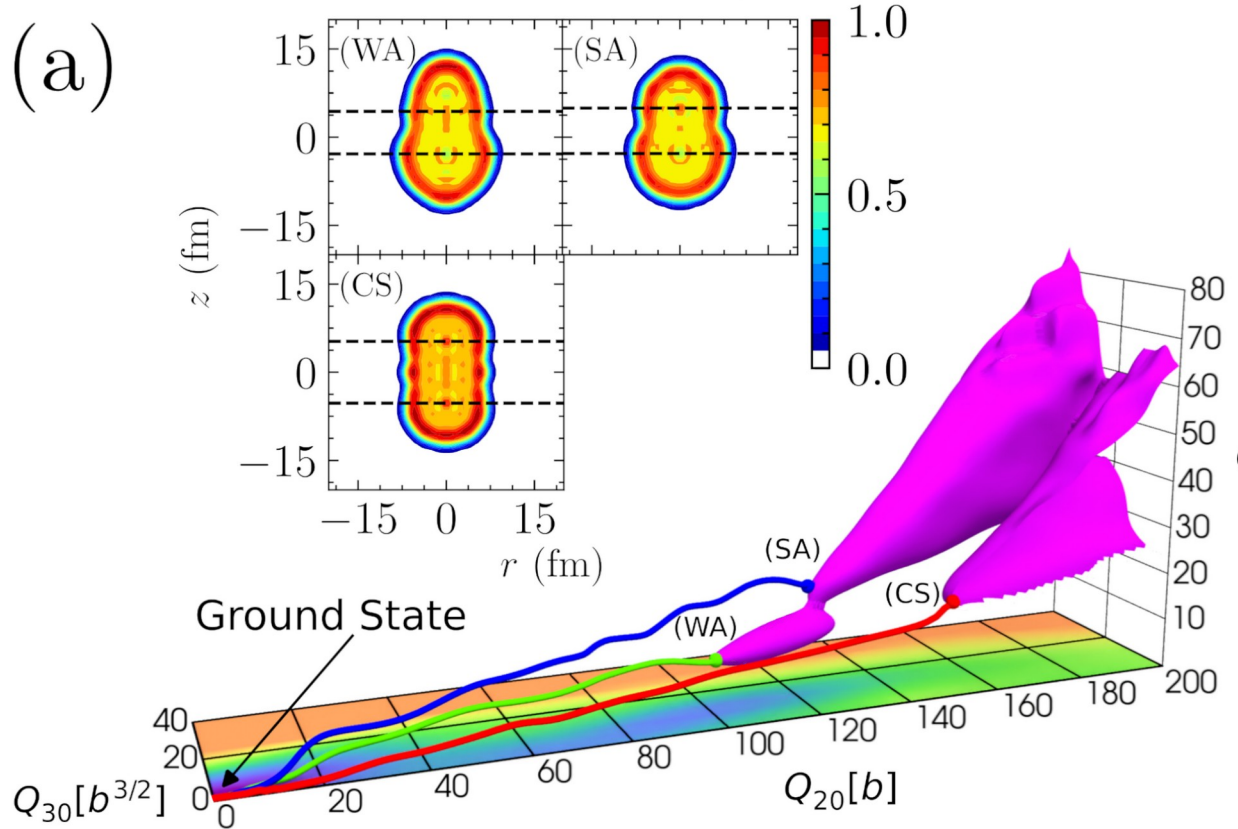
# Additional Slides



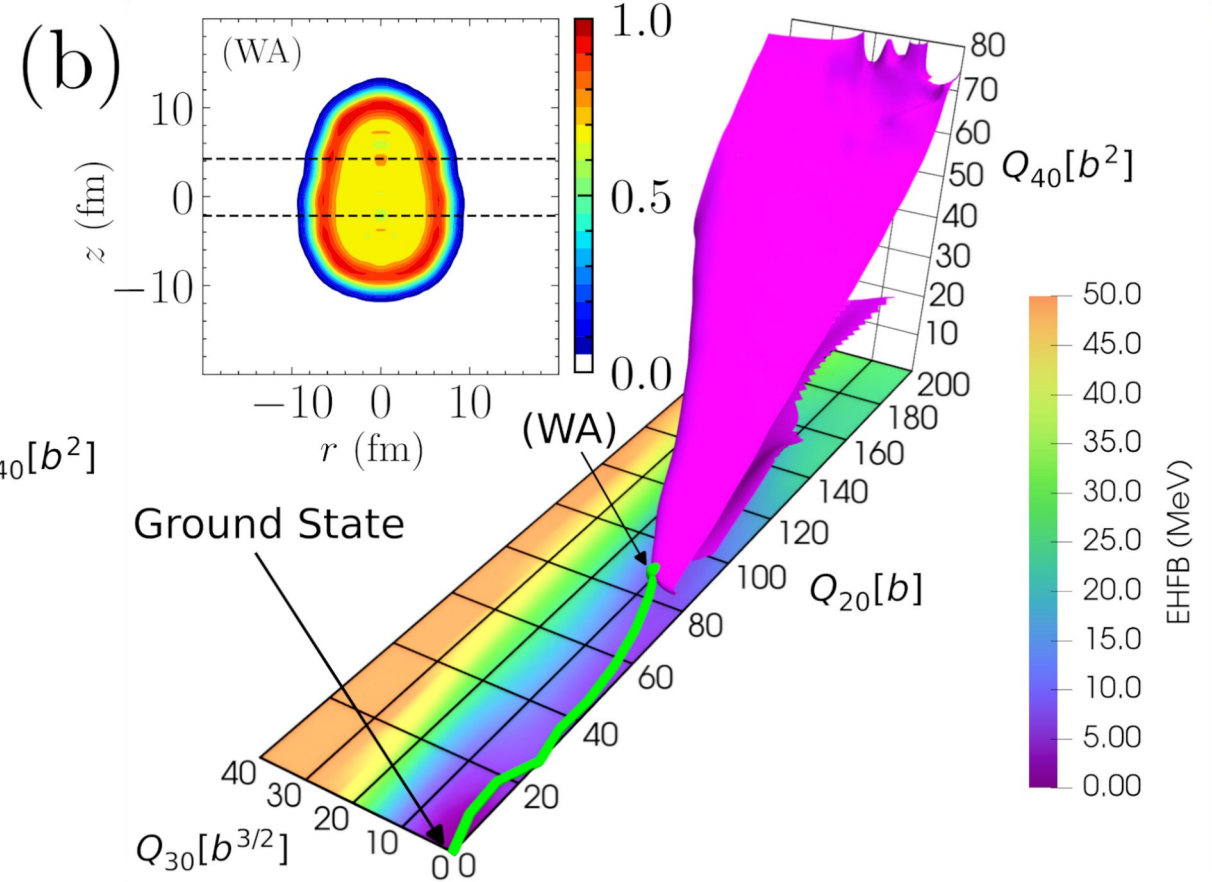
# Superheavy Results



# $^{306}\text{122}$ in 3-D

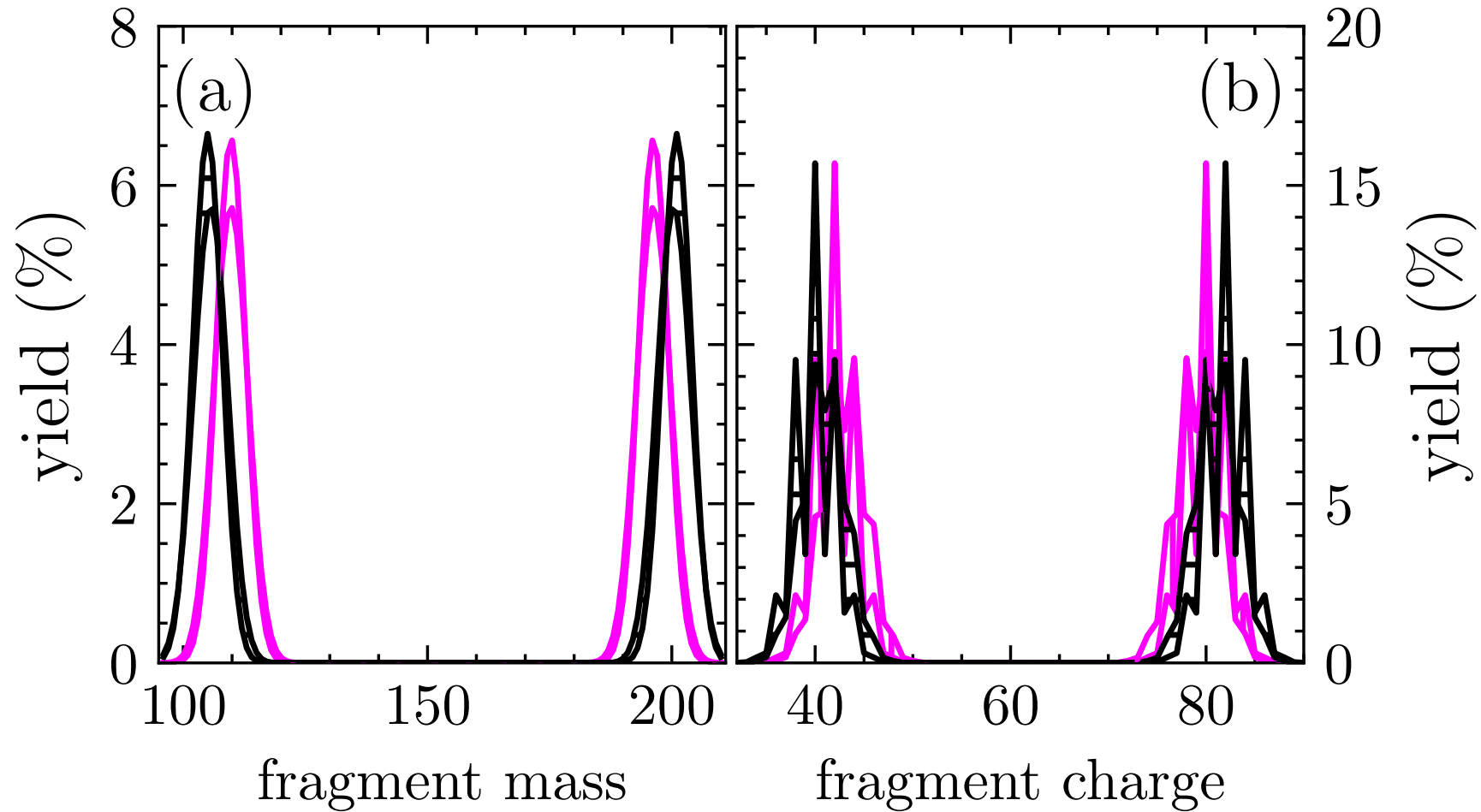


SkM\*



UNEDF1

# $^{306}\text{122}$ Yields



# Future Directions: Instantons

$$\left( -\hbar \frac{\partial}{\partial \beta} - h[\sigma] \right) \phi_{p_k}(x, \beta) = \varepsilon_{p_k}[\sigma, \beta_T] \phi_{p_k}(x, \beta)$$

$$\left( \hbar \frac{\partial}{\partial \beta} - h[\sigma] \right) \phi_{p_k}(x, -\beta) = \bar{\varepsilon}_{p_k}[\sigma, \beta_T] \phi_{p_k}(x, -\beta)$$