

Global theories of nuclear structure

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Hirscheegg
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1. HFB theory and its extensions
2. CEA/DAM survey of ground state properties and even-parity excited states

J.-P. Delaroche, M. Girod, J. Libert, H. Goutte, S. Hilaire, S. Peru, N. Pillet, and G.F. Bertsch,
Phys. Rev. C 81 014303 (2010)

Other related work:

D. Vretenar

P.-G. Reinhard

M. Bender, P.-H. Heenen

From model to theory

Characteristics of good theories

- need only a small set of parameters
- have wide predictive power
- have intrinsic criteria for limits of validity

Goals in this work

- apply theory globally (but with cuts generated by internal criteria)
- quantitative assessment of performance
- for theorists: report weaknesses as well as strengths
- for experimenters: predictions to be tested

The Hamiltonian

We would like to have an effective Hamiltonian:

$$H = ta^\dagger a + v^{(2)} a^\dagger a^\dagger a a + v^{(3)} a^\dagger a^\dagger a^\dagger a a a$$

but all we actually have to work with is an energy functional:

$$H = ta^\dagger a + v^{(2)} (r - r') a_r^\dagger a_{r'}^\dagger a_{r'} a_r + t_3 \rho(r)^{1/3} v^{(3)} a_r^\dagger a_r^\dagger a_r a_r$$

Self-consistent Mean-Field Theory

See the textbook (Ring and Schuck, 1980).

Example of Pb-208

²⁰⁸ Pb energy	Ska	D1S
Kinetic	3863	3920
Coulomb direct/exchange	831/-31	832/-31
Spin-orbit	-97	-105
Central 2B	-12480	-12783
t_3	6274	6530
Total	-1640	-1637

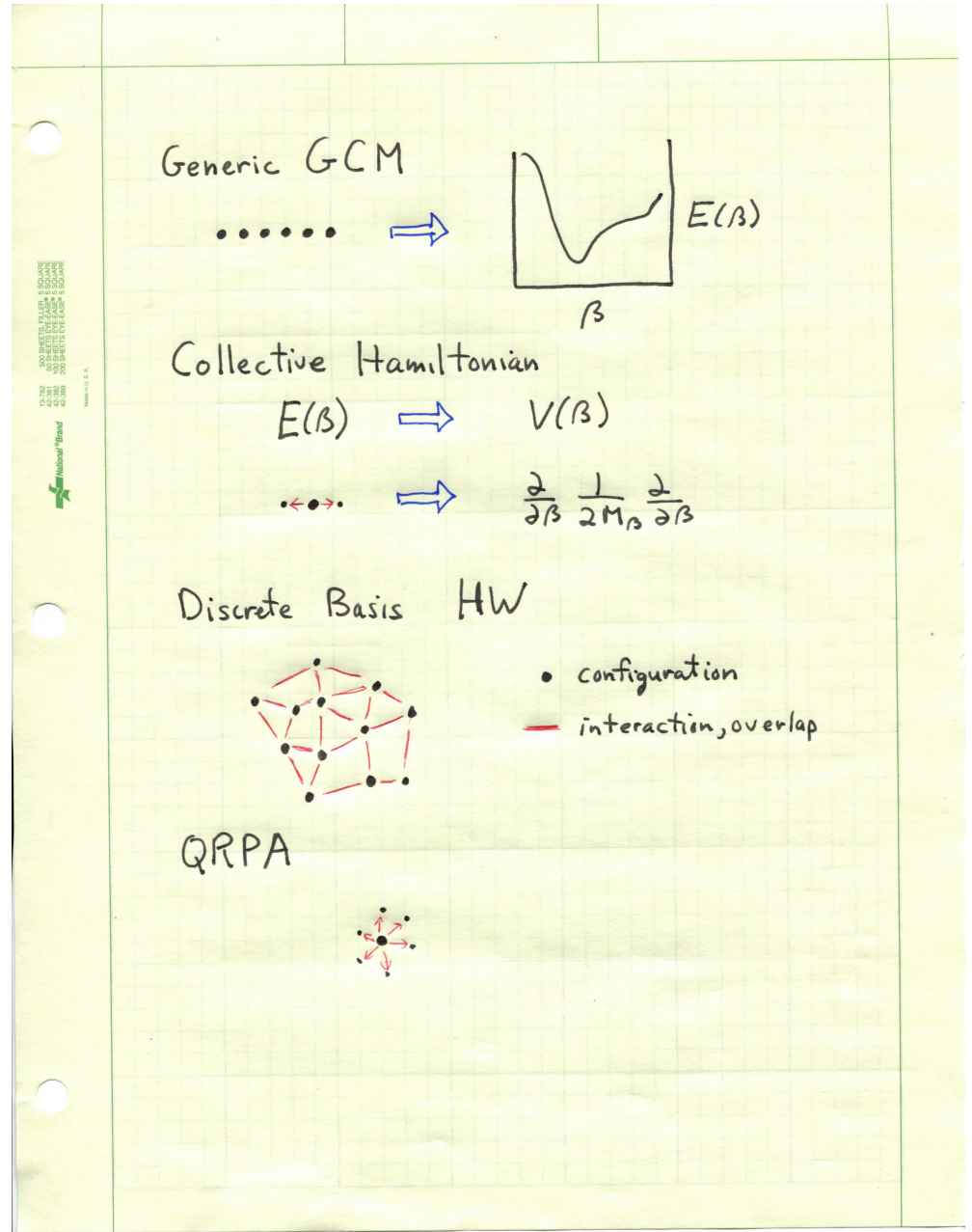
Extensions of self-consistent mean-field theory for spectroscopy

Generator Coordinate Methods

Collective Hamiltonian
GOA

Discrete-basis Hill-Wheeler

Quasiparticle RPA



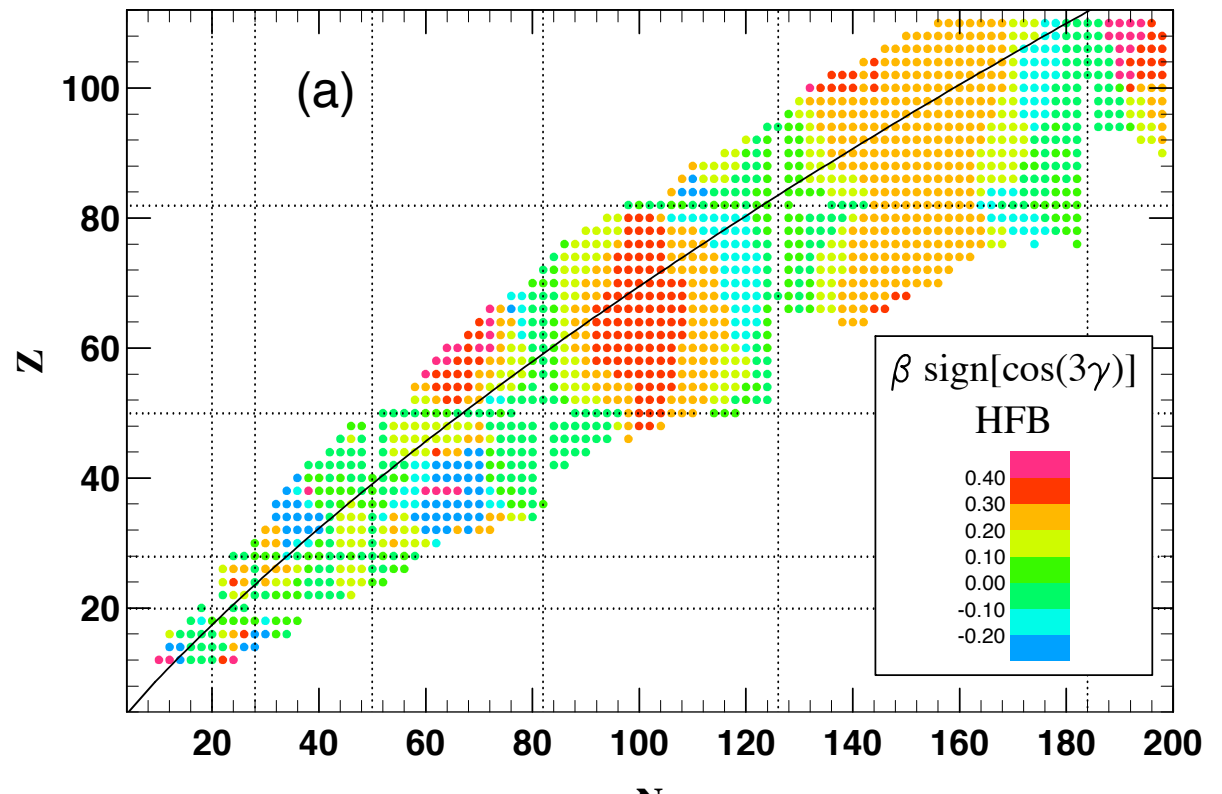
The CEA/DAM global survey of even-even spectroscopy

Mapped collective Hamiltonian method

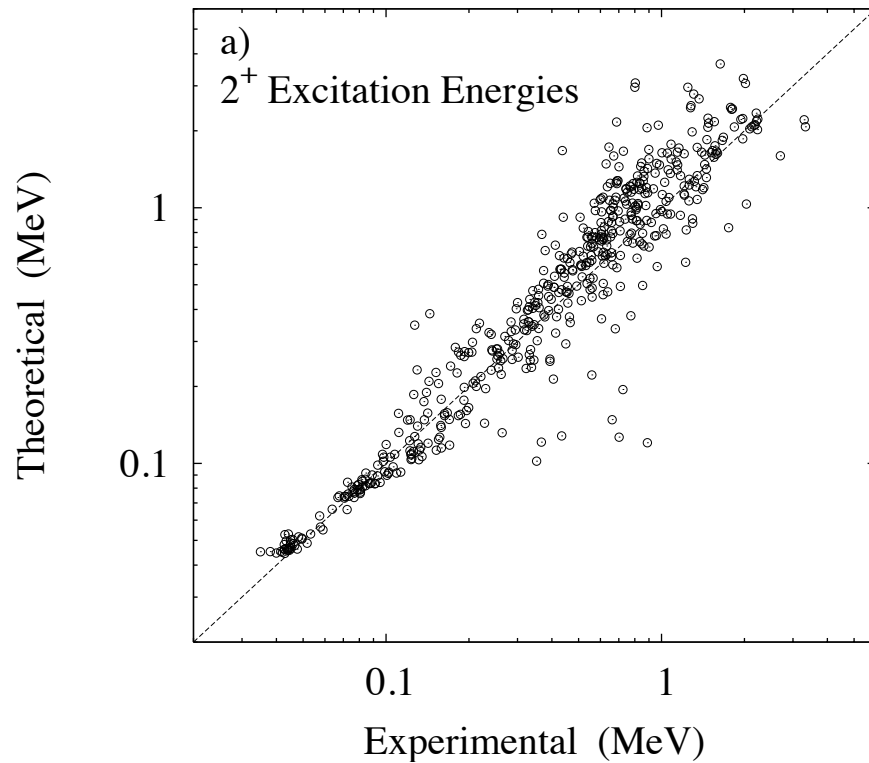
Gogny DIS interaction: CPC 63 (1991), 13 parameters

Computed spectroscopic observables for 1712 nuclei:

- yrast energies up to $J=6$
- excited 0^+ , first and second yrare $J=2$
- $B(E2)$ values for many of the transitions
- $E0$ matrix elements
- deformations, including triaxiality



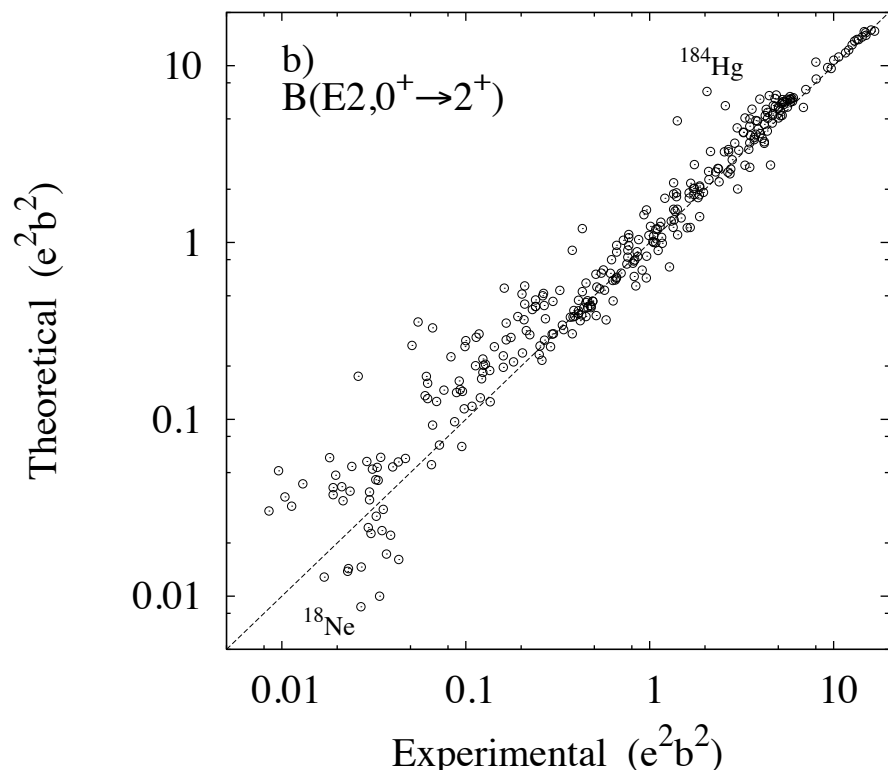
The first excited 2+ state



$$E_{theory} = (1.12 \pm 40\%) \times E_{exp} \quad \text{over 2 orders of magnitudes}$$

Dispersion from $\langle \log(E_{t}/E_{x}) \rangle$

Transition strengths



$$B(E2)_{theory} = (1.20 \pm 45\%) \times B(E2)_{exp}$$

An indicator of deformation: R_{42}

$$R_{42} = \frac{E(4_1^+)}{E(2_1^+)}.$$

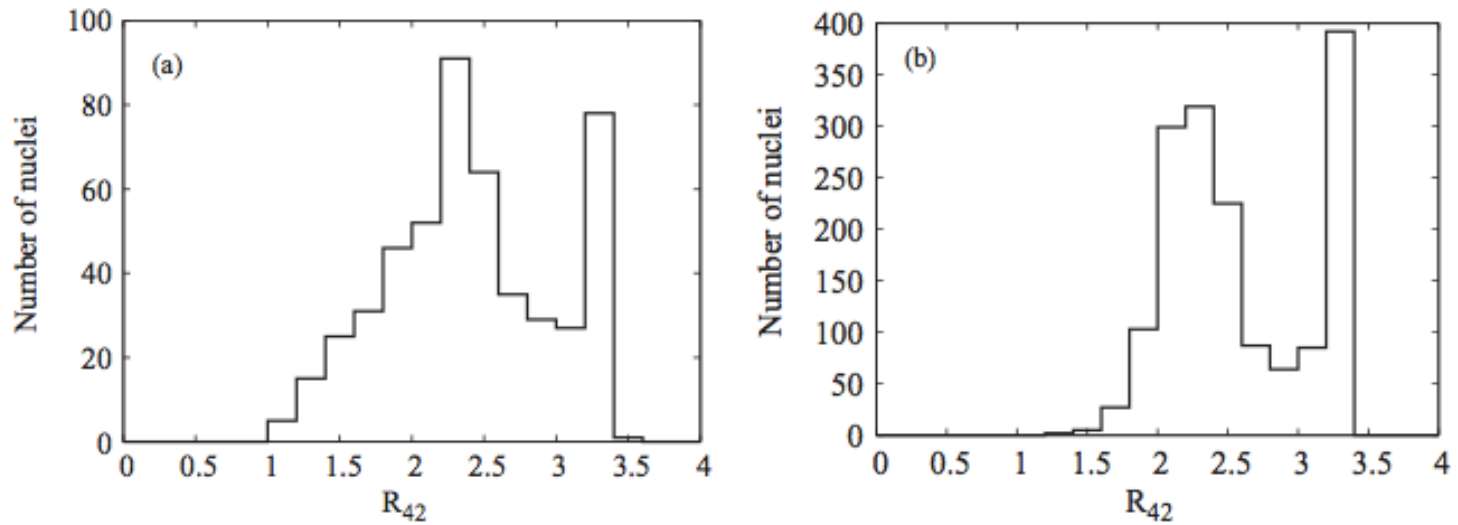


FIG. 12. (a) Histogram of experimental R_{42} ratios, Eq. (26), for 501 even-even nuclei, with data from Ref. [24]. (b) Histogram of calculated R_{42} ratios for 1609 even-even nuclei calculated in the CHFb+5DCH theory.

$$\frac{R_{42}(theory)}{R_{42}(exp.)} = 1.03 \pm 0.15$$

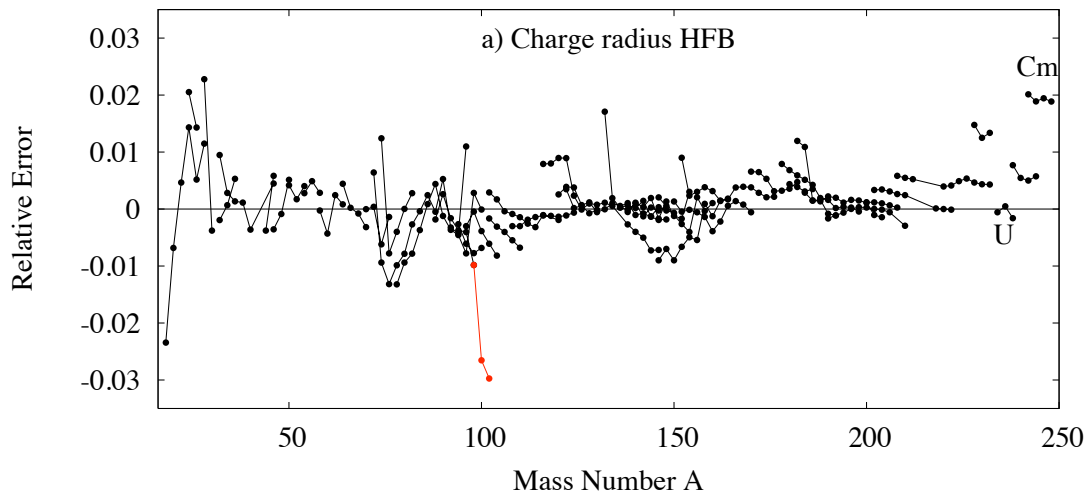
Predictive power for deformed nuclei

$$E_{theory} = (1.00 \pm 11\%) \times E_{exp} \quad 95 \text{ nuclei}$$

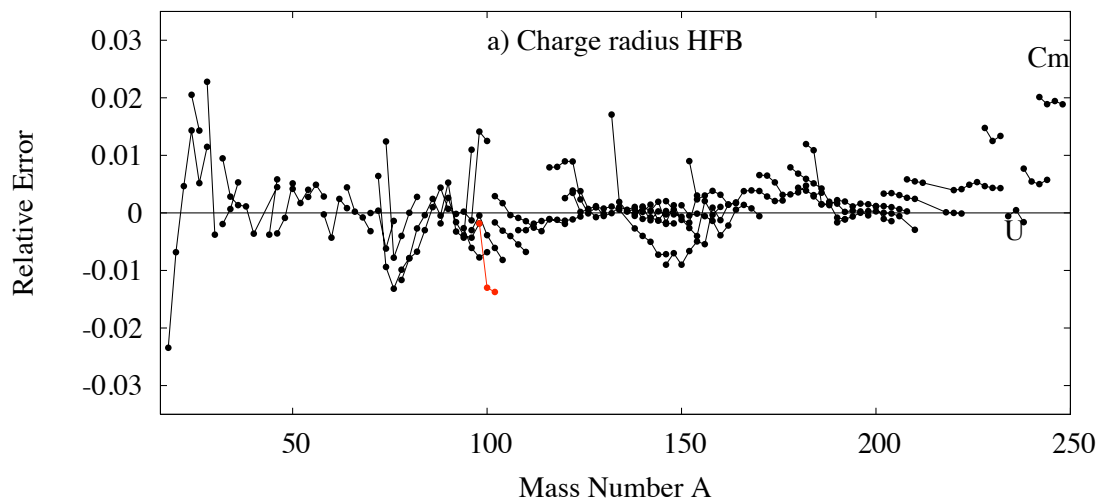
$$B(E2)_{theory} = (1.10 \pm 14\%) \times B(E2)_{exp} \quad 59 \text{ nuclei}$$

Charge radii

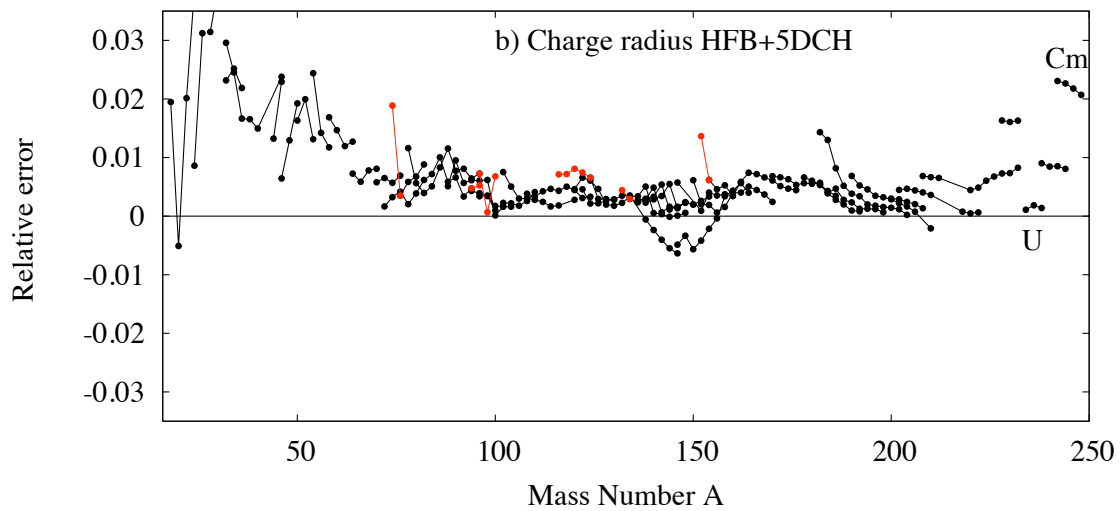
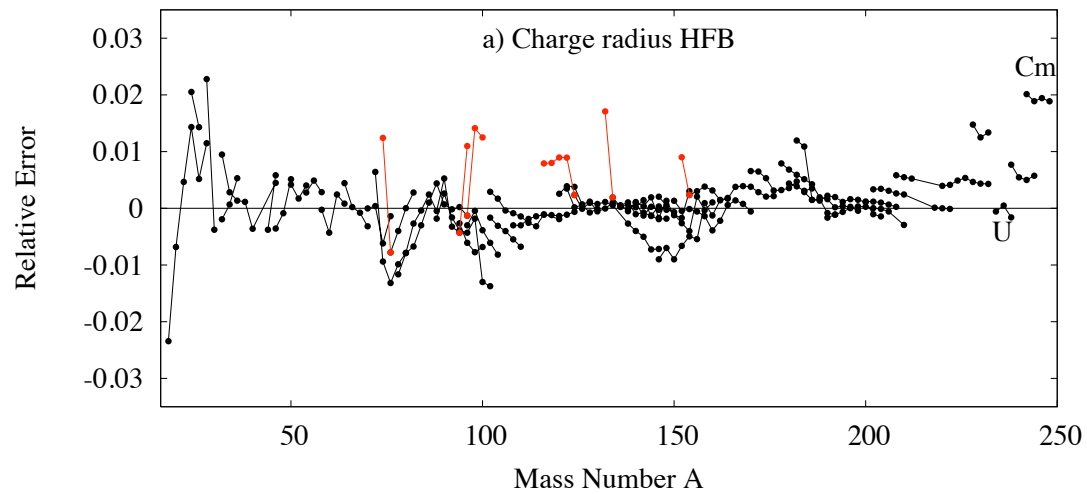
Experimental data from Angeli, ADNDT 87 (2004)



HFB



5DCH



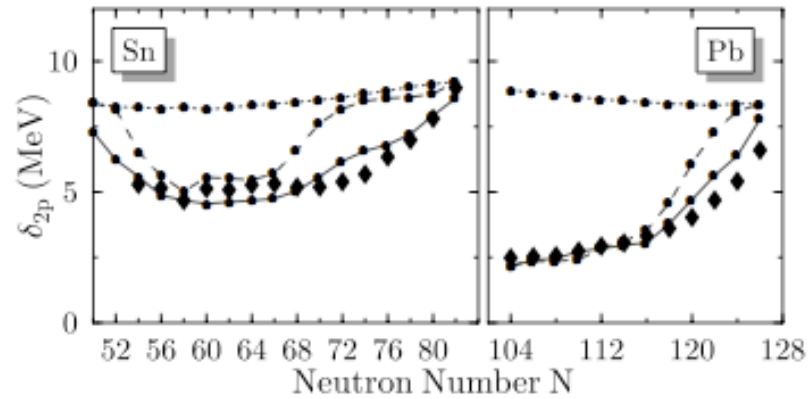
Performance on charge radius

TABLE II. Comparison of calculated charge radii with experiment: $\bar{\epsilon}$ is the mean of ϵ [see Eq. (18)]; σ is its rms dispersion about the average. Three hundred thirteen nuclear radii were included in the comparison as in Fig. 6. In the column “HFB (new)” we use the modern value $r_p = 0.875$ fm for the proton charge radius [48].

Theory	$\bar{\epsilon}$	σ
HFB	0.001	0.006
HFB (new)	0.005	0.007
CHFb+5DCH	0.006	0.007
Finite surface	0.0000	0.012

“Mutually enhanced magicity”

See Lunney, et al. RMP 75 (2003)

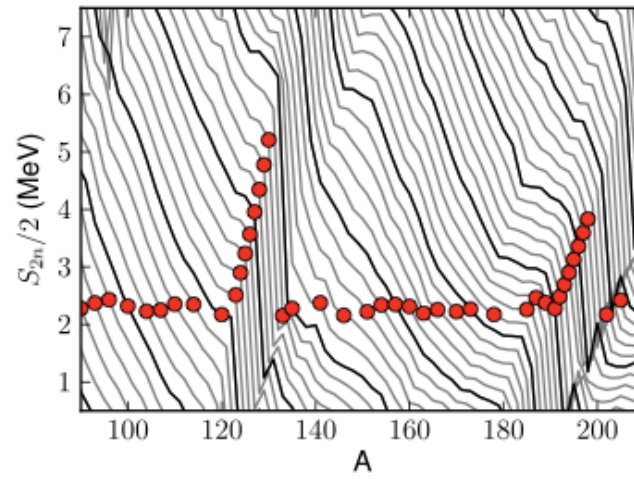
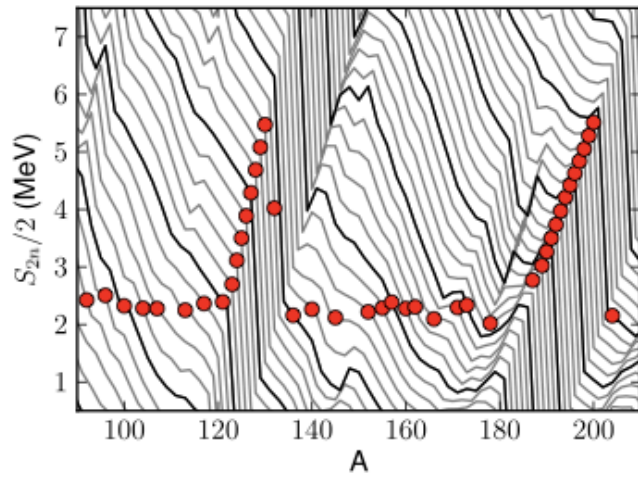


Bender, et al. PRL 94 (2005)

FIG. 2. Two-proton gaps, Eq. (3), for Pb and Sn isotopic chains. Theoretical curves are the following: spherical mean field (short dashed lines); mean field allowing for static deformations (long dashed lines); present theory (solid lines). Experimental values [1] are shown as diamonds.

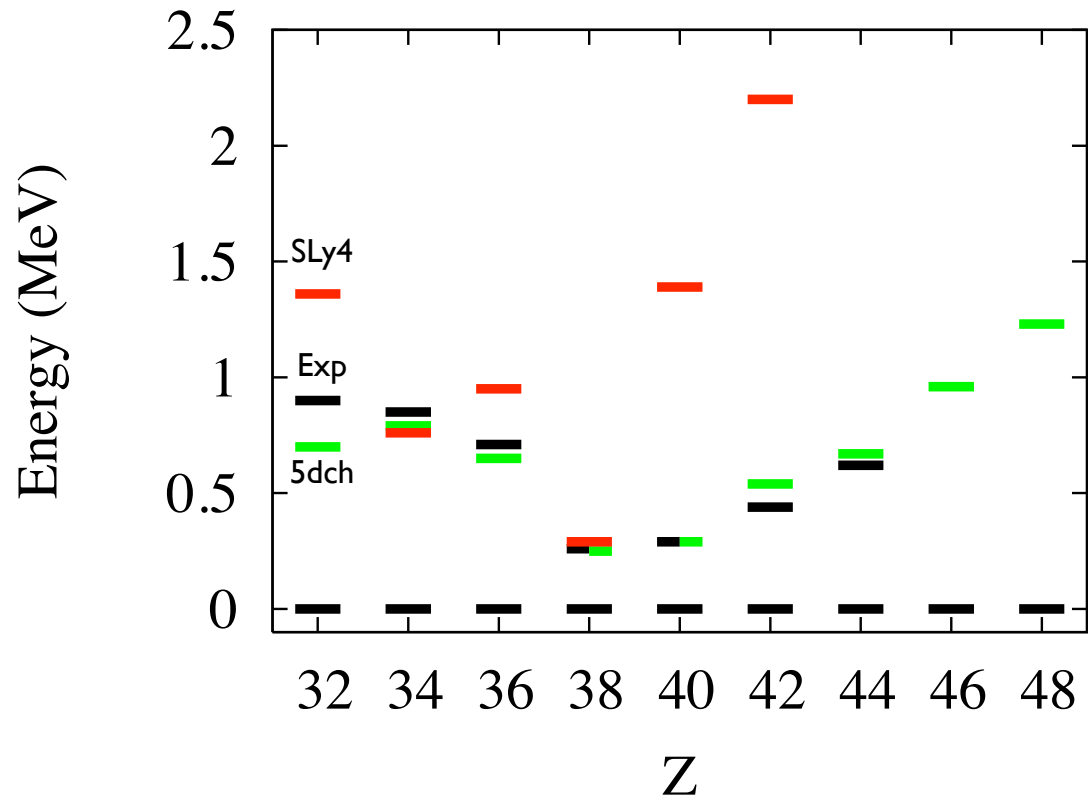
r-Process nucleosynthesis

Arcones and Martinez-Pinedo, PRC 83



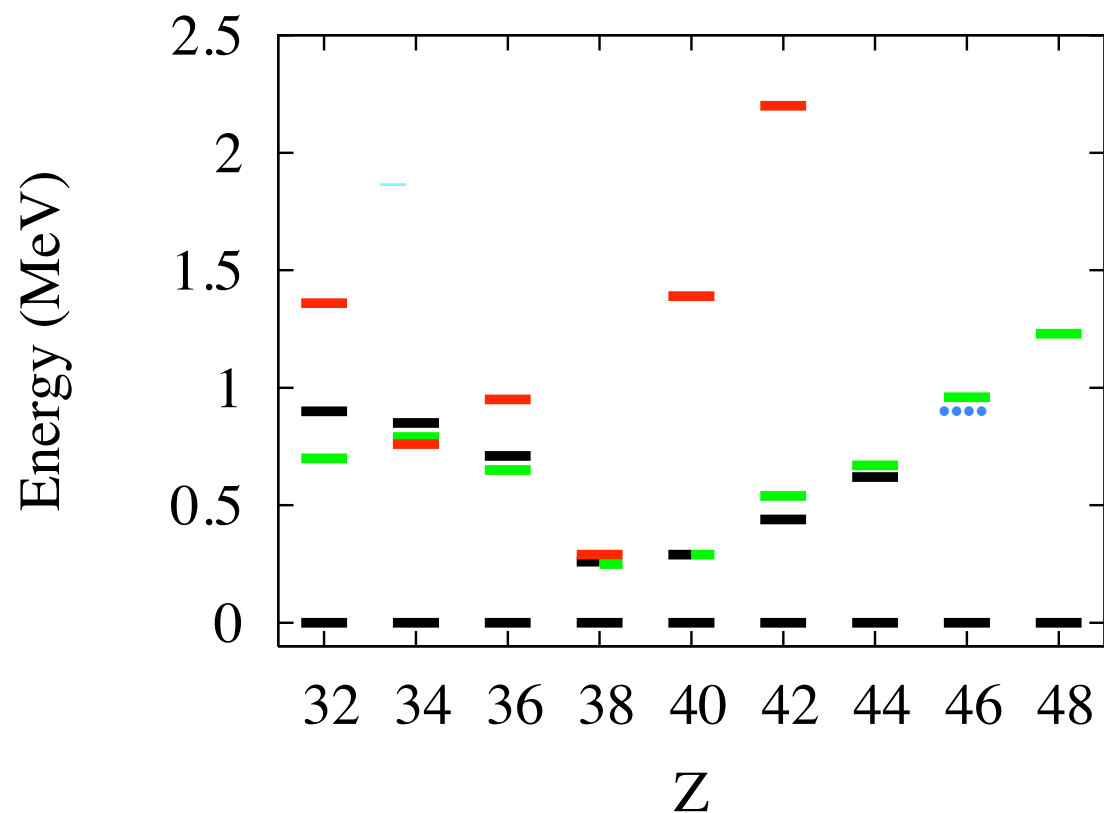
The CEA/DAM global survey, on the N=Z line

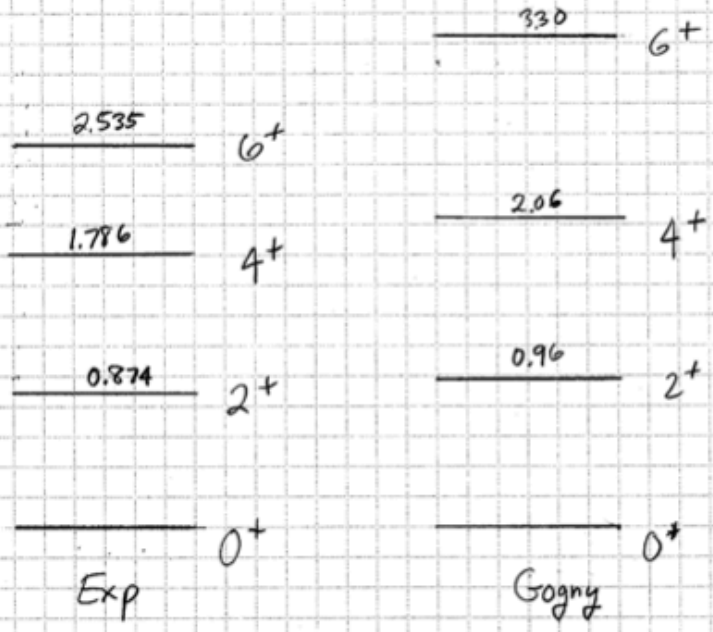
2+ Excitation energies



Evidence for a spin-aligned neutron-proton paired phase from the level structure of ^{92}Pd

B. Cederwall¹, F. Ghazi Moradi¹, T. Bäck¹, A. Johnson¹, J. Blomqvist¹, E. Clément², G. de France², R. Wadsworth³, K. Andgren¹, K. Lagergren^{1,4}, A. Dijon², G. Jaworski^{5,6}, R. Liotta¹, C. Qi¹, B. M. Nyakó⁷, J. Nyberg⁸, M. Palacz², H. Al-Azri³, A. Algora⁹, G. de Angelis¹⁰, A. Ataç¹¹, S. Bhattacharyya^{2,†}, T. Brock³, J. R. Brown³, P. Davies³, A. Di Nitto¹², Zs. Dombrádi⁷, A. Gadea⁹, J. Gál⁷, B. Hadinia¹, F. Johnston-Theasby³, P. Joshi³, K. Juhász³, R. Julin¹⁴, A. Jungclaus¹⁵, G. Kalinka⁷, S. O. Kara¹¹, A. Khaplanov¹, J. Kownacki⁵, G. La Rana¹², S. M. Lenzi¹⁶, J. Molnár⁷, R. Moro¹², D. R. Napoli¹⁰, B. S. Nara Singh³, A. Persson¹, F. Recchia¹⁶, M. Sandzelius^{1,†}, J.-N. Scheurer¹⁷, G. Sletten¹⁸, D. Sohler⁷, P.-A. Söderström⁸, M. J. Taylor³, J. Timár⁷, J. J. Valiente-Dobón¹⁰, E. Vardaci¹² & S. Williams¹⁹





92Pd

Phys. Rev. C 81, 014303 (2010) [23 pages]

Structure of even-even nuclei using a mapped collective Hamiltonian and the D1S Gogny interaction

Abstract

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The CEA/DAM global survey (HFB/GCM/5DCH)

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Computed spectroscopic observables for 1712 nuclei:

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- deformations, including triaxiality

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44	102	-1001.112	0.290	24.	4.84	4.81	5.32	-1004.872	-3.760	(
46	42	-699.649	0.000	0.	4.37	4.31	4.19	-701.419	-1.770	(
46	44	-730.992	0.000	0.	4.37	4.32	4.22	-733.013	-2.021	(
46	46	-761.066	0.106	0.	4.38	4.33	4.26	-763.027	-1.961	(
46	48	-789.873	0.000	0.	4.38	4.33	4.28	-790.963	-1.090	(
46	50	-817.521	0.000	0.	4.38	4.33	4.30	-816.719	0.802	(
46	52	-836.390	0.000	0.	4.40	4.35	4.34	-838.888	-2.498	(

How to do GCM

\hat{H} = the many-body Hamiltonian, usually approximated by an EDF.

\hat{Q}_i = a set of one-body operators

I) Minimize $\langle \psi_\lambda | \hat{H} - \sum_i \lambda_i \hat{Q}_i | \psi_\lambda \rangle$ to find ψ_λ

II find expectation values $q_i = \langle \psi_\lambda | \hat{Q}_i | \psi_\lambda \rangle$

$$V(q) \equiv V(\lambda(q)) = \langle \psi_\lambda | \hat{H} | \psi_\lambda \rangle$$

This is the potential energy surface.

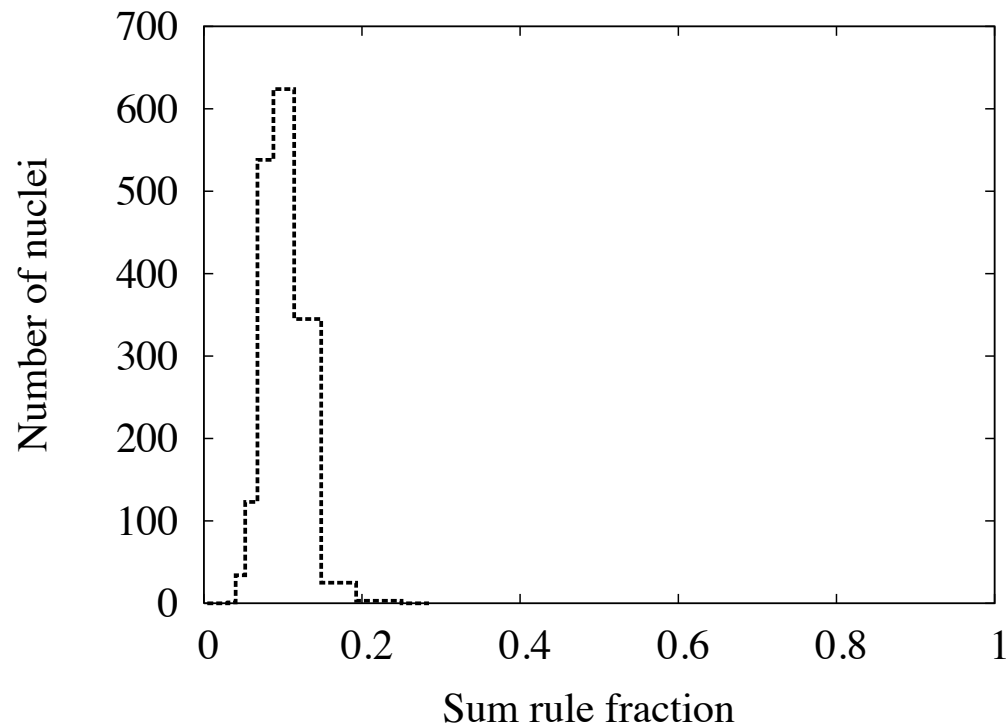
5DCH work: $\hat{N}, \hat{Z}, r^2 Y_{20}, r^2 Y_{22}, \hat{J}_x$

octupole study: $\hat{N}, \hat{Z}, r^2 Y_{20}, r^3 Y_{30}$

Sum rule

$$S = \sum_i E(2_i^+) B(E2; 0_1^+ \rightarrow 2_i^+) = \frac{25}{4\pi} \left(\frac{\hbar^2}{m} \right) Z^2 \langle r^2 \rangle$$

The fraction of the sum rule in the lowest excitation is $\sim 10\%$.



Separation Energies

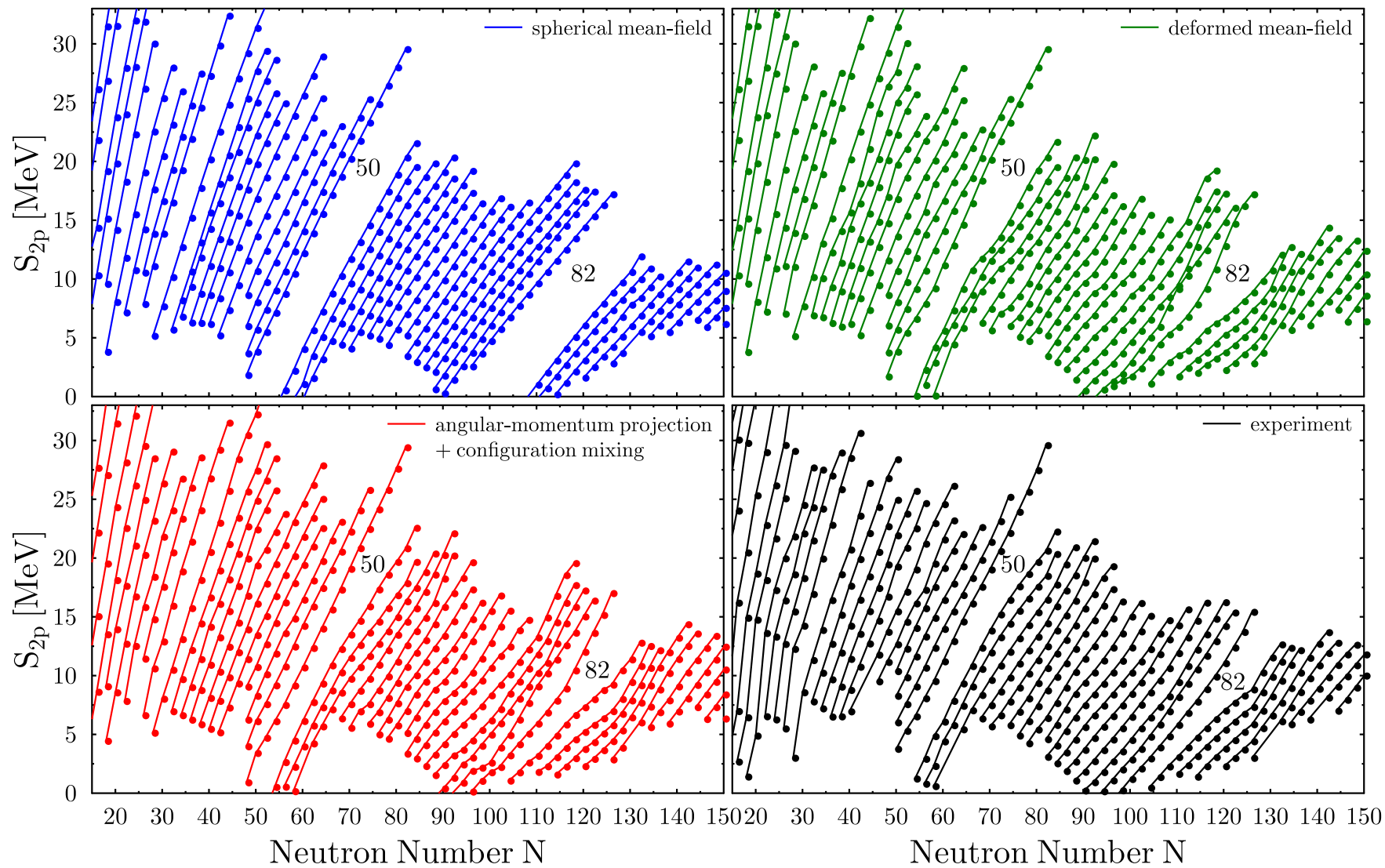


FIG. 2: Two-proton separation energies for isotonic chains.