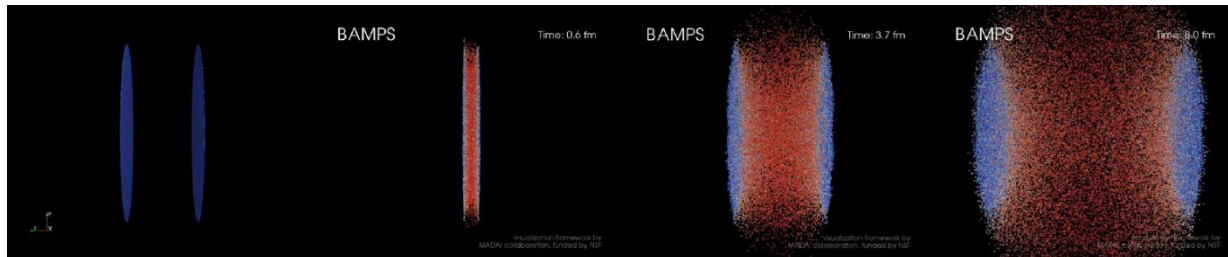


Heavy flavor at RHIC and LHC in a partonic transport model

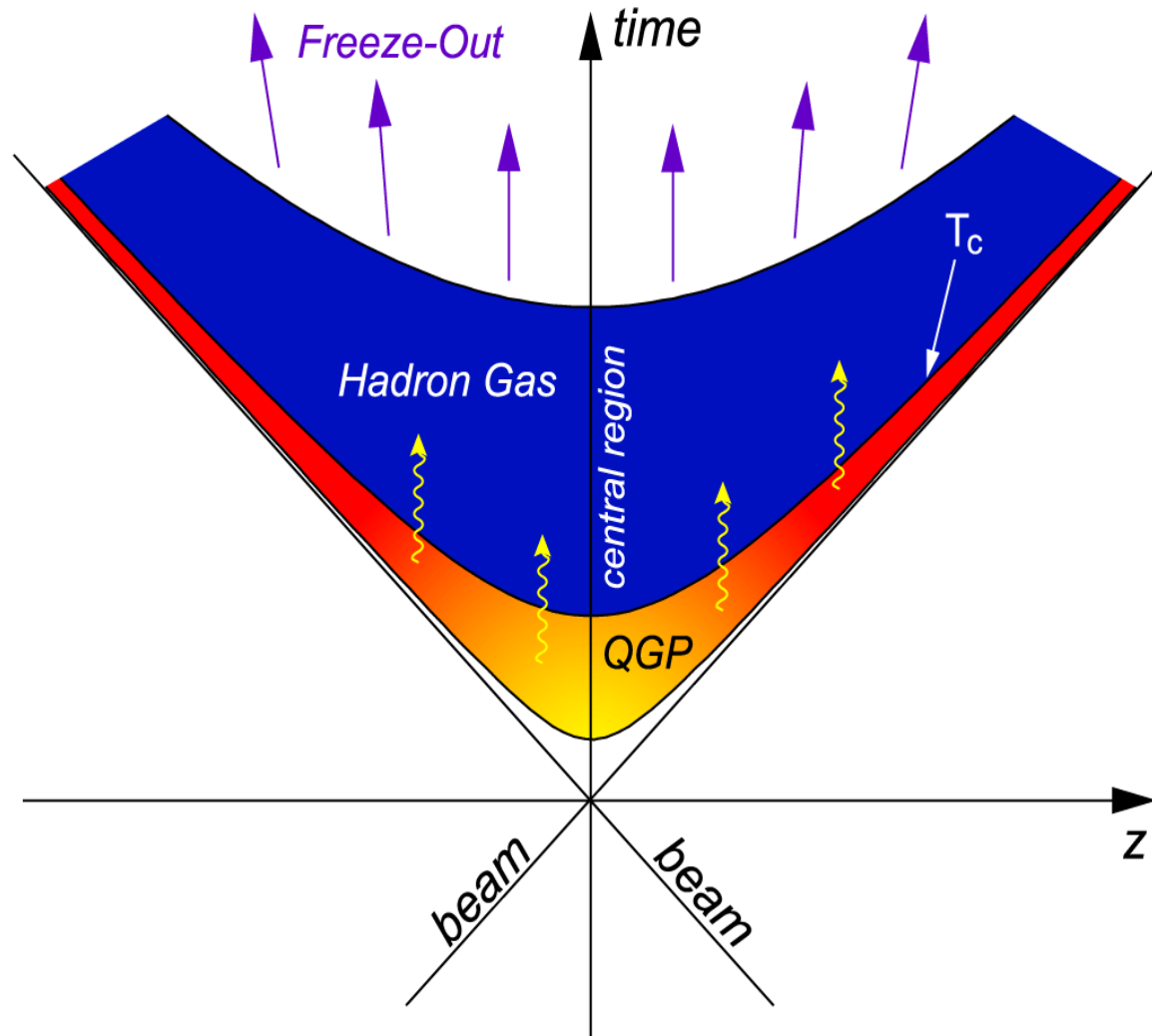
Jan Uphoff

with O. Fochler, K. Zhou, Z. Xu and C. Greiner

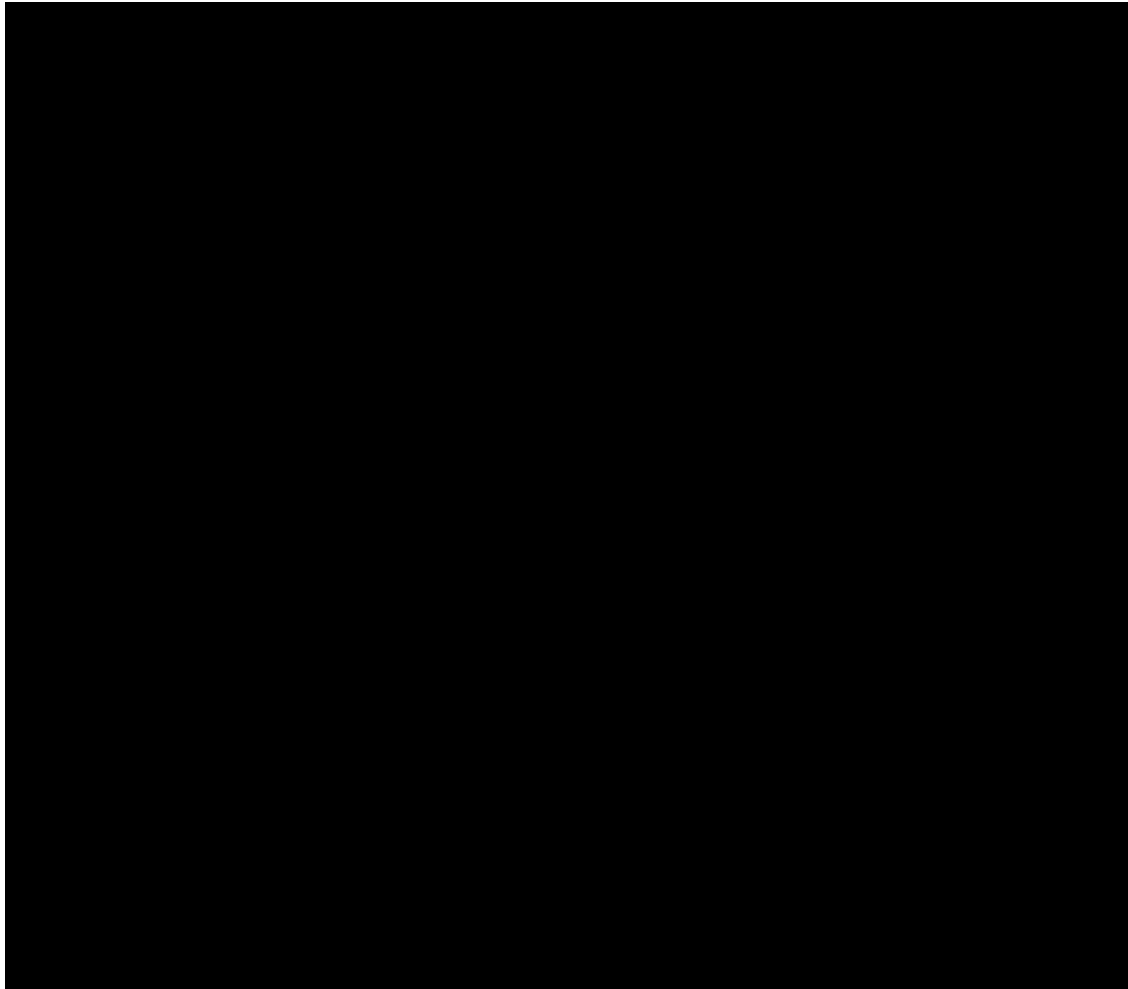
Based on arXiv:1104.2295 and 1104.2437



Ultrarelativistic heavy ion collision



BAMPS simulation of QGP phase at LHC at $s_{NN} = 2.76$ TeV



Visualization framework
courtesy MADAI
collaboration, funded by
the NSF under grant# NSF-
PHY-09-41373

Heavy quarks in heavy-ion collisions

Large heavy quark mass

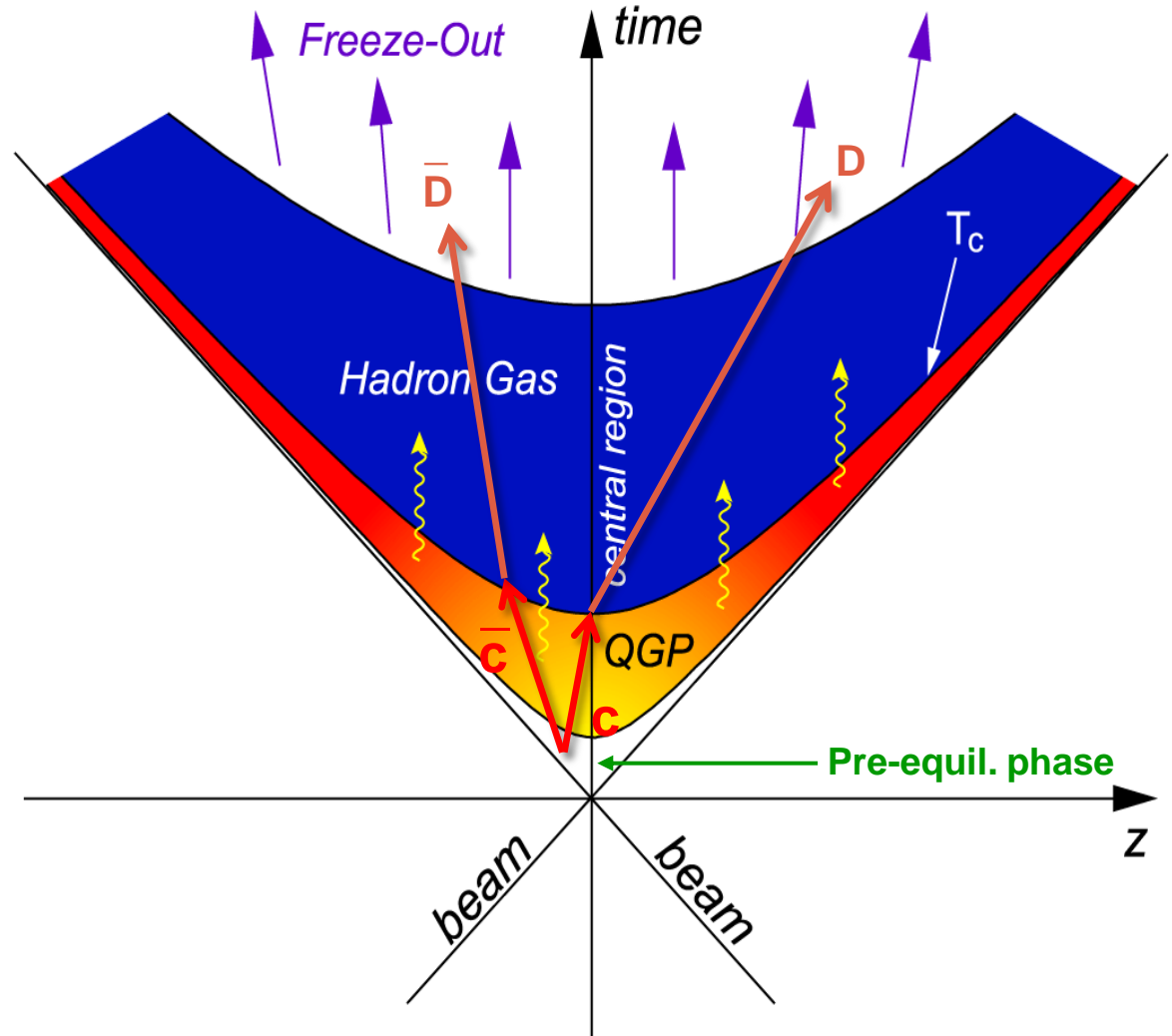
$$\gg \Lambda_{\text{QCD}}$$

Charm: $M_c \approx 1.5 \text{ GeV}$

Bottom: $M_b \approx 4.75 \text{ GeV}$

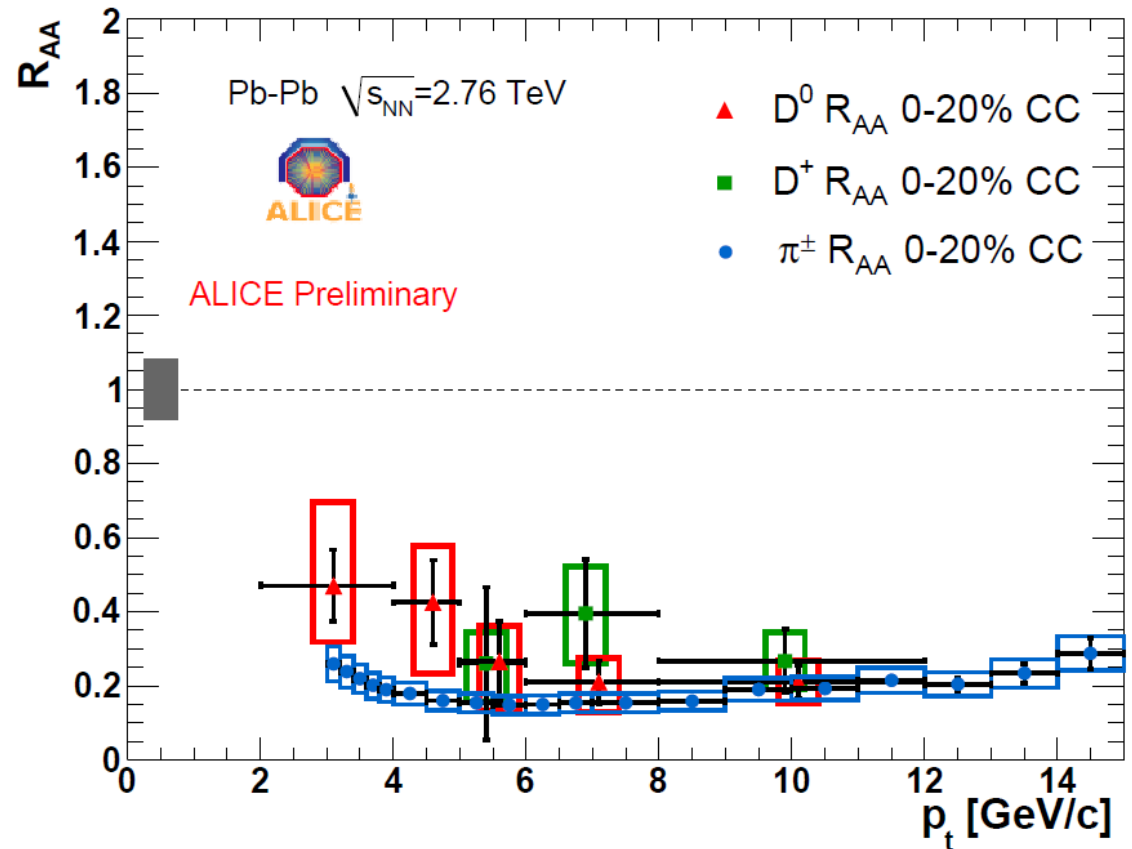
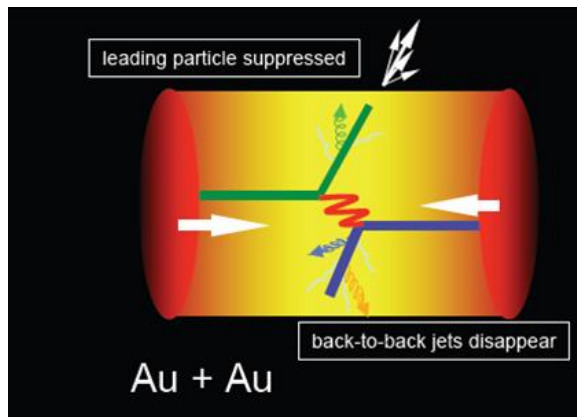
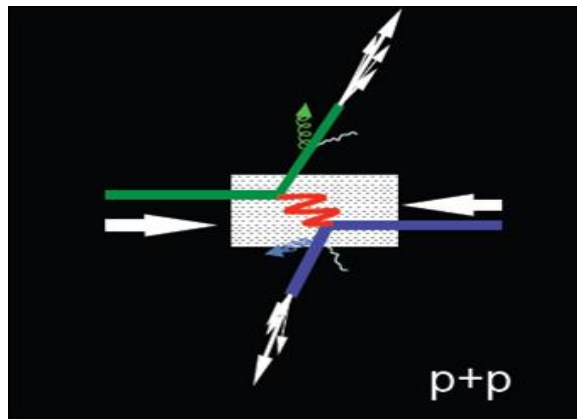
➔ Heavy quark production at early stage of collision

➔ ideal probe for this stage



Nuclear modification factor

$$R_{AA} = \frac{dN/dp_T dy|_{A+A}}{N_{\text{bin}} dN/dp_T dy|_{p+p}}$$

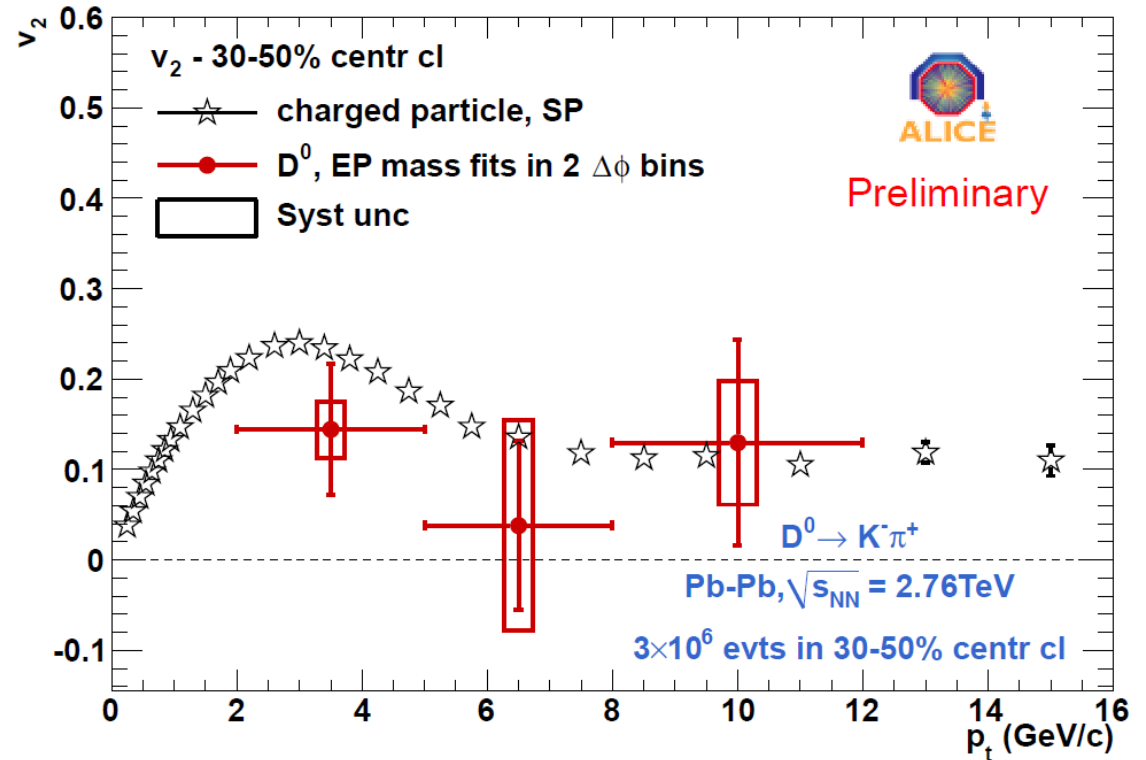
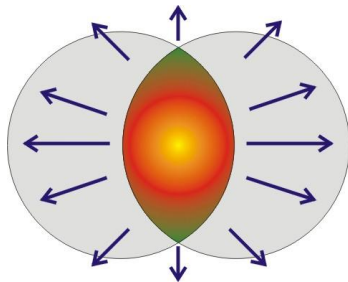
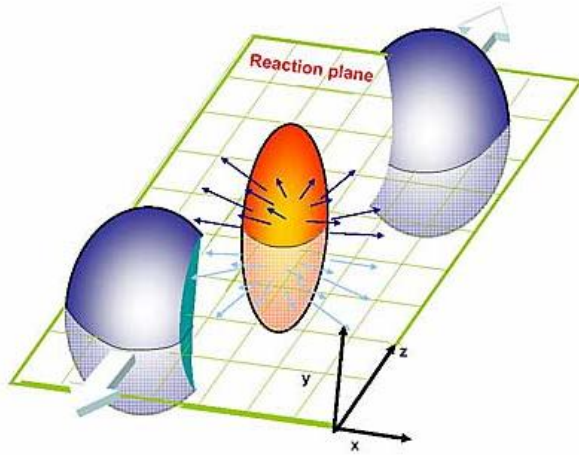


Motivation

Elliptic flow

$$v_2 = \left\langle \frac{p_x^2 - p_y^2}{p_T^2} \right\rangle$$

$$\frac{d^3N}{p_T dp_T dy d\phi}(p_T, y, \phi) = \frac{1}{2\pi} \frac{d^2N}{p_T dp_T dy} [1 + 2v_2(p_T, y) \cos(2\phi) + \dots]$$



BAMPS: Boltzmann Approach of MultiParton Scatterings

- 3+1 dimensional, fully dynamic parton transport model
- solves the Boltzmann equations for on-shell partons with pQCD interactions

$$\left(\frac{\partial}{\partial t} + \frac{\mathbf{p}_i}{E_i} \frac{\partial}{\partial \mathbf{r}} \right) f_i(\mathbf{r}, \mathbf{p}_i, t) = \mathcal{C}_i^{2 \rightarrow 2} + \mathcal{C}_i^{2 \leftrightarrow 3} + \dots$$

Z. Xu & C. Greiner,
Phys. Rev. C71 (2005)
Phys. Rev. C76 (2007)

Implemented processes:

$$g + g \rightarrow g + g$$

$$g + g \rightarrow g + g + g$$

$$g + g + g \rightarrow g + g$$

Light quarks have been implemented but are not included in the present calculation

$$g + g \rightarrow Q + \bar{Q}$$

$$Q + \bar{Q} \rightarrow g + g$$

$$g + Q \rightarrow g + Q$$

$$g + \bar{Q} \rightarrow g + \bar{Q}$$

$$g + J/\psi \rightarrow c + \bar{c}$$

$$c + \bar{c} \rightarrow g + J/\psi$$

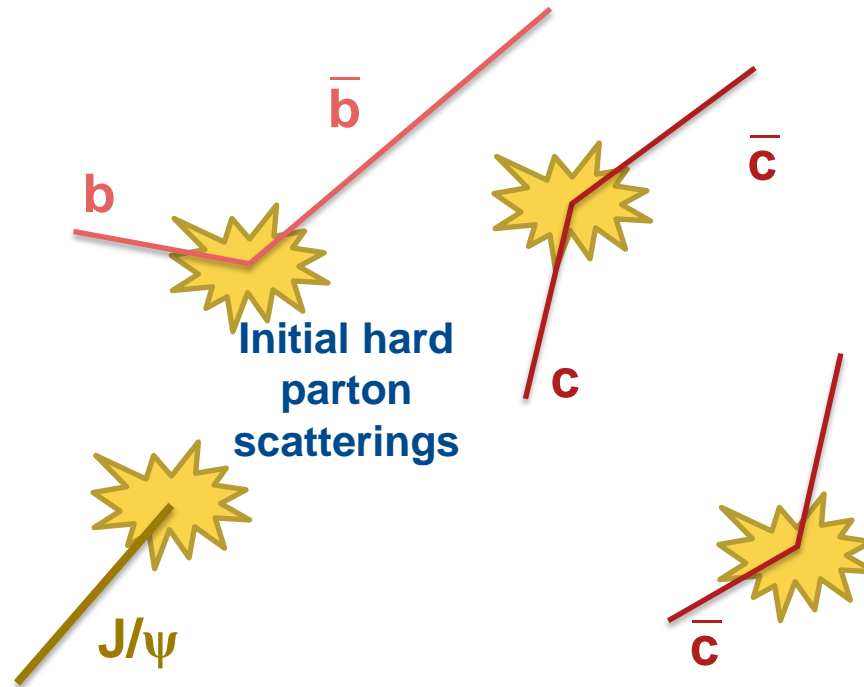
Sketch of heavy-ion collision in BAMPS

Heavy flavor in BAMPS



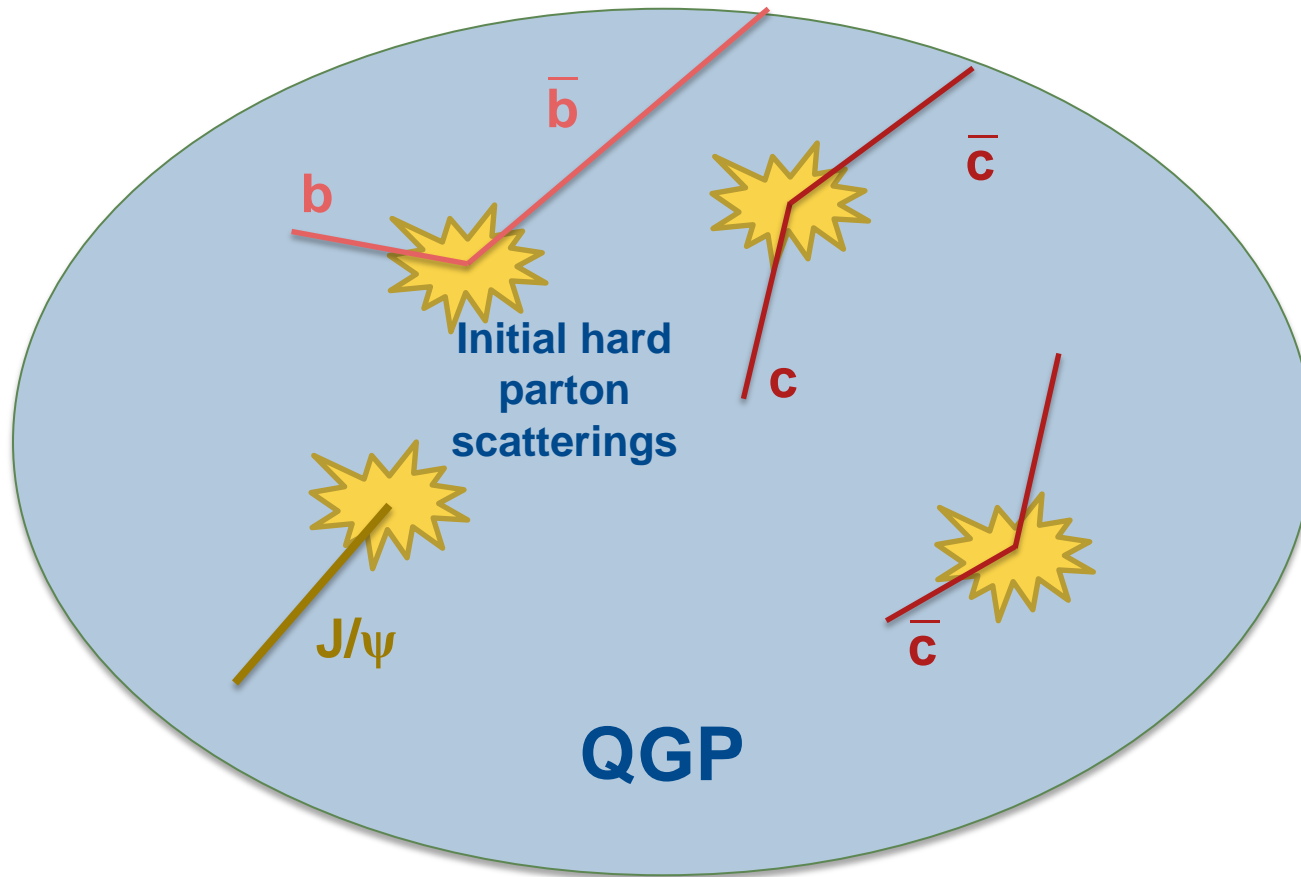
Sketch of heavy-ion collision in BAMPS

Heavy flavor in BAMPS



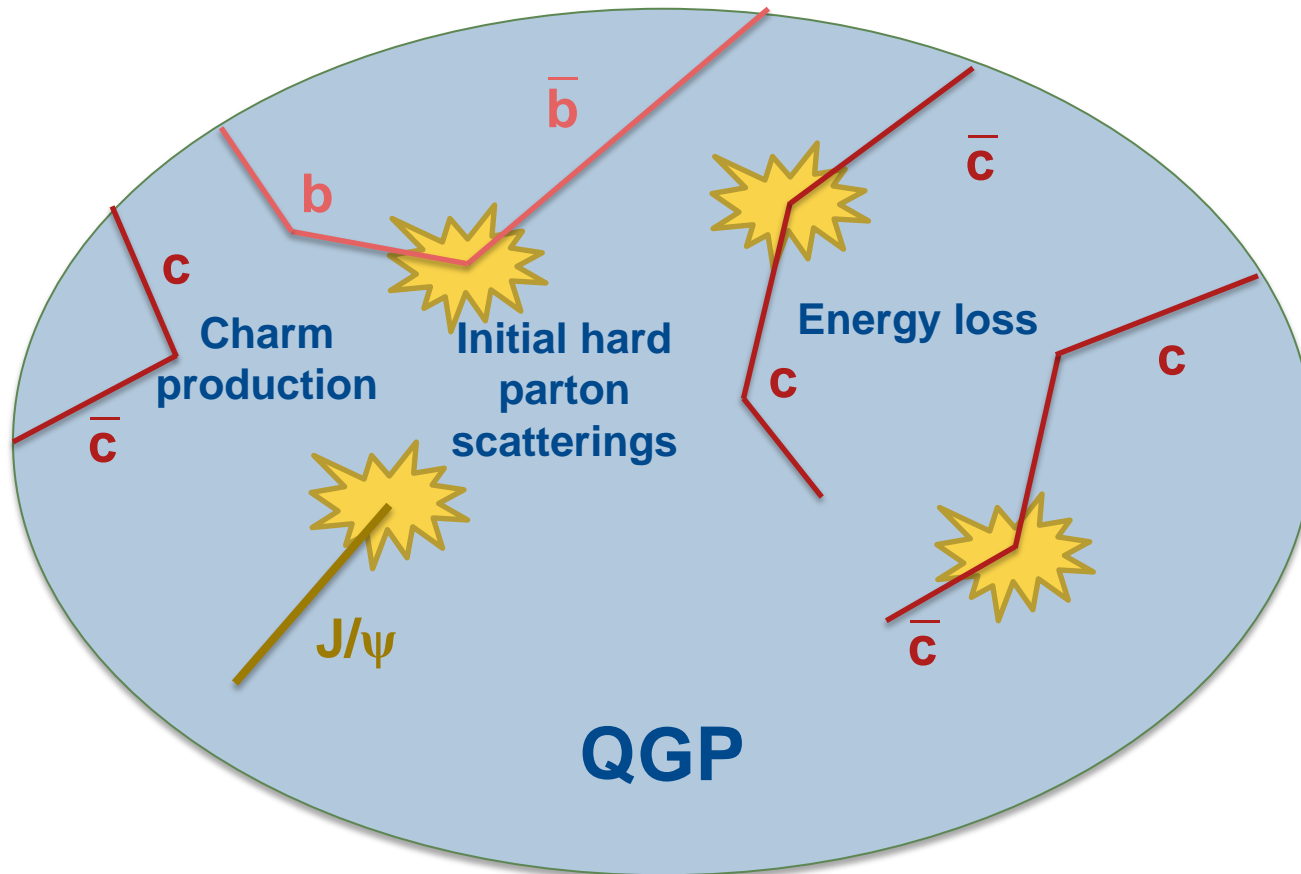
Sketch of heavy-ion collision in BAMPS

Heavy flavor in BAMPS



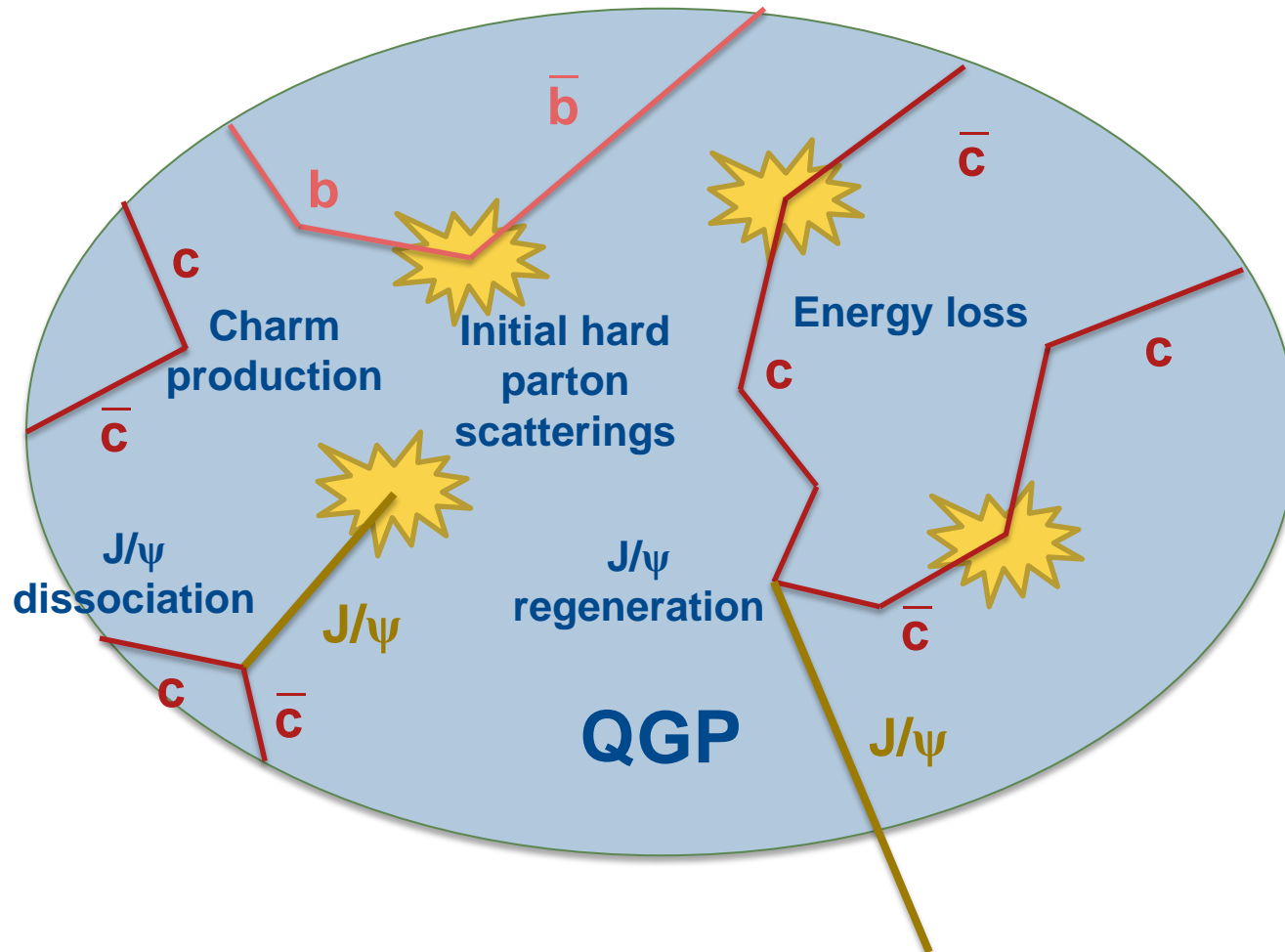
Sketch of heavy-ion collision in BAMPS

Heavy flavor in BAMPS



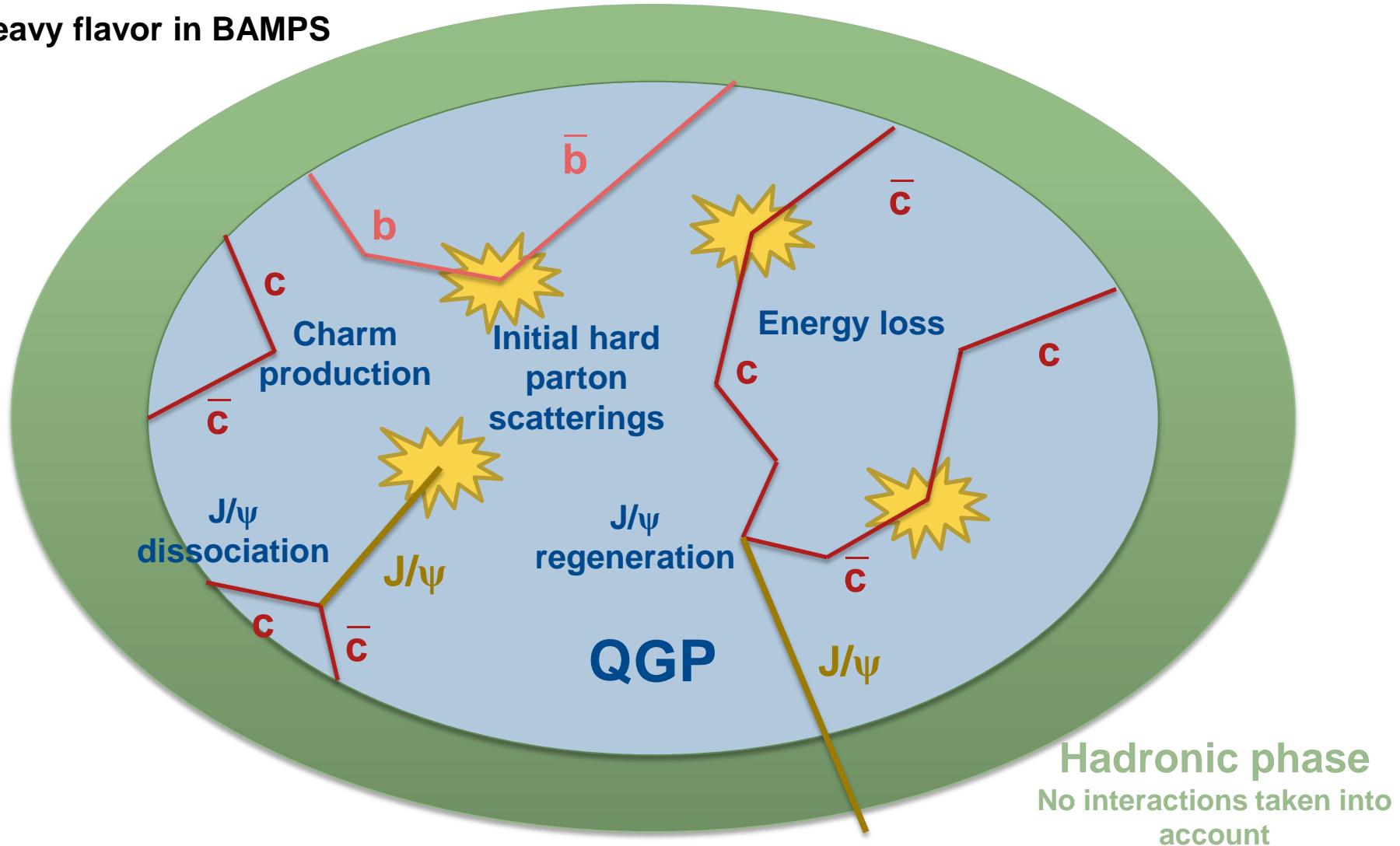
Sketch of heavy-ion collision in BAMPS

Heavy flavor in BAMPS



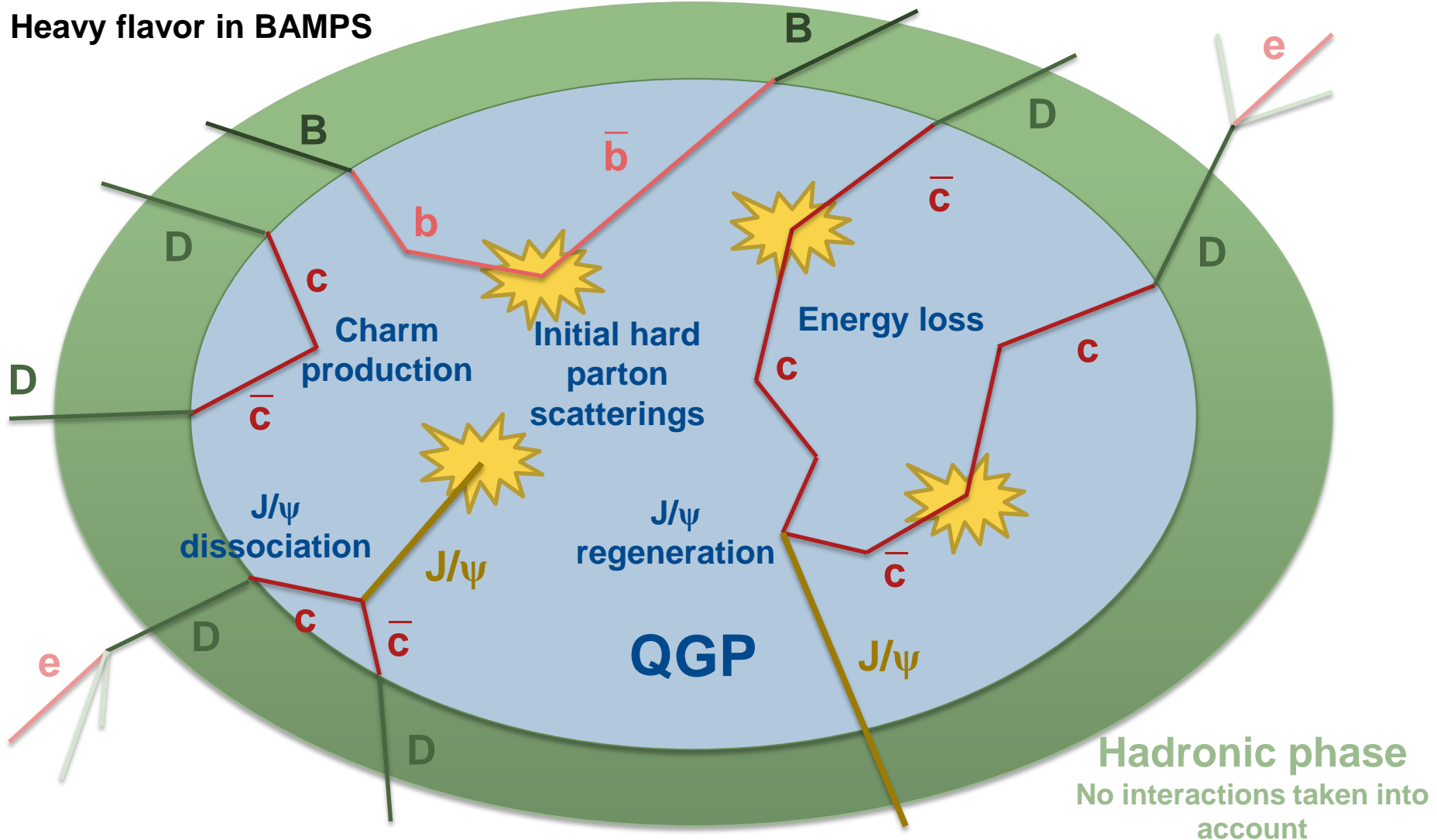
Sketch of heavy-ion collision in BAMPS

Heavy flavor in BAMPS



Sketch of heavy-ion collision in BAMPS

Heavy flavor in BAMPS



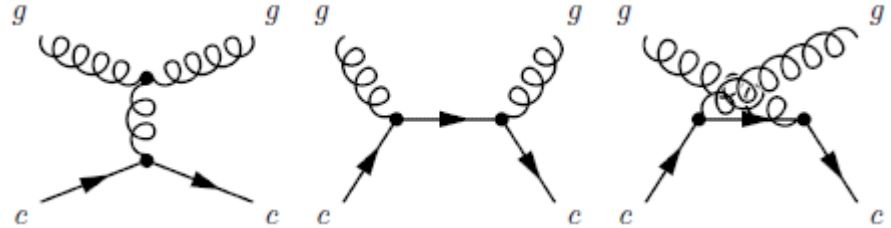
Heavy quark scattering

Leading order perturbative QCD:

$$g + Q \rightarrow g + Q$$

$$g + \bar{Q} \rightarrow g + \bar{Q}$$

t channel is divergent for small t



$$\frac{1}{t} \rightarrow \frac{1}{t - \kappa m_D^2}$$

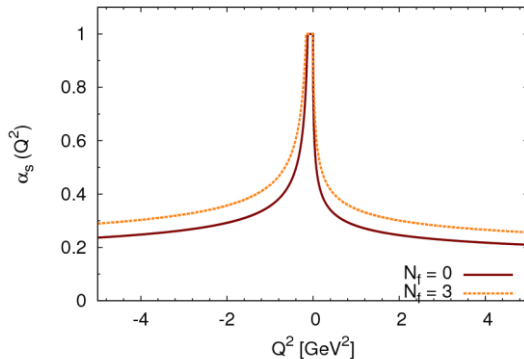
κ can be fixed to

$$\kappa = \frac{1}{2e} \approx 0.184 \approx 0.2$$

by comparing dE/dx to
HTL result beyond
logarithmic accuracy

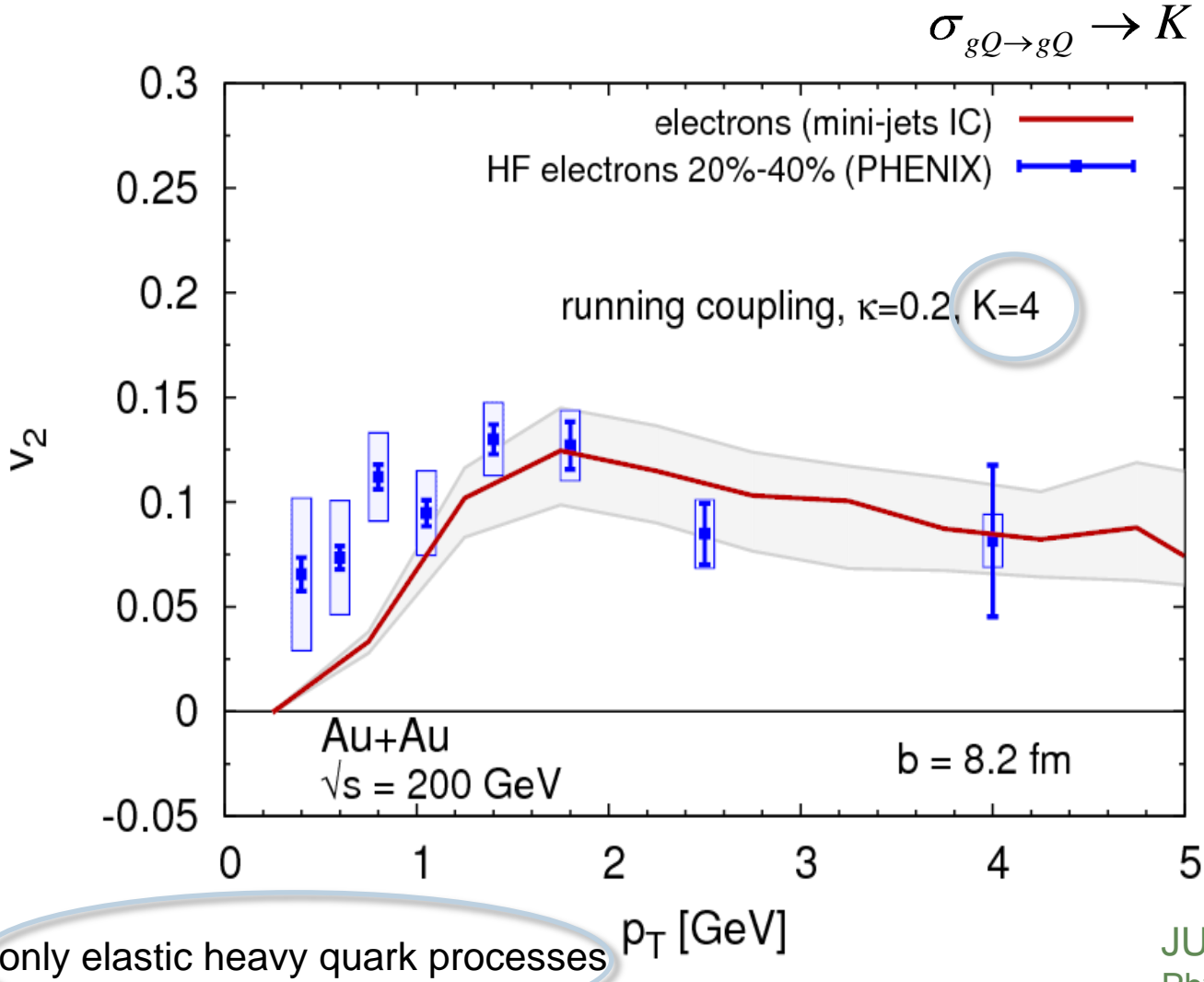
A. Peshier,
arXiv:0801.0595
[hep-ph]

P.B. Gossiaux,
J. Aichelin,
Phys.Rev.C78 (2008)



**Introduce a running coupling
constant for all channels**

Heavy quark elliptic flow v_2 at RHIC



RHIC

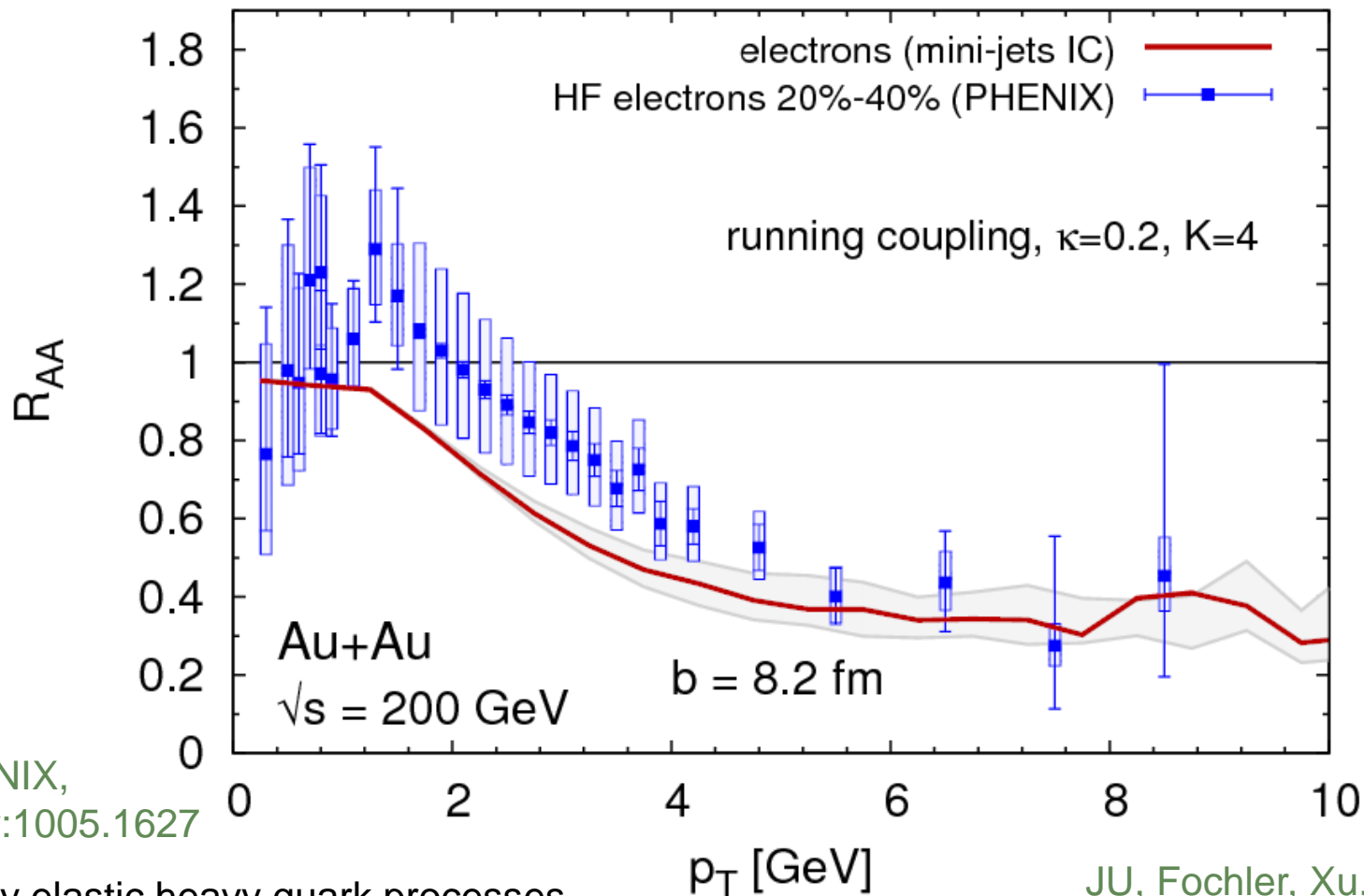
- What is missing:
- Radiative contributions
 - Quantum statistics

PHENIX,
arXiv:1005.1627

JU, Fochler, Xu, Greiner
Phys. Rev. C 84 (2011)

Heavy quark R_{AA} at RHIC

$$\sigma_{gQ \rightarrow gQ} \rightarrow K \sigma_{gQ \rightarrow gQ}$$



RHIC

PHENIX,
arXiv:1005.1627

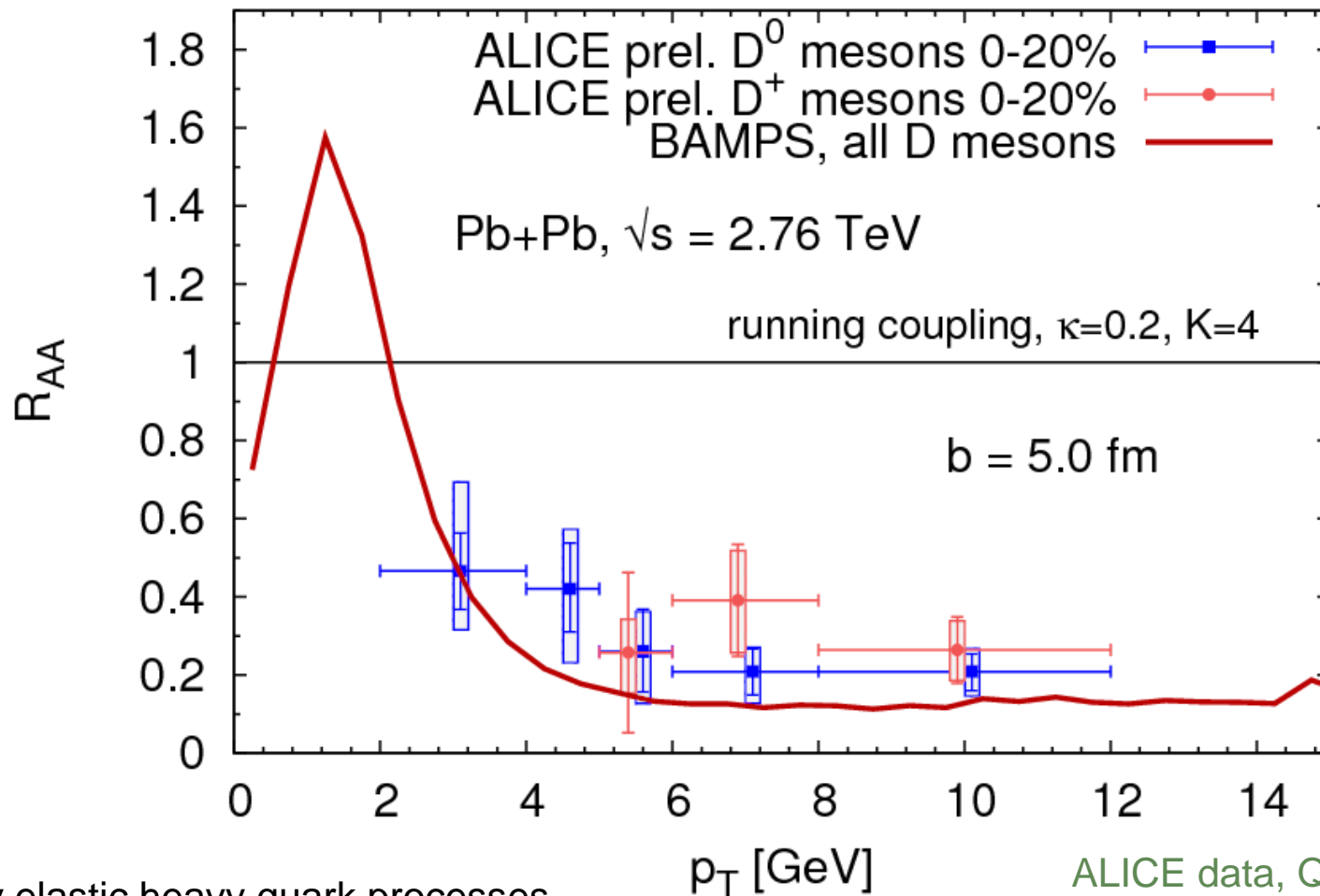
only elastic heavy quark processes

JU, Fochler, Xu, Greiner
Phys. Rev. C 84 (2011)

D meson R_{AA} at LHC

$$\sigma_{gQ \rightarrow gQ} \rightarrow K \sigma_{gQ \rightarrow gQ}$$

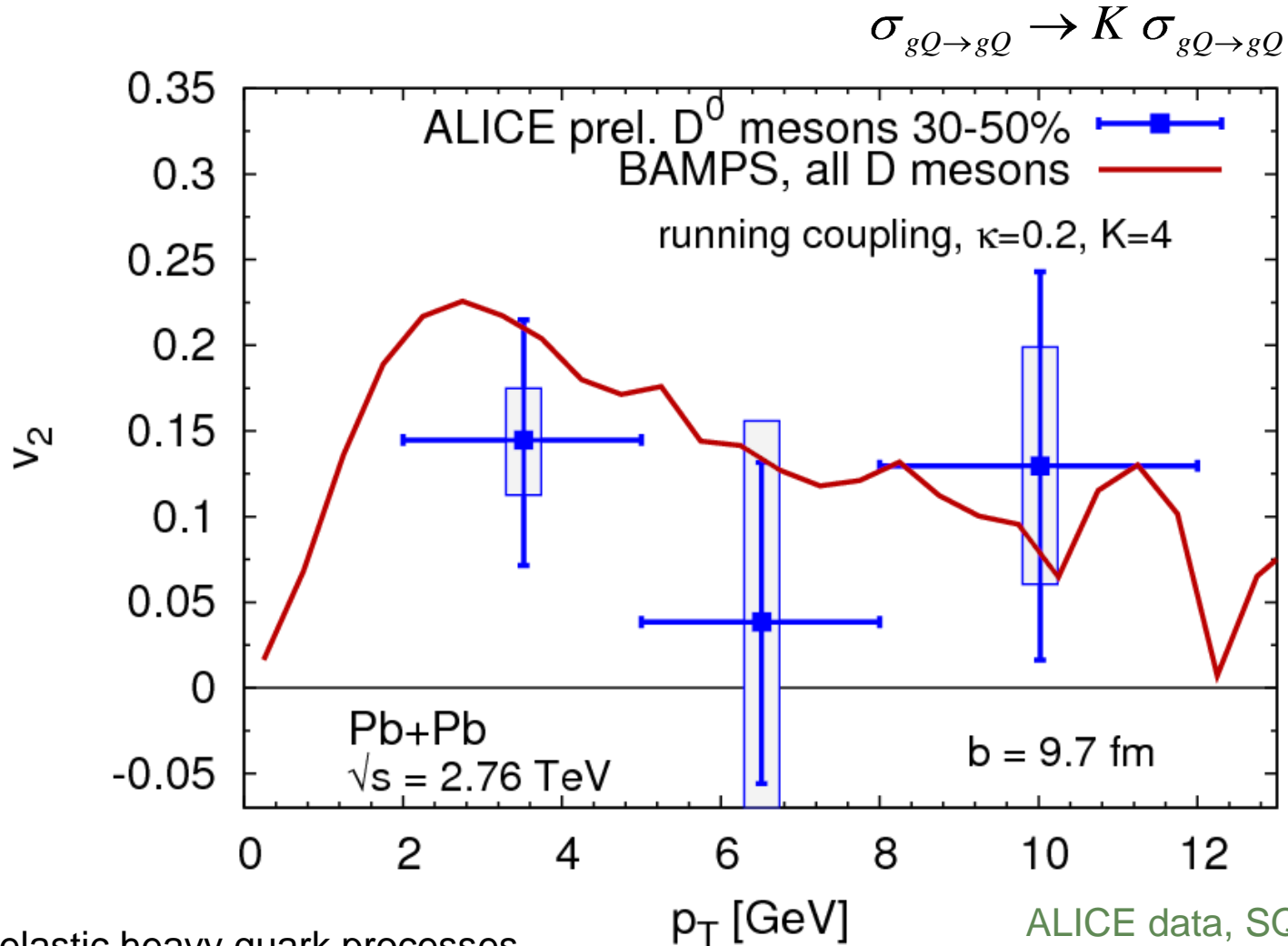
LHC



only elastic heavy quark processes

ALICE data, QM 2011,
arXiv:1106.5931

D meson v_2 at LHC



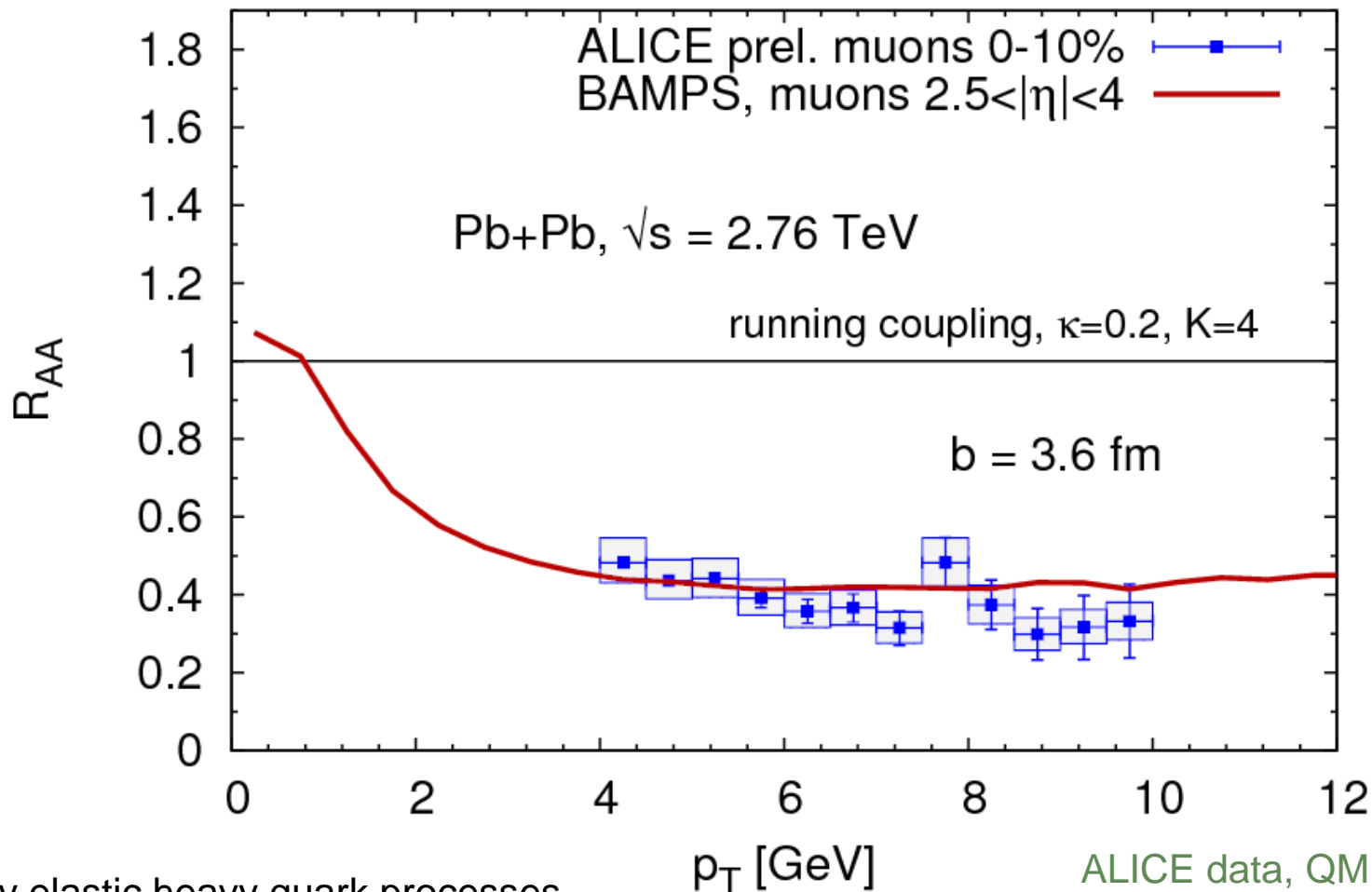
only elastic heavy quark processes

ALICE data, SQM 2011

Muon R_{AA} at forward rapidity at LHC

preliminary

$$\sigma_{gQ \rightarrow gQ} \rightarrow K \sigma_{gQ \rightarrow gQ}$$



LHC

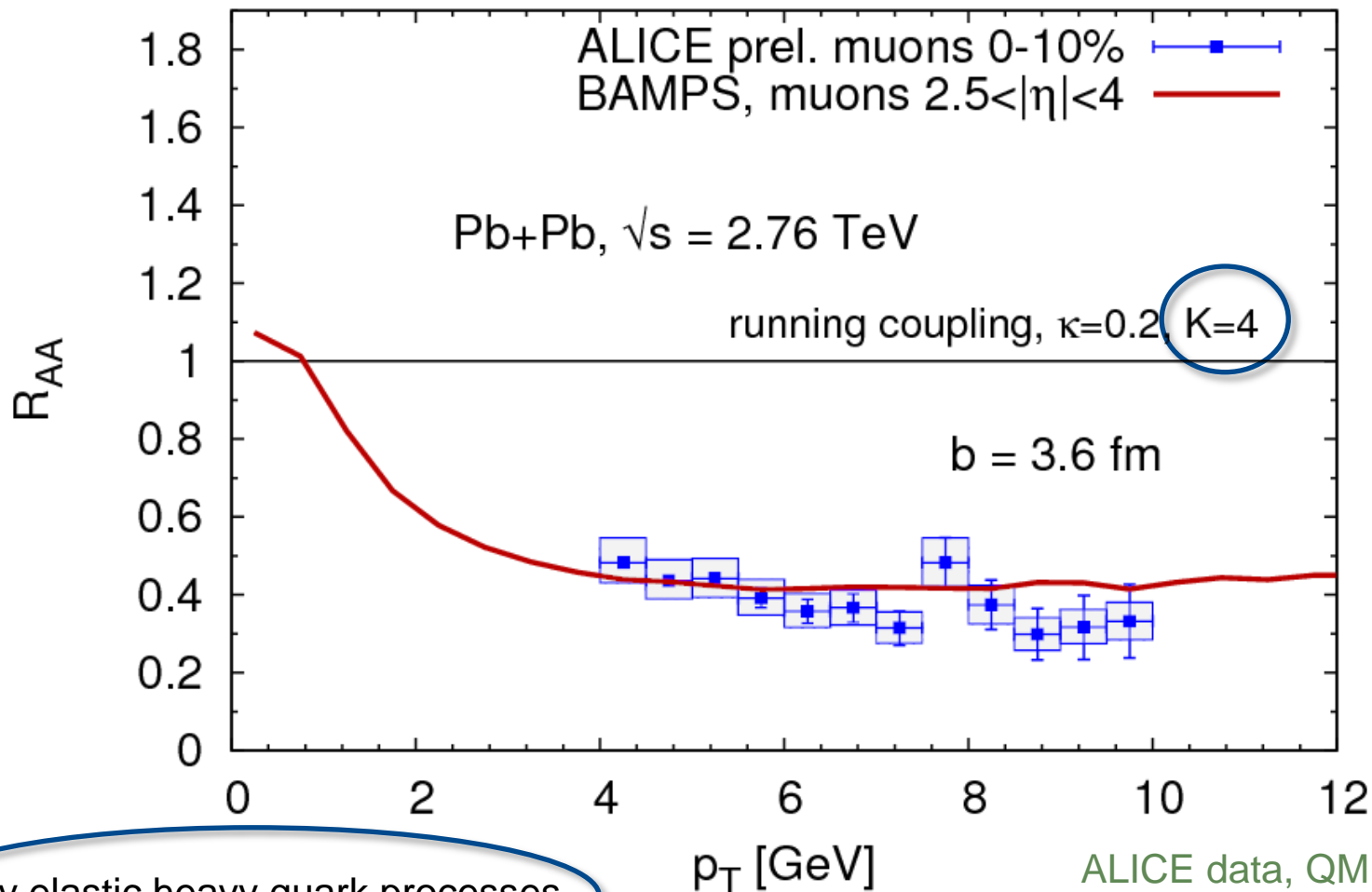
only elastic heavy quark processes

ALICE data, QM 2011,
arXiv:1106.4042

Muon R_{AA} at forward rapidity at LHC

preliminary

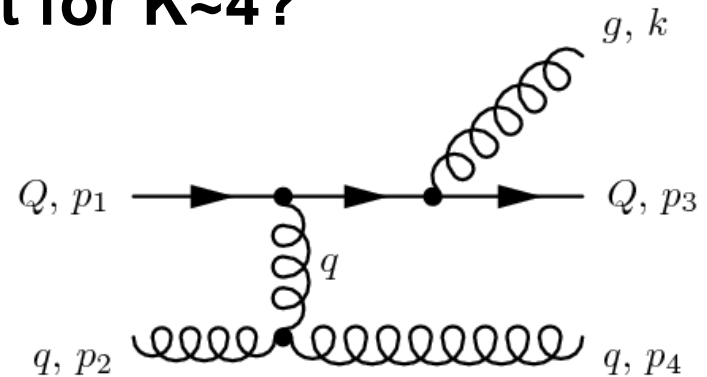
$$\sigma_{gQ \rightarrow gQ} \rightarrow K \sigma_{gQ \rightarrow gQ}$$



ALICE data, QM 2011,
arXiv:1106.4042

Can radiative processes account for $K \sim 4$?

$$g + Q \rightarrow g + Q + g$$

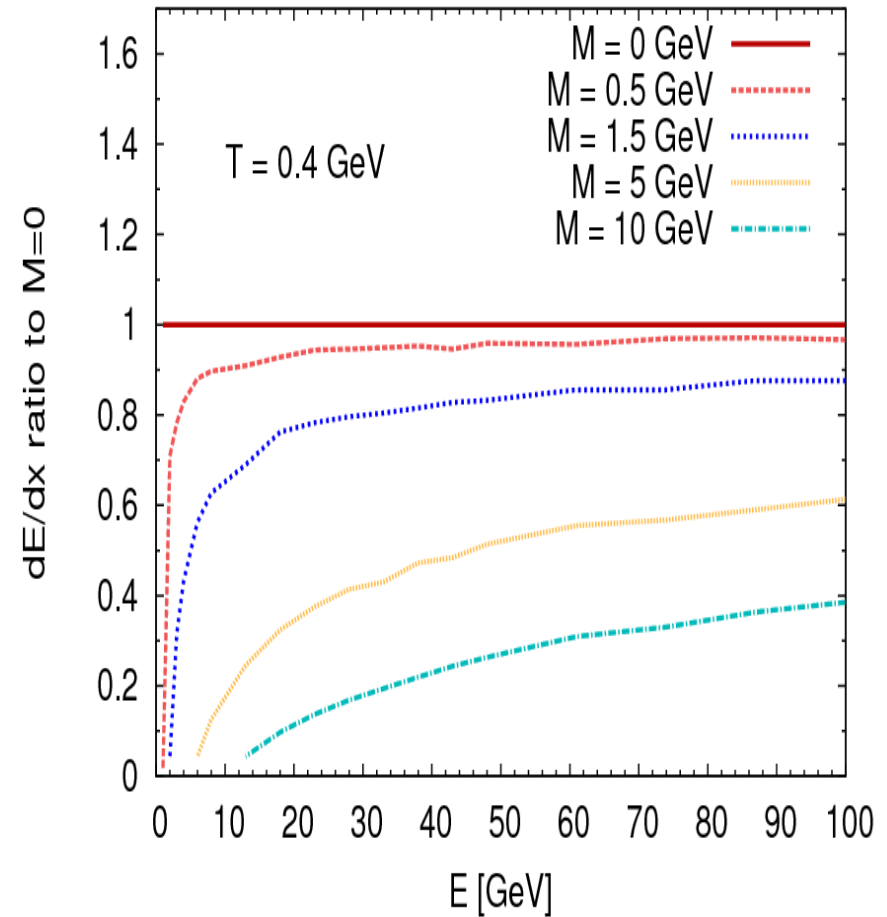
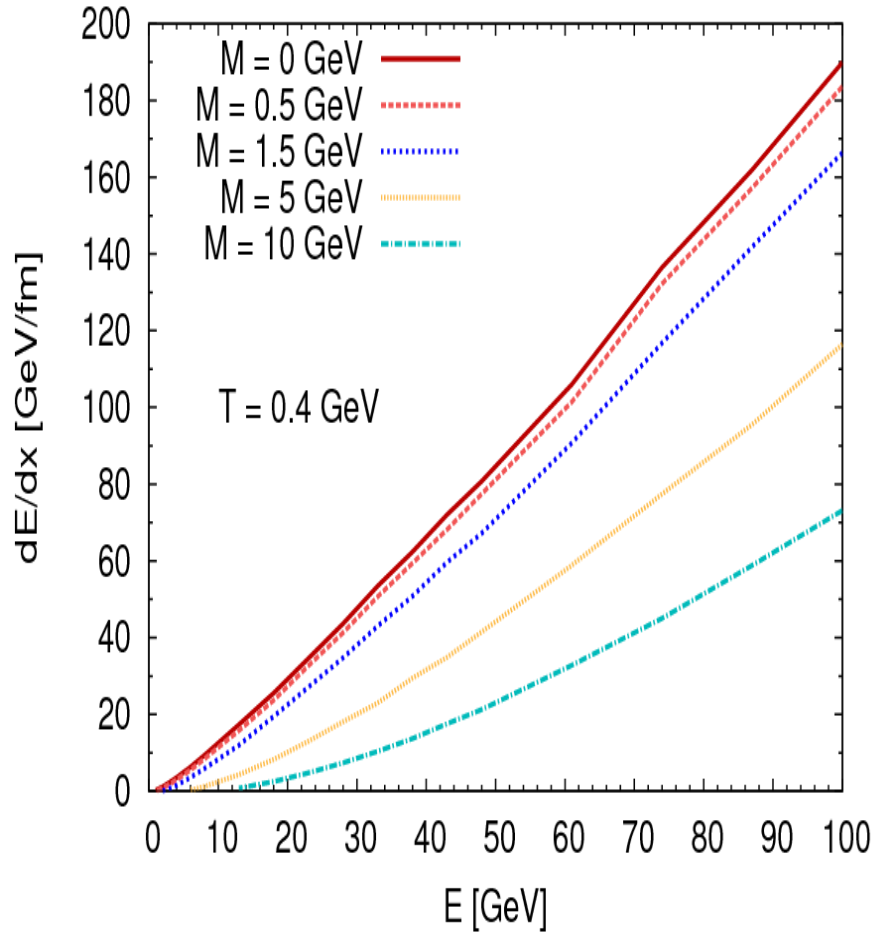


Gunion-Bertsch matrix element generalized to heavy quarks:

$$|\overline{\mathcal{M}}_{gQ \rightarrow gQg}|^2 = 12g^2 \left| \overline{\mathcal{M}}_0^{gQ} \right|^2 \left[\frac{\mathbf{k}_\perp}{k_\perp^2 + x^2 M^2} + \frac{\mathbf{q}_\perp - \mathbf{k}_\perp}{(\mathbf{q}_\perp - \mathbf{k}_\perp)^2 + x^2 M^2} \right]^2$$

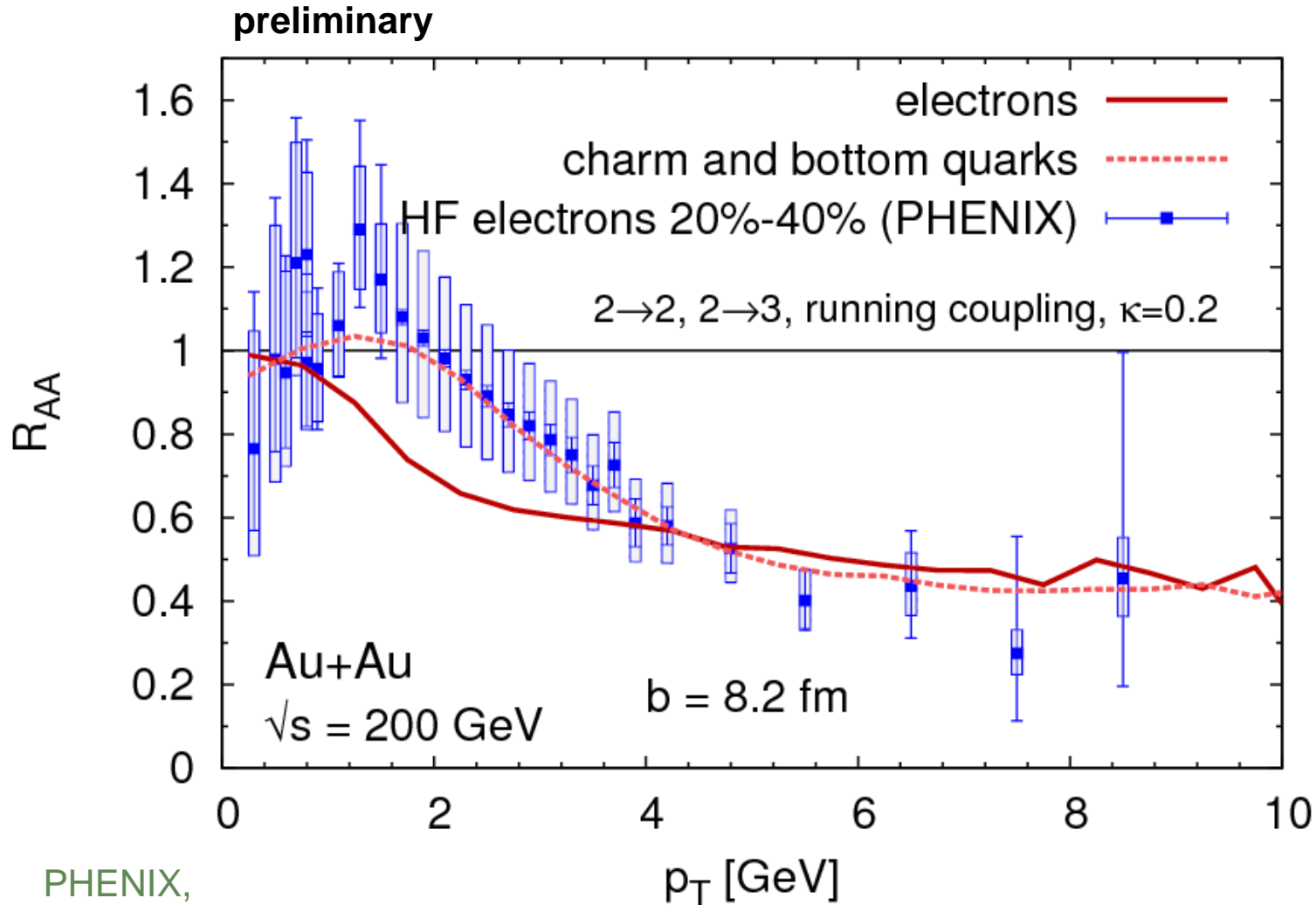
In accordance to scalar QCD result from
Gossiaux, Aichelin, Gousset, Guiho, J.Phys.G37 (2010)

Energy loss in static medium



Fixed coupling, without LPM effect

Heavy quark R_{AA} at RHIC with 2- \rightarrow 3



RHIC

PHENIX,
arXiv:1005.1627

Full space-time evolution of QGP with charm and bottom quarks

- Running coupling and improved Debye screening yield results that can explain experimental v_2 and R_{AA} at RHIC and LHC if $K=4$ is introduced
- Importance of $2 \rightarrow 3$ processes estimated in energy loss calculations in static medium
- Preliminary results with $2 \rightarrow 3$ in full cascade are promising

Further details on [arXiv:1104.2295](https://arxiv.org/abs/1104.2295) and [1104.2437](https://arxiv.org/abs/1104.2437)

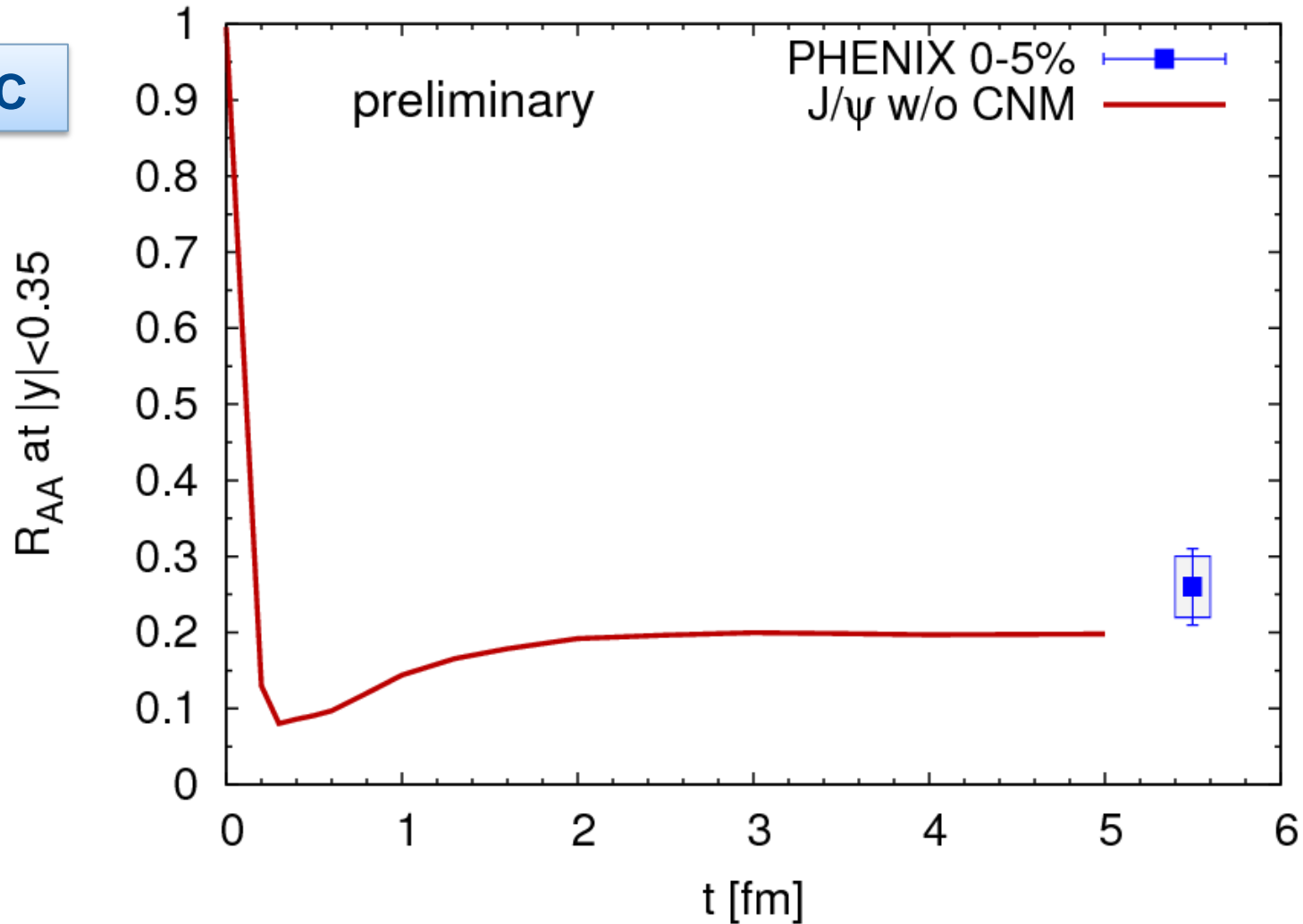
Future tasks:

- Further study of radiative heavy quark scattering in full cascade
- Light quark interactions with heavy quarks
- Further study of J/ψ suppression, also at LHC

Thank you for your attention.

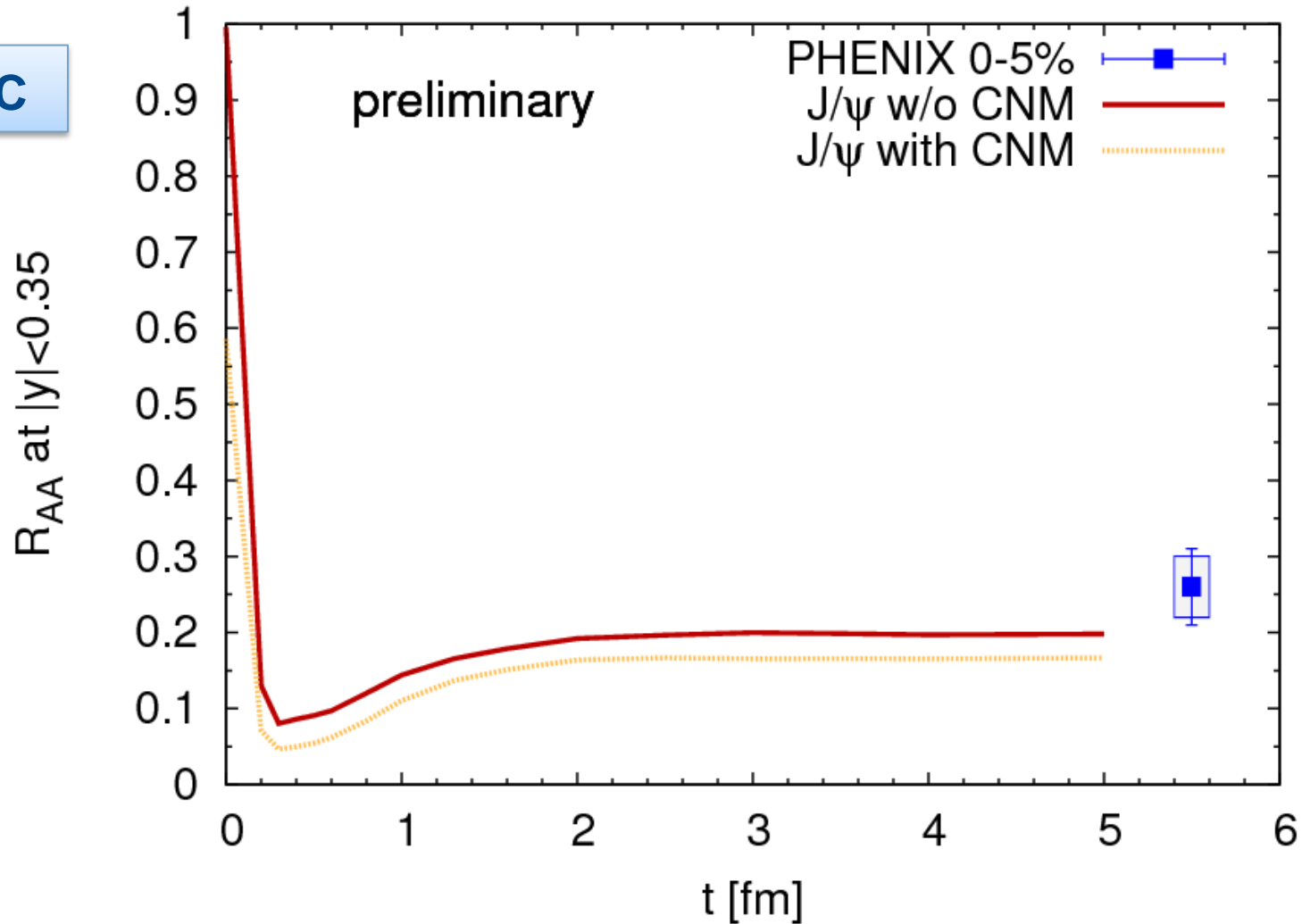
J/ ψ production

RHIC



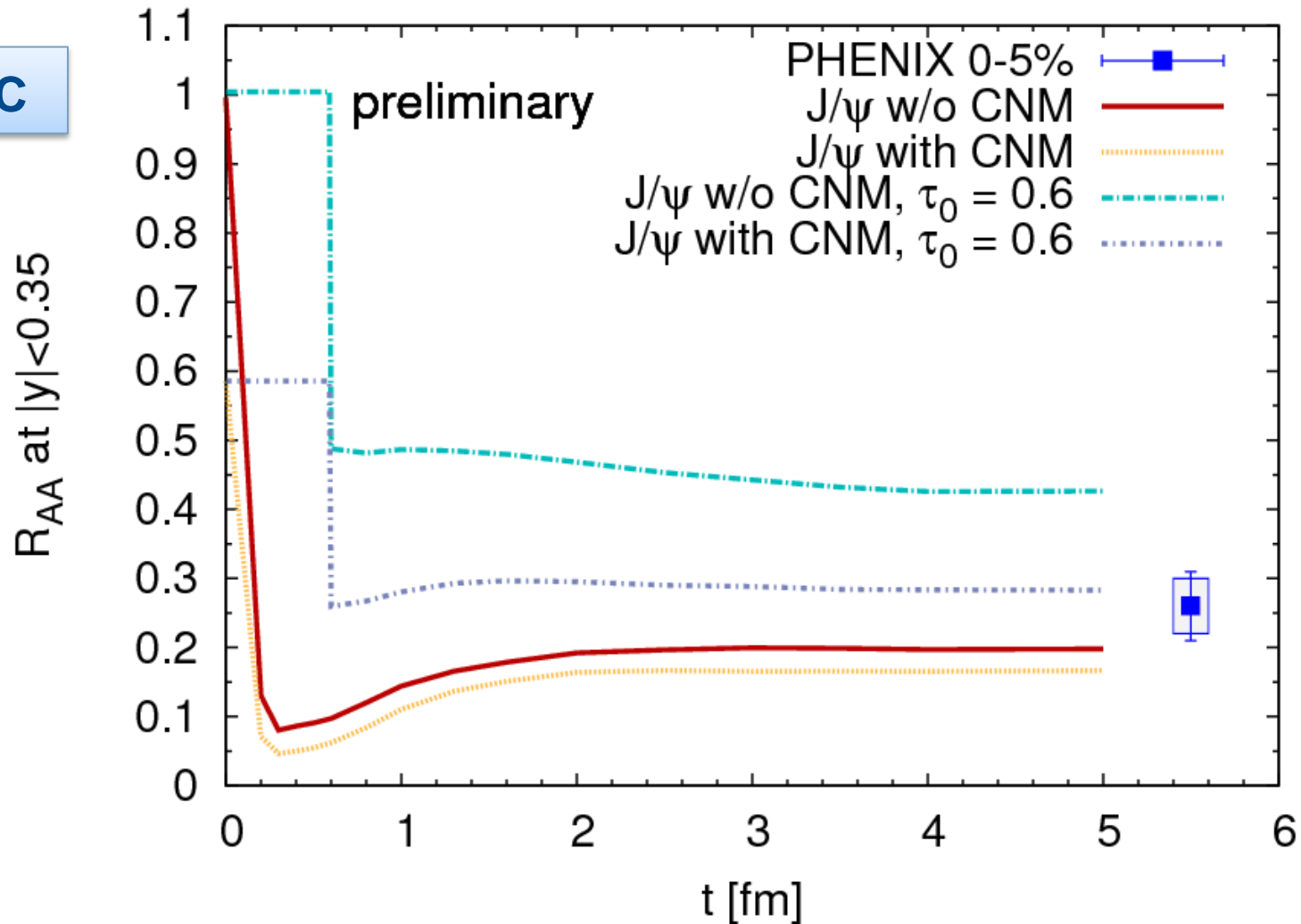
J/ ψ production

RHIC



J/ ψ production

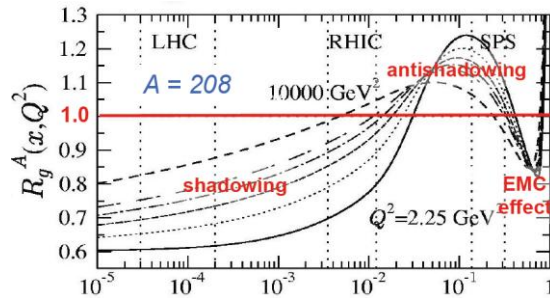
RHIC



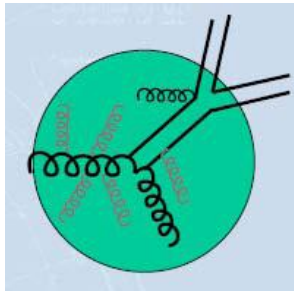
J/ψ suppression

Cold nuclear matter effects

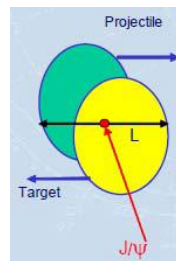
Shadowing



Cronin effect



Nuclear absorption



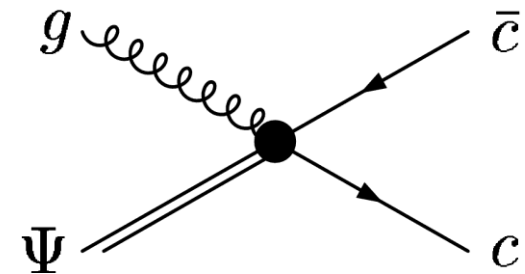
Hot nuclear matter effects

J/psi dissociation

- If $T > T_d = 2 T_c = 330 \text{ MeV}$
- Via $J/\psi + g \rightarrow c + \bar{c}$

J/psi regeneration

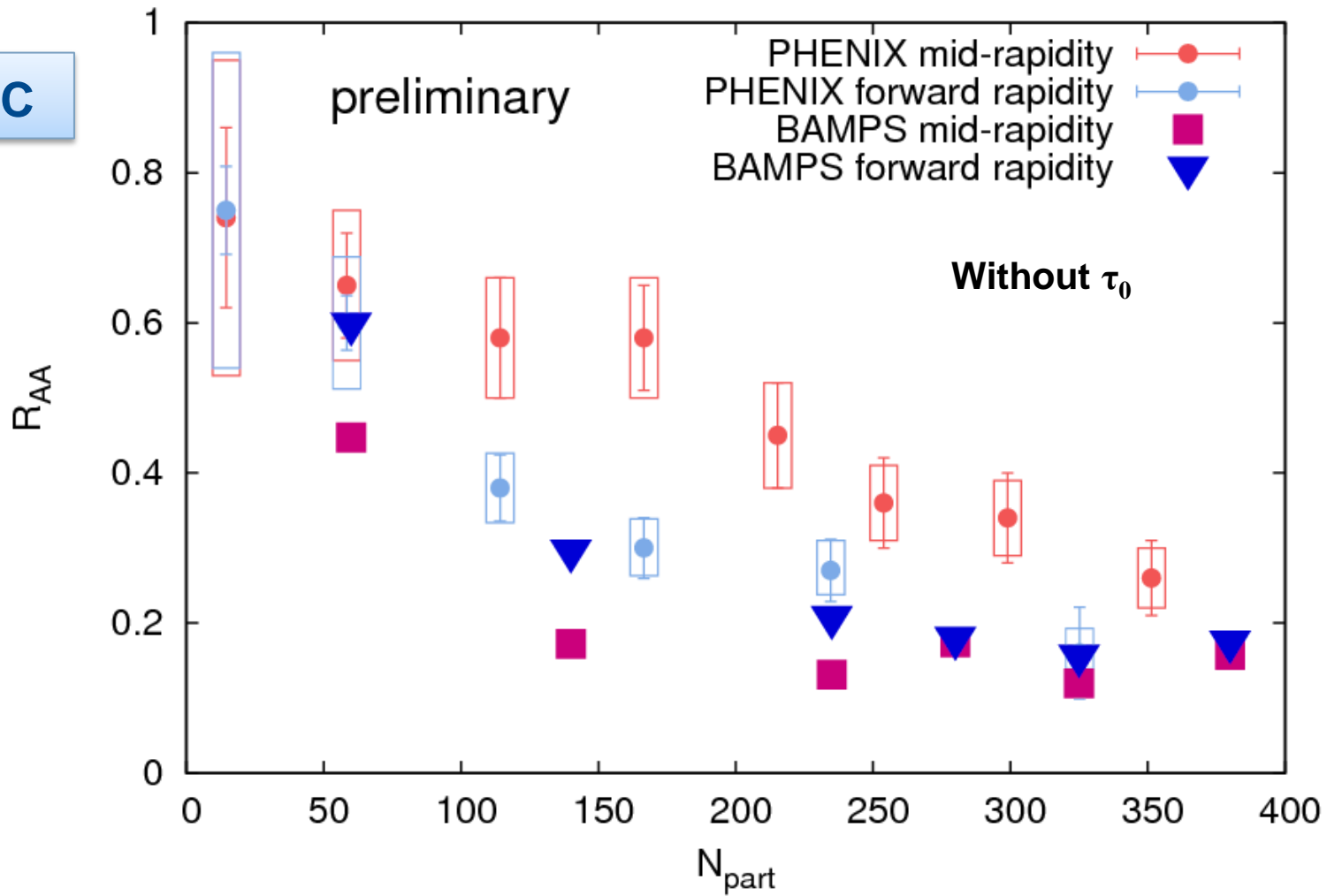
via $c + \bar{c} \rightarrow J/\psi + g$



Bhanot+Peskin 79

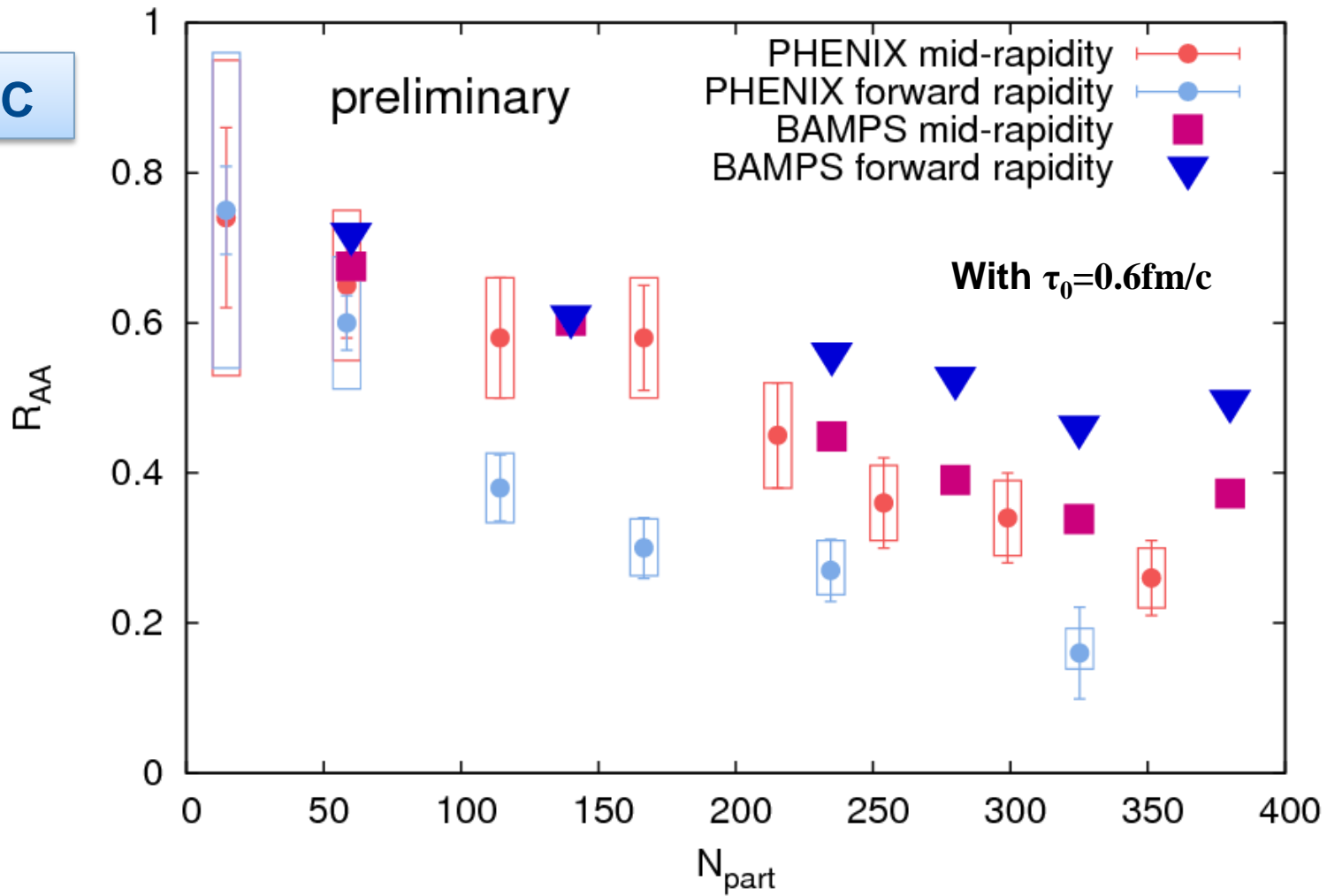
J/ ψ production

RHIC

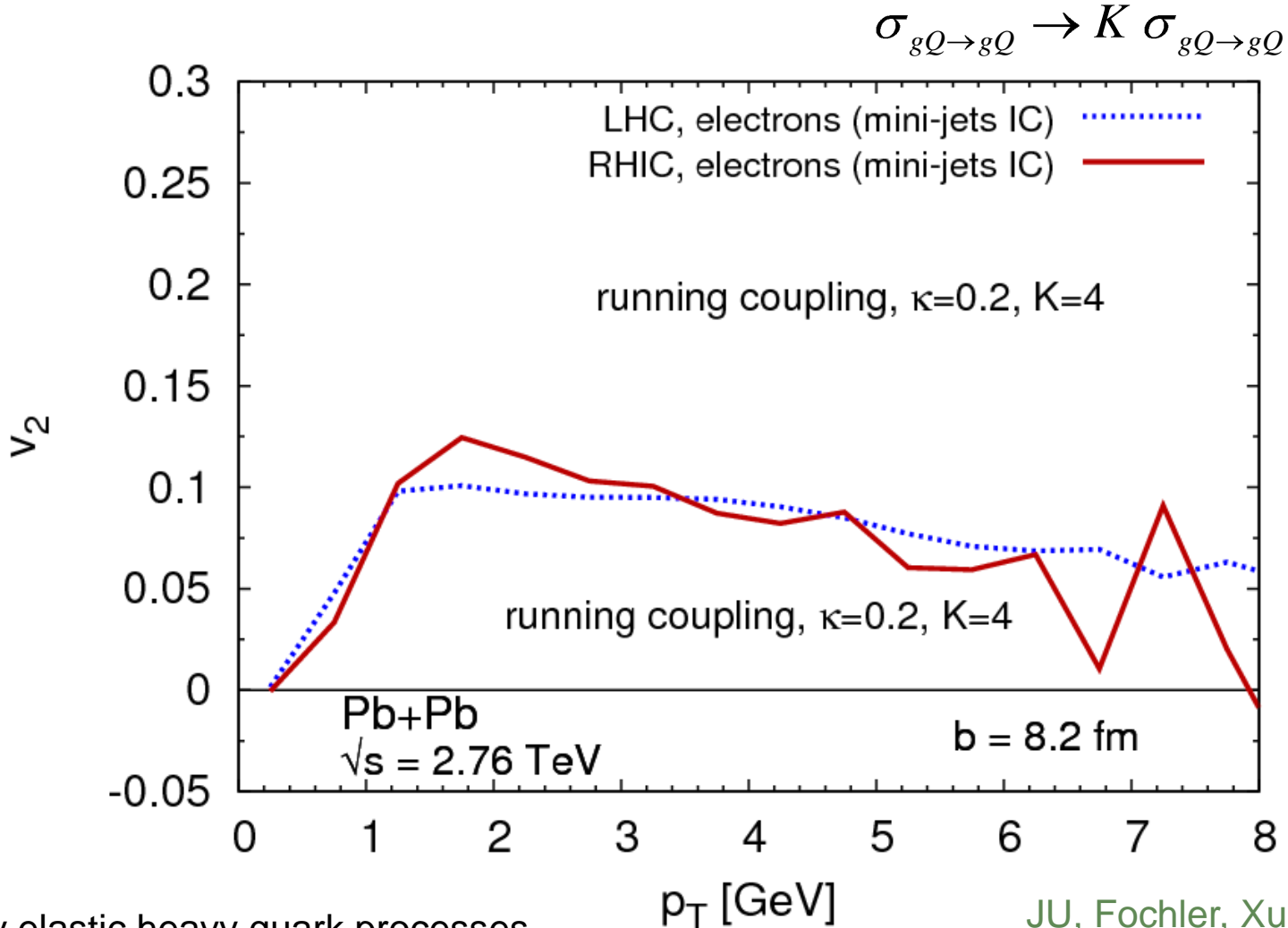


J/ ψ production

RHIC



Heavy quark elliptic flow v_2 at LHC



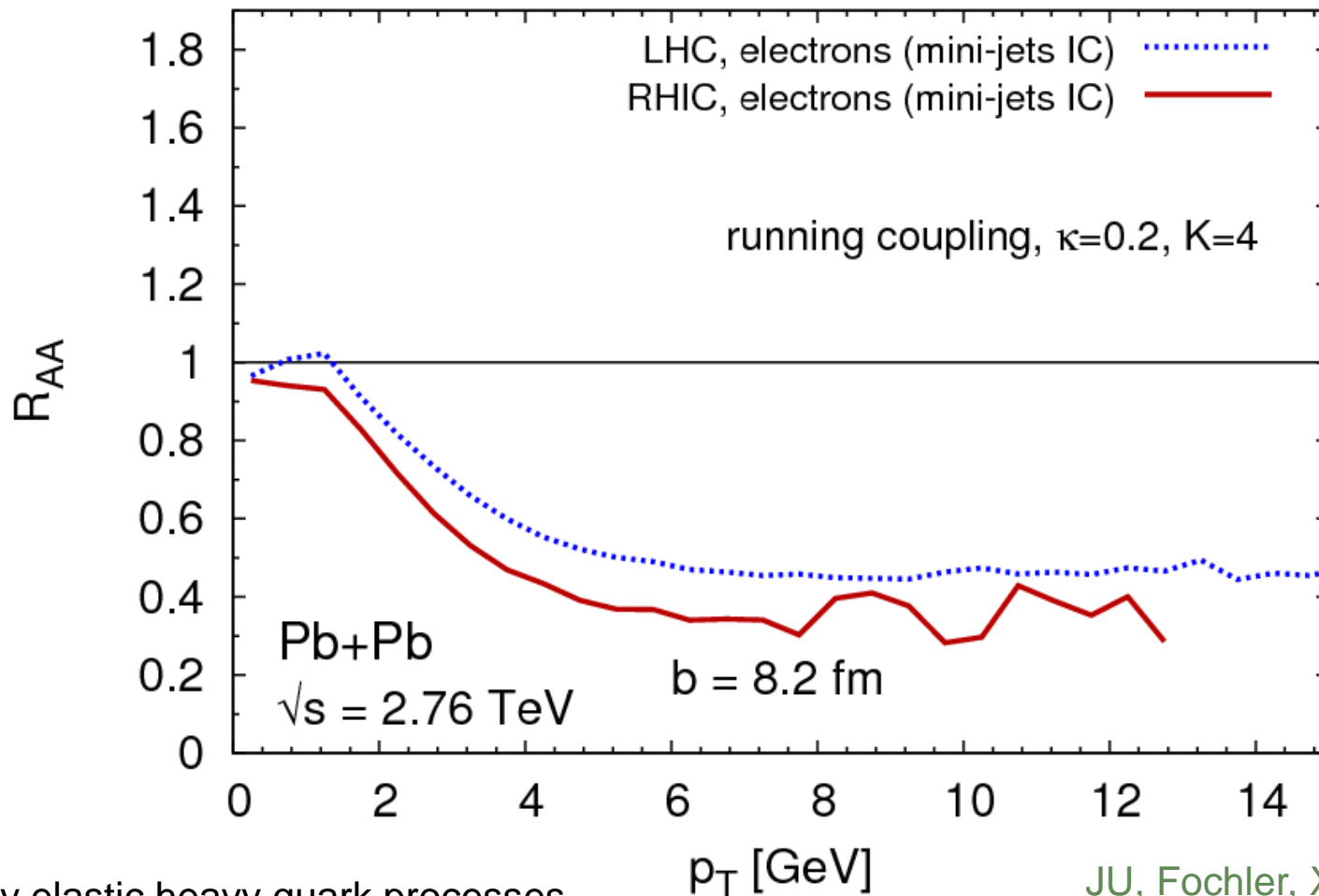
LHC

only elastic heavy quark processes

JU, Fochler, Xu, Greiner
Phys. Rev. C 84 (2011)

Heavy quark R_{AA} at LHC

$$\sigma_{gQ \rightarrow gQ} \rightarrow K \sigma_{gQ \rightarrow gQ}$$



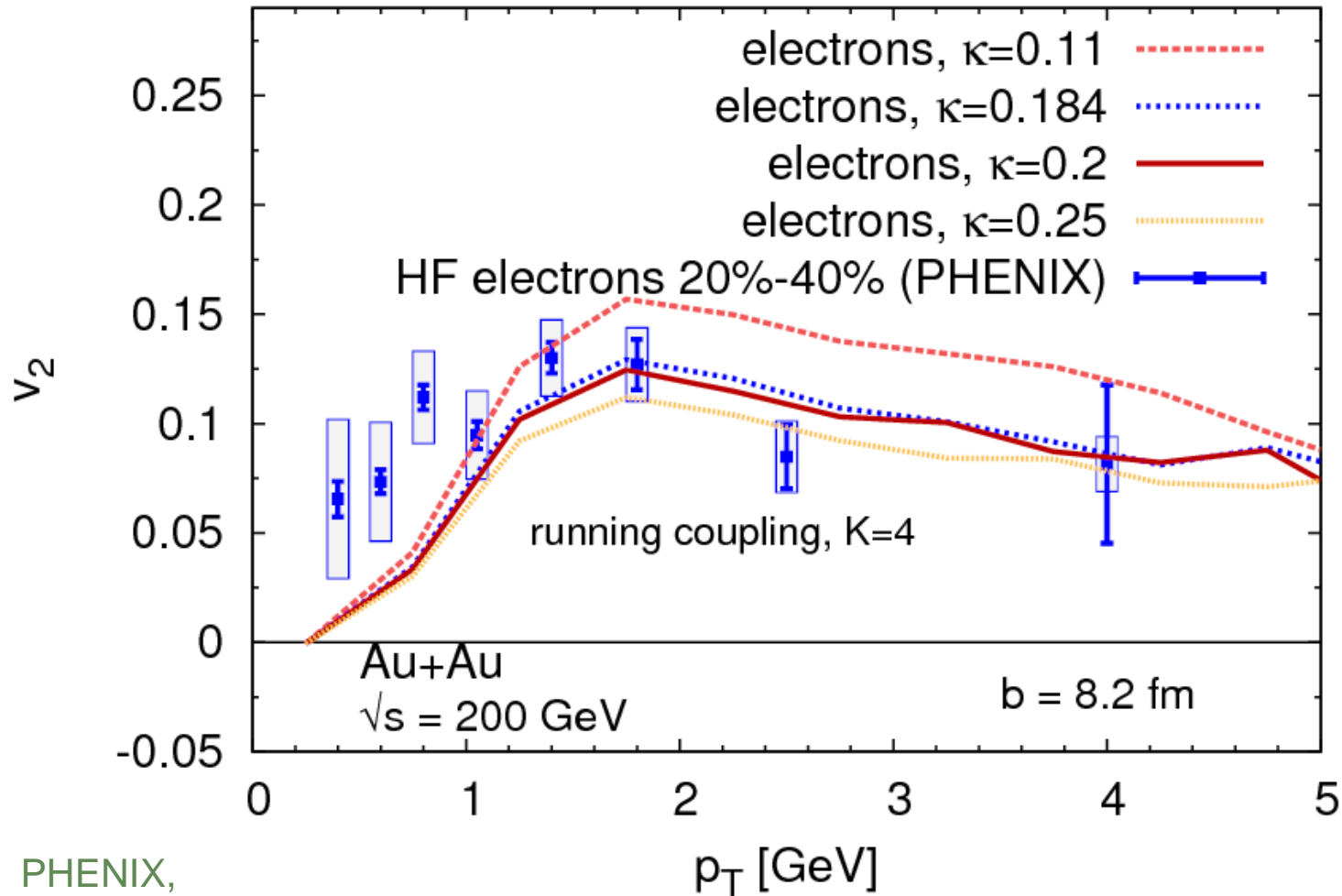
LHC

only elastic heavy quark processes

JU, Fochler, Xu, Greiner
Phys. Rev. C 84 (2011)

Heavy quark elliptic flow v_2 at RHIC

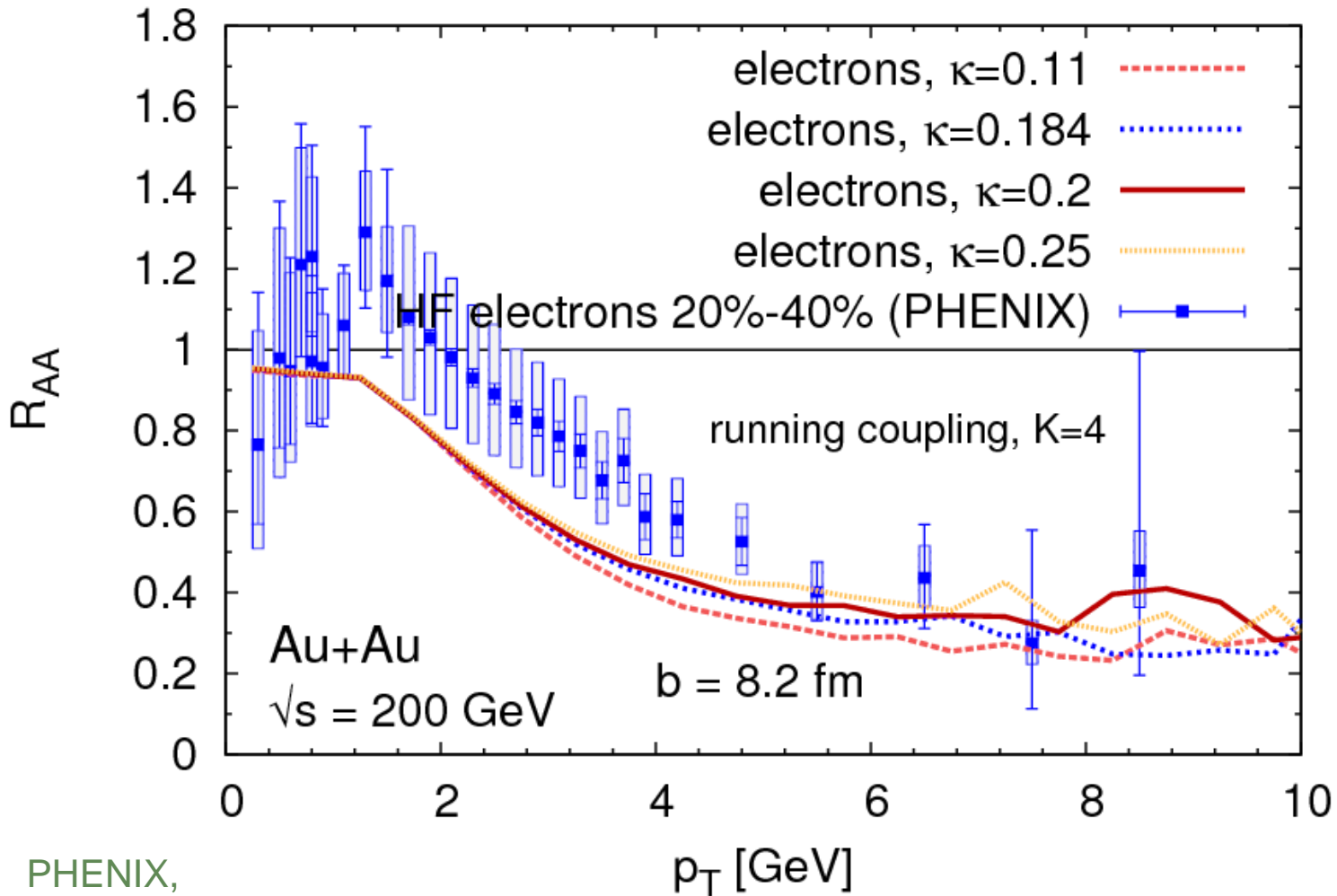
RHIC



PHENIX,
arXiv:1005.1627

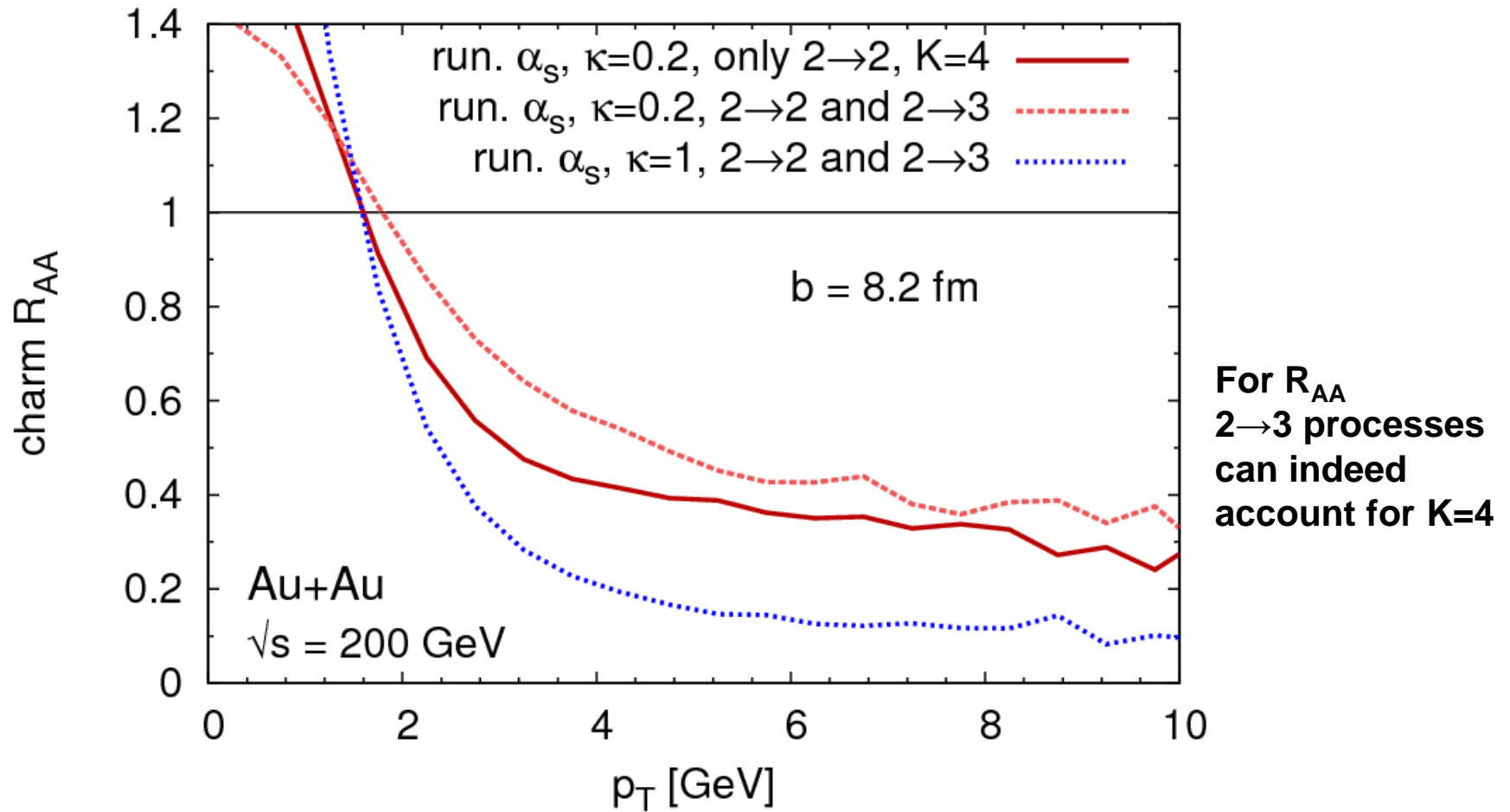
Heavy quark elliptic flow v_2 at RHIC

RHIC



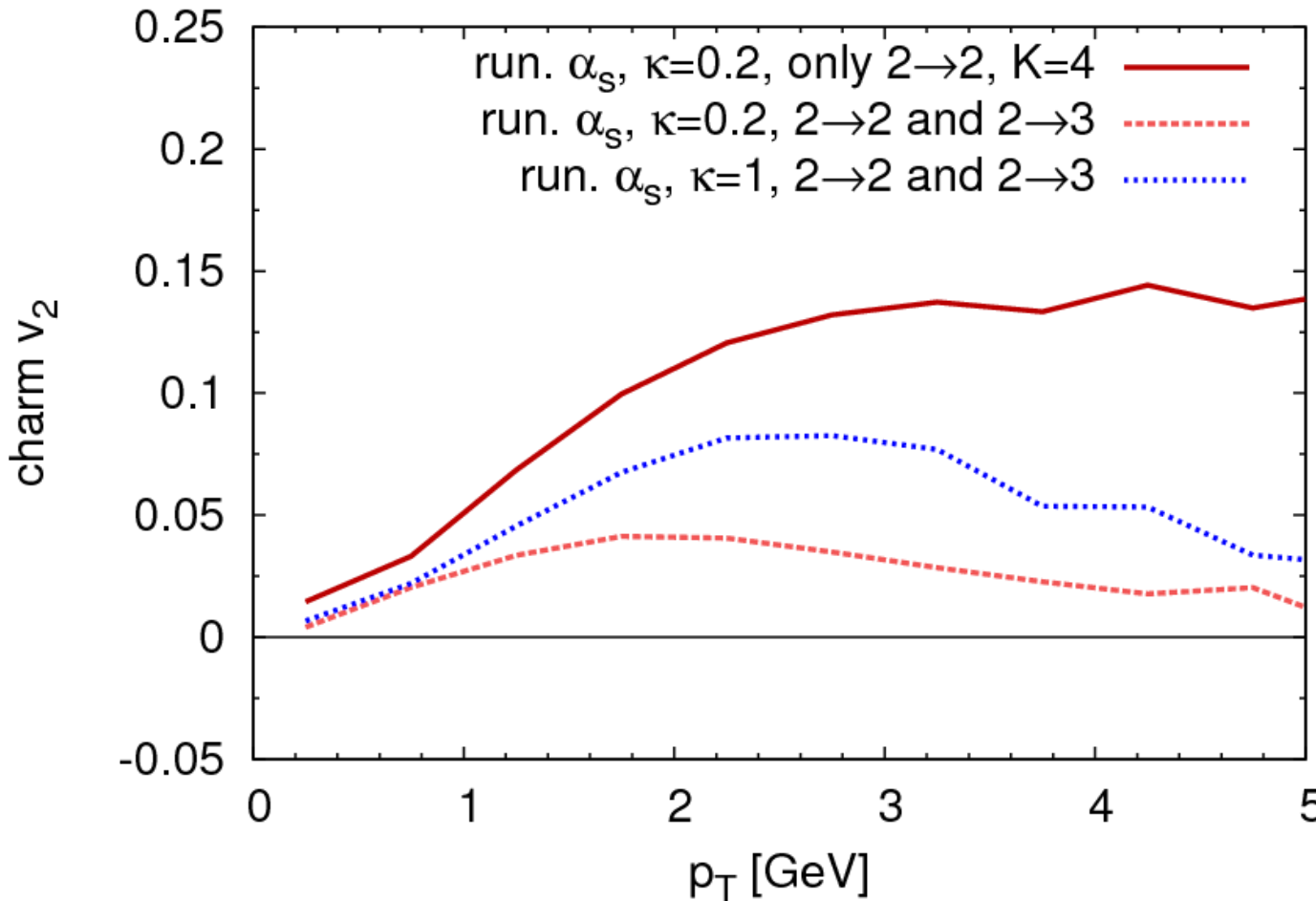
PHENIX,
arXiv:1005.1627

Charm R_{AA} at RHIC



Only charm quarks (no heavy flavor electrons!) for better comparison

Charm elliptic flow v_2 at RHIC



For v_2
 $2 \rightarrow 3$ processes
cannot explain
missing factor
 $K=4$

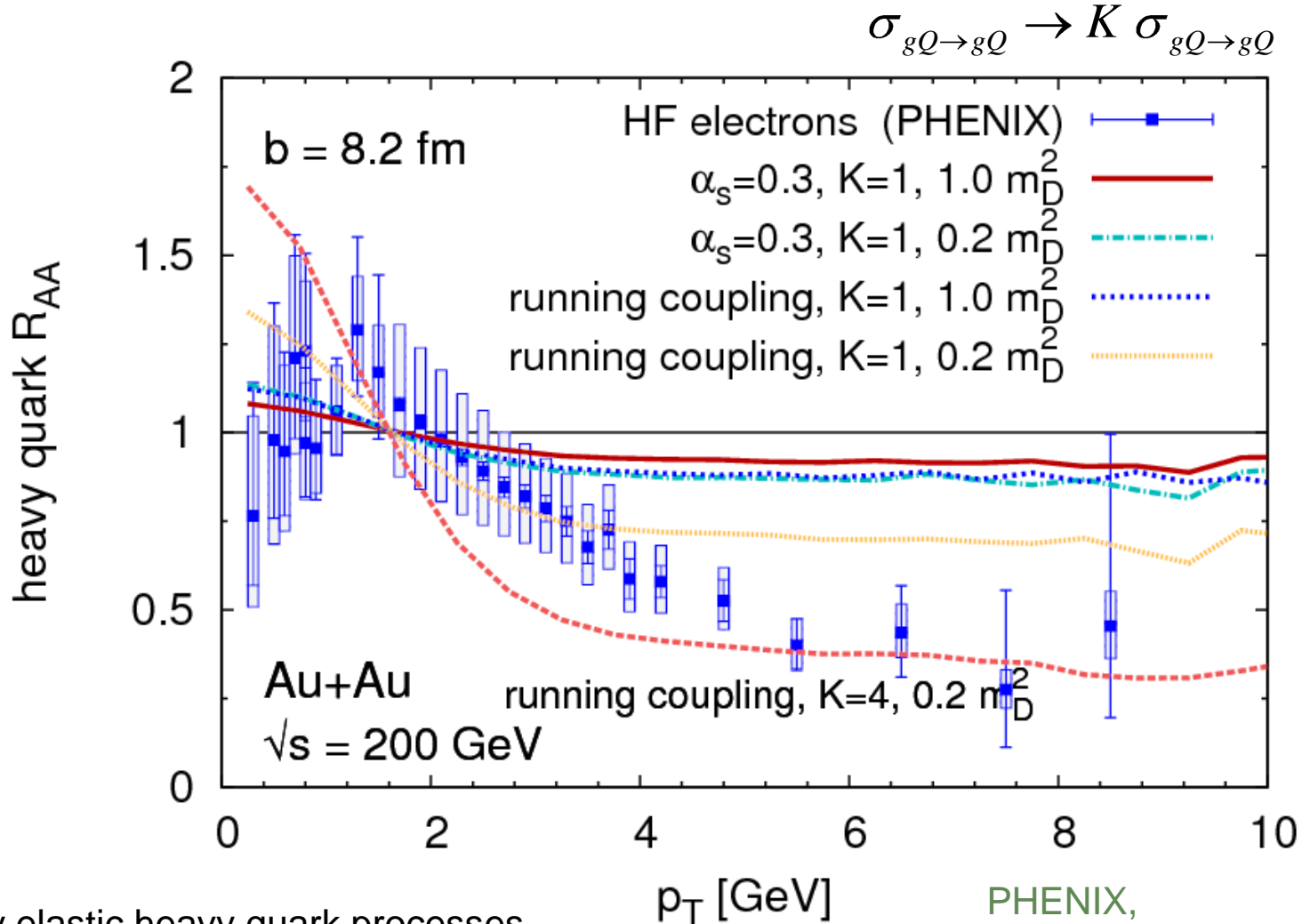
Different impact
of $2 \rightarrow 3$
processes on v_2
and R_{AA}

Reason:
LPM effect

$\kappa=1$ is even
better since $2 \rightarrow 3$
processes more
important due to
LPM effect

Only charm quarks (no heavy flavor electrons!) for better comparison

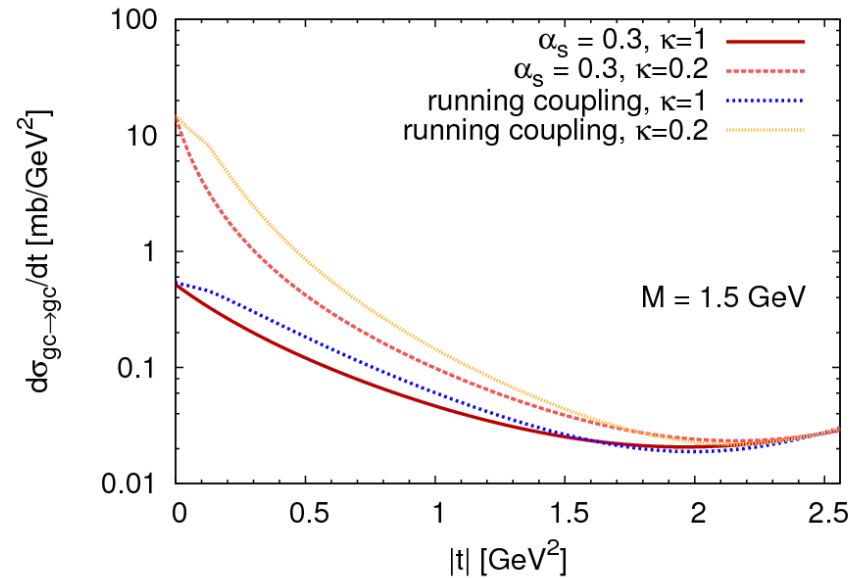
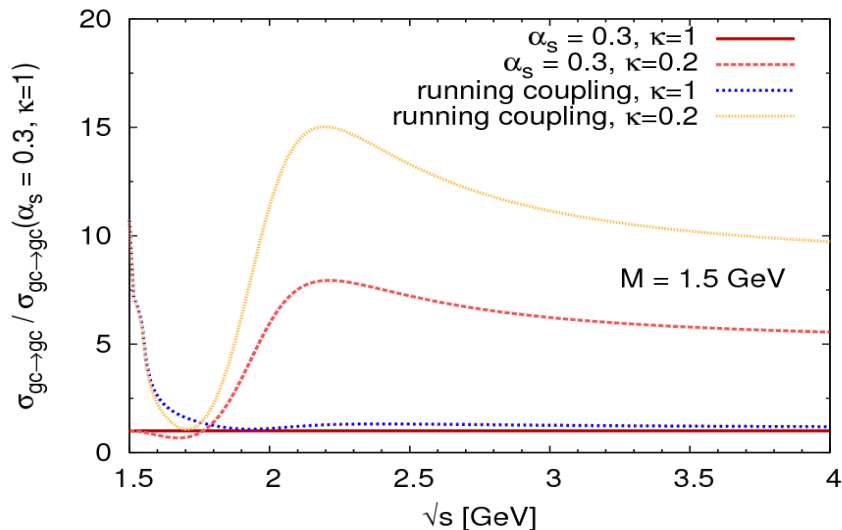
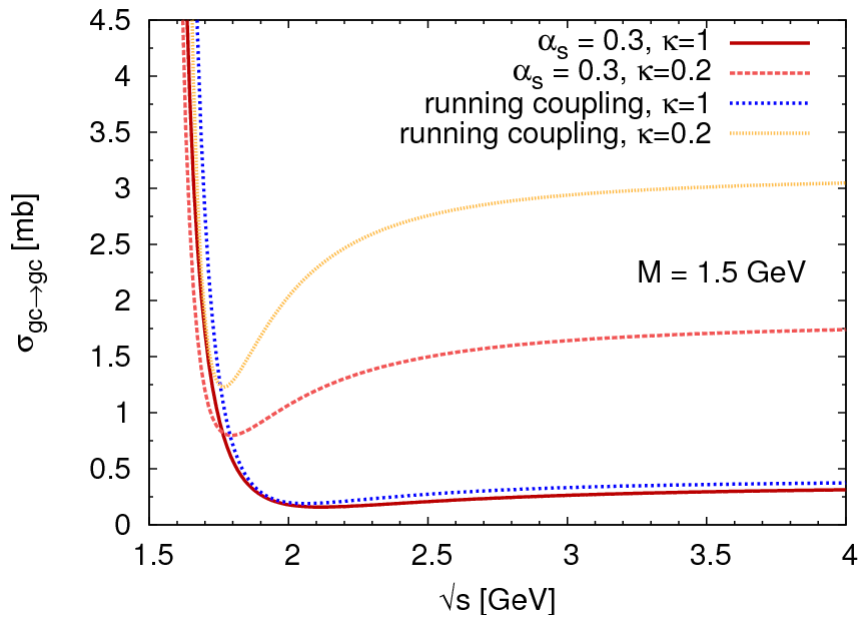
Heavy quark R_{AA} at RHIC



only elastic heavy quark processes

PHENIX,
arXiv:1005.1627

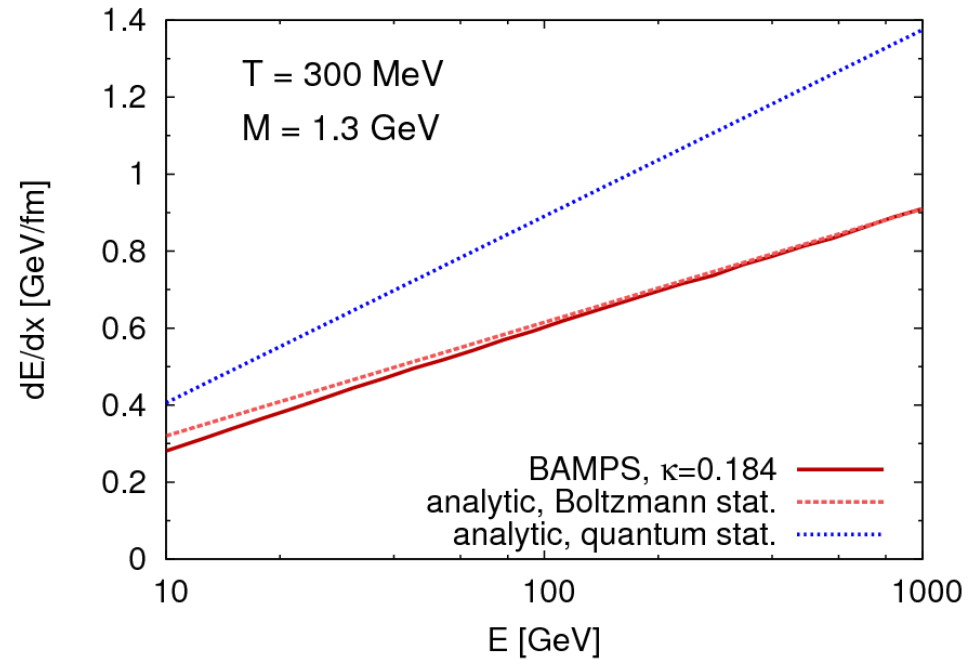
Heavy quark scattering cross section



Heavy quark scattering

Compare to analytic formula

$$\frac{dE}{dx} = \frac{8\alpha_s^2 T^2}{\pi} \left[\left(1 + \frac{n_f}{3}\right) \ln \frac{ET}{m_D^2} + \frac{2}{9} \ln \frac{ET}{M^2} \right. \\ \left. + \left(\ln 2 - \frac{1}{4} - \frac{\gamma}{3}\right) n_f \right. \\ \left. + \frac{31}{9} \ln 2 - \frac{101}{108} - \frac{11\gamma}{9} \right]$$



Fragmentation and Decay

- Peterson fragmentation

Peterson et al., Phys. Rev. D27 (1983)

$$D_{H/Q}(z) = \frac{N}{z \left(1 - \frac{1}{z} - \frac{\epsilon_Q}{1-z}\right)^2}$$

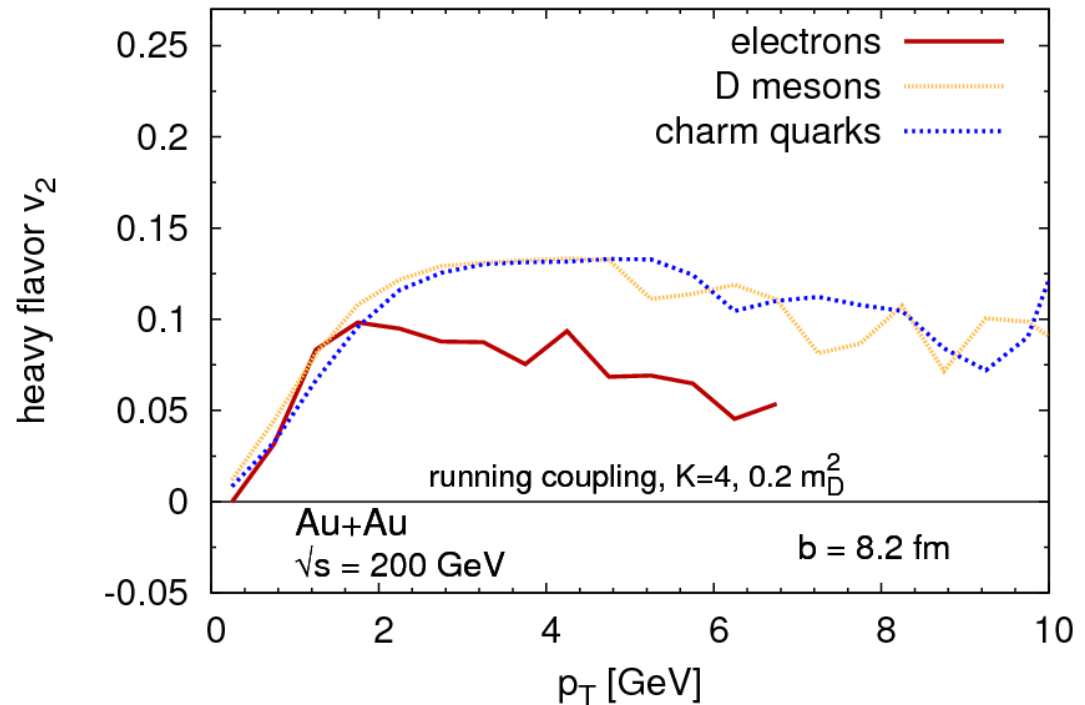
$$z = \frac{|\vec{p}_H|}{|\vec{p}_Q|}$$

$$\epsilon_c = 0.05$$

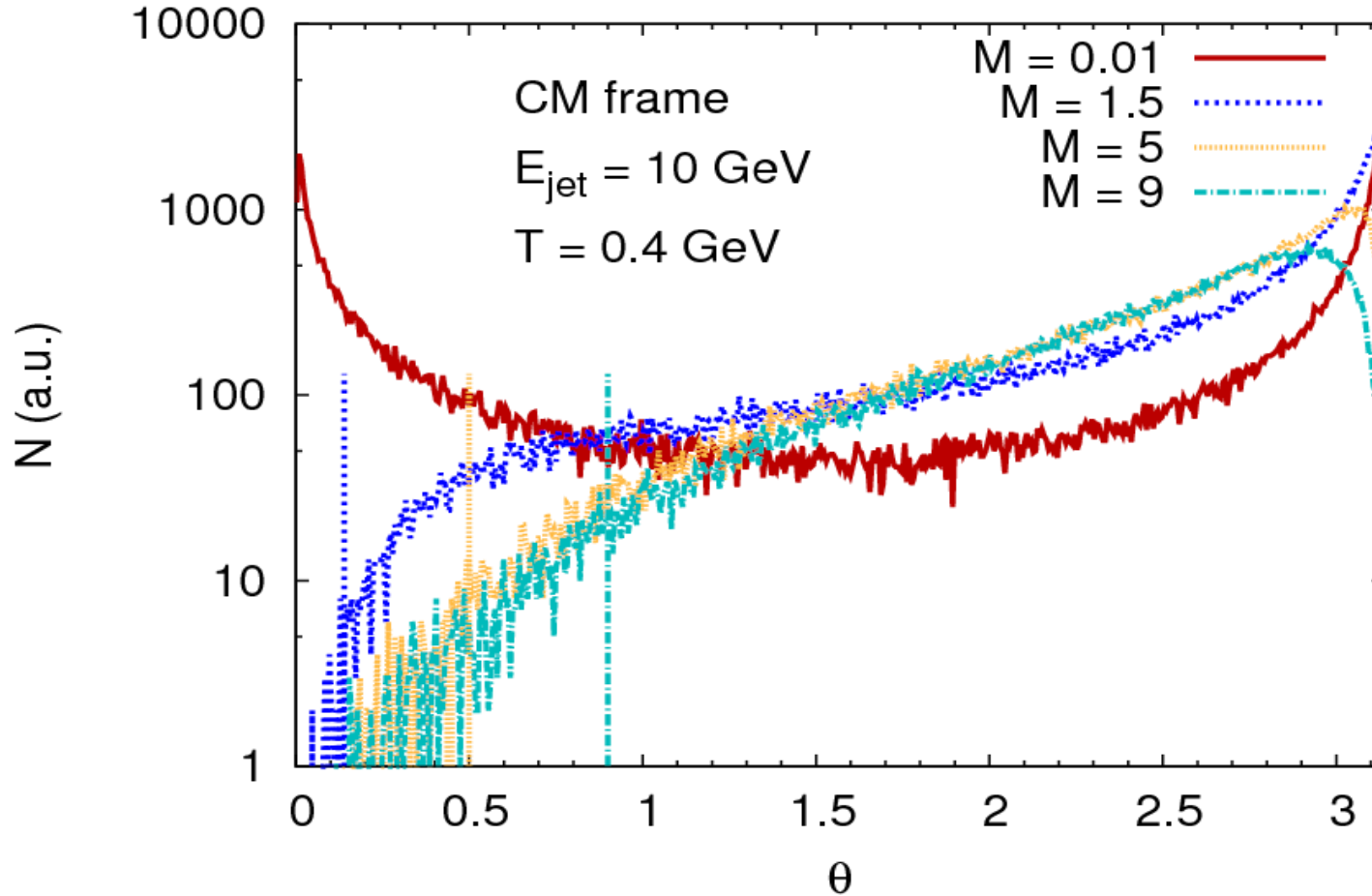
$$\epsilon_b = 0.005$$

- Decay to electrons with PYTHIA

Impact of hadronization and decay small

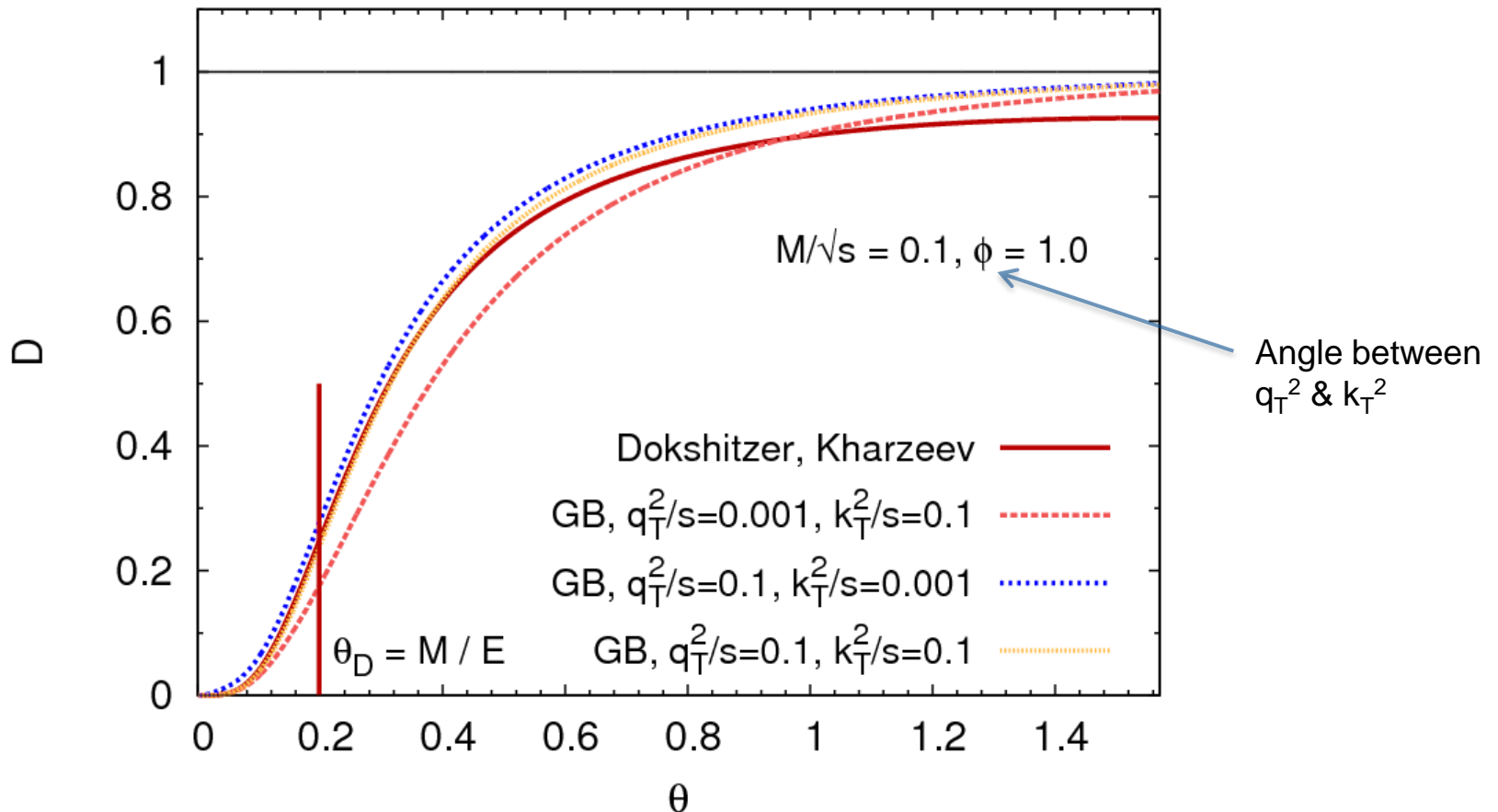


Θ dependence in static medium

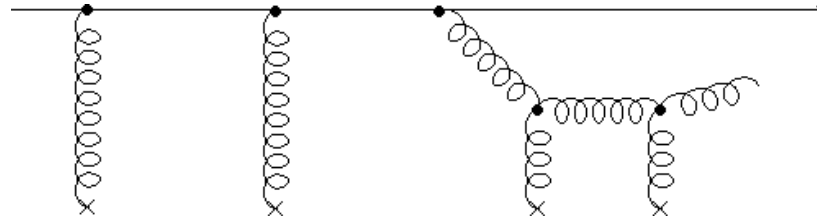


Dead cone effect

$$|\overline{\mathcal{M}}_{gQ \rightarrow gQg}|^2 = 12g^2 |\overline{\mathcal{M}}_0^{gQ}|^2 \left[\frac{\mathbf{k}_\perp}{k_\perp^2 + x^2 M^2} + \frac{\mathbf{q}_\perp - \mathbf{k}_\perp}{(\mathbf{q}_\perp - \mathbf{k}_\perp)^2 + x^2 M^2} \right]^2$$



LPM effect

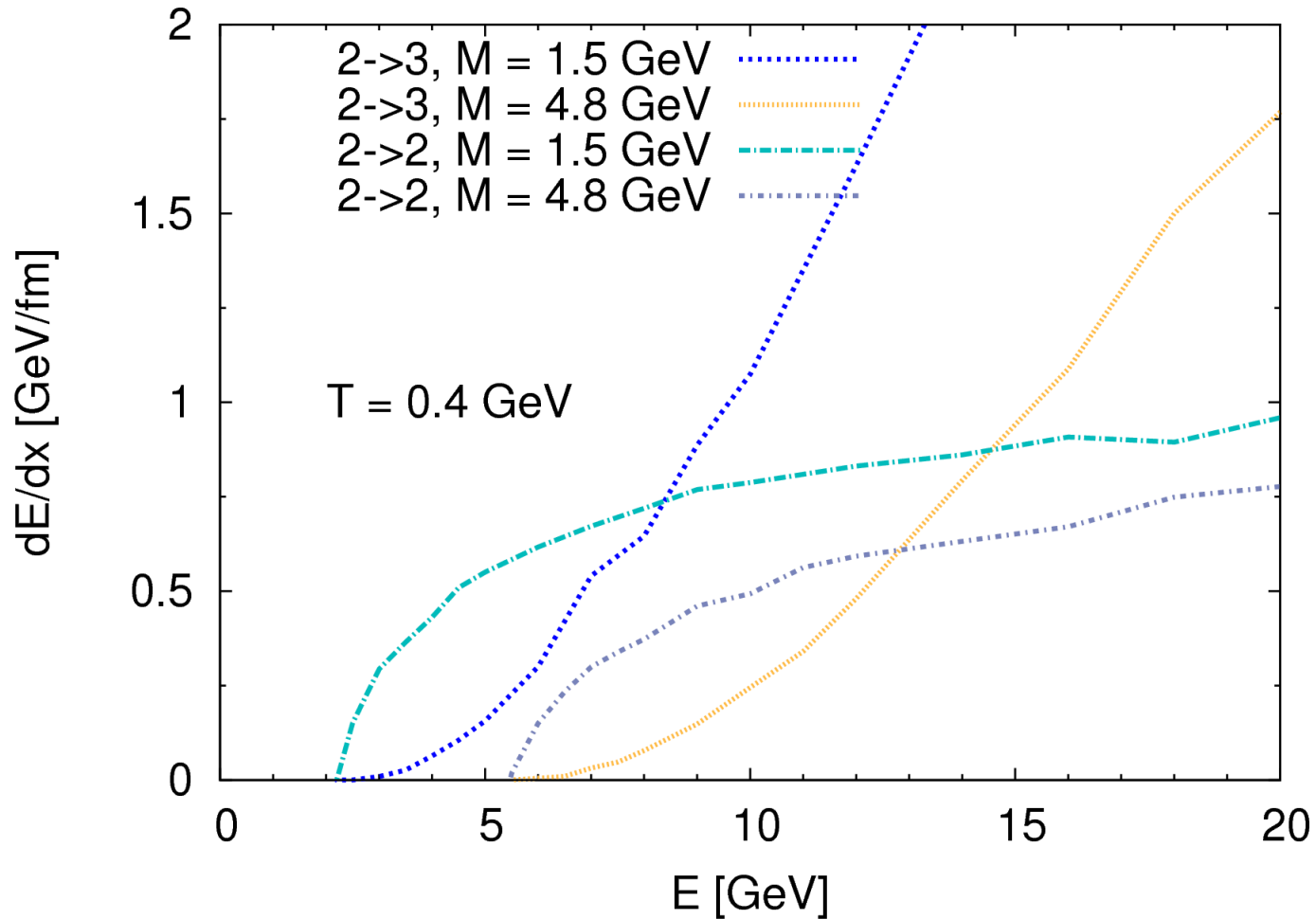


$$\lambda > \tau$$

2 \rightarrow 3 only allowed if mean free path of jet larger than formation time of radiated gluon

➔ Bethe-Heitler regime, independent scatterings

Energy loss in static medium



Running coupling, with LPM effect

Initial conditions

Gluons:

- **PYTHIA**
scaling to heavy-ion collisions with Glauber model (considering shadowing) and energy conservation
- **Minijets**
(low p_T cut-off at 1.4 GeV)
- **Color glass condensate**
H.J. Drescher & Y. Nara, Phys. Rev. C75 (2007)

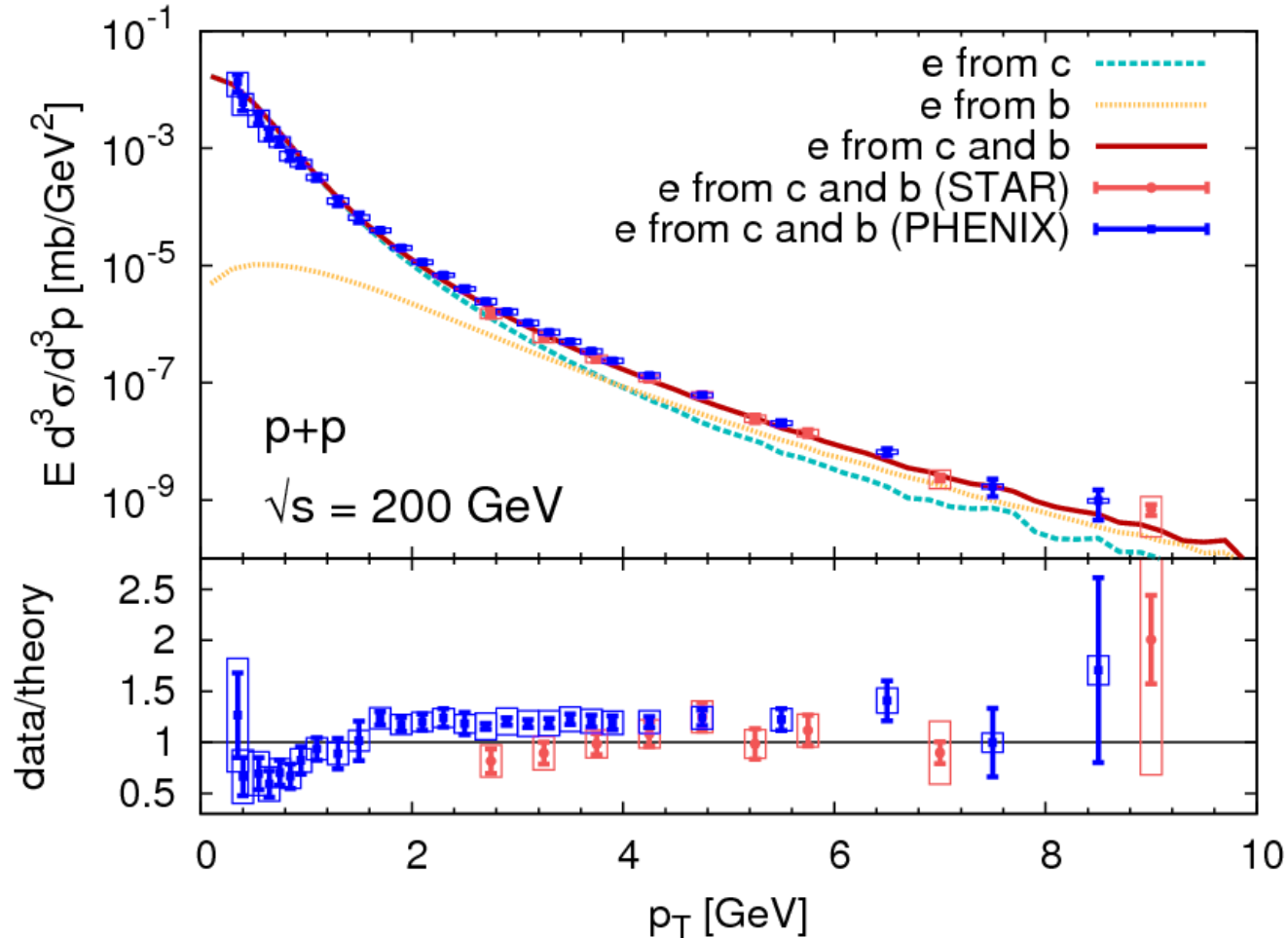
Heavy quarks:

- **PYTHIA**
Monte Carlo Event Generator for nucleon-nucleon collisions



- **NLO pQCD**
Distributions from R. Vogt
- **MC@NLO**
Next-to-leading order matrix elements

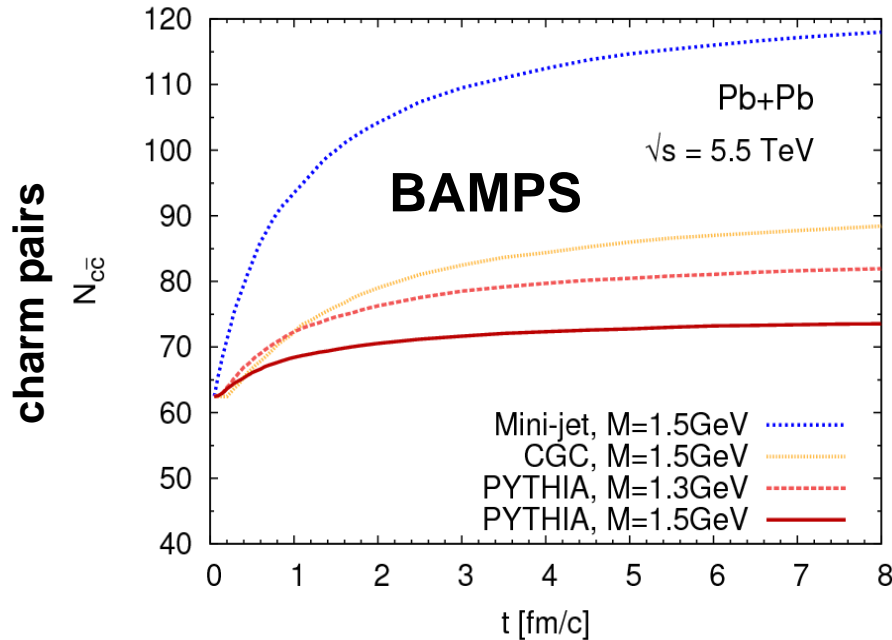
Initial heavy quark distribution



$$\mu_F = \mu_R = 0.65 \sqrt{p_T^2 + M_c^2} \text{ for charm } (M_c = 1.3 \text{ GeV})$$

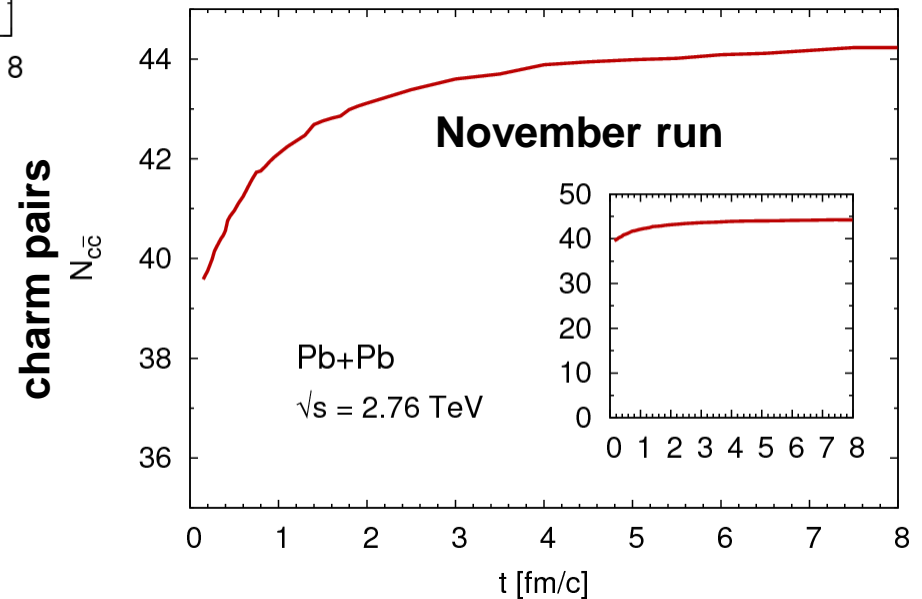
$$\mu_F = \mu_R = 0.4 \sqrt{p_T^2 + M_b^2} \text{ for bottom quarks } (M_b = 4.6 \text{ GeV})$$

Charm production in the QGP at LHC



LHC

Large secondary production
→ **Can even be comparable to initial production**



JU, Fochler, Xu, Greiner
Phys. Rev. C 82 (2010)