

Hadron spectroscopy results from



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(on behalf of the LHCb collaboration)

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European
Research
Council

THE UNIVERSITY OF
WARWICK

Overview

- Introduction to LHCb
- Experimental results
 - Light mesons
 - Charm
 - Charmonium(-like)
 - Beauty
 - Baryons
- Future prospects
- Conclusion

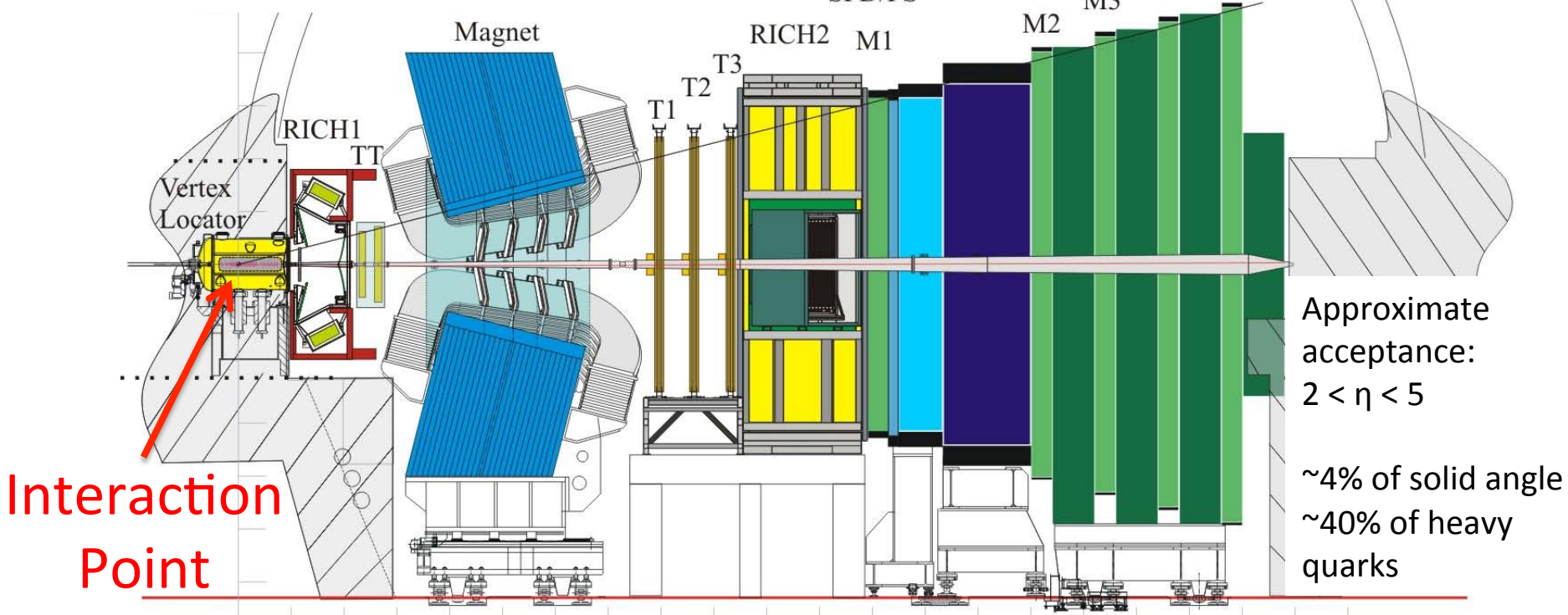
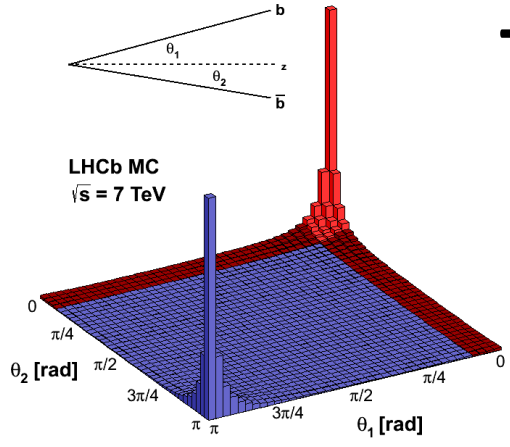


The experiment

- Primary goal of LHCb is to make precision tests of the Standard Model in order to discover or exclude **New Physics** models
- Physics programme is broad:
 - Core measurements of **CP-violation** and **rare decays**
 - Plus **spectroscopy** & **hadronic** physics, EW physics, etc.
- Knowledge from latter parts, which are very interesting in themselves(!), are also key to better precision/accuracy in core programme

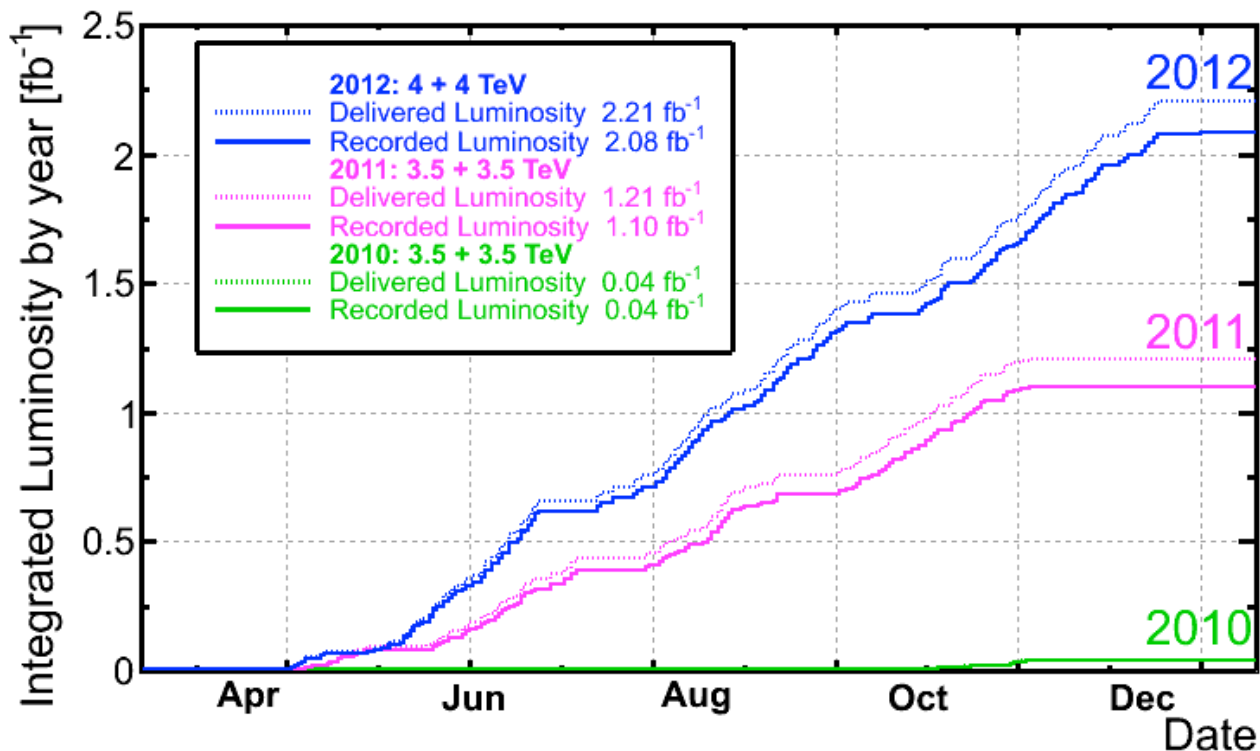
The detector

LHCb $\sigma(pp \rightarrow H_b X) = (75 \pm 5 \pm 13) \mu\text{b}$
 [Phys. Lett. B 694, 209-216 (2010)]

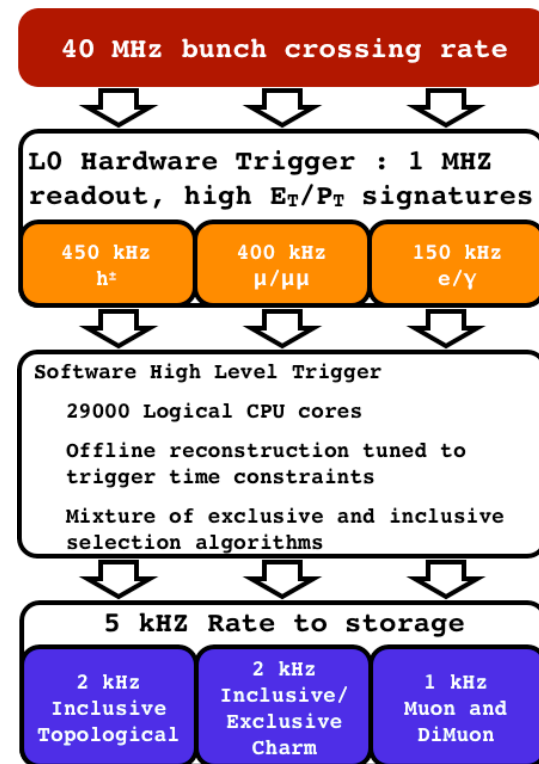




data samples



- 2011: $\sqrt{s} = 7 \text{ TeV}$, $L \approx 1 \text{ fb}^{-1}$
- 2012: $\sqrt{s} = 8 \text{ TeV}$, $L \approx 2 \text{ fb}^{-1}$

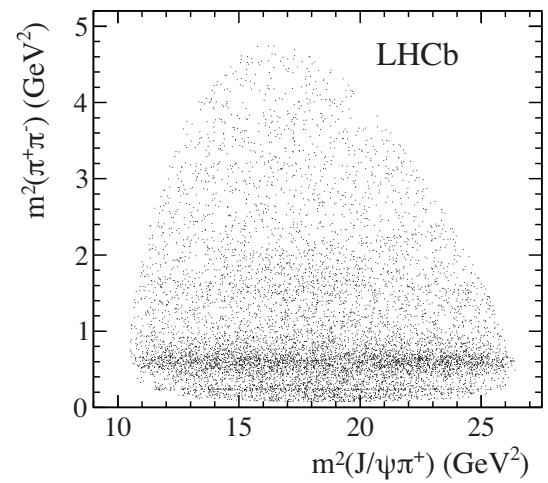
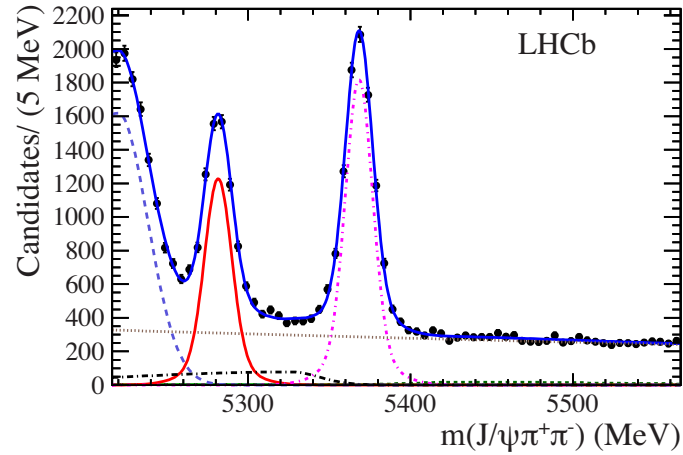


- ☐ Trigger Efficiencies:
 - ☐ ~30% efficient for multi-body hadronic
 - ☐ ~90% efficient for di-muons

Light mesons

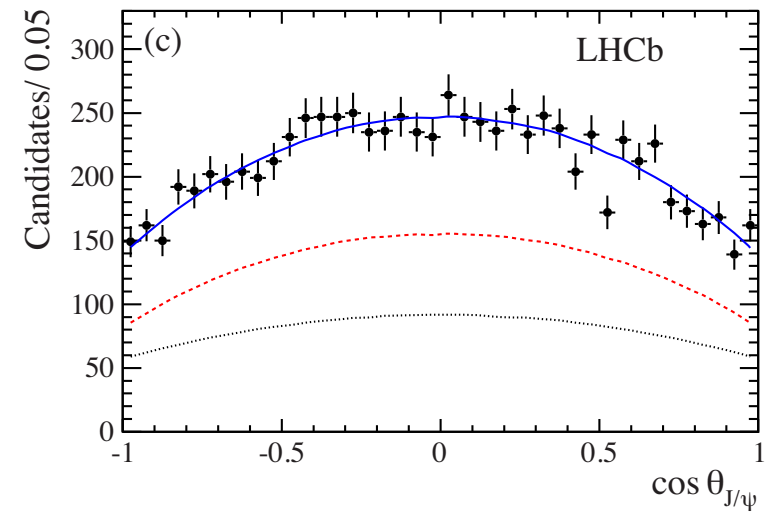
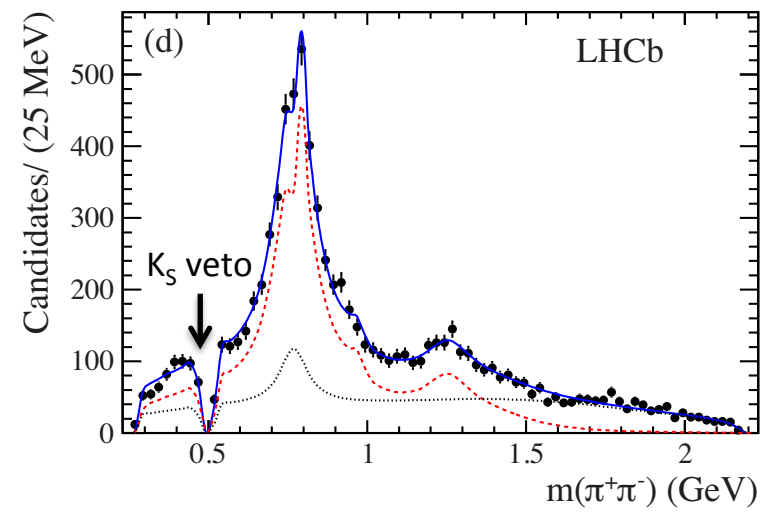
Mixing angle of $f_0(980)$

- Analysis of the resonant structures in $B^0 \rightarrow J/\psi \pi^+ \pi^-$ decays
- Modified Dalitz-plot fit including also the J/ψ decay angle and accounting for different helicity states of intermediate resonances
- Model includes contributions from $\rho(770)$, $\omega(782)$, $f_0(980)$, $f_2(1270)$, $\rho(1450)$ and $f_0(500)$
- Efficiency variation taken into account as a function of DP and J/ψ decay angle, including correlations



Mixing angle of $f_0(980)$

- Largest contribution from $\rho(770)$ [63%]
- Smaller but significant contributions from $f_0(500)$ [16%] and $f_2(1270)$ [9%]
- Evidence of $\omega(782)$ [0.6%] and $\rho(1450)$ [5.3%]
- Significance of $f_0(980)$ [1.5%] only 2.5σ



Mixing angle of $f_0(980)$

- Interpreting $f_0(500)$ and $f_0(980)$ as mixtures of underlying states:

$$|f_0(980)\rangle = \sin \varphi_m \frac{1}{\sqrt{2}} (|u\bar{u}\rangle + |d\bar{d}\rangle) + \cos \varphi_m |s\bar{s}\rangle$$

$$|f_0(500)\rangle = \cos \varphi_m \frac{1}{\sqrt{2}} (|u\bar{u}\rangle + |d\bar{d}\rangle) - \sin \varphi_m |s\bar{s}\rangle$$

- In this decay, only the $d\bar{d}$ part contributes, so

$$\tan^2 \varphi_m = \frac{\text{BF}(B^0 \rightarrow J/\psi f_0(980)) \Phi(500)}{\text{BF}(B^0 \rightarrow J/\psi f_0(500)) \Phi(980)}$$

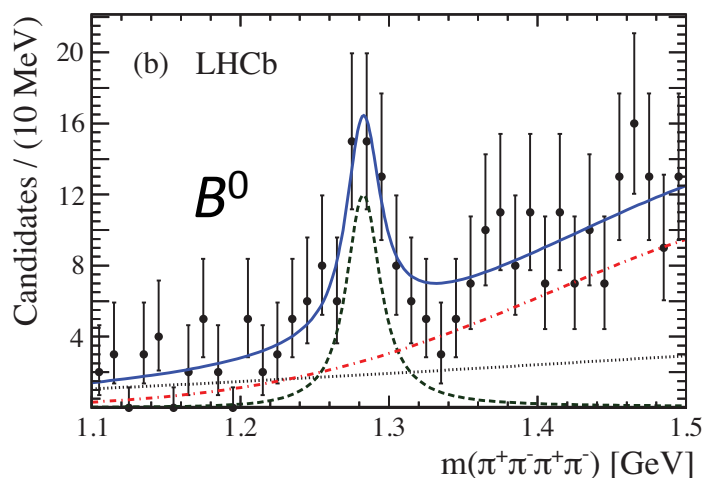
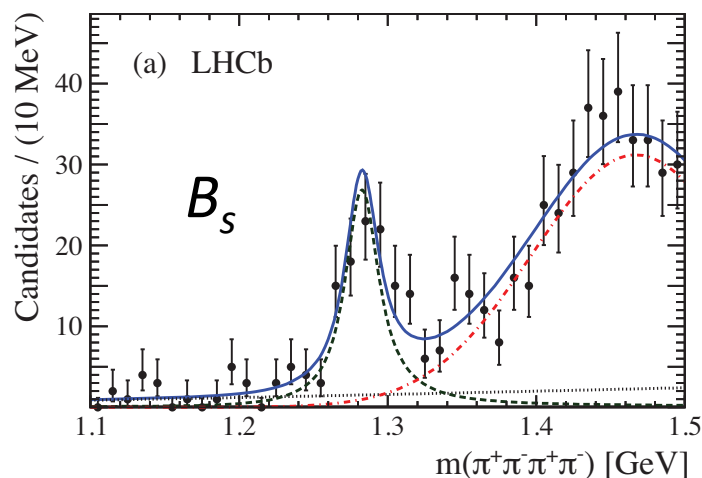
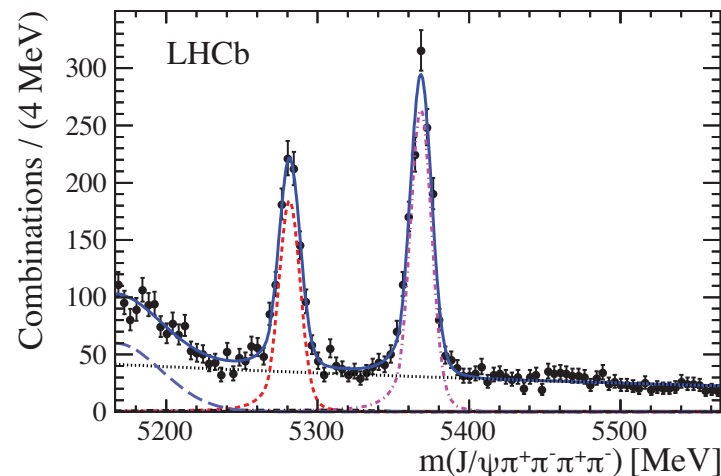
- Where the last term is a ratio of phase space factors
- Using the information from the fit and external information on the f_0 decay BFs we find

$$\tan^2 \varphi_m = \frac{\text{BF}(B^0 \rightarrow J/\psi f_0(980)) \Phi(500)}{\text{BF}(B^0 \rightarrow J/\psi f_0(500)) \Phi(980)} < 0.35 @ 90\% \text{CL}$$

$$|\varphi_m| < 31^\circ @ 90\% \text{CL}$$

Mixing angle of $f_1(1285)$

- Combined analysis of the B^0 and B_s decays to $J/\psi 4\pi$
- Normalise wrt $B \rightarrow J/\psi \pi\pi$
- Fit 4π invariant mass to determine $f_1(1285)$ contribution in each case



Mixing angle of $f_1(1285)$

- Similar interpretation as with $f_0(500)$ and $f_0(980)$
 - $f_1(1285)$ and as yet unidentified mixing partner
- Here use $f_1(1285)$ information from B^0 and B_s decays:

$$\tan^2 \varphi_m = \frac{1}{2} \frac{\text{BF}(B_s^0 \rightarrow J/\psi f_1(1285)) \tau_0 |V_{cd}|^2 \Phi_0}{\text{BF}(B^0 \rightarrow J/\psi f_1(1285)) \tau_s |V_{cs}|^2 \Phi_s} = 0.1970 \pm 0.053^{+0.014}_{-0.012}$$

- The resulting mixing angle is:

$$|\varphi_m| = \left(24.0^{+3.1}_{-2.6} \begin{matrix} +0.6 \\ -0.8 \end{matrix} \right)^\circ$$

- The observed ratio of branching fractions disfavors the tetraquark interpretation of the $f_1(1285)$

[Phys. Rev. Lett. 111 (2013) 062001]

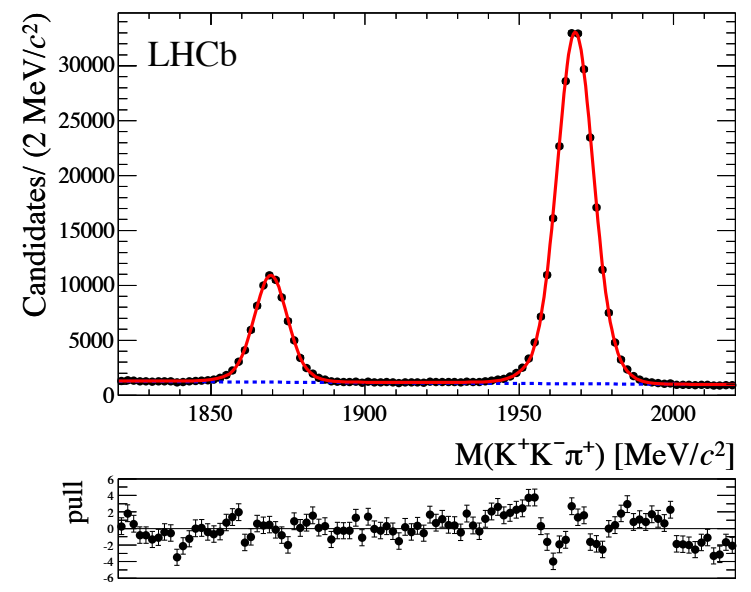
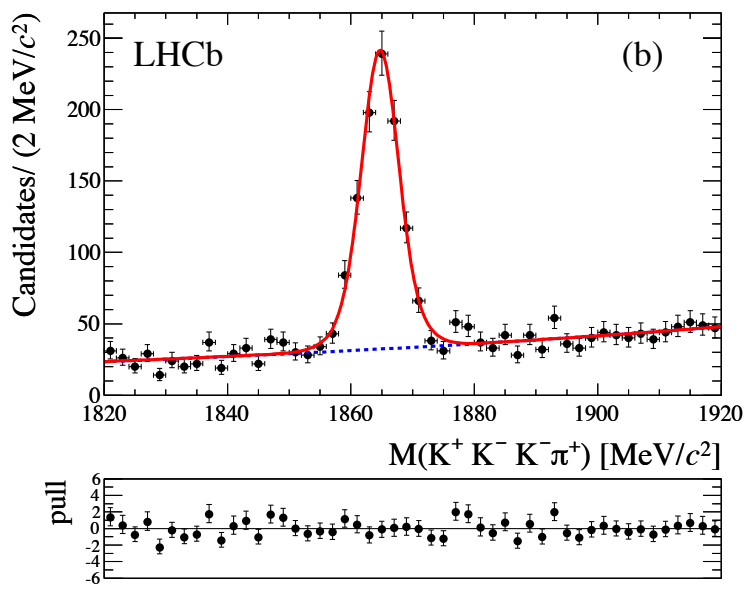
Charm

D meson masses

- D meson masses not so precisely known as B 's
- Precision on world averages of $m(D^0)$, $m(D^+)$ and $m(D_s)$ are 0.2, 0.4 and 1.4 MeV/c² [PDG 2012 Averages]
- Mass splittings slightly better known:
 - $m(D^+) - m(D^0) = 4.76 \pm 0.28$ MeV/c²
 - $m(D_s) - m(D^+) = 98.85 \pm 0.25$ MeV/c²
- Much interest in more precise measurements:
 - Nature of $X(3872)$ – is its mass above or below $m(D^0) + m(D^{*0})$ threshold
 - Largest systematic uncertainty on the B_c mass measurement in decay to $J/\psi D_s$

D meson masses

- Use $D^0 \rightarrow 3K\pi$ to measure D^0 mass
 - Low Q-value, reduced systematics due to detector momentum-scale
- Use $D_{(s)}^+ \rightarrow KK\pi$ (and $D^0 \rightarrow KK\pi\pi$) to measure mass splittings
 - The different D^0 decay is used due to its more similar Q-value to the $D_{(s)}^+$ decay



D meson masses

- Results for mass and mass splittings are:

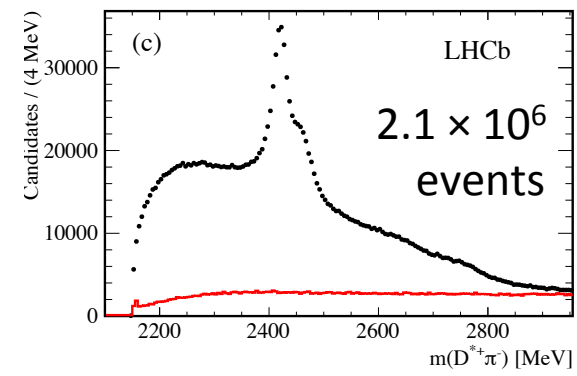
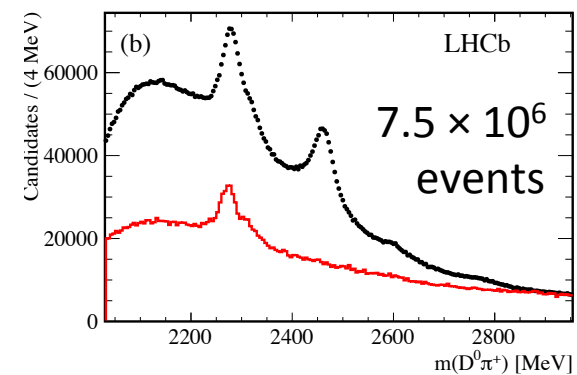
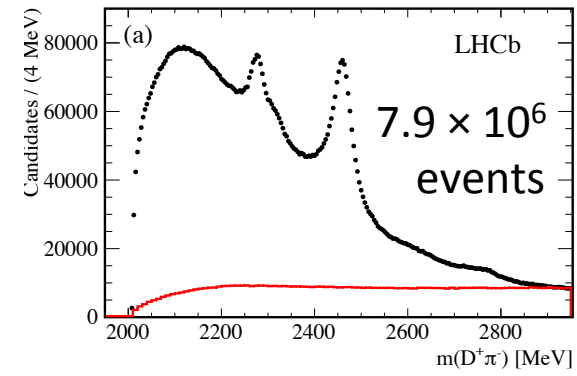
$$\begin{aligned}
 M(D^0) &= 1864.75 \pm 0.15 \text{ (stat)} \pm 0.11 \text{ (syst)} \text{ MeV}/c^2, \\
 M(D^+) - M(D^0) &= 4.76 \pm 0.12 \text{ (stat)} \pm 0.07 \text{ (syst)} \text{ MeV}/c^2, \\
 M(D_s^+) - M(D^+) &= 98.68 \pm 0.03 \text{ (stat)} \pm 0.04 \text{ (syst)} \text{ MeV}/c^2.
 \end{aligned}$$

- Dominant systematic uncertainty is from the knowledge of the detector momentum scale
- Combining with PDG values, can determine more precise value for D_s mass:

$$M(D_s^+) = 1968.19 \pm 0.20 \pm 0.14 \pm 0.08 \text{ MeV}/c^2,$$

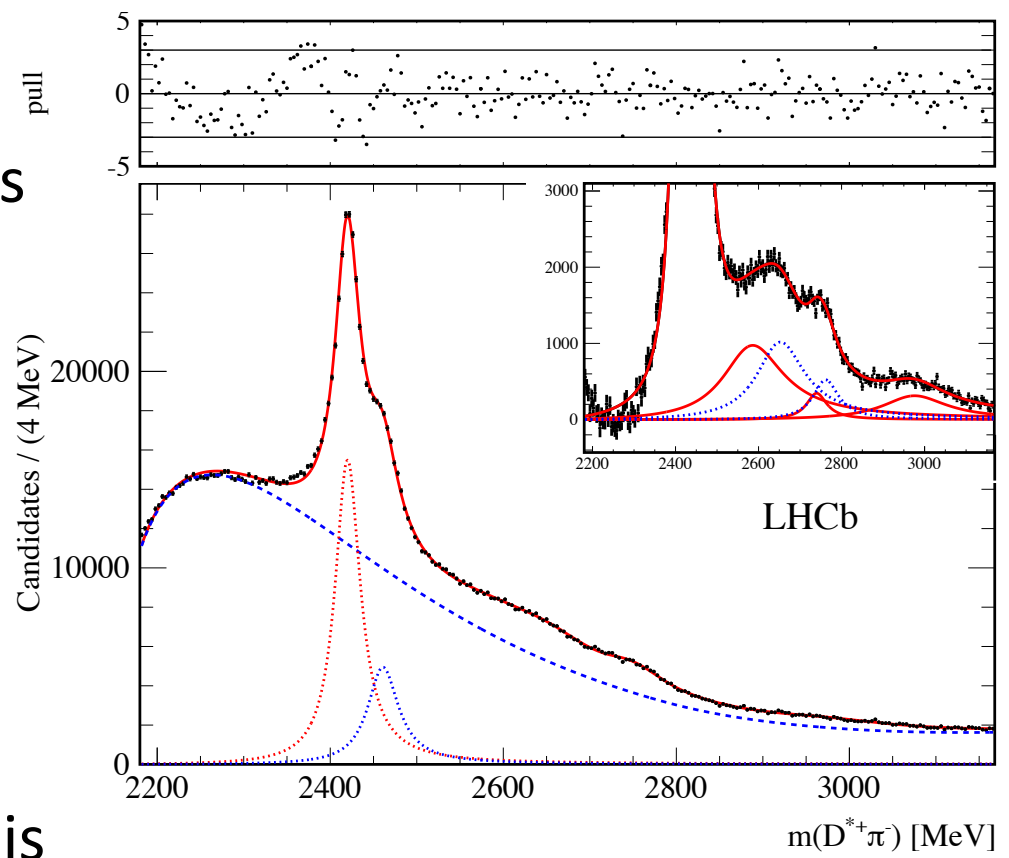
Excited D meson states

- Search for D_J states decaying to $D^+\pi^-$, $D^0\pi^+$ and $D^{*+}\pi^-$
 - $D^+ \rightarrow K^-\pi^+\pi^+$, $D^0 \rightarrow K^-\pi^+$ and $D^{*+} \rightarrow D^0\pi^+$
- D candidates required to be produced at primary collision vertex to reduce contributions from b -hadron decays
- $D^+\pi^-$ and $D^0\pi^+$ spectra contain cross feed from the decay of $D_1(2420)$ or $D_2^*(2460)$ to $D^*\pi$ where the π^0 or γ from the D^* decay is lost
- The wrong-sign spectra are also studied as cross-checks
 - (Shown in red)



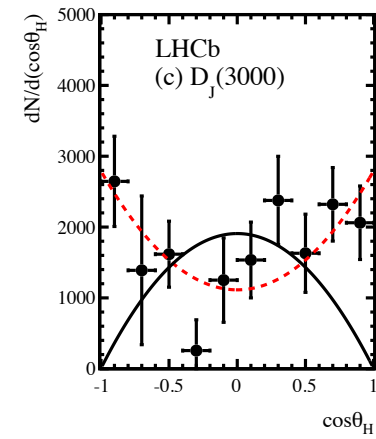
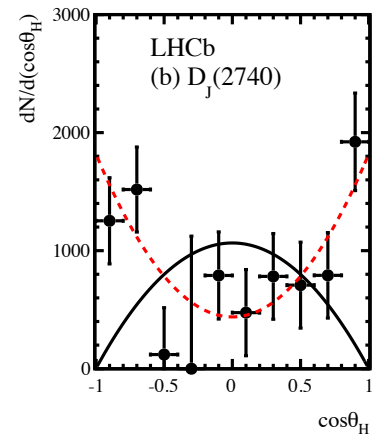
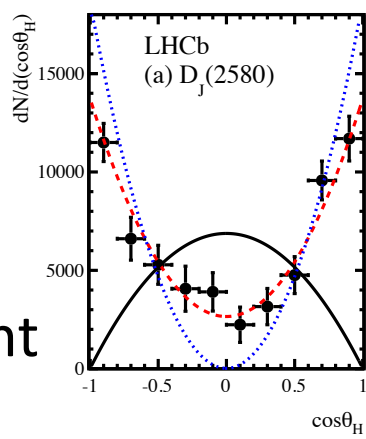
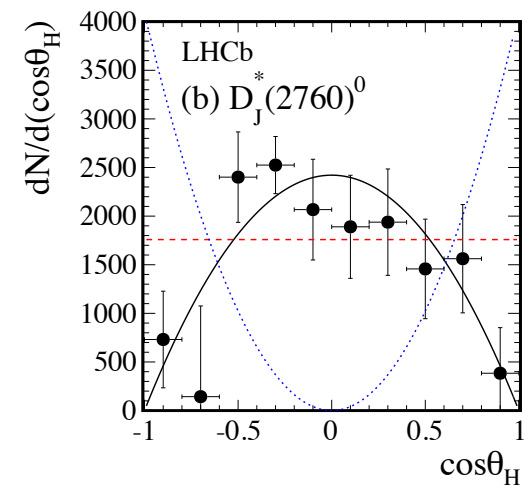
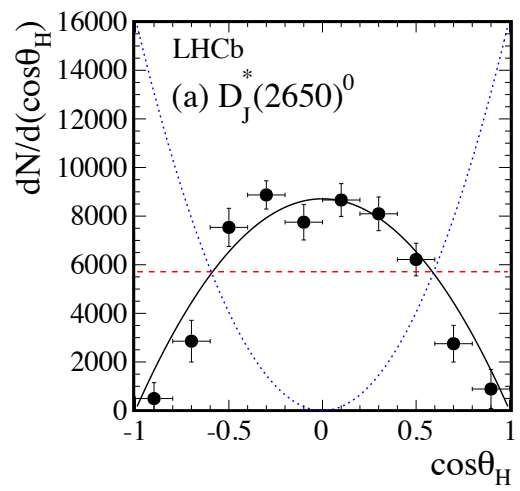
Excited D meson states

- First examine the $D^* \pi$ spectrum
- Fit includes contributions from well established $D_1(2420)$ and $D_2^*(2460)$ plus 5 further states:
 - $D_J(2580)$, $D_J(2650)$,
 $D_J(2740)$, $D_J(2760)$,
 $D_J(3000)$
- Three-body final state allows spin-parity analysis

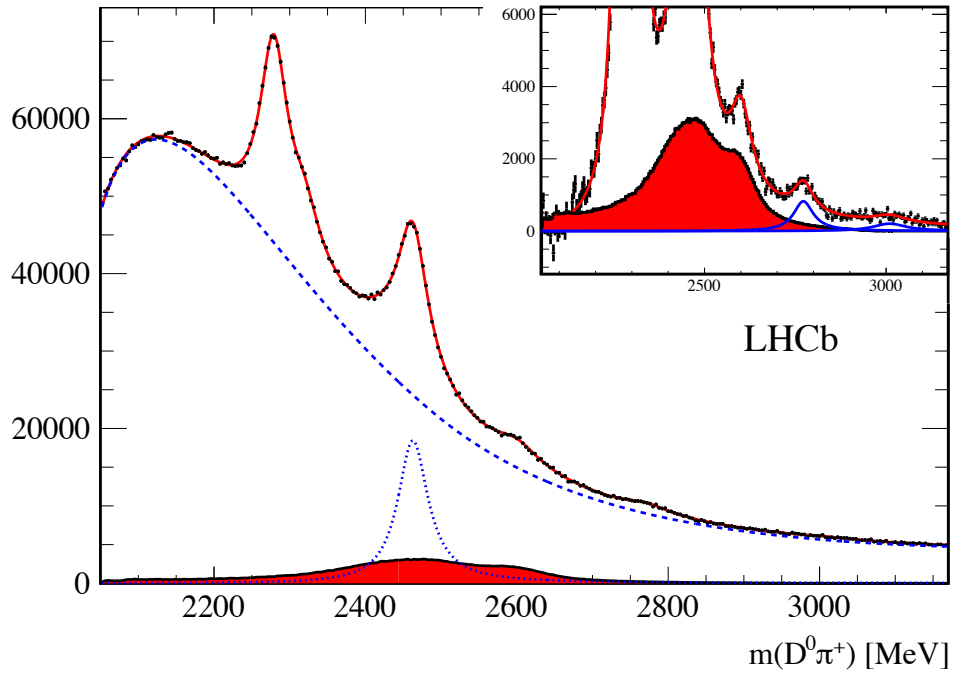
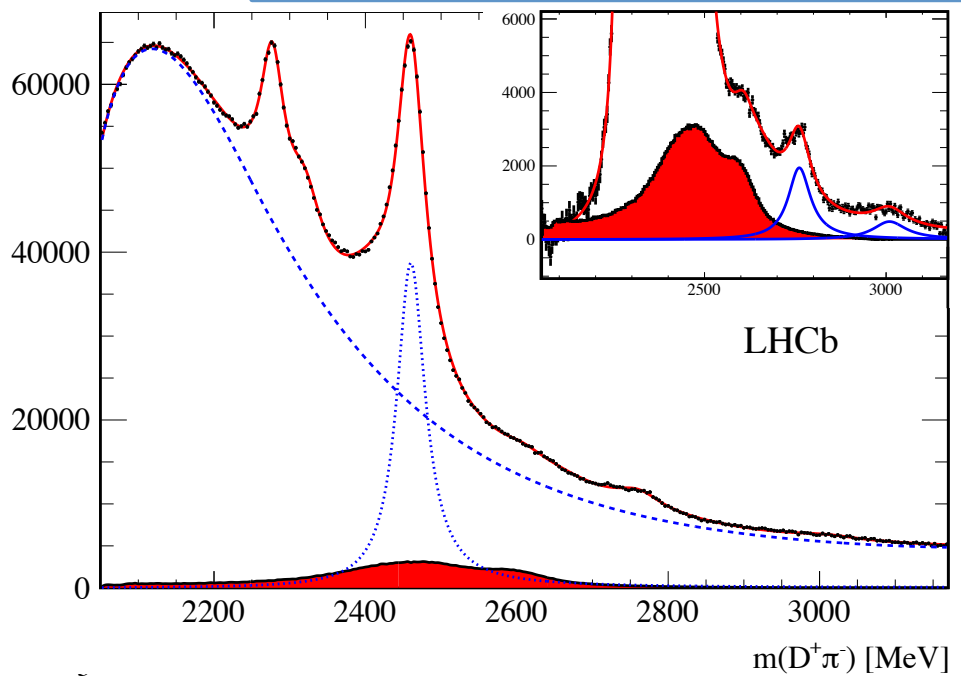


Excited D meson states

- Use angle between two pions in the rest frame of the $D^* \pi$ system (θ_H)
 - Natural parity:
 - $\sin^2(\theta_H)$
 - Unnatural parity:
 - $1 + h \cos^2(\theta_H)$
- Fit mass spectrum in bins of $\cos(\theta_H)$ to determine yields of each contribution
- Use the $\cos(\theta_H)$ distribution to make spin-parity assignment



- Use the $D^*\pi$ results to include a cross feed component into the $D\pi$ spectra (solid red histo)
- Expect contributions from the natural parity states $D_J(2650)$ and $D_J(2760)$
- The $D_J(2760)$ is clearly present but the fit to the $D_J(2650)$ is complicated by the cross feeds
- Mass and width measurements of $D_J(2650)$ are therefore reported only from the $D^*\pi$ fit
- Fits also require a broad structure around 3000 MeV



Excited *D* meson states

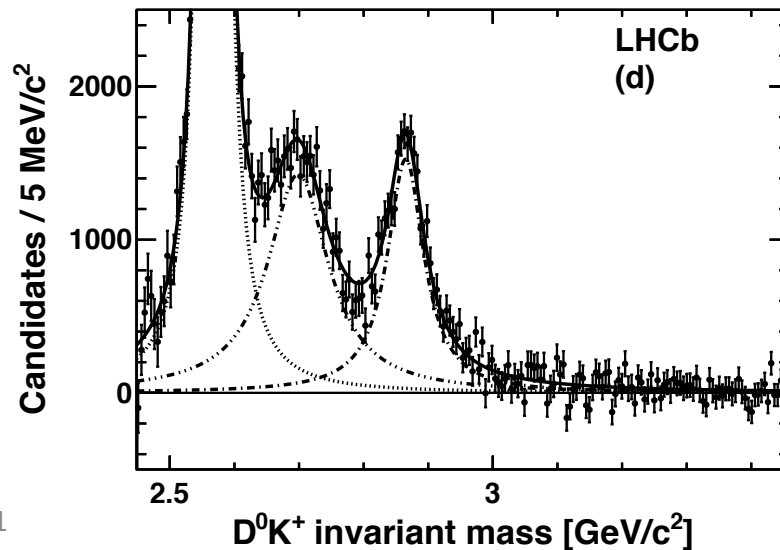
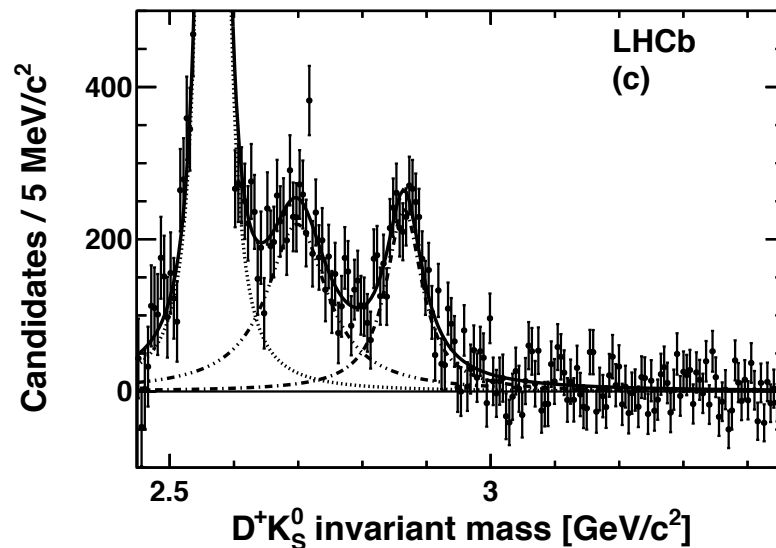
Resonance	Final state	Mass (MeV)	Width (MeV)	Yields × 10 ³	Significance (σ)
<i>D</i> ₁ (2420) ⁰	<i>D</i> ⁺⁺ π ⁻	2419.6 ± 0.1 ± 0.7	35.2 ± 0.4 ± 0.9	210.2 ± 1.9 ± 0.7	
<i>D</i> ₂ [*] (2460) ⁰	<i>D</i> ⁺⁺ π ⁻	2460.4 ± 0.4 ± 1.2	43.2 ± 1.2 ± 3.0	81.9 ± 1.2 ± 0.9	
<i>D</i> _{<i>J</i>} [*] (2650) ⁰	<i>D</i> ⁺⁺ π ⁻	2649.2 ± 3.5 ± 3.5	140.2 ± 17.1 ± 18.6	50.7 ± 2.2 ± 2.3	24.5
<i>D</i> _{<i>J</i>} [*] (2760) ⁰	<i>D</i> ⁺⁺ π ⁻	2761.1 ± 5.1 ± 6.5	74.4 ± 3.4 ± 37.0	14.4 ± 1.7 ± 1.7	10.2
<i>D</i> _{<i>J</i>} (2580) ⁰	<i>D</i> ⁺⁺ π ⁻	2579.5 ± 3.4 ± 5.5	177.5 ± 17.8 ± 46.0	60.3 ± 3.1 ± 3.4	18.8
<i>D</i> _{<i>J</i>} (2740) ⁰	<i>D</i> ⁺⁺ π ⁻	2737.0 ± 3.5 ± 11.2	73.2 ± 13.4 ± 25.0	7.7 ± 1.1 ± 1.2	7.2
<i>D</i> _{<i>J</i>} (3000) ⁰	<i>D</i> ⁺⁺ π ⁻	2971.8 ± 8.7	188.1 ± 44.8	9.5 ± 1.1	9.0
<i>D</i> ₂ [*] (2460) ⁰	<i>D</i> ⁺ π ⁻	2460.4 ± 0.1 ± 0.1	45.6 ± 0.4 ± 1.1	675.0 ± 9.0 ± 1.3	
<i>D</i> _{<i>J</i>} [*] (2760) ⁰	<i>D</i> ⁺ π ⁻	2760.1 ± 1.1 ± 3.7	74.4 ± 3.4 ± 19.1	55.8 ± 1.3 ± 10.0	17.3
<i>D</i> _{<i>J</i>} [*] (3000) ⁰	<i>D</i> ⁺ π ⁻	3008.1 ± 4.0	110.5 ± 11.5	17.6 ± 1.1	21.2
<i>D</i> ₂ [*] (2460) ⁺	<i>D</i> ⁰ π ⁺	2463.1 ± 0.2 ± 0.6	48.6 ± 1.3 ± 1.9	341.6 ± 22.0 ± 2.0	
<i>D</i> _{<i>J</i>} [*] (2760) ⁺	<i>D</i> ⁰ π ⁺	2771.7 ± 1.7 ± 3.8	66.7 ± 6.6 ± 10.5	20.1 ± 2.2 ± 1.0	18.8
<i>D</i> _{<i>J</i>} [*] (3000) ⁺	<i>D</i> ⁰ π ⁺	3008.1 (fixed)	110.5 (fixed)	7.6 ± 1.2	6.6

- Tentative spin parity assignments made comparing measured parameter values and results of spin-parity analysis
- *D*_{*J*}(2580): 2S state with J^P = 0⁻
- *D*_{*J*}(2650): 2S state with J^P = 1⁻
- *D*_{*J*}(2760): 1D state with J^P = 1⁻
- *D*_{*J*}(2740): 1D state with J^P = 2⁻ (although the measured and predicted mass do not agree well)

Excited D_s meson states

- Similar study of D_{sJ} states decaying to $D^+K_S^0$ and D^0K^+
- Confirms the existence of $D_{s1}^*(2700)$ and $D_{sJ}^*(2860)$ states seen by the B -factory analyses
 - First observations of these states in hadronic interactions
- Measure their parameters to be consistent with B -factory measurements:

$$\begin{aligned}
 m(D_{s1}^*(2700)^+) &= 2709.2 \pm 1.9(\text{stat}) \pm 4.5(\text{syst}) \text{ MeV}/c^2, \\
 \Gamma(D_{s1}^*(2700)^+) &= 115.8 \pm 7.3(\text{stat}) \pm 12.1(\text{syst}) \text{ MeV}/c^2, \\
 m(D_{sJ}^*(2860)^+) &= 2866.1 \pm 1.0(\text{stat}) \pm 6.3(\text{syst}) \text{ MeV}/c^2, \\
 \Gamma(D_{sJ}^*(2860)^+) &= 69.9 \pm 3.2(\text{stat}) \pm 6.6(\text{syst}) \text{ MeV}/c^2.
 \end{aligned}$$



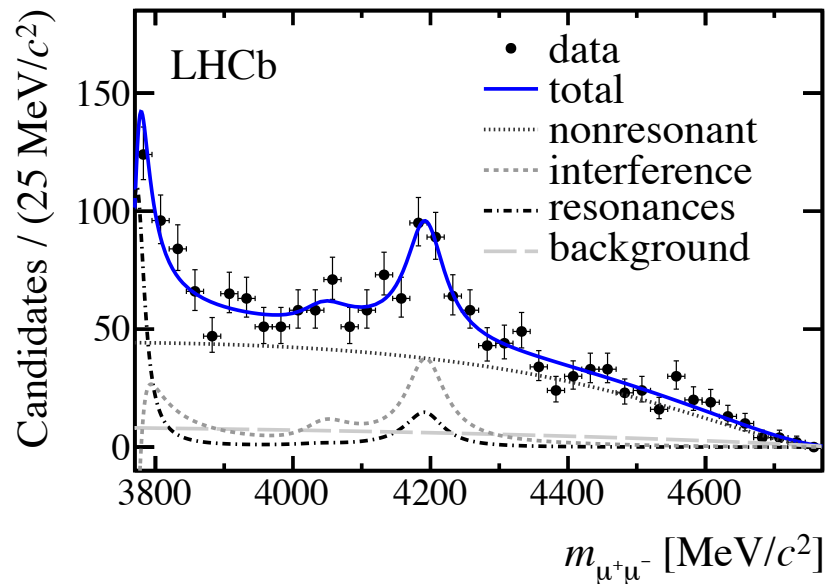
Charmonium(-like)

$$B^+ \longrightarrow \psi(4160) K^+$$

- $B^+ \longrightarrow K^+ \mu^+ \mu^-$ decays used to probe contributions to the loop processes from new physics amplitudes
- Important to understand the tree-level contributions via vector charmonium resonances
- Dedicated analysis of the decays $B^+ \longrightarrow K^+ \mu^+ \mu^-$ above open-charm threshold
- Require highly displaced vertex to remove combinatorial backgrounds
- Backgrounds from B decays reduced with strong particle identification criteria and vetoes on misidentified $\mu^- K^+$ combinations
- Fit to the B -candidate invariant mass used to estimate background level

$B^+ \rightarrow \psi(4160) K^+$

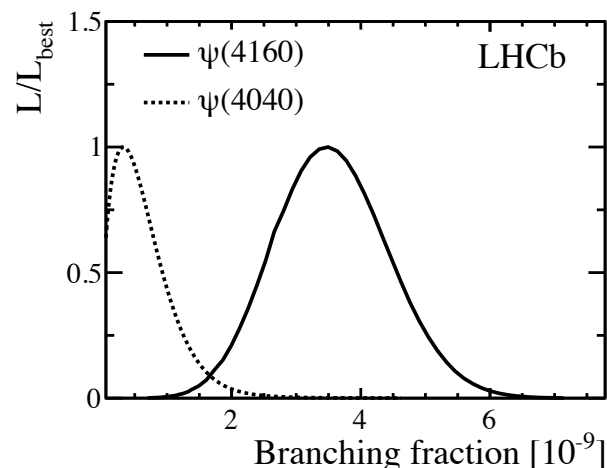
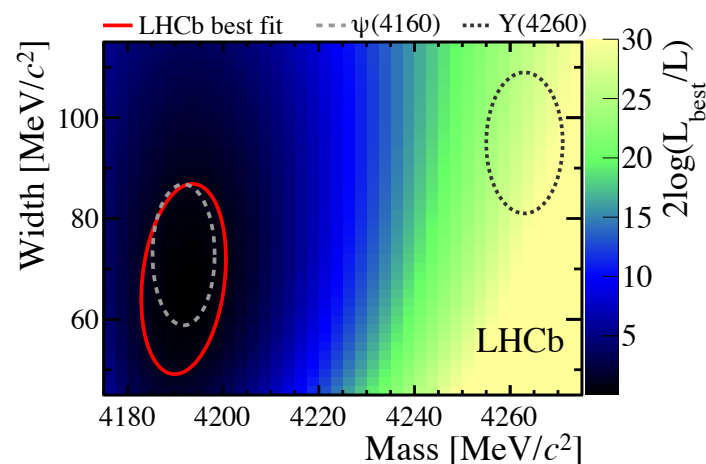
- $\mu^+\mu^-$ invariant mass resolution improved using constraints on the vertex fit:
 - Candidate invariant mass constrained to nominal B mass
 - Candidate momentum must point back to associated primary vertex
- Fit includes terms for:
 - Vector and axial vector nonresonant
 - Vector resonances (including $\psi(3770)$ – on far left of plot)
 - Interference between the resonances and vector NR



- Contributions for $\psi(4160)$ and $\psi(4040)$ are included
- $\psi(4160)$ is observed with significance $> 6\sigma$
- Signal yield for $\psi(4040)$ is not significant

$B^+ \rightarrow \psi(4160) K^+$

- First observations of the decays $B^+ \rightarrow \psi(4160)K^+$ and $\psi(4160) \rightarrow \mu^+\mu^-$
- Mass and width are consistent with previous measurements
- Branching fraction: $(5.1 \pm 1.3 \pm 3.0) \times 10^{-4}$ assuming lepton universality
- Contribution from $\psi(4160)$ is $\sim 20\%$ of total signal in low-recoil region – larger than theoretical estimates

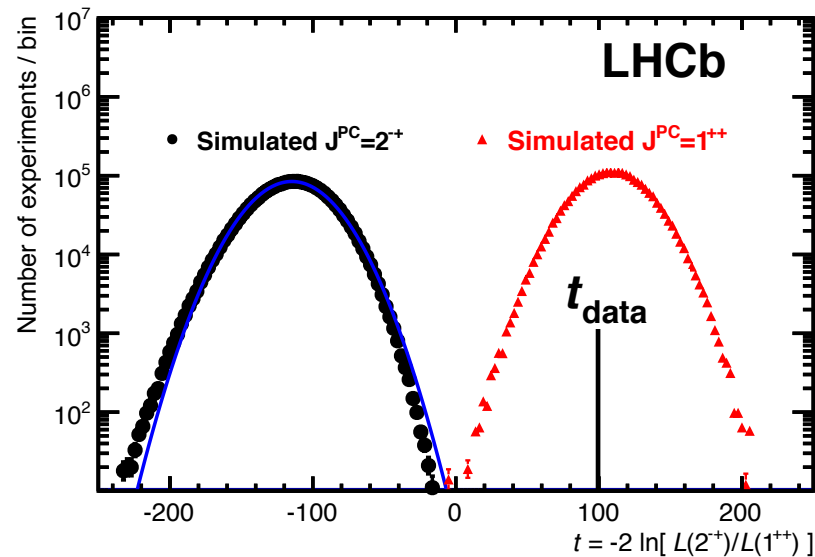
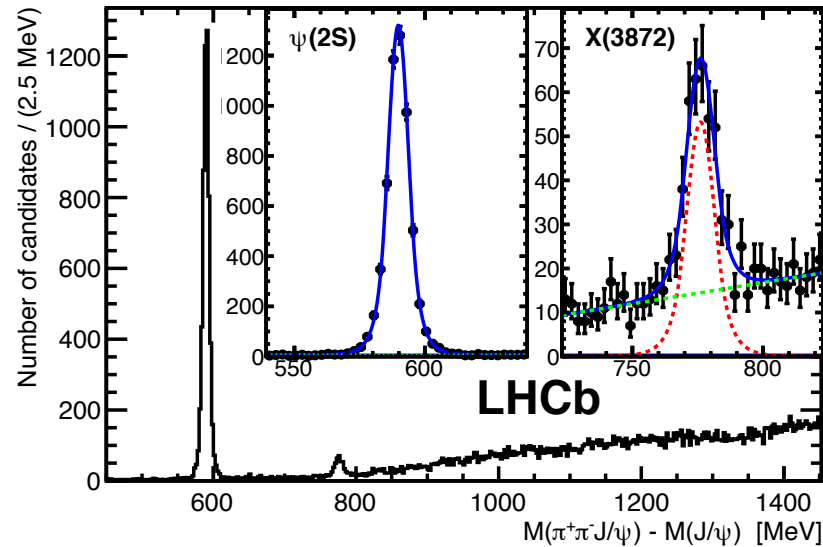


Properties of $X(3872)$

- The nature of the $X(3872)$ particle remains unclear over 10 years after its discovery by Belle
- Possibilities include conventional charmonium or exotic states like tetra-quarks or D^*D molecules
 - Including latest $m(D^0)$ result (presented earlier) in determination of $X(3872) = D^*D$ molecule binding energy gives $E_B = 0.09 \pm 0.28 \text{ MeV}/c^2$
 - i.e. if it is a molecular bound state it is very loosely bound
- Its **quantum numbers** are important inputs to the puzzle
 - Possible values reduced to $J^{PC} = 1^{++}$ or 2^{-+} by a 3D angular analysis by the CDF collaboration
- LHCb have performed the first **5D angular analysis** of the decay chain $B^+ \rightarrow X(3872)K^+$, $X(3872) \rightarrow J/\psi\pi\pi$, $J/\psi \rightarrow \mu\mu$
 - The similar decay substituting the $\psi(2S)$ for the $X(3872)$ is used as a control sample for many aspects of the analysis

Properties of X(3872)

- Likelihood ratio constructed using the PDFs of the 5D angular space:
 - $X, J/\psi$ and $\pi\pi$ helicity angles
 - angle between X and J/ψ decay planes
 - angle between X and $\pi\pi$ decay planes
- Angular distributions are constructed in the helicity formalism
 - 2 free parameters are floated in the fit to data
- **2⁻⁺ hypothesis is excluded at 8.4σ**
- Angular parameters agree well with Belle values and values obtained from fitting 1⁺⁺ simulation
- This result **rules out** the X(3872) being the $\eta_{c2}(1^1D_2)$ conventional charmonium state

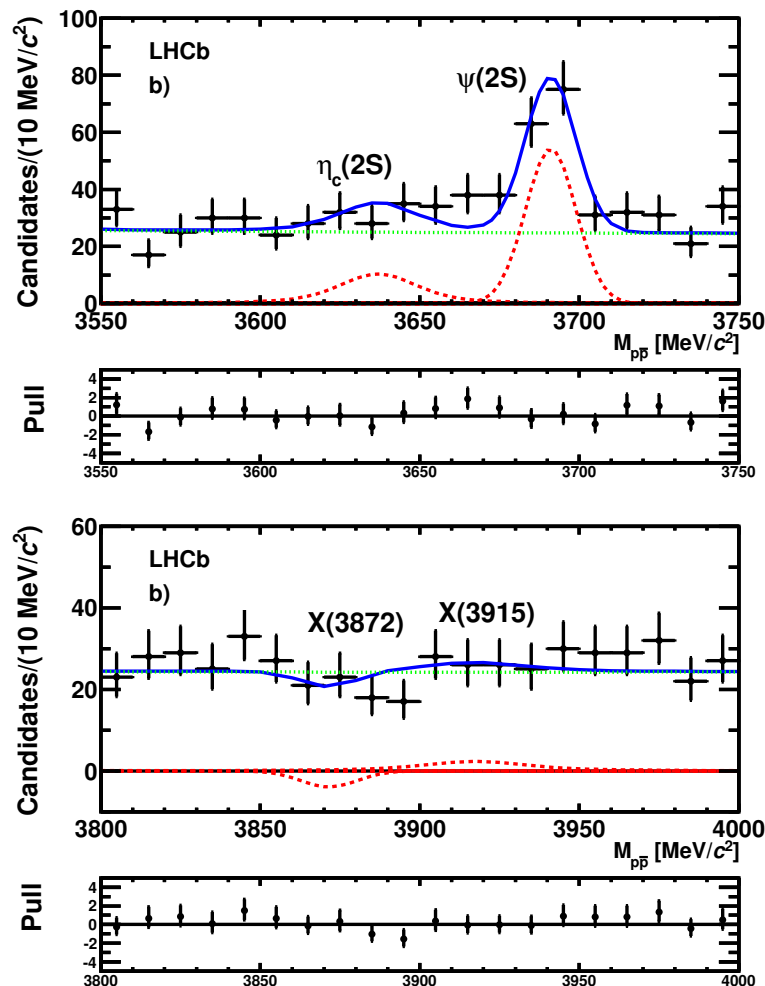


Properties of X(3872)

- A separate analysis of $B^+ \rightarrow p\bar{p}K^+$ decays places limits on branching fraction of $X(3872) \rightarrow p\bar{p}$
- Fits are performed to the $p\bar{p}$ invariant mass spectrum, including background, nonresonant and resonant terms
- No significant signal is observed for the $X(3872)$ decay and a limit is placed on the BF of:

$$\frac{\mathcal{B}(X(3872) \rightarrow p\bar{p})}{\mathcal{B}(X(3872) \rightarrow J/\psi\pi^+\pi^-)} < 2.0 \times 10^{-3}$$

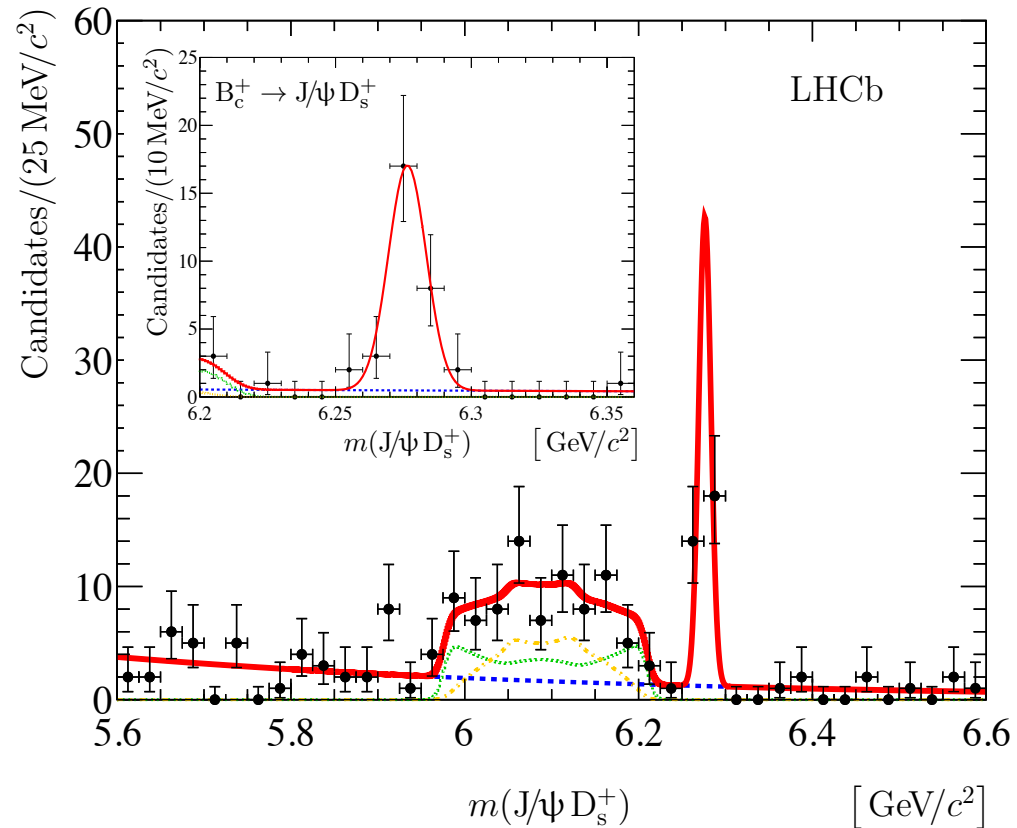
- This value is in tension with some predictions based on the molecular interpretation of the state [[Phys. Rev. D77 \(2008\) 034019](#)]



B_c properties

Mass of the B_c meson

- Properties of B_c meson are still rather poorly determined
- Lifetime is rather shorter than other B mesons
 - Indicates importance of charm quark in the weak decay
 - See also recent LHCb observation of $B_c \rightarrow B_s \pi$
- [Phys. Rev. Lett. 111 (2013) 181801]
- Recent observation of B_c decays to $J/\psi D_s$ provide a way to reduce the systematic uncertainty on the mass measurement
- Obtains a yield of 29 signal events



Mass of the B_c meson

- Mass determination makes use of recent improvement in D_s mass measurement (discussed earlier)
- Also helped by low Q-value for the decay
- Dominant systematics from knowledge of momentum scale and detector material
- Mass determined to be:

$$m(B_c) = 6276.28 \pm 1.44 \pm 0.36 \text{ MeV}/c^2$$

- Uncertainties on momentum scale and D_s mass largely cancel in the mass difference:

$$m(B_c) - m(D_s) = 4307.97 \pm 1.44 \pm 0.20 \text{ MeV}/c^2$$

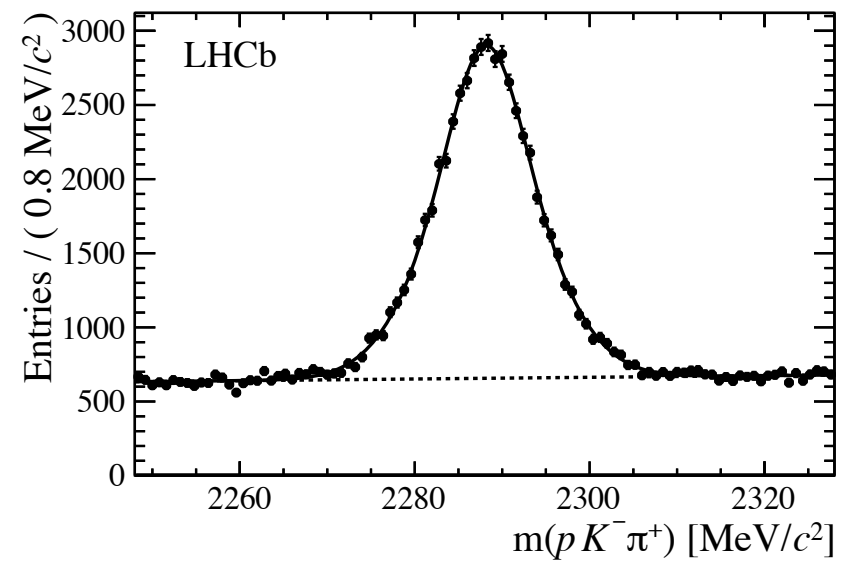
Baryons

Search for Ξ_{cc}^+ baryon

- Signals claimed by SELEX collaboration for Ξ_{cc}^+ decaying to both $\Lambda_c^+ K^- \pi^+$ and $D^+ p K^-$
- In contrast, *B*-factories and FOCUS have found no evidence
- Search for Ξ_{cc}^+ decaying to $\Lambda_c^+ K^- \pi^+$
- Production cross-section determined relative to that of Λ_c^+

$$R \equiv \frac{\sigma(\Xi_{cc}^+) \text{BF}(\Xi_{cc}^+ \rightarrow \Lambda_c^+ K^- \pi^+)}{\sigma(\Lambda_c^+)}$$

- Both the Ξ_{cc}^+ and Λ_c^+ candidates are required to have vertices displaced from the primary collision vertex

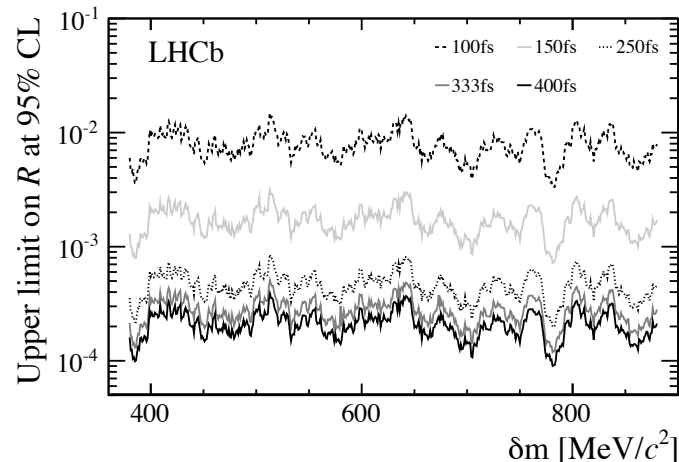
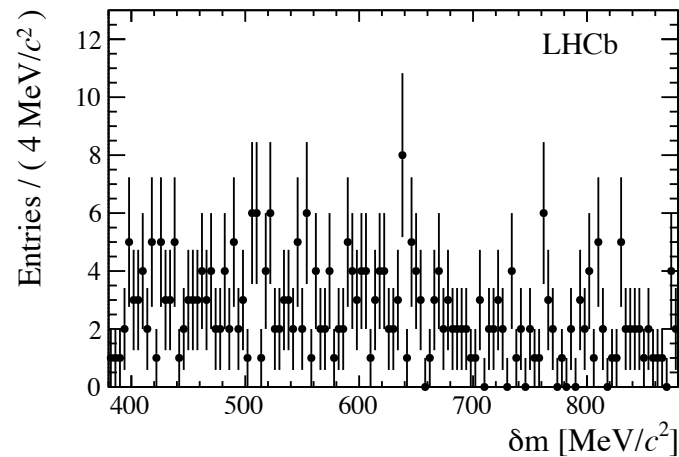


- The Λ_c^+ is reconstructed decaying to $p K^- \pi^+$ and same Λ_c^+ selection is applied in both signal & normalisation modes
- Λ_c^+ candidate invariant mass required to be within 2273 – 2303 MeV/c² window

Search for Ξ_{cc}^+ baryon

- Search is performed at 1 MeV steps in 10 MeV windows of the mass difference:

$$\delta m = m([pK^-\pi^+] K^-\pi^+) - m(pK^-\pi^+) - m(K^-) - m(\pi^+)$$
- Events outside window used to subtract background (two different methods used as cross-check)
- Ratios of efficiencies calculated using simulation and correcting for known data/simulation differences
 - Variation as a function of lifetime and δm determined by reweighting
- **No significant signals** seen and upper limits on R are determined as a function of lifetime and δm
- The current results neither confirm nor exclude the SELEX observation



Masses of Λ_b , Ξ_b^- and Ω_b^-

- Previous measurements of the Ω_b^- mass by CDF and D0 differ by more than 6σ
- The b-baryons are reconstructed in their decays to $J/\psi \Lambda$, $J/\psi \Xi^-$ and $J/\psi \Omega^-$, respectively
- The subsequent decays $\Xi^- \rightarrow \Lambda \pi^-$ and $\Omega^- \rightarrow \Lambda K^-$ are considered
- These decay chains contain several long-lived particles and hence have a characteristic event topology

Masses of Λ_b^- , Ξ_b^- and Ω_b^-

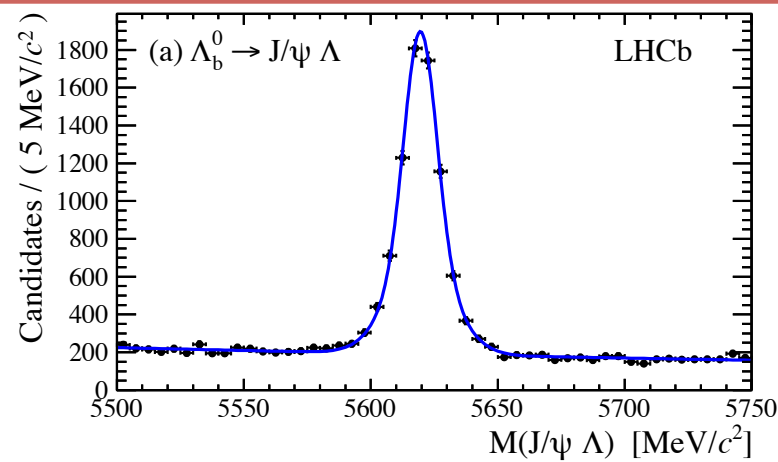
- Masses measured with much improved precision
- Measurement of $M(\Omega_b^-)$ in agreement with CDF but not D0
- Mass differences also measured:

$$M(\Xi_b^-) - M(\Lambda_b^0) = 176.2 \pm 0.9 \text{ (stat)} \pm 0.1 \text{ (syst)} \text{ MeV}/c^2,$$

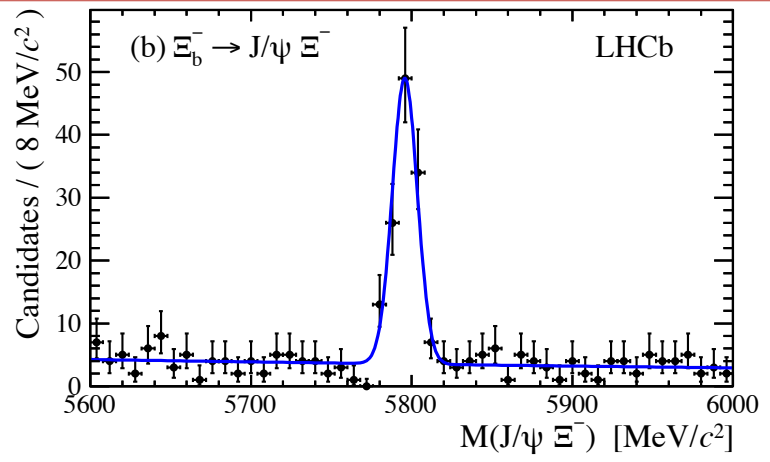
$$M(\Omega_b^-) - M(\Lambda_b^0) = 426.4 \pm 2.2 \text{ (stat)} \pm 0.4 \text{ (syst)} \text{ MeV}/c^2.$$

- Largest systematic uncertainties from momentum scale mostly cancel but those from hyperon masses remain

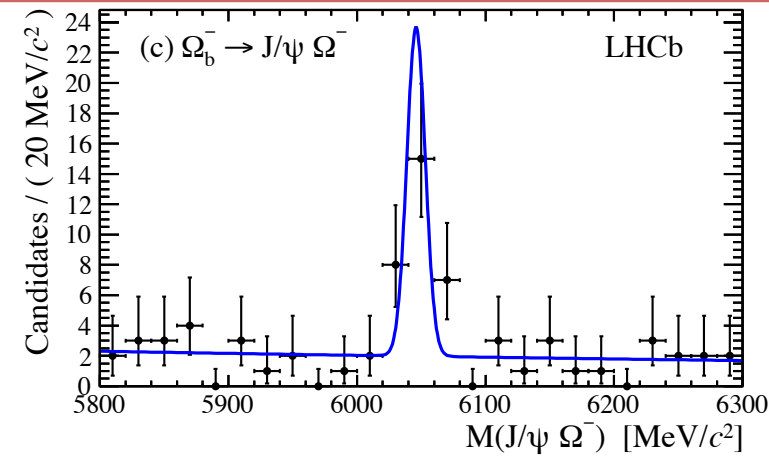
$$M(\Lambda_b^0) = 5619.53 \pm 0.13 \text{ (stat)} \pm 0.45 \text{ (syst)} \text{ MeV}/c^2$$



$$M(\Xi_b^-) = 5795.8 \pm 0.9 \text{ (stat)} \pm 0.4 \text{ (syst)} \text{ MeV}/c^2$$

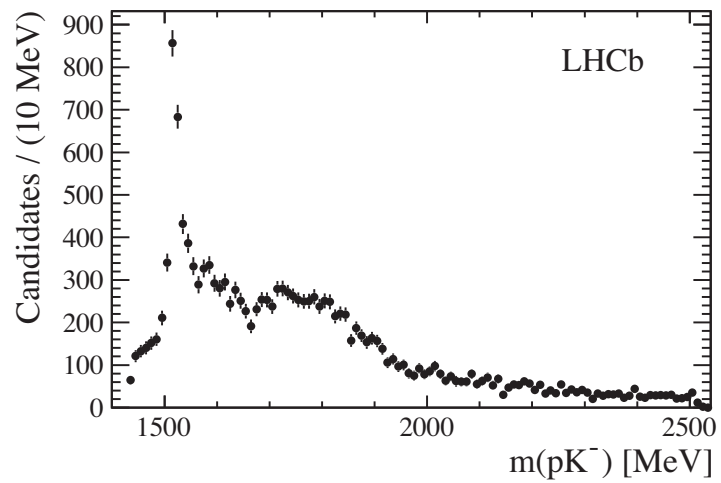
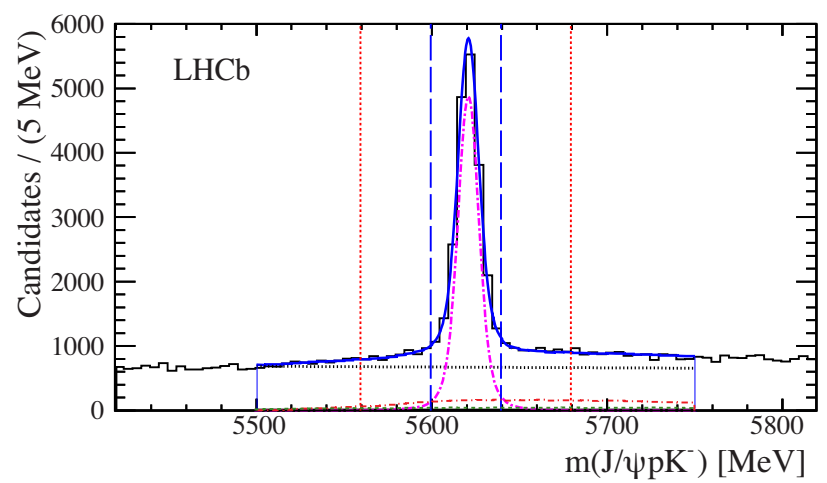


$$M(\Omega_b^-) = 6046.0 \pm 2.2 \text{ (stat)} \pm 0.5 \text{ (syst)} \text{ MeV}/c^2$$



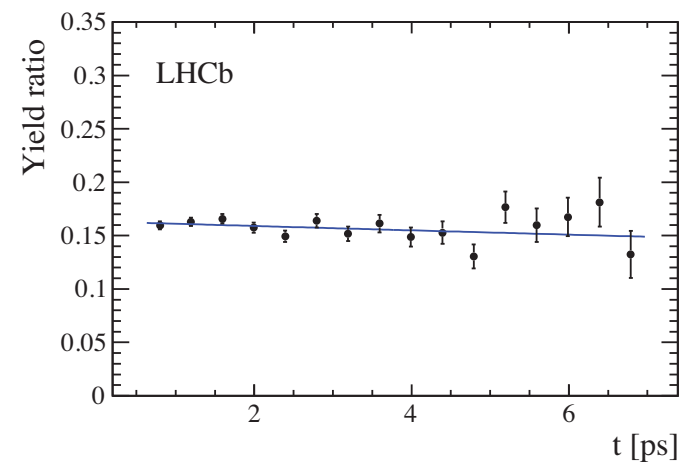
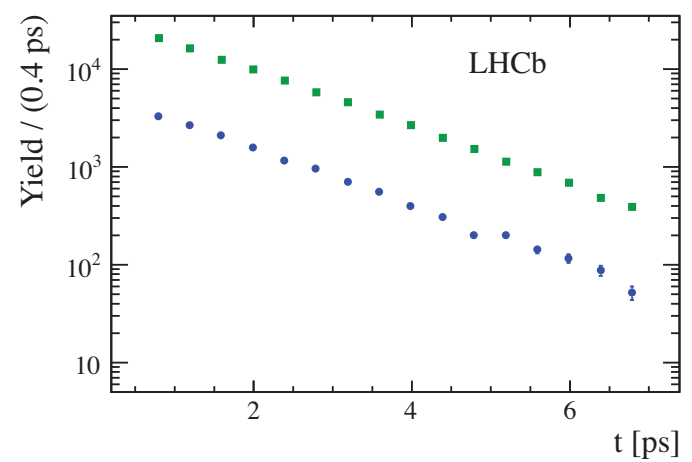
Lifetime of Λ_b

- Previous measurements of the Λ_b lifetime found much smaller values than expected from HQE
- More recent measurements by ATLAS, CMS, CDF and D0 indicate larger values
- A precision measurement is needed to resolve the discrepancy
- This analysis uses decay $\Lambda_b \rightarrow J/\psi p K^-$
- Measurement is made relative to B^0 lifetime in its decay to $J/\psi \pi^+ K^-$ (in region of $K^*(892)$ resonance)
- Topological and kinematic similarities allows cancellation of several systematic uncertainties
- Simulation reweighted to account for dynamics of Λ_b decay



Lifetime of Λ_b

- Signal yield determined in bins of decay time distribution for both $\Lambda_b \rightarrow J/\psi p K^-$ and $B^0 \rightarrow J/\psi \pi^+ K^-$
- Yield ratio fitted allowing for differences in acceptance as a function of decay time
- Lifetime ratio:
 $0.976 \pm 0.012 \pm 0.006$
- Using the world average B^0 lifetime (1.519 ± 0.007 ps):
 $\tau(\Lambda_b) = 1.482 \pm 0.018 \pm 0.012$ ps
- Central value of ratio differs from unity by only few %
 - Indeed, is less than 2σ from unity
 - In agreement with HQE

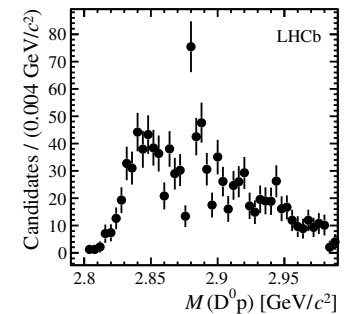
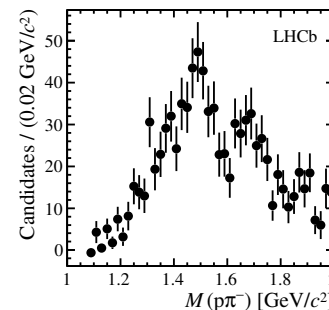
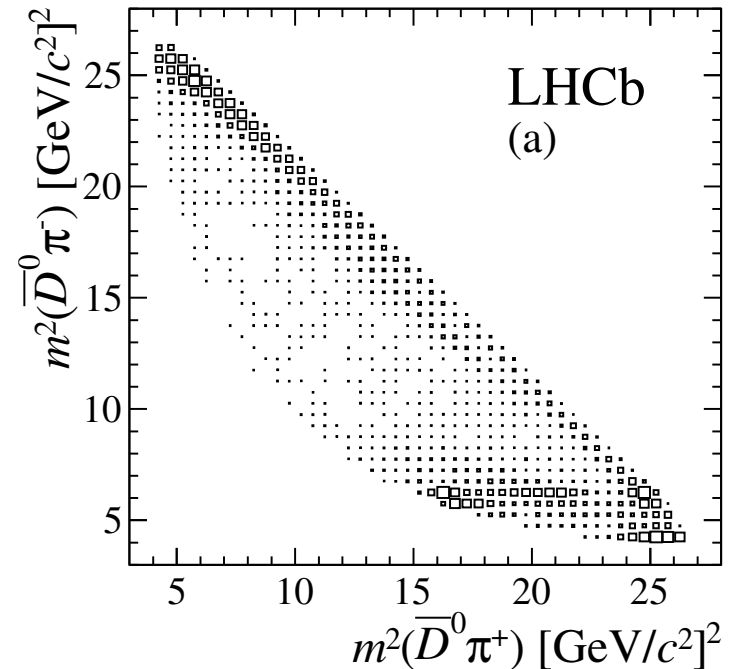


Conclusion

Future prospects

- Input on the f_0 mixing angle can also come from final states such as $D^0\pi\pi$
- Dalitz plot analyses of $D^{(*)}hh$ final states can also provide input on excited charm states, in particular determining their spin
- Charmless b-hadron decays can provide much input into light meson/baryon spectroscopy
 - In particular $B^+ \rightarrow 3h$ and $p\bar{p}h$
 - Understanding the large CP-violation will require precise understanding of Dalitz-plot structure
- Charm baryons can also be studied in B decays such as those to $Dp\pi$ and DpK
- B_c and b -baryon lifetime measurements in the pipeline

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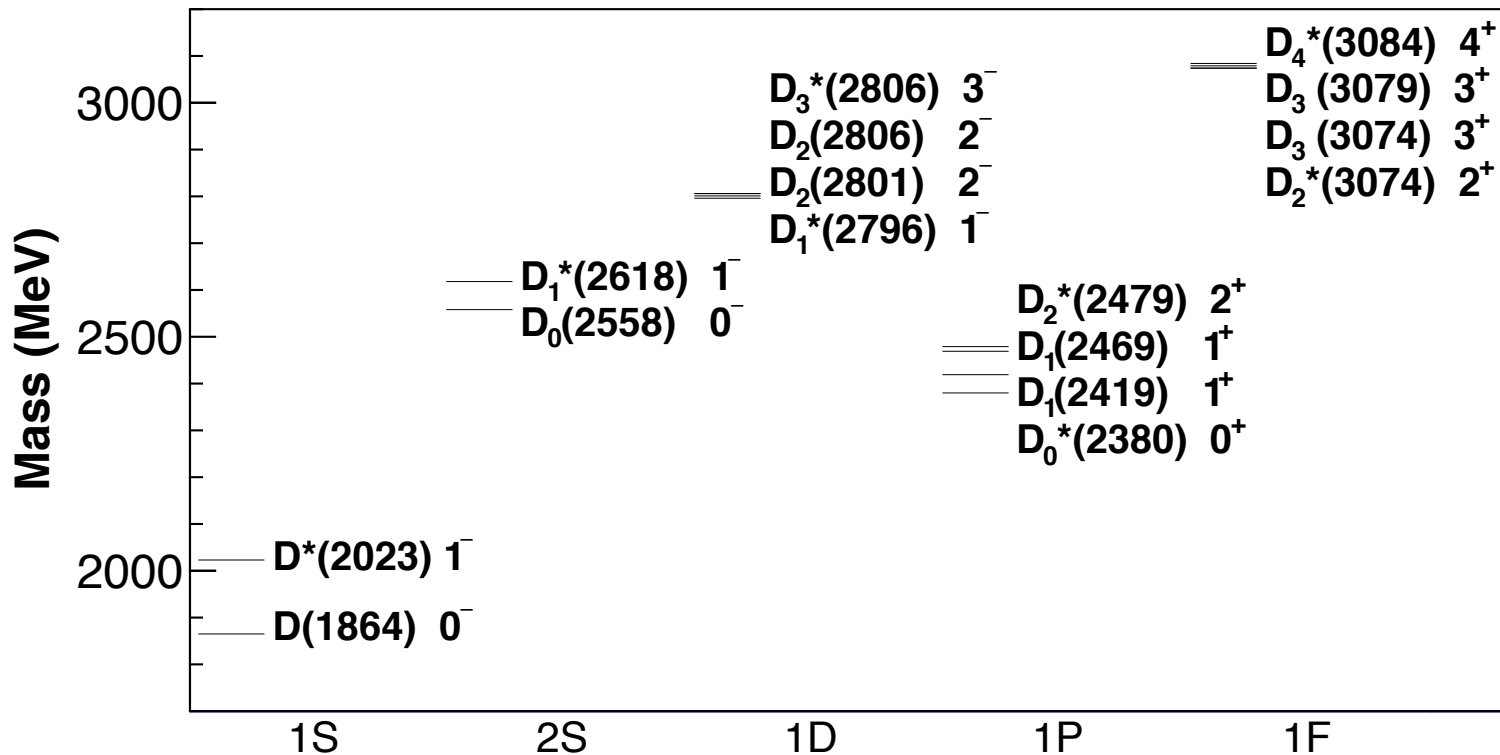


Summary

- A wealth of hadron spectroscopy results from LHCb in the last year
- Improved precision possible for some of the measurements using the combined 2011 + 2012 ($1 \text{ fb}^{-1} + 2 \text{ fb}^{-1}$) dataset
- Many more exciting results in the pipeline
- Watch this space!

Backup Slides

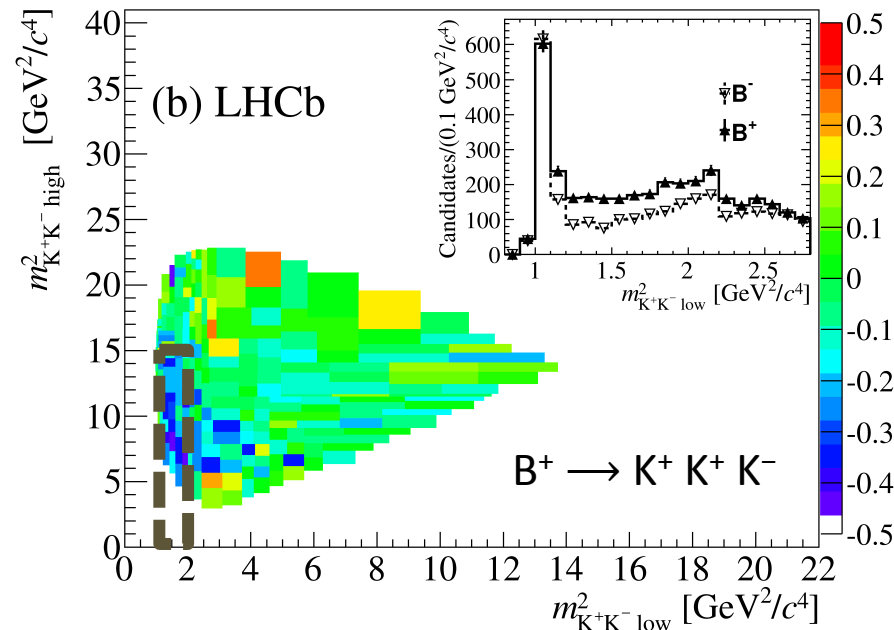
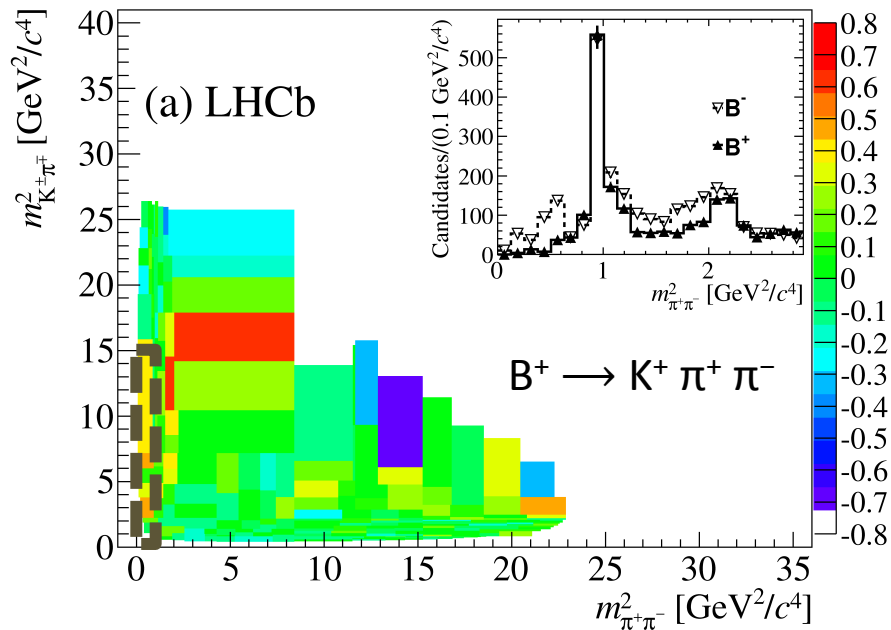
Excited D meson states



Modified Godfrey-Isgur mass predictions [[Phys. Rev. D32 \(1985\) 189](#)].

The spectrum is scaled such that the ground state coincides with the D^0 mass.

CPV in B⁺ → K⁺ h⁺ h⁻



$$A_{CP}(\text{local}) = 0.678 \pm 0.078 \pm 0.032 \pm 0.007$$

$$A_{CP}(\text{local}) = -0.226 \pm 0.020 \pm 0.004 \pm 0.007$$

- Study variation of A_{raw} over Dalitz plot
- Some areas of phase space have very large asymmetries, e.g. region around ρ^0 resonance in $B^+ \rightarrow K^+ \pi^+ \pi^-$ but also regions not clearly associated with a resonance, particularly in $B^+ \rightarrow K^+ K^+ K^-$

CPV in B⁺ → π⁺ h⁺ h⁻

- Again, asymmetries have opposite sign between the two modes
- Large localised asymmetries, not clearly associated with a narrow resonance

$$A_{CP}(\text{local}) = 0.584 \pm 0.082 \pm 0.027 \pm 0.007$$

$$A_{CP}(\text{local}) = -0.648 \pm 0.070 \pm 0.013 \pm 0.007$$

