

PROPERTIES OF STRONGLY COUPLED QCD FROM RENORMALIZATION FLOWS

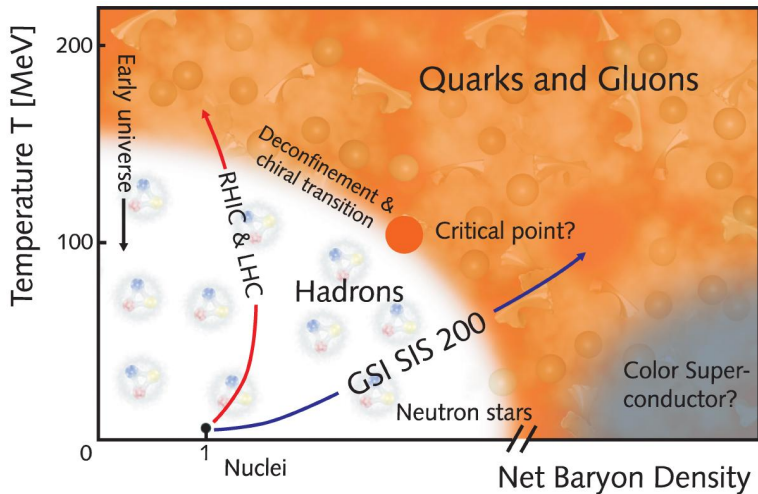
Holger Gies

Universität Heidelberg

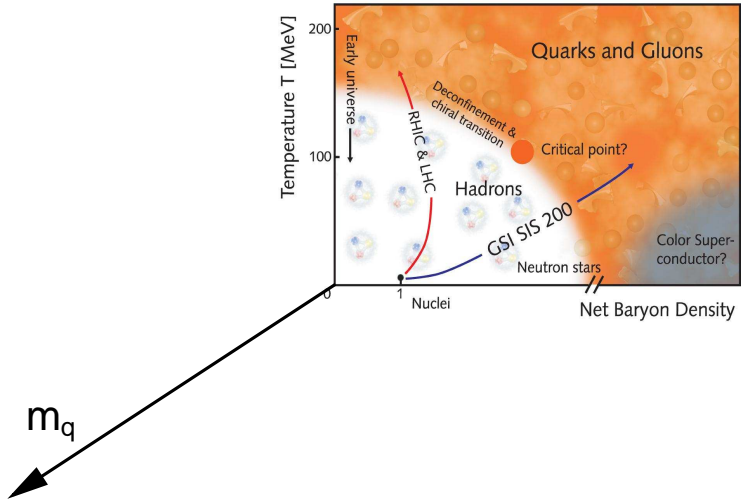


& J. Braun, J. Jaeckel, J.M. Pawłowski, C. Wetterich

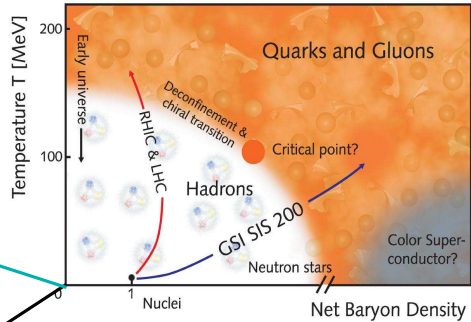
QCD PHASE DIAGRAM



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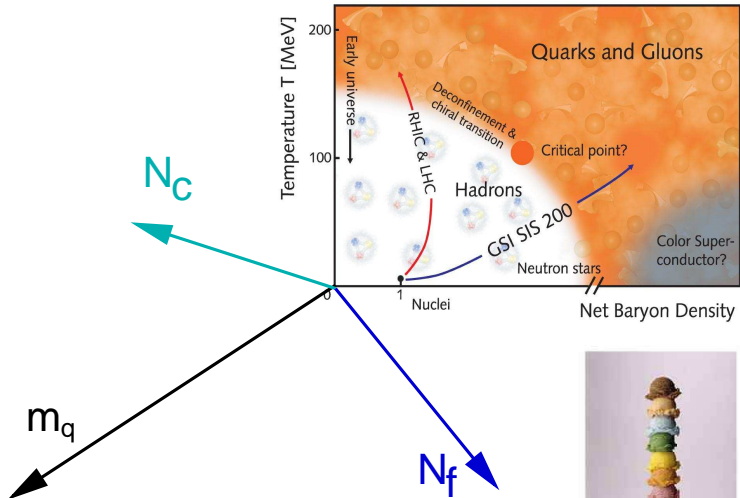
QCD PHASE DIAGRAM



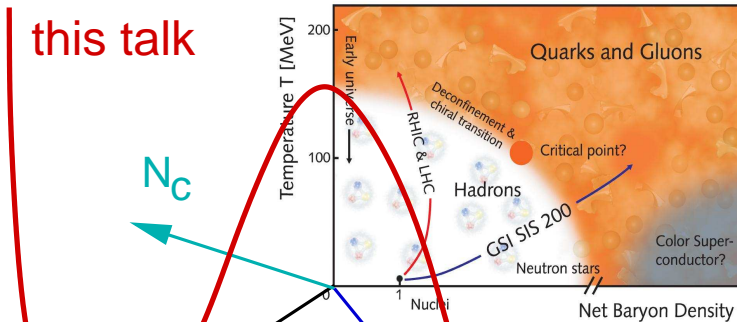
N_c

m_q

QCD PHASE DIAGRAM



QCD PHASE DIAGRAM



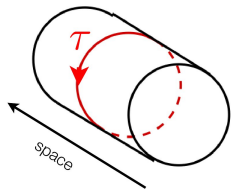
(DE-)CONFINEMENT IN GAUGE THEORIES

(DE-)CONFINEMENT ORDER PARAMETER

- ▷ order parameter: **Polyakov loop**

(POLYAKOV'78; SUSSKIND'79)

$$L(\mathbf{x}) = \frac{1}{N_c} \mathcal{P} \exp \left(i \int_0^\beta \tau A_0(\mathbf{x}) \right)$$



- ▷ heavy-quark free energy

$$\langle \text{tr}_F L(\mathbf{x}) \rangle \sim e^{-\beta \mathcal{F}}$$

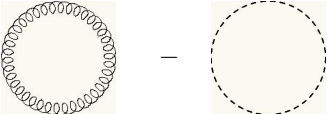
confinement: $\mathcal{F} \rightarrow \infty \iff \langle L \rangle = 0$ SYM Z_{N_c}

deconfinement: $\mathcal{F} < \infty \iff \langle L \rangle \neq 0$ SSB Z_{N_c}

[TALK BY F. MARHAUSER]

PERTURBATIVE ORDER-PARAMETER POTENTIAL

▷ e.g., perturbation theory in background-field gauge:

$$\begin{aligned} V_{\text{pert}}(A_0) &= \frac{1}{2} \text{Tr}_{\text{cLx}} \ln G_{\text{pert,gluon}}^{-1}[A_0] - \text{Tr}_{\text{cx}} \ln G_{\text{pert,ghost}}^{-1}[A_0] \\ &= \frac{1}{2} \text{---} \text{---} \text{---} - \text{---} \end{aligned}$$


▷ perturbative propagators

$$G_{\text{pert,gluon}}^{-1}, G_{\text{pert,ghost}}^{-1} \sim p^2$$

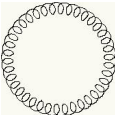
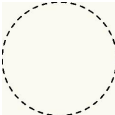
▷ A_0 background

$$\implies G_{\text{pert,gluon}}^{-1}(A_0), G_{\text{pert,ghost}}^{-1}(A_0) \sim -D^2[A_0]$$

PERTURBATIVE ORDER-PARAMETER POTENTIAL

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 V_{\text{pert}}(A_0) &= \frac{1}{2} \text{Tr}_{\text{cLx}} \ln G_{\text{pert,gluon}}^{-1}[A_0] - \text{Tr}_{\text{cx}} \ln G_{\text{pert,ghost}}^{-1}[A_0] \\
 &= \frac{1}{2} \left(\text{Diagram 1} - \text{Diagram 2} \right)
 \end{aligned}$$


-


▷ background field in Polyakov gauge

$$A_0 = A_0(\mathbf{x}) \in \text{Cartan}, \quad \text{e.g. for SU(2): } A_0^a = A_0 \delta^{a3}$$

▷ e.g., SU(2) order parameter

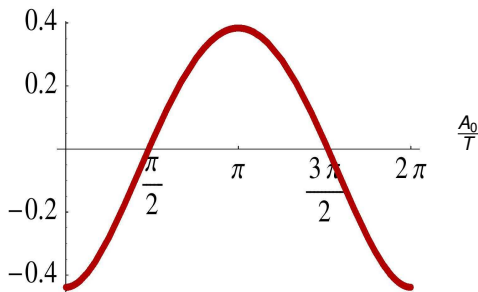
$$\text{tr}_F L(\mathbf{x}) = \cos \frac{A_0(\mathbf{x})}{2T} \quad L[\langle A_0 \rangle] \geq \langle L[A_0] \rangle \quad \text{confinement: } \left\langle \frac{A_0}{T} \right\rangle = \pi$$

PERTURBATIVE ORDER-PARAMETER POTENTIAL

▷ e.g., perturbation theory in background-field gauge for $A_0 = \text{const.}$

$$V_{\text{pert}}(A_0) = -\frac{4}{\pi^2} T^4 \sum_{n=1}^{\infty} \frac{\cos\left(2\pi n \frac{A_0}{2\pi T}\right)}{n^4}$$

(WEISS'81)

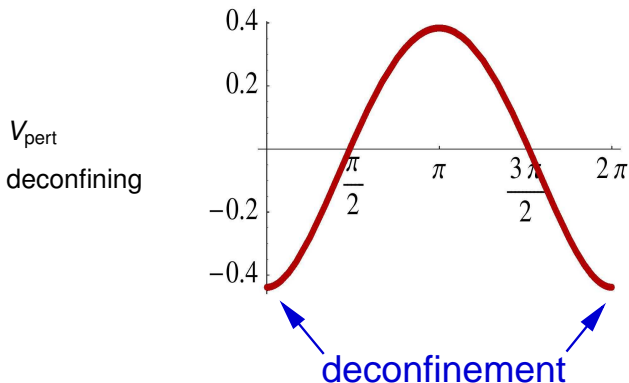


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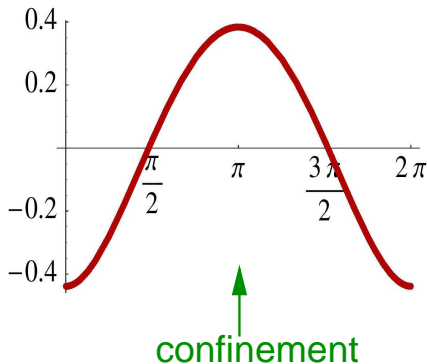


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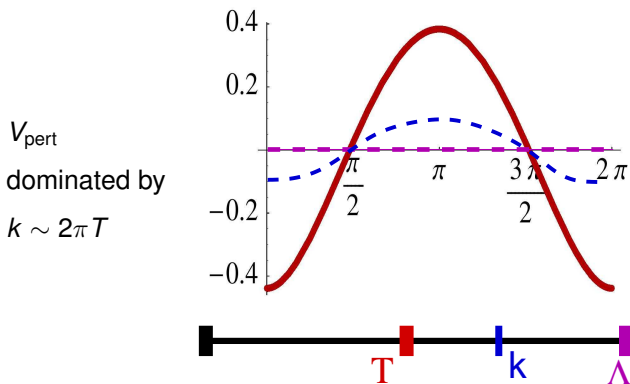


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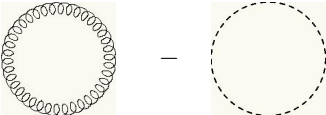
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PERTURBATIVE ORDER-PARAMETER POTENTIAL

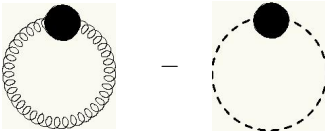
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PERTURBATIVE ORDER-PARAMETER POTENTIAL

▷ “improved” potential

$$V_{\text{impr}}(A_0) = \frac{1}{2} \text{Tr}_{\text{cLx}} \ln G_{\text{gluon}}^{-1}[A_0] - \text{Tr}_{\text{cx}} \ln G_{\text{ghost}}^{-1}[A_0]$$

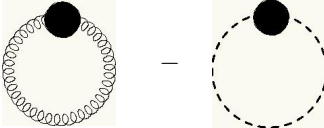
$$= \frac{1}{2} \left(\text{Diagram 1} \right) - \left(\text{Diagram 2} \right)$$


The diagram shows two terms in parentheses, separated by a minus sign. The first term is a loop of gluons (wavy lines) with a black dot on top. The second term is a loop of ghosts (dashed lines) with a black dot on top.

PERTURBATIVE ORDER-PARAMETER POTENTIAL

▷ “improved” potential

$$V_{\text{impr}}(A_0) = \frac{1}{2} \text{Tr}_{\text{cLx}} \ln G_{\text{gluon}}^{-1}[A_0] - \text{Tr}_{\text{cx}} \ln G_{\text{ghost}}^{-1}[A_0]$$

$$= \frac{1}{2} \left(\text{Diagram 1} \right) - \left(\text{Diagram 2} \right)$$


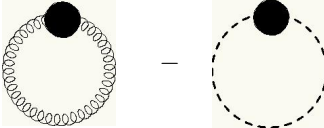
▷ $V_{\text{impr}}(A_0)$ dominated by modes $k \sim T$

UV: $G_{\text{pert,gluon}}^{-1}, G_{\text{pert,ghost}}^{-1} \sim p^2$

PERTURBATIVE ORDER-PARAMETER POTENTIAL

▷ “improved” potential

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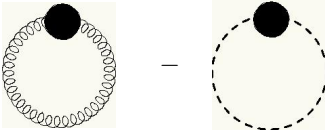
▷ $V_{\text{impr}}(A_0)$ dominated by modes $k \sim T$

IR ($k \ll \Lambda_{\text{QCD}}$): $G_{\text{gluon}}^{-1} \sim (p^2)^{1+\kappa_A}$, $G_{\text{ghost}}^{-1} \sim (p^2)^{1+\kappa_C}$

PERTURBATIVE ORDER-PARAMETER POTENTIAL

▷ “improved” potential

$$V_{\text{impr}}(A_0) = \frac{1}{2} \text{Tr}_{\text{cLx}} \ln G_{\text{gluon}}^{-1}[A_0] - \text{Tr}_{\text{cx}} \ln G_{\text{ghost}}^{-1}[A_0]$$

$$= \frac{1}{2} \left(\text{Diagram 1} \right) - \left(\text{Diagram 2} \right)$$


▷ $V_{\text{impr}}(A_0)$ dominated by modes $k \sim T$

$$\text{IR } (k \ll \Lambda_{\text{QCD}}) : G_{\text{gluon}}^{-1} \sim (p^2)^{1+\kappa_A}, \quad G_{\text{ghost}}^{-1} \sim (p^2)^{1+\kappa_C}$$

▷ A_0 background

$$\ln(-D^2[A_0])^{1+\kappa} = (1 + \kappa) \ln(-D^2[A_0])$$

ORDER-PARAMETER POTENTIAL

▷ low-energy effective potential

(BRAUN,HG,PAWLOWSKI'07)

$$V_{\text{IR}}(A_0) \simeq \left\{ \frac{d-1}{2}(1 + \kappa_A) + \frac{1}{2} - (1 + \kappa_C) \right\} \frac{1}{\Omega} \text{Tr} \ln (-D^2[A_0])$$

ORDER-PARAMETER POTENTIAL

- ▶ low-energy effective potential

(BRAUN, HG, PAWLOWSKI'07)

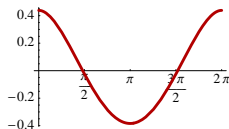
$$V_{\text{IR}}(A_0) = \left\{ \underbrace{\frac{d-1}{2}(1+\kappa_A)}_{\text{transv. gluons}} + \underbrace{\frac{1}{2}}_{\text{long. gluons}} \underbrace{-(1+\kappa_C)}_{\text{ghosts}} \right\} \underbrace{\frac{1}{\Omega} \text{Tr} \ln(-D^2[A_0])}_{\sim V_{\text{pert}}}$$

- ▶ confinement criterion (Landau gauge)

$$d - 2 + (d - 1)\kappa_A - 2\kappa_C < 0$$

- ▶ $d = 4$:

$$3\kappa_A - 2\kappa_C < -2$$



- ▶ quark confinement induced by:

IR gluon suppression and/or ghost enhancement

CONFINEMENT CRITERION

(TAYLOR'71)

(ZWANZIGER'02; LERCHE, VON SMEKAL'02)

(SCHLEIFENBAUM, MAAS, WAMBACH, ALKOEFER'05)

(CUCCHIERI, MAAS, MENDES'08)

▷ Landau-gauge sum rule

$$0 = \kappa_A + 2\kappa_C - \frac{d-4}{2}$$

▷ quark confinement

(BRAUN, HG, PAWLOWSKI'07)

$$\kappa \equiv \kappa_C > \frac{d-3}{4}$$

▷ $d = 4$:

$$\implies \kappa > \frac{1}{4}$$

CONFINEMENT CRITERION ($d = 4$)

- ▶ quark confinement

(BRAUN,HG,PAWLOWSKI'07)

$$\kappa > \frac{1}{4}$$

- ▶ Kugo-Ojima color confinement

(KUGO,OJIMA'79)

$$\kappa > 0$$

- ▶ Gribov-Zwanziger color confinement

(GRIBOV'78, ZWANZIGER'94,'04)

$$\kappa > \frac{1}{2} \quad (\text{horizon condition})$$

FUNCTIONAL RG

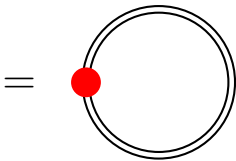
FUNCTIONAL RG FLOW EQUATION

IR: $k \rightarrow 0$



UV: $k \rightarrow \Lambda$

$$\partial_t \Gamma_k \equiv k \partial_k \Gamma_k = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)} + R_k} \partial_t R_k$$

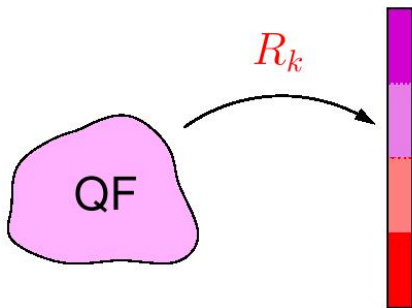


(WILSON'71; WEGNER&HOUGHTON'73; POLCHINSKI'84; WETTERICH'93)

FUNCTIONAL RG FLOW EQUATION

$$\partial_t \Gamma_k = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)} + R_k} \partial_t R_k$$

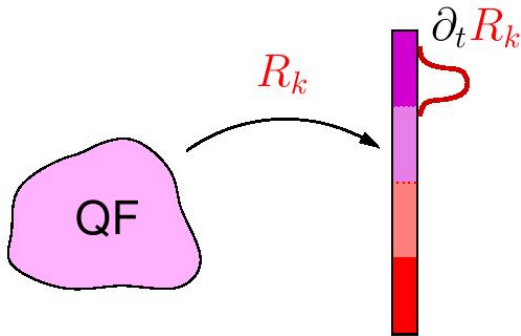
▷ quantum fluctuations:



FUNCTIONAL RG FLOW EQUATION

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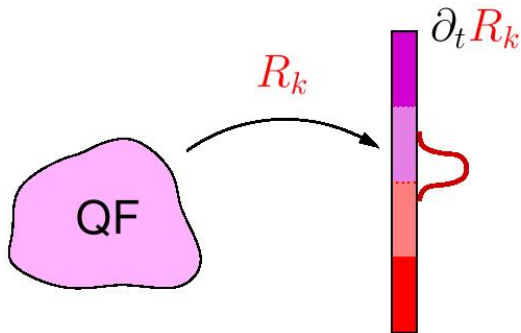
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FUNCTIONAL RG FLOW EQUATION

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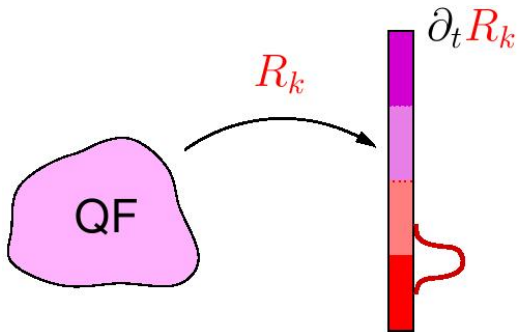
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FUNCTIONAL RG FLOW EQUATION

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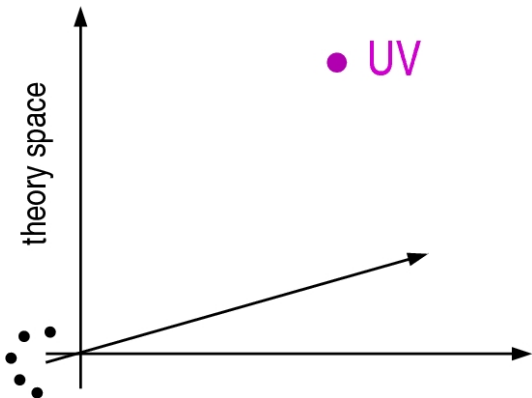
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FUNCTIONAL RG FLOW EQUATION

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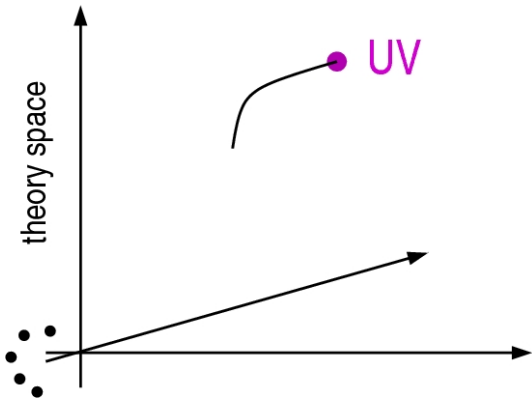
▷ RG trajectory: $\Gamma_{k=\Lambda} = S_{\text{bare}} = \int \frac{1}{4} F_{\mu\nu}^z F_{\mu\nu}^z + \bar{\psi} (i\not{\partial} + gA) \psi$



FUNCTIONAL RG FLOW EQUATION

$$\partial_t \Gamma_k = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)} + R_k} \partial_t R_k$$

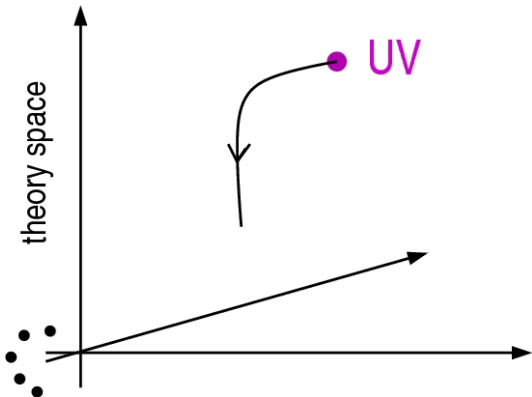
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FUNCTIONAL RG FLOW EQUATION

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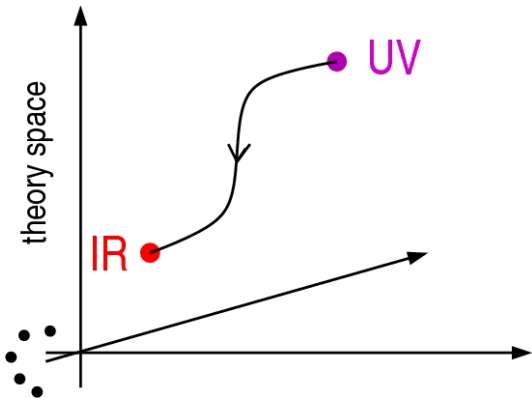


FUNCTIONAL RG FLOW EQUATION

$$\partial_t \Gamma_k = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)} + R_k} \partial_t R_k$$

▷ RG trajectory:

$$\Gamma_{k \rightarrow 0} = \Gamma$$

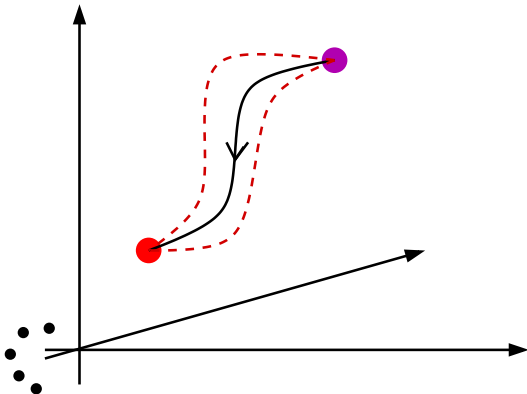


FUNCTIONAL RG FLOW EQUATION

$$\partial_t \Gamma_k = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)} + R_k} \partial_t R_k$$

▷ RG trajectory:

R_k scheme independence

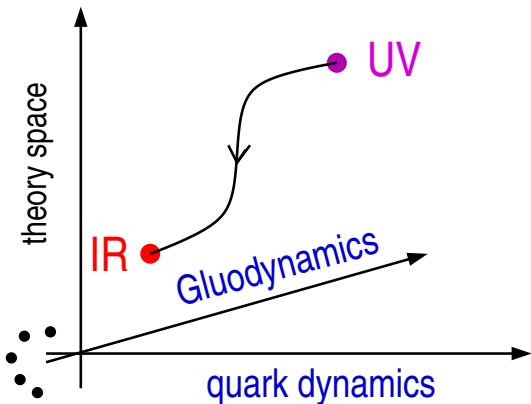


FUNCTIONAL RG FLOW EQUATION

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▷ RG trajectory:

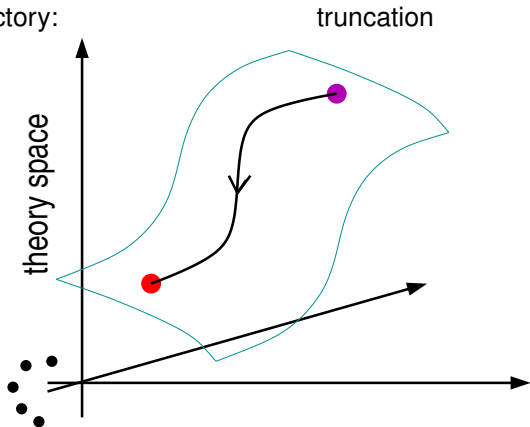
truncation



FUNCTIONAL RG FLOW EQUATION

$$\partial_t \Gamma_k = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)} + R_k} \partial_t R_k$$

▷ RG trajectory:



ORDER PARAMETER POTENTIAL FROM FUNCTIONAL RG

▷ flow equation

$$\partial_t \Gamma_k = \frac{1}{2} \text{Tr} \partial_t R_k [\Gamma_k^{(2)} + R_k]^{-1}$$

$$\Rightarrow \Gamma = \frac{1}{2} \text{Tr} \ln \Gamma^{(2)} - \frac{1}{2} \int_0^\infty \frac{dk}{k} \text{Tr} \frac{1}{\Gamma_k^{(2)} + R_k} \partial_t \Gamma_k^{(2)} + \text{c.t.}$$

▷ A_0 potential:

$$\Gamma[A_0] = \int d^d x V(A_0) + Z(A_0) \partial_\mu A_0 \partial_\mu A_0 \dots$$

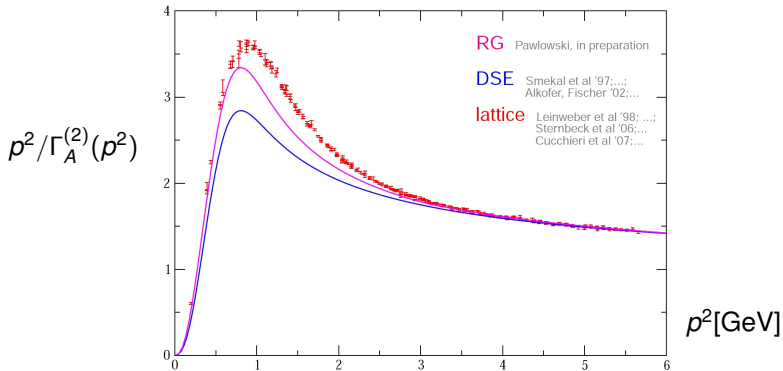
▷ $G \cdot \Gamma^{(2)} \equiv \mathbb{1}$

$$V(A_0) = \frac{1}{2\Omega} \text{Tr} \ln G^{-1} + \mathcal{O}(\partial_t \Gamma_k^{(2)})$$

DECONFINEMENT PHASE TRANSITION

▷ INPUT:

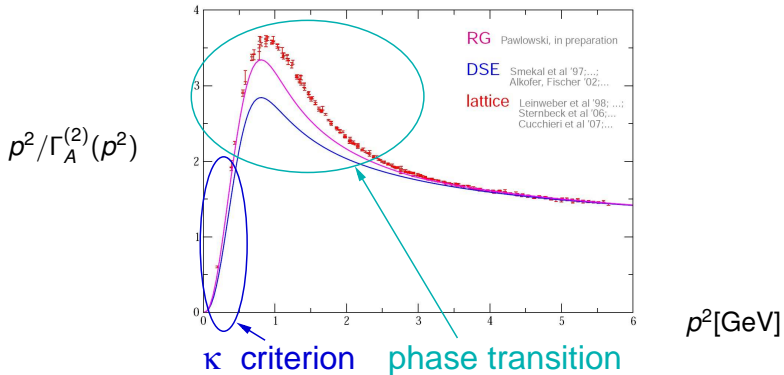
$$\text{Landau-gauge } \Gamma^{(2)} \equiv G^{-1} \Big|_{T=0}$$



DECONFINEMENT PHASE TRANSITION

▷ INPUT:

$$\text{Landau-gauge } \Gamma^{(2)} \equiv G^{-1} \Big|_{T=0}$$

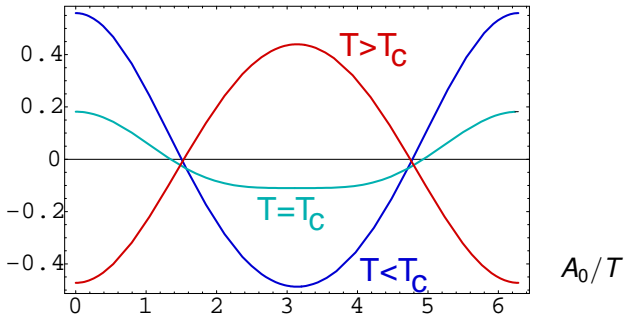


⇒ mid-momentum regime is decisive

DECONFINEMENT PHASE TRANSITION

▷ SU(2) A_0 potential

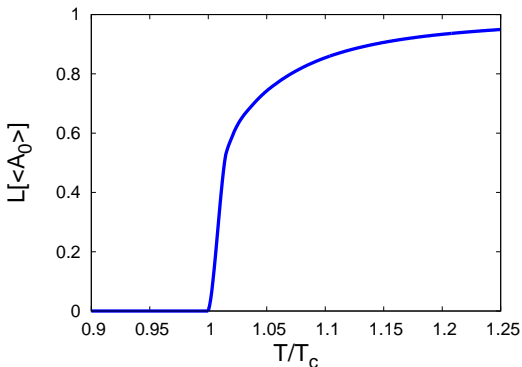
(BRAUN, HG, PAWLOWSKI'07)



DECONFINEMENT PHASE TRANSITION

▷ SU(2): 2nd order phase transition

(BRAUN, HG, PAWLOWSKI '07)



$$T_c/\sqrt{\sigma} = 0.614 \pm 0.023,$$

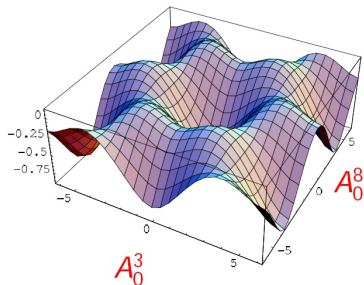
$$\text{cf. lattice: } T_c/\sqrt{\sigma} \simeq 0.709$$

(KACZMAREK ET AL. '02)

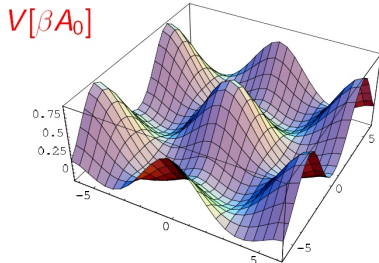
DECONFINEMENT PHASE TRANSITION

▷ SU(3) A_0 potential

(BRAUN, HG, PAWLOWSKI'07)



$T > T_c$

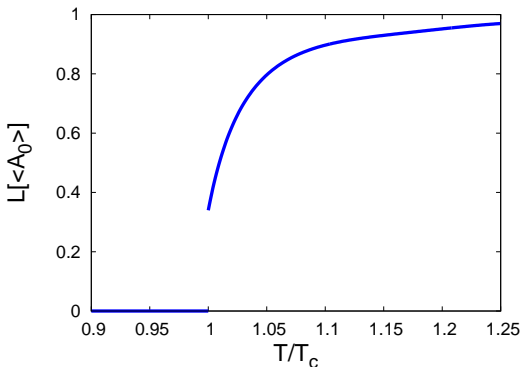


$T < T_c$

DECONFINEMENT PHASE TRANSITION

▷ SU(3): 1st order phase transition

(BRAUN, HG, PAWLOWSKI'07)



$$T_c/\sqrt{\sigma} = 0.646 \pm 0.023 \quad \Rightarrow \quad T_c \simeq 284\text{MeV}, \quad \text{cf. Lattice: } T_c/\sqrt{\sigma} \simeq 0.646$$

(KACZMAREK ET AL.'02)

ERROR ESTIMATE

$$\Gamma = \frac{1}{2} \text{Tr} \ln \Gamma^{(2)} - \frac{1}{2} \int_0^\infty \frac{dk}{k} \text{Tr} \frac{1}{\Gamma_k^{(2)} + R_k} \partial_t \Gamma_k^{(2)}$$

lattice (mid momentum!)
@ T=0

RG flow

▷ A_0 fluctuations neglected: $\Gamma_k^{(2)} \sim -\partial^2 + V_k''(A_0)$

relevant
less relevant

SU(2): Ising universality class

SU(3): 1st order

RG Flow towards the Chiral Transition

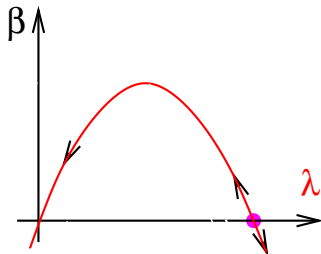
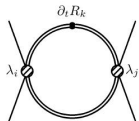
RG FLOW OF THE CHIRAL SECTOR

▷ effective action:

$$\Gamma_k = \int \frac{1}{2} \frac{\lambda_\sigma}{k^2} [(\bar{\psi}^a \psi^b)^2 - (\bar{\psi}^a \gamma_5 \psi^b)^2]$$

▷ RG flow

$$\partial_t \lambda_\sigma = 2\lambda_\sigma - \frac{N_c}{4\pi^2} \lambda_\sigma^2$$



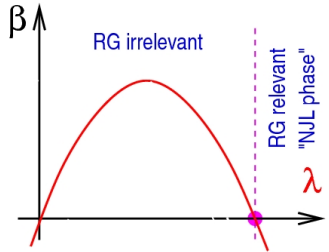
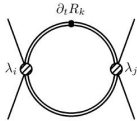
RG FLOW OF THE CHIRAL SECTOR

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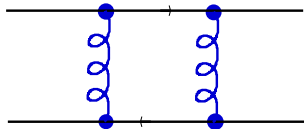
RG FLOW OF THE CHIRAL SECTOR

▷ effective action:

$$\Gamma_k = \int \frac{1}{4} F_{\mu\nu}^z F_{\mu\nu}^z + \dots + \bar{\psi} (i\cancel{\partial} + \bar{g}A) \psi + \frac{1}{2} \frac{\lambda_\sigma}{k^2} [(\bar{\psi}^a \psi^b)^2 - (\bar{\psi}^a \gamma_5 \psi^b)^2]$$

▷ RG flow

$$\begin{aligned} \partial_t \lambda_\sigma &= 2\lambda_\sigma - \frac{N_c}{4\pi^2} \lambda_\sigma^2 \\ &\quad - \frac{3}{8\pi^2} \frac{N_c^2 - 1}{N_c} g^2 \lambda_\sigma \\ &\quad - \frac{9}{256\pi^2} \frac{3N_c^2 - 8}{N_c} g^4 \end{aligned}$$



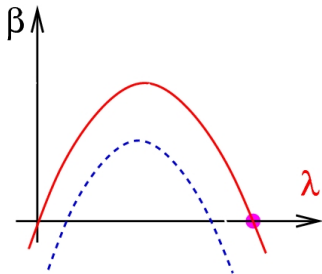
RG FLOW OF THE CHIRAL SECTOR

▷ effective action:

$$\Gamma_k = \int \frac{1}{4} F_{\mu\nu}^z F_{\mu\nu}^z + \dots + \bar{\psi} (i\cancel{\partial} + \bar{g}A) \psi + \frac{1}{2} \frac{\lambda_\sigma}{k^2} [(\bar{\psi}^a \psi^b)^2 - (\bar{\psi}^a \gamma_5 \psi^b)^2]$$

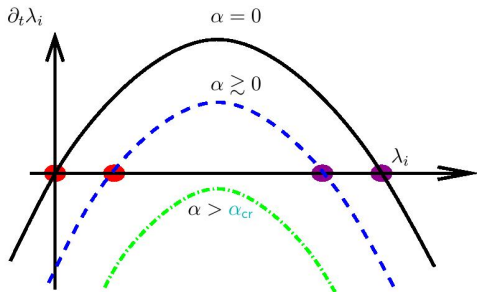
▷ RG flow

$$\partial_t \lambda_\sigma = 2\lambda_\sigma - \frac{N_c}{4\pi^2} \lambda_\sigma^2 - \frac{3}{8\pi^2} \frac{N_c^2 - 1}{N_c} g^2 \lambda_\sigma - \frac{9}{256\pi^2} \frac{3N_c^2 - 8}{N_c} g^4$$



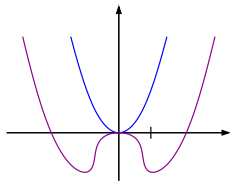
CHIRAL CRITICALITY

▷ critical gauge coupling α_{cr} :



⇒ bosonization $\rightarrow \chi$ SB:

$$\text{if } \alpha > \alpha_{cr} : \quad \lambda \sim \frac{1}{m_\phi^2} \rightarrow \infty$$



[TALK BY B.-J. SCHAEFER]

RG FLOW OF THE CHIRAL SECTOR

▷ effective action: $SU(N_c)$, $SU(N_f)_L \times SU(N_f)_R$

$$\begin{aligned} \Gamma_k = & \int \frac{Z_F}{4} F_{\mu\nu}^z F_{\mu\nu}^z + \dots + \bar{\psi} (iZ_\psi \not{\partial} + Z_1 \bar{g} A) \psi \\ & + \frac{1}{2} \frac{\lambda_\sigma}{k^2} (\text{S-P}) + \frac{1}{2} \frac{\lambda_{VA}}{k^2} [2(\text{V-A})^{\text{adj.}} + (1/N_c)(\text{V-A})] \\ & + \frac{1}{2} \frac{\lambda_+}{k^2} (\text{V+A}) + \frac{1}{2} \frac{\lambda_-}{k^2} (\text{V-A}) \end{aligned}$$

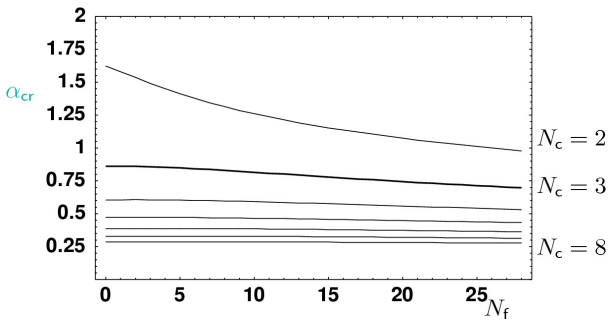
▷ RG flow, e.g.,

$$\begin{aligned} \partial_t \lambda_\sigma = & 2\lambda_\sigma - \frac{1}{4\pi^2} I_1^{(F)} [R_k] \left\{ 2N_c \lambda_\sigma^2 - 2\lambda_- \lambda_\sigma - 2N_f \lambda_\sigma \lambda_{VA} - 6\lambda_+ \lambda_\sigma \right\} \\ & - \frac{1}{8\pi^2} I_{1,1}^{(FB)} [R_k] \left[3 \frac{N_c^2 - 1}{N_c} g^2 \lambda_\sigma - 6g^2 \lambda_+ \right] \\ & - \frac{3}{128\pi^2} I_{1,2}^{(FB)} [R_k] \frac{3N_c^2 - 8}{N_c} g^4 \end{aligned}$$

(HG, JAECKEL, WETTERICH'04)

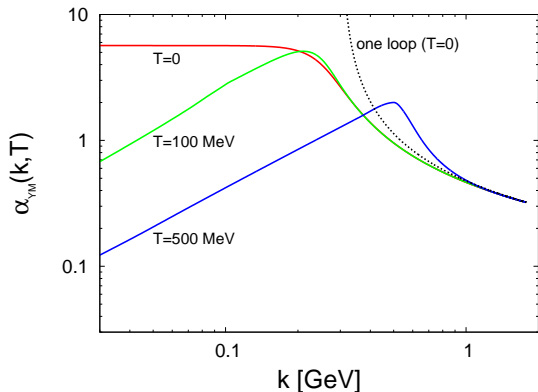
χ SB CRITICAL COUPLING

(HG, JAECKEL'05)



e.g., for $N_c = 3 = N_f$: $\alpha_{cr} \simeq 0.85$

RUNNING GAUGE COUPLING AT FINITE T



Background gauge:

(HG'02)

(BRAUN, HG'05)

cf. Landau gauge:

(v.SMEKAL, ALKOFE, HAUCK'97)

(LEINWEBER ET AL'98)

(LERCHE, v.SMEKAL'02)

(FISCHER, ALKOFE'02)

(ZWANZIGER'02)

(PAWLOWSKI, LITIM, NEDELKO, v.SMEKAL'03)

(FISCHER, HG'04)

(OLIVEIRA, SILVA'04)

(BLOCH, CUCCHIERI, LANGFELD, MENDES'04)

(STERNBECK ET AL.'06)

(MAAS'07)

(CUCCHIERI, MENDES, OLIVEIRA, SILVA'07)

▷ $T/k \rightarrow \infty$: strongly interacting 3D theory

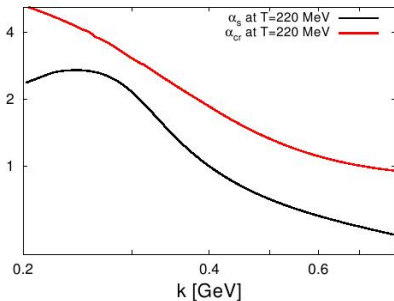
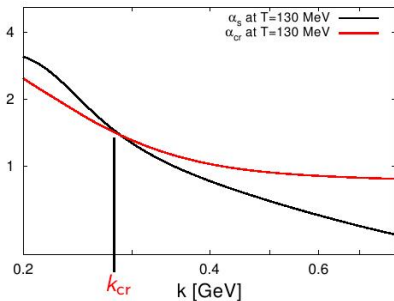
$$\alpha \rightarrow \frac{k}{T} \alpha_{3D}, \quad \alpha_{3D} \rightarrow \alpha_{3D,*} \simeq 2.7$$

cf. lattice: (CUCCHIERI, MAAS, MENDES'07)

CHIRAL PHASE TRANSITION



$\alpha(k, T)$ vs. $\alpha_{cr}(T/k)$



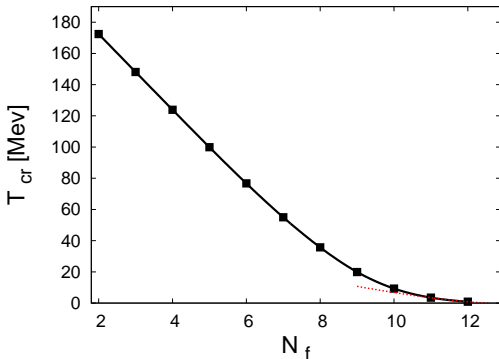
⇒ χ SB triggered by α_s

single input: $\alpha_s(m_\tau) = 0.322$

T_c [MeV]	RG (BRAUN, HG'05)
$N_f=2$	172 ± 37
$N_f=3$	148 ± 32

T_c [MeV]	Lattice (BI) (CHEN ET AL.'06)	Lattice (W) (AOKI ET AL.'06)
$N_f=2+1$	$192(7)(4)$	$151(3)(3)$

CHIRAL PHASE BOUNDARY $T - N_F$

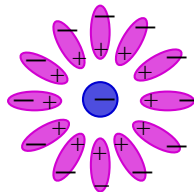


(BRAUN, HG'05,'06)

▷ small N_f : fermionic screening, $\beta_{\text{quark}} \simeq \frac{2}{3} N_f \frac{g^4}{8\pi^2}$

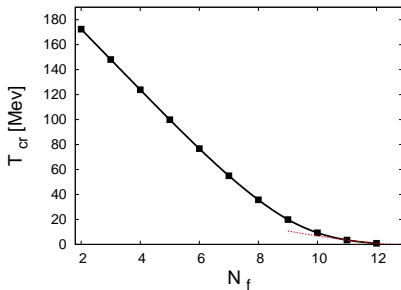
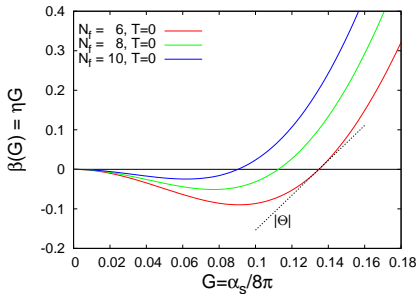
▷ critical flavor number:

$$N_f^{\text{cr}} \simeq 12$$



(CF. APPELQUIST ET AL.'96; MIRANSKI, YAMAWAKI'96; HG, JAECKEL'05)

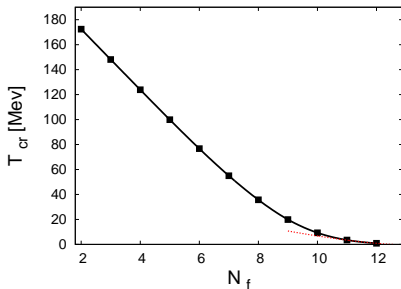
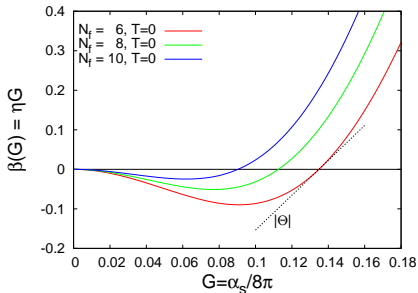
CHIRAL PHASE BOUNDARY $T - N_F$



▷ fixed-point regime: **critical exponent Θ**

$$\beta_{g^2} \simeq -\Theta (g^2 - g_*^2)$$

CHIRAL PHASE BOUNDARY $T - N_F$



- ▷ fixed-point regime: **critical exponent Θ**

$$\beta_{g^2} \simeq -\Theta (g^2 - g_*^2)$$

- ▷ shape of the phase boundary for $N_f \simeq N_f^{cr}$:

(BRAUN, HG'05,'06)

$$T_{cr} \sim k_0 |N_f - N_f^{cr}|^{\frac{1}{|\Theta|}}, \quad \Theta \simeq -0.71$$

CONCLUSIONS

- ▷ A_0 potential: quark confinement from color confinement

$$3\kappa_A - 2\kappa_C < -2$$

quark confinement from **IR** gluon suppression / ghost enhancement

- ▷ functional RG for $\Gamma[\phi]$

- **systematic** and **consistent** expansion schemes for QCD
- chiral symmetry ✓
- calculations “from **first principles**”

- ▷ Many-flavor QCD: relation among universal aspects:

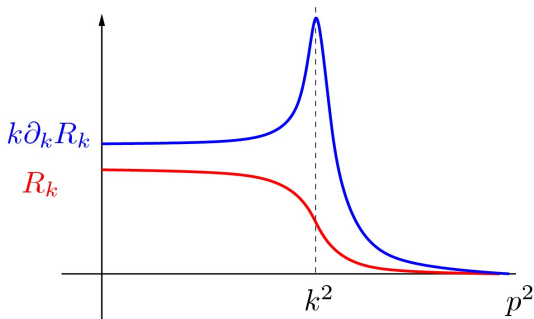
shape of the phase boundary \iff **IR** critical exponent

Appendix

FUNCTIONAL RG FLOW EQUATION

$$\partial_t \Gamma_k = \frac{1}{2} \text{Tr} \frac{1}{\Gamma_k^{(2)} + R_k} \partial_t R_k$$

▷ regulator

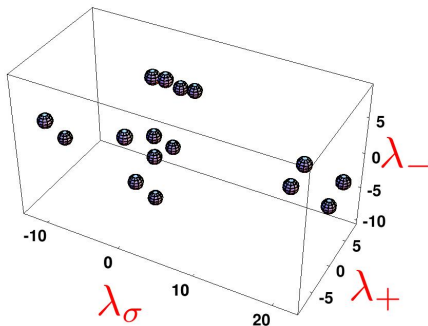


RG FLOW OF THE CHIRAL SECTOR

▷ 2 fixed points per $\partial_t \lambda$

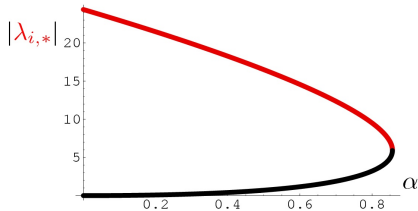
⇒ $2^4 = 16$ fixed points

▷ in general: 2^n FP's
for $n = \#$ of λ 's



▷ fixed-point annihilation

e.g., $N_c = N_f = 3$



CHIRAL CRITICALITY AT FINITE TEMPERATURE

▷ quark modes:

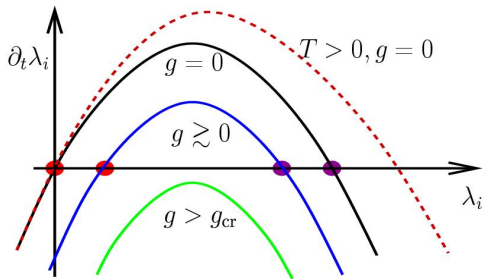
$$m_T^2 = m_f^2 + (2\pi T(n + \frac{1}{2}))^2$$

⇒ T -dependent

critical coupling:

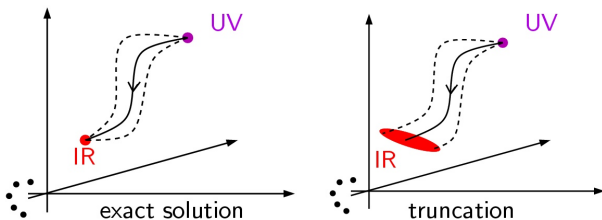
$$\alpha_{\text{cr}}(T) \gtrsim \alpha_{\text{cr}} \simeq 0.85$$

(BRAUN, HG'05)



ERROR ESTIMATE

- ▷ regulator dependence



- ▷ fermion sector: “optimized” regulator vs. “sharp cutoff”

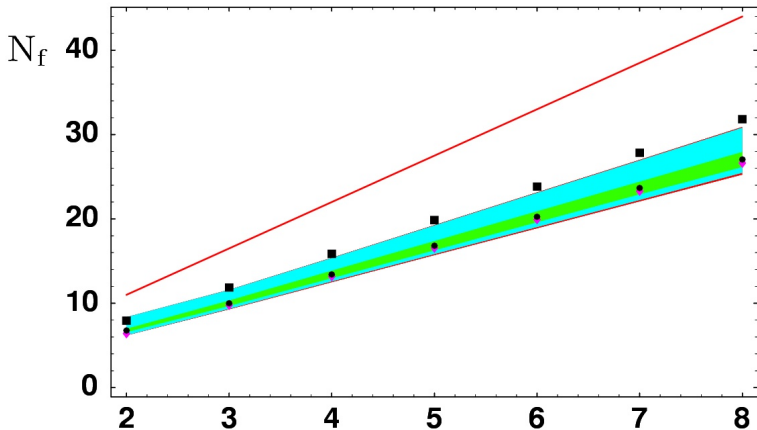
(LITIM'01)

$$l_1^{(F),4} = \frac{1}{2}, \quad l_{1,1}^{(FB),4} = 1, \quad l_{1,2}^{(FB),4} = \frac{3}{2} \quad \text{vs.} \quad l_1^{(F),4} = l_{1,1}^{(FB),4} = l_{1,2}^{(FB),4} = 1$$

- ▷ anomalous dimensions, momentum dependencies, higher-order operators $\sim \psi^8$, etc. ...
- ▷ gauge sector: 2-loop, 3-loop, 4-loop β function

$\overline{\text{MS}}$ scheme vs. RG scheme ($\sim 10, 30, 50$ % variation (?))

χ SB CRITICAL COUPLING



▷ SU(3) “conformal phase” for

N_c

$$N_{f,cr} = 10.0 \pm 0.29(\text{fermion}) \begin{matrix} +1.55 \\ -0.63 \end{matrix}(\text{gluon}) \lesssim N_f < 16.5$$

Lessons to be learned for “real QCD”

- fermionic screening is rather weak
- fermionic truncation (surprisingly) stable in χ symmetric phase
- phase boundary detectable with fermionic “derivative expansion”
- “real QCD” requires nonperturbative estimate of β_{g^2}