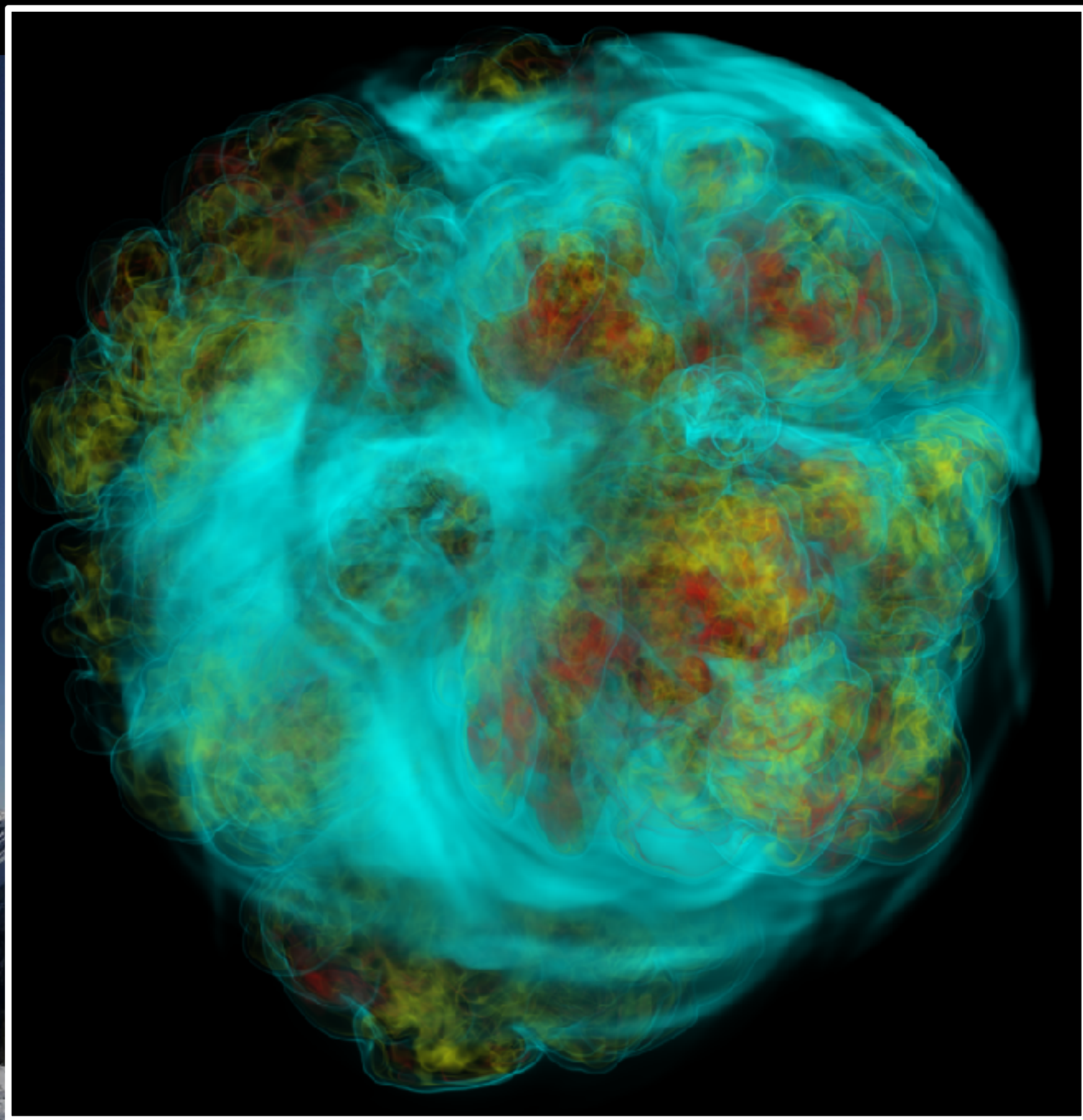
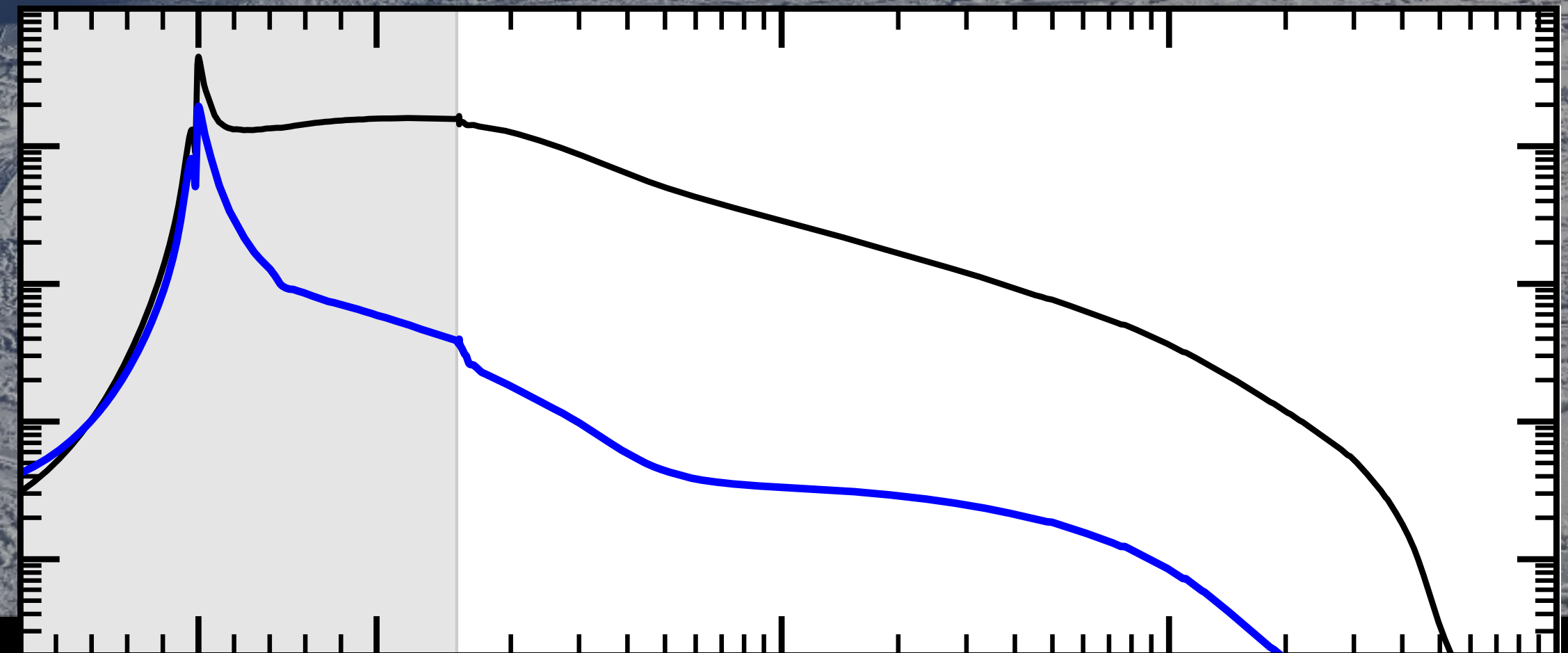


The impact of the nuclear EOS on neutrino emission from proto-neutron stars



Luke Roberts
Michigan State University



Overview

- Progenitor dependence of 3D models of CCSNe
- EOS sensitivity of CCSN explosion mechanism
- Some EOS sensitivities of long term CCSN neutrino emission

Core Collapse

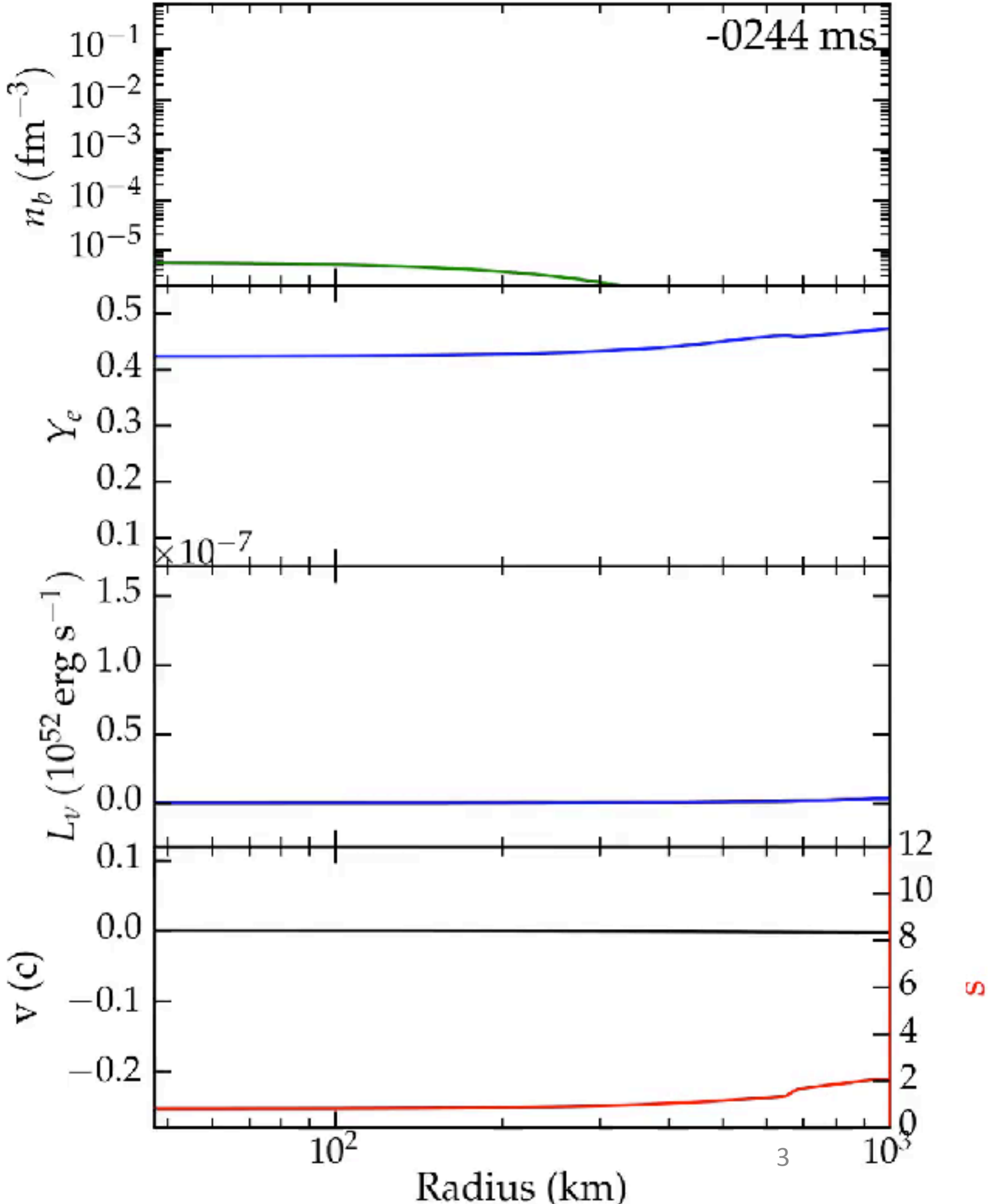
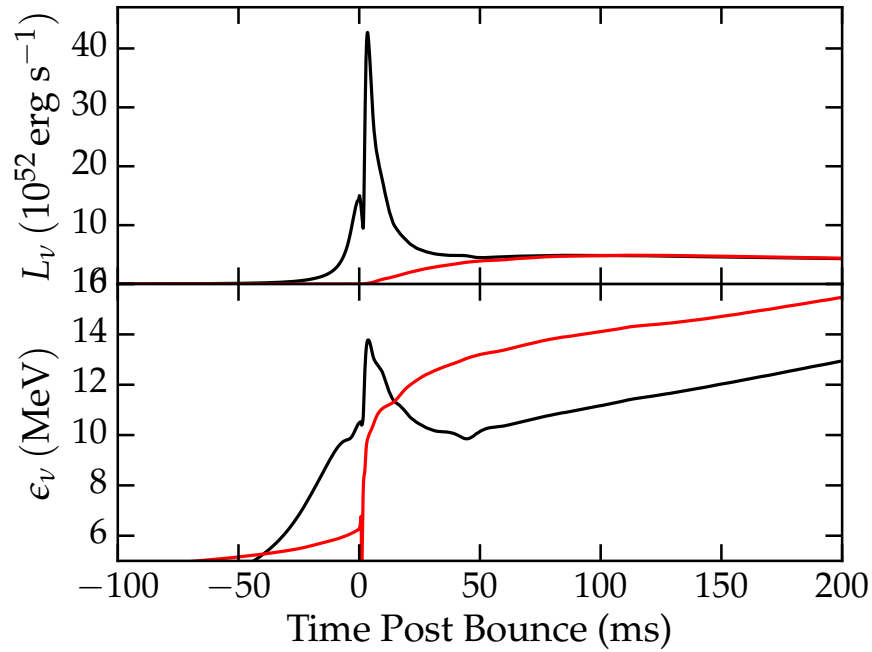
- Stars with $M > \sim 9 M_{\text{sun}}$ burn their core to Fe
- Core exceeds a Chandrasekhar mass supersonic collapse outside of homologous core \rightarrow bounce shock after $\sim 2 \times$ saturation density

- Gravitational binding energy of compact remnant:

$$\frac{GM_{NS}^2}{R_{NS}} \sim 3 \times 10^{53} \text{ erg}$$

- Binding energy of stellar envelope:

$$\sim 10^{51} \text{ erg}$$



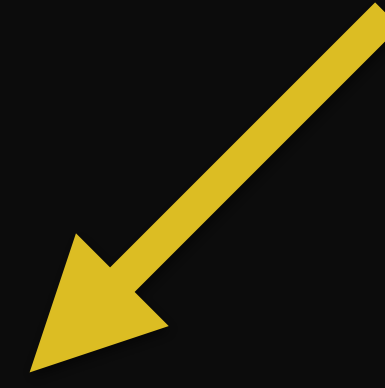
Simulating the Neutrino Mechanism

Hydrodynamics
+
General Relativity
+
Neutrino Transport
+
Microphysics
(EoS, ν -opacities)

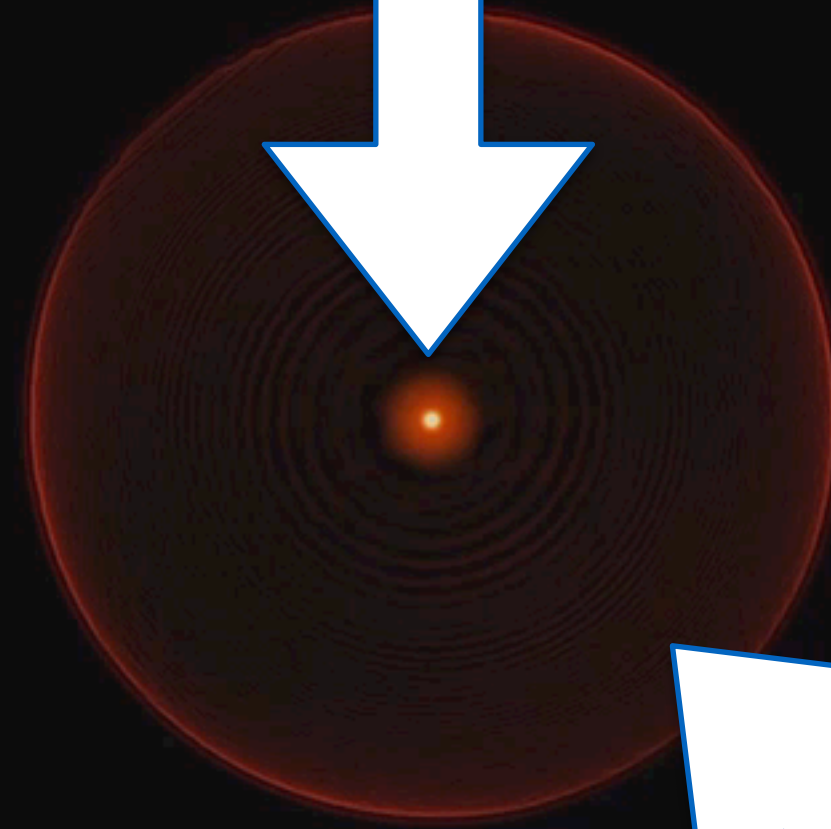
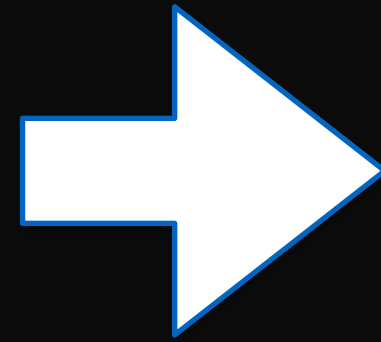
Entropy



Protoneutron Star



Shock



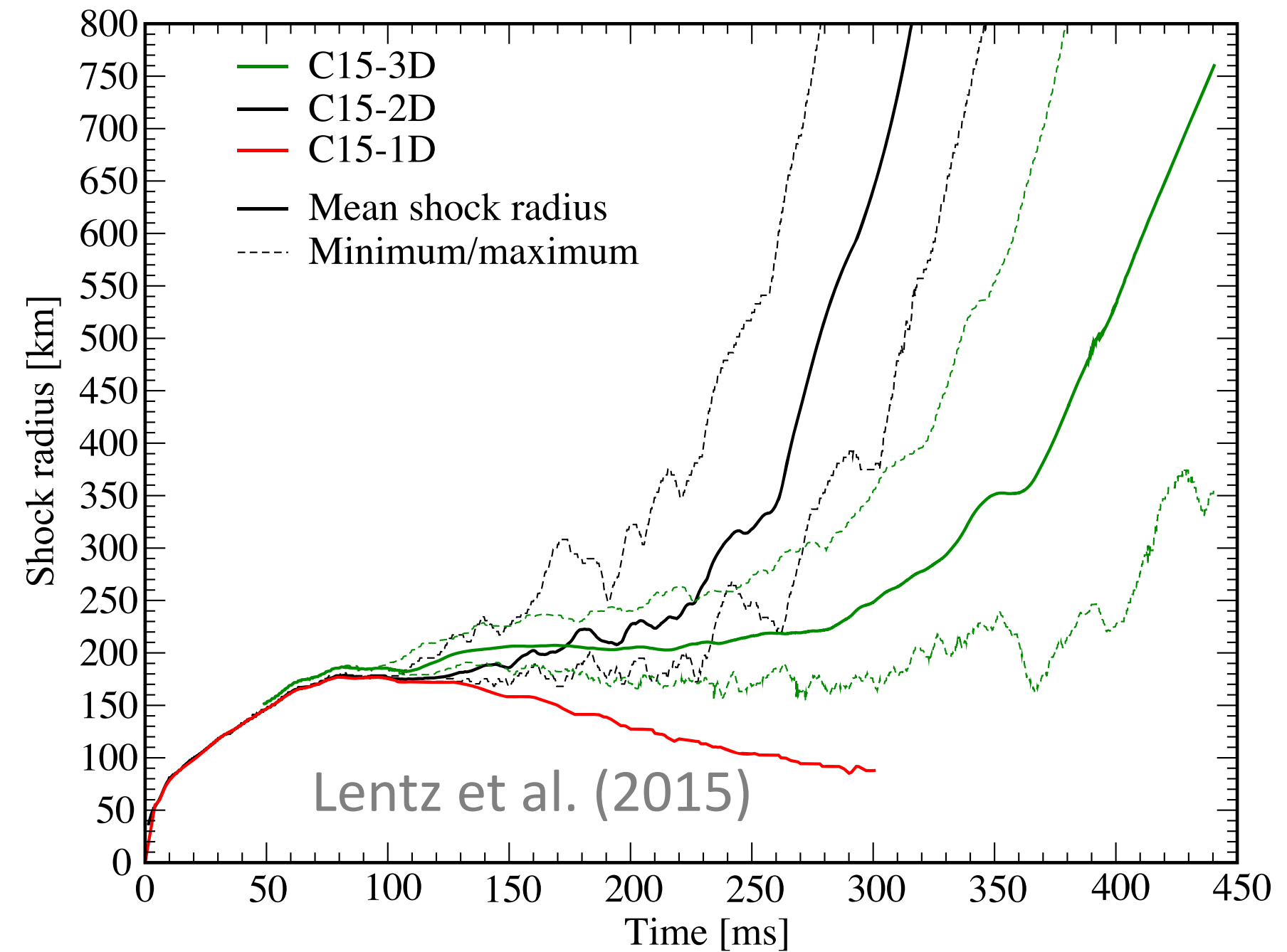
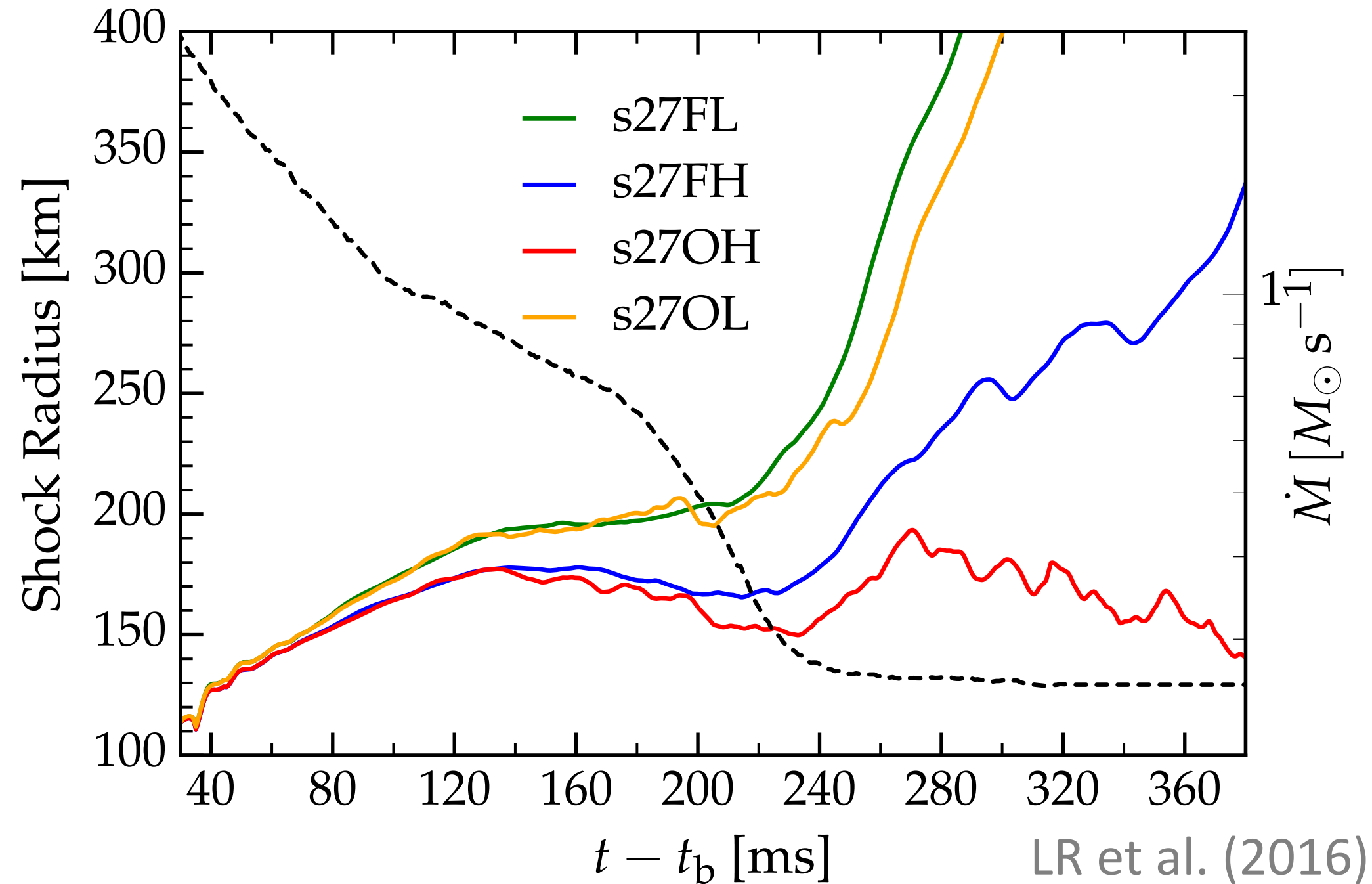
Gain Region



100 km

3D Explosion Models

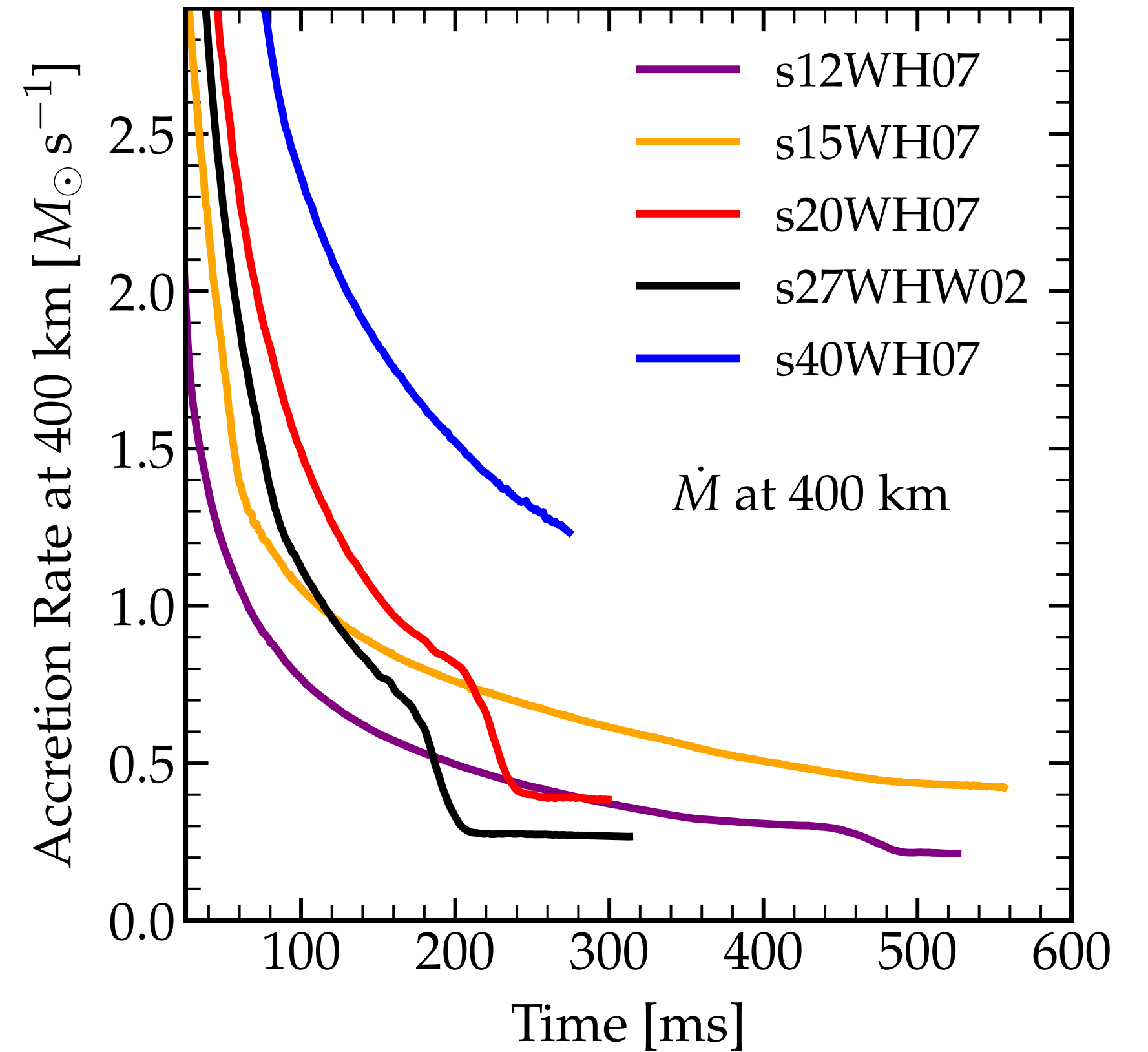
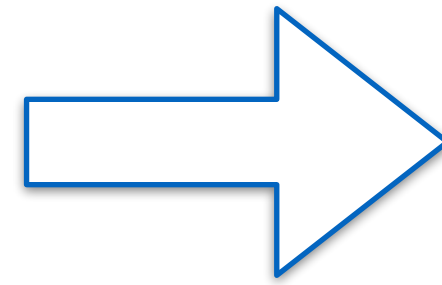
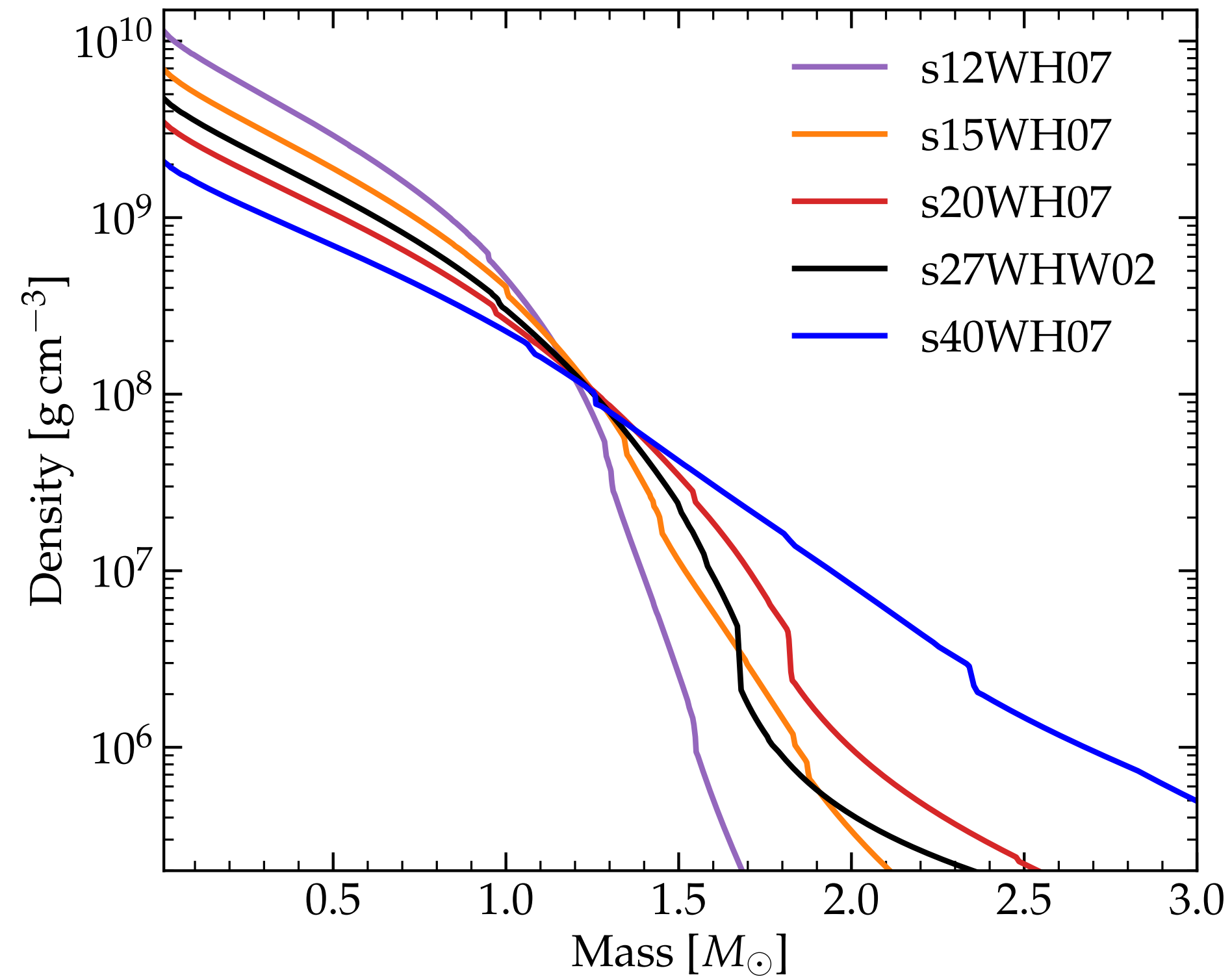
Takiwaki et al. '12, Melson '15, Lentz '15, LR et al. '16, Takiwaki et al. '16, Ott et al. '18, O'Connor and Couch '18, Burrows et al. '19



- Many groups are seeing shock runaway (although not all), but maybe not quantitative agreement
- Sensitive to input physics (e.g. Melson et al. '15) and resolution
- Nevertheless, things look relatively positive for 3D shock runaway

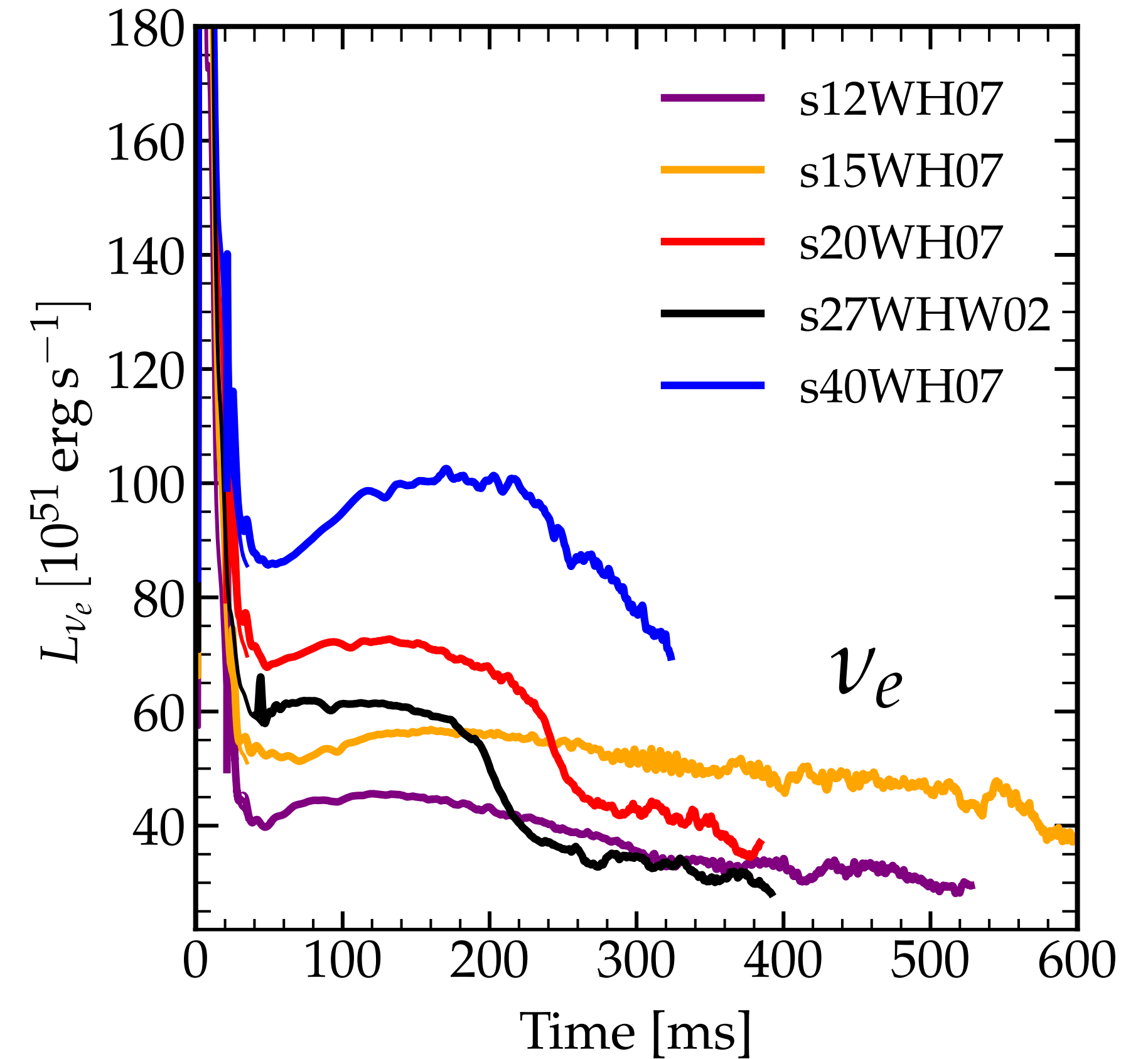
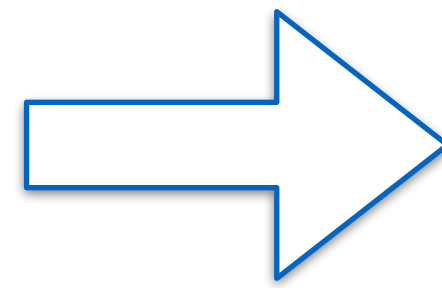
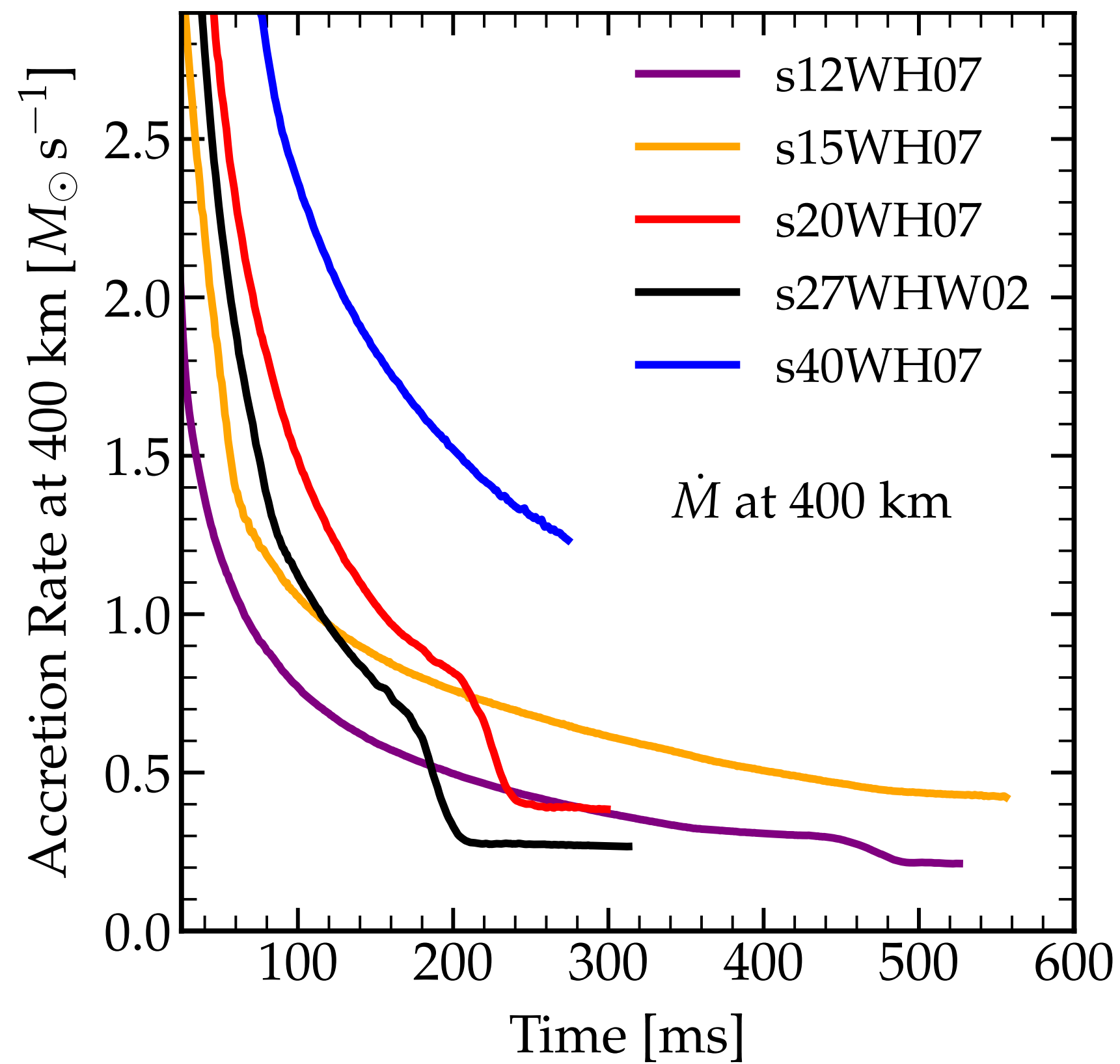
Stellar Progenitor Dependence

Ott, LR et al. (2018)



Stellar Progenitor Dependence

Ott, LR et al. (2018)



$$P_{\text{ram}} = \rho v^2 \propto \dot{M} M_{\text{PNS}}^{1/2} r_s^{-5/2}$$

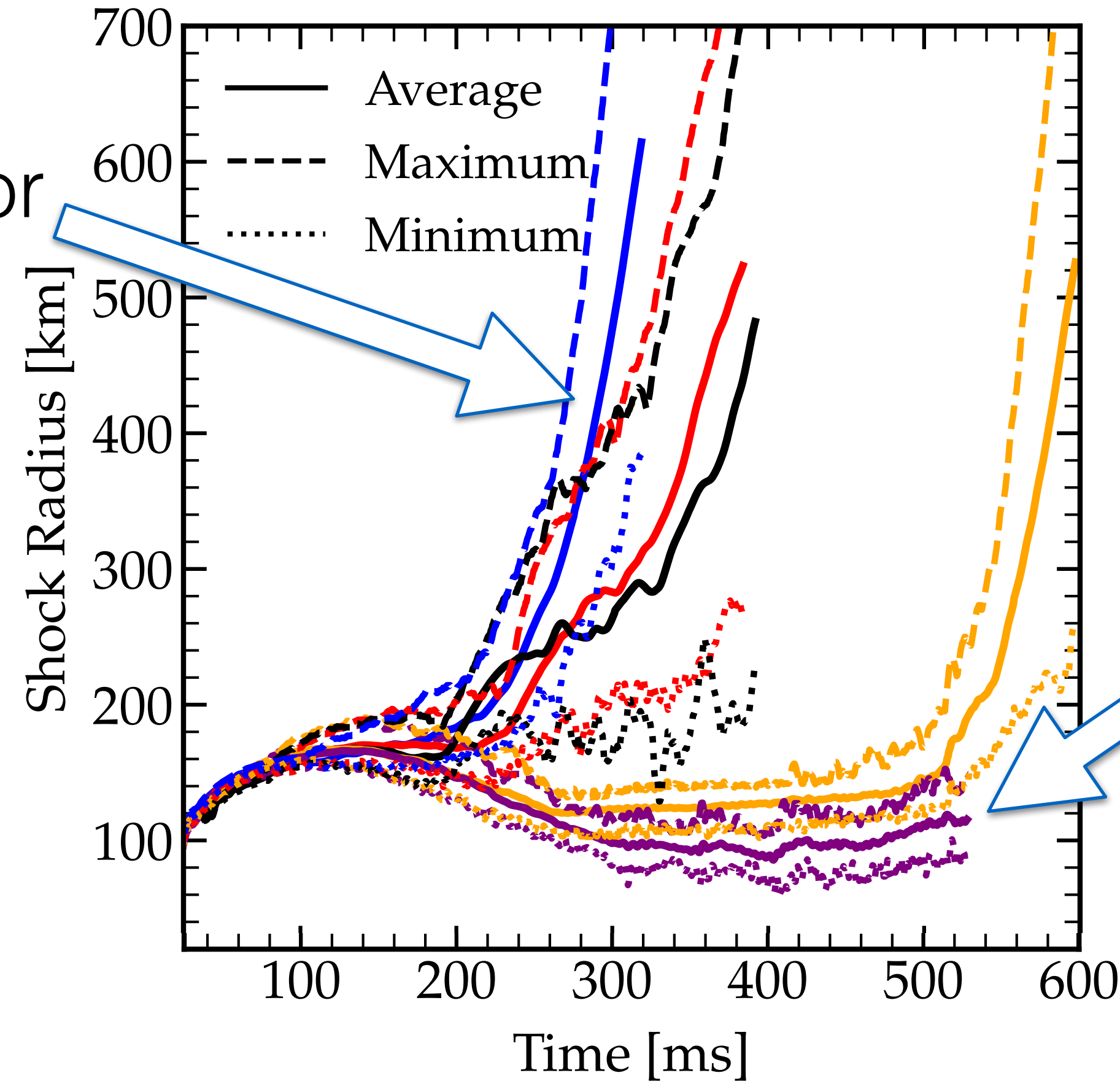
but also

$$L_{\text{acc}} \propto \dot{M} M_{\text{PNS}} R_{\text{PNS}}^{-1}$$

Stellar Progenitor Dependence

Ott, LR et al. (2018)

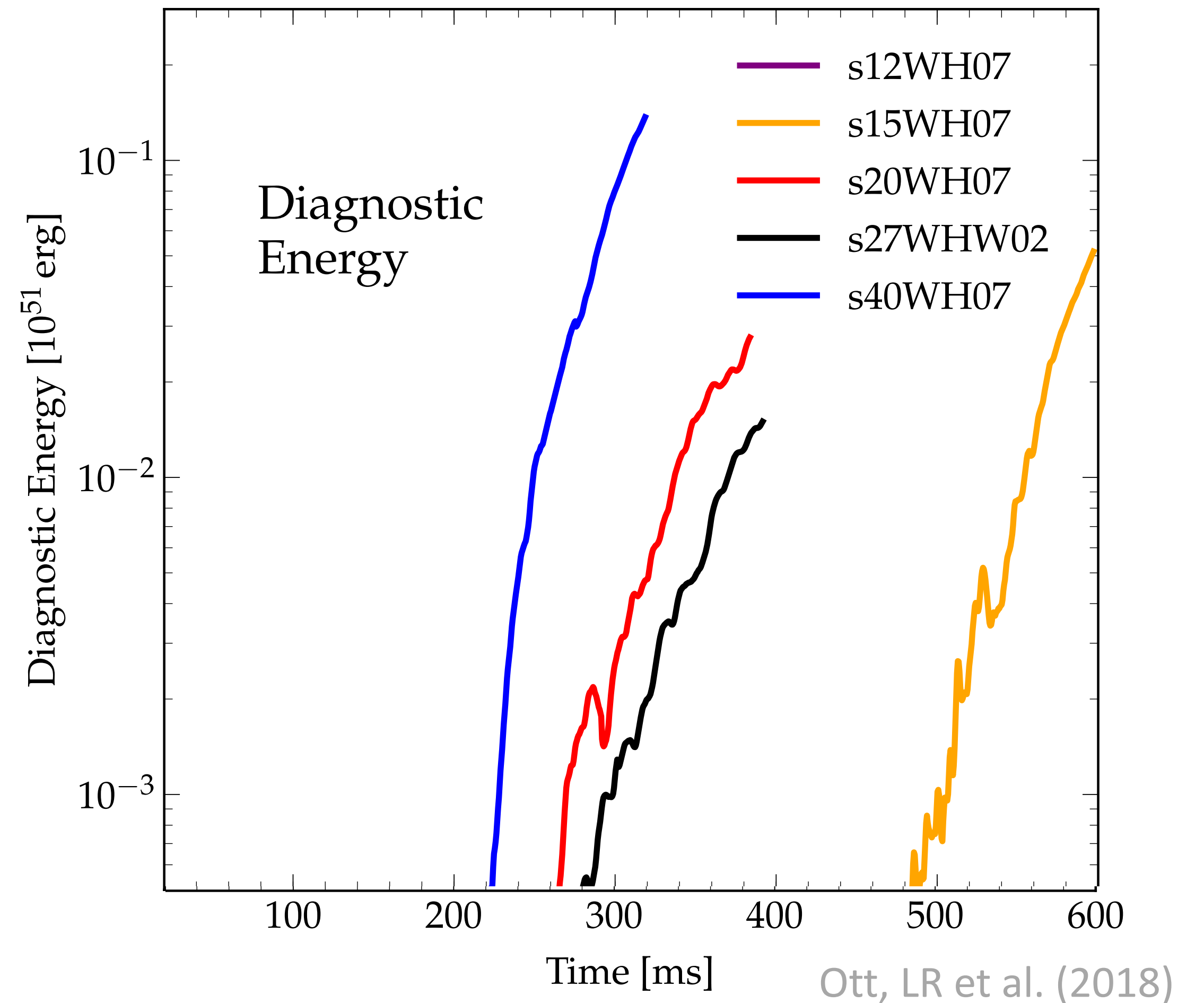
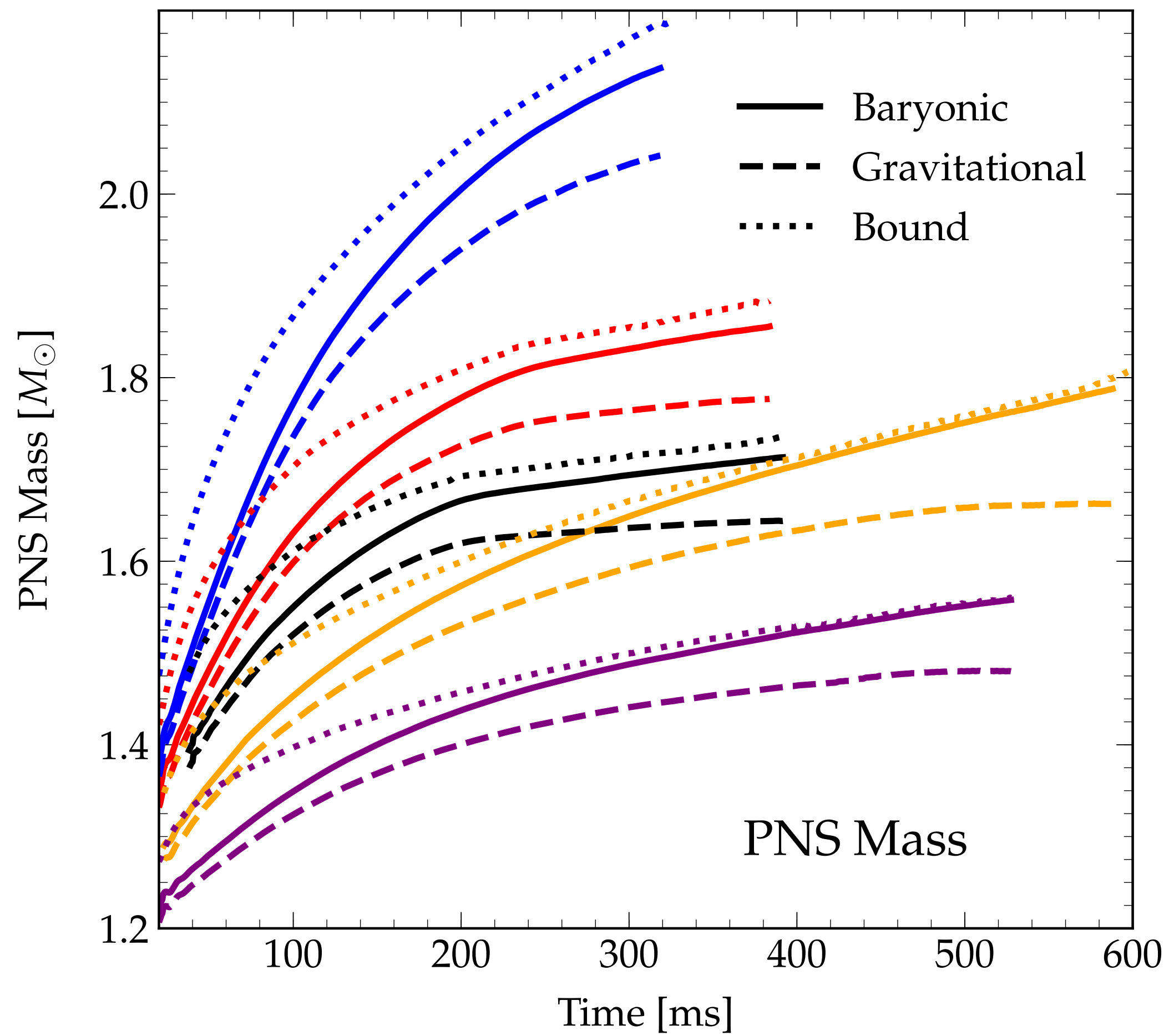
Highest mass progenitor



Lowest mass progenitor

$$P_{\text{ram}} = \rho v^2 \propto \dot{M} M_{\text{PNS}}^{1/2} r_s^{-5/2} \quad \text{but also} \quad L_{\text{acc}} \propto \dot{M} M_{\text{PNS}} R_{\text{PNS}}^{-1}$$

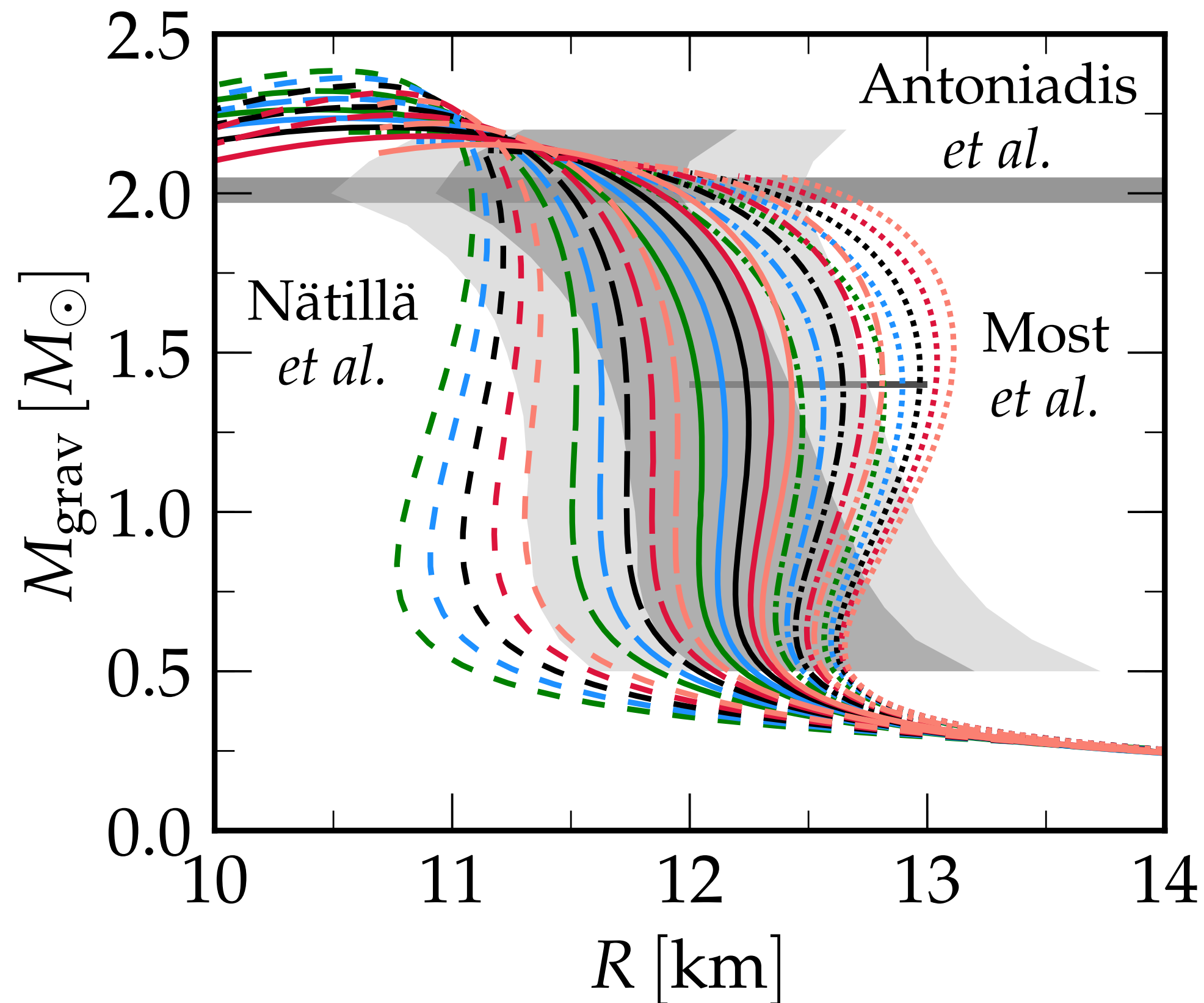
Final States



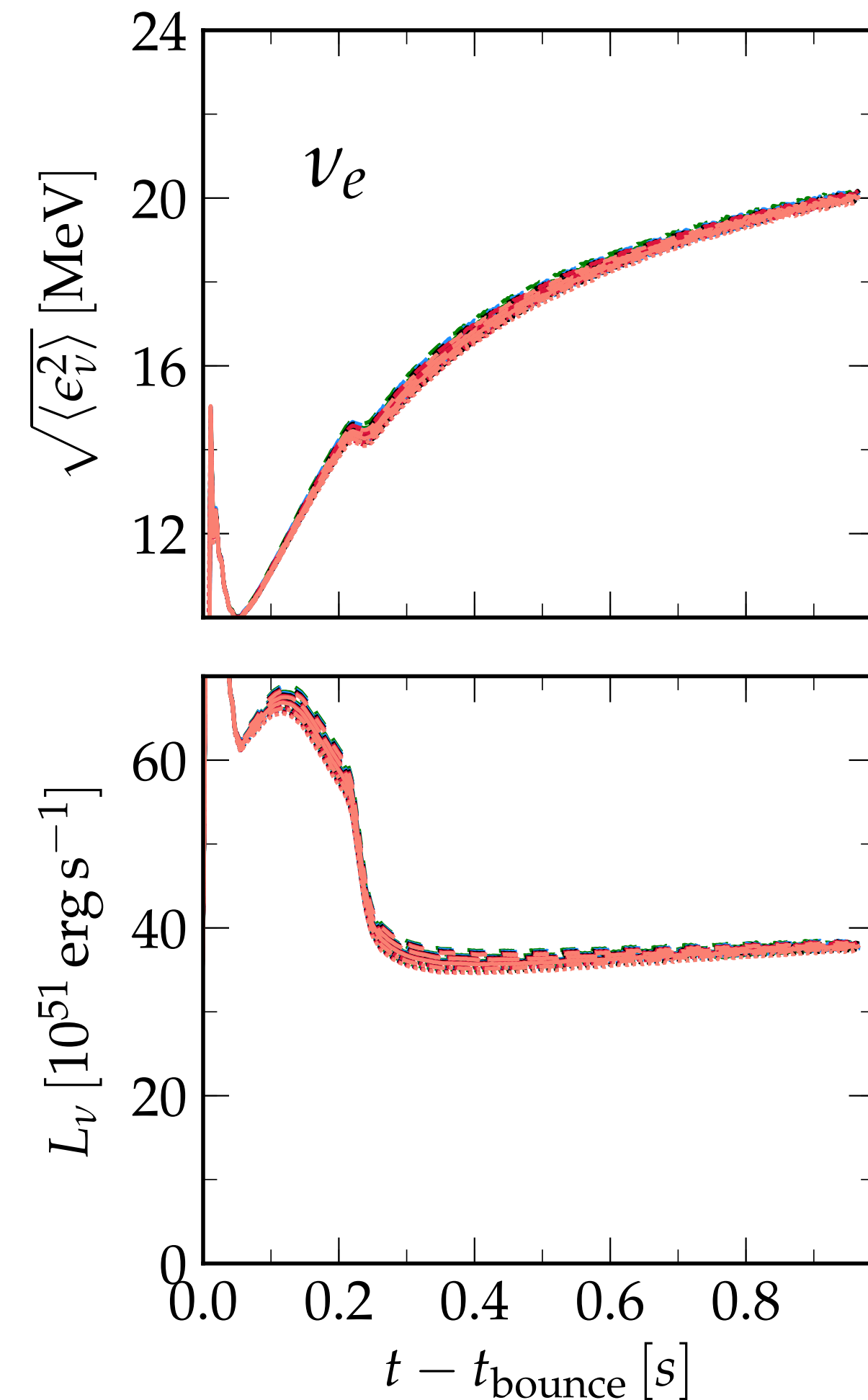
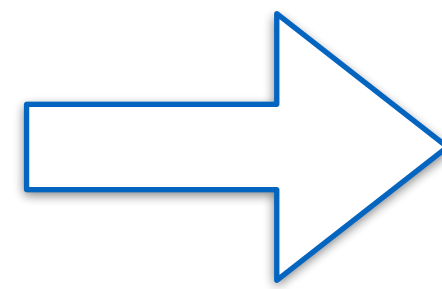
Equation of State Impacts

Melson et al. '15, Bollig et al., '17, Yassin et al. '19

$K_{\text{sat}}[\text{MeV}]$: — 200 — 215 — 230 — 245 — 260
 $K_{\text{sym}}[\text{MeV}]$: - - -300 - - -200 — -100 - - - 0 100



Schneider, LR, et al. (2019)



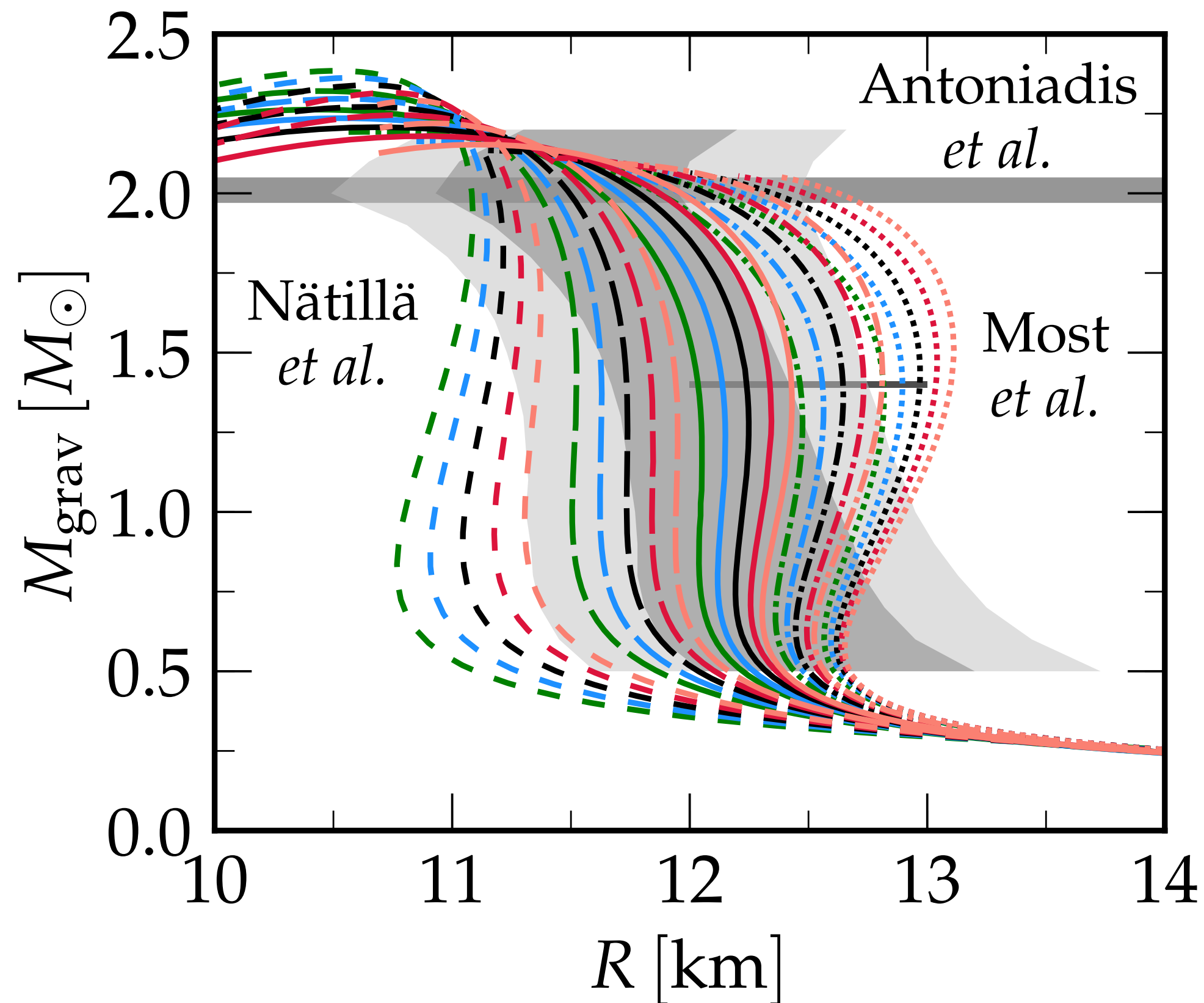
Vary within experimental constraints using Schneider '16 EoS framework:

s_M	m^*
	Δm^*
<hr/>	
—	n_{sat}
	ϵ_{sat}
<hr/>	
s_S	ϵ_{sym}
	L_{sym}
<hr/>	
s_K	K_{sat}
	K_{sym}
<hr/>	
s_P	$P_{\text{SNM}}^{(4)}$
	$P_{\text{PNM}}^{(4)}$

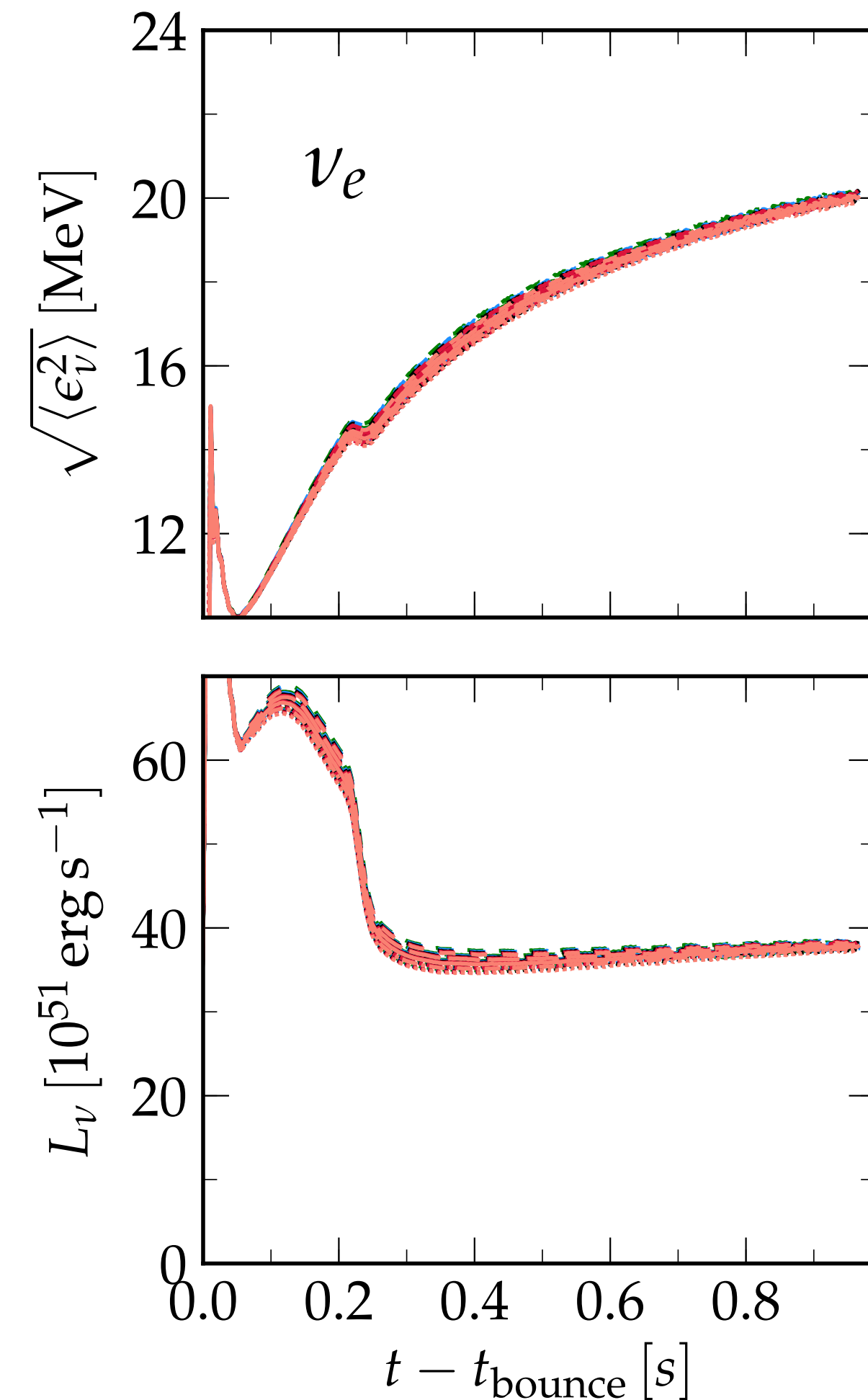
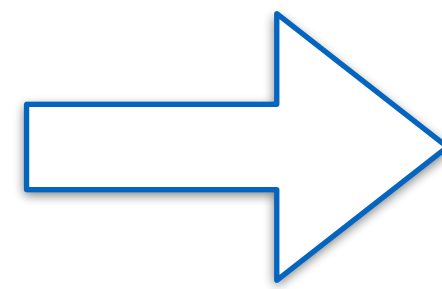
Equation of State Impacts

Melson et al. '15, Bollig et al., '17, Yassin et al. '19

$K_{\text{sat}}[\text{MeV}]$: — 200 — 215 — 230 — 245 — 260
 $K_{\text{sym}}[\text{MeV}]$: - - -300 - - -200 — -100 - ···· 0 ····· 100



Schneider, LR, et al. (2019)

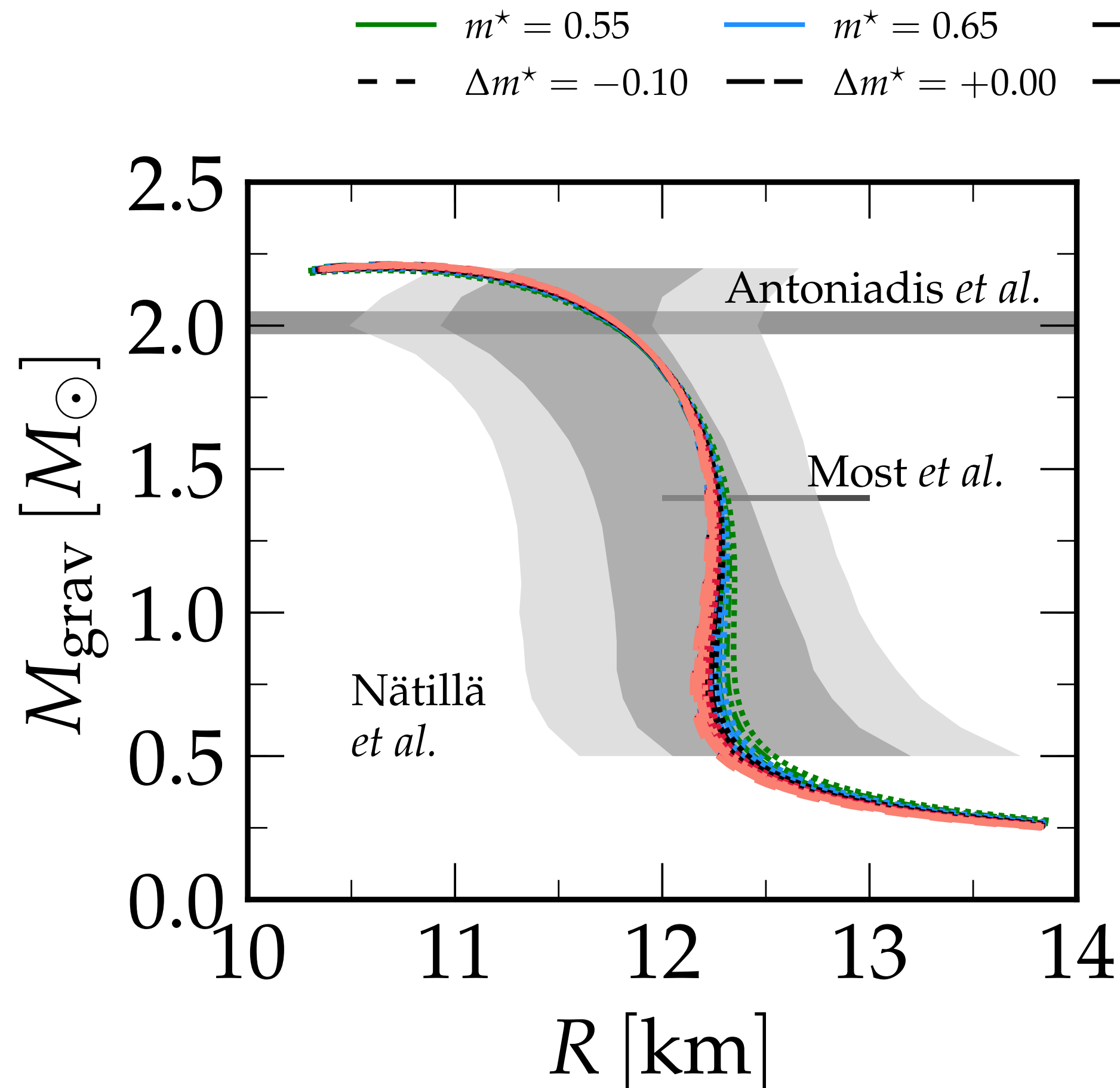


Similarly small impacts for:

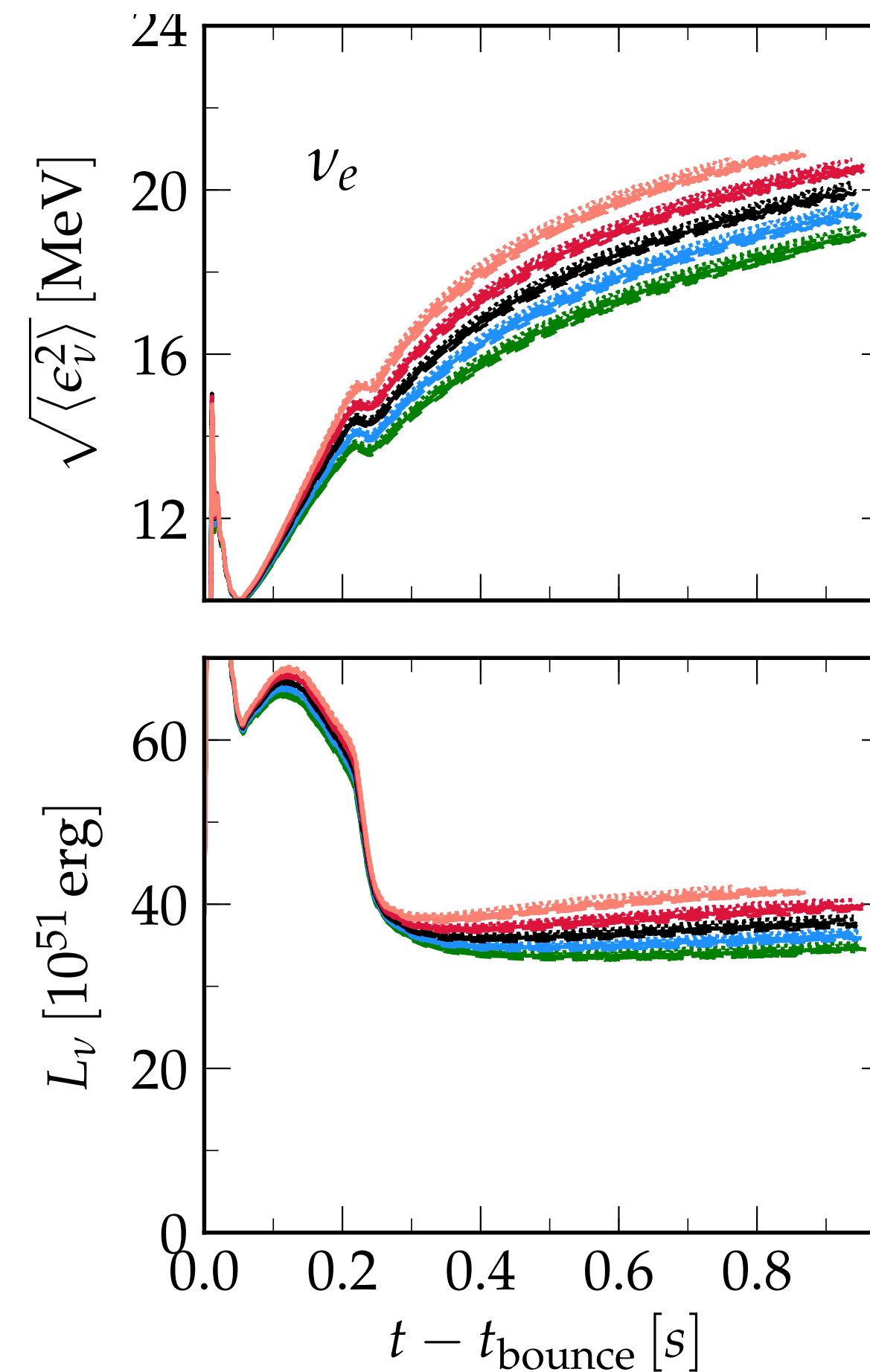
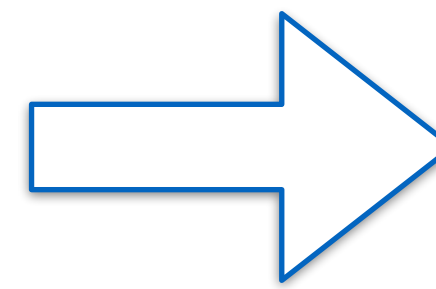
- ϵ_{sym}
- L_{sym}
- K_{sat}
- K_{sym}
- $P_{\text{SNM}}^{(4)}$
- $P_{\text{PNM}}^{(4)}$

Equation of State Impacts

Melson et al. '15, Bollig et al., '17, Yassin et al. '19



Schneider, LR, et al. (2019)

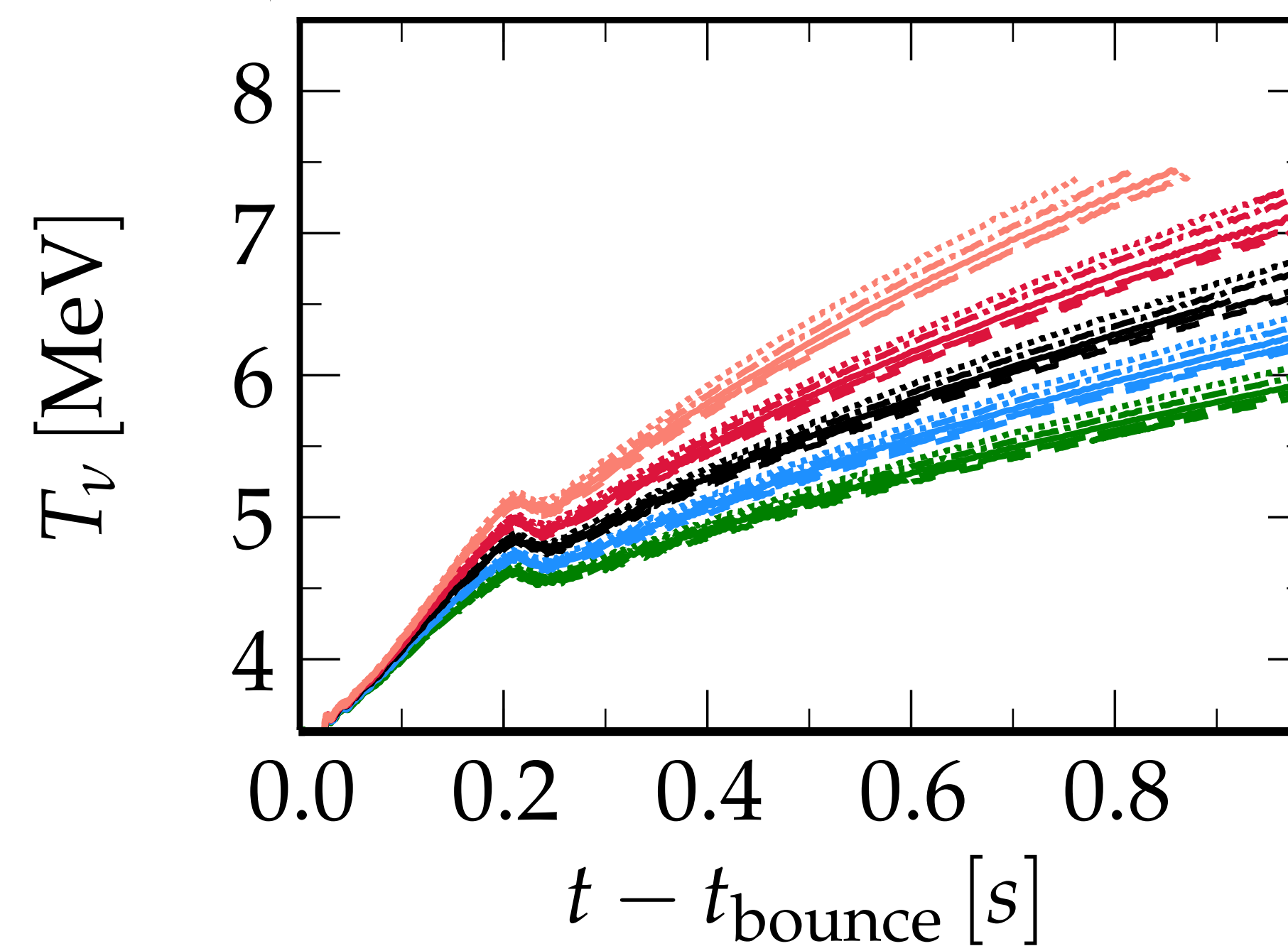
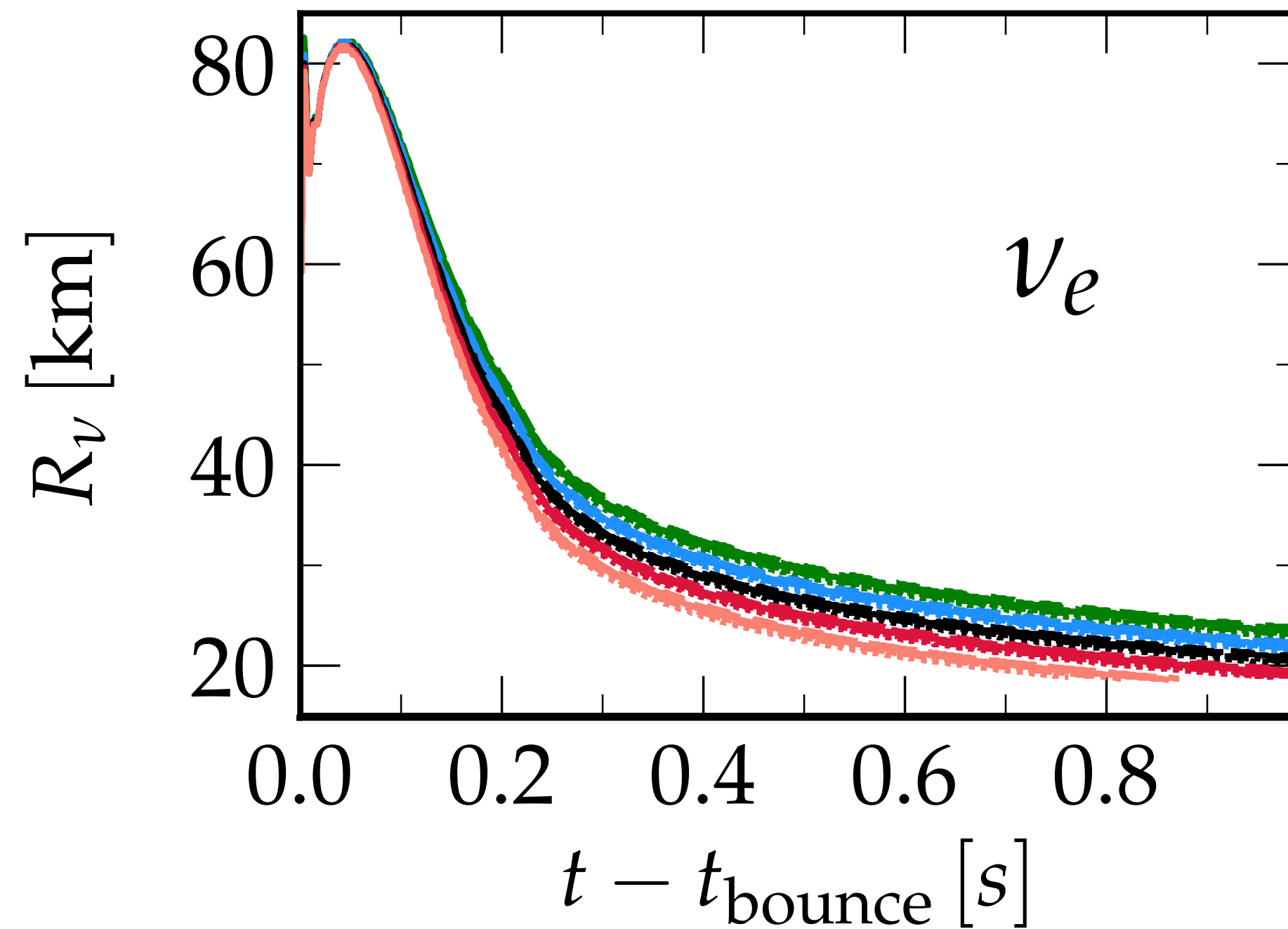
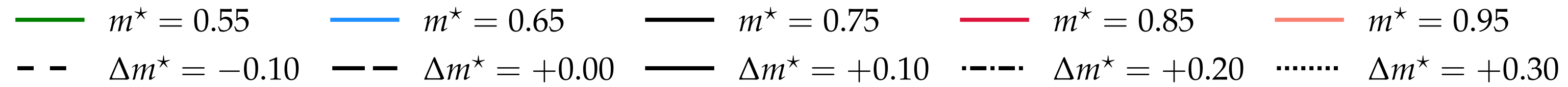


Vary the effective mass:

s_M	m^*
—	Δm^*
—	n_{sat}
—	ϵ_{sat}
s_S	ϵ_{sym}
—	L_{sym}
s_K	K_{sat}
—	K_{sym}
s_P	$P_{\text{SNM}}^{(4)}$
—	$P_{\text{PNM}}^{(4)}$

Equation of State Impacts

Melson et al. '15, Bollig et al., '17, Yassin et al. '19

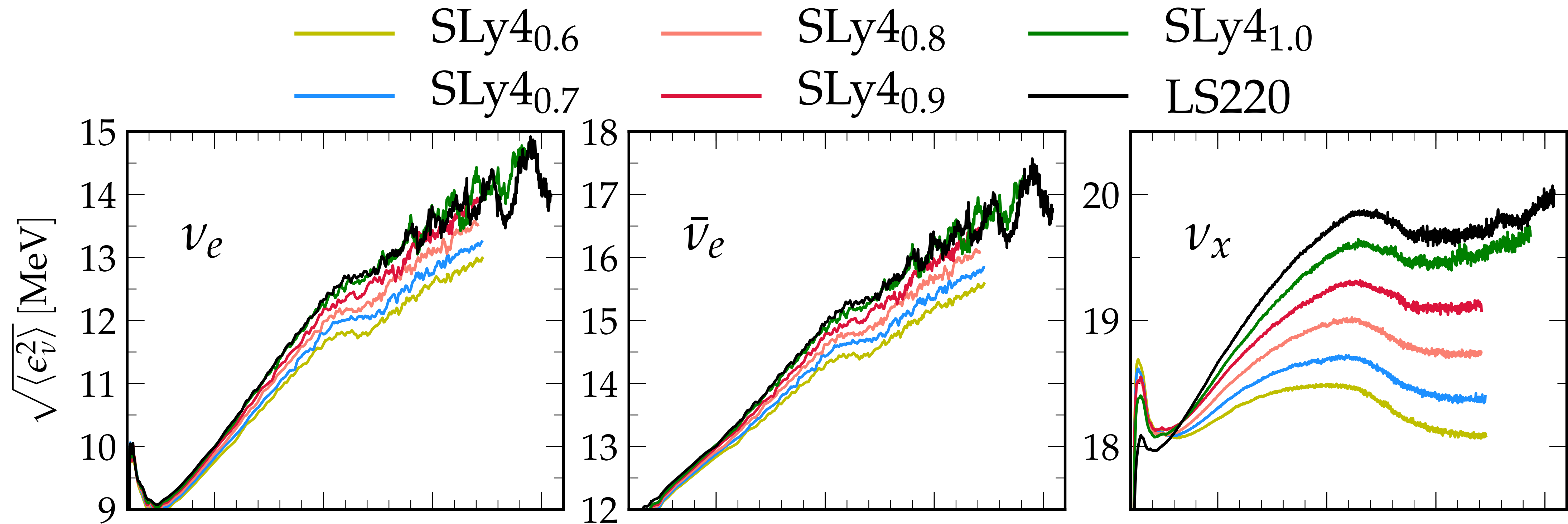


Schneider, LR, et al. (2019)

Impact on neutrino sphere properties

$$P_{th,nuc} \propto \frac{1}{m^*}$$

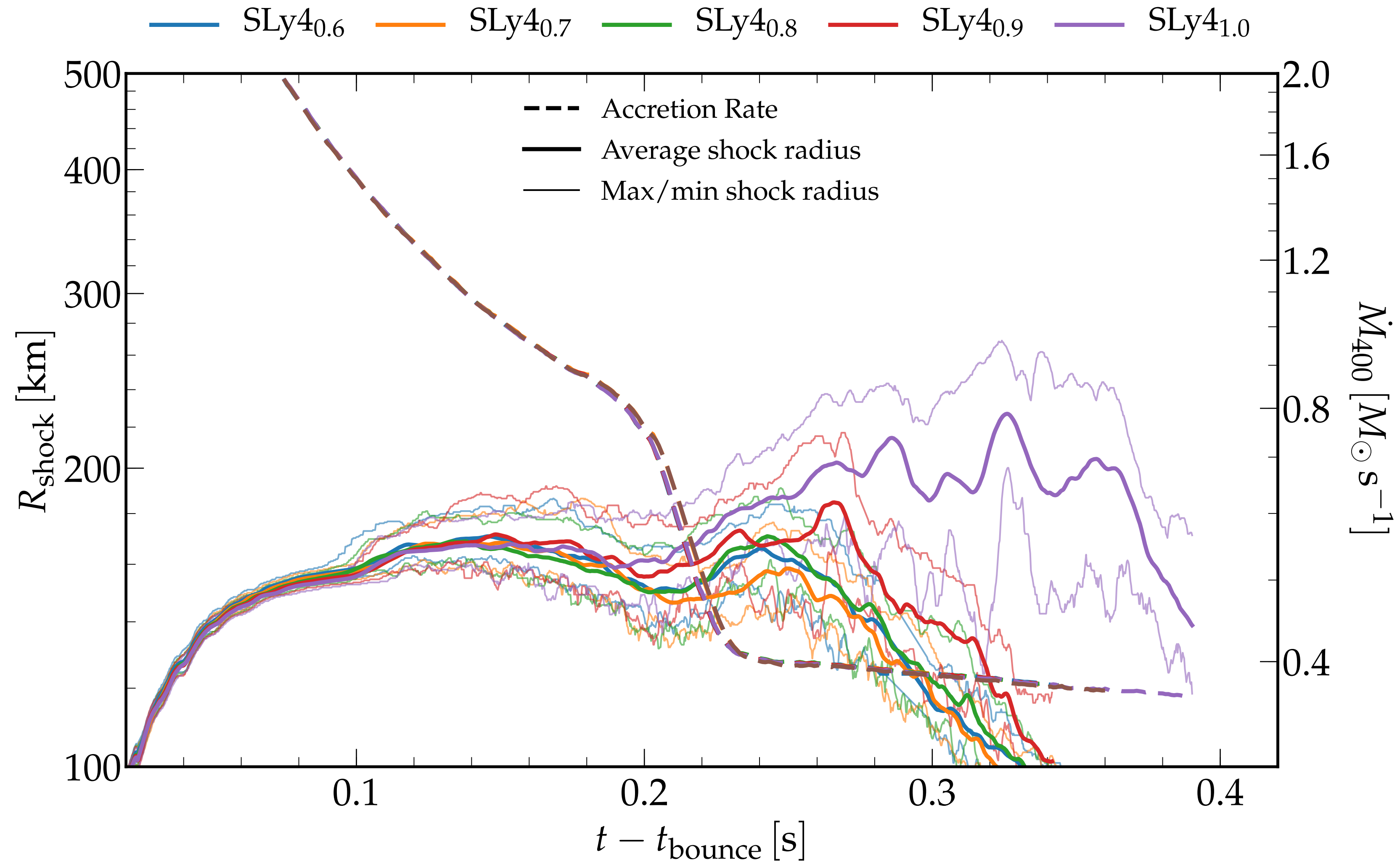
Equation of State Impacts



Schneider, LR, et al. (2019)

3D Octant simulations see similar impact of effective mass to spherically symmetric simulations

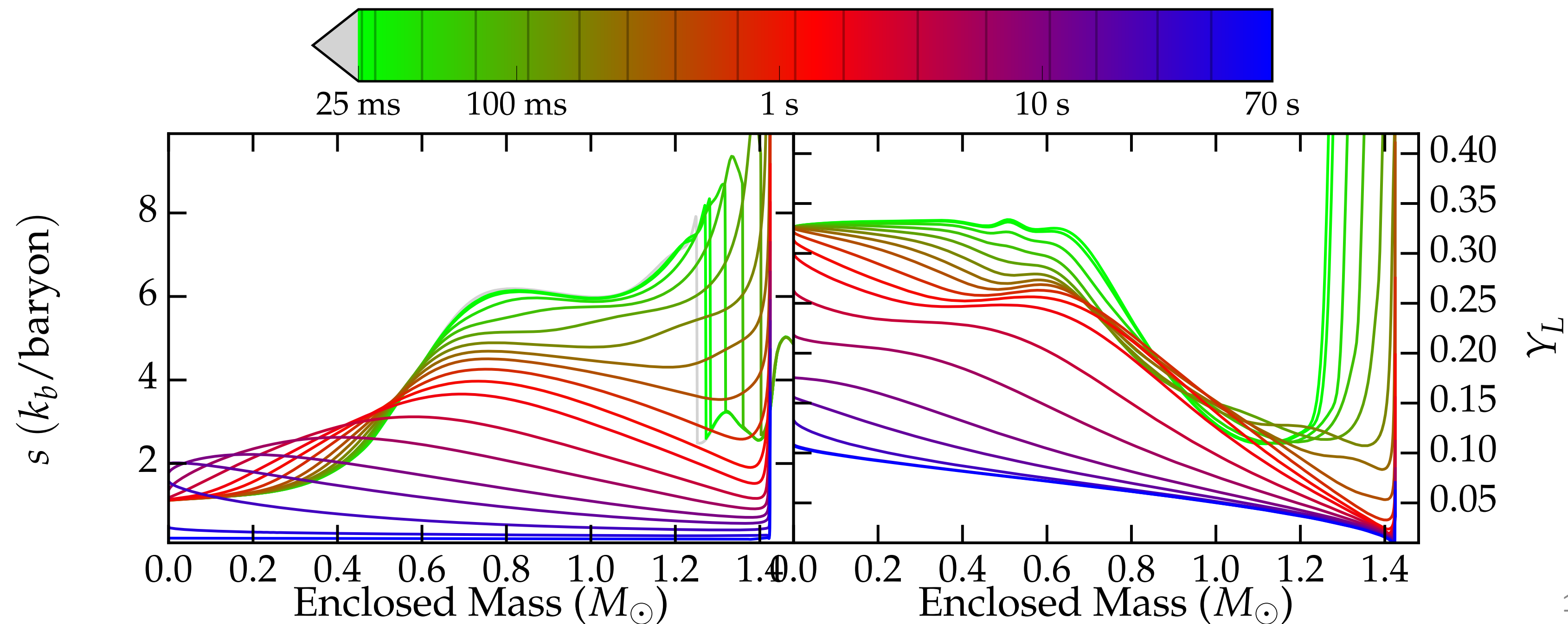
Equation of State Impacts



Late Time Neutrino Emission

See e.g. Burrows & Lattimer '86, Pons et al. '99, Huedepohl et al. '10, Fischer et al. '10, LR '12, Nakazato '13

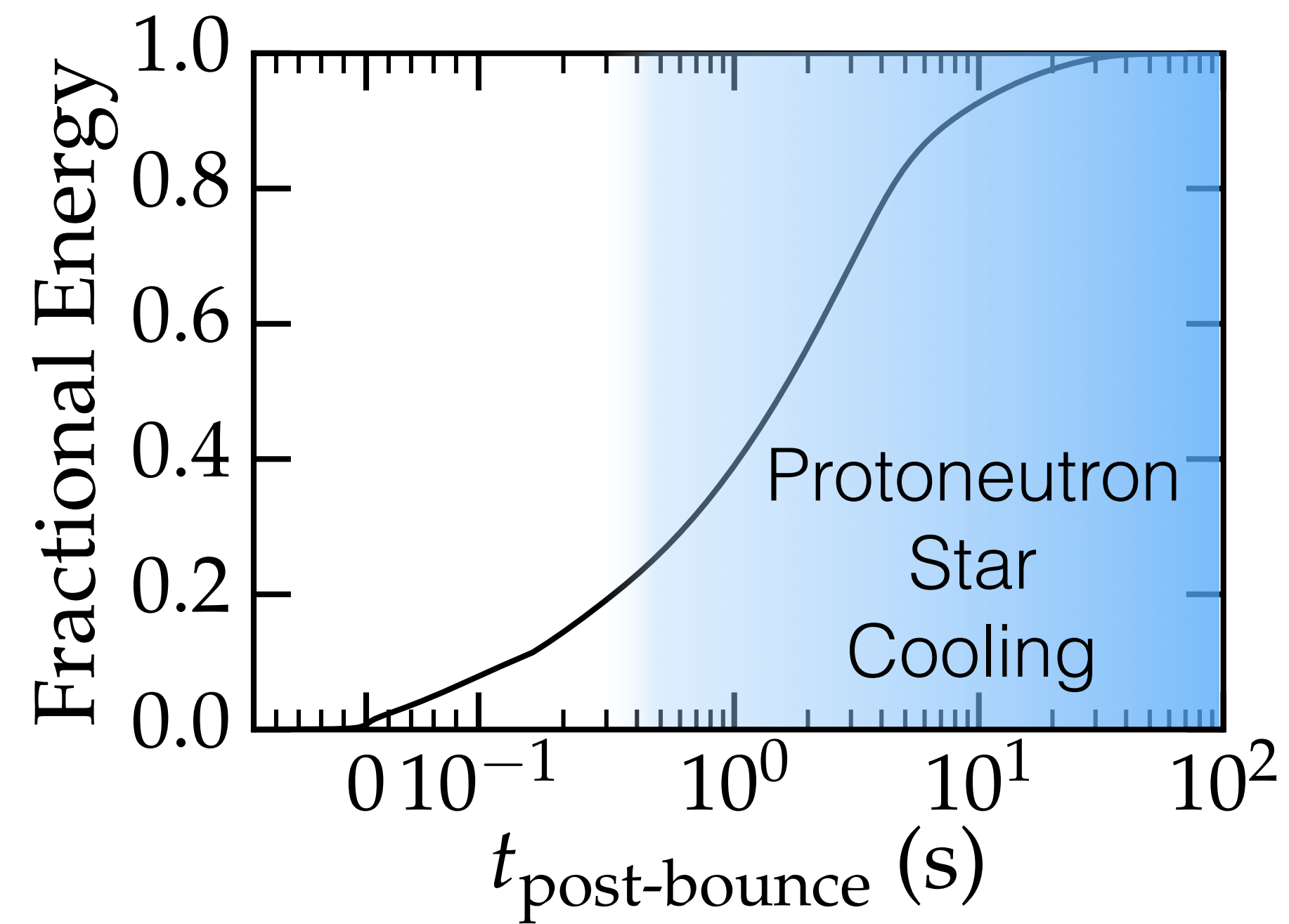
- Driven by cooling and deleptonization of the remnant
- Coupled neutron star structure and neutrino transport
- Sensitive to dense matter equation of state, neutrino opacities



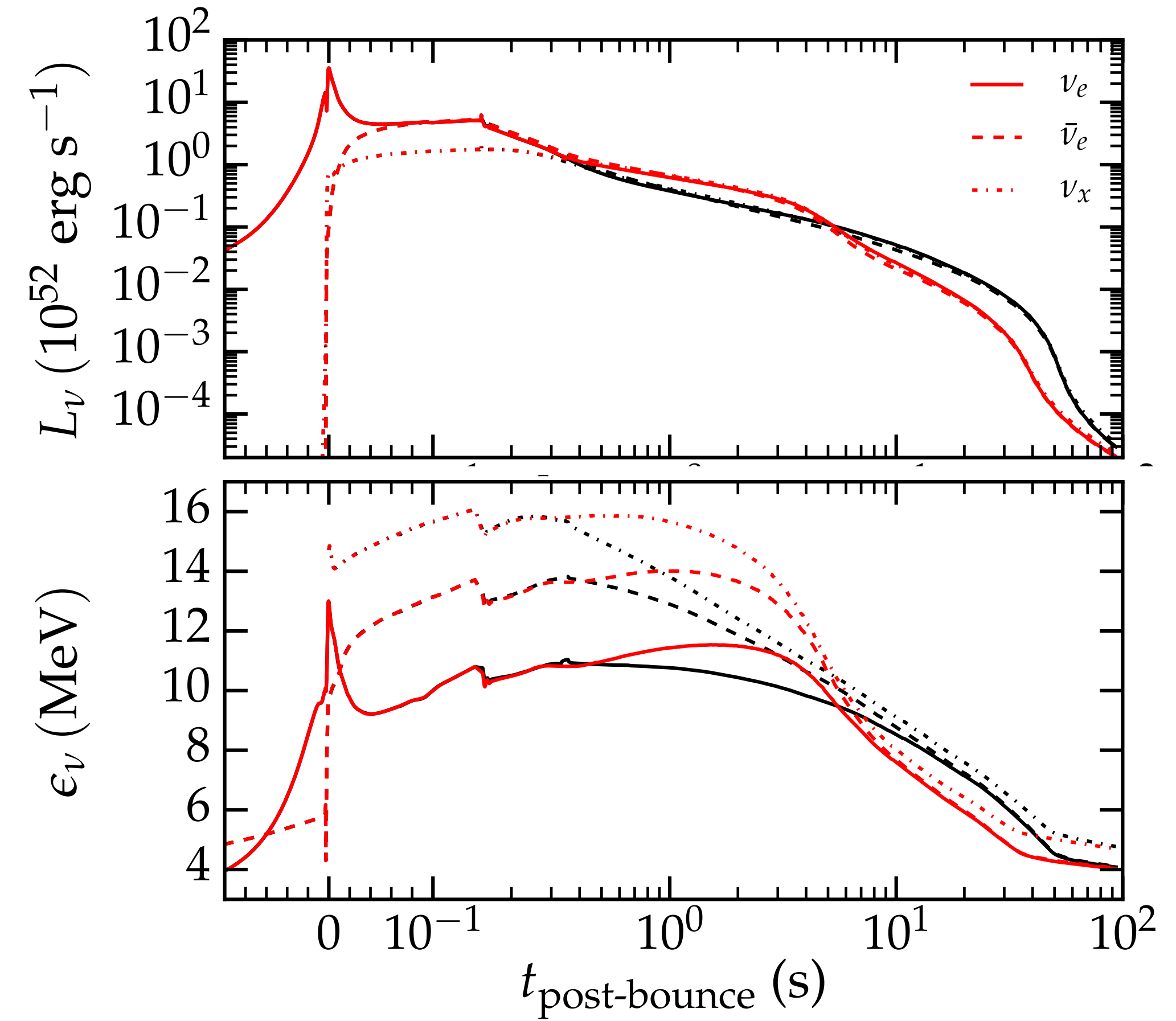
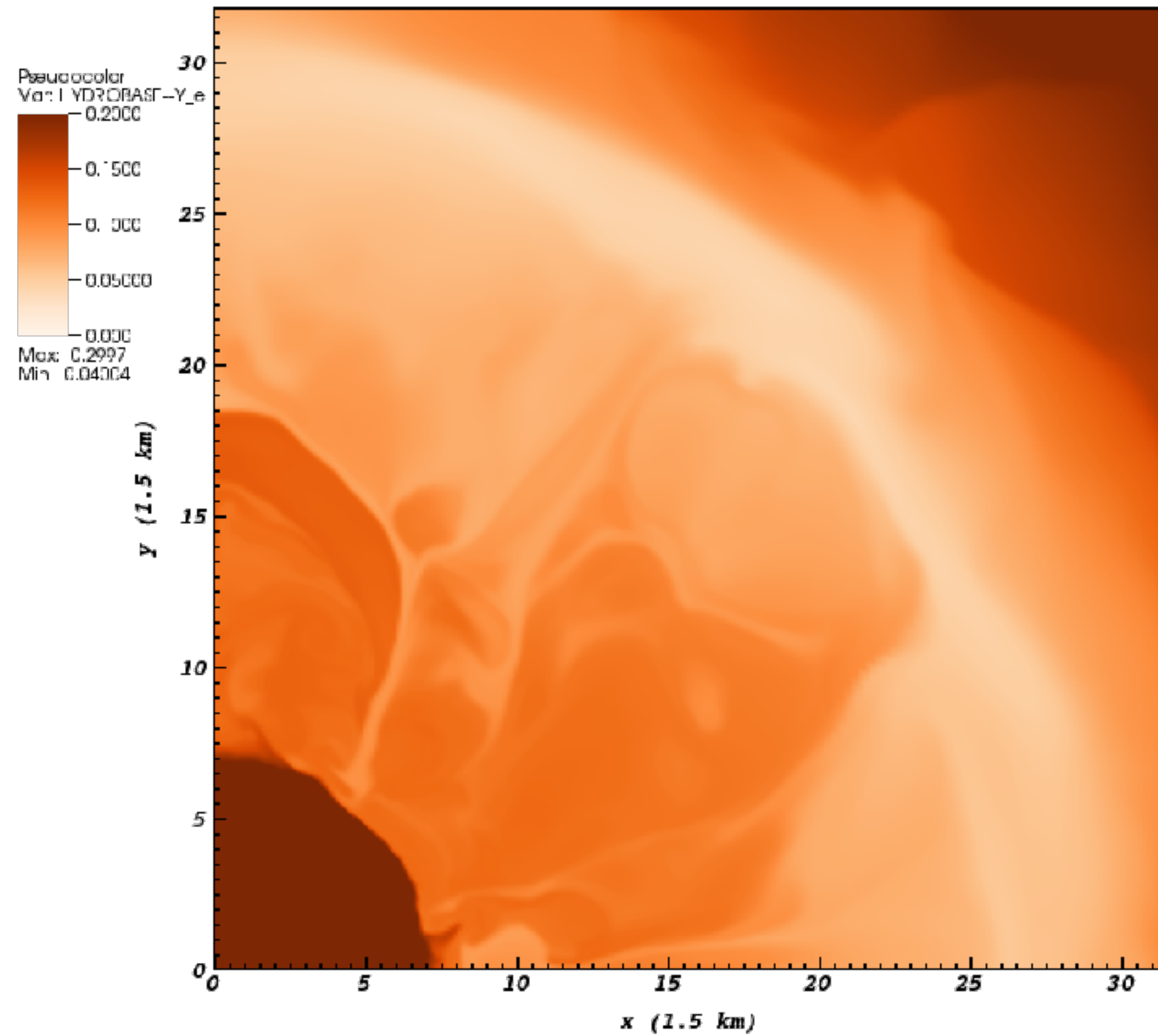
Late Time Neutrino Emission

See e.g. Burrows & Lattimer '86, Pons et al. '99, Huedepohl et al. '10, Fischer et al. '10, LR '12, Nakazato '13

- Driven by cooling and deleptonization of the remnant
- Coupled neutron star structure and neutrino transport
- Sensitive to dense matter equation of state, neutrino opacities



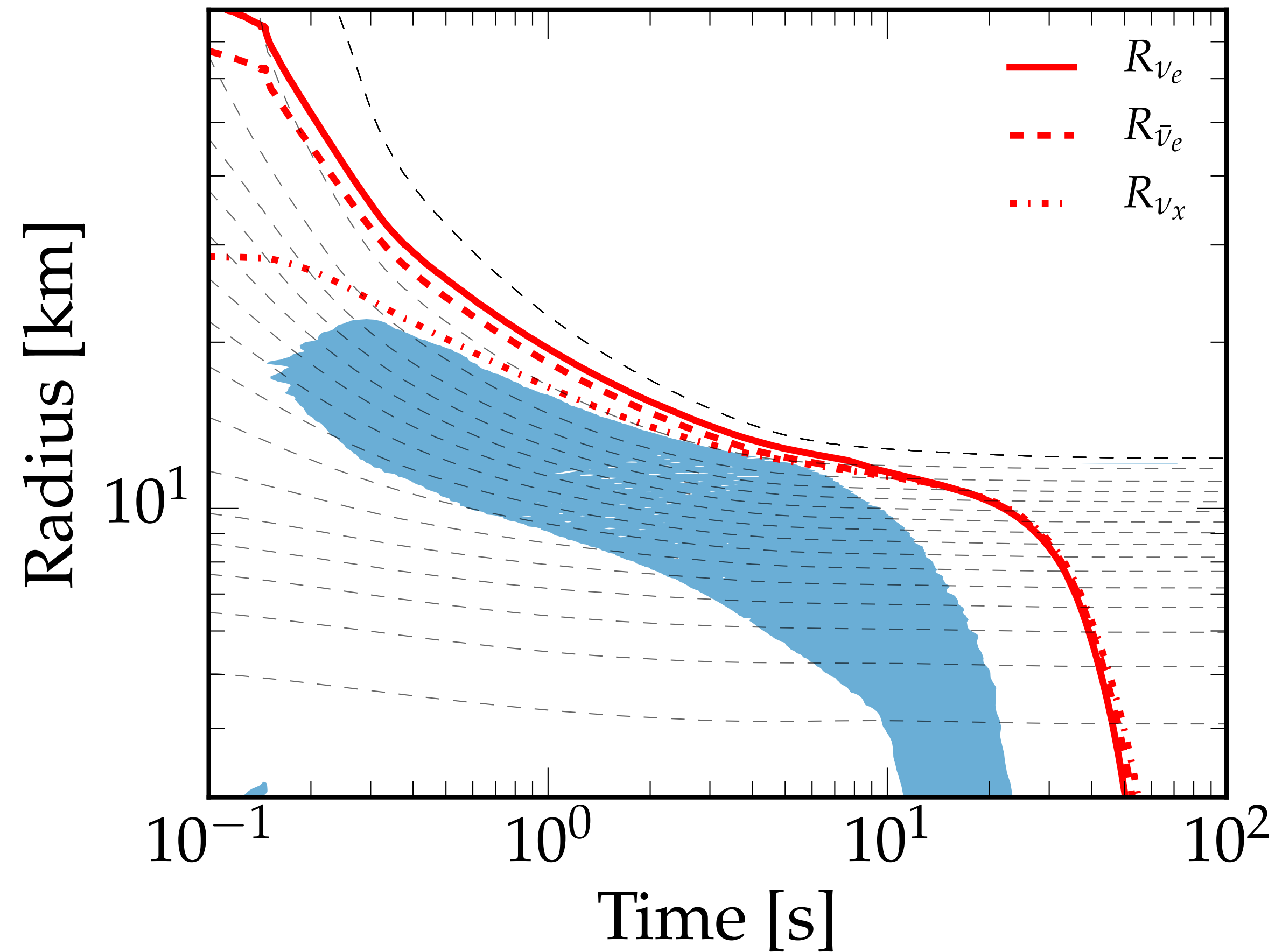
Proto-Neutron Star Convection



Red: Convection

Black: No Convection

Proto-Neutron Star Convection

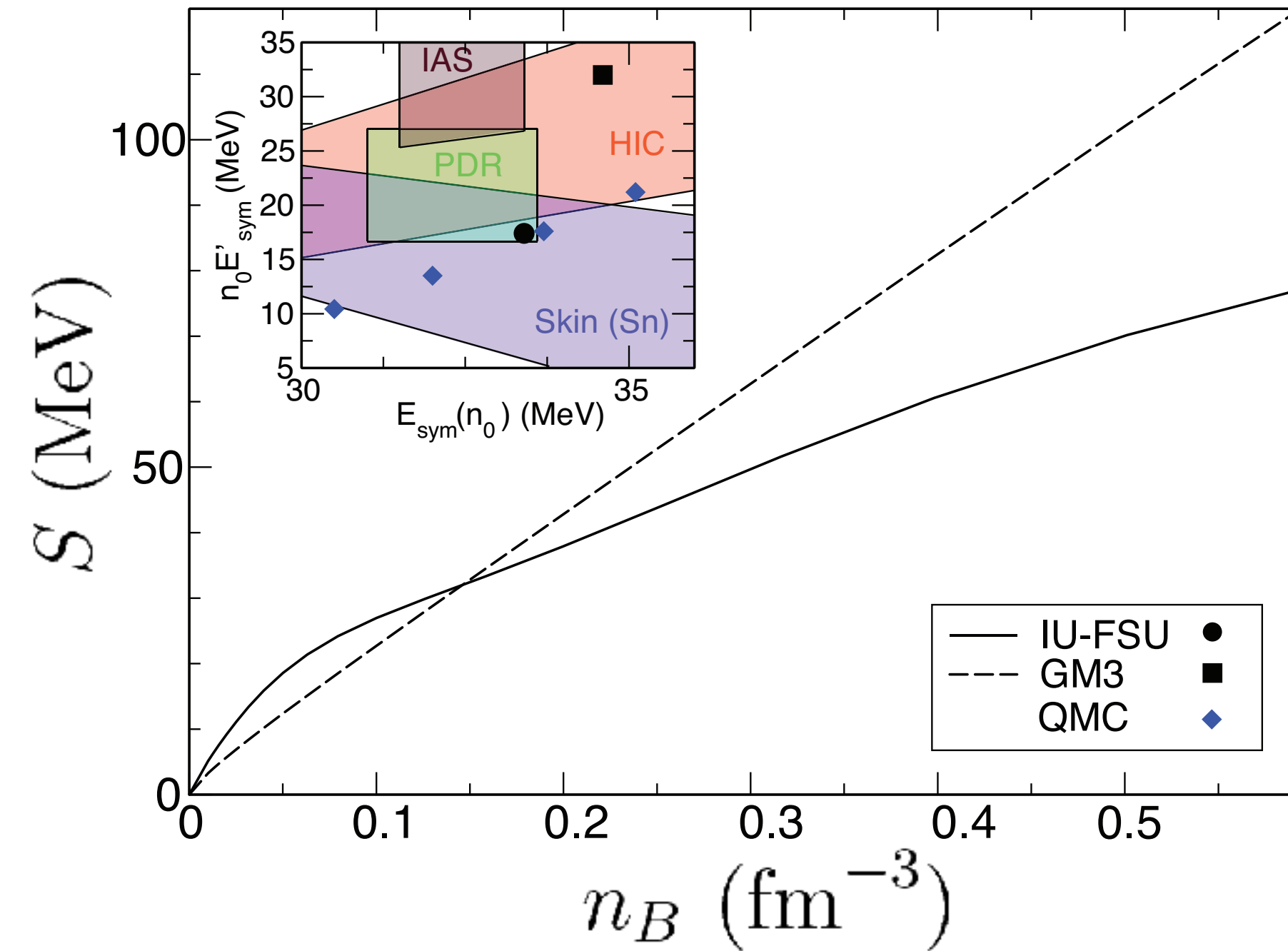


Region of convective instability determined by the Ledoux Criterion:

$$C_L = - \left(\frac{\partial P}{\partial s} \right)_{n, Y_l} \frac{ds}{dr} - \left(\frac{\partial P}{\partial Y_l} \right)_{n, s} \frac{dY_l}{dr} > 0$$

Proto-Neutron Star Convection

Dependence on the EoS



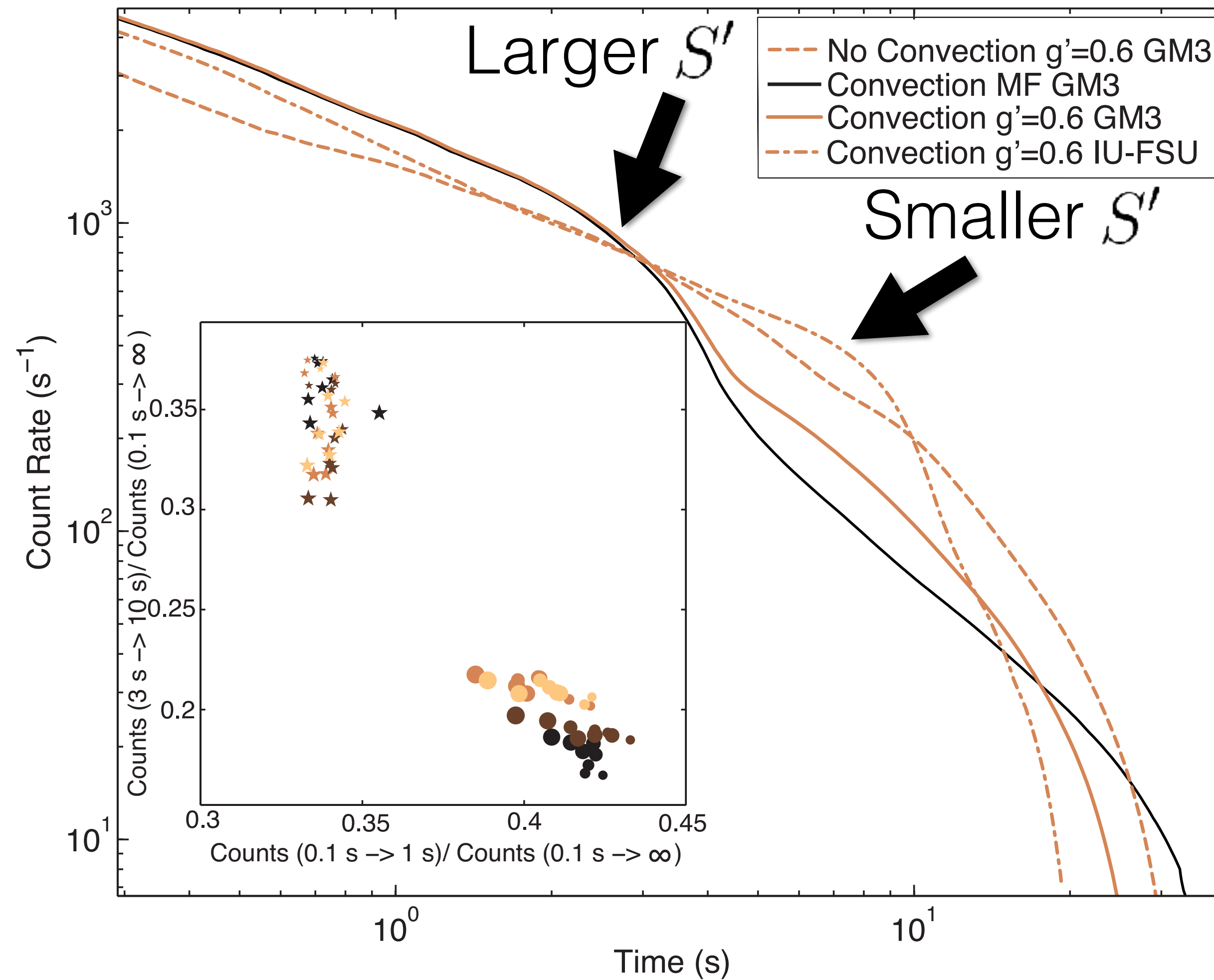
Pressure derivatives are sensitive to the symmetry energy derivative:

$$\epsilon(n_B, Y_e) = \epsilon(n_B, 1/2) + S(n_B)(1 - 2Y_e)^2$$

➔

$$\left(\frac{dP}{dY_L} \right)_{n_B, s} \approx n_B^{4/3} Y_e^{1/3} - 4n_B^2 S'(1 - 2Y_e)$$

Proto-Neutron Star Convection



LR et al. (2012)

Conclusions

- 3D explosion models becoming available for a range of progenitors
- Early-time neutrino emission is sensitive to finite temperature properties of nuclear equation of state, larger nucleon effective masses result in conditions more favorable for explosion
- PNS convection significantly impacts the neutrino cooling timescale, produces a break in the neutrino emission, sensitive to the nuclear EoS through the density dependence of the symmetry energy
- ~~Nuclear correlations and nuclear pasta have possibility of impacting late time cooling, but the critical temperature for pasta formation may be too low and neutrino opacity in the relevant range suppressed~~

Set	Quantity	Range	This work	Units
s_M	m^*	0.75 ± 0.10	0.75 ± 0.10	m_n
	Δm^*	0.10 ± 0.10	0.10 ± 0.10	m_n
—	n_{sat}	0.155 ± 0.005	0.155	fm^{-3}
	ϵ_{sat}	-15.8 ± 0.3	-15.8	MeV baryon^{-1}
s_S	ϵ_{sym}	32 ± 2	32 ± 2	MeV baryon^{-1}
	L_{sym}	60 ± 15	45 ± 7.5	MeV baryon^{-1}
s_K	K_{sat}	230 ± 20	230 ± 15	MeV baryon^{-1}
	K_{sym}	-100 ± 100	-100 ± 100	MeV baryon^{-1}
s_P	$P_{\text{SNM}}^{(4)}$	100 ± 50	125 ± 12.5	MeV fm^{-3}
	$P_{\text{PNM}}^{(4)}$	160 ± 80	200 ± 20	MeV fm^{-3}