## 1. Basics of nonrelativistic quantum mechanics

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very brief repitition of some basic elements of nonrelativistic quantum mechanics you probably know already from the introductory courses
- Goals:
- general introduction to this course
- among other things: preparation for the relativistic quantum mechanics


### 1.1 Heuristic motivation of the Schrödinger equation from wave-particle duality

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- Wave-Particle-Duality

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- energy $E=\hbar \omega$ (Einstein 1905; photoelectric effect),
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This motivates the introduction of hermitian operators

- $\hat{E}=i \hbar \frac{\partial}{\partial t} \quad \Rightarrow \quad \hat{E} \psi=\hbar \omega \psi$
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- nonrelativistic energy-momentum relation: $\quad E=\frac{\vec{\rho}^{2}}{2 m}+V$
$\Rightarrow \quad$ Schrödinger equation: $i \hbar \frac{\partial}{\partial t} \psi=\left(-\frac{\hbar^{2}}{2 m} \vec{\nabla}^{2}+V\right) \psi \equiv H \psi$
- Hamiltonian: $H=-\frac{\hbar^{2}}{2 m} \vec{\nabla}^{2}+V$


### 1.2 Probabalistic interpretation and continuity equation

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Obviously, this also works with a real potential $V$.

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$\Rightarrow \quad \frac{\partial \rho}{\partial t}+\vec{\nabla} \cdot \vec{j}=0 \quad$ continuity equation
probability density: $\quad \rho=\psi^{*} \psi=|\psi|^{2}$
probability current:

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\vec{j}=\frac{\hbar}{2 m i}\left(\psi^{*} \vec{\nabla} \psi-\psi \vec{\nabla} \psi^{*}\right) \equiv \frac{\hbar}{2 m i} \psi^{*}(\vec{\nabla}-\overleftarrow{\nabla}) \psi
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- example: $\quad \psi(\vec{r}, t)=\mathcal{N} e^{i(\vec{k} \cdot \vec{r}-\omega t)}$
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variation of charge in $\mathcal{V}$ with time $=-$ current through the surface of $\mathcal{V}$
- Quantum mechanics: $Q=$ probability to find a particle in $\mathcal{V}$
(if correctly normalized as $\int_{\mathbb{R}^{3}} d^{3} r|\psi|^{2}=1$ )


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(ii) Observables are represented by hermitian operators $\hat{O}$.
(iii) Possible measurements correspond to the eigenvalues of $\hat{O}$.
(iv) The corresponding eigenstates $|n\rangle$, i.e., $\hat{O}|n\rangle=\lambda_{n}|n\rangle$, form a complete orthonormal basis of the Hilbert space.

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\langle m \mid n\rangle=\delta_{m n}, \quad \sum_{n}|n\rangle\langle n|=11
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(For continuous spectra one can generalize this to $\delta$-functions and integrals.)
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$=$ sum over the possible measurements, weighted by their probability $\checkmark$


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- coordinate-space rep.: $\quad \psi(\vec{r}, t)=\langle\vec{r} \mid \psi(t)\rangle$,
$|\vec{r}\rangle$ : eigenstate of the coordinate operator


## Time evolution

(vi) Immediately after a measurement of the value $\lambda_{n}$, the system is in the eigenstate $|n\rangle$.
(vii) As long as no measurement is performed, the time evolution of $|\psi\rangle$ is determined by the Schrödinger equation:

$$
i \hbar \frac{\partial}{\partial t}|\psi(t)\rangle=\hat{H}|\psi(t)\rangle
$$

Formal solution for time independent Hamiltonians:

$$
|\psi(t)\rangle=\exp \left(-\frac{i}{\hbar} \hat{H} t\right)|\psi(0)\rangle \equiv \sum_{n=0}^{\infty} \frac{1}{n!}\left(-\frac{i}{\hbar} \hat{H} t\right)^{n}|\psi(0)\rangle
$$

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- coordinate-space rep.: $\quad \psi(\vec{r}, t)=\langle\vec{r} \mid \psi(t)\rangle$,
$|\vec{r}\rangle$ : eigenstate of the coordinate operator
- momentum-space rep.: $\tilde{\psi}(\vec{p}, t)=\langle\vec{p} \mid \psi(t)\rangle$,
$|\vec{p}\rangle$ : eigenstate of the momentum operator


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- Time derivative of the Heisenberg operator: $\frac{d \hat{O}_{H}}{d t}=\frac{1}{i \hbar}\left[\hat{O}_{H}, \hat{H}\right]$ Heisenberg's equation of motion

