

► Lagrangedichte:

$$\mathcal{L} = (D_\mu \phi)^* (D^\mu \phi) + \mu^2 \phi^* \phi - \lambda (\phi^* \phi)^2 - \frac{1}{4} F_{\mu\nu} F^{\mu\nu}$$

- $\phi = \frac{1}{\sqrt{2}}(\phi_1 + i\phi_2)$ komplexes Skalarfeld
- $D_\mu \phi = (\partial_\mu - igA_\mu)\phi$, A_μ : abelsches Eichfeld, $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$
- invariant unter $\phi(x) \rightarrow e^{i\alpha(x)}\phi(x)$, $A_\mu(x) \rightarrow A_\mu(x) + \frac{1}{g}\partial_\mu\alpha(x)$

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▶ spontane Symmetriebrechung:

- ▶ $V(\phi) = -\mu^2|\phi|^2 + \lambda|\phi|^4$ hat Minimum bei $|\phi| = \frac{v}{\sqrt{2}}$, $v \equiv \frac{\mu}{\sqrt{\lambda}}$
- ▶ wähle Vakuum $\phi_1 = v$, $\phi_2 = 0 \rightarrow$ verschobene Felder: $\phi'_1 = \phi_1 - v$, $\phi'_2 = \phi_2$
- ▶ ohne Eichfeld: $\phi'_2 =$ Goldstone-Boson

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▶ kovariante Ableitung: $(D_\mu \phi)^* (D^\mu \phi) = \dots - gv\partial^\mu \phi'_2 A_\mu + \frac{1}{2}g^2 v^2 A^\mu A_\mu$

- ▶ Massenterm für das Eichboson: $M = gv$
- ▶ A^μ und ϕ'_2 mischen

► neue reelle Felder: $\phi(x) = \frac{1}{\sqrt{2}}[v + \eta(x)] e^{i\xi(x)/v}$

► Eichtransformation:

$$\phi(x) \rightarrow \tilde{\phi}(x) = e^{-i\xi(x)/v} \phi(x) = \frac{1}{\sqrt{2}}[v + \eta(x)]$$

$$A_\mu(x) \rightarrow B_\mu(x) = A_\mu - \frac{1}{g_v} \partial_\mu \xi(x)$$

$$\Rightarrow D_\mu \phi = e^{i\frac{\xi}{v}} \frac{1}{\sqrt{2}} [\partial_\mu \eta - ig B_\mu (v + \eta)] \rightarrow (D_\mu \phi)^* (D_\mu \phi) \text{ unabhängig von } \xi$$

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▶ Endergebnis:

$$\mathcal{L} = \frac{1}{2} (\partial_\mu \eta)^2 - \mu^2 \eta^2 - \frac{1}{4} (\partial_\mu B_\nu - \partial_\nu B_\mu)^2 + \frac{1}{2} g^2 v^2 B_\mu B^\mu + \mathcal{L}_{\text{WW}}$$

▶ Vektorboson mit Masse $M = gv$ Higgs-Mechanismus

▶ kein Goldstone-Boson

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▶ Vektorboson mit Masse $M = gv$ **Higgs-Mechanismus**

▶ kein Goldstone-Boson

Freiheitsgrade: 1 reelles Skalarfeld + 3 Polarisationen des Vektorbosons

→ „would-be-Goldstone-Boson“ wurde von den Vektorbosonen „gegessen“

► Systematischere Vorgehensweise:

Eliminiere Mischterm mittels Eichfixierung

$$\begin{aligned}\mathcal{L}_{gf} &= -\frac{1}{2\xi}(\partial_\mu A^\mu + \xi M\phi_2')^2, \quad \text{'t Hooft-Eichung, } M = gv, \\ &= -\frac{1}{2\xi}(\partial_\mu A^\mu)^2 \underbrace{-M\phi_2' \partial_\mu A^\mu}_{\text{hebt Mischterm weg}} - \frac{1}{2}\xi M^2 \phi_2'^2\end{aligned}$$

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► Eichabhängige Goldstone-Boson-Masse: $m'_2 = \sqrt{\xi}M$

- Landau-Eichung: $\xi = 0 \Rightarrow m'_2 = 0$
- Feynman-Eichung: $\xi = 1 \Rightarrow m'_2 = M$
- unitäre Eichung: $\xi \rightarrow \infty \Rightarrow m'_2 \rightarrow \infty \rightarrow \phi'_2$ entkoppelt