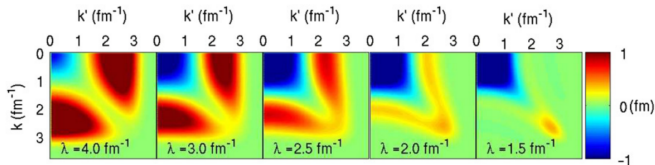


Similarity Renormalization Groups (SRG) for nuclear forces

Nuclear structure and nuclear astrophysics

Philipp Dijkstal

12.5.2016



Content

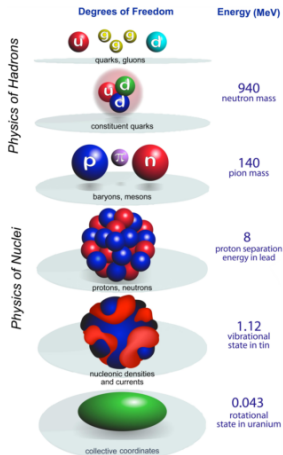
Nuclear Forces

SRG

Applications of SRG

Conclusion

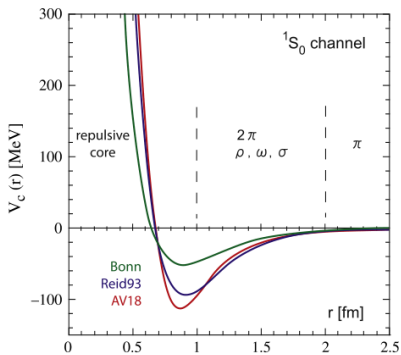
Resolution scale in nuclear physics



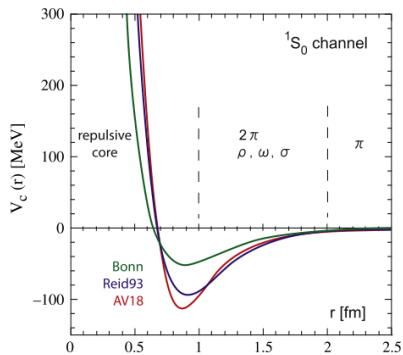
Separation of scales

- ▶ Systematic restriction on relevant degrees of freedom
 \hookrightarrow p, n
- ▶ Typical momenta within large nuclei: $200 \text{ MeV} \approx 1 \text{ fm}^{-1}$

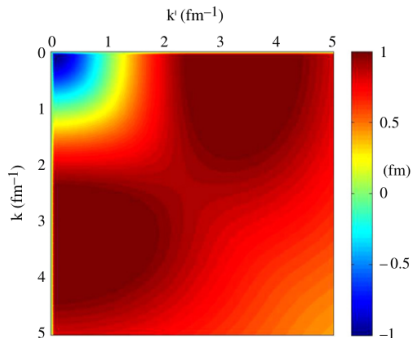
Phenomenological NN potentials (purely local)



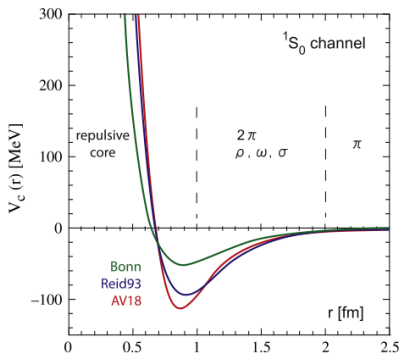
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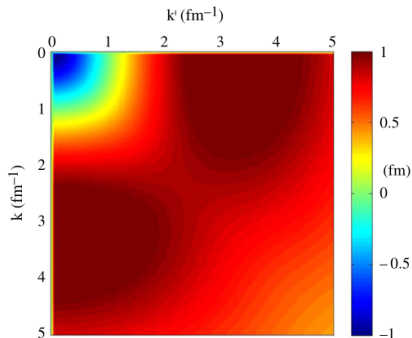
Bogner et. al., PPNP 65 (2010),94-147



Phenomenological NN potentials (purely local)



Bogner et. al., PPNP 65 (2010),94-147



Coupling in Lippmann-Schwinger equation for T matrix:

$$T_l(k, k', E) = V_l(k, k') + \frac{2}{\pi} \int q^2 dq \frac{V_l(k, q) T_l(q, k', E)}{E - q^2/2\mu + i\epsilon}$$

Similarity Renormalization Groups

Unitary transformation

$$\begin{aligned}U^\dagger U &= \mathbb{1} \\E &= \langle \Psi | H | \Psi \rangle \\&= (\langle \Psi | U^\dagger) U H U^\dagger (U | \Psi \rangle) \\&= \langle \tilde{\Psi} | \tilde{H} | \tilde{\Psi} \rangle\end{aligned}$$

Similarity Renormalization Groups

Unitary transformation

$$\begin{aligned}
 U^\dagger U &= \mathbb{1} \\
 E &= \langle \Psi | H | \Psi \rangle \\
 &= (\langle \Psi | U^\dagger) U H U^\dagger (U | \Psi \rangle) \\
 &= \langle \tilde{\Psi} | \tilde{H} | \tilde{\Psi} \rangle
 \end{aligned}$$

Desired properties of \tilde{H} :

- ▶ Dependent on resolution scale with parameter s or $\lambda = 1/s^{1/4}$
- ▶ Steady evolution

$$U = U_s, \quad \tilde{H} = H_s$$

$$U_0 = \mathbb{1}$$

Parameter flow equation

Differential equation to evolve V :

$$H_s = U_s H U_s^\dagger$$

Parameter flow equation

Differential equation to evolve V :

$$H_s = U_s H U_s^\dagger$$

$$\begin{aligned} \frac{dH_s}{ds} &= \frac{dH_s}{ds} = \frac{dU_s}{ds} H U_s^\dagger + U_s H \frac{dU_s^\dagger}{ds} \\ &= \frac{dU_s}{ds} \underbrace{U_s^\dagger U_s}_{\mathbb{1}} H U_s^\dagger + \underbrace{U_s H U_s^\dagger}_{H_s} U_s \frac{dU_s^\dagger}{ds} \\ &= \eta_s H_s + H_s \eta_s^\dagger = [\eta_s, H_s] \\ \eta_s &= \frac{dU_s}{ds} U_s^\dagger = -\eta_s^\dagger \end{aligned}$$

Parameter flow equation

Generator G_s for η_s :

$$\begin{aligned}\eta_s &= [G_s, H_s] \\ \frac{dH_s}{ds} &= [[G_s, H_s], H_s] \\ &= G_s H_s H_s - 2H_s G_s H_s + H_s H_s G_s\end{aligned}$$

Common choice: $G_s = T$ with $T|k\rangle = \epsilon_k|k\rangle$

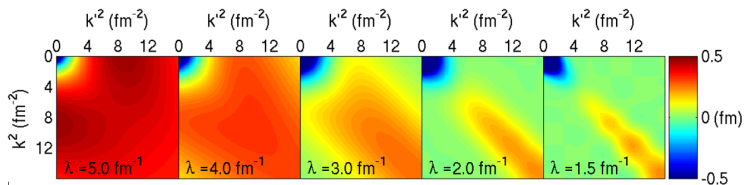
Parameter flow equation

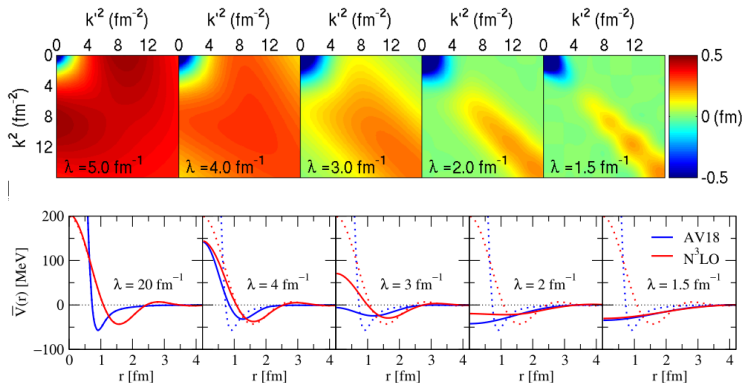
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Common choice: $G_s = T$ with $T|k\rangle = \epsilon_k|k\rangle$

$$\begin{aligned}\frac{dV_s}{ds} &= \frac{dH_s}{ds} = 2(TV_s T - V_s TV_s) + V_s^2 T + TV_s^2 - V_s T^2 - T^2 V_s \\ \frac{dV_s}{ds}(k, k') &= \langle k' | V_s | k \rangle \\ &= -(\epsilon_k'^2 - \epsilon_k)^2 V(k, k') \\ &\quad + \frac{2}{\pi} \int_0^\infty q^2 dq (\epsilon_k + \epsilon_{k'} - 2\epsilon_q) V_s(k, q) V_s(q, k')\end{aligned}$$

SRG Examples - Argonne ν_{18} 

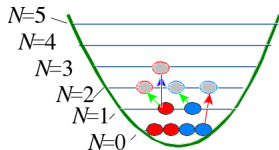
SRG Examples - Argonne ν_{18} 

Furnstahl, Nucl. Phys. B. (Proc. Suppl.) 228 (2012) 139-175

- ▶ 3S_1 channel
- ▶ Non-local evolved potentials

Matrix dimensions in nuclear physics calculations

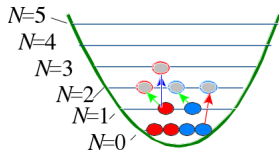
No Core Shell Model calculations



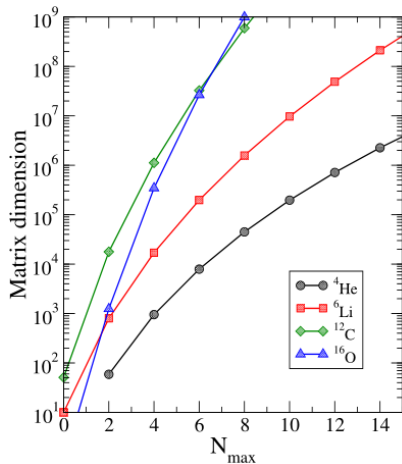
- ▶ H in HO basis with N_{\max} shells
- ▶ Numerical diagonalization
- ▶ Max matrix dim. $\approx 10^9$

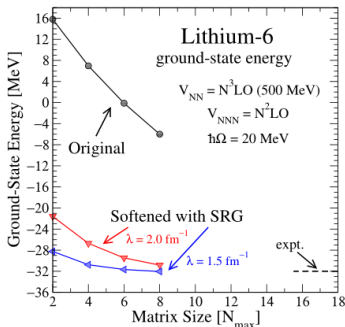
Matrix dimensions in nuclear physics calculations

No Core Shell Model calculations



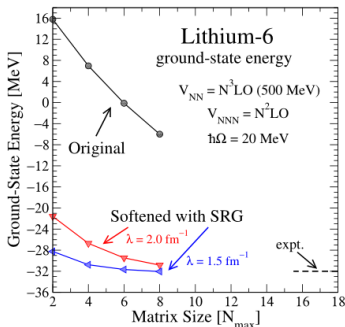
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SRG Examples - ${}^6\text{Li}$ 

Furnstahl, Nuc. Phys. B (Proc. Suppl.) 228 (2012) 139-175

- ▶ NCFC calculations
- ▶ Based of EFT potential
- ▶ Convergence possible with SRG

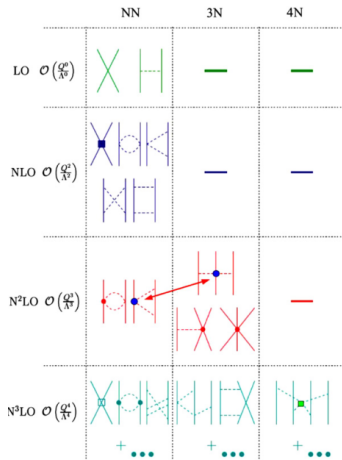
SRG Examples - ${}^6\text{Li}$ 

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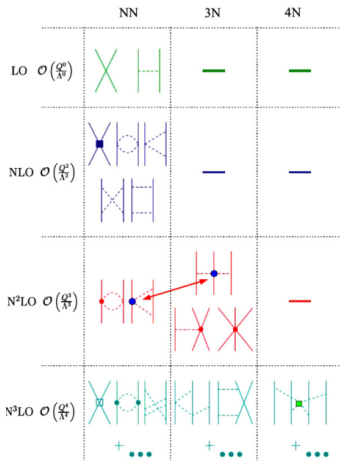
↔ 3-body forces?

Hierarchy of forces in chiral Effective Field Theory



Bogner et al., Progress in Particle and Nucl. Phys. 65
(2010) 94-147

Hierarchy of forces in chiral Effective Field Theory



► Chiral EFT: Pions are included

► Expansion parameter : $(\frac{Q}{\Lambda})$

↔ Multi-body forces arise naturally in EFT

↔ Consequence of neglecting higher order interaction processes

N-body forces from SRG

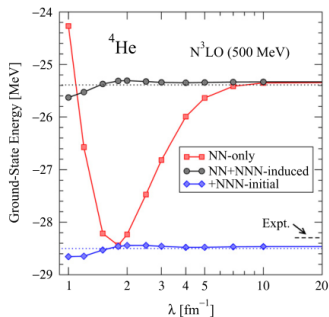
Why do 3-body forces and higher increase?

- ▶ Consider mixing terms in the 2nd quantized flow equation
- ▶ Original V_{NN} only contains 2-body forces
- ▶ T is one-body operator

$$\begin{aligned}
 \frac{dV_s}{ds} &= [[T, V_{NN}], H_{NN}] \\
 &= [[\sum a^\dagger a, \sum a^\dagger a^\dagger aa], \sum a^\dagger a^\dagger aa] \\
 &= \underbrace{\sum a^\dagger a^\dagger aa}_{2 \text{ body}} + \underbrace{\sum a^\dagger a^\dagger a^\dagger aaa}_{3 \text{ body}} + \dots
 \end{aligned}$$

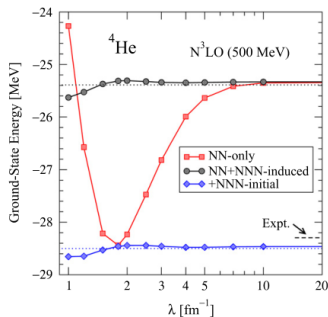
$$V_{s'} = V_s + \frac{dV_s}{ds} \cdot \Delta s$$

Importance of N-body forces



Furnstahl, Nuc. Phys. B (Proc. Suppl.) 228 (2012)
139-175

Importance of N-body forces



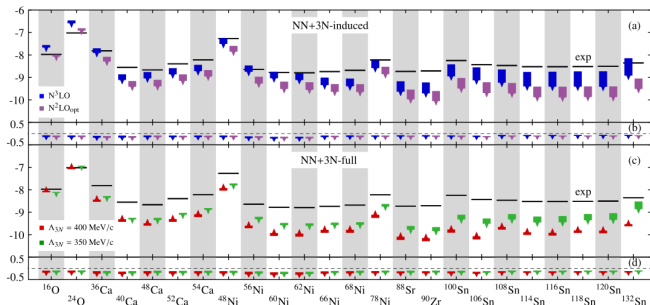
Furnstahl, Nuc. Phys. B (Proc. Suppl.) 228 (2012)
139-175

- ▶ Binding energy in ^4He
- ▶ Evolution from right to left
- ▶ **NN-only fails**

↪ 3-body forces and higher
increase from SRG

↪ 4-body contribution
significant

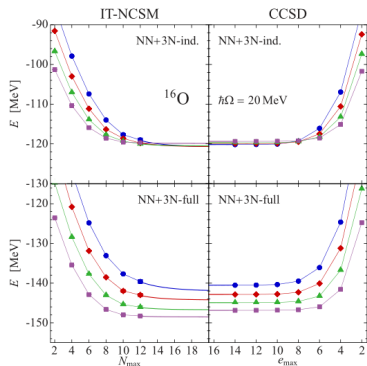
Heavy nuclei



Binder et al., Phys. Lett. B 736 (2014) 119-123

- ▶ Coupled Cluster
- ▶ Ground state energy per N
- ▶ (b) and (d): Triple correction

Comparison of CC and NCSM



Furnstahl and Hebeler, Rep. Prog. Phys. 76 (2013)

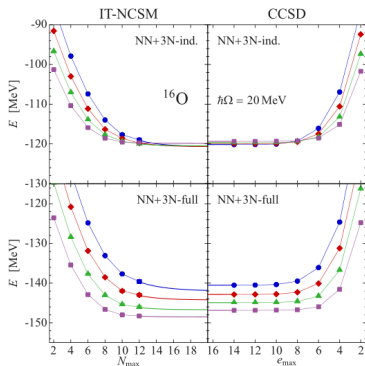
$$s = 0.04 \text{ fm}^4$$

$$s = 0.05 \text{ fm}^4$$

$$s = 0.06 \text{ fm}^4$$

$$s = 0.08 \text{ fm}^4$$

Comparison of CC and NCSM



Furnstahl and Hebeler, Rep. Prog. Phys. 76 (2013)

- ▶ CC completely different from NCSM
- ▶ Still great agreement

$$s = 0.04 \text{ fm}^4$$

$$s = 0.05 \text{ fm}^4$$

$$s = 0.06 \text{ fm}^4$$

$$s = 0.08 \text{ fm}^4$$

Conclusion

- ▶ Basics of low-energy nuclear physics
 - ▶ Separation of scales
 - ▶ Hierarchy of forces
- ▶ Phenomenological potentials
 - ▶ Repulsive core
 - ▶ coupling
- ▶ SRG
 - ▶ Flow equations
 - ▶ decoupling
 - ▶ Higher order forces
 - ▶ Applications
- ▶ Open Questions
 - ▶ 4-body forces
 - ▶ Alternative Generators