

Reminder 1: lectures & hw solutions are password protected. See Chat!!

Recordings (2 books)

lecture notes

2: wednesday help

lectures

lecture slides (after)

↳ See Chat.

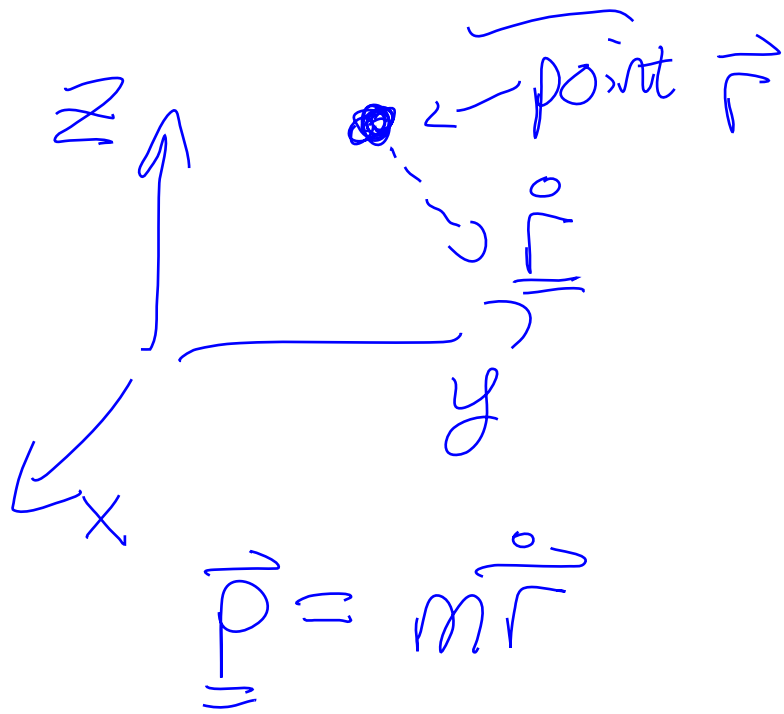
Recordings ← PW prot.

HW

Solutions ← PW prot.

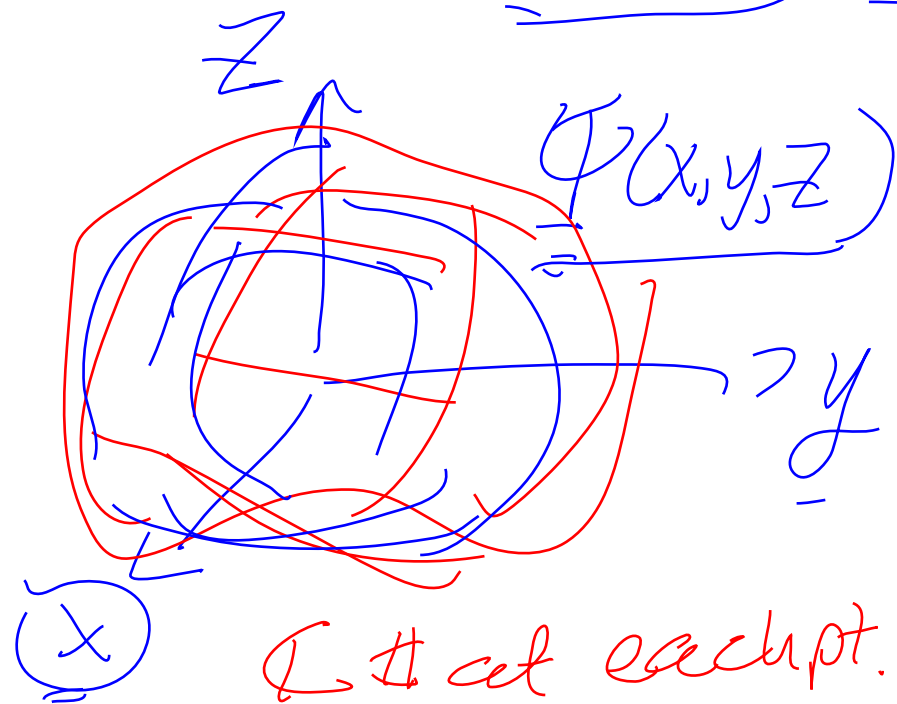
Quantum Field Theory

Classical



$$\vec{p} = m\vec{v}$$

Quantum Mech



\otimes at each pt. in sp.

Main: operator $\hat{p}_j = -i \frac{\partial}{\partial x_j}$

" \hat{p} info. is already stored in Ψ

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Classical Field Thry

State $\varphi(x)$ value at each pt in space

Canonical $\pi(x) = \dot{\varphi}(x)$
how fast φ is changing

Kronecker

$$\frac{\partial}{\partial x_i} x_j = \delta_{ij}$$

Dirac \neq

$$\int \varphi(y) = \int (x-y)$$

\mathbb{C} functional

Quantum Field Thry.

State-space:
all possible $\varphi(x)$ choices

$$\Psi[\varphi(x)]$$

Every poss. class field config, \mathbb{C} #

$$\pi(x) = -i \int \frac{\delta}{\delta \varphi(x)} \Psi[\varphi(x)]$$

Ψ is a function over

functions in contin space

2nd Quantization

Ask a Mathematician: Does this make sense?

Answer: no. Or we don't think so.

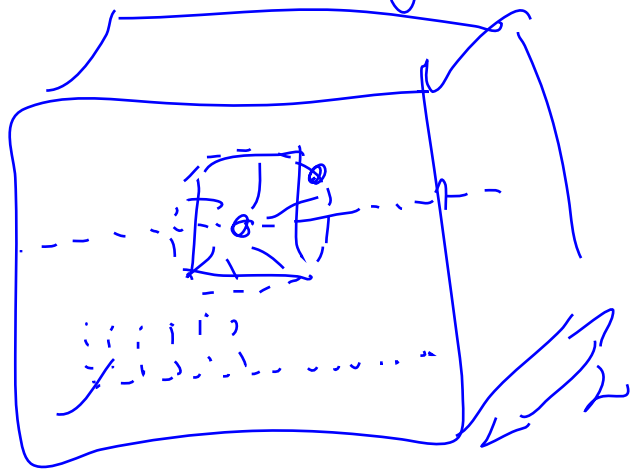
But: I took a crystal lattice: finite # of DOF (atoms)

Built a long-dist descrip. \rightarrow Field Thry
pretend contin space - and it worked.

Shortest probed dist? highest E : $\frac{10 \text{ TeV}}{hc} \rightarrow 10^{-20} \text{ m.}$

Assume space discrete 10^{-24} m (really: 10^{-35} m.)

(Regularization Regularization)



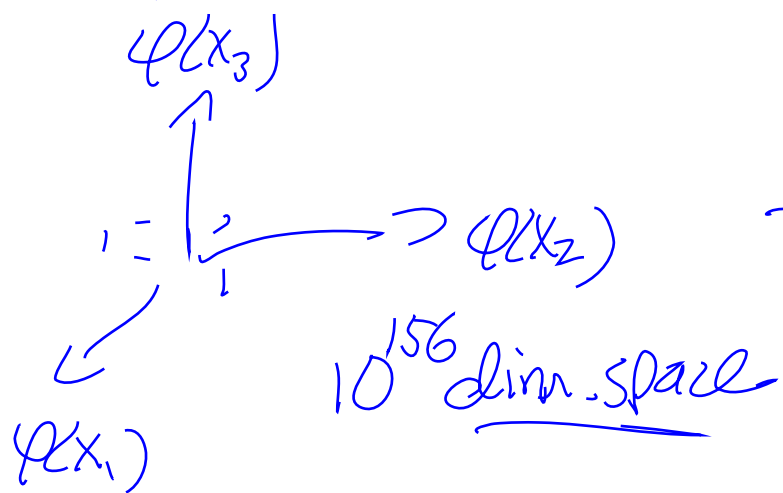
box, $L \times L \times L$ $L \approx 10^{12} \text{ light yrs}$
 $\approx 10^{28} \text{ m}$

NO consequence

10⁵² points across latt.

Grid with $(10^{52})^3 = \underline{10^{156}}$ points (or more)

$$\varphi(x) = (\varphi(x_1), \varphi(x_2), \varphi(x_3), \dots, \varphi(x_{10^{156}}))$$



Wave function on 10^{156} dim space.

Now it's well defined

Size L : $L/a > 10^{52}$...

Spacing a

How do answers depend on a ?

Take $\lim_{a \rightarrow 0}$. Version 3 - real version

To define ψ , need "cutoff" removing arb. short dist.
Hope/need/expect $a \rightarrow 0$ lim. works exists.
Totally tricky nontrivial Group.

Simplest theory: 1 real scalar field.

$\phi(x)$ func. over \vec{x} in 3+1 Dim (class.)

||
$$\mathcal{L} = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{m^2}{2} \phi^2$$

NO ϕ , ϕ^3 , ϕ^4 today!
 (can make $m^2=0$)

$$L = \int d^3x \mathcal{L} \rightarrow \underbrace{a^3}_{\substack{n_1 \\ n_2 \\ n_3 \\ \text{integers}}} \sum \left[\frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{m^2}{2} \phi^2 \right]$$



n_1
 n_2
 n_3 integers

huh?? $\partial_0 \phi = \dot{\phi}$ is OK

$$\vec{\nabla}_i \phi = \phi(x + a \hat{i}) - \phi(x)$$

$\vec{x} = a \vec{n}$
 vector of ints

space latt
time continuous

what is $\phi(x)$?? An operator (like x in QM)

$\langle \Phi | \phi(x) | \Psi \rangle$ or $|\Phi\rangle$

$$\langle \underbrace{\Phi}_{\text{state where } \Phi=0} | \phi(x) | \Psi \rangle = \langle \Phi | \phi(x) | \Psi \rangle$$

except on specific class. conf. where ϕ finite.

$$L = a^3 \sum_{\vec{n} \in \mathbb{Z}^3} \frac{1}{2} \dot{\varphi}^2(x) - \frac{|\nabla \varphi|^2}{2} - \frac{m^2 \varphi^2}{2}$$

$\vec{x} = a\vec{n}$ $\pi^2(x)$ canon momentum?? (class. \mathbb{R}^0)

$\pi(x)$ also an op. Class. $\pi(\vec{x}) = \frac{\partial L}{\partial \dot{\varphi}} = a^3 \frac{\partial}{\partial \dot{\varphi}} \varphi(\vec{x})$
 $= a^3 \vec{n} \cdot \nabla \varphi(\vec{x})$

$$H = \sum_{\vec{n}} \pi \dot{\varphi} - L = \sum_{\vec{n}} \left[\frac{a^3}{2} \pi(x) \pi(x) + \frac{a^3}{2} |\nabla \varphi|^2 + \frac{a^3 m^2}{2} \varphi^2(x) \right]$$

$$\hat{\pi}(\vec{x}) = -i \vec{\partial} \varphi(x)$$

in "standard" "coord" basis

$$\underline{\Psi}(\vec{\varphi}(x)) = \langle \underline{\vec{\varphi}}(x) | \underline{\Psi} \rangle$$

$$\langle \underline{\vec{\varphi}}(x) | \underline{\pi}(x) | \underline{\Psi} \rangle = -i \underline{\partial} \underline{\Psi}(\vec{\varphi}(x))$$

$$H = \int d^3x \left[\frac{1}{2} \dot{\pi}^2(x) + \frac{1}{2} (\nabla \hat{\varphi})^2(x) + \frac{m^2}{2} \hat{\varphi}^2(x) \right]$$

$\int_{x=\vec{n}a}$
 $n \in \mathbb{D}^3$

$$([\pi(x), \varphi(y)] = -i \delta^3(x-y))$$

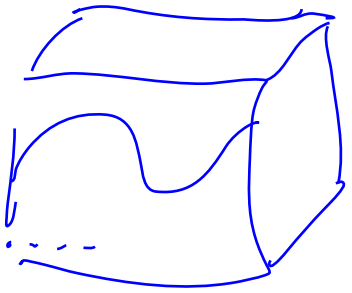
$$-i\rho$$

↑

H not diag in (φ, π) basis because $\nabla_i \varphi(x) = \frac{\varphi(x+a e_i) - \varphi(x)}{a}$

Fourier techniques $\tilde{\varphi}(\rho) = \int d^3x e^{+i\rho \cdot \vec{x}} \varphi(x)$

$\rho = \vec{n} \frac{2\pi}{L}$



$$\tilde{\varphi}(\rho)$$

$$\varphi(x) = \sum_{\rho = \vec{n} \frac{2\pi}{L}} \tilde{\varphi}(\rho) e^{+i\rho \cdot \vec{x}}$$

L, a finite

$$\vec{n}_{max} = \frac{L}{2a}$$

$\tilde{\pi}(\rho)$, similarly.

$\tilde{\varphi}(p)$ complex while $\varphi(x)$ real

$$\tilde{\varphi}(-p) = (\tilde{\varphi}(p))^*$$

$$\int d^3x \frac{\overline{\psi(x)} \psi(x)}{2} = \sum_{\mathbf{p}_1 = \vec{n}_1 \frac{2\pi}{L}} \sum_{\mathbf{p}_2 = \vec{n}_2 \frac{2\pi}{L}} L^{-6} \int d^3x e^{i\vec{p}_1 \cdot x} e^{i\vec{p}_2 \cdot x} \frac{\tilde{\psi}(\mathbf{p}_1) \tilde{\psi}(\mathbf{p}_2)}{2}$$

no x

$$= \sum_{\mathbf{p} = \vec{n} \frac{2\pi}{L}} L^{-3} \frac{\tilde{\psi}(\mathbf{p}) \tilde{\psi}^*(\mathbf{p})}{2}$$

$$\frac{\tilde{\psi}_r^2(\mathbf{p}) + \tilde{\psi}_i^2(\mathbf{p})}{2}$$

$$L^3 \int_{\mathbf{p}_1 + \mathbf{p}_2 = 0} d^3p_1 d^3p_2 \tilde{\psi}(\mathbf{p}_1) \tilde{\psi}(\mathbf{p}_2)$$

$$\tilde{\psi}(\mathbf{p}) = \tilde{\psi}_r(\mathbf{p}) + i\tilde{\psi}_i(\mathbf{p})$$

$$H = \sum_{\vec{p} = \vec{n} \frac{2\pi}{L}} L^{-3} \int \frac{\tilde{\pi}(\vec{p}) \tilde{\pi}(-\vec{p})}{2} + \underbrace{\left(m^2 + \frac{4}{c^2} \sin^2\left(\frac{p_0 a}{2}\right) \right)}_{\sum_{\vec{l}} \frac{4}{c^2} \sin^2\left(\frac{\vec{l} \cdot \vec{p} a}{2}\right) = \vec{p}^2} \tilde{\varphi}(\vec{p}) \tilde{\varphi}(-\vec{p})$$

SHO's? Yes! $\omega_{\vec{p}} = \sqrt{p^2 + m^2}$ $\omega_{\vec{p}} = |\vec{p}|$ if $m=0$

$$a_{\vec{p}} = \frac{1}{\sqrt{2L^3}} \left(\sqrt{\omega_{\vec{p}}} \tilde{\varphi}(\vec{p}) + \frac{i}{\sqrt{\omega_{\vec{p}}}} \tilde{\pi}(\vec{p}) \right)$$

note $\varphi^*(-\vec{p}) = \varphi(\vec{p})$


$$a_{\vec{p}}^\dagger = \frac{1}{\sqrt{2L^3}} \left(\sqrt{\omega_{\vec{p}}} \tilde{\varphi}(-\vec{p}) - \frac{i}{\sqrt{\omega_{\vec{p}}}} \tilde{\pi}(-\vec{p}) \right)$$

really is \neq of $a_{\vec{p}}$

Much work: $[a_{\vec{p}}^\dagger, a_{\vec{q}}] = \int_{\vec{p}, \vec{q}} \delta_{\vec{p}, -\vec{q}}$ Kronecker

$$H = \sum_{\mathbf{p} = \vec{n} \frac{2\pi}{L}} \left(\frac{a_{\mathbf{p}} a_{\mathbf{p}}^\dagger}{2} + \frac{a_{-\mathbf{p}}^\dagger a_{-\mathbf{p}}}{2} \right) \omega_{\mathbf{p}} = \sum_{\mathbf{p} = \vec{n} \frac{2\pi}{L}} \left(a_{\mathbf{p}}^\dagger a_{\mathbf{p}} + \frac{1}{2} \right) \omega_{\mathbf{p}}$$

10¹⁵⁶ independent simple harm. osc.

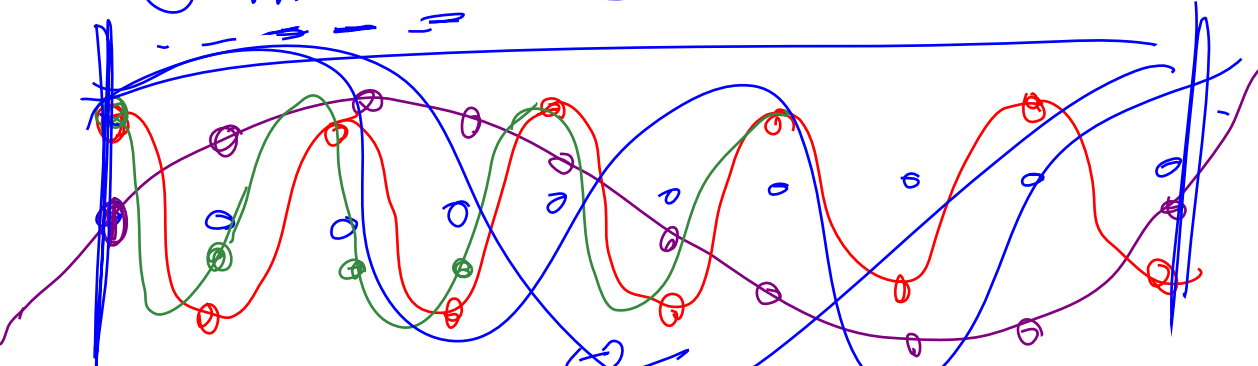
If $a \rightarrow 0$ I can - \vec{n} sum runs over \mathbb{Z}^3
 no just a box 

$$\tilde{\varphi}(\mathbf{p}) = \sqrt{\frac{L^3}{2\omega_{\mathbf{p}}}} (a_{\mathbf{p}} + a_{-\mathbf{p}}^\dagger)$$

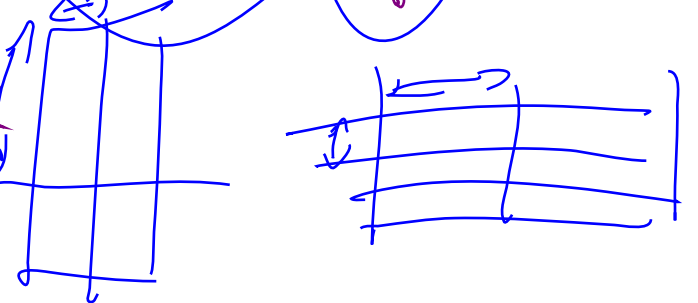
$$\varphi(\mathbf{x}) = \sum_{\mathbf{p} = \vec{n} \frac{2\pi}{L}} \frac{\sqrt{\omega_{\mathbf{p}}}}{\sqrt{2\omega_{\mathbf{p}} L^3}} \left(e^{i\vec{p} \cdot \vec{x}} a_{\mathbf{p}} + e^{-i\vec{p} \cdot \vec{x}} a_{\mathbf{p}}^\dagger \right)$$

$\varphi(\mathbf{x})$ both hermitian

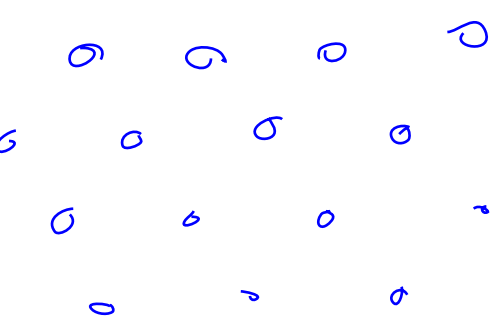
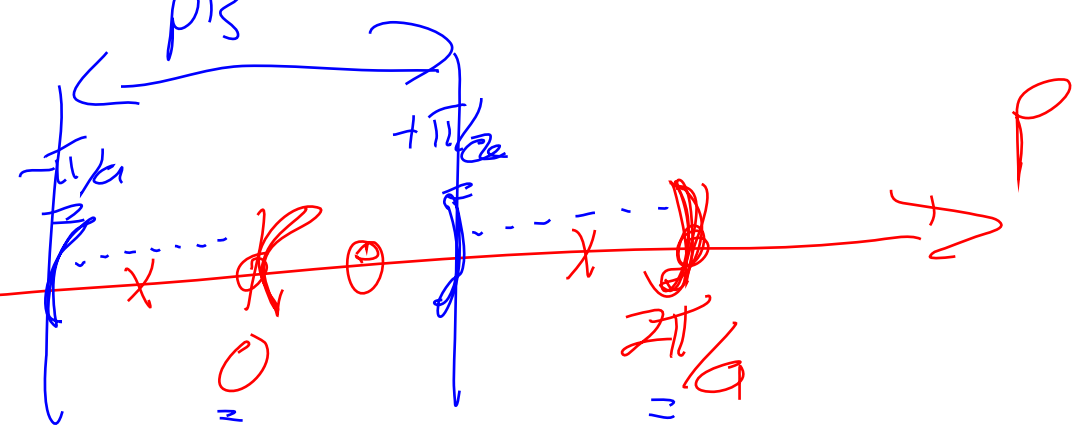
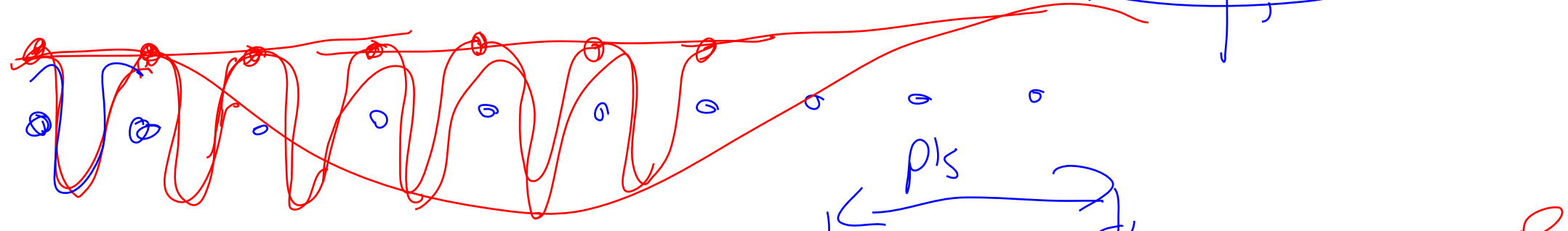
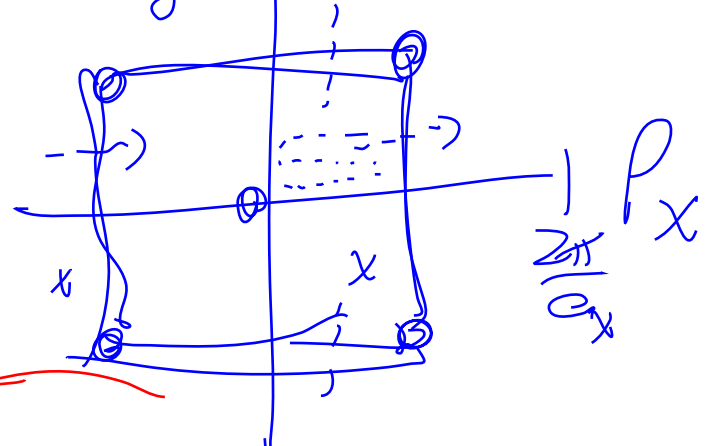
Brillouin Zone



$p \ll \frac{1}{a}$ is OK



\vec{p} is same as $p + \frac{2\pi}{a}$



finite box
 ↳ discrete p's
 discrete box
 ↳ finite p-region

