

Free  $\varphi$  theory  $\mathcal{L}(\varphi, d\varphi) = \frac{1}{2} \partial_\mu \varphi \partial^\mu \varphi - \frac{m^2}{2} \varphi^2$

$|n_1 n_2 n_3 \dots\rangle$  generic

occ. of SHO's  $P_1 P_2$  all possible modes in box

$+E_0$

Vacuum  $|0 0 0 0 \dots\rangle = |0\rangle$

$|0 0 0 0 1 0 0 0 \dots\rangle = |p\rangle$  1 in "p" SHO

↑ which p.

$|0 0 \dots 1 0 0 0 \dots\rangle = |P_1 P_2\rangle$

$\hat{H} |0\rangle = E_0 |0\rangle = \sum_p \left( a_p^\dagger a_p + \frac{1}{2} \right) \omega_p |0\rangle = \left( \frac{1}{2} \sum_p \omega_p \right) |0\rangle$

$E_0 \sim \frac{L^3}{(2\pi)^3} \int d^3k \frac{c \hbar k}{2}$

Actual physical value  $E_0 = L^3 (2.4 \text{ meV})^4$   
Cosmological Const. Problem

$$\text{II } \mathcal{L} = \frac{1}{2} \partial_\mu \varphi \partial^\mu \varphi - \frac{m^2}{2} \varphi^2 - \frac{\lambda}{24} \varphi^4 \quad \lambda \text{ pure \# (or } h^{-1}\text{)}$$

$$\mathcal{L}(\varphi) = \mathcal{L}(\varphi) + E_0$$

- Is this also a bunch of Sturm. Osc.? NO Is there vacuum  $|\Omega\rangle$ ?
- Are there particles? Probably - See next Q. yes but
- Does this theory exist? It's complicated.

If  $\lambda$  small, exists if  $|a \neq 0|$  even if  $a$  is supersuper small. But formal  $a \rightarrow 0$  limit is a problem:

$\lambda \rightarrow 0$  for deep reasons

$$\lambda < \frac{24}{\log(\bar{a}')} \quad \text{with } \bar{a}'$$

.....

How to figure this out?

$$\underbrace{\varphi(x)}_{=} |0\rangle = \underbrace{|\varphi(x)\rangle}_{=}$$

What is this state like?

Ask about amplitudes

$$\langle \varphi(y) | \varphi(x) \rangle$$

$$\langle 0 | \varphi(x) | 0 \rangle = 0$$

$$\underline{G^2}(y-x) = \langle 0 | \varphi(y) \varphi(x) | 0 \rangle$$

Translation inv:  $\checkmark$   
 $\langle 0 | \varphi(y+\xi) \varphi(x+\xi) | 0 \rangle$   
 $= \langle 0 | \varphi(y) \varphi(x) | 0 \rangle^*$

In free thy, this func. tells us about particles.  
 $\underline{G^2}(p) = 2\pi \delta(\underline{p}^2 - m^2) \Theta(\underline{p}^0)$

Pick  $\xi = -x$

Shows th $\checkmark$  has particles mass  $m$

What about in general?

We have  $p^\mu$  as conserved currents.  $p^\mu = \int d^3x T^{0\mu}(x)$

$$\frac{d}{dt} p^\mu = 0 \quad [\hat{p}^\mu, \hat{H}] = 0 \quad [\hat{p}^\mu, \hat{p}^\nu] = 0$$

Translations commute.

$p^\mu$  are set of Op's which commute with  $H$  & each other.

Seek simultaneous eig. states of  $p^\mu$ .

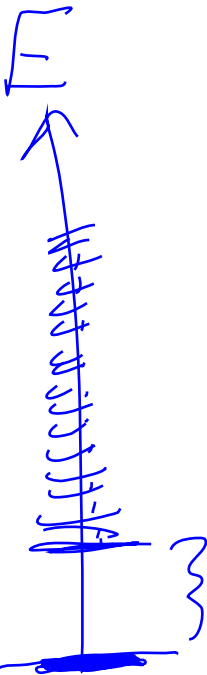
$$\hat{p}^\mu |n\rangle = p_n^\mu |n\rangle \quad \text{eigenstate} \sim \text{eig. val. } p_n^\mu$$

Basis of states  $\mathbb{1} = \sum_n |n\rangle\langle n|$

Choose  $p^\mu |0\rangle = 0 |0\rangle$

Assume Mass Gap: next-smallest  $E$  is  $M$ , with  $M > 0$

Mass Gap



$$G(\vec{y}) = \sum_n \langle 0 | \varphi(\vec{y}) | n \rangle \langle n | \varphi(0) | 0 \rangle$$

$$\varphi(\vec{y}, t) = e^{+iHt} \varphi(\vec{y}) e^{-iHt} \quad H = P^0 \text{ gen. of time translations}$$

$$\varphi(\vec{y}) = e^{-i\vec{y} \cdot \vec{P}} \varphi(0) e^{+i\vec{y} \cdot \vec{P}} \quad \vec{P} \text{ operator gen. of time-trans.}$$

$$\varphi(y^\mu) = e^{+i\hat{P}_\mu y^\mu} \varphi_0 e^{-i\hat{P}_\mu y^\mu}$$

$$G(\vec{y}) = \sum_n \langle 0 | e^{+i\hat{P}_\mu y^\mu} \varphi(0) e^{-i\hat{P}_\mu y^\mu} | n \rangle \langle n | \varphi(0) | 0 \rangle$$

$$G(\vec{y}) = \sum_n e^{+i0} e^{-i\hat{P}_\mu y^\mu} \underbrace{\langle 0 | \varphi(0) | n \rangle}_{C_n^*} \underbrace{\langle n | \varphi(0) | 0 \rangle}_{C_n} = \sum_n C_n^* C_n e^{-i\hat{P}_\mu y^\mu}$$

Hook?? Interpret answer??

Write in momentum space

Note  
for any  
 $p_n^{\mu}$

$$1 = \int \frac{ds}{2\pi} \int \frac{d^4 p}{(2\pi)^4} (2\pi) \delta(p^2 - s) \delta(p) (2\pi)^4 \delta(p^\mu - p_n^\mu)$$

$$p_n^\mu p_{n\mu} \geq M^2$$

$$G^2(y) = \int \frac{ds}{2\pi} \int \frac{d^4 p}{(2\pi)^4} e^{-i p_n y^\mu} \delta(p) \delta(p^2 - s) \delta(p^\mu - p_n^\mu) C_n C_n^*$$

$G^2(p)$

$\parallel$   
 $\mathcal{O}(p^2)$

Depends on  $p^\mu$  but only through  $p^2$

$p^\mu$  and  $\Lambda^\mu_{\nu} p^\nu$  gave diff't answers  
 $\rightarrow$  break ~~the~~ Lorentz inv.

$$\underline{G^2(y)} = \int \frac{d^4 p}{(2\pi)^4} e^{-i p \cdot y}$$

$$\underline{G^2(p)} \rightarrow \text{Feynman: } 2\pi \int (\rho^2 - m^2) \Theta(\rho^0)$$

$$\hookrightarrow \int_{m^2}^{\infty} ds \int (\rho^2 - s) \Theta(\rho^0) \underline{\underline{\rho(s)}}$$

$$= \int \frac{ds}{(2\pi)} \rho(s) \underline{\underline{\Delta(x,s)}}$$

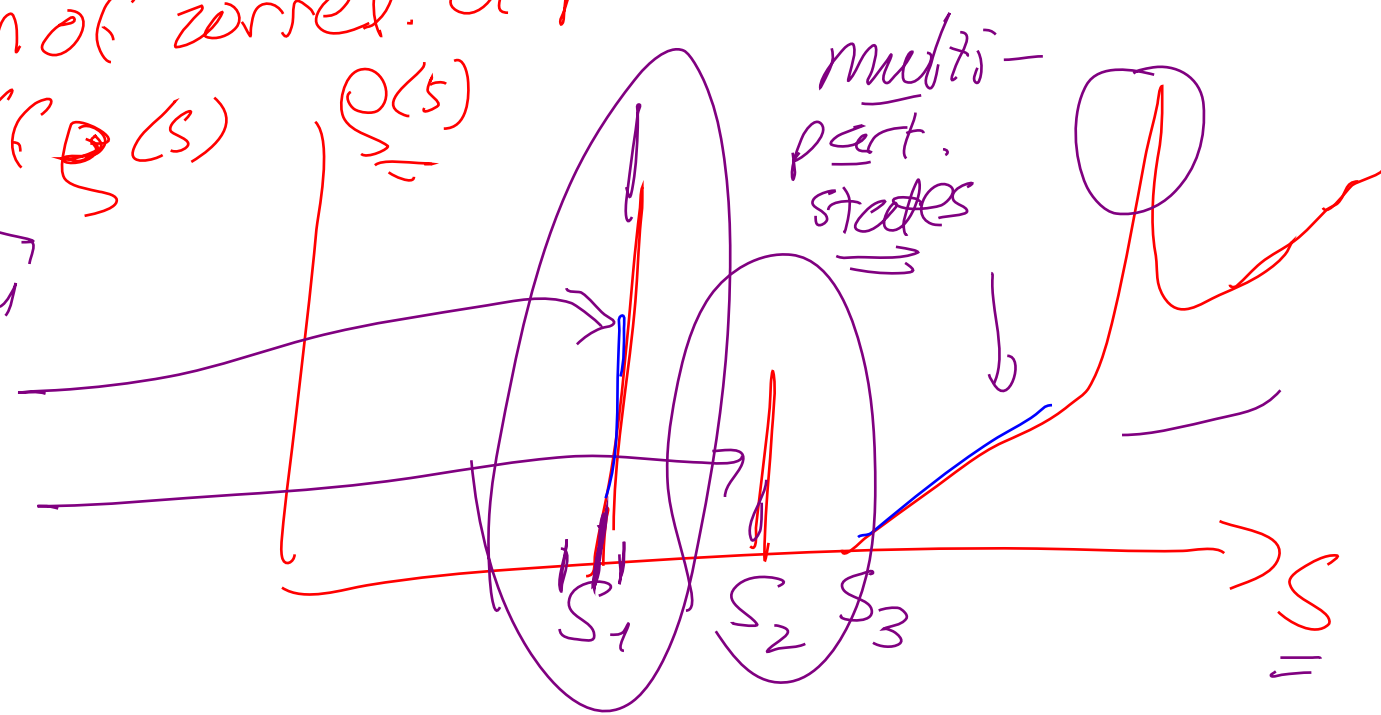
↑ Feynman prop. from last time  
but with  $m^2 \rightarrow s$

Correl. func.  $\rightarrow$  sum of correl. of parts w. mass  $^2 = s$

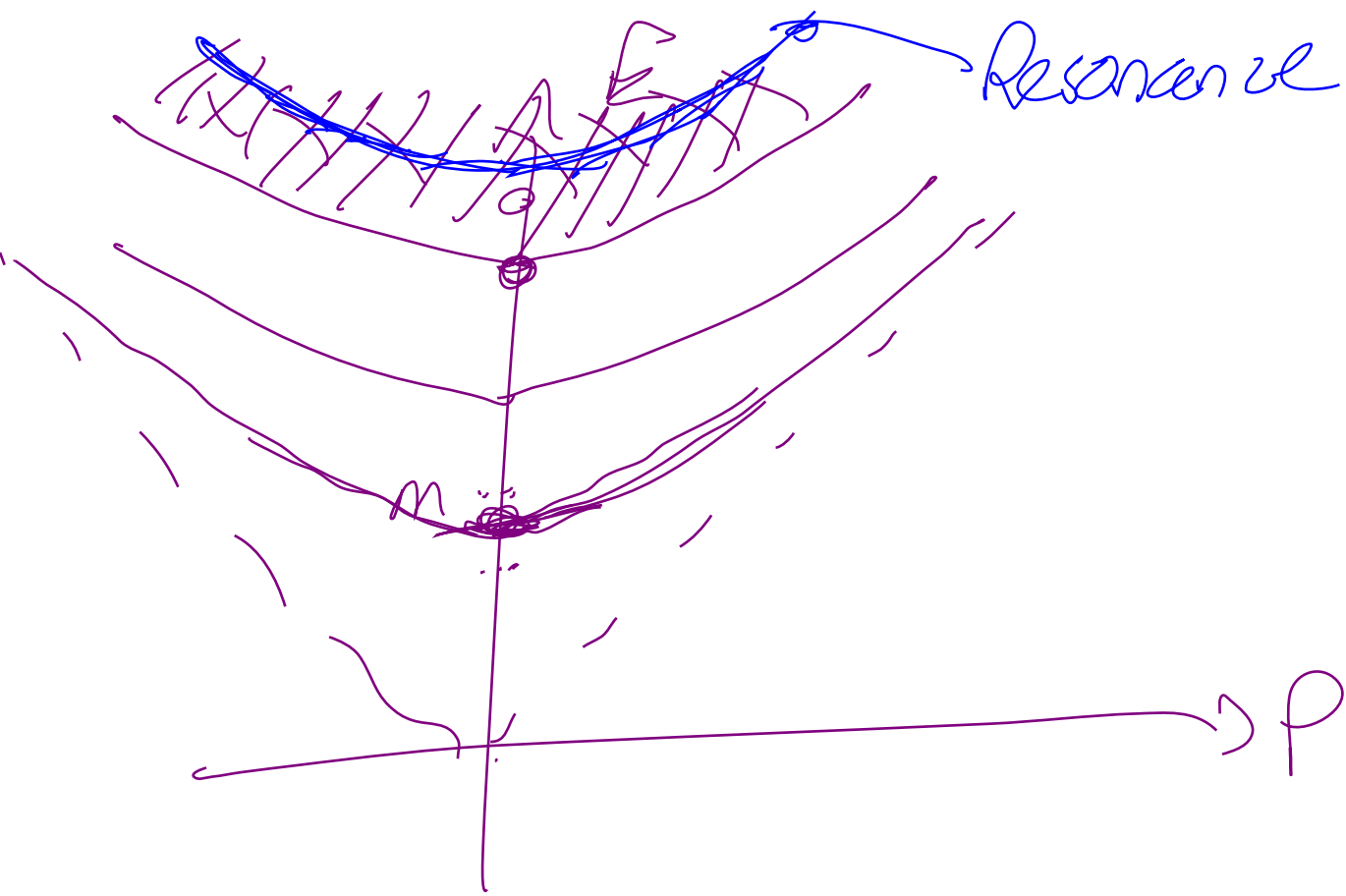
integr. over  $s$  with weight  $\rho(s)$

Particle  $m^2 = s_1$  or  $m = \sqrt{s_1}$

Part.  $m^2 = s_2$



Resonance





Correl. func's

$$\underline{\underline{G^2(y)}} = \int \frac{d^4 p}{(2\pi)^4} e^{-i p \cdot y} G^2(p)$$

$$G^2(p) = \int \frac{ds}{2\pi} \Theta(p^0) 2\pi \delta(p^2 - s) \underline{\underline{\rho(s)}}$$

theory info

Lehmann Represent. of propagator

$$\rho(s) = \sum_n \delta(p_n^2 - s) \underbrace{C_n^* C_n}_{\text{positive def.}}$$

$\rho(s)$  is positive  
(semi)def.

~~The~~ stuff prop. as sum of prop. rules for part's.  
with positive weights.

Other interesting correl. funcs:

$$G^<(x) = G^>(-x) = \langle 0 | \varphi(0) \varphi(x) | 0 \rangle$$

Answer =  $G^>$  answer but with  $\Theta(-x^0)$

$$G^>(x) \rightsquigarrow K_1(rm) \text{ spacelike} \quad G^< \rightsquigarrow K_1(rm)$$

$$H_1(tm) \text{ timelike} \quad \underline{H_1^*}$$

$$\underline{D}(x) = G^>(x) - G^<(x) = \langle 0 | \underline{[\varphi(x), \varphi(0)]} | 0 \rangle$$

Knows about causality:

$$D(x) : x^2 < 0 \Rightarrow 0$$

spacelike

$$\underline{G_R}(x) \text{ Retarded Func.} \Rightarrow D(x) \Theta(x^0)$$



$$\mathcal{L} \rightarrow \mathcal{L} - \underline{\underline{\varphi(x) J(x)}}$$

$\underline{\underline{J(x)}}$  is "external"  
you choose it func.

Modify in this way!

Interaction picture  $-\underline{\underline{\varphi J}}$  Int., regard  $\mathcal{L}$  as  $\underline{\underline{\mathcal{L}_0}}$

$$\underline{\underline{H = H_0 + \varphi J}}$$

What is  $\langle 0 | e^{+iH(t-t_0)} \underline{\underline{\varphi(y)}}$   $e^{-iH(t-t_0)} | 0 \rangle$

distant past distant past

to linear order  
in  $\underline{\underline{J}}$

$$\mathcal{U}(t, t_0) \approx 1 - i \int_{t_0}^t dt' \int d^3x J(x, t') \varphi(x, t')$$

$$\mathcal{U}^{\dagger} \approx 1 + \dots$$

Linear order

$$\langle 0 | \phi(x) \phi(y) | 0 \rangle = \langle 0 | \left[ 1 + i \int d^4x' \cancel{T(x')} \phi(x') \right] \phi(y) \left[ 1 - i \int d^4x' \cancel{T(x')} \phi(x') \right] | 0 \rangle$$

$$\langle \phi \rangle = \langle 0 | \phi | 0 \rangle$$

$$= \langle 0 | \int d^4x T(x) \delta(y-x^0) (i\phi(x)\phi(y) - i\phi(y)\phi(x)) | 0 \rangle$$

$$= i \int d^4x T(x) G_R(y-x)$$

New  $G_R$  includes this  $i$

$G^>$ : are fields correlated?

$G_R$ : if you "put" on field, how does it respond?

$\square$ : include causality

What's  $\underline{G_R(p^\mu)}$ ? =  $\int_{\underline{=}} d^4x \underline{e}^{i p_\mu x^\mu} \underline{G_R(x)}$

Note: as  $x^0 > 0$ , I am allowed to take  $p^0$  to have  $\geq \underline{Im}$

Im  $p^0 > 0$  not just  $\geq 0$ .

Find in Free Thy!  $\langle 0 | \varphi(y) \varphi(0) | 0 \rangle = \int d^4q e^{i p_\mu y^\mu} \frac{(\theta(q^0) - \theta(-q^0))}{2\pi\delta(q^2 - m^2)} - \varphi(0) \varphi(y)$

$$G_R(p) = \int d^4y \theta(y^0) e^{i p_\mu y^\mu} \int \frac{d^4q}{(2\pi)^4} \left[ -i 2\pi\delta(q^2 - m^2) (\theta(q^0) - \theta(-q^0)) \right] e^{i q_\mu y^\mu}$$

$\int d^4y \rightarrow \int dt d^3\vec{y}$   
 $e^{i p_\mu y^\mu} \rightarrow e^{i p^0 y^0} e^{-i \vec{p} \cdot \vec{y}}$   
 $\int \frac{d^4q}{(2\pi)^4} \rightarrow \int d^3\vec{q} \frac{dq^0}{2\pi}$

$e^{-i q^0 y^0} e^{i \vec{q} \cdot \vec{y}}$

$(2\pi)^3 \int^3 (\vec{p} - \vec{q}) \vec{q} \rightarrow \vec{p}$

$$G_R(p) = \int dy^0 \underline{\Theta}(y^0) \int \frac{d\vec{y}^0}{2\pi} (-2\pi i) e^{i(p^0 - q^0)y^0} (\underline{\Theta}(q^0) - \underline{\Theta}(-q^0))$$

$$\omega_p^2 = \vec{p}^2 + m^2$$

$$\int_0^\infty dy^0 \underline{\Theta}(y^0) (-2\pi i) \left[ e^{-\frac{i(p^0 - \omega_p)y^0}{2\omega_p}} - e^{-\frac{i(p^0 + \omega_p)y^0}{2\omega_p}} \right] \frac{1}{2\omega_p}$$

$$\int_0^\infty dy^0 e^{-Ay^0} = \frac{1}{A}$$

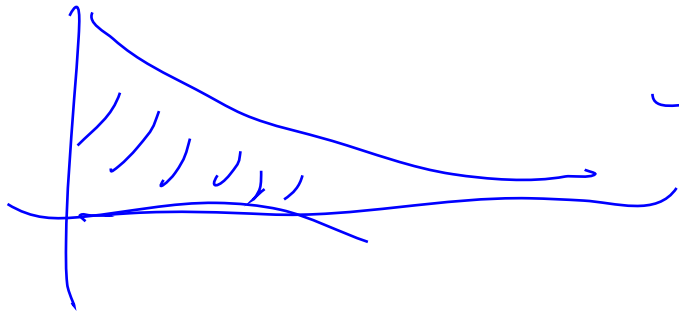
$$\frac{(-i)}{2\omega_p} \left[ \frac{-i}{i(\omega_p - p^0)} + \frac{-i}{i(\omega_p + p^0)} \right]$$

$$\frac{(-1)}{2\omega_p} \left( \frac{\omega_p + p^0 + \omega_p - p^0}{\omega_p^2 - p^0^2} \right) = \frac{2\omega_p}{2\omega_p(p^0^2 - \omega_p^2)} = \frac{1}{p^2 - m^2}$$

$$\underline{G_2(p^\mu)} = \frac{1}{p^2 - m^2} = \frac{1}{(p_0)^2 - \vec{p}^2 - m^2}$$

For  $p^0$  with  $\pm$  in part

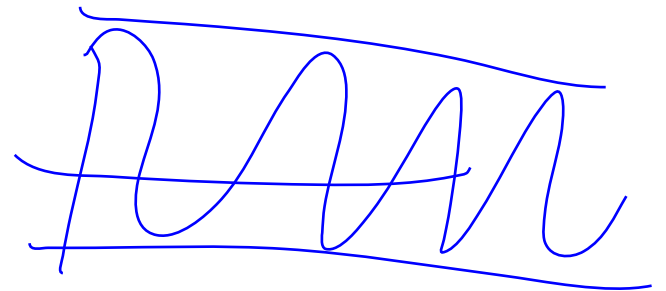
$$\int e^{-At}$$



$$\int e^{-At} e^{ipt}$$



$$\int e^{ipt}$$



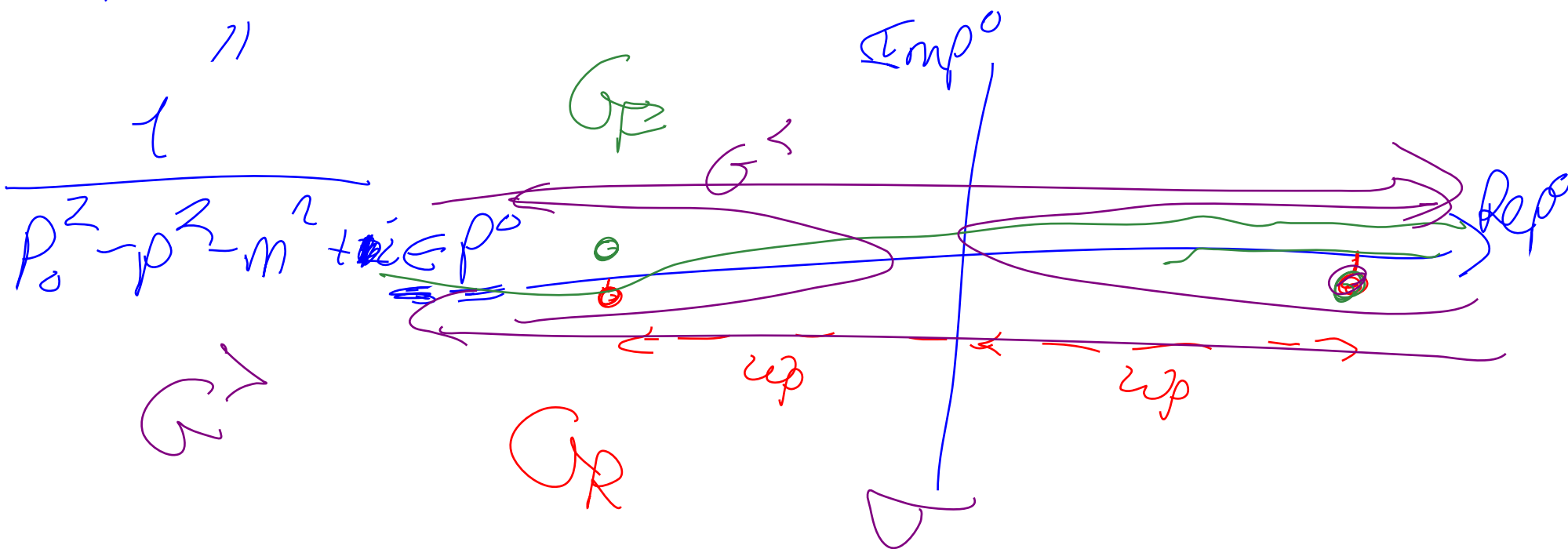
If  $p^0$  has  $\pm$  in part ✓

IF not  $\lim_{\epsilon \rightarrow 0} \frac{1}{(p_0 + i\epsilon)^2 - \vec{p}^2 - m^2}$

$$G_F(x) = -i \langle 0 | \left( \underbrace{\varphi(x) \varphi(0)}_{=} \oplus \underbrace{(x^0)}_{=} + \underbrace{\varphi(0) \varphi(x)}_{=} \oplus \underbrace{(-x^0)}_{=} \right) | 0 \rangle$$

T-ordered correl. func.

$$G_R = \frac{1}{(p_0 + i\epsilon)^2 - p^2 - m^2} \quad G_F = \frac{1}{p_0^2 - p^2 - m^2 + i\epsilon}$$





$G_R(p)$



$$= \int_{-\infty}^{\infty} \frac{ds}{2\pi} \rho(s)$$



$$\frac{1}{(p_0 + i\epsilon)^2 - p^2 - s}$$

