

Dirac Equation

Last remarks about  $\gamma$ -matrices

- 1) Sometimes I only have 1 2-comp. field (Standard Model)  
(neutrinos)

Then define  $\Psi = \begin{bmatrix} \psi_2 \\ \psi_1 \end{bmatrix}$  same field top & bottom

Some possible terms now equivalent, must be careful. Majorana  
(Notation) (field)

- 2) I am always free to make

Similarity Transformation  $\Psi \rightarrow S\Psi$   
(basis change)

$$\bar{\Psi} \rightarrow \bar{\Psi} \cancel{S^\dagger} \cancel{\gamma^0} \Psi S^{-1}$$

$$\gamma^\mu \rightarrow S \gamma^\mu S^{-1} \quad S \text{ unitary } 4 \times 4$$

gives all same physics.

Our basis: Chiral basis - best for very relativistic settings

Other bases better, eg, for  $v \ll 1$ .

In all cases,  $\{\gamma^\mu, \gamma^\nu\} = 2g^{\mu\nu} \mathbb{1}_{4 \times 4}$  "Clifford Algebra"

enough do derive all Lorentz properties.

Let's build a Lagrangian. Consider 1 Dirac field

L16 P2

$$\mathcal{L} = \underbrace{-m \bar{\psi} \psi}_{\text{scalar}} + i \underbrace{\bar{\psi} \gamma^\mu \partial_\mu \psi}_{\text{as } \bar{\psi} \gamma^\mu \psi \text{ 4-vector, cancels 4-vector index of } \partial_\mu}$$

Hermitian?  $(\psi)^\dagger = \psi^\dagger$ , Recall,  $\bar{\psi} = \psi^\dagger \gamma^0$ , Oh yes,  $\gamma^{0\dagger} = \gamma^0$

$$(\bar{\psi})^\dagger = \gamma^{0\dagger} \psi = \gamma^0 \psi \quad \gamma^{i\dagger} = -\gamma^i$$

$$(\bar{\psi} \psi)^\dagger = (\psi^\dagger \gamma^0 \psi)^\dagger$$

$$\psi^\dagger \cancel{\gamma^0} \psi = \bar{\psi} \psi \text{ is Hermitian } \checkmark$$

(in general  $\gamma^0 \gamma^{\mu\dagger} = \gamma^\mu \gamma^0$ )

$$\gamma^{\mu\dagger} \gamma^0 = \gamma^0 \gamma^\mu$$

$$\bar{\psi} \gamma^\mu \psi = \psi^\dagger \gamma^0 \gamma^\mu \psi$$

$$(\bar{\psi} \gamma^\mu \psi)^\dagger = \psi^\dagger \gamma^{\mu\dagger} \gamma^0 \psi = \psi^\dagger \gamma^0 \gamma^\mu \psi = \bar{\psi} \gamma^\mu \psi \text{ Hermitian}$$

But

$$A \partial_\mu B = -B \partial_\mu A + \underbrace{\partial_\mu (AB)}_{\text{total deriv.}}$$

$$(\bar{\psi} \gamma^\mu \partial_\mu \psi)^\dagger = (\partial_\mu \bar{\psi}) \gamma^\mu \psi = -\bar{\psi} \gamma^\mu \partial_\mu \psi + \text{tot deriv.}$$

Hence I need an  $i$  in that part of  $\mathcal{L}$ .

Technically I should probably write  $\mathcal{L} = \frac{i}{2} (\bar{\psi} \gamma^\mu \partial_\mu \psi - (\partial_\mu \bar{\psi}) \gamma^\mu \psi)$

$$\equiv i \bar{\psi} \gamma^\mu \overleftrightarrow{\partial}_\mu \psi$$

$$A \overleftrightarrow{\partial}_\mu B \equiv \frac{1}{2} (A \partial_\mu B - (\partial_\mu A) B)$$

Some people do this & it can matter in curved space. But let's not get fussy.

It could be  $-m^2 \bar{\psi} \psi + A i \bar{\psi} \gamma^\mu \partial_\mu \psi$  but then  $\psi_{\text{new}} = \sqrt{A} \psi$

$$m_{\text{new}} = m/A \dots \text{get rid of } A.$$

Dimensions:  $[\psi] = 3/2$

Interacting theory? Since  $\{4\} = 3/2$ , not much room to build dim  $\leq 4$  objects!

$$\mathcal{L} = -m\bar{\psi}\psi + i\bar{\psi}\gamma^\mu\partial_\mu\psi + y\bar{\psi}\psi\phi + eA_\mu\bar{\psi}\gamma^\mu\psi$$

appears to be only possibilities (for one  $\psi$ . Could have, eg,  $\phi_a\bar{\psi}_i T_{ij}^a\psi_j, A_\mu^a\bar{\psi}_i T_{ij}^a\psi_j$  for multi-imp  $\phi, A, \psi$ )

$y\bar{\psi}\psi\phi$ : "Yukawa coupling"

$eA_\mu\bar{\psi}\gamma^\mu\psi$ : "gauge coupling" same as  $i\bar{\psi}\gamma^\mu(\partial_\mu - ieA_\mu)\psi$  turns out to be better way to understand it.

Let's look at free theory and ask (possibly stupid) question

"What's classical field dynamics of  $\psi$  field??"

$$\mathcal{L} = -m\bar{\psi}\psi + i\bar{\psi}\gamma^\mu\partial_\mu\psi$$

Apply  $\frac{\delta}{\delta\psi(x)}$   $S=0 \implies -m\psi + i\gamma^\mu\partial_\mu\psi = 0$  Dirac Equation  
 $(i\gamma^\mu\partial_\mu - m)\psi = 0$  what are solutions?

Well they are also solutions to  $(i\gamma^\nu\partial_\nu + m)(i\gamma^\mu\partial_\mu - m)\psi = 0$

$$[-\gamma^\mu\gamma^\nu\partial_\mu\partial_\nu - m^2]\psi = 0$$

Oh,  $\partial_\mu\partial_\nu = \partial_\nu\partial_\mu$  so I can  $\gamma^\mu\gamma^\nu \rightarrow \frac{1}{2}\{\gamma^\mu, \gamma^\nu\} = g^{\mu\nu}$

$$(-\partial_\mu\partial^\mu - m^2)\psi = 0 \text{ Klein-Gordon equation}$$

Momentum-space:  $(p^2 - m^2)\psi = 0$  or  $E^2 = p^2 + m^2$  good.

So  $\psi$  must obey  $\psi(x) = \int d^4p e^{ip \cdot x} \psi(p)$

L36P4

with  $(p^2 - m^2) \psi(p) = 0$

But that's necessary, not sufficient constraint on  $\psi(p)$ .

Applying  $(i\gamma^\mu \partial_\mu + m)$  may have sneaked in spurious solutions.

~~But before trying to do better, let's ask about interacting WEM field,~~

~~eg,  $A^\mu \neq 0$~~

~~$$(i(\not{\partial} - ieA^\mu) \gamma_\mu + m)(i(\not{\partial} - ieA^\nu) \gamma_\nu - m) \psi = 0$$~~

~~$$[-\gamma^\mu \gamma^\nu (\underbrace{\not{\partial}_\mu - ieA^\mu}_\mu) (\underbrace{\not{\partial}_\nu - ieA^\nu}_\nu) - m^2] \psi = 0$$~~

~~do not commute.~~