

Signs and Fermions

Order of operators matters in fermionic system.

But sign of external state also not always uniquely defined.

eg, if in-state has two fermions ψ_1, ψ_2 is it

$$|\psi_1(p_1) \psi_2(p_2)\rangle = \int \bar{\psi}_1(x_1) \bar{\psi}_2(x_2) |0\rangle e^{i(p_1 \cdot x_1 + p_2 \cdot x_2)} d^4x_1 d^4x_2$$

or

$$= \int \bar{\psi}_2(x_2) \bar{\psi}_1(x_1) |0\rangle$$

I have to choose one by convention.

Overall sign in \mathcal{M} .

Then when I take $\mathcal{M}^* \mathcal{M}$ to get rate \rightarrow doesn't matter!!

But relative signs can matter. Here is where:

If ~~as~~ two processes connect fermions together differently

If there is an "internal loop" of fermions

Illustration: Yukawa theory $\mathcal{L} = \frac{1}{2} (\partial_\mu \psi)^2 + \bar{\psi} (i\not{\partial} - m) \psi$
 $- g \bar{\psi} \psi \phi - \frac{\lambda}{24} \phi^4 - \frac{m_\phi^2}{2} \phi^2$

Consider 3 processes

$\psi\psi \rightarrow \psi\psi$ scatt (for practice)

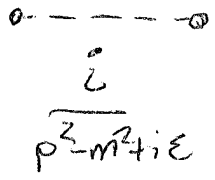
$\psi\psi \rightarrow \psi\psi$ scatt (exchange-signs)

$\psi\psi \rightarrow \psi\psi$ if $g^4 \sim \lambda$ (you'll see why!)

$\phi\phi \rightarrow \phi\phi$

Feynman rules

$$i\mathcal{L} = i\int d^4x \left[\frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{im^2}{2} \phi^2 + i\bar{\psi}(\not{\partial} - m)\psi - iy\phi\bar{\psi}\psi \right]$$



inverse of thing
between ϕ 's
 $\phi(-i\not{\partial} - m)\phi$
 $= \phi(i(\not{p}^2 - m^2))\phi$



$$i \frac{(\not{p} + m)}{p^2 - m^2 + i\epsilon} = \frac{i}{\not{p} - m}$$

inverse of thing
between $\bar{\psi}$ and ψ

$$- \frac{i\lambda}{24} \phi^4$$

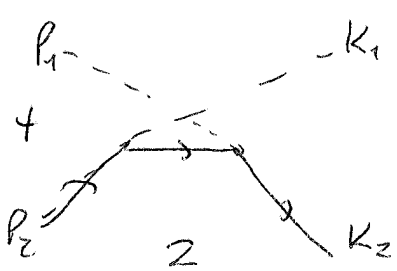
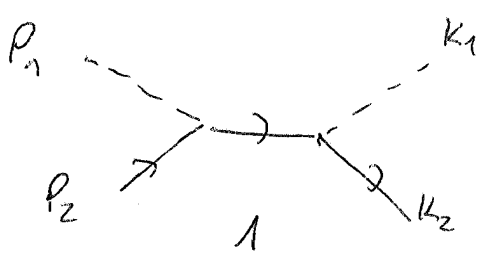
$$-iy \rightarrow i$$

$$- \frac{i\lambda}{24} \times$$

not symm factor

And overall i
from sign convention
of \mathcal{L} .

$$i^*(\epsilon) \mathcal{M} = -i(G's) \langle \phi\phi \rangle \dots$$

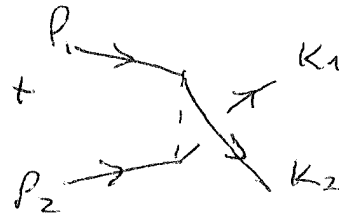
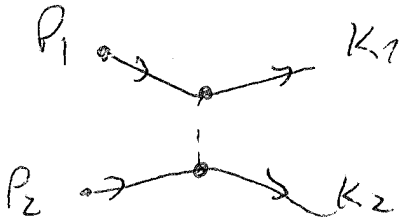


$$M_1 = \bar{u}(\sigma_2, k_2) (-iy) \frac{i(\not{p}_1 + \not{p}_2 + m)}{(p_1 + p_2)^2 - m^2} (-iy) u(\sigma_1, p_2) = -iy^2 \frac{\bar{u}(\sigma_2, k_2) (\not{p}_1 + \not{p}_2 + m) u(\sigma_1, p_2)}{s - m^2}$$

$$M_2 = \bar{u}(\sigma_2, k_2) (-iy) \frac{i(\not{p}_2 - \not{k}_1 + m)}{(p_2 - k_1)^2 - m^2} (-iy) u(\sigma_1, p_2) = -iy^2 \frac{\bar{u}(\sigma_2, k_2) (\not{p}_2 - \not{k}_1 + m) u(\sigma_1, p_2)}{u - m^2}$$

Relative sign is same - as they happen in same fermionic order
evaluation - trace methods (wing it)

$$\psi\psi \rightarrow \psi\psi$$



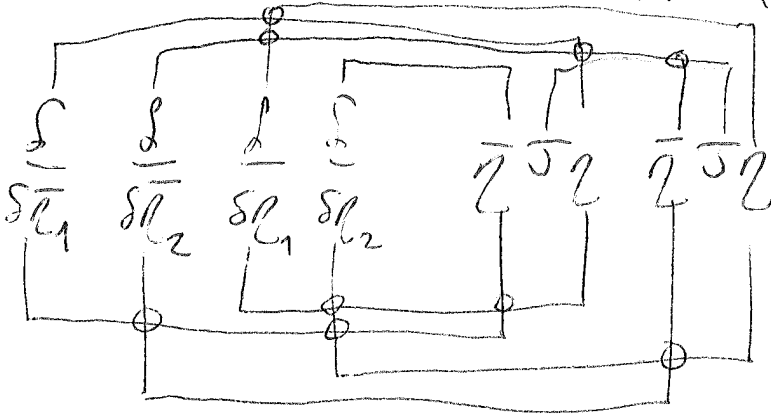
Beware!

Suppose I choose convention, in state = $b_{p_2}^\dagger b_{p_1}^\dagger |0\rangle$

but state = $\langle 0 | b_{k_1} b_{k_2}$

Then I want (up to i 's)

$$\langle 0 | T(\psi(k_1) \psi(k_2) \bar{\psi}(p_2) \bar{\psi}(p_1)) | 0 \rangle$$



4 order changes for diagram 2

5 order changes for diagram 1

Different signs! so $M_1 + M_2 = \frac{(i)^2 \bar{u}(k_1) u(p_1) \bar{u}(k_2) u(p_2)}{(p_1 - k_1)^2 - m^2 + i\epsilon}$

relative - sign is physical.

$$- \frac{(i)^2 \bar{u}(k_1) u(p_2) \bar{u}(k_2) u(p_1)}{(p_1 - k_2)^2 - m^2 + i\epsilon}$$

Arises when assignment of incoming fermion lines to outgoing involves an odd permutation

$$\begin{pmatrix} 1 \\ 2 \end{pmatrix} \rightarrow \begin{pmatrix} 1 \\ 2 \end{pmatrix} \text{ vs } \begin{pmatrix} 1 \\ 2 \end{pmatrix} \rightarrow \begin{pmatrix} 2 \\ 1 \end{pmatrix}$$

└── odd permutation

evaluation: $M_1^2 = \text{Tr } u_{k_1} \bar{u}_{k_1} u_{p_1} \bar{u}_{p_1} \text{Tr } u_{k_2} \bar{u}_{k_2} u_{p_2} \bar{u}_{p_2} = \text{Tr}(k_1 - m) \text{Tr}(p_1 - m) \text{Tr}(k_2 - m) \text{Tr}(p_2 - m)$

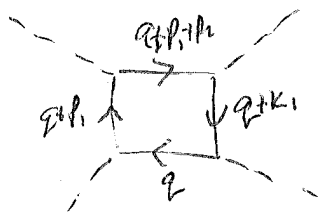
$$p_1^2 - 2k_1 \cdot p_1 + k_1^2 = (p_1 - k_1)^2 = t$$

$$2m^2 - 2k_1 \cdot p_1 = t \quad k_1 \cdot p_1 = \frac{-t + m^2}{2}$$

$$\rightarrow \frac{4(t - 2m^2)^2}{(t - m^2)^2}$$

while $\psi\psi \rightarrow \psi\psi$ at $\mathcal{O}(y^4)$

L20 P4



+ 5 permutations

$$\int_{S_1} \dots \int_{S_4} (iy)^4 \int_{S_1} \int_{S_2} \int_{S_3} \int_{S_4} \int_{S_1} \int_{S_2} \int_{S_3} \int_{S_4} \int_{S_1} \int_{S_2} \int_{S_3} \int_{S_4} \int_{S_1} \int_{S_2} \int_{S_3} \int_{S_4}$$

pair them off: $\int_{S_1} \int_{S_2} e^{+i\vec{q} \cdot S_2} = i5$

but one is $\int_{S_2} \int_{S_1} = - \int_{S_1} \int_{S_2}$

but $\int_{S_2} \int_{S_1} e^{i\vec{q} \cdot S_1} = -i5$ sign!

Always a - sign.

One - sign per fermion loop because of this.

$$i \int \frac{d^4 q}{(2\pi)^4} (iy)^4 i^4 (-) \text{Tr} \left(\frac{(q+m)(q+k_1+m)(q+k_1+k_2+m)(q+k_2+m)}{(q^2-m^2)(k_1^2-m^2)(k_1+k_2)^2-m^2)(k_2^2-m^2)} \right)$$

asthene a loop

Oh - and this looks like, at large q , it will

$$\rightarrow -iy \int \frac{d^4 q}{q^2 q^2 q^2 q^2} \sim \int \frac{d^4 q}{q^4} \text{ diverges logarithmically for large } q? \text{ Is that a problem?}$$

Takeaway: - signs when:

$(-)^{\# \text{ of fermion loops}}$

relative - between \mathcal{M} -contrib's with odd permutation of in \leftrightarrow out fermion lines.

That's it. Memorize & stop thinking 😊

γ^5 aside note $\gamma^0 \gamma^1 = \begin{bmatrix} -\sigma_1 & 0 \\ 0 & \sigma_1 \end{bmatrix} \rightarrow K_A$ [L20P5]

$\gamma^2 \gamma^3 = \begin{bmatrix} i\sigma_1 & 0 \\ 0 & -i\sigma_1 \end{bmatrix} \rightarrow J_A$

$\gamma^0 \gamma^1 \gamma^2 \gamma^3 = +i \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$

$i \gamma^0 \gamma^1 \gamma^2 \gamma^3 = -1 \begin{bmatrix} 1 & 0 \\ 0 & +1 \end{bmatrix}$

opposite sign for R (\mathbb{F}) as for L (+)

and that's invariant as $K_A J_A = \pm i 11$ depending on R, L for any Weyl.

More Lorentz way to write

$$\gamma^5 = \frac{i}{24} \epsilon_{\mu\nu\rho\sigma} \gamma^\mu \gamma^\nu \gamma^\rho \gamma^\sigma \quad (\epsilon_{0123} = 1)$$

$$\frac{11 - \gamma^5}{2} \text{ projects out R-comp.}$$

$$\frac{11 + \gamma^5}{2} \text{ projects out L-comp}$$

$\gamma^5 = \gamma_5 = \gamma_5$ in diff⁴
people's notation