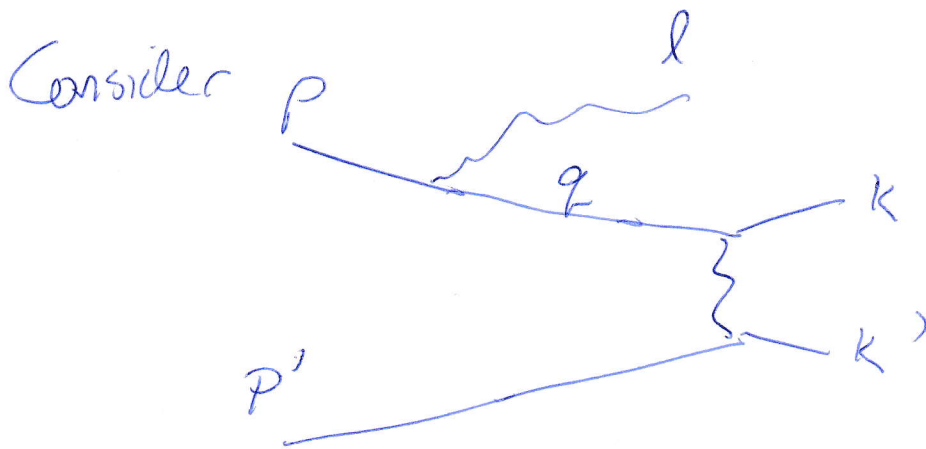


Finish Splitting

QCDG P1



Break q propagator into $\frac{u\bar{u}(q)}{q^2-m^2} = u \frac{1}{q^2-m^2} \bar{u}$

part of part of

Next time: $\int \frac{d^4 l}{(2\pi)^4} 2\pi \delta(l^2) = \int \frac{dl^0}{2\pi} \frac{dl_z}{2\pi} \int \frac{d^2 l_\perp}{(2\pi)^2} 2\pi \delta(l_0^2 - l_z^2 - l_\perp^2)$

$\frac{dl_z}{p} = \int dz =$ momentum fraction $SO \rightarrow \int_0^1 \frac{dz}{z} \frac{1}{4\pi} \frac{1}{4\pi} \int d(l_\perp^2)$

$\frac{1}{2l_0} = \frac{1}{2zP}$

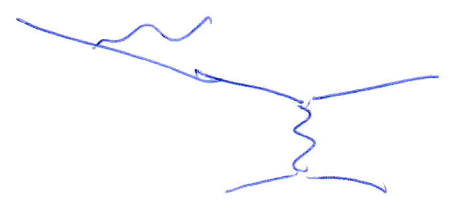
Next time: produced $|M|^2 = \frac{2e^2 l_\perp^2 (1+(1-z)^2)}{(1-z)z^2}$

$\left(\frac{1}{q^2-m^2}\right)^2 = \frac{z^2}{(l_\perp^2)^2}$ ~~AA~~

Finally, $\frac{1}{2p \cdot 2p'} |M|_{p' \rightarrow k'}^2 \rightarrow \frac{1}{2p \cdot 2p'} |M|_{2p' \rightarrow k k'}^2 = \frac{(1-z)}{2q^0 2p'^0} |M|_{2p' \rightarrow k k'}^2$

~~AA~~

Combining



emitting phase space

Matrix element propagator

matrix element

$$\int \frac{d^3 k_\perp d^2 k_\parallel}{(2\pi)^4} z \pi \delta(k_\perp^2)$$

↓

$$\int \frac{dz}{z} \frac{1}{16\pi^2} \int d^2(k_\perp^z) \frac{z^2(1+(1-z)^2) k_\perp^2}{z^2(1-z) (k_\perp^z)^2} (1-z) \times (\text{usual } 2 \rightarrow 2 \text{ calc.})$$

$$= \int \frac{dz}{z} \int \frac{d^2(k_\perp^z) k_\perp^2}{(k_\perp^z)^2} \left(\frac{e^2}{8\pi^2}\right) (1+(1-z)^2) \times \int_{(1-z)p, p' \rightarrow kk'}$$

Do first. Max value of k_\perp^2 is zS

min value is $z m_e^2$ - the $\frac{1}{k_\perp^2}$ ignored mass effects which emerge here.

→ $\ln\left(\frac{S}{m_e^2}\right) + \mathcal{O}(1)$. We have to think about meaning

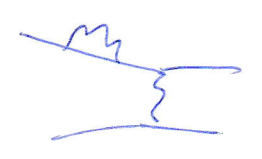
$$\int_{1-\delta \text{ emit}}^{\text{scatt}} = \left(\frac{\alpha_{EM}}{2\pi}\right) \left(\ln \frac{S}{m_e^2}\right) \int_0^1 dz \left(\frac{1+(1-z)^2}{z}\right) \int_{\text{Energy } (1-z)p}^{\text{scatt, no } \delta}$$

Meanwhile

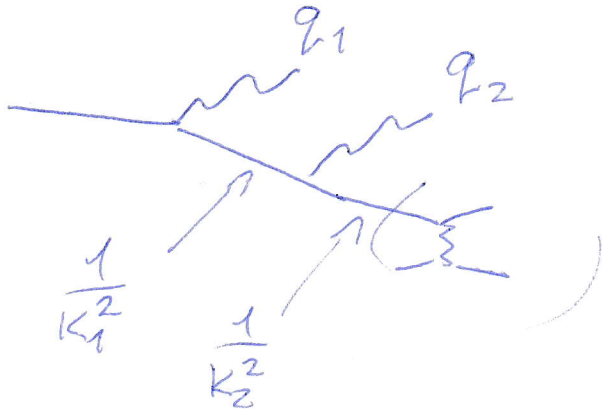
$\int_{\text{no } \delta - 1 \text{ loop correction}}^{\text{scatt}} = \dots$

||

$\int_{\text{Energy } p}^{\text{scatt, no } \delta}$



What about 2 γ emit?



Claim: $\frac{1}{k_1^2} \ll \frac{1}{k_2^2}$ eg if $q_{1\perp}^2 \ll q_{2\perp}^2$
 \int get iteration of last calculation

$$\int \frac{d^2 q_{1\perp}}{m^2 (2\pi)^2} \frac{1}{q_{1\perp}^2} \int \frac{d^2 q_{2\perp}}{q_{2\perp}^2 (2\pi)^2} \frac{1}{q_{2\perp}^2}$$

But if the $q_{2\perp}^2 \ll q_{1\perp}^2$ part of \int does not contribute much

why? if $q_{1\perp}^2 \ll q_{2\perp}^2$, e^-

mayas well be on-shell
 same calc applies.

opposite case - it's not almost on-shell - say, $k_1^0 = \sqrt{k_{1\perp}^2} + \delta k_1^0$

then $\int d^4 k_2 \int d^4 q^0 d^2 q_z d^2 q_{1\perp} \delta(k_1 - q_2 - k_2) \delta(q_2^2)$

has $\int d^4 q^0 \delta(q_0^2 - q_z^2 - q_{1\perp}^2)$ set $q^0 = q_z + \frac{q_{1\perp}^2}{2q_z}$ as before

then $k_2^0 = k_1^0 - q^0 = \delta k_1^0 + k_z - q_z + (\text{small})$

... k^0 not

So enhancement from propagator
 doesn't make up for q_{\perp}^2 phase space

$$\frac{\int d^2 q_{\perp 2} q_{\perp 2}^2}{(q_{\perp 1}^2)^2} \quad \text{for } q_{\perp 1}^2 \gg q_{\perp 2}^2 \quad \text{no large contrib.}$$

With this in mind, we can solve probs to have energies from
 "outside - in"

$$\sigma_{\text{scatt}} = \int_0^1 dz \left(\frac{d \text{Prob to have } e^- \text{ or } \gamma \text{ w.}}{dz} \text{ fraction } z \right) \sigma_{\text{ep}} \dots$$

for 1 emission we found this was $\frac{\alpha}{2\pi} \left(\frac{1+(1-z)^2}{z} \right) \ln \frac{s}{m_e^2}$

for photon, $z \leftrightarrow (1-z)$ for electron ...
 0 emission: $\int (z-1) [1 - \text{chance to have 1 emission}]$


§. Name these probs $e_e(z), \gamma_e(z)$ Also depend on S .
 chance to have e or γ
 if you start with e
 and look at mom frac z

Call intermediate quantity: $e_e(z)$ counting emissions with q_{\perp}^2 up
 to some value μ^2 ,

$e_e(z, \mu^2)$ book uses f_e

Then way $\int d^2 q_{\perp}^2$ int. worked, and nested nature of ints,

tells us that

As you increase μ^2 ,  process should turn e 's into $e\gamma$ pairs.

Evolving w μ^2 (like RG) captures iterated emissions?

Pick μ_2^2 slightly larger than μ_1^2

$$e_e(z, \mu_2^2) = e_e(z, \mu_1^2) + \ln \frac{\mu_2^2}{\mu_1^2} * \left(\text{coeff. of } \int \frac{d^2 q_{\perp}}{q_{\perp}^2} \text{ of all} \right)$$

(ways $z^- e^-$ could turn into $z^- e^-$)

$-\ln \left(\frac{\mu_2^2}{\mu_1^2} \right) * \text{coeff. of } \int \frac{d^2 q_{\perp}}{q_{\perp}^2}$ & all ways $e^-(z)$ could turn into γ, e of smaller man fac.

$$\mu^2 \frac{d}{d\mu^2} e_e(x, \mu^2) =$$

$$= \int_0^1 \frac{dz}{z} \left(\frac{1+z^2}{1-z} \right) \frac{\alpha}{2\pi} e_e\left(\frac{x}{z}, \mu^2\right)$$

ways someone else makes me

$$- \int_0^1 dz \left(\frac{1+(1-z)^2}{z} \right) \frac{\alpha}{2\pi} e_e(x, \mu^2)$$

ways of $\rightarrow e\gamma$

$$\mu^2 \frac{d}{d\mu^2} \gamma_e^e(x, \mu) = \int_x^1 \frac{dz}{z} \left(\frac{1+(1-z)^2}{z} \right) \frac{\alpha}{2\pi} e_e\left(\frac{x}{z}, \mu^2\right)$$

Subtlety 1: canceling divergences

$$\int_x^1 \frac{dz}{z} \left(\frac{1+z^2}{1-z} \right) \frac{d}{2\pi} \rho_e \left(\frac{x}{z}, \mu^2 \right)$$

diverges as $z \rightarrow 1$. where $\rho_e \left(\frac{x}{z} \rightarrow x \right)$ same as in other term!


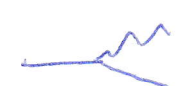
Define $\frac{1}{(1-x)_+}$ as " $\frac{1}{1-x}$ but with negative δ -func. at $x=1$ "

such that $\int_0^1 dx \frac{f(x)}{(1-x)_+} = \int_0^1 dx \frac{f(x) - f(1)}{1-x}$ is well defined.

(or $\lim_{\epsilon \rightarrow 0} \frac{1}{1-x} \Theta(1-x-\epsilon) - \int_0^{1-\epsilon} dx / (1-x)$)

note, $\int_0^1 dx \frac{1+x^2}{(1-x)_+} = -3/2$ (show!)

Now $\mu^2 \frac{d}{d\mu^2} \rho_e(x, \mu^2) = \int_x^1 \frac{dz}{z} \left[\frac{1+z^2}{(1-z)_+} + \frac{3}{2} \delta(1-z) \right] \frac{d}{2\pi} \rho_e \left(\frac{x}{z}, \mu^2 \right)$

Next: realize that  also possible. μ^2 by crossing from 

$\mu^2 \frac{d}{d\mu^2} \rho_e(x, \mu^2) = \frac{d}{2\pi} \int_x^1 \frac{dz}{z} \left[P_{\leftarrow e}(z) \rho_e \left(\frac{x}{z}, \mu^2 \right) + P_{\leftarrow \gamma}(z) \gamma_e \left(\frac{x}{z}, \mu^2 \right) \right]$

similar for γ

$P_{\leftarrow e} = \frac{1+z^2}{(1-z)_+} + \frac{3}{2} \delta(1-z)$ $P_{\leftarrow \gamma} = z^2 \sqrt{1-z}^2$