

Consider massless theory, eg,

$$\mathcal{L}(\psi, \phi) = \bar{z}_1 \bar{\psi} (\not{\partial} + \cancel{m} + sm) \psi + \frac{z_2}{2} \partial_\mu \phi \partial^\mu \phi - \frac{(\lambda + s\lambda)}{4!} z_2^2 \phi^4 - (y + sy) z_1 z_2^{3/2} \phi \bar{\psi} \psi - sM^2 \phi^2$$

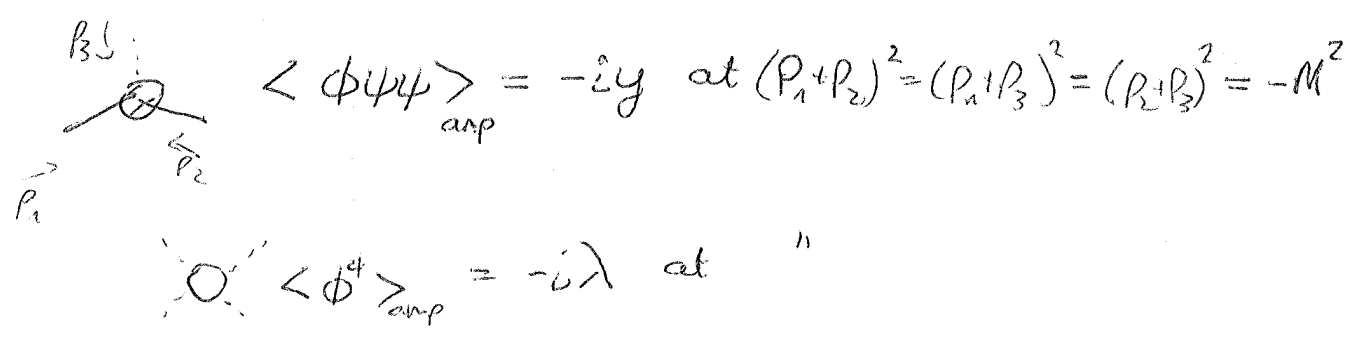
with  $\psi, \lambda, y, \phi$  renormalized quantities,  $z_1, z_2, s\lambda, sy, sM^2, sm$  renorm. const.

Choose  $z, s\lambda, etc$  based on renorm. prescription

①  $\Sigma = \frac{d}{dp} \Sigma = 0$  at  $p^2 = -M^2$

②  $\Pi = \frac{d}{dp^2} \Pi = 0$  at  $p^2 = -M^2$

note, all at spacelike external momenta, described by introduced scale  $M$ .



Unique mapping between bare  $\phi_0, \psi_0, \lambda_0, y_0 \rightarrow$  ren  $z_1, z_2, \lambda, y$   
 But mapping is  $M$  dependent!

How must  $z, \lambda, y$  change if I change  $M$ ?

Holding what fixed? Physics  $\rightarrow$  bare quantities & their correl. funcs.

what happens when I change scale  $M$ ? Renorm.  $\lambda, y, z_1, z_2$  change.

specifically consider bare correl. func.

L8P2

$$G_0^{(n)} \equiv \langle 0 | T [\phi_0(x_1) \dots \phi_0(x_n)] | 0 \rangle_{\text{conn}} \quad \text{or some } \psi\text{'s in there,}$$

eg,  $G_{\psi\psi\psi}$  or  $G_{\bar{\psi}\psi\psi}$  ...

$$G^{(n)} = \langle 0 | T (\phi(x_1) \dots \phi(x_n)) | 0 \rangle_{\text{conn}} = Z^{-n/2} G_0^{(n)}$$

or  $G^{(n_1)(n_2)} = Z_1^{-n_1/2} Z_2^{-n_2/2} G_0^{(n)}$

~~Now~~ Now  $G_0^{(n)}$  is  $M$ -independent but it looks  $M$ -dependent from logs and chosen values of counterterms, eg,

$$X + \mathcal{K} = \lambda - \frac{\lambda^2}{32\pi^2} \left( \frac{3}{\epsilon} + \ln \frac{\mu^6}{stu} + c \right) + \text{c.t.} = \lambda \text{ for } s=t=u=M^2$$

so c.t. =  $\frac{\lambda^2}{32\pi^2} \left( \frac{3}{\epsilon} + \ln \frac{\mu^6}{M^6} + c \right)$

$$X + \mathcal{K} = \lambda - \frac{\lambda^2}{32\pi^2} \left( \ln \frac{M^6}{stu} \right) \text{ renormalized expression with our renorm. prescription.}$$

Explicitly  $M$ -dep. from counterterm choice

Implicitly  $M$ -dep. from  $\lambda = \lambda(M)$  depends on renorm. scale

But  $G_0^{(n)}$  has no actual  $M$ -dependence.

$$M \frac{d}{dM} G_0^{(n)} = 0 = \left( M \frac{\partial}{\partial M} + \frac{M \partial \lambda}{\partial M} \frac{\partial}{\partial \lambda} + \frac{M \partial y}{\partial M} \frac{\partial}{\partial y} + \frac{M \partial Z_1}{\partial M} \frac{\partial}{\partial Z_1} + \frac{M \partial Z_2}{\partial M} \frac{\partial}{\partial Z_2} \right) G_0^{(n)}$$

$$Z_1^{-n_1/2} Z_2^{-n_2/2} G^{(n_1, n_2)}$$

Define  $\beta_\lambda = \frac{M \partial \lambda}{\partial M}$

$$0 = \left( M \frac{\partial}{\partial M} + \beta_\lambda \frac{\partial}{\partial \lambda} + \beta_y \frac{\partial}{\partial y} + n_1 \delta_\psi + n_2 \delta_{\bar{\psi}} \right) G^{(n_1, n_2)}$$

Callan-Symanzik Eq

How to compute? For me, Wick rotate first (Euclidean methods)

$$\langle \phi(x) \phi(0) \rangle = \int \frac{d^4 p}{(2\pi)^4} e^{ip \cdot x} \frac{1}{p^2} \quad \text{no } i\text{'s!!} \quad p^2 = \vec{p}^2 + p_0^2$$

positive metric  $g = \int_{\mu}^{\nu}$

Our renorm. convention defined here anyways.

$$\int \mathcal{D}\phi \bar{\phi} \psi \exp - \int d^4 x \left[ \bar{\psi} (\not{\partial}) \psi + \frac{1}{2} \partial_\mu \phi \partial^\mu \phi + \frac{\lambda}{4!} \phi^4 + \gamma \phi \bar{\psi} \psi \right]$$

every term positive - positive action for convergence  
no  $i$  if  $\{x^\mu, y^\nu\} = g^{\mu\nu} = \delta^{\mu\nu}$ .

~~$\times$~~   $\begin{matrix} | \\ \diagdown \\ \diagup \\ \times \end{matrix}$  due to this - sign.  
 $-\lambda$   $-\gamma$

Consider first  $Z_1$ :

$$\frac{1}{i\phi} + \frac{p_0}{i\phi} \frac{1}{(2\pi)^D} \int d^D q \frac{q+p}{(q^2 + (q+p)^2)} \frac{1}{i\phi} = \frac{1}{i(\phi - \Sigma(p))}$$

from  $i^2$  sorry.

$$\Sigma = \frac{-\gamma^2}{16\pi^2} \int_0^1 dx \left( \frac{\mu^2}{4\pi} \right)^{4-D} \int_0^1 dq^2 \frac{(q+p)^2}{(q^2 + x(1-x)p^2)^2}$$

$$= \phi \left( \frac{-\gamma^2}{32\pi^2} \right) \left( \frac{1}{\epsilon} + 2 \int_0^1 dx \times \ln \left( \frac{\mu^2}{x(1-x)p^2} \right) \right) = 2 + \ln \frac{\mu^2}{p^2}$$

Prescription:  $p^2 = m^2$  above = 0.  
Same as  $\mu^2 = M^2 e^{-2}$

$$\Sigma = \frac{-\gamma^2}{32\pi^2} \ln \frac{m^2}{p^2}$$

$$\langle \psi_r \bar{\psi}_r \rangle = \frac{1}{\phi + \frac{\gamma^2}{32\pi^2} \ln \frac{m^2}{p^2} \phi} = \frac{-1}{Z_1(m)} \langle \psi_0 \bar{\psi}_0 \rangle$$

(since  $Z_1 \bar{\psi} \psi = \bar{\psi}_0 \psi_0$ )

$$\langle \Phi_0 \Phi_0 \rangle = \frac{Z(M)}{Z(0)}$$

$$\rho \left( 1 + \frac{y^2}{32\pi^2} \ln \frac{M^2}{\rho^2} \right)$$

$$\frac{\partial Z}{\partial M} = 0 = \frac{\partial}{\partial M} \left( \frac{M^2 Z}{Z \partial M} + \frac{y^2}{16\pi^2} \right)$$

" $\equiv$ "  $2\delta_\phi$

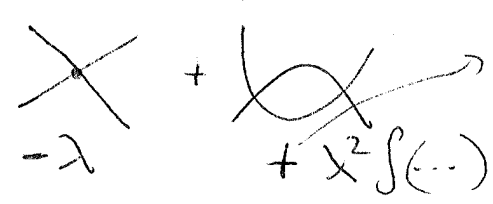
$$\delta_\phi = + \frac{y^2}{32\pi^2}$$

opposite sign on  $\phi$  in  $\Sigma$

But wait: there's also  $\frac{\partial Z}{\partial M} y$  ... but that's higher order.

Similarly  $\bigcirc$  gives  $\delta_\phi = y^2$  (Find it!) Return to...

If I take  $Z(\Lambda_{\text{cutoff}}) = 1$ ,  $\frac{1}{Z} \frac{\partial Z}{\partial M} = \delta_\phi < 0$  then  $Z(\text{low scale}) > 1$

Finally  ~~$\langle \phi^4 \rangle$~~   $\langle \phi_{\text{ren}}^4 \rangle_{\text{conn}} =$  

$$-G^{(4)} = \left( \frac{1}{\rho^2 (1 + \# y^2 \ln \frac{M^2}{\rho^2})} \right) \text{for each leg} \times \left( 1 - \frac{3\lambda^2}{32\pi^2} \left( \frac{1}{\epsilon} + \ln \frac{M^2}{(\text{str})^{4/3}} \right) \right)$$

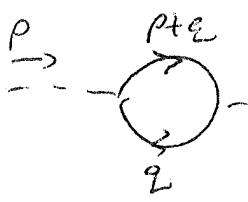
$$\left[ \frac{\partial}{\partial M} + \beta_\lambda \frac{\partial}{\partial \lambda} + 4\delta_\phi \right] = 0$$

on  $M^2$  in propagators

$$\left( \frac{1}{\pi \rho^2} \right) \left( \frac{-3\lambda^2}{16\pi^2} \right) + \frac{1}{\pi \rho^2} \beta_\lambda + 4\delta_\phi \frac{\lambda}{\pi \rho^2} - 8\delta_\phi \frac{\lambda}{\pi \rho^2} = 0$$

$$\beta_\lambda = + \frac{3\lambda^2}{16\pi^2} + 4\lambda\delta_\phi$$

→ gives rise to  $\lambda y^2$  term



has lessons

why?

$$\Pi(p^2) = -y^2 \int \frac{d^D q}{(2\pi)^D} \frac{\text{Tr } \not{q}(\not{p}+\not{q})}{i \not{q}^2 i (\not{p}+\not{q})^2}$$

fermion loop

another sign. But  $\not{q}\not{q} = +q^2$ .  
OR no i's but  $\not{q}\not{q} = -q^2$ .

Tr  $\rightarrow 4(q^2 + p \cdot q)$

shift:  $\int dx \frac{q^2 + p \cdot q}{(xq^2 + (1-x)(q^2 + 2p \cdot q + p^2))^2}$

shift

$q \rightarrow q - (1-x)p$   
 $q^2 + p \cdot q \rightarrow q^2 + (2x-1)p \cdot q + (1-x)^2 p^2$

$-x(1-x)p^2$   
 $-(1-x)p^2$

$-\frac{4y^2}{(4\pi)^{D/2}} \int dx \int (q^2)^{\frac{D-2}{2}} dq^2 \frac{q^2 + (2x-1)p \cdot q + (1-x)^2 p^2}{(q^2 + x(1-x)p^2)^2}$

$\Gamma(\epsilon-1) = \frac{1}{\epsilon-1} \Gamma(\epsilon)$   
 $= (-1-\epsilon) \Gamma(\epsilon)$   
 $= (\frac{1}{\epsilon} - \epsilon)$

Careful:  $\int \frac{q^2}{(q^2 + A)^2} (q^2)^{\frac{D}{2}-1} dq^2 = (A)^{1-\epsilon} \frac{\Gamma(\epsilon-1)\Gamma(3-\epsilon)}{\Gamma(2)} \approx 2$

$q^2$  term  $-\cancel{x(1-x)p^2}^{1-\epsilon} (-\frac{1}{\epsilon} + \gamma_E - 1)(2 + \mathcal{O}(\epsilon))$

$q$ -integral  $\rightarrow (\frac{1}{\epsilon} + \text{const.}) (-x(1-x)^2 p^2 - 2x(1-x)p^2) (1 + \epsilon \ln \frac{\mu^2}{p^2})$

$\int_0^1 dx (1-x)^2 = 1/3$   
 $x(1-x) = 1/6$   
 $-\frac{1}{2} p^2$

$= -\frac{y^2 p^2}{16\pi^2} \left( \frac{1}{\epsilon} + \ln \frac{\mu^2}{p^2} + \text{const.} \right)$

Note:  $q^2$  term also turned into  $p^2$  effect. No  $M^2$  generated.

same opposite signs as  $\Sigma$ .  $\gamma_\phi = +\frac{3y^2}{16\pi^2}$

$\beta_\lambda = \frac{3\lambda^2}{16\pi^2} + \frac{8y^2\lambda}{16\pi^2}$

Also: badway

$$q^0(p+q) = \left( \frac{1}{2}q^2 + p \cdot q + \frac{p^2}{2} \right) + \frac{1}{2}q^2 - \frac{1}{2}p^2$$

$$= \frac{(q+p)^2 + q^2 - p^2}{2}$$

$$\frac{q^0(p+q)}{q^2(p+q)^2} = \frac{(p+q)^2}{2q^2(p+q)^2} + \frac{q^2}{2q^2(p+q)^2} - \frac{p^2}{2q^2(p+q)^2}$$

Now  $\int \frac{d^D q}{(2\pi)^D} \frac{1}{(p+q)^2} = \int \frac{d^D q}{(2\pi)^D} \frac{1}{q^2}$  shift. IN MS, this is zero!  
Because, dimensions but no dimensionful quantity!

So, left with  $-\frac{p^2}{2} \int \frac{d^D q}{(2\pi)^D} \frac{1}{q^2(p+q)^2}$

$$= -\frac{p^2}{2} \frac{1}{(16\pi^2)} \left( \frac{1}{\epsilon} + \int_0^1 \ln \frac{\mu^2}{x(1-x)p^2} dx = 2 + \ln \frac{\mu^2}{p^2} \right)$$

$$= -\frac{p^2}{2} \left( \frac{1}{16\pi^2} \right) \left( \ln \frac{M^2}{p^2} \right) \text{ after counterterm removal.}$$

$$\overline{\pi} = \frac{-4y^2}{16\pi^2} p^2 \ln \frac{M^2}{p^2}$$

Also is forgot!



let's just find  $1/\epsilon$  - enough to get anom. dim etc.

$$-\lambda + (-)(y)^4 \int \frac{d^D q}{(2\pi)^D} \text{Tr} \left( \left( \frac{1}{i\cancel{q}} \right)^4 \right) \approx -\lambda - 4y^4 \int \frac{d^D q}{(2\pi)^D} (q^2 \dots)^2 \rightarrow \frac{1}{16\pi^2} \left( \frac{1}{\epsilon} + \ln \frac{M^2}{p^2} \right)$$

opposite sign as  $\lambda^2$  term we found!!

$$\beta_\lambda = \frac{+3\lambda^2}{16\pi^2} - \frac{4y^4}{16\pi^2} + \frac{8\lambda y^2}{16\pi^2} \quad \text{hmmm.}$$