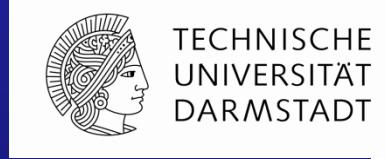


The Similarity Renormalization Group



- Renormalization Group methods in nuclear physics

by Timon Dörnfeld

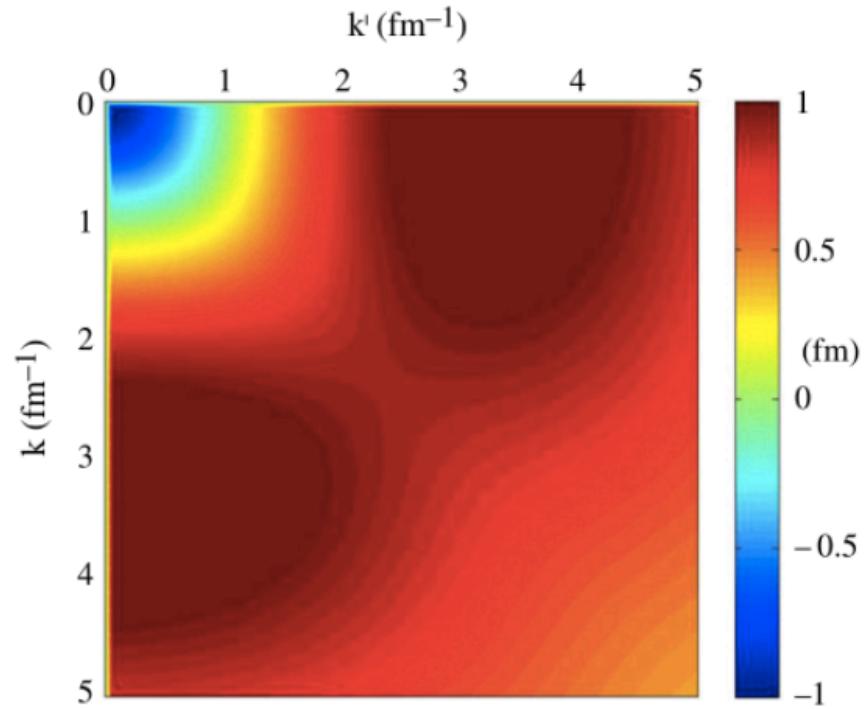
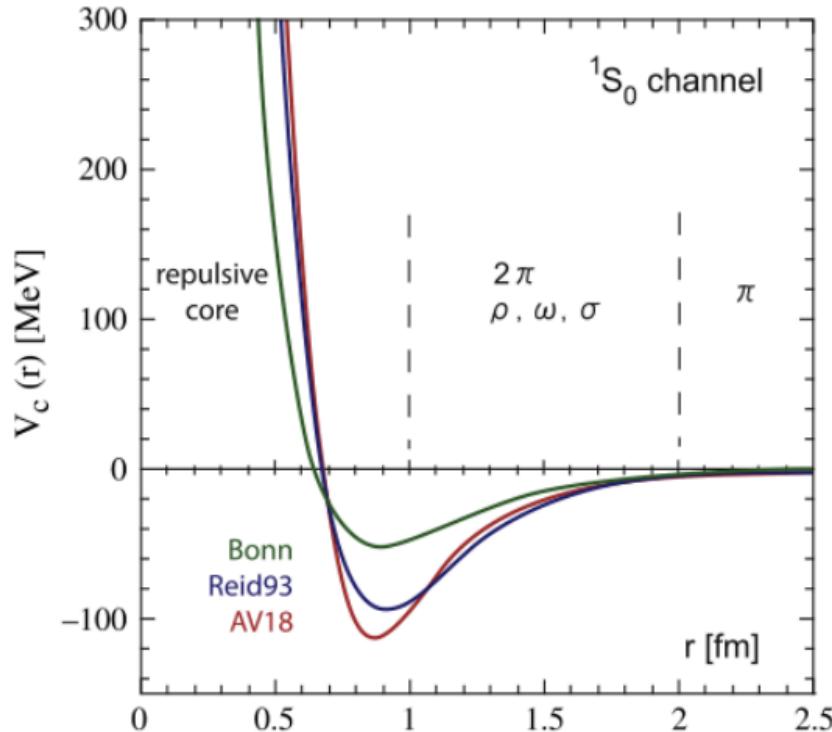
Topics



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- Introduction
 - Flow equation
 - Decoupling
 - NN-interaction-evolution & Universality
 - 3N-interaction-evolution
 - Many-body-interaction-evolution & In-Medium SRG
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Introduction

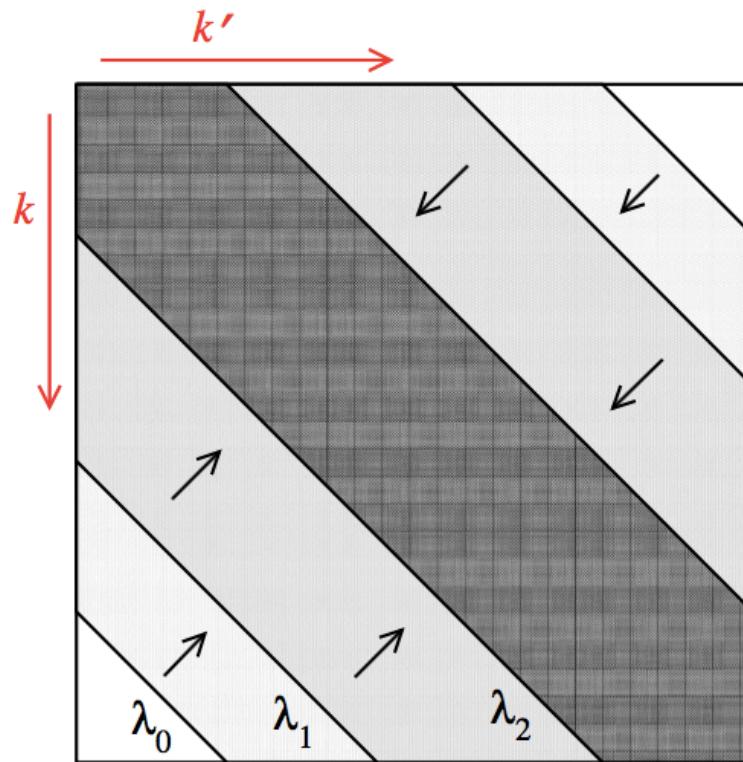


Left panel: three phenomenological potentials as function of interparticle distance that accurately describe proton-neutron scattering up to laboratory energies of 300 MeV. Right panel: alternative momentum space representation of the Av18 potential in the 1S_0 channel [2].



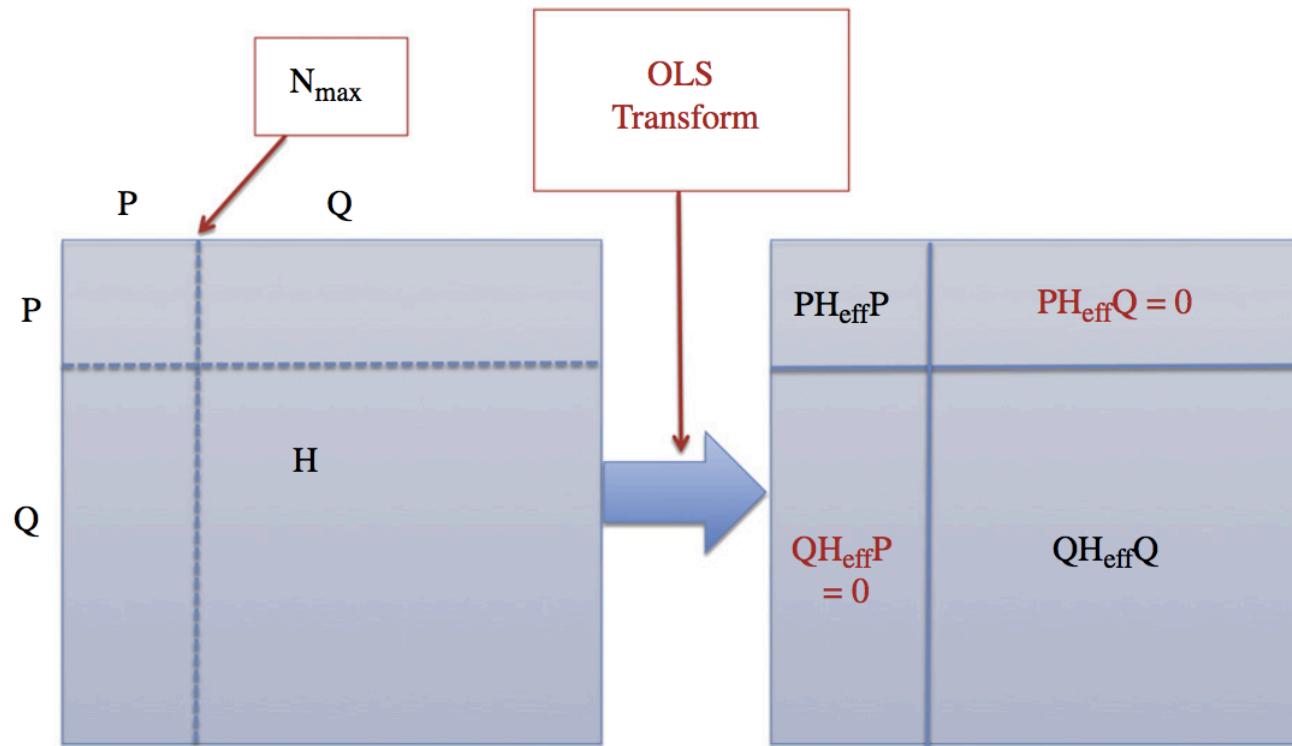
- coupling of low & high momenta (repulsive core)
 - slowly converging system (harmonic oscillator basis)
 - expansive numerical calculations of energies
- decoupling of low & high momenta
- faster converging hamiltonian

- *Similarity Renormalization Group (SRG) methode:*



Introduction

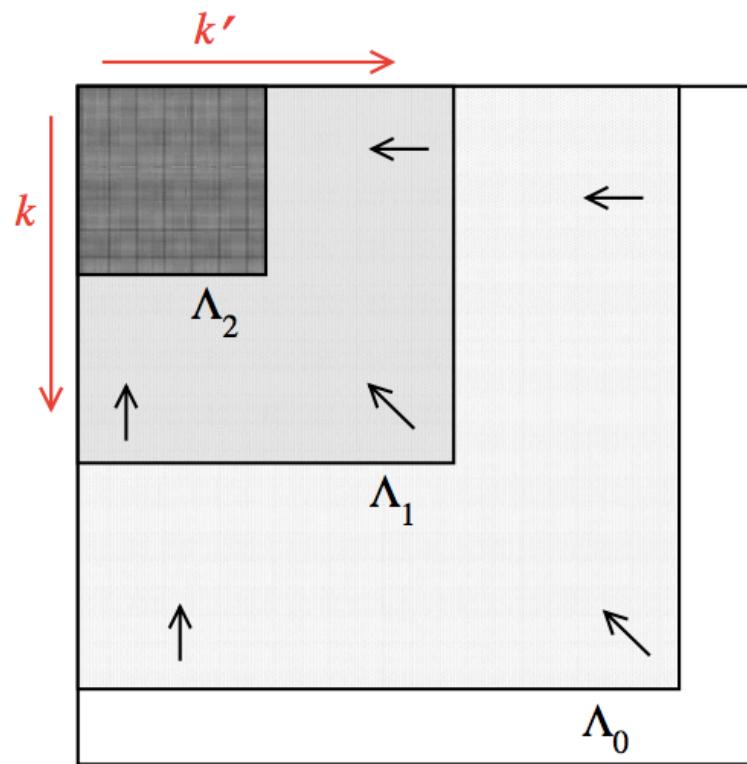
- *Okubo-Lee-Suzuki similarity transformation (OLS) methode:*



Introduction



- $V_{\text{low},k}$ methode:





Flow equation

$$H_s = U_s H U_s^\dagger = T + V_s$$

$$\text{with } T = \frac{kk'}{2\mu} \cdot \delta^3(k' - k) = \kappa^2$$

$$\frac{dH_s}{ds} = [\eta_s, H_s] = [[G_s, H_s], H_s]$$

$$\text{with the anti-hermitian generator } \eta_s = \frac{dU_s}{ds} U_s^\dagger$$

Decoupling



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$$\frac{dH_s}{ds} = G_s H_s H_s - 2H_s G_s H_s + H_s H_s G_s$$

$$\frac{dV_s}{ds} = 2(TV_s T - V_s TV_s) - T^2 V_s - V_s T^2 + TV_s^2 + V_s^2 T$$

inserting an $1 = \frac{2}{\pi} \int dq q^2$ to all V_s^2 -terms

$$\begin{aligned} \frac{dV_s}{ds}(\kappa, \kappa') &= - (\kappa^2 - \kappa'^2)^2 V_s(\kappa, \kappa') \\ &\quad + \frac{2}{\pi} \int dq q^2 (\kappa^2 + \kappa'^2 - 2q^2) V_s(\kappa, q) V_s(q, \kappa') \end{aligned}$$

Decoupling



$$\frac{dV_s}{ds}(\kappa, \kappa') = - (\kappa^2 - \kappa'^2)^2 V_s(\kappa, \kappa') + \sigma(V_s^2(\kappa, \kappa'))$$

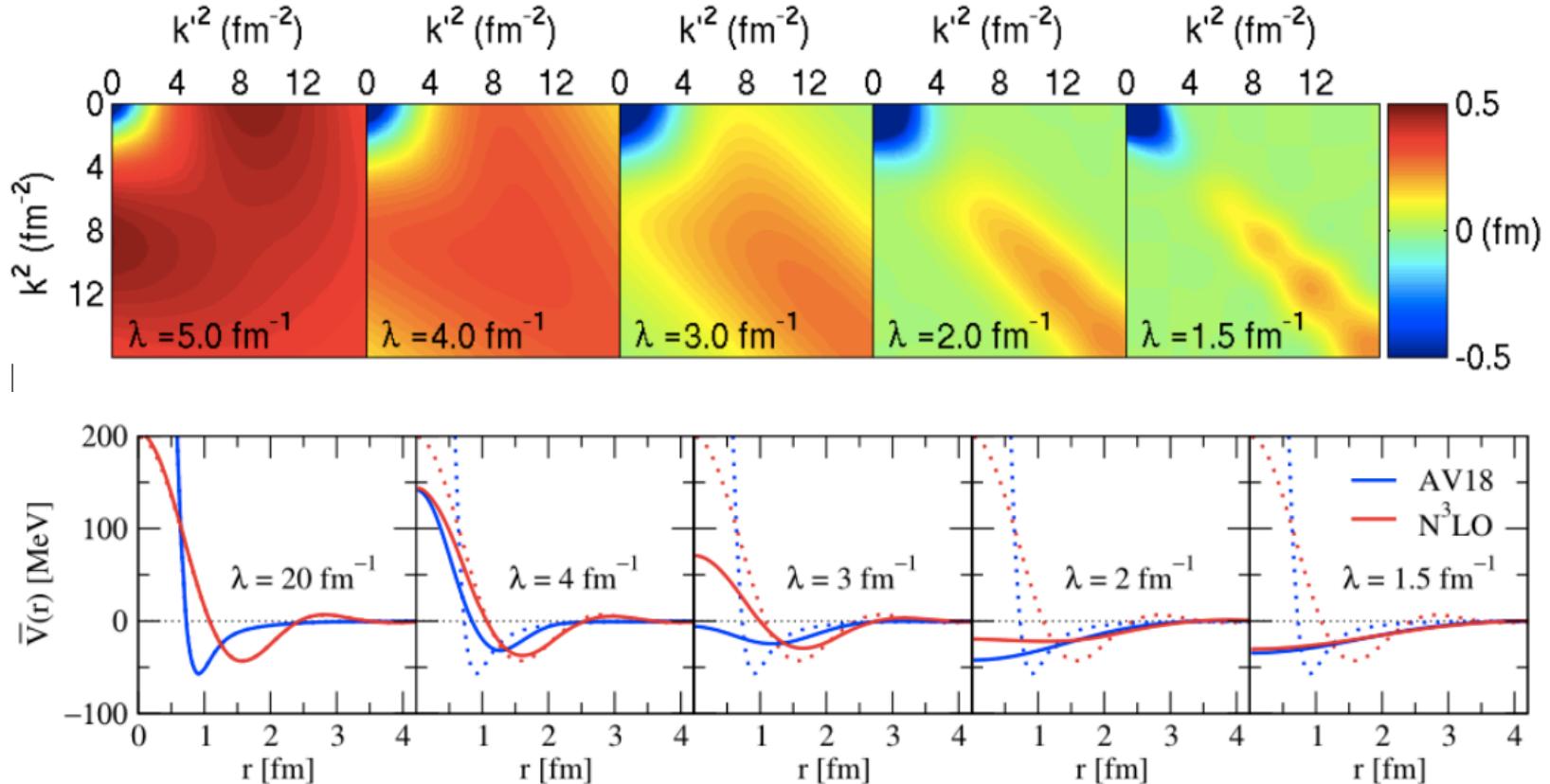
$$\Rightarrow V_s(\kappa, \kappa') \approx V_{s=0}(\kappa, \kappa') \cdot \exp\left(-s(\kappa^2 - \kappa'^2)^2\right)$$

$$(s : 0 \rightarrow \infty)$$

$$\Rightarrow V_\lambda(\kappa, \kappa') \approx V_{\lambda=\infty}(\kappa, \kappa') \cdot \exp\left(-\frac{1}{\lambda^4}(\kappa^2 - \kappa'^2)^2\right)$$

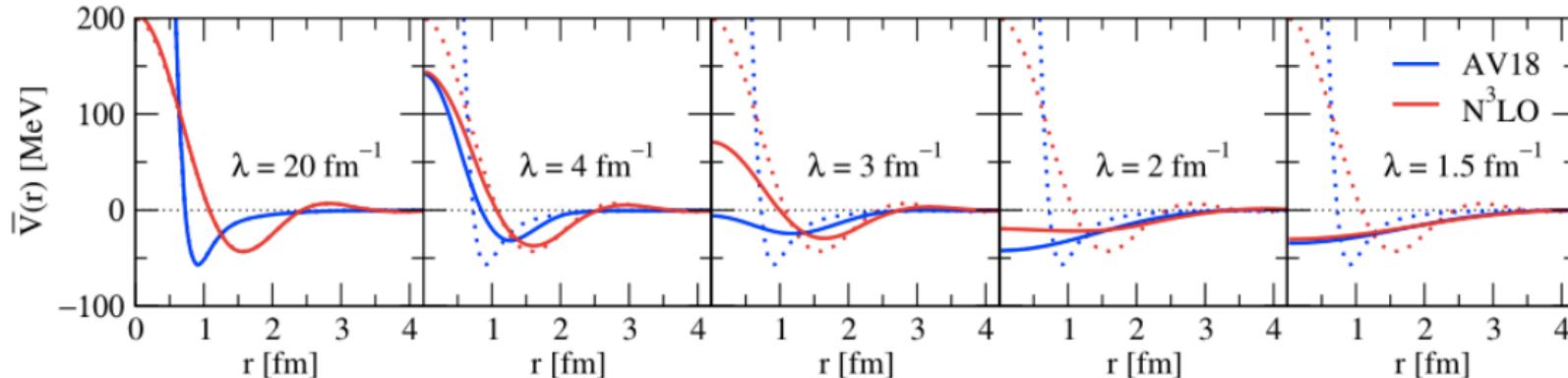
$$\left(\lambda = s^{-\frac{1}{4}} : \infty \rightarrow 0\right)$$

Decoupling & SRG-evolution



Top: The Av18 potential in the 3S_1 channel is evolved with SRG. Bottom: The local projections of Av18 and N^3LO (500 MeV) potential in 3S_1 channel at different resolutions [2].

NN-interaction-evolution & Universality



The local projections of Av18 and N^3LO (500 MeV) potential in 3S_1 channel at different resolutions [2].

- dissolution of the repulsive core to more attractive form
- shift of the long-range tail to a more attractive form
- NN-potentials converge to less repulsive forms
- evolved potentials vary hardly

- Types of 3N-interaction SRG-treatments:

(i) full evolution of NN- & 3N-forces

(ii) evolution of NN-interaction & additional N2LO 3N-forces

(iii) evolution of 3N-forces in a continuous plane wave basis

3N-interaction-evolution

$$H_\lambda = T + \sum_{i,j=1}^3 V_{ij} + V_{123} = T + V_\lambda$$

$$\Rightarrow d_\lambda V_\lambda = d_\lambda V_{12} + d_\lambda V_{23} + d_\lambda V_{13} + d_\lambda V_{123}$$

$$= [[T, V_\lambda], H_\lambda]$$

$$\Rightarrow d_\lambda V_{123} = [[T_{12}, V_{12}], V_{13} + V_{23} + V_{123}]$$

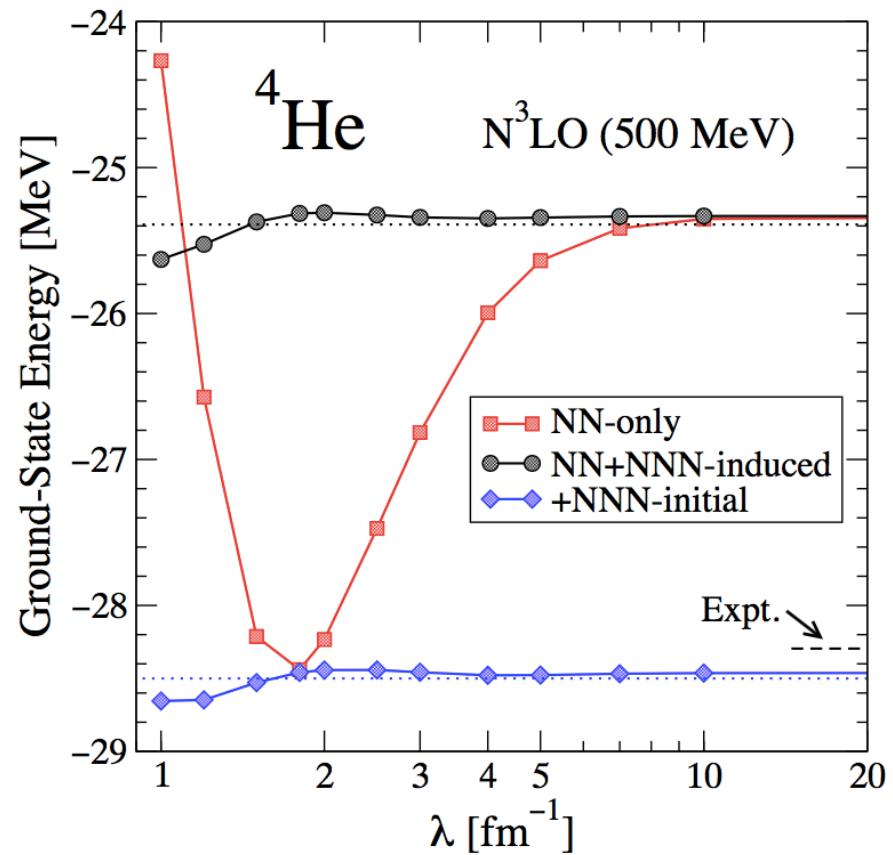
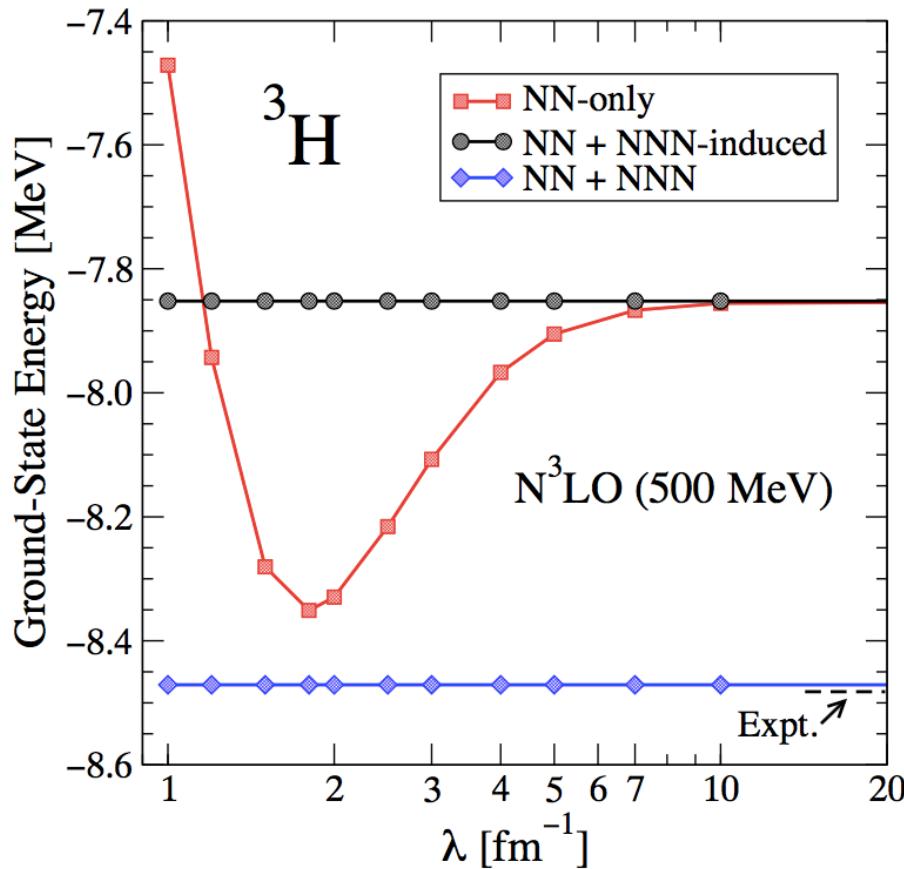
$$+ [[T_{13}, V_{13}], V_{12} + V_{23} + V_{123}]$$

$$+ [[T_{23}, V_{23}], V_{12} + V_{13} + V_{123}]$$

$$+ [[T, V_{123}], H_\lambda]$$

$$\Rightarrow d_\lambda V_{123} |_{V_{123}=0} \neq 0$$

3N-interaction-evolution



Ground-state energy of 3H (left) and 4He (right) as a function of the SRG evolution parameter λ for three different SRG flows [4].

Many-body-interaction-evolution & In medium-SRG

$$H_\lambda = \langle T \rangle a^\dagger a + \langle V_\lambda^{(2)} \rangle a^\dagger a^\dagger aa + \langle V_\lambda^{(3)} \rangle a^\dagger a^\dagger a^\dagger aaa + \sigma(V_\lambda^{(4)})$$

- SRG-evolution will induce A-body-forces, if applied in an A-body subsystem
- reference to a state of finite density instead of the vacuum
- 3- to A-body-forces by using 2-body machineries

References



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- [4] E. D. Jurgenson, P. Navrátil, R. J. Furnstahl - *Evolution of Nuclear Many-Body Forces with Similarity Renormalization Group*, Physical Review Letter (2009), LLNL-JRNL-412933
- [5] B. R. Barrett, P. Navrátil, J. P. Vary - Ab initio no core shell model, Progress in Particle and Nuclear Physics 69 (2013) 131-181