

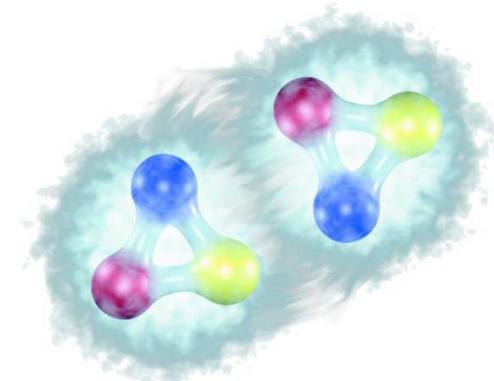
New Frontiers in Nuclear Structure Theory

From Realistic Interactions to the Nuclear Chart



Robert Roth

Institut für Kernphysik
Technical University Darmstadt



Overview

- Motivation
- Nucleon-Nucleon Interactions
- Solving the Many-Body Problem
- Correlations & Unitary Correlation Operator Method
- Applications

Nuclear Structure in the 21st Century

**new frontiers in
nuclear structure physics**

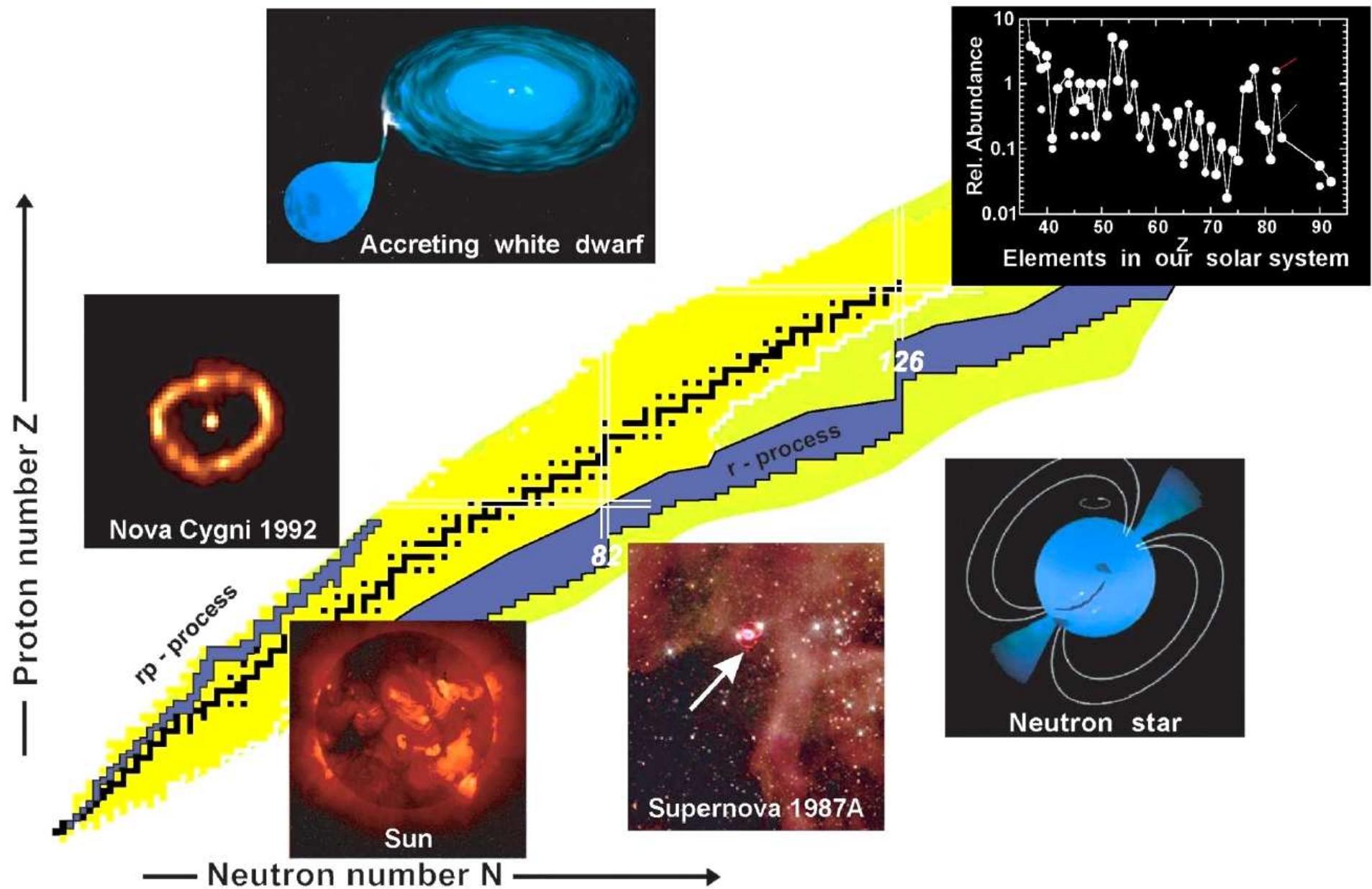
Experiment

- fundamental astrophysical questions need nuclear input
- possibilities to investigate nuclei far off stability
- new nuclear structure facilities: RIA, FAIR@GSI,...

Theory

- improved understanding of fundamental degrees of freedom / QCD
- high-precision realistic nucleon-nucleon potentials
- *ab initio* treatment of the many-body problem

Astrophysical Challenges



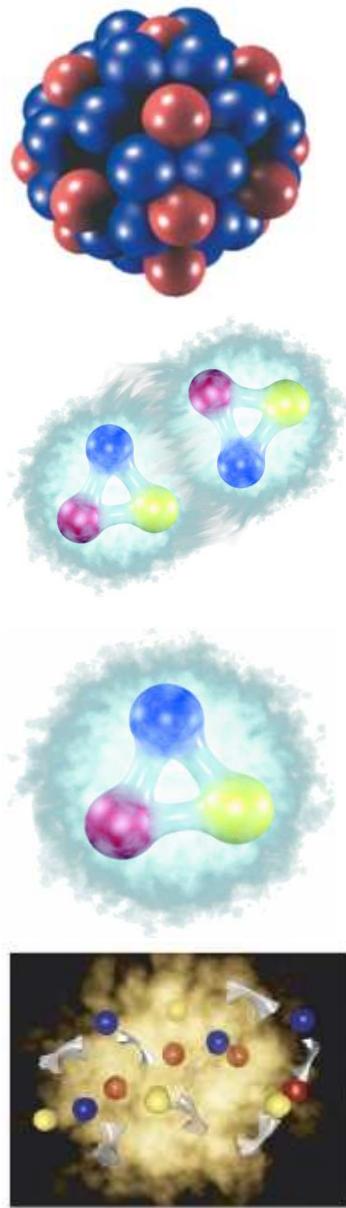
Theoretical Context

better resolution / more fundamental



Quantum Chromo Dynamics

Nuclear Structure



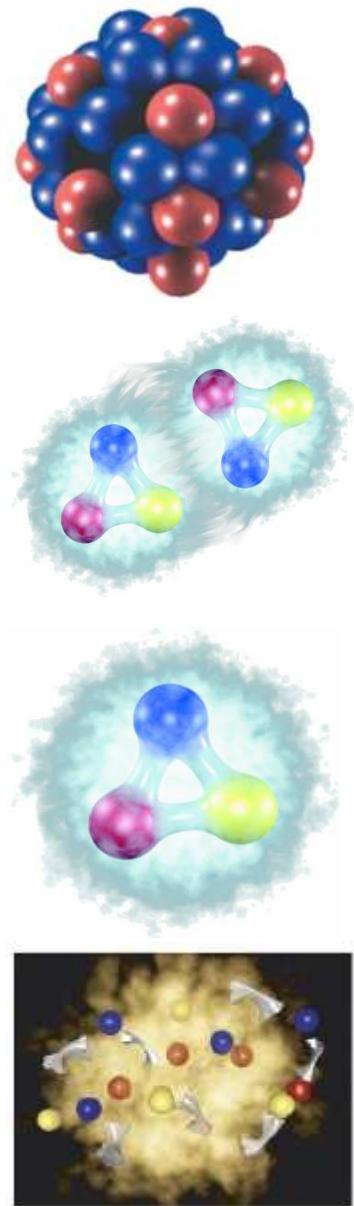
- finite nuclei
- few-nucleon systems
- nucleon-nucleon interaction
- hadron structure
- quarks & gluons
- deconfinement

Theoretical Context

better resolution / more fundamental

Quantum Chromo Dynamics

Nuclear Structure

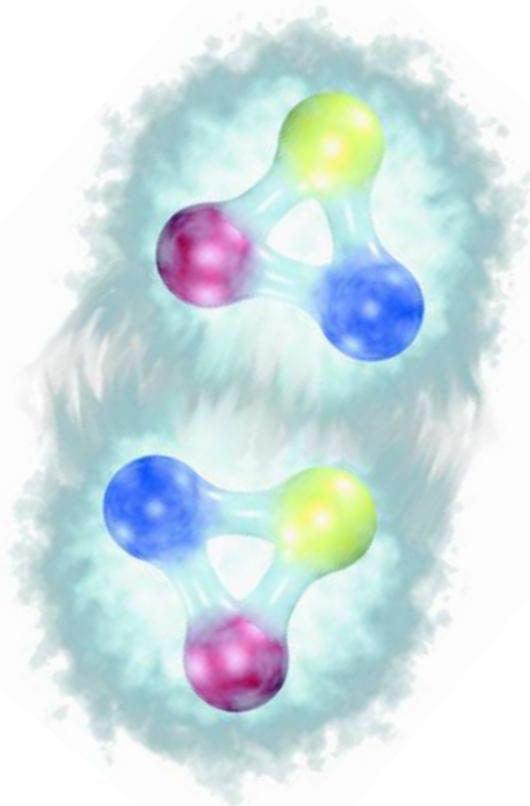


“solve” the quantum
many-body problem with
this interaction

“derive” a realistic
nucleon-nucleon
interaction from QCD

Realistic Nucleon-Nucleon Potentials

Nature of the NN-Interaction



~ 1.6 fm

$$\rho_0^{-1/3} = 1.8 \text{ fm}$$

- NN-interaction is **not fundamental**
- induced via mutual **polarization** of quark & gluon distributions
- analogous to **van der Waals** interaction between neutral atoms
- **short-ranged**: acts only if the nucleons overlap
- genuine **NNN-interaction** is important

How to Construct the NN-Potential?

■ QCD input

- symmetries
- meson-exchange picture
- chiral perturbation theory

■ short-range phenomenology

- ansatz for short-range behavior

■ experimental two-body data

- scattering phase-shifts & deuteron properties
- reproduced with $\chi^2/\text{datum} \approx 1$

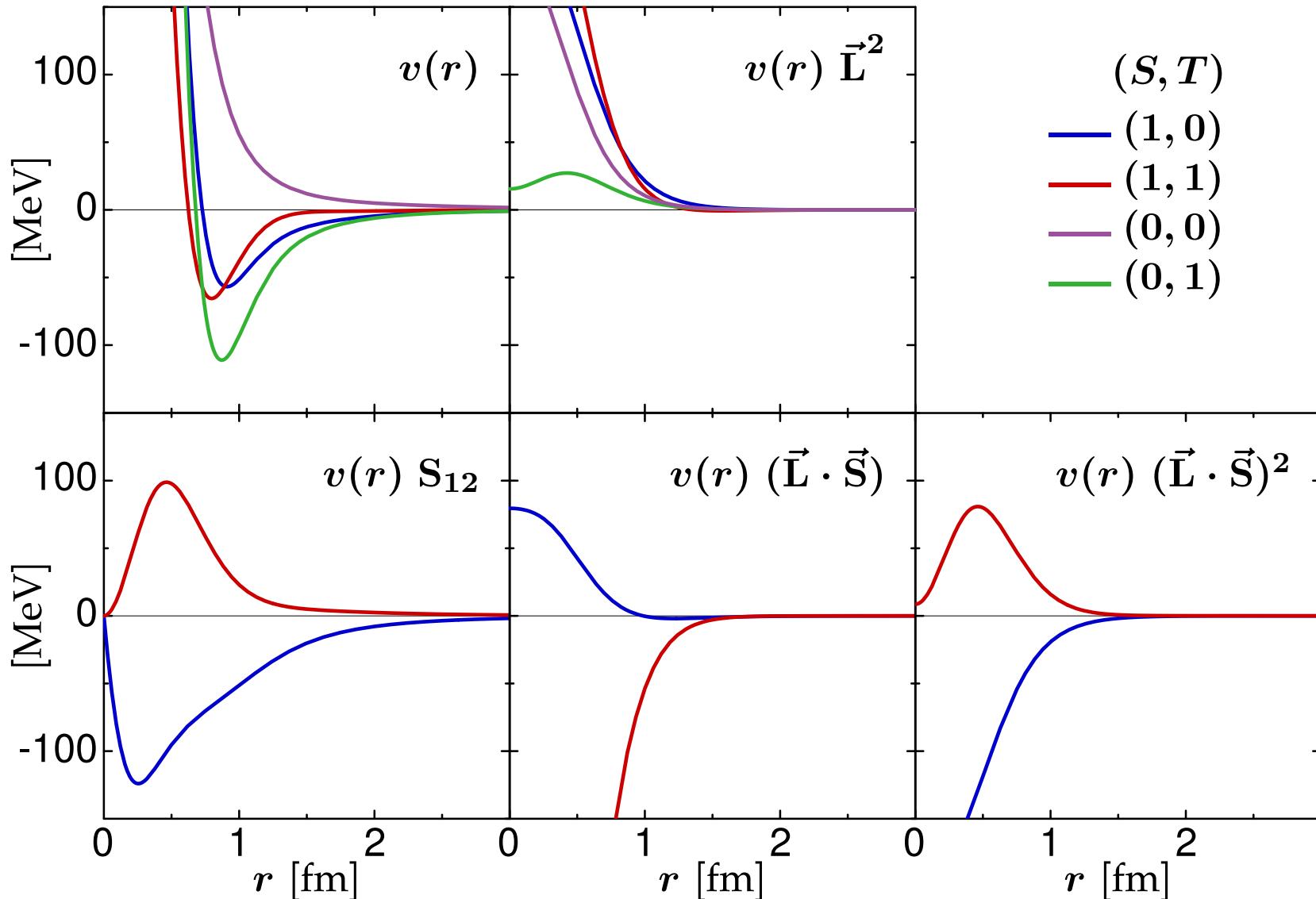
Argonne
V18

CD Bonn

Nijmegen
I/II

Chiral...

Argonne V18 Potential



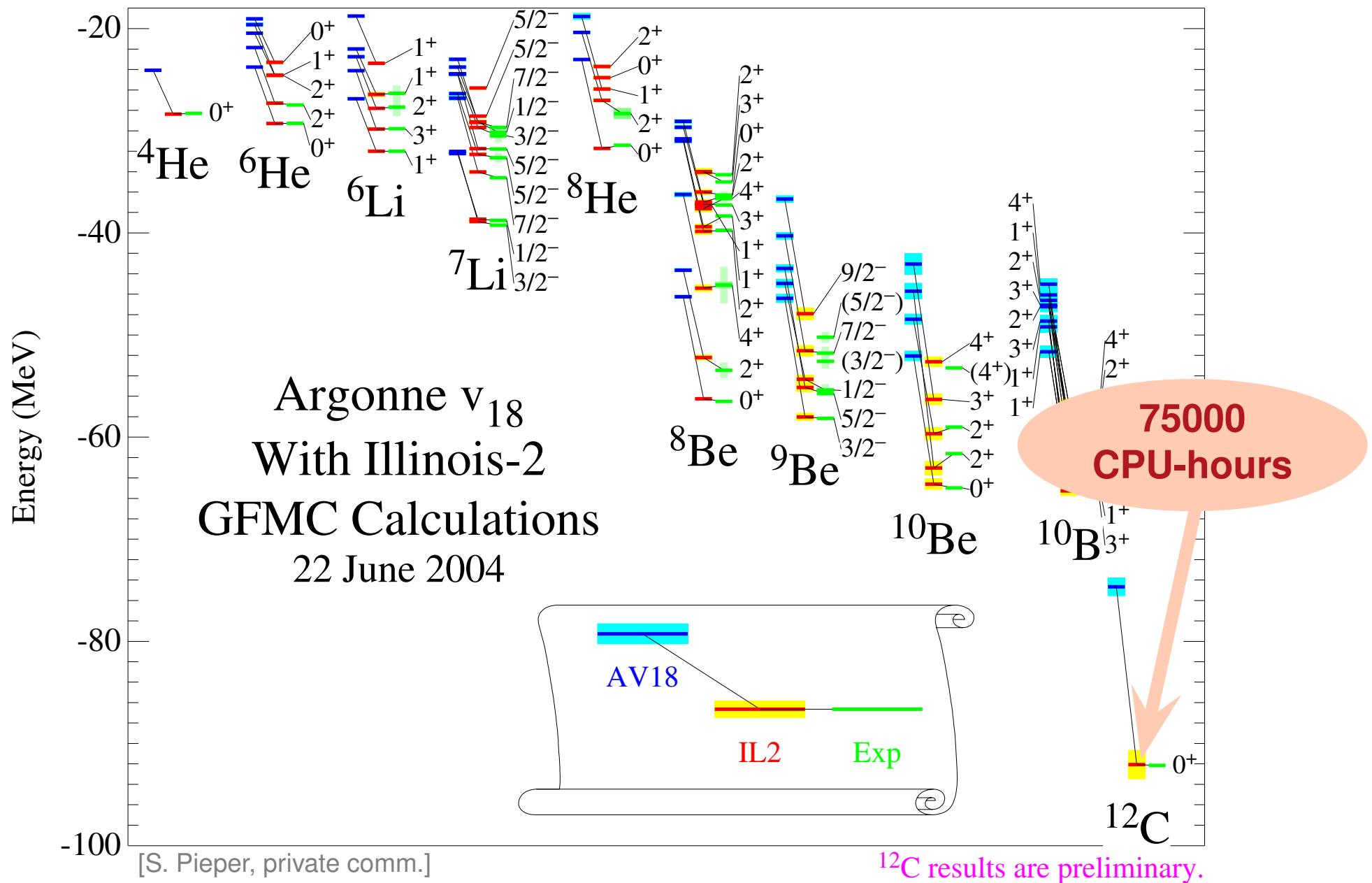
Nuclear Many-Body Problem

Ab initio Calculations

solve the quantum many-body
problem for A nucleons interacting
via a realistic NN-potential

- exact numerical solution possible for small systems at an enormous computational cost
- **Green's Function Monte Carlo**: Monte Carlo sampling of the A -body wave function in coordinate space; imaginary time cooling
- **No-Core Shell Model**: large-scale diagonalization of the Hamiltonian in a harmonic oscillator basis

Green's Function Monte Carlo



Our Goal

nuclear structure calculations
across the whole nuclear chart
based on realistic NN-potentials

bound to **simple**
Hilbert spaces for large
particle numbers

need to deal with
strong **interaction-**
induced correlations

Correlations in Nuclei

What are Correlations?

correlations

=

everything beyond the
independent particle picture

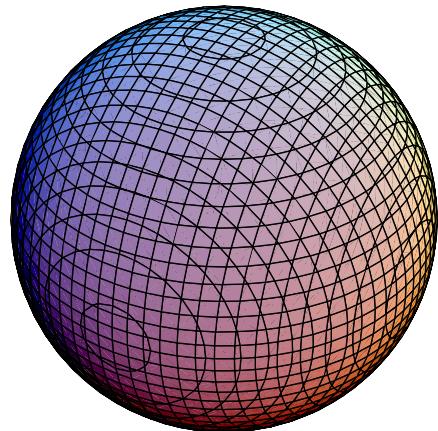
- the quantum state of A independent (non-interacting) fermions is a **Slater determinant**

$$|\psi\rangle = \mathcal{A}(|\phi_1\rangle \otimes |\phi_2\rangle \otimes \cdots \otimes |\phi_A\rangle)$$

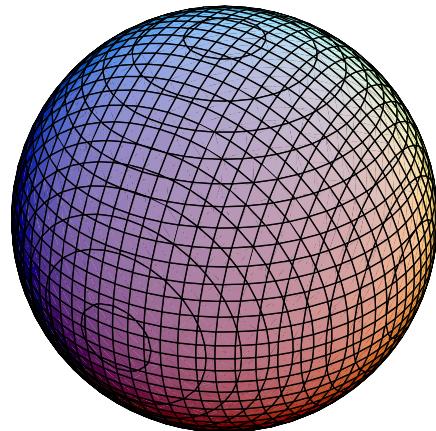
- Slater determinants **cannot describe correlations** by definition

Deuteron: Manifestation of Correlations

$$M_S = 0$$
$$\frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle)$$



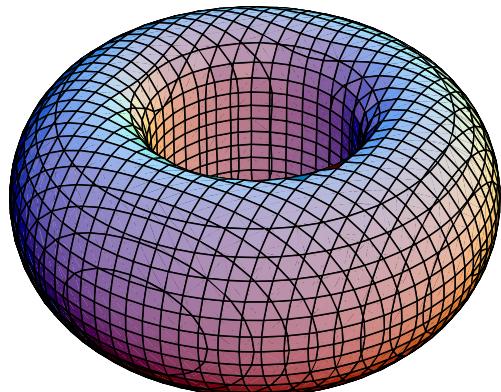
$$M_S = \pm 1$$
$$|\uparrow\uparrow\rangle, |\downarrow\downarrow\rangle$$



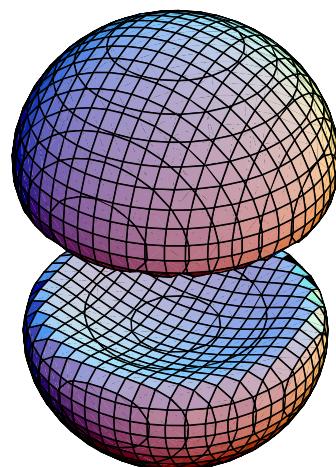
- spin-projected two-body density $\rho_{1,M_S}^{(2)}(\vec{r})$
- **uncorrelated** two-body state

Deuteron: Manifestation of Correlations

$$M_S = 0 \\ \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle)$$



$$M_S = \pm 1 \\ |\uparrow\uparrow\rangle, |\downarrow\downarrow\rangle$$



- spin-projected two-body density $\rho_{1,M_S}^{(2)}(\vec{r})$
- **exact deuteron solution** for Argonne V18 potential

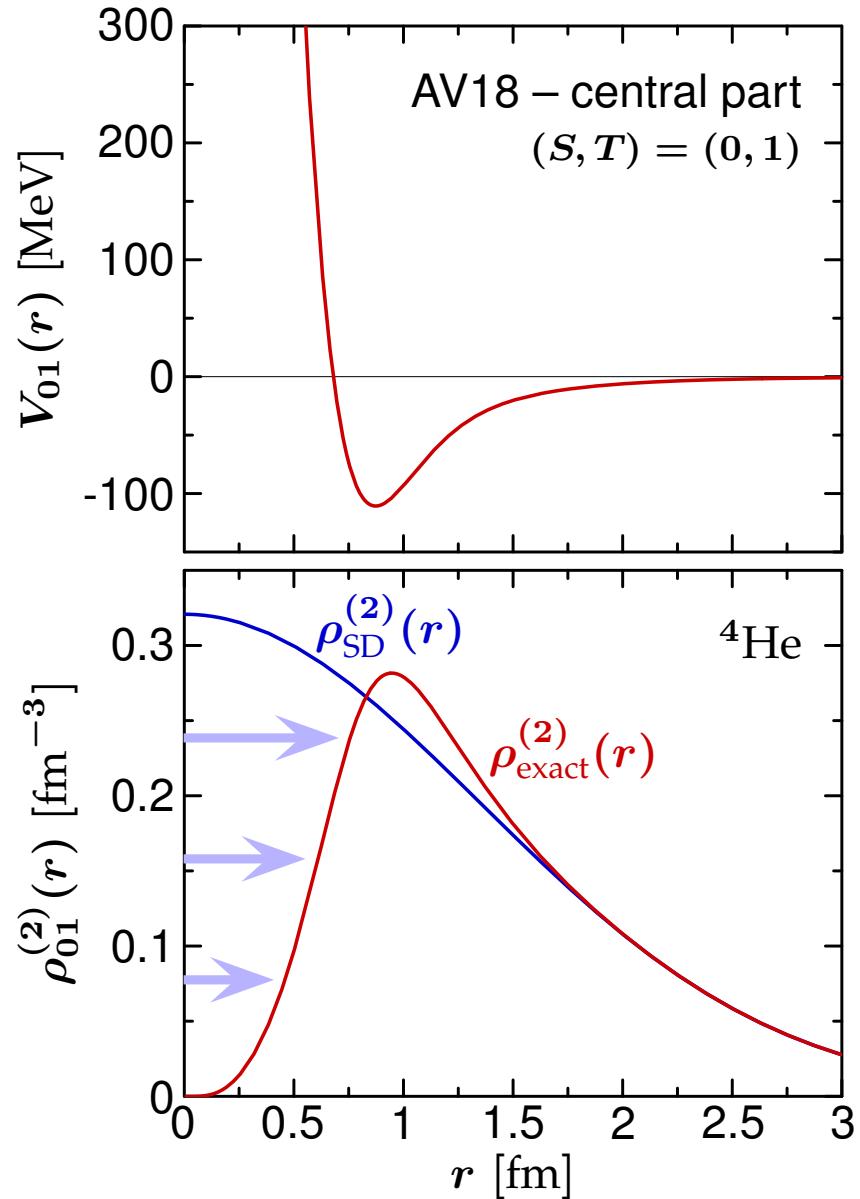
two-body density fully suppressed at small particle distances $|\vec{r}|$

central correlations

angular distribution depends strongly on relative spin orientation

tensor correlations

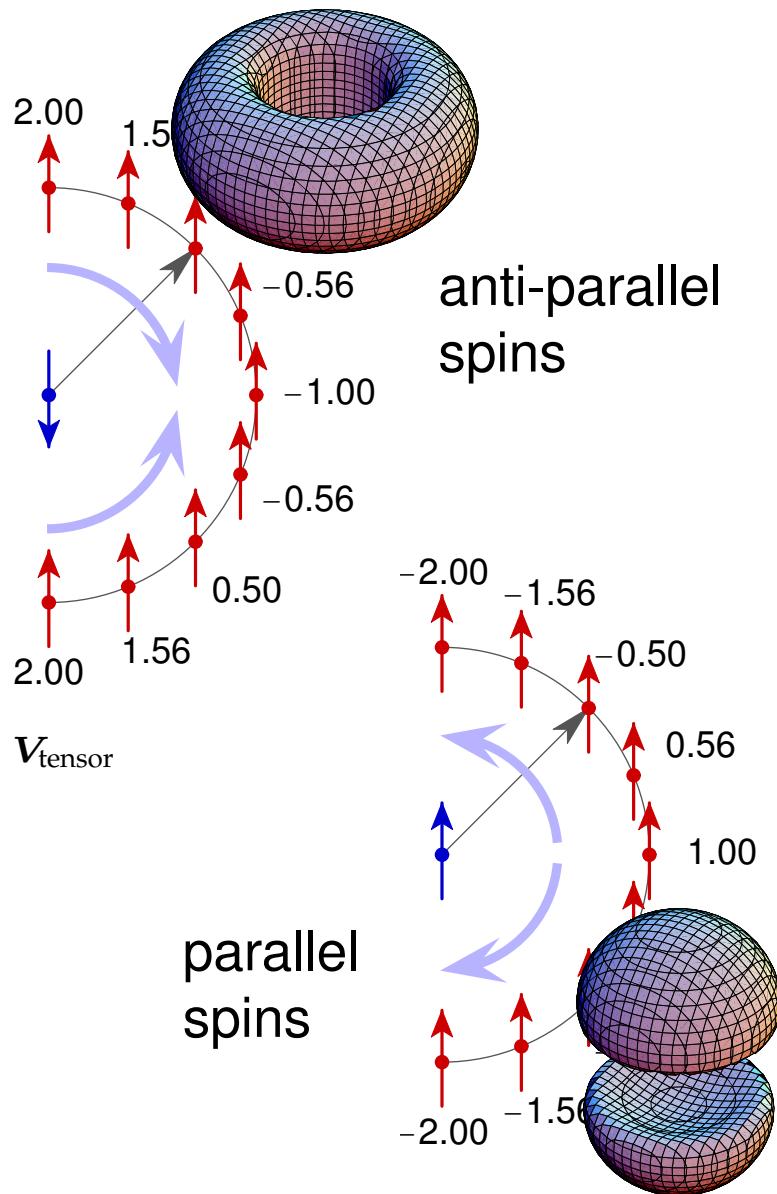
Central Correlations



- strong repulsive core in central part of realistic interactions
- suppression of the probability density for finding two nucleons within the core region → **central correlations**

can be described by
**“shifting” the
nucleons out of the
core region**

Tensor Correlations



- analogy with dipole-dipole interaction

$$V_{\text{tensor}} \sim - \left(3 \frac{(\vec{\sigma}_1 \vec{r})(\vec{\sigma}_2 \vec{r})}{r^2} - \vec{\sigma}_1 \vec{\sigma}_2 \right)$$

- couples the relative spatial orientation of two nucleons with their spin orientation → **tensor correlations**

can be described by
“rotating” nucleons
towards pole or equator
depending on spin

Unitary Correlation Operator Method (UCOM)

Unitary Correlation Operator Method

Correlation Operator

introduce correlations by means of an unitary transformation with respect to the relative coordinates of all pairs

$$\mathbf{C} = \exp[-i G] = \exp\left[-i \sum_{i < j} g_{ij}\right]$$
$$g = g(\vec{r}, \vec{q}; \vec{\sigma}_1, \vec{\sigma}_2, \vec{\tau}_1, \vec{\tau}_2)$$

$$G^\dagger = G$$
$$C^\dagger C = 1$$

Correlated States

$$|\hat{\psi}\rangle = \mathbf{C} |\psi\rangle$$

Correlated Operators

$$\hat{O} = \mathbf{C}^\dagger O \mathbf{C}$$

$$\langle \hat{\psi} | O | \hat{\psi}' \rangle = \langle \psi | \mathbf{C}^\dagger O \mathbf{C} | \psi' \rangle = \langle \psi | \hat{O} | \psi' \rangle$$

Central and Tensor Correlators

$$\mathbf{C} = \mathbf{C}_\Omega \mathbf{C}_r$$

Central Correlator \mathbf{C}_r

- radial distance-dependent shift in the relative coordinate of a nucleon pair

$$g_r = \frac{1}{2} [s(r) \mathbf{q}_r + \mathbf{q}_r s(r)]$$

$$\mathbf{q}_r = \frac{1}{2} [\vec{\mathbf{r}} \cdot \vec{\mathbf{q}} + \vec{\mathbf{q}} \cdot \vec{\mathbf{r}}]$$

Tensor Correlator \mathbf{C}_Ω

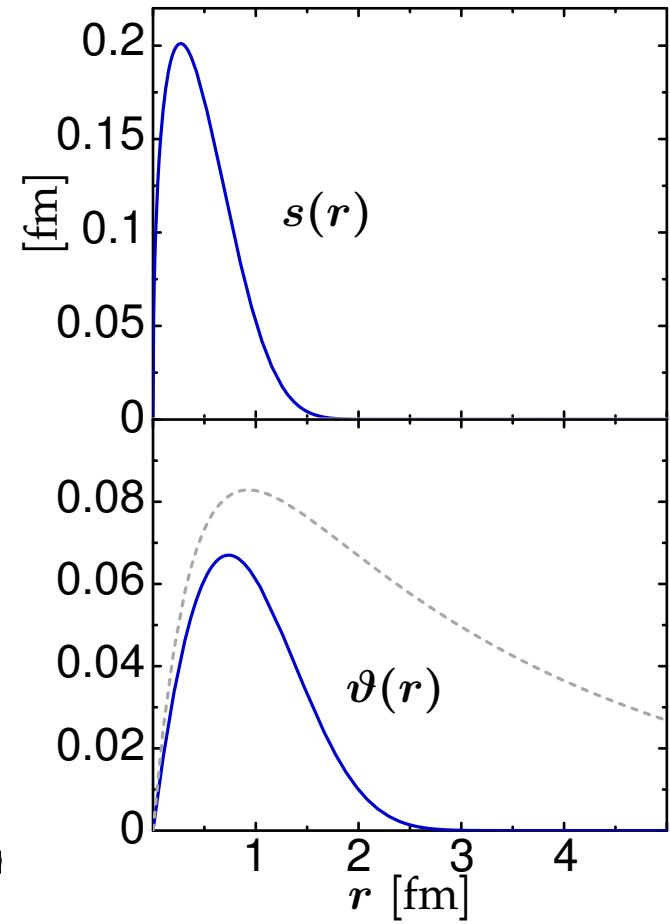
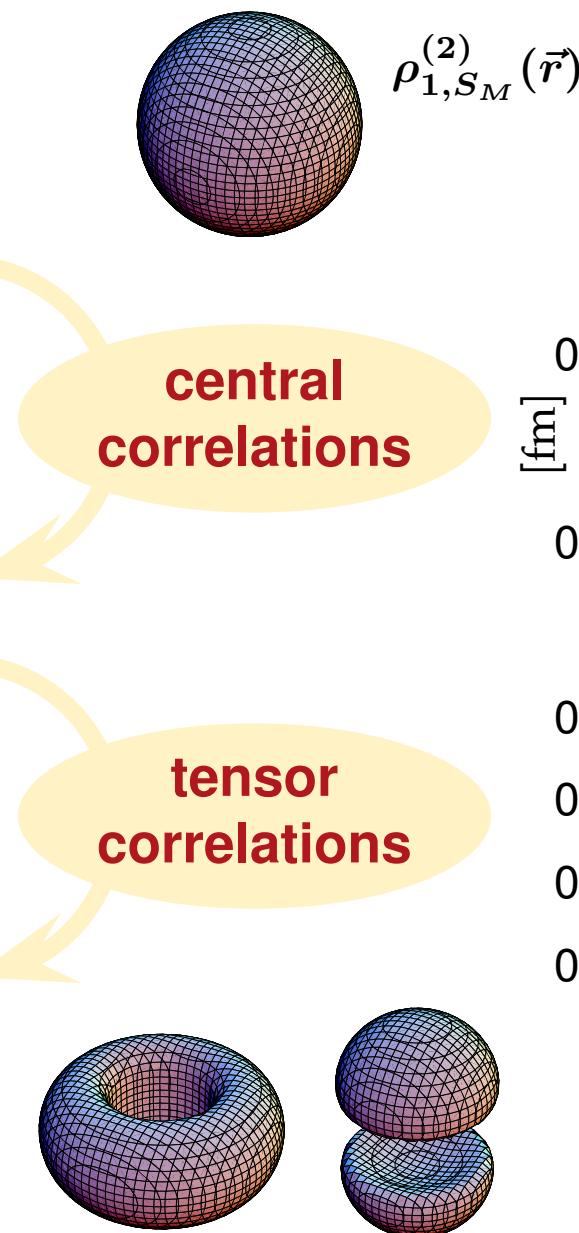
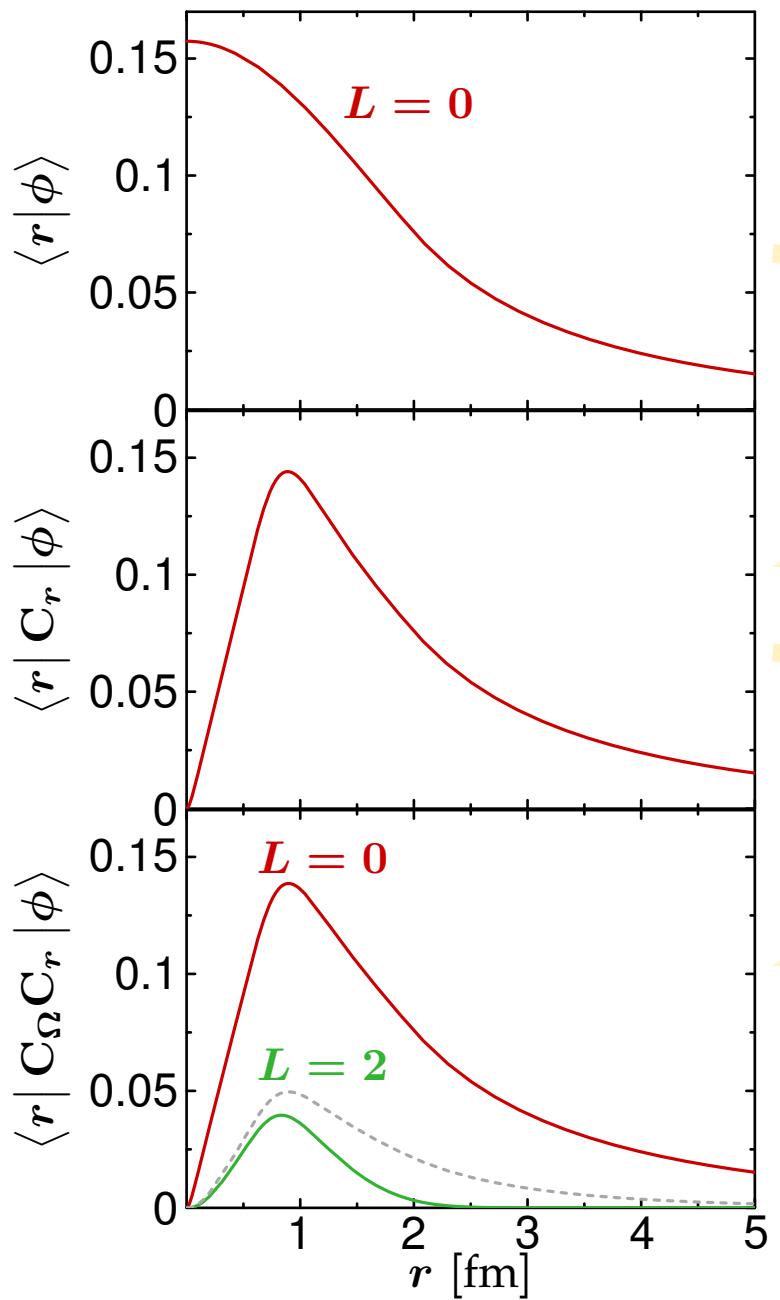
- angular shift depending on the orientation of spin and relative coordinate of a nucleon pair

$$g_\Omega = \frac{3}{2} \vartheta(r) [(\vec{\sigma}_1 \cdot \vec{\mathbf{q}}_\Omega)(\vec{\sigma}_2 \cdot \vec{\mathbf{r}}) + (\vec{\mathbf{r}} \leftrightarrow \vec{\mathbf{q}}_\Omega)]$$

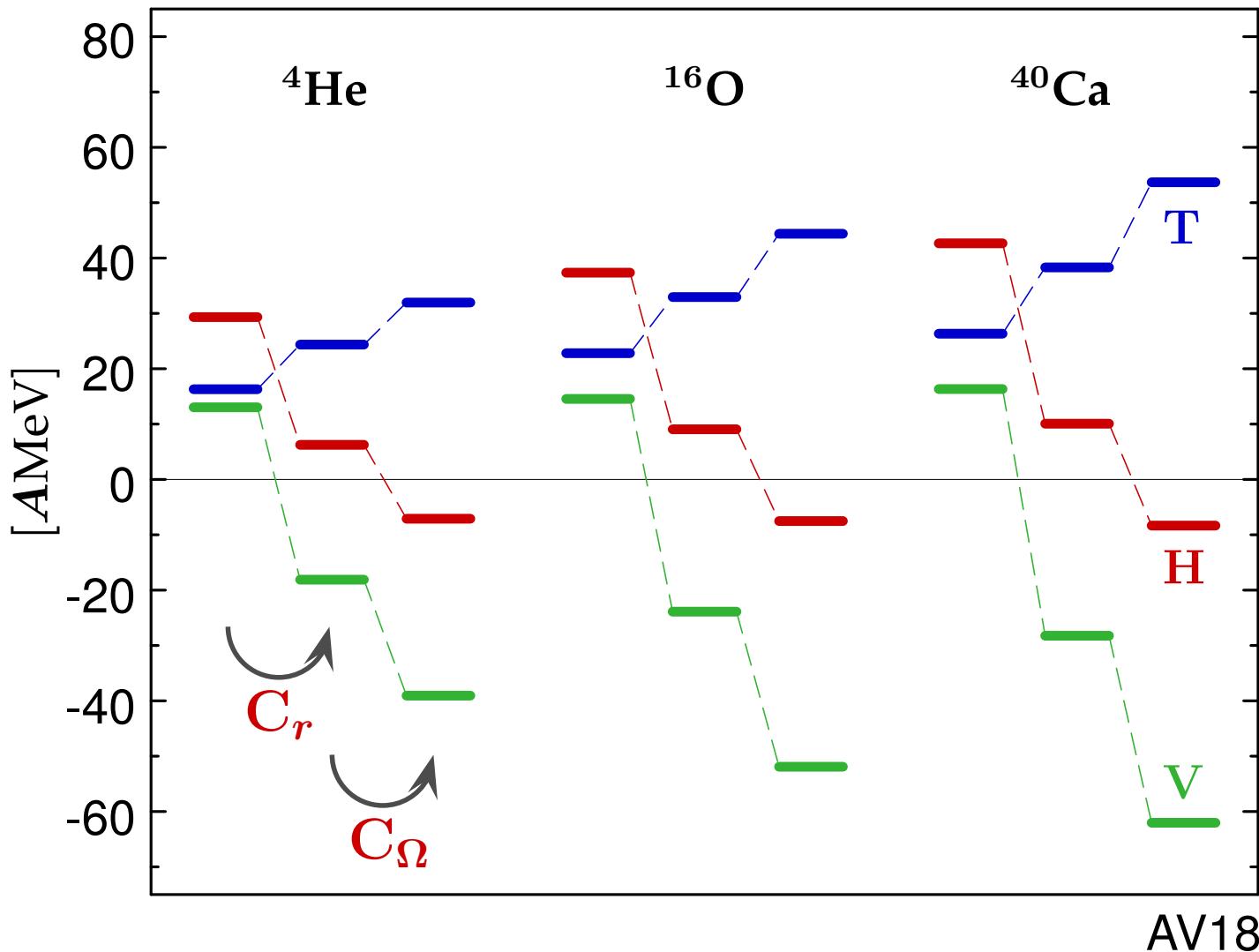
$$\vec{\mathbf{q}}_\Omega = \vec{\mathbf{q}} - \frac{\vec{\mathbf{r}}}{r} \mathbf{q}_r$$

$s(r)$ and $\vartheta(r)$
encapsulate the physics of
short-range correlations

Correlated States



Simplistic “Shell-Model” Calculation



- expectation values for harmonic osc. Slater determinant
- nuclei unbound without inclusion of correlations
- central and tensor correlations essential to obtain bound system

Application I:

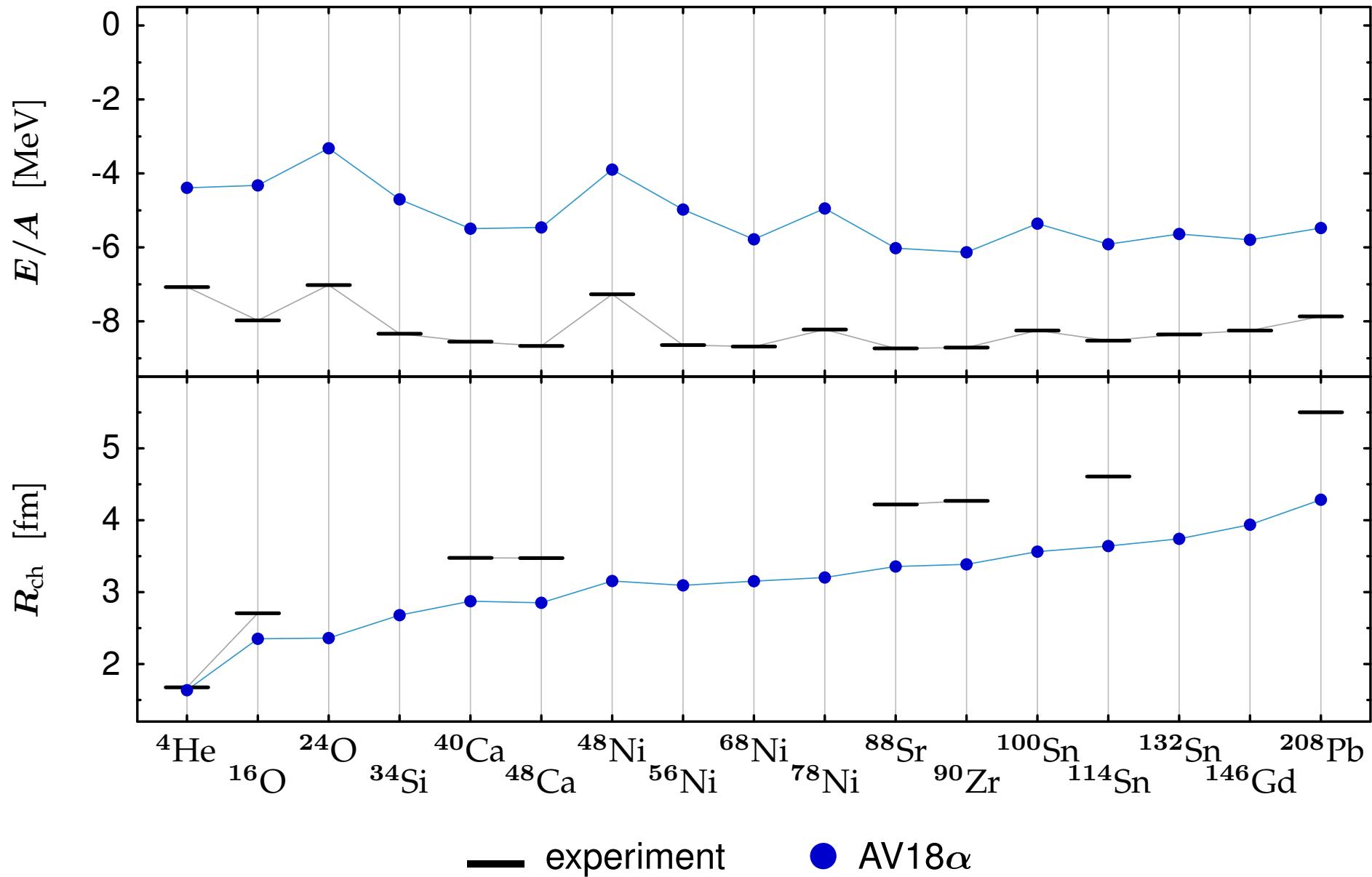
Hartree-Fock Calculations

UCOM-Hartree-Fock Approach

Standard Hartree-Fock
+
Matrix Elements of Correlated
Realistic NN-Interaction V_{UCOM}

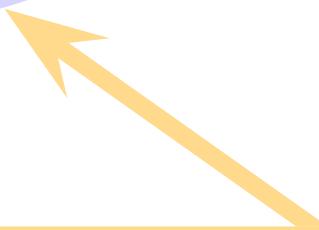
- many-body state is a **Slater determinant** of single-particle states obtained by energy minimization
- **correlations cannot be described** by Hartree-Fock states
- bare realistic NN-potential leads to **unbound nuclei**

Correlated Argonne V18



Missing Pieces

long-range correlations



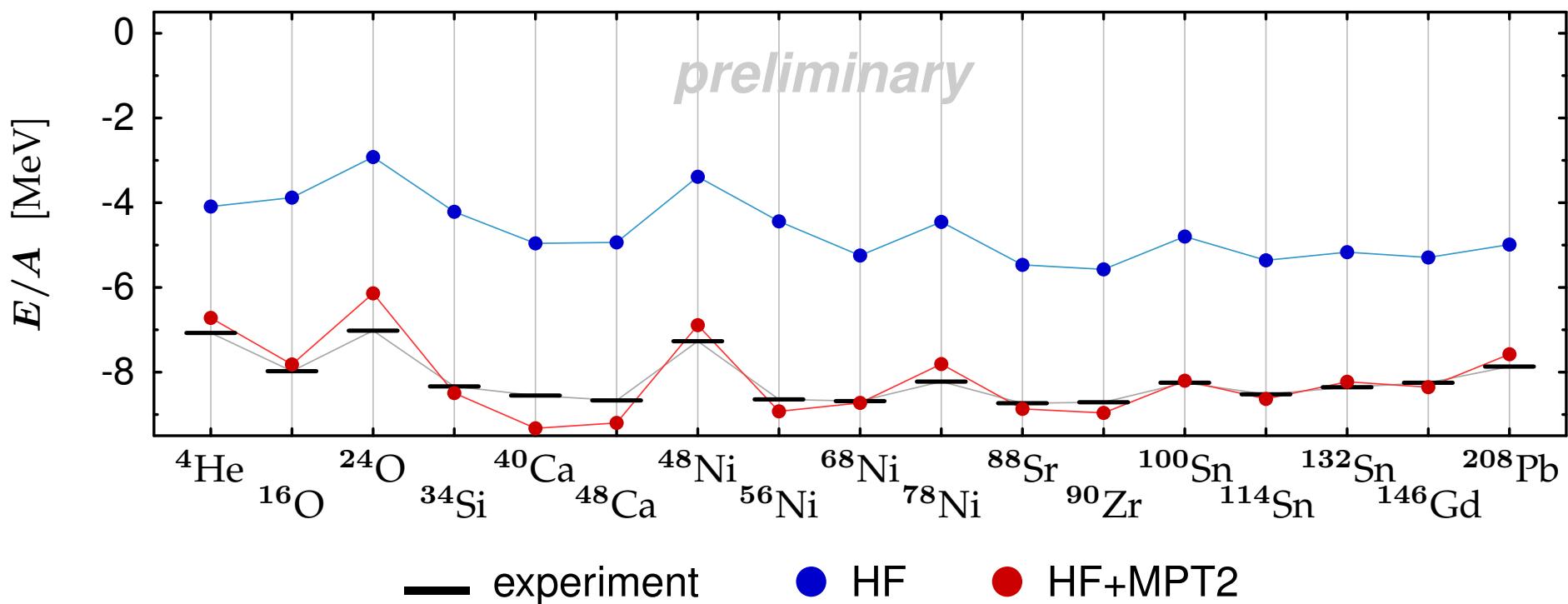
Ab Initio Strategy

- improve many-body states such that long-range correlations are included
- many-body perturbation theory (MPT), configuration interaction (CI), coupled-cluster (CC),...

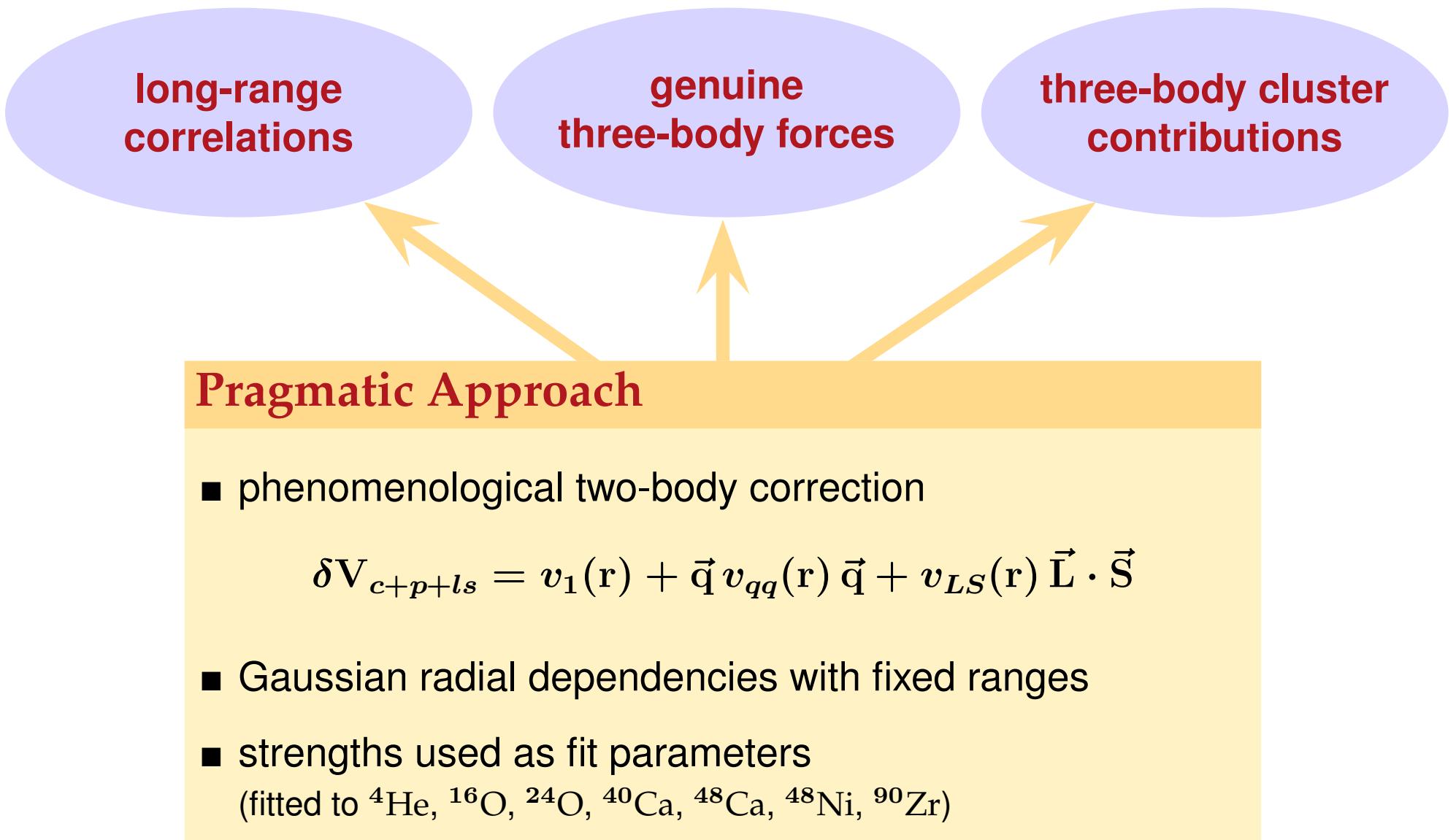
Long-Range Correlations

- **many-body perturbation theory**: second-order energy shift gives estimate for influence of long-range correlations

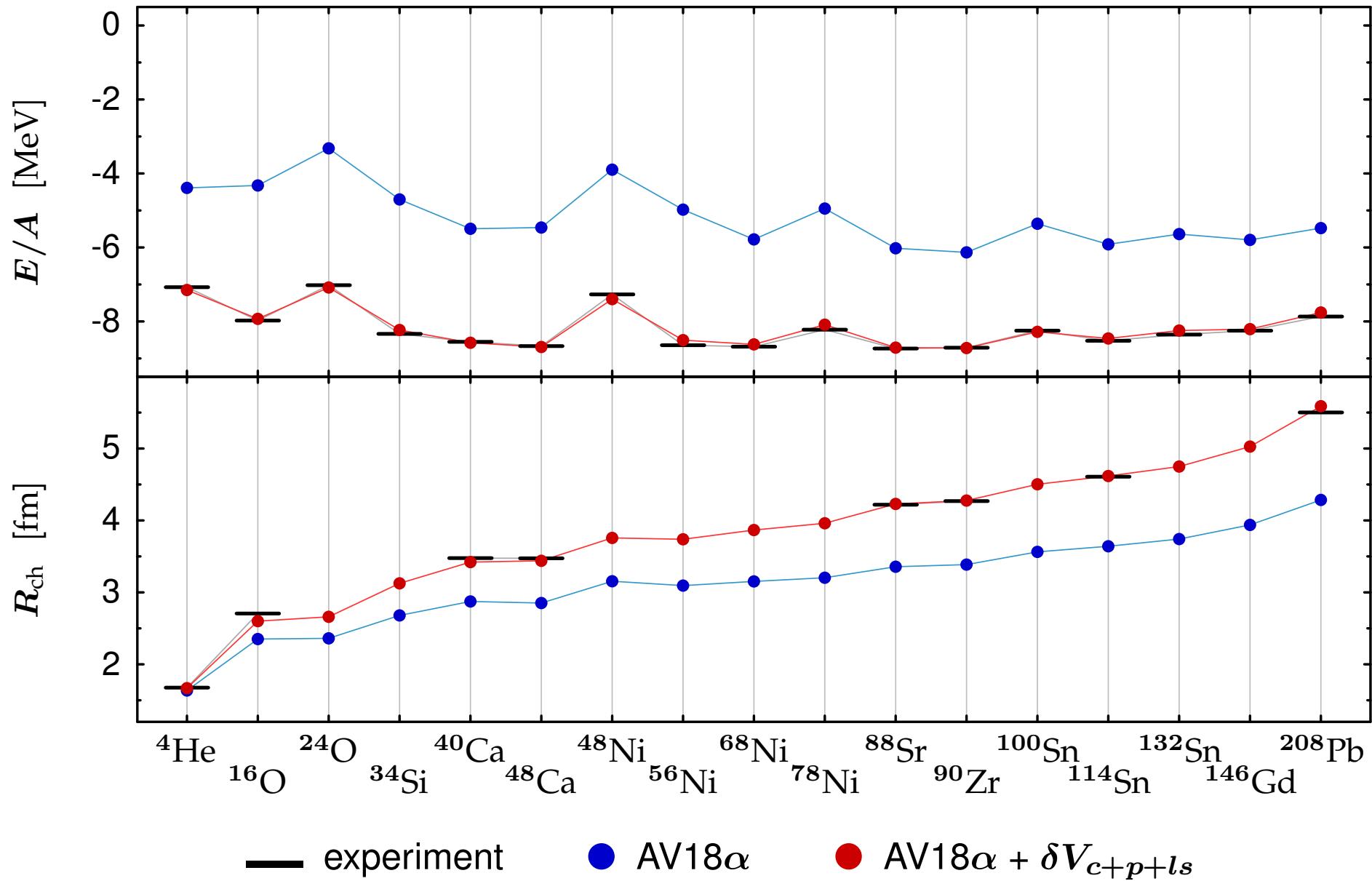
$$\Delta E^{(2)} = -\frac{1}{4} \sum_{i,j}^{\text{occu. unoccu.}} \sum_{a,b} \frac{|\langle \phi_a \phi_b | V_{\text{UCOM}} | \phi_i \phi_j \rangle|^2}{\epsilon_a + \epsilon_b - \epsilon_i - \epsilon_j}$$



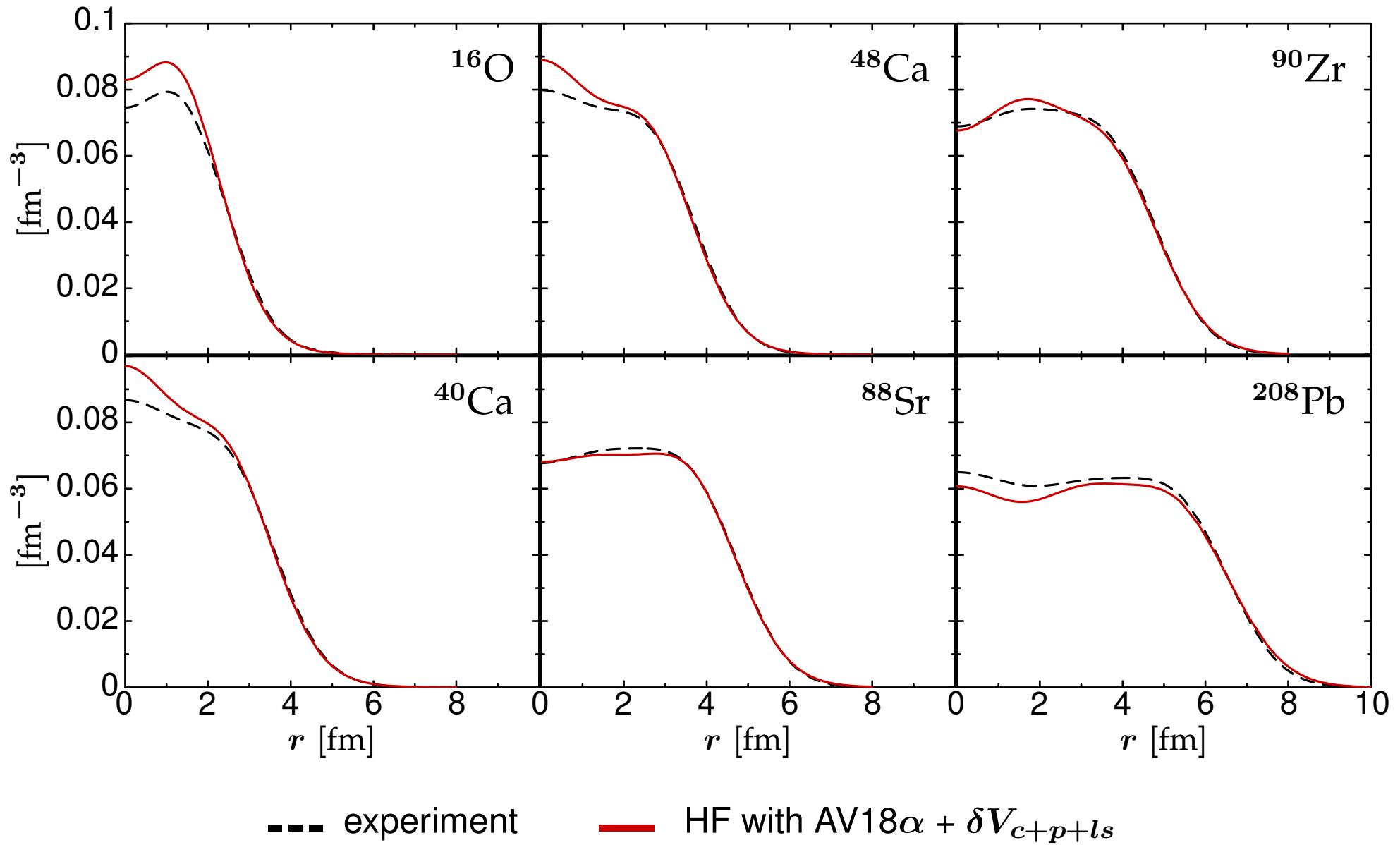
Missing Pieces



Correlated Argonne V18 + Correction



Charge Distributions



Application II

Fermionic Molecular Dynamics (FMD)

UCOM-FMD Approach

Gaussian Single-Particle States

$$|q\rangle = \sum_{\nu=1}^n \mathbf{c}_{\nu} \ |a_{\nu}, \vec{b}_{\nu}\rangle \otimes |\chi_{\nu}\rangle \otimes |m_t\rangle$$

$$\langle \vec{x} | a_{\nu}, \vec{b}_{\nu} \rangle = \exp \left[- \frac{(\vec{x} - \vec{b}_{\nu})^2}{2 a_{\nu}} \right]$$

a_{ν} : complex width

χ_{ν} : spin orientation

\vec{b}_{ν} : mean position & momentum

Variation

$$\frac{\langle Q | \hat{H} - T_{cm} | Q \rangle}{\langle Q | Q \rangle} \rightarrow \min$$

Slater Determinant

$$|Q\rangle = \mathcal{A} (|q_1\rangle \otimes |q_2\rangle \otimes \cdots \otimes |q_A\rangle)$$

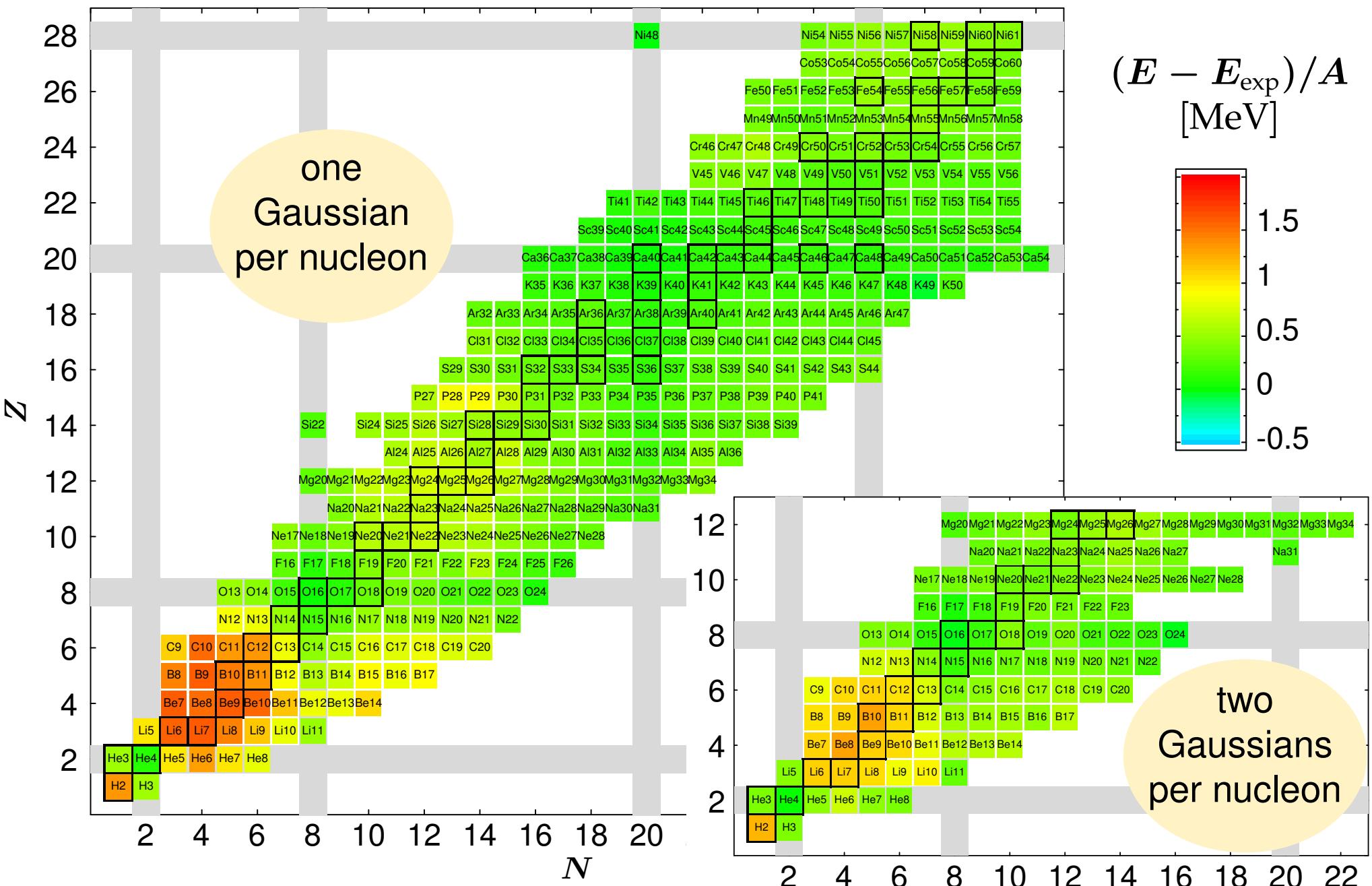
Diagonalization

in sub-space spanned
by several non-ortho-
gonal Slater deter-
minants $|Q_i\rangle$

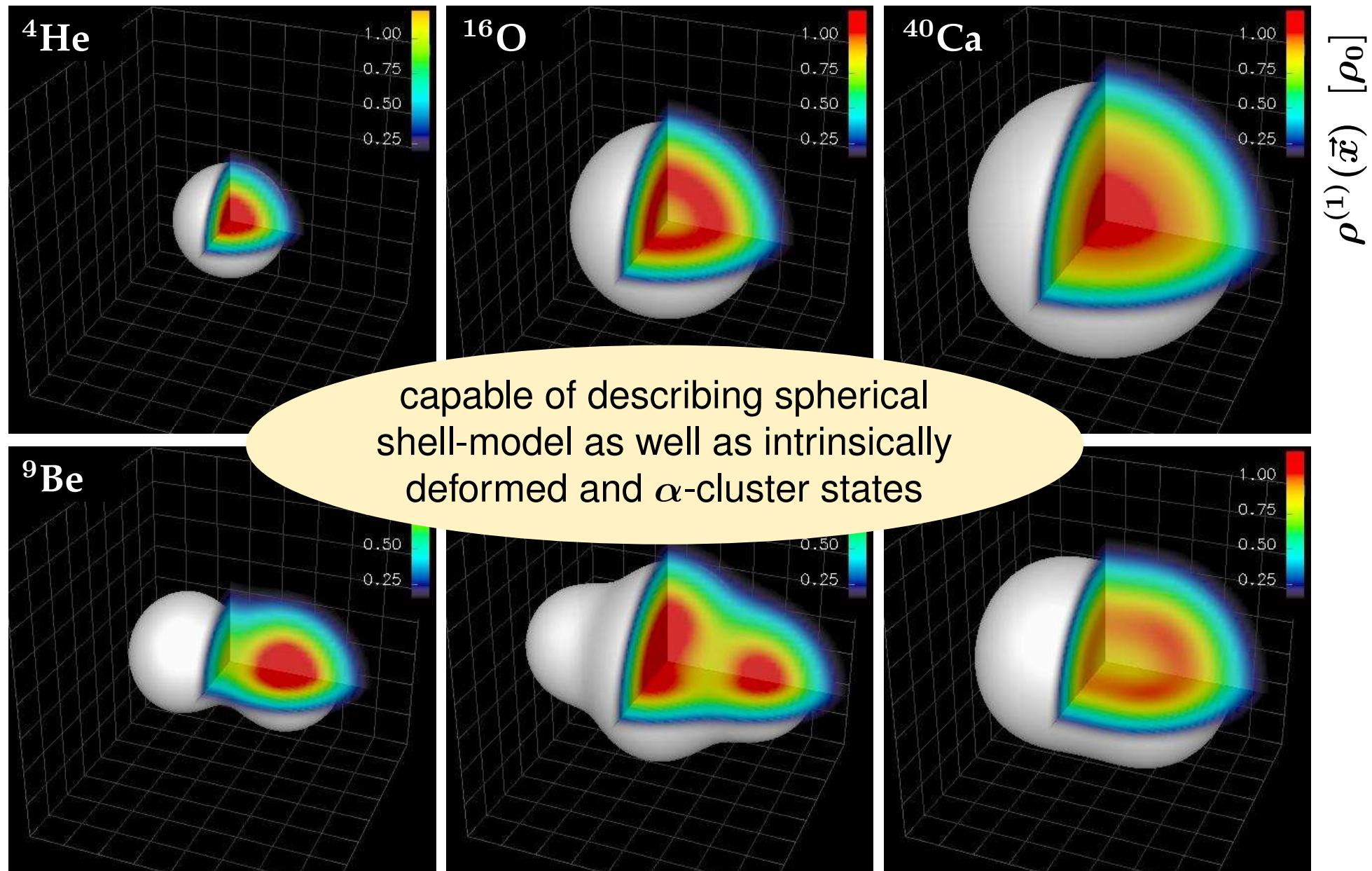
Correlated Hamiltonian

$$\hat{H} = T + V_{UCOM} [+ \delta V_{c+p+ls}]$$

Variation: Chart of Nuclei



Intrinsic One-Body Density Distributions



Beyond Simple Variation

■ Projection after Variation (PAV)

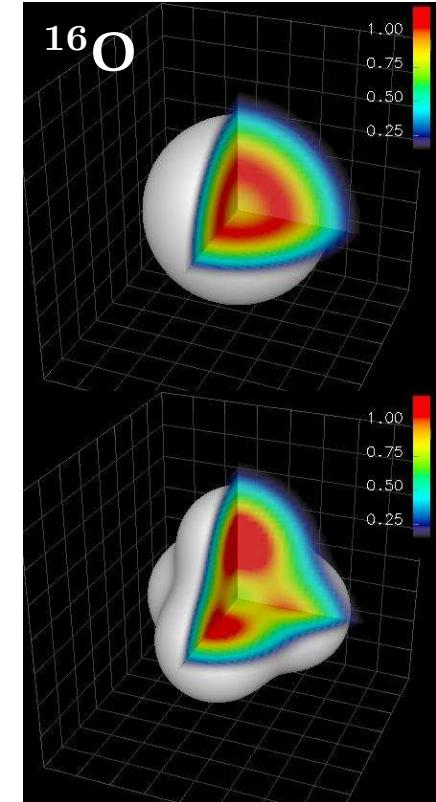
- restore inversion and rotational symmetry by angular momentum projection

■ Variation after Projection (VAP)

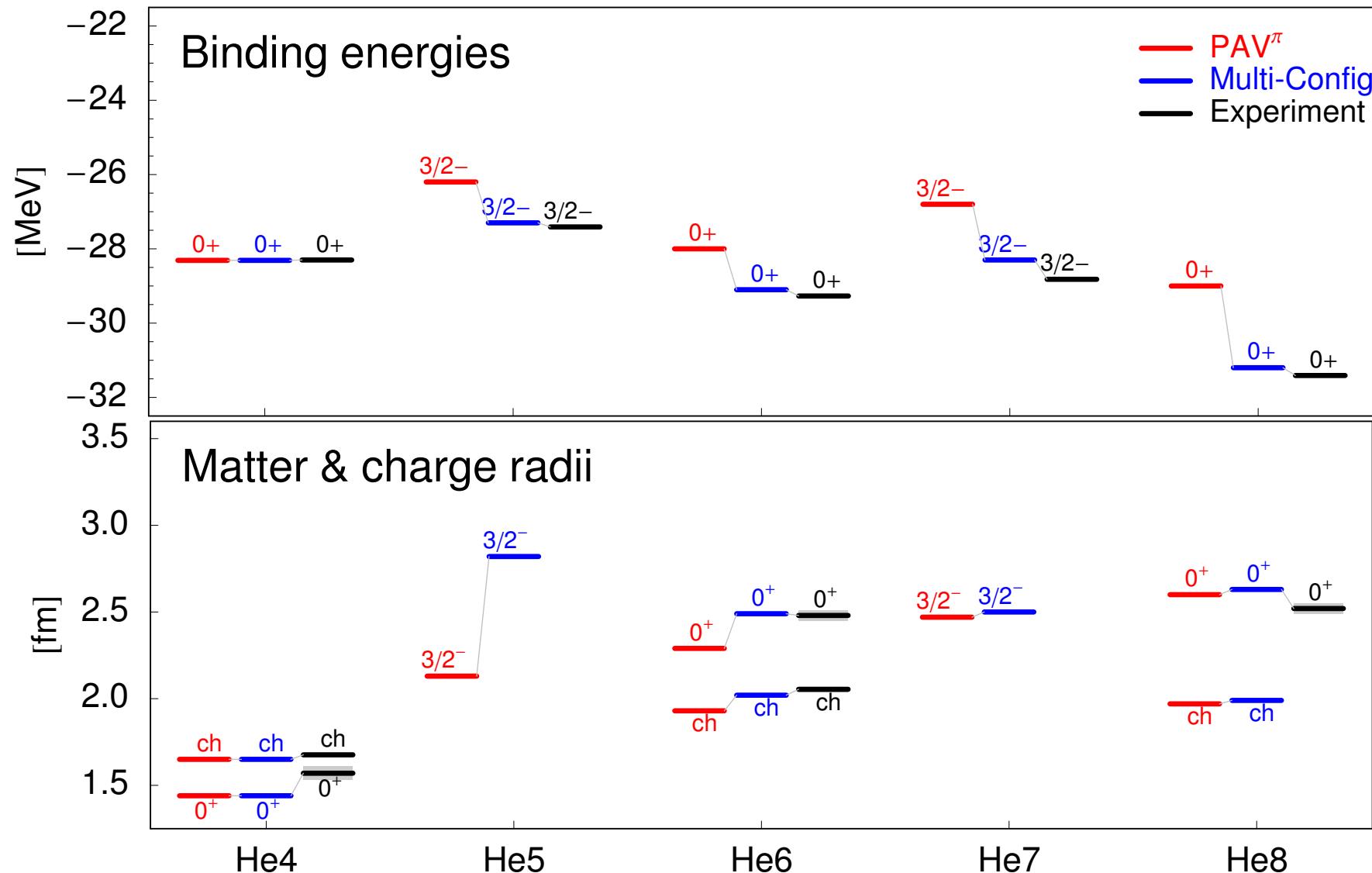
- find energy minimum within parameter space of parity and angular momentum projected states
- implementation via generator coordinate method (constraints on multipole moments)

■ Multi-Configuration

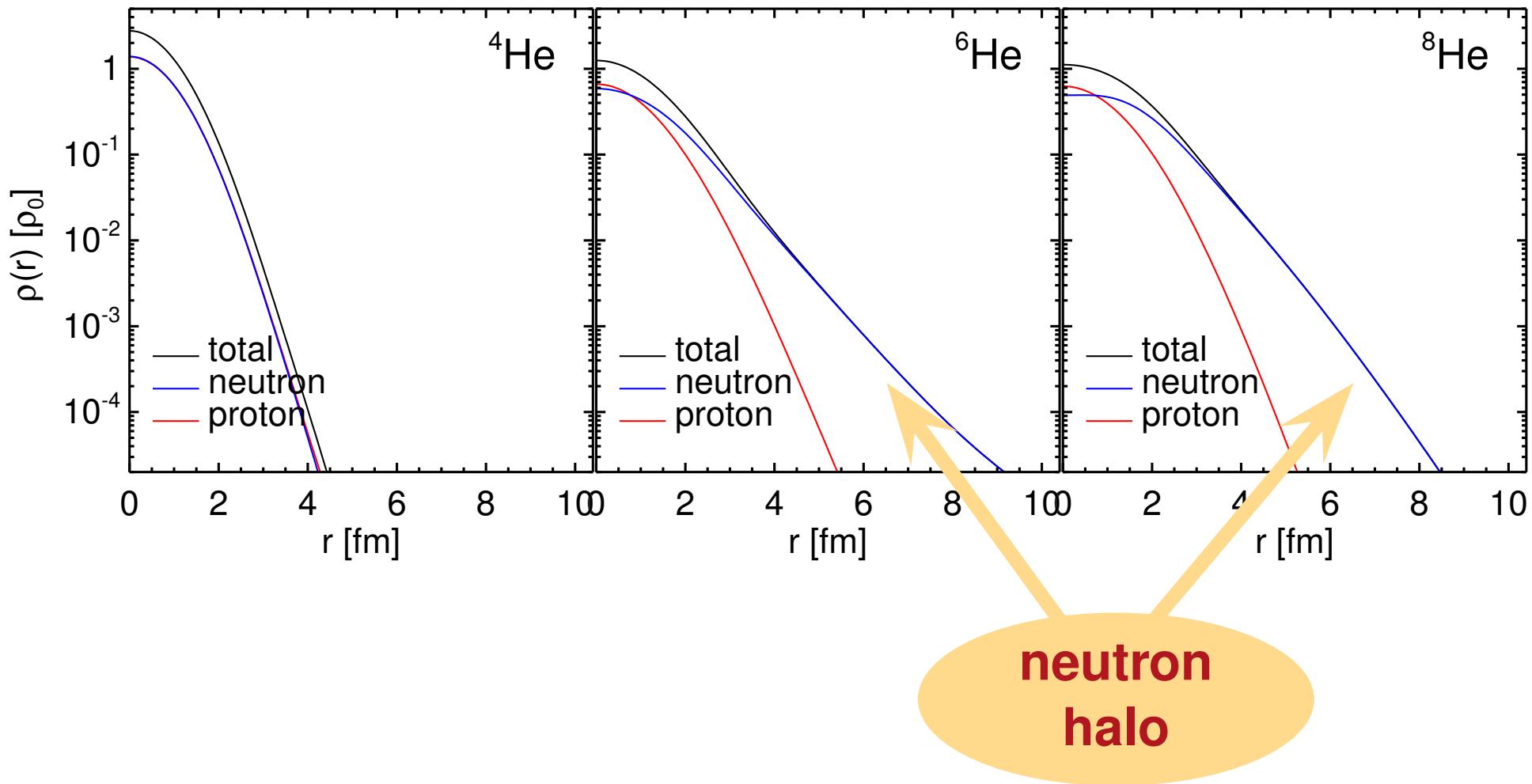
- diagonalization within a set of different Slater determinants



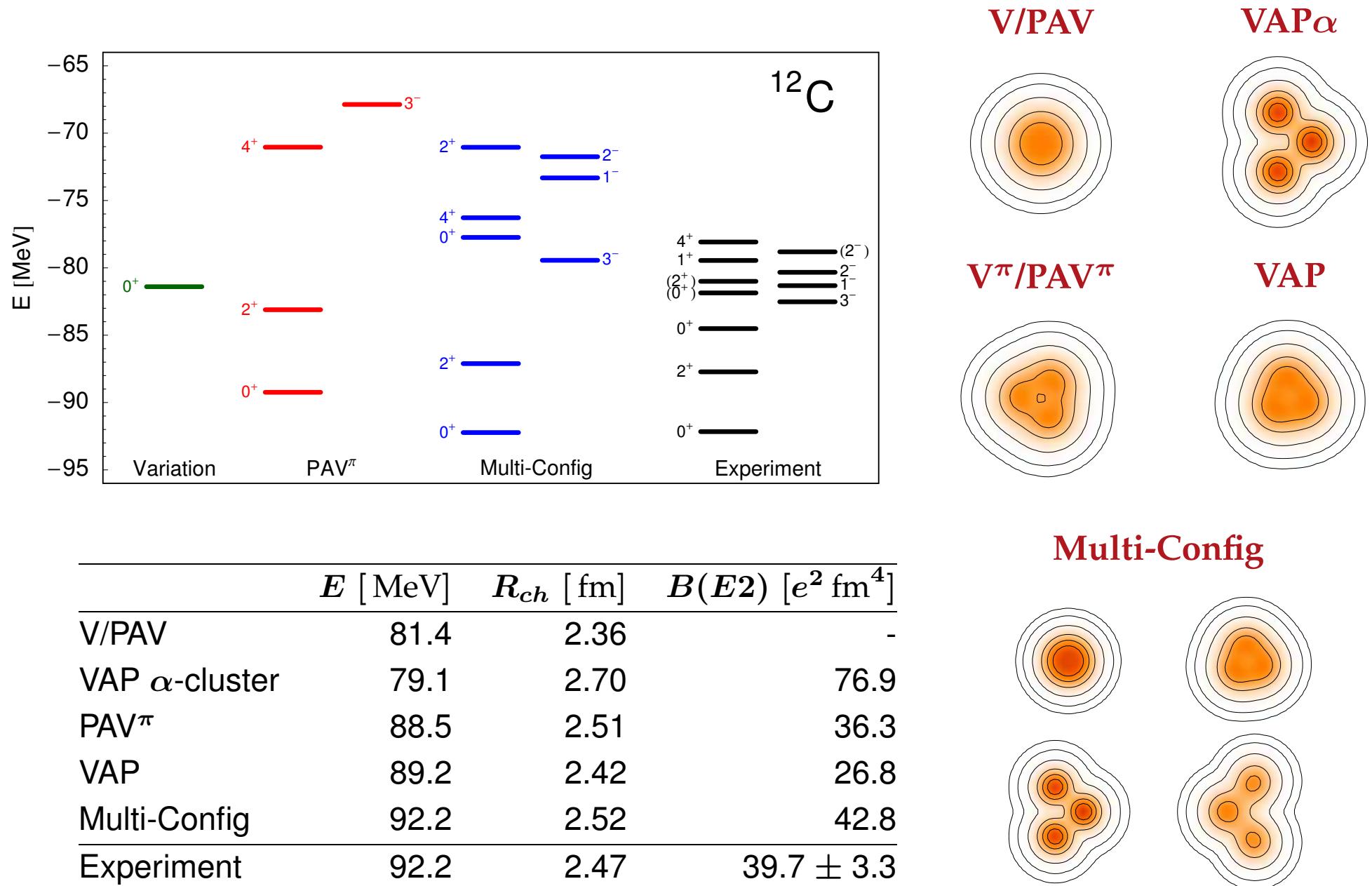
Helium Isotopes: Energies & Radii



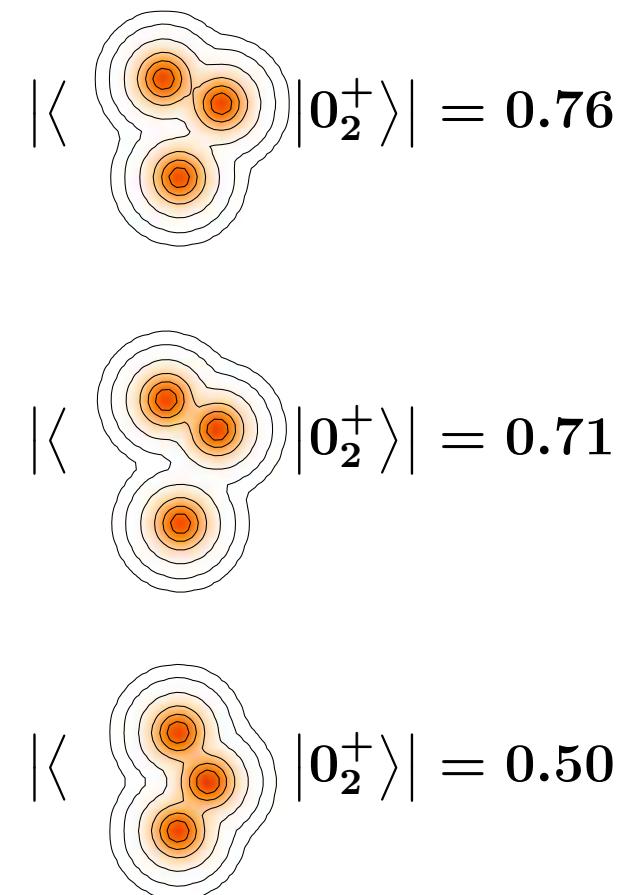
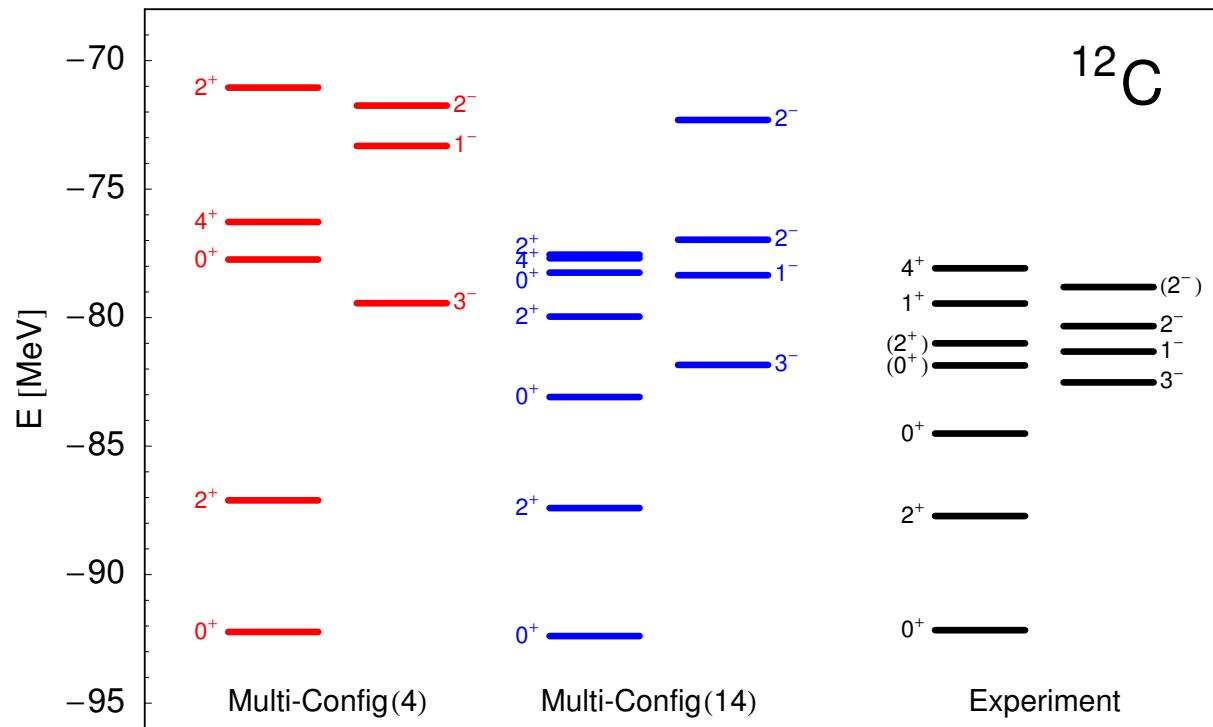
Helium Isotopes: Density Profiles



Structure of ^{12}C



Structure of ^{12}C — Hoyle State



	Multi-Config	Experiment
E [MeV]	92.4	92.2
R_{ch} [fm]	2.52	2.47
$B(E2, 0_1^+ \rightarrow 2_1^+) [e^2 \text{ fm}^4]$	42.9	39.7 ± 3.3
$M(E0, 0_1^+ \rightarrow 0_2^+) [\text{fm}^2]$	5.67	5.5 ± 0.2

Conclusions

- exciting times for nuclear structure physics!
- realistic NN-potentials & *ab initio* calculations
- systematic schemes to derive effective (correlated / low-momentum) interactions
- innovative ways to treat the many-body problem

**unified description of nuclear
structure across the whole
nuclear chart is within reach**

Epilogue

■ thanks to my group & my collaborators

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