Hartree-Fock and Hartree-Fock-Bogoliubov with Modern Effective Interactions

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# Overview

#### Motivation

- Modern Effective Interactions
  - Unitary Correlation Operator Method
  - Similarity Renormalization Group
- Few-Body Systems
- Many-Body Systems
  - Hartree-Fock, Many-Body Perturbation Theory & Beyond
  - Hartree-Fock-Bogoliubov
- Conclusions

#### **Nuclear Structure**



#### **Nuclear Structure**



- chiral interactions: consistent NN & 3N interaction derived within *x*EFT
- traditional NN-interactions: Argonne V18, CD Bonn,...
- reproduce experimental NN phaseshifts with high precision
- induce strong short-range central & tensor correlations

#### **Nuclear Structure**

#### Exact / Approx. Many-Body Methods

- 'exact' solution of the many-body problem for light and intermediate masses (GFMC, NCSM, CC,...)
- controlled approximations for heavier nuclei (HF & MBPT,...)
- rely on restricted model spaces of tractable size
- not suitable for the description of short-range correlations



#### **Nuclear Structure**

#### Exact / Approx. Many-Body Methods

#### Modern Effective Interactions

#### Realistic Nuclear Interactions

Low-Energy QCD

#### adapt realistic potential to the available model space

- tame short-range correlations
- improve convergence behavior
- conserve experimentally constrained properties (phase shifts)
  - generate new realistic interaction
- provide consistent effective interaction & effective operators
- unitary transformations most convenient

# Deuteron: Manifestation of Correlations

#### **Argonne V18 Deuteron Solution**



short-range repulsion supresses wavefunction at small distances *r* **central correlations**  tensor interaction generates D-wave admixture in the ground state tensor correlations

Modern Effective Interactions I

Unitary Correlation Operator Method (UCOM)

## Unitary Correlation Operator Method

#### **Correlation Operator**

define an unitary operator C to describe the effect of short-range correlations

$$\mathbf{C} = \exp[-\mathrm{i}\,\mathrm{G}] = \expigg[-\mathrm{i}\sum_{i < j}\mathrm{g}_{ij}igg]$$

#### **Correlated States**

imprint short-range correlations onto uncorrelated many-body states

$$\left| \widetilde{\psi} 
ight
angle = {f C} \; \left| \psi 
ight
angle$$

#### **Correlated Operators**

adapt Hamiltonian and all other observables to uncorrelated many-body space

 $\widetilde{\mathbf{O}} = \mathbf{C}^{\dagger} \; \mathbf{O} \; \mathbf{C}$ 

$$\left\langle \widetilde{\psi} \right| \mathbf{O} \left| \widetilde{\psi'} \right\rangle = \left\langle \psi \right| \mathbf{C^{\dagger}} \mathbf{O} \mathbf{C} \left| \psi' \right\rangle = \left\langle \psi \right| \widetilde{\mathbf{O}} \left| \psi' \right\rangle$$

### Unitary Correlation Operator Method

explicit ansatz for the correlation operator motivated by the **physics of short-range central and tensor correlations** 

#### **Central Correlator** C<sub>r</sub>

 radial distance-dependent shift in the relative coordinate of a nucleon pair

$$\begin{split} \mathbf{g}_r &= \frac{1}{2} \big[ s(\mathbf{r}) \; \mathbf{q}_r + \mathbf{q}_r \; s(\mathbf{r}) \big] \\ \mathbf{q}_r &= \frac{1}{2} \big[ \frac{\vec{\mathbf{r}}}{\mathbf{r}} \cdot \vec{\mathbf{q}} + \vec{\mathbf{q}} \cdot \frac{\vec{\mathbf{r}}}{\mathbf{r}} \big] \end{split}$$

#### **Tensor Correlator** $C_{\Omega}$

 angular shift depending on the orientation of spin and relative coordinate of a nucleon pair

$$egin{aligned} &\mathbf{g}_\Omega = rac{3}{2} artheta(\mathbf{r}) ig[ (ec{\sigma}_1 \!\cdot ec{\mathbf{q}}_\Omega) (ec{\sigma}_2 \!\cdot ec{\mathbf{r}}) + (ec{\mathbf{r}} \!\leftrightarrow \!ec{\mathbf{q}}_\Omega) ig] \ & ec{\mathbf{q}}_\Omega = ec{\mathbf{q}} - rac{ec{\mathbf{r}}}{ec{\mathbf{r}}} \,\, \mathbf{q}_r \end{aligned}$$

• s(r) and  $\vartheta(r)$  for given potential determined by energy minimization in the two-body system (for each S, T)

### Correlated States: The Deuteron



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# Correlated Interaction: $V_{\rm UCOM}$



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Modern Effective Interactions II Similarity Renormalization Group (SRG)

# Similarity Renormalization Group

unitary transformation of the Hamiltonian to a band-diagonal form with respect to a given uncorrelated many-body basis

**Flow Equation for Hamiltonian** 

evolution equation for Hamiltonian

$$\widetilde{\mathrm{H}}(lpha) = \mathrm{C}^\dagger(lpha) \, \mathrm{H}\, \mathrm{C}(lpha) \hspace{0.5cm} 
ightarrow \hspace{0.5cm} rac{\mathrm{d}}{\mathrm{d} lpha} \widetilde{\mathrm{H}}(lpha) = ig[\eta(lpha), \widetilde{\mathrm{H}}(lpha)ig]$$

 dynamical generator defined as commutator with the operator in whose eigenbasis H shall be diagonalized

$$\eta(\alpha) = \left[\mathrm{T}_{\mathrm{int}}, \widetilde{\mathrm{H}}(\alpha)\right] \stackrel{\mathrm{2B}}{=} rac{1}{2\mu} \left[ \vec{\mathrm{q}}^2, \widetilde{\mathrm{H}}(\alpha) 
ight]$$

[Bogner et al., PRC75 061001(R) (2007); Hergert & Roth, PRC75 051001(R) (2007)]

### The SRG Generator: A Closer Look

■ typical *NN* interaction operators:

$$\mathcal{O}_p \in \{\mathbb{1}, \vec{\sigma}_1 \cdot \vec{\sigma}_2, \vec{\mathfrak{l}}^2, \vec{\mathfrak{l}} \cdot \vec{\mathfrak{s}}, \mathfrak{s}_{12}(\vec{\mathfrak{r}}, \vec{\mathfrak{r}}), \ldots\} \otimes \{\mathbb{1}, \vec{\tau}_1 \cdot \vec{\tau}_2, \ldots\}$$

#### **Radial Kinetic Energy**

$$\eta_r(0) \sim \left[\mathbf{q}_r^2, \mathbf{V}\right] = \sum_p \left[\mathbf{q}_r^2, v_p(\mathbf{r})\mathbf{O}_p\right] = \sum_p \left(\mathbf{q}_r v_p'(\mathbf{r})\mathbf{O}_p + \mathbf{O}_p v_p'(\mathbf{r})\mathbf{q}_r\right)$$

#### **Angular Kinetic Energy**

$$\eta_{\Omega}(0) \sim \left[\vec{\mathbf{l}}^2, \mathbf{V}\right] = \left[\vec{\mathbf{l}}^2, v_t(\mathbf{r})\mathbf{s}_{12}(\vec{\mathbf{r}}, \vec{\mathbf{r}})\right] = -4iv_t(\mathbf{r})\mathbf{s}_{12}(\vec{\mathbf{r}}, \vec{\mathbf{q}}_{\Omega})$$

 $\Im \eta(0)$  has the same structure as the UCOM generators  $\mathrm{g}_r$  and  $\mathrm{g}_\Omega$ 



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# Few-Body Systems

# <sup>4</sup>He: Convergence



NCSM code by P. Navrátil [PRC 61, 044001 (2000)]

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NCSM code by P. Navrátil [PRC 61, 044001 (2000)]

## Three-Body Interactions — Strategies

**Correlated Hamiltonian in Many-Body Space** 

$$\begin{split} \widetilde{H} &= C^{\dagger} (T + V_{NN} + V_{3N}) C \\ &= \widetilde{T}^{[1]} + (\widetilde{T}^{[2]} + \widetilde{V}^{[2]}_{NN}) + (\widetilde{T}^{[3]} + \widetilde{V}^{[3]}_{NN} + \widetilde{V}^{[3]}_{3N}) + \cdots \\ &= T + V_{UCOM} + V^{[3]}_{UCOM} + \cdots \end{split}$$

strategies for treating the three-body contributions:

include full V<sup>[3]</sup><sub>UCOM</sub> consisting of genuine and induced 3N terms
 replace V<sup>[3]</sup><sub>UCOM</sub> by "phenomenological" three-body force
 minimize V<sup>[3]</sup><sub>UCOM</sub> by proper choice of unitary transformation



Tjon-line: E(<sup>4</sup>He) vs. E(<sup>3</sup>H) for phase-shift equivalent NNinteractions



Tjon-line: E(<sup>4</sup>He) vs. E(<sup>3</sup>H) for phase-shift equivalent NNinteractions



- Tjon-line: E(<sup>4</sup>He) vs. E(<sup>3</sup>H) for phase-shift equivalent NNinteractions
- $\blacksquare$  use lpha / range of  $\mathrm{C}_\Omega$  to
  - $\bullet$  test dependence of  $\mathbf{V}_{\alpha}$  or  $V_{UCOM}$
  - tune contributions of net 3N force



# Many-Body Systems

# Hartree-Fock with V<sub>UCOM</sub>



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# Hartree-Fock with V<sub>UCOM</sub>



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### Hartree-Fock with SRG Potentials



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# Hartree-Fock with V<sub>UCOM</sub>



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# Perturbation Theory with V<sub>UCOM</sub>



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# RPA Ring Summation with $V_{UCOM}$



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### 3N Forces: Energies & Radii



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# 3N Forces: HF Single-Particle Energies



# 3N Forces: HF Single-Particle Energies



# Beyond Hartree-Fock

perturbation theory, Brueckner-HF , RPA ... Hartree-Fock (Mean-Field)

Long-Range Correlations Hartree-Fock-Bogoliubov

Pairing Correlations

# HFB Theory Overview

### **Bogoliubov Transformation**

$$egin{aligned} eta_k^\dagger &= \sum_q U_{qk} \mathrm{c}_q^\dagger + V_{qk} \mathrm{c}_q \ eta_k &= \sum_q U_{qk}^* \mathrm{c}_q + V_{qk}^* \mathrm{c}_q^\dagger \end{aligned}$$

where

$$egin{aligned} &\{eta_k,eta_{k'}\} \stackrel{!}{=} \{eta_k^\dagger,eta_{k'}^\dagger\} \stackrel{!}{=} 0 \ &\{eta_k,eta_{k'}^\dagger\} \stackrel{!}{=} \delta_{kk'} \end{aligned}$$

#### **HFB Densities & Fields**

$$egin{aligned} &
ho_{kk'} \equiv ig\langle \Psi ig| \, \mathbf{c}_{k'}^{\dagger} \mathbf{c}_{k} \, ig| \Psi ig
angle = (V^{*}V^{T})_{kk'} \ &\kappa_{kk'} \equiv ig\langle \Psi ig| \, \mathbf{c}_{k'} \mathbf{c}_{k} \, ig| \Psi ig
angle = (V^{*}U^{T})_{kk'} \ &\Gamma_{kk'} = \sum_{qq'} \left( rac{2}{A} ar{\mathbf{t}}_{\mathrm{rel}} + ar{\mathbf{v}} 
ight)_{kq',k'q} 
ho_{qq'} \ &\Delta_{kk'} = \sum_{qq'} \left( rac{2}{A} ar{\mathbf{t}}_{\mathrm{rel}} + ar{\mathbf{v}} 
ight)_{kk',qq'} \kappa_{qq'} \end{aligned}$$

Energy

$$E[
ho,\kappa,\kappa^*] = rac{ig\langle\Psiigert \, {f H} igert \Psiig
angle}{ig\langle\Psiigert \Psiig
angle} \equiv rac{1}{2} \left({
m tr} \ \Gamma
ho - {
m tr} \ \Delta\kappa^*
ight)$$

### **HFB Equations**

$$\left(\mathcal{H}-\lambda\mathcal{N}
ight)egin{pmatrix} U\V\end{pmatrix}\equiv egin{pmatrix} \Gamma-\lambda&\Delta\ -\Delta^*&-\Gamma^*+\lambda\end{pmatrix}egin{pmatrix} U\V\end{pmatrix}=Eegin{pmatrix} U\V\end{pmatrix}$$

# Particle Number Projection

#### **Variation of Projected Energy**

$$\delta E(N_0) = rac{1}{2\pi ig\langle \mathbf{P}_{N_0} ig
angle} \int_0^{2\pi} d\phi ig\langle e^{i\phi(\mathbf{N}-N_0)} ig
angle \left\{ \delta ig\langle \mathbf{H} ig
angle_{\phi} - \left( E(N_0) - ig\langle \mathbf{H} ig
angle_{\phi} 
ight) \delta \log ig\langle e^{i\phi\mathbf{N}} ig
angle 
ight\}$$

$$\left< \mathrm{H} \right>_{\phi} \equiv \left< \mathrm{H} e^{i \phi \mathrm{N}} \right> / \left< e^{i \phi \mathrm{N}} \right>$$

### Lipkin-Nogami + PAV

- power series expansion
- expansion coefficients not varied
- indeterminate / numerically unstable at shell closures
- exact PNP after variation

#### VAP

- higher (but managable) computational effort
- implement with care: subtle cancellations between divergences of direct, exchange, and pairing terms

#### Structure of **HFB equations is preserved** by both methods!

Flocard & Onishi, Ann. Phys. 254, 275 (1997, approx. PNP); Sheikh et al., Phys. Rev. C66, 044318 (2002, exact PNP)

# Sn Isotopes: Binding & Pairing Energies



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# Sn Isotopes: Binding & Pairing Energies



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# Density-Dependent Force: Pairing



• linear density dependence,  $t_0 = \frac{1}{6}C_{3N}$ :

$${
m V}_
ho = t_0 \left( {1 + {
m P}_\sigma } 
ight) 
ho \left( {rac{1}{2} ({ec {
m r}_1} + {ec {
m r}_2} )} 
ight) \delta^3 \left( {ec {
m r}_1} - {ec {
m r}_2} 
ight)$$

- mixed density for PAV/VAP:  $\rho(\vec{R}) \longrightarrow \rho_{\phi_p,\phi_n}(\vec{R})$
- phenomenological VAP calculations:  $E_{pair} \simeq 10 20 \text{ MeV}$ (Stoitsov et al., nucl-th/0610061; Anguiano et al., Phys. Lett. B545 (2002), 62)

"correct" order of magnitude with realistic NN int.

### Density-Dependent Force: Pairing



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# Conclusions

# Modern Effective Interactions

### Status

- treatment of **short-range central** and **tensor correlations** by unitary transformations:
  - Unitary Correlation Operator Method
  - Similarity Renormalization Group
- universal phase-shift equivalent correlated interaction  $V_{UCOM}$

### Outlook

- connections between UCOM and SRG
- inclusion & treatment of 3N Forces, in particular...
- chiral interactions

# HF, HFB, and Beyond



# HF, HFB, and Beyond

### Status

- fully consistent HFB calculations with particle number projection
- inclusion of 3*N*-forces: contact & finite range matrix elements for HF, density-dependent force for HFB, RPA, ...
- Like-particle- & *pn*-QRPA (benchmarked)

### Outlook

- *pn* **pairing**, Isobaric Analog & Gamow-Teller Resonances
- **deformation** and **symmetry restoration** by projection (isospin, parity, angular momentum)
- caveat: analytic structure of density-dependent forces is very important in symmetry-projected HFB (J. Dobaczewski, arXiv: 0708.0441)

# Perspectives

### Modern Effective Interactions

- treatment of short-range central and tensor correlations by unitary transformations: UCOM, SRG, Lee-Suzuki,...
- $\bullet$  phase-shift equivalent correlated interaction  $V_{\text{UCOM}}$
- universal input for...

### Innovative Many-Body Methods

- No-Core Shell Model,...
- Importance Truncated NCSM, Coupled Cluster Method,...
- Hartree-Fock plus MBPT, Padé Resummed MBPT, BHF, HFB, RPA,...
- Fermionic Molecular Dynamics,...

### Last Words...

### **My Collaborators**

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- http://crunch.ikp.physik.tu-darmstadt.de/tnp/

# Appendix

# <sup>10</sup>B: Hallmark of a 3N Interaction?



# <sup>10</sup>B: Hallmark of a 3N Interaction?



### RPA + Realistic Interactions



- fully self-consistent RPA based on the Hartree-Fock orbits using the same V<sub>UCOM</sub>
- recovering sum rules with high precision, spurious center-ofmass mode fully decoupled at  $\sim 10 \text{ keV}$
- Extended-RPA and Second-RPA to include effects of ground state correlations and complex configurations

# RPA with $V_{UCOM}$



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# SRPA: Complex Configurations



# SRPA: Complex Configurations



# Outlook: RPA with Three-Body Forces



- Iong-range tensor correlator & repulsive three-body contact interaction
- systematic improvement of
  - → rms-radii
  - → single-particle spectra
  - → strength distributions

- **—** standard  $V_{\text{UCOM}}$ 
  - V<sub>UCOM</sub> with long-range tensor & three-body contact interaction