

Ultracold Atomic Gases in 1D Optical Lattices

DMRG Method in Inhomogeneous Lattice Topologies

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Summary

- cold atomic gases on optical lattices are a perfect experimental realization of the Hubbard Modell [1,2]
- all relevant parameters can be precisely controlled in the experiment allowing for unique comparison with theory throughout the whole phase-diagram [8]
- the Hubbard Hamiltonian facilitates a straight forward growing of the system, which makes this an ideal case for DMRG calculations [6,7]
- nevertheless DMRG involves a truncation procedure, which should be tested in all different regimes of the phase-diagram
- investigate how DMRG performs in spatially inhomogeneous lattices
- append lattice sites with different on-site energies, e.g. apply a superlattice, to check for the effect of localization or disorder on the spectrum of the reduced density-matrix and thus on the target state

Bose-Hubbard Model

- 1D optical lattice with I lattice sites and N bosonic particles
- restriction to the first energy-band, $T = 0$, nearest neighbor tunneling, and an on-site two-particle contact interaction
- additional on-site potential can map arbitrary lattice topologies

$$\hat{H} = -J \sum_{i=1}^I (\hat{a}_{i+1}^\dagger \hat{a}_i + \hat{a}_i^\dagger \hat{a}_{i+1}) \quad \text{tunneling}$$

$$+ \frac{V}{2} \sum_{i=1}^I \hat{n}_i (\hat{n}_i - 1) \quad \text{interaction}$$

$$+ \Delta \sum_{i=1}^I \epsilon_i \hat{n}_i \quad \text{superlattice potential}$$

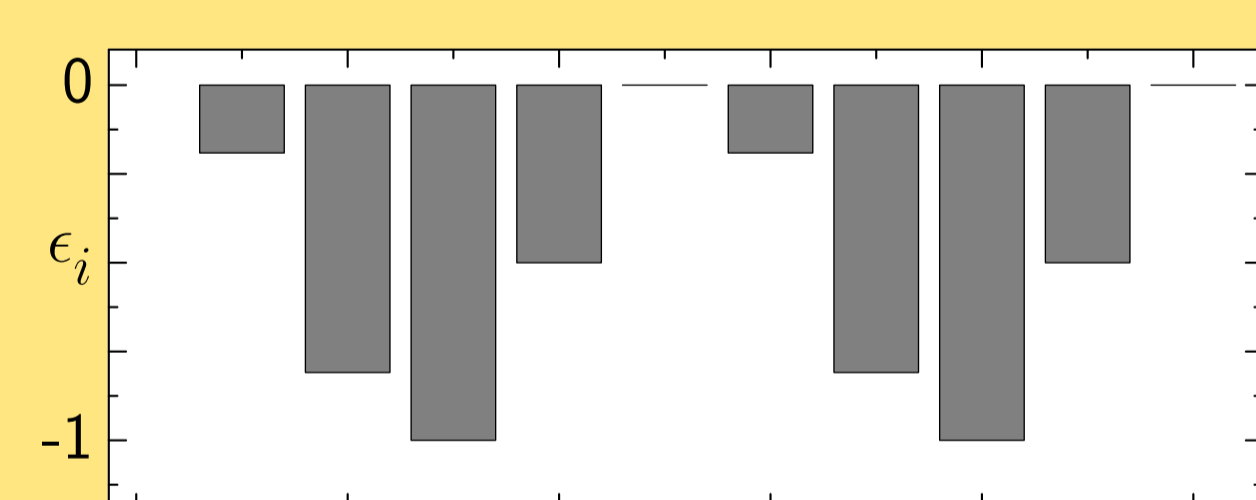
$\hat{a}_i^\dagger, \hat{a}_i, \hat{n}_i$ creation, annihilation, occupation-number operators
 J tunneling matrix element
 V two particle interaction energy
 Δ strength of the superlattice potential
 ϵ_i topology of the superlattice potential

- use an occupation-number representation spanning the Fock-space \mathcal{H} to formulate a matrix representation of the Hamiltonian

$$\mathcal{H} = \text{span} \{ | \{n_1, \dots, n_I\}_\alpha \rangle \}$$

Two-Color Superlattice

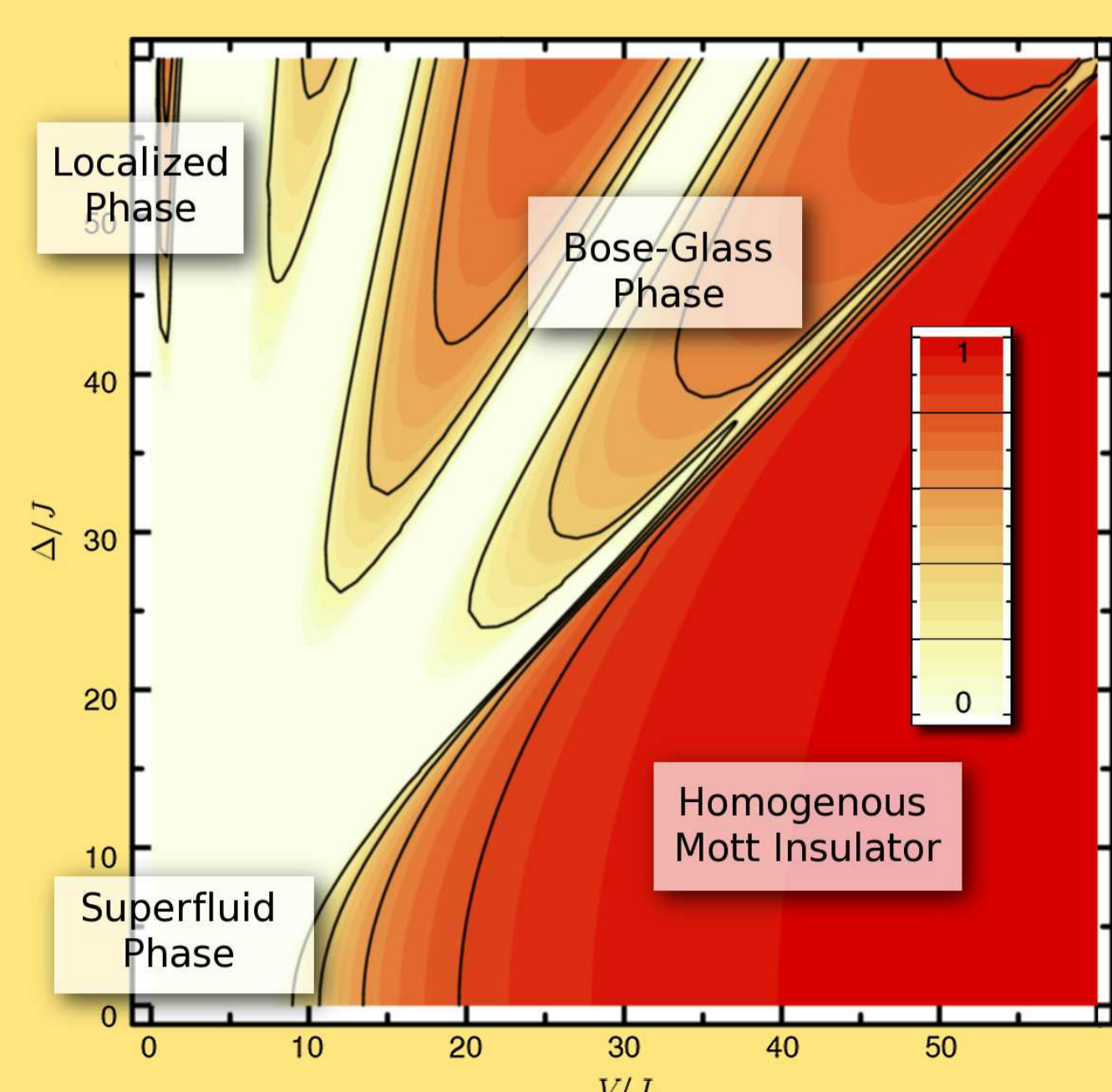
- superposition of two laser beams with slightly different wavelengths [5] form a superlattice potential
- Δ/J is the on-site energy of the deepest superlattice well



Phase Diagram

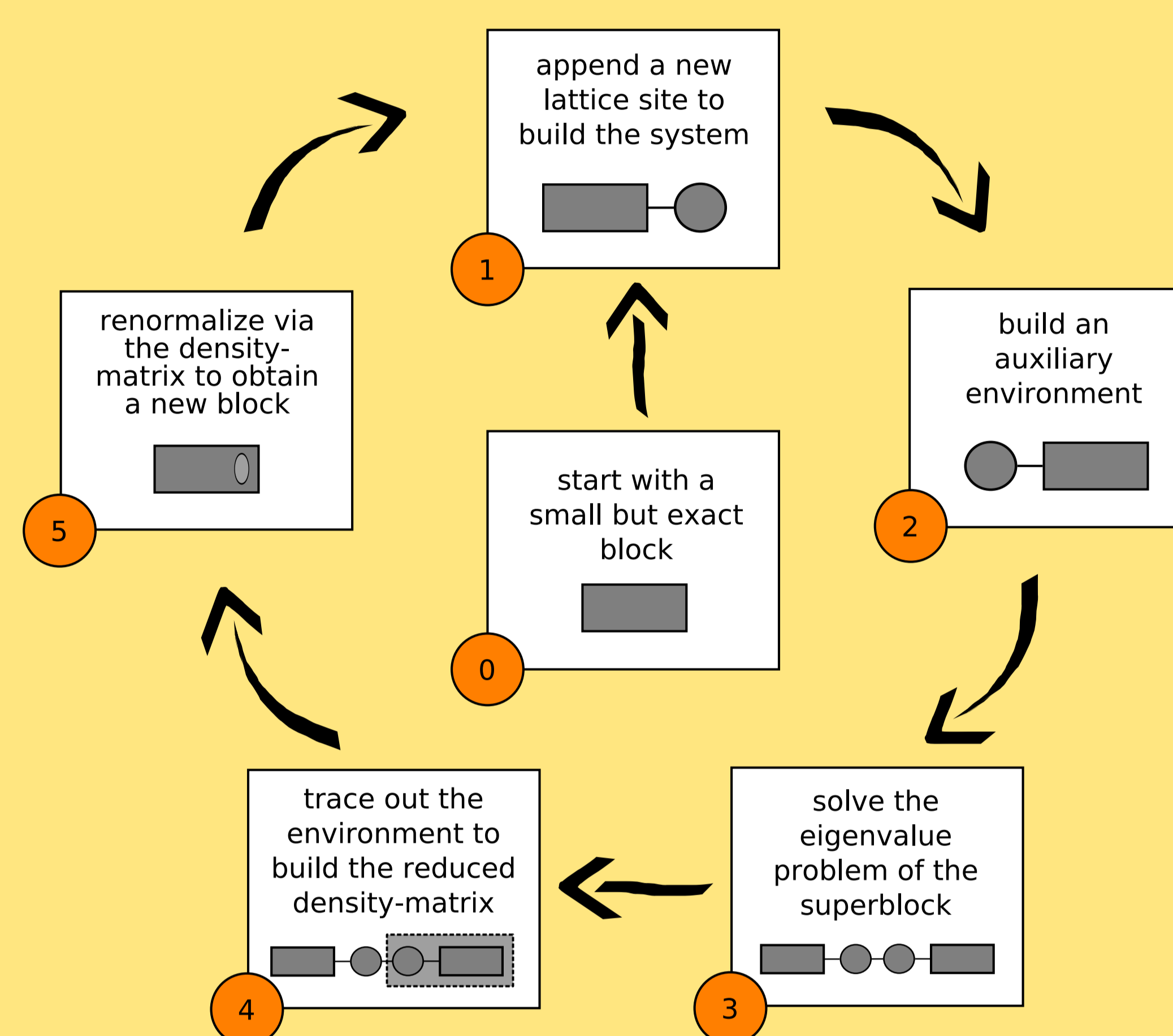
- exact diagonalization for a moderate system size ($I = N = 10$) with periodic boundary conditions yields ground states

$$| \psi^{(0)} \rangle = \sum_{\alpha=1}^D C_\alpha^{(0)} | \{n_1, \dots, n_I\}_\alpha \rangle \quad \hat{H} | \psi^{(0)} \rangle = E_0 | \psi^{(0)} \rangle$$



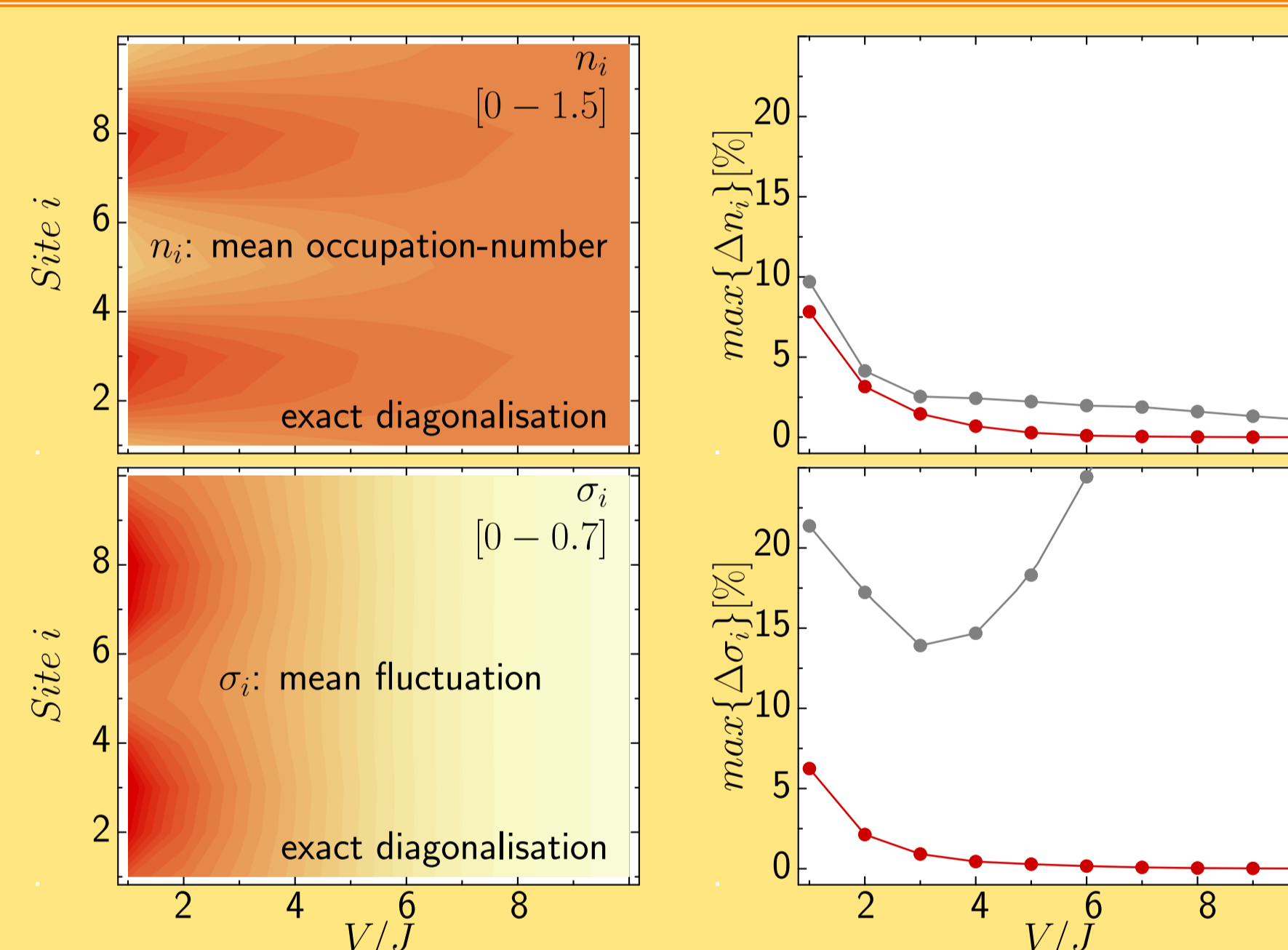
- $\max\{|C_\alpha^{(0)}|^2\}$ reflects the composition of the ground state
- commensurate systems exhibit a rich phase-diagram [3,4]
- subtle interplay between the different energy scales (J, V and Δ)
- test DMRG calculations in the transition regimes of the phase diagram

DMRG Algorithm



1. $\mathcal{H}_{system} = \mathcal{H}_{block} \otimes \mathcal{H}_{site}$
2. $\mathcal{H}_{env} = \tilde{\mathcal{H}}_{Block} \otimes \tilde{\mathcal{H}}_{Site}$
3. $\mathcal{H}_{super} = \mathcal{H}_{system} \otimes \mathcal{H}_{env}$
 $\hat{H}_{super} | \psi^{(0)} \rangle = E_0 | \psi^{(0)} \rangle$
4. $\hat{\rho}_{red} = \text{Tr}_{env} (| \psi^{(0)} \rangle \langle \psi^{(0)} |)$
5. $\hat{\rho}_{red} | \omega_i \rangle = \omega_i | \omega_i \rangle$
 $\sum_i \omega_i = 1 \quad \omega_i \geq \omega_{i+1}$
 $\mathcal{O} = (\vec{w}_1, \dots, \vec{w}_{D_b})$
 $H_{block(L)} = \mathcal{O}^T H_{system(L)} \mathcal{O}$

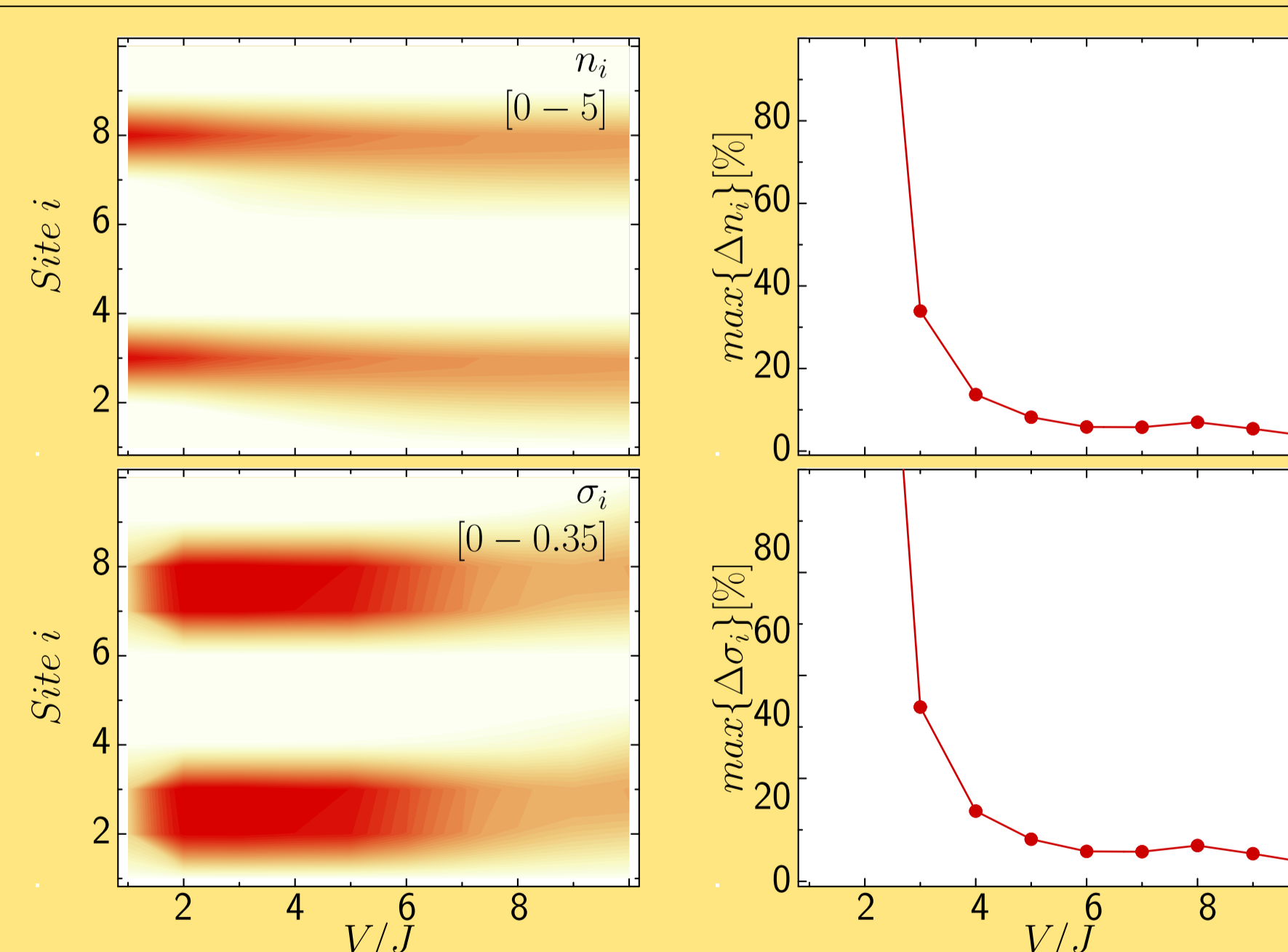
Results



- 2D plots show the maximum relative error of the DMRG calculations
- gray: DMRG without sweep
- red: DMRG with one sweep

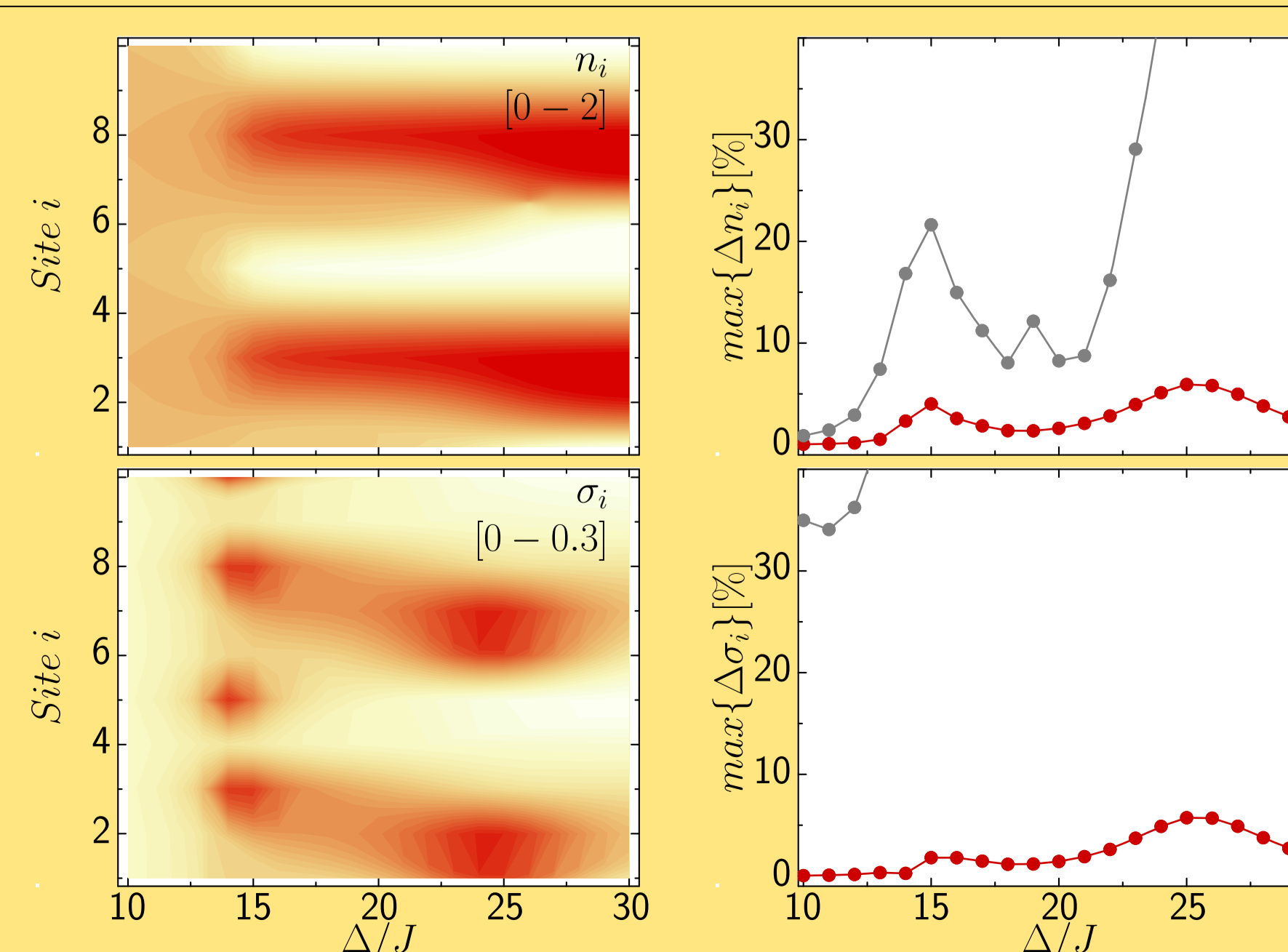
Superfluid to Mott-Insulator ($\Delta/J = 1$)

- lattice sites are almost equal, ground state is nearly homogenous
- good results even without sweeping



Localized to Bose-Glass ($\Delta/J = 50$)

- ground state is strongly inhomogeneous
- DMRG needs at least the average number of particles per unit-cell to reproduce the localized phase
- sweeping leads to good results as soon as delocalization of particles occurs



Mott-Insulator to Bose-Glass ($V/J = 15$)

- intermediate inhomogeneity
- redistribution of particles is well reproduced by DMRG calculations with sweeps
- exact calculations involve a $D \approx 10^5$ eigenvalue problem
- DMRG calculations involve a $D \approx 400$ eigenvalue problem