

# Chiral 3N Forces in Ab-Initio Nuclear Spectroscopy & Reactions

Joachim Langhammer



TECHNISCHE  
UNIVERSITÄT  
DARMSTADT

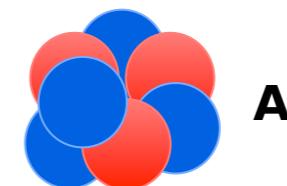
# Outline

- Motivation
- Ingredients from Three-Body Technology
- Nuclear Structure within the IT-NCSM
  - Sensitivity of Spectra of  $p$ -Shell Nuclei on LECs & Cutoff
  - The Oxygen Isotopic Chain
- Nuclear Reactions within the NCSM/RGM approach
  - Formalism Including 3N forces
  - $n-^4\text{He}$  Scattering
  - Scattering Involving Heavier Nuclei
- Conclusions

# What we are aiming for...

**Realistic ab-initio description of light nuclei**

Bound states  
& spectroscopy



**Using NN+3N forces that are solidly rooted in QCD**

**Importance-Truncated NCSM**

Ab-initio description  
of nuclear clusters

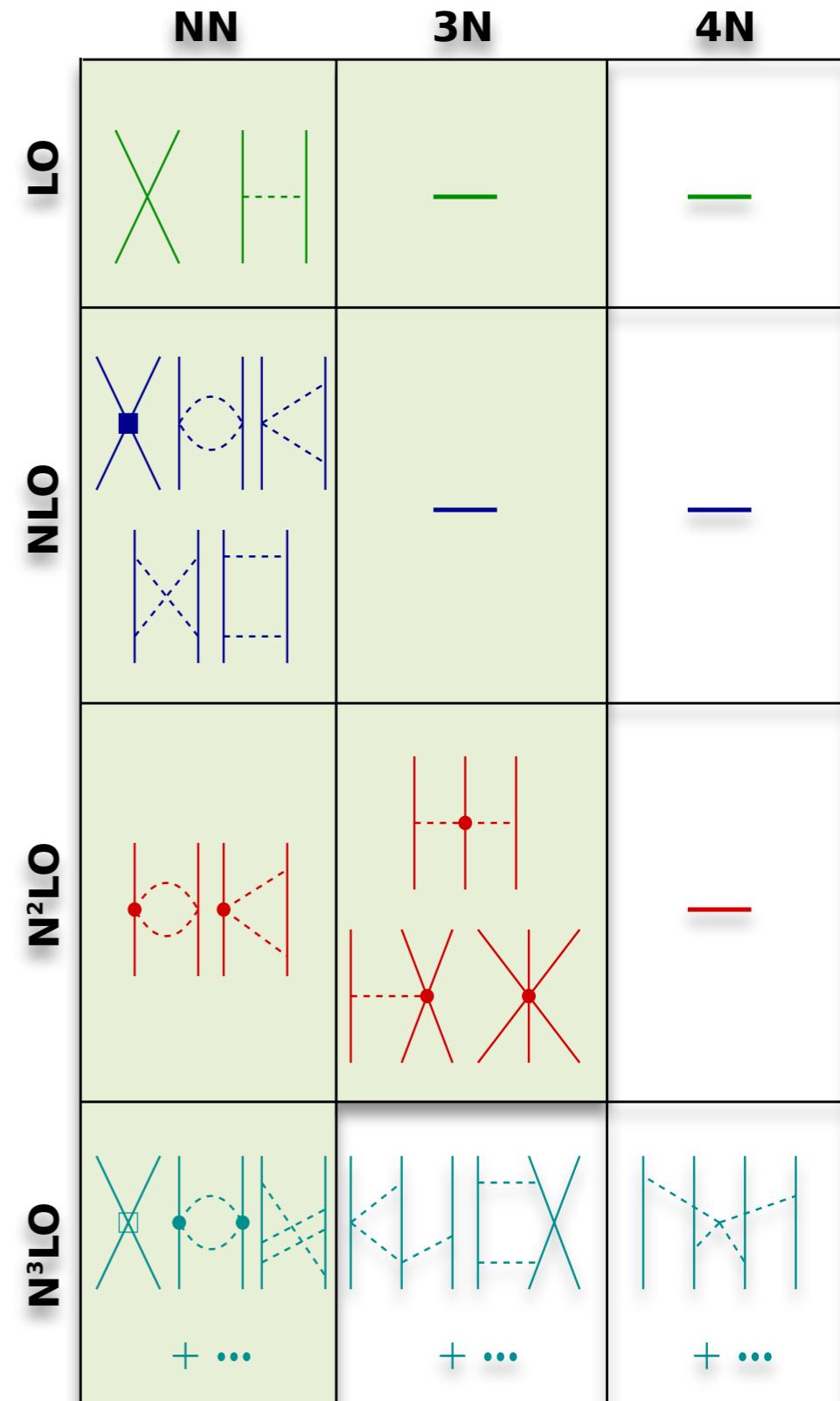
All relevant observables  
accessible from the eigenstates

# Ingredients from Three-Body Technology

# The Chiral NN+3N Hamiltonian

Weinberg, van Kolck, Machleidt, Entem, Meißner, Epelbaum, Krebs, Bernard, Skibinski, Golak...

- Hierarchy of consistent nuclear NN, 3N,... forces (and currents)
- NN interaction @ N<sup>3</sup>LO ( $\Lambda=500\text{MeV}$ )  
[Entem, Machleidt, Phys.Rev C **68**, 041001(R) (2003)]
- Standard Hamiltonian
  - 3N interaction @ N<sup>2</sup>LO( $\Lambda=500\text{MeV}$ )
    - Local form by Navrátil
    - LECs  $c_D, c_E$  fitted to  $\beta$ -decay halflife & binding energy of <sup>3</sup>H  
[Gazit et.al., Phys.Rev.Lett. **103**, 102502 (2009)]
- Reduced-Cutoff Hamiltonian
  - 3N interaction @ N<sup>2</sup>LO( $\Lambda=400\text{MeV}$ )
    - $c_D=-0.2, c_E$  fitted to <sup>4</sup>He



# The Similarity Renormalization Group

Wegner, Glazek, Wilson, Perry, Bogner, Furnstahl, Hergert, Calci, Langhammer, Roth, Jurgenson, Navrátil,...

...yields an evolved  
Hamiltonian with **improved convergence properties** in  
many-body calculations

- Unitary transformation of Hamiltonian  $H_\alpha = U_\alpha^\dagger H U_\alpha$

## Different SRG-Evolved Hamiltonians

- **NN-only**: start with NN initial Hamiltonian and keep two-body terms only
- **NN+3N-induced**: start with NN initial Hamiltonian and keep two- and three-body terms
- **NN+3N-full**: start with NN+3N initial Hamiltonian and keep two- and three-body terms

# Nuclear Structure within the IT-NCSM

# Sensitivity on Chiral 3N Forces

- LECs from  $\pi n$  vertices have sizable variations

$$c_1 = (-1.23 \dots - 0.76) \text{ GeV}^{-1}$$

$$c_3 = (-5.94 \dots - 3.2) \text{ GeV}^{-1}$$

$$c_4 = (3.4 \dots 5.4) \text{ GeV}^{-1}$$

- Uncertainties may propagate into nuclear structure observables
  - ⇒ **quantify theoretical uncertainties**
  - ⇒ **provide constraints for next generation chiral forces**

- Some N<sup>3</sup>LO contributions can be absorbed into the N<sup>2</sup>LO diagrams by shifting the  $c_i$  constants:

$$\tilde{c}_1 = c_1 - \frac{g_A^2 M_\pi}{64\pi F_\pi^2}$$

$$\tilde{c}_3 = c_3 - \frac{g_A^4 M_\pi}{16\pi F_\pi^2}$$

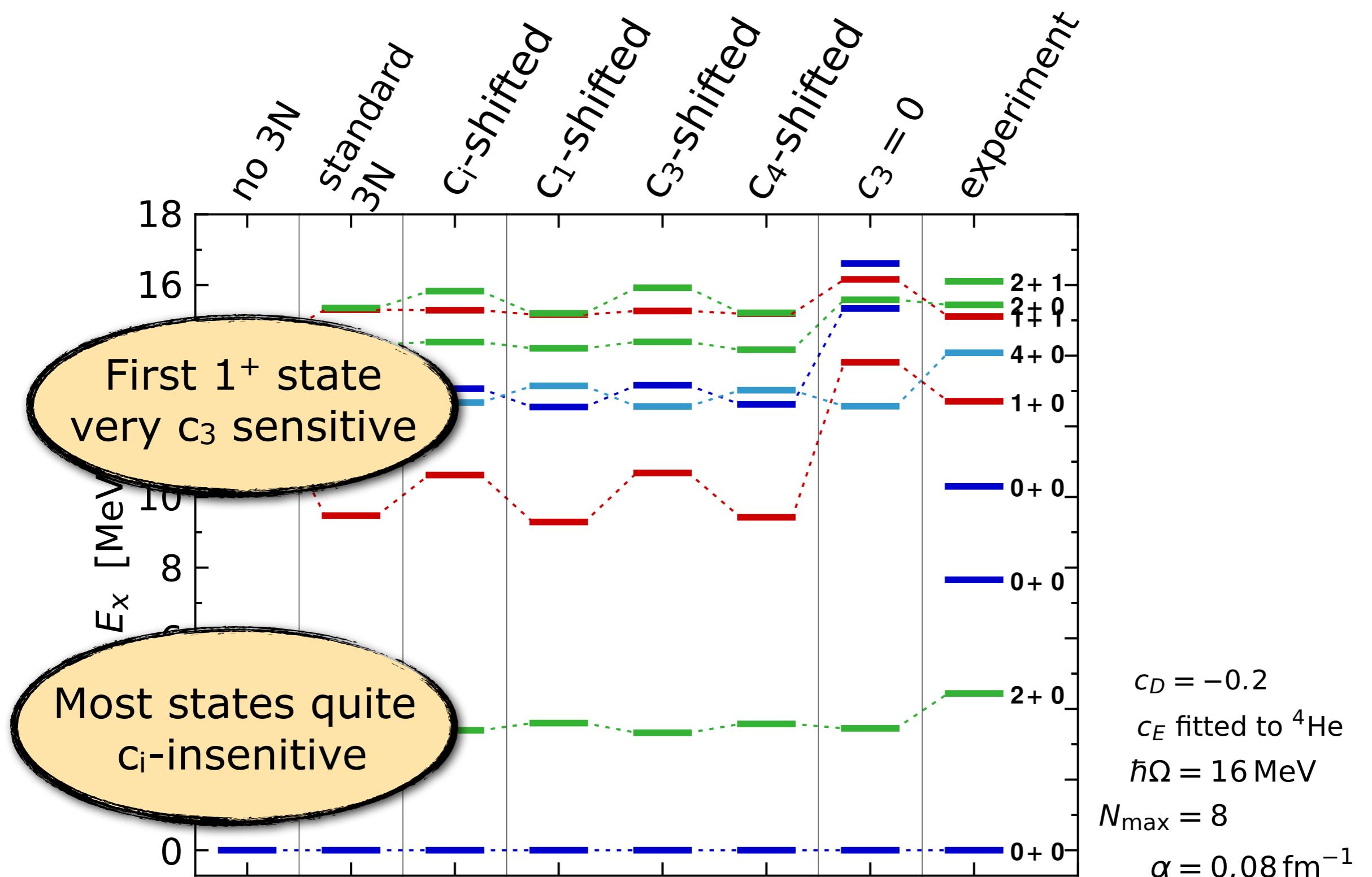
$$\tilde{c}_4 = c_4 - \frac{g_A^4 M_\pi}{16\pi F_\pi^2}$$

*10-20% change*

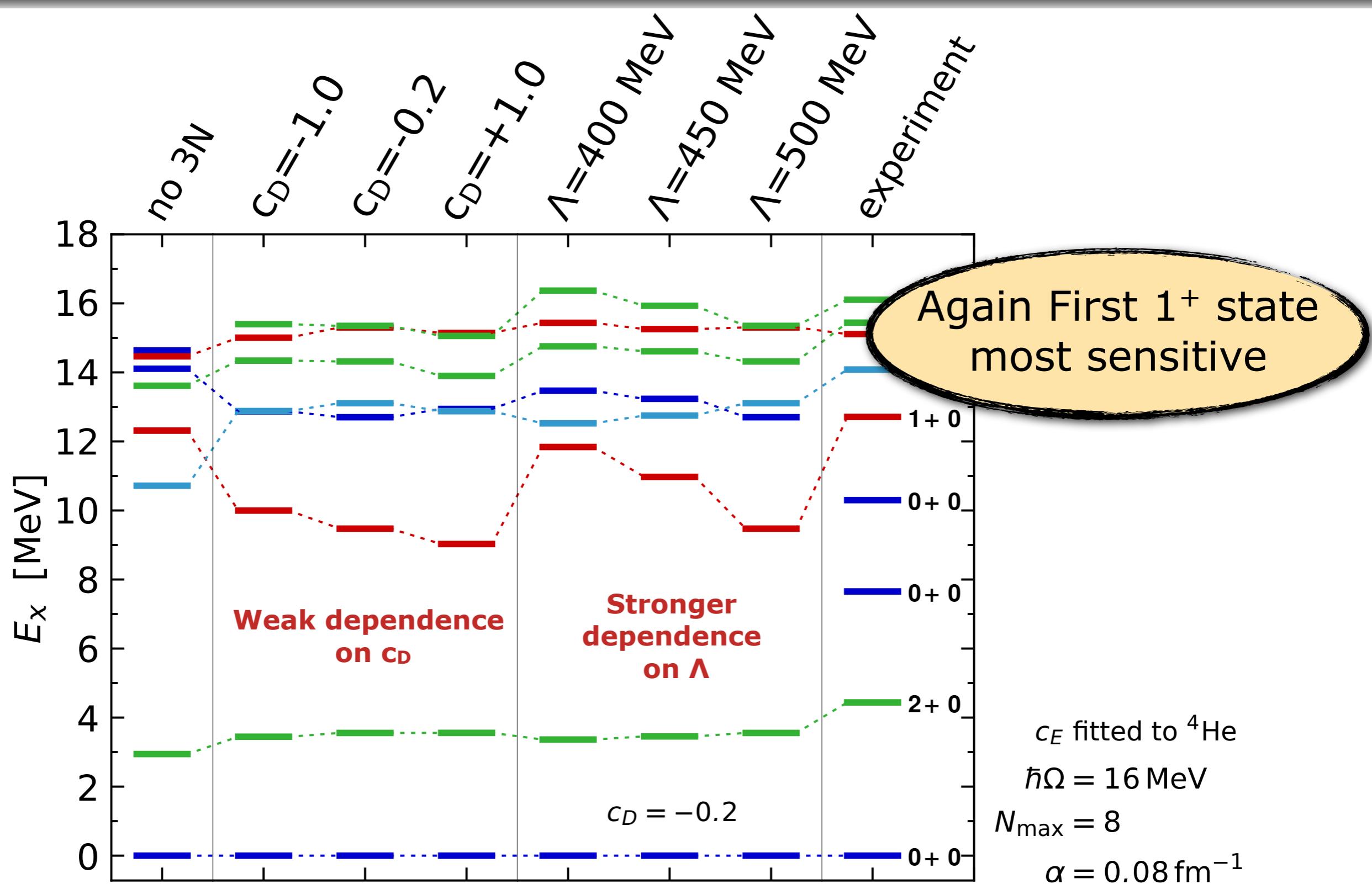
- Does the chosen cutoff  $\Lambda$  affect nuclear structure observables?

Analyze **sensitivity** of spectra on LECs and  
**chiral cutoff** of the chiral 3N force at N<sup>2</sup>LO

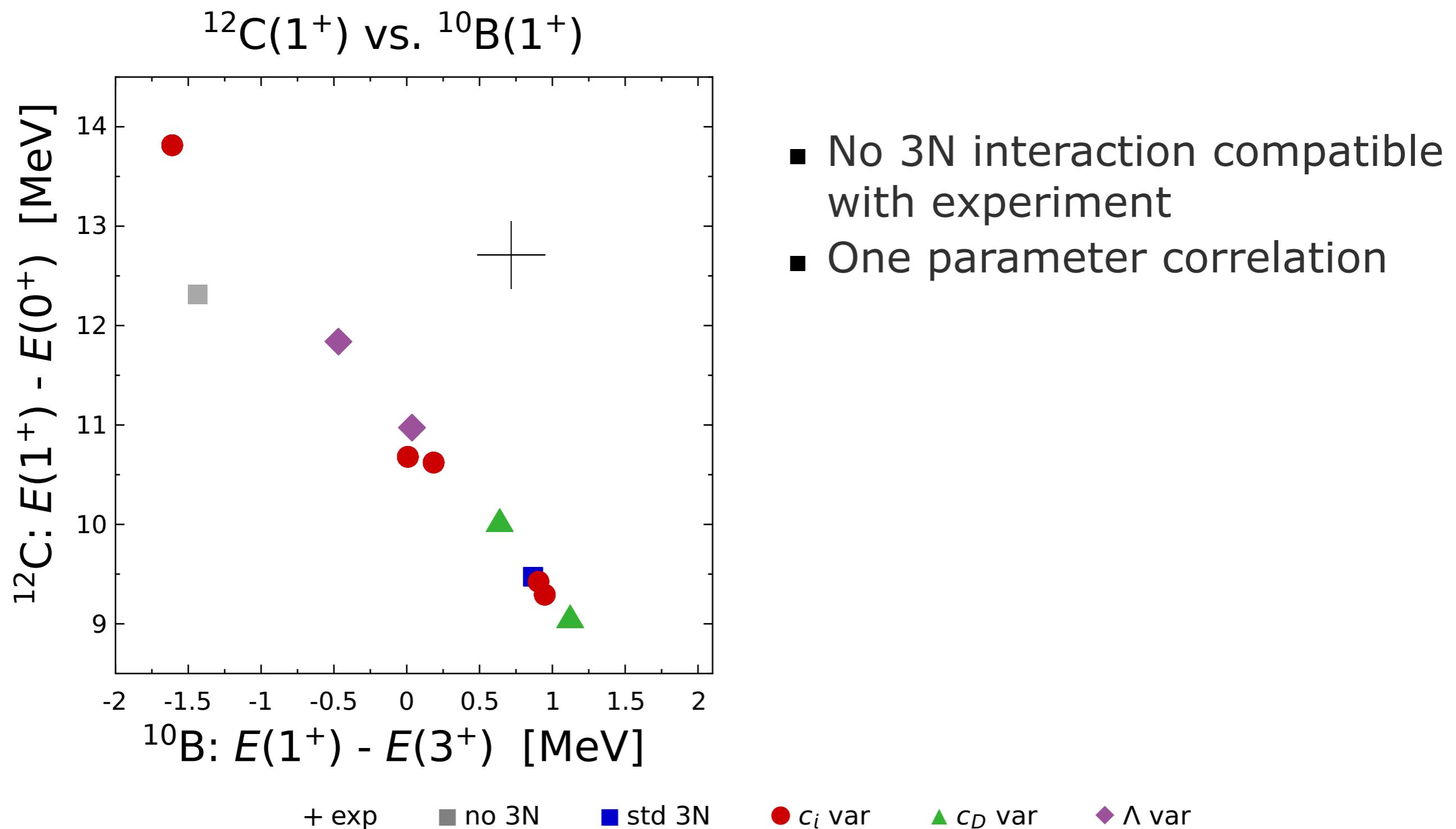
# $^{12}\text{C}$ : Sensitivity on $c_i$



# $^{12}\text{C}$ : Sensitivity on $c_D$ and Cutoff

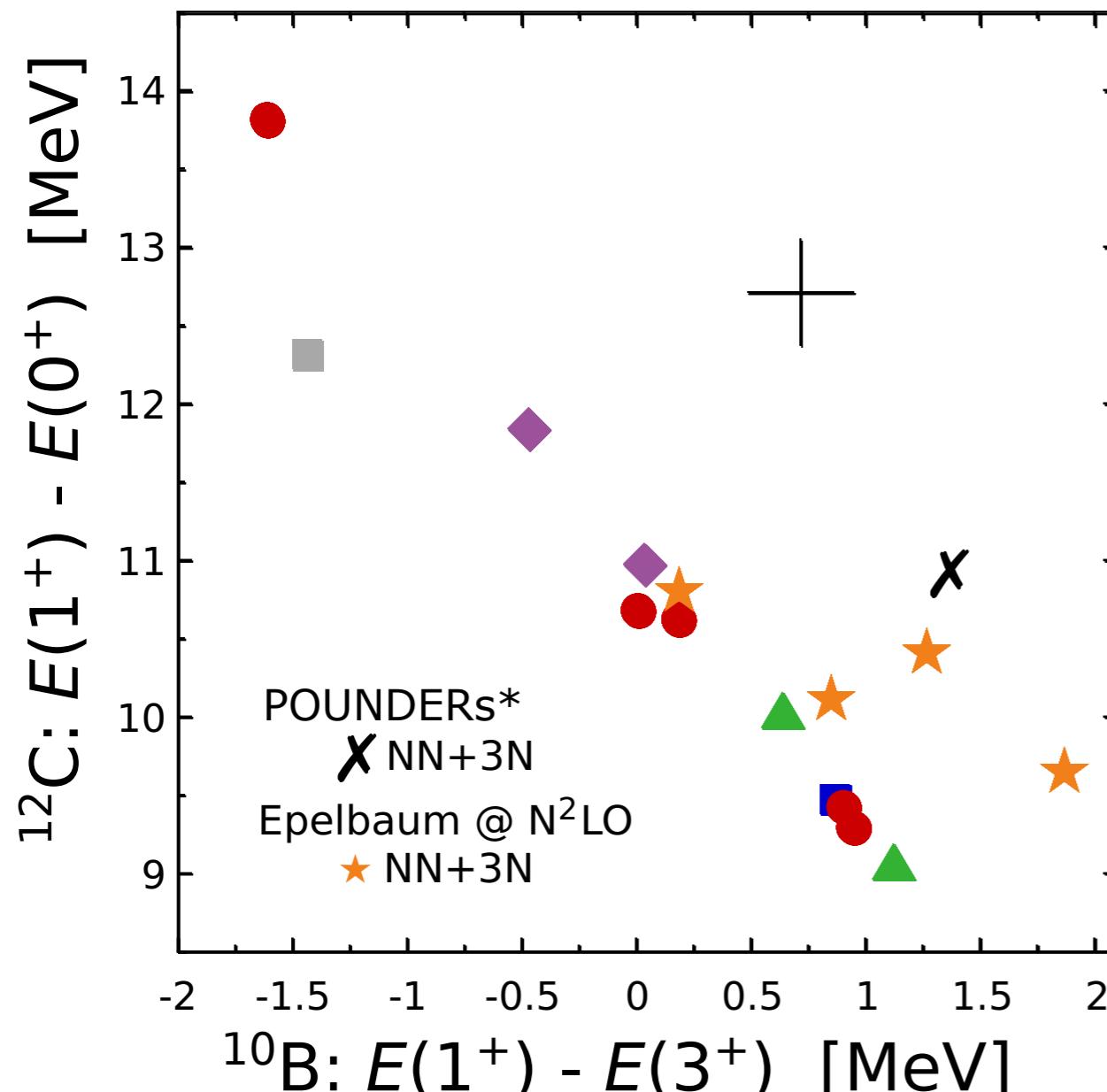


# Correlation Analysis



# Correlation Analysis

$^{12}\text{C}(1^+) \text{ vs. } ^{10}\text{B}(1^+)$



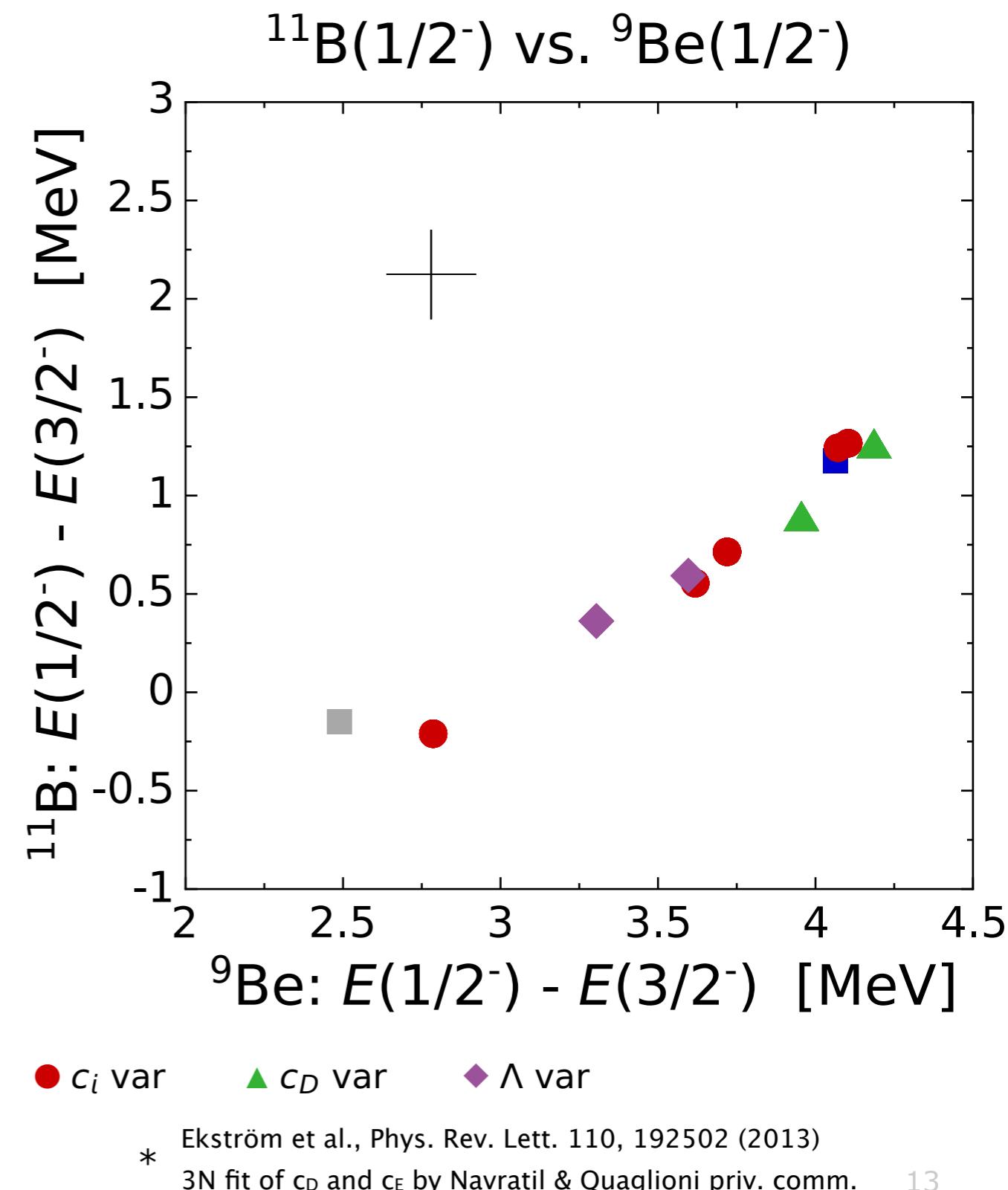
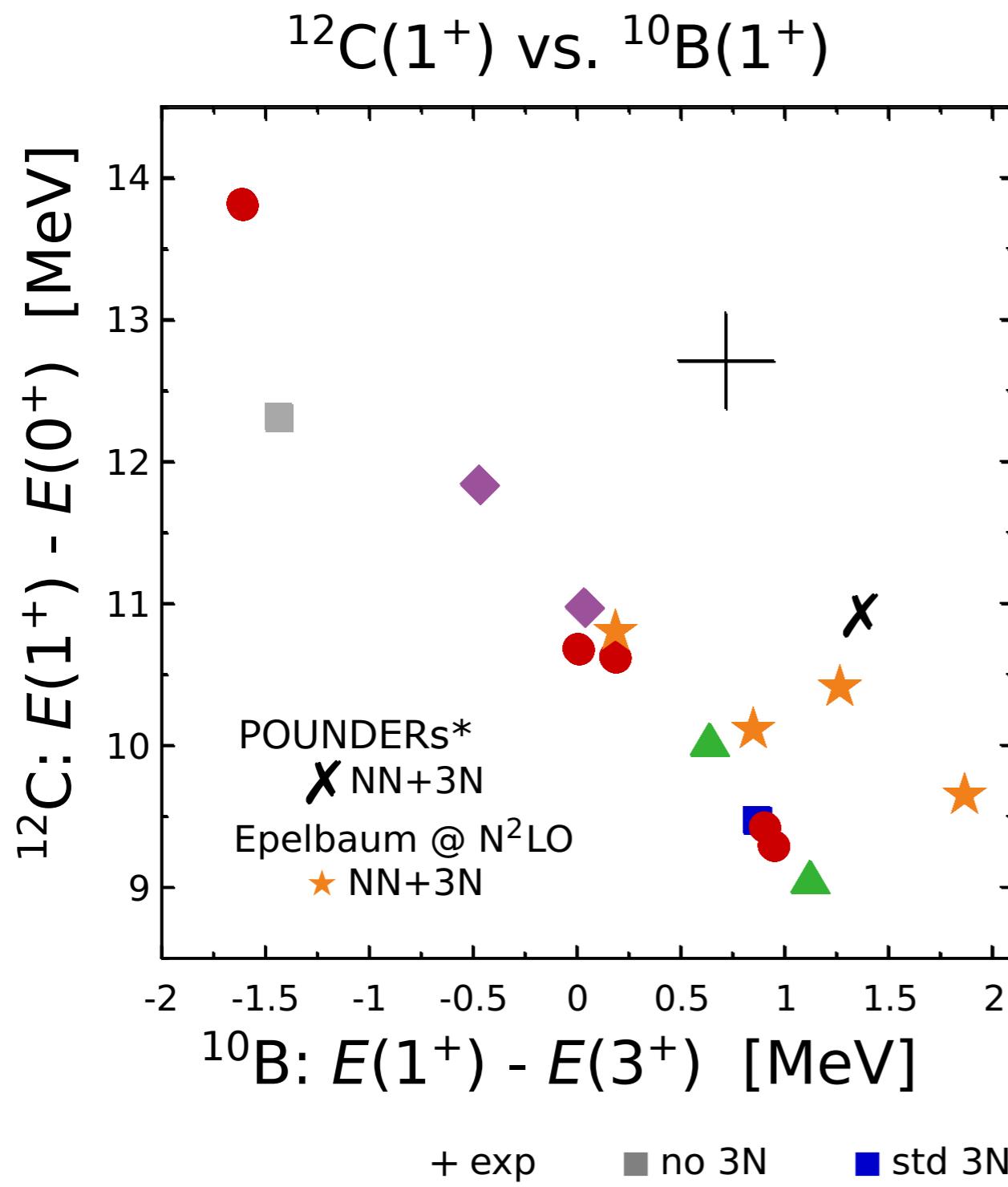
- No 3N interaction compatible with experiment
- One parameter correlation

Epelbaum & POUNDERS  
@ N<sup>2</sup>LO show interesting deviations

+ exp      ■ no 3N      ■ std 3N      ●  $c_i$  var      ▲  $c_D$  var      ◆  $\Lambda$  var

\* Ekström et al., Phys. Rev. Lett. 110, 192502 (2013)  
3N fit of  $c_D$  and  $c_E$  by Navratil & Quaglioni priv. comm.

# Correlation Analysis

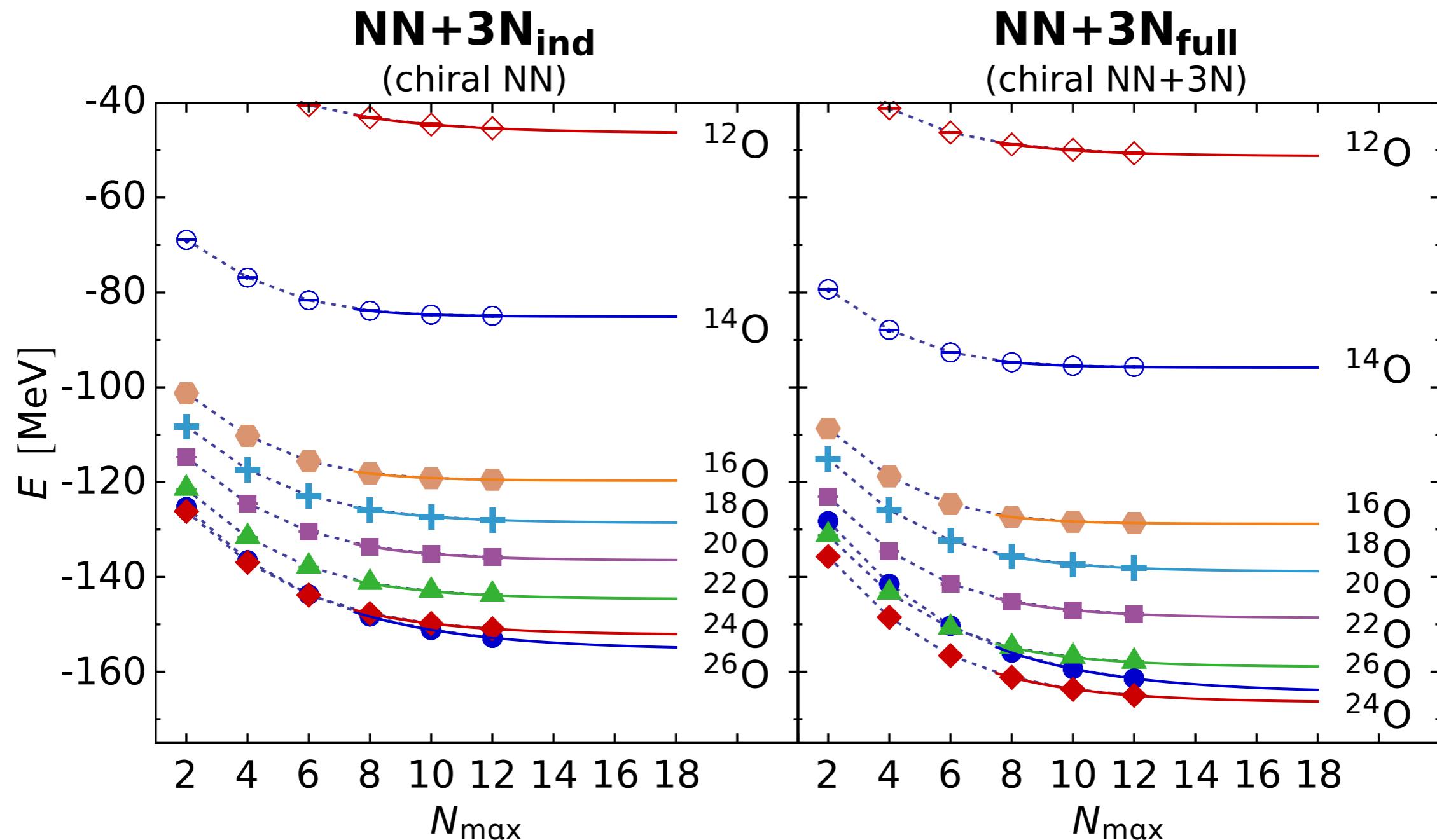


# Ab-Initio Study of the Oxygen Chain

Hergert, Binder, Calci, Langhammer, Roth ----- Phys. Rev. Lett. 110 242501 (2013)

# Ground-States of Even Oxygen Isotopes

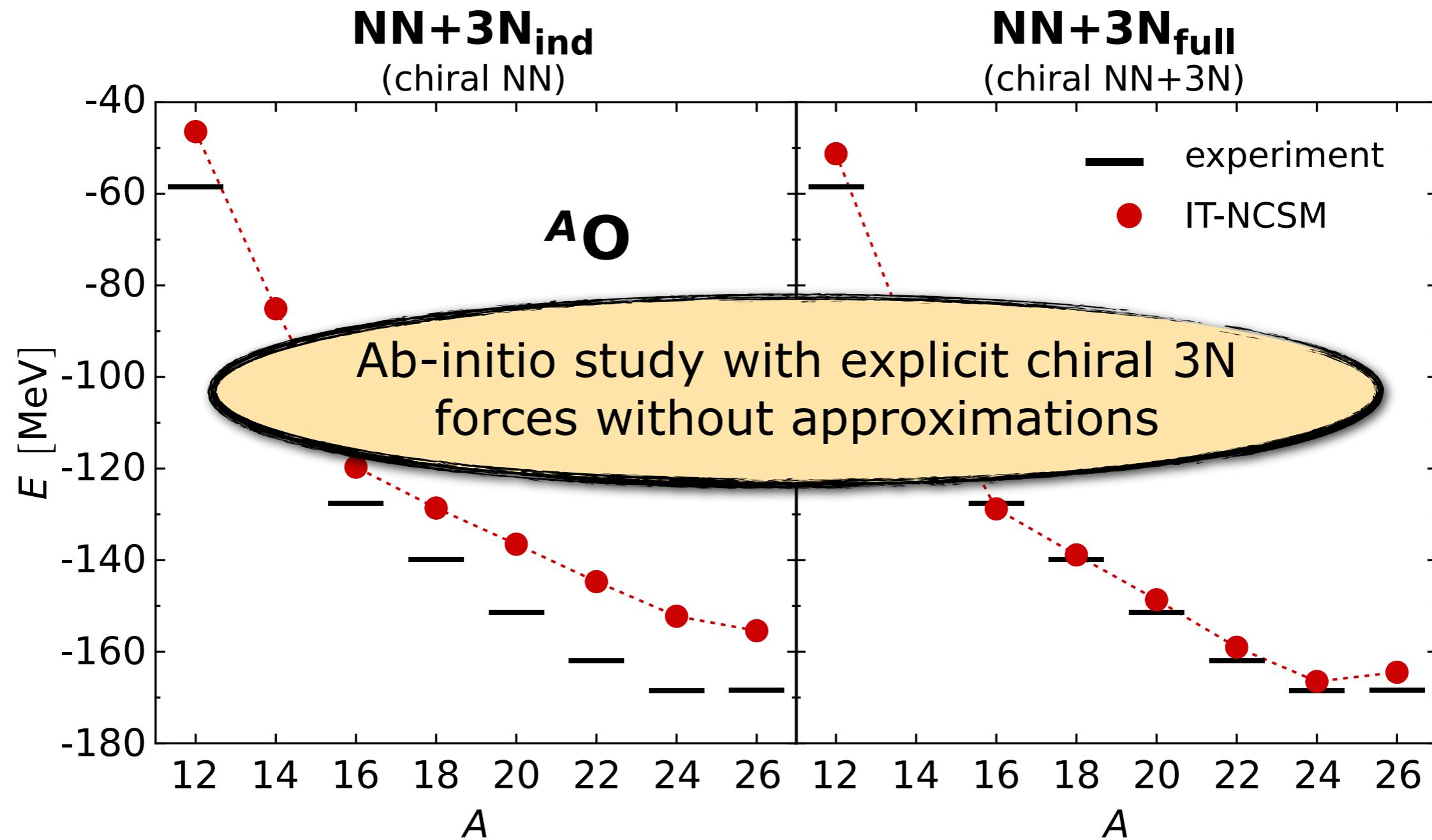
Hergert, Binder, Calci, JL, Roth, PRL **110** 242501 (2013)



$$\Lambda_{3N} = 400 \text{ MeV}, \quad \alpha = 0.08 \text{ fm}^4, \quad E_{3\max} = 14, \quad \text{optimal } \hbar\Omega$$

# Ground-States of Even Oxygen Isotopes

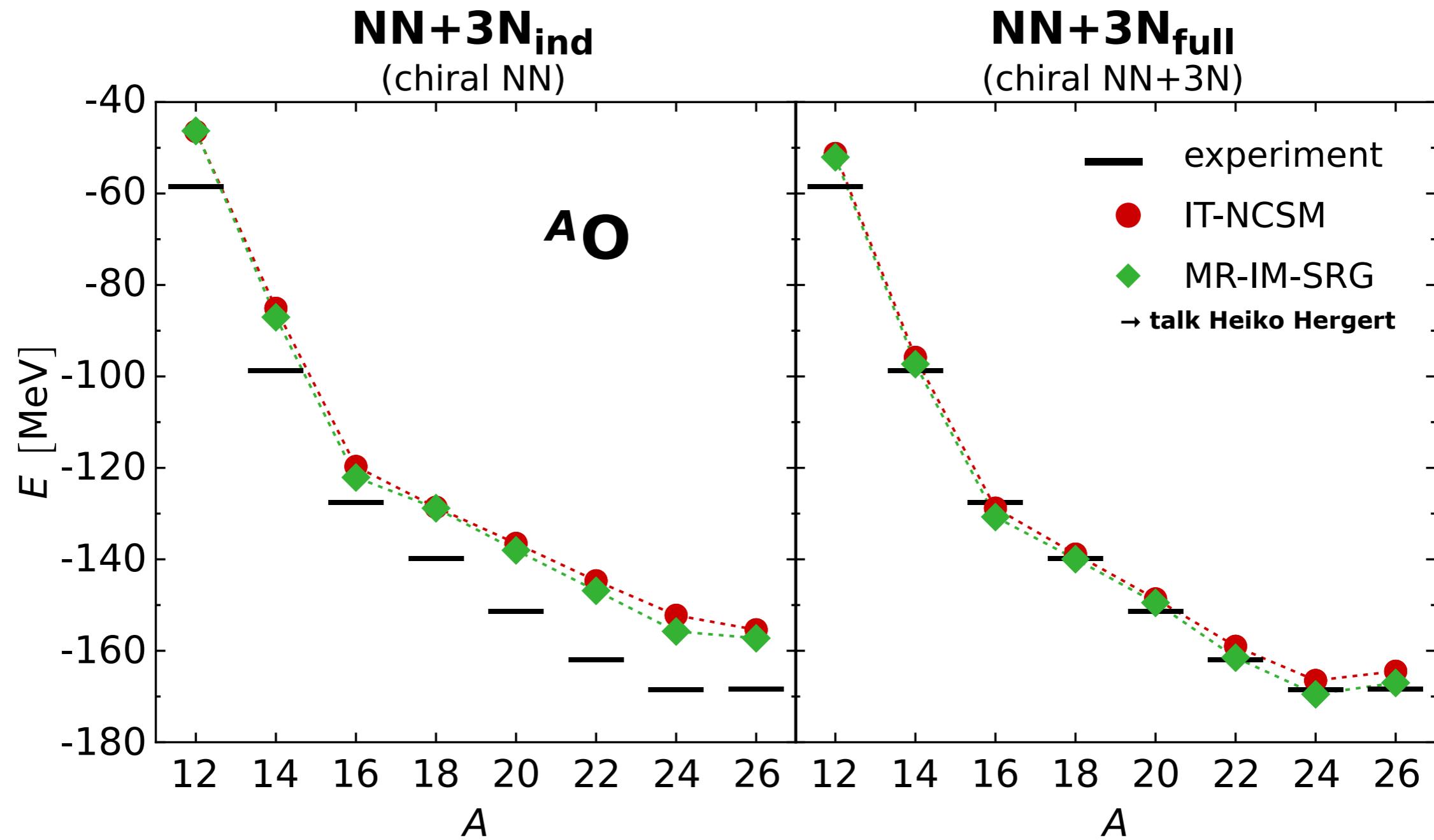
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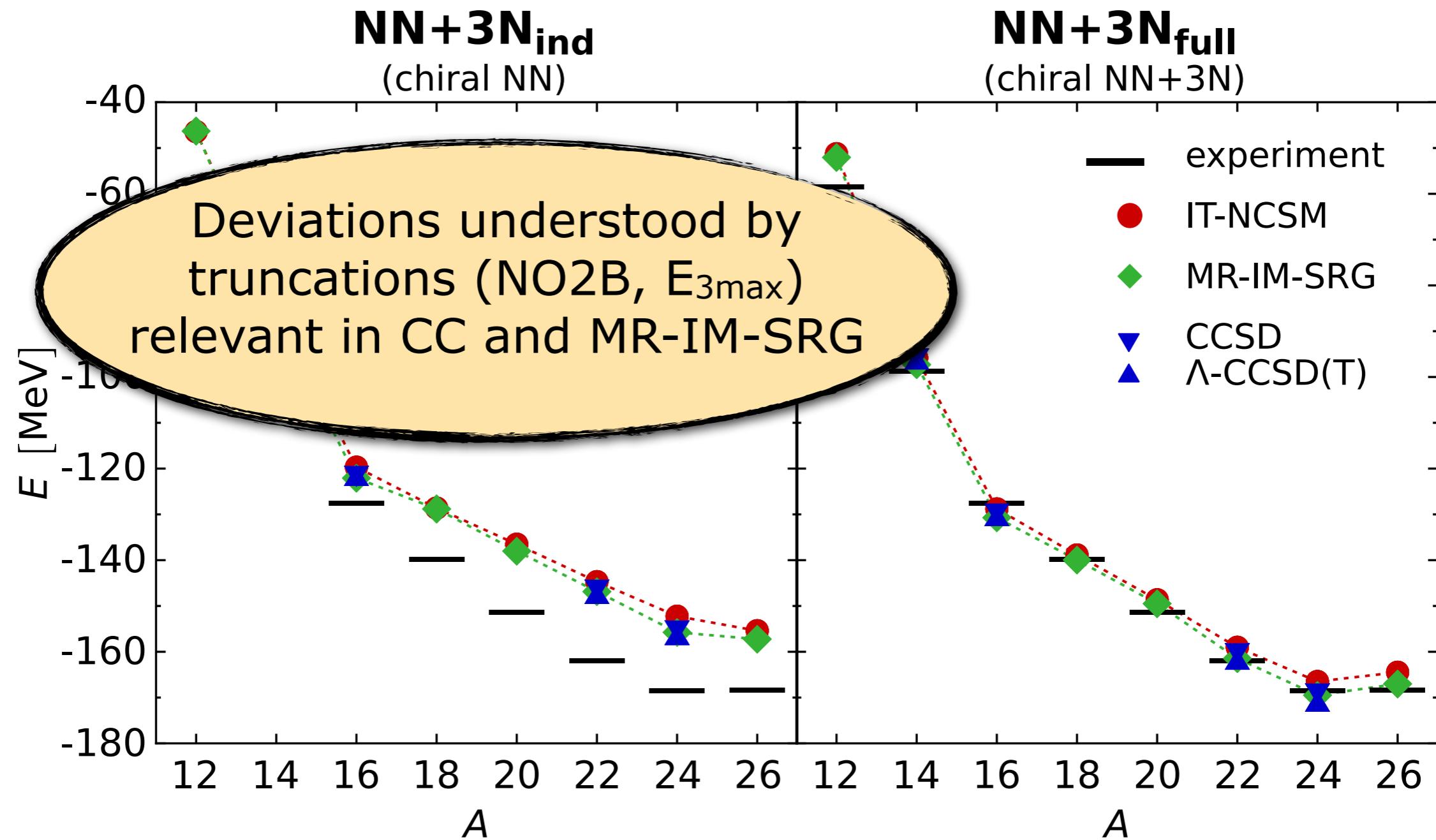
Hergert, Binder, Calci, JL, Roth, PRL **110** 242501 (2013)



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# Ground-States of Even Oxygen Isotopes

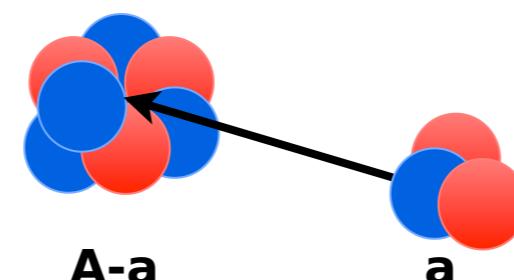
Hergert, Binder, Calci, JL, Roth, PRL **110** 242501 (2013)



# What we are aiming for...

## Realistic ab-initio description of light nuclei

Bound states  
& spectroscopy



Resonances  
& scattering states

**(IT-)NCSM**  
Ab-initio description of  
nuclear clusters

(IT-)NCSM/RGM  
approach

**RGM**  
Describing relative  
motion of clusters

Successfully applied with NN interactions  
Now: Inclusion of 3N Forces

# Nuclear Reactions within the NCSM/RGM approach

In collaboration with  
**G. Hupin, S. Quaglioni, P. Navrátil & R. Roth**

- S. Quaglioni and P. Navrátil ----- Phys. Rev. Lett. 101, 092501 (2008)  
P. Navrátil, R. Roth and S. Quaglioni ----- Phys. Rev. C 82, 034609 (2010)  
S. Quaglioni, P. Navrátil, G. Hupin, J. Langhammer et al. ----- Few-Body Syst. DOI 10.1007/s00601-012-0505-0 (2012)  
S. Quaglioni, P. Navrátil, R. Roth, W. Horiuchi ----- J.Phys.Conf.Ser. 402 (2012)

# General Approach of NCSM/RGM

Wildermuth, Thompson, Tang, ..., Navrátil, Quaglioni, Roth, Hupin, Langhammer,...

- Represent  $H |\psi^{J\pi T}\rangle = E |\psi^{J\pi T}\rangle$  using the **over-complete basis**

$$|\psi^{J\pi T}\rangle = \sum_{\nu} \int dr r^2 \frac{g_{\nu}^{J\pi T}(r)}{r} \mathcal{A}_{\nu} |\phi_{\nu r}^{J\pi T}\rangle \quad g_{\nu}^{J\pi T}(r) \text{ unknown}$$

with the binary-cluster channel states

$$|\phi^{J\pi T}\rangle = \left\{ |\Phi^{(A-a)}\rangle \, |\Phi^{(a)}\rangle \right\}^{J\pi T} \frac{\delta(r-r_{A-a,a})}{rr_{A-a,a}}$$

NCSM delivers  
 $|\Phi^{(A-a)}\rangle$  and  $|\Phi^{(a)}\rangle$

- Solve **generalized eigenvalue** problem

$$\sum_{\nu} \int dr r^2 [\mathcal{H}_{\nu,\nu'}^{J\pi T}(r',r) - E \mathcal{N}_{\nu,\nu'}^{J\pi T}(r,r')] \frac{g_{\nu r}^{J\pi T}}{r} = 0$$

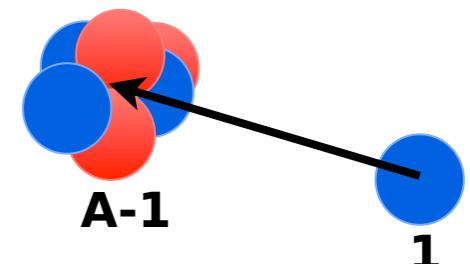
Hamiltonian kernel  
 $\langle \phi_{\nu' r'}^{J\pi T} | \mathcal{A}_{\nu'} H \mathcal{A}_{\nu} | \phi_{\nu r}^{J\pi T} \rangle$

Norm kernel  
 $\langle \phi_{\nu' r'}^{J\pi T} | \mathcal{A}_{\nu'} \mathcal{A}_{\nu} | \phi_{\nu r}^{J\pi T} \rangle$

# Towards Inclusion of Full 3N Forces

- Derive expressions for 3N-interaction kernel

$$\begin{aligned}
 \langle \phi_{\nu' r'}^{J\pi T} | V_{3N} A^2 | \phi_{\nu r}^{J\pi T} \rangle &= \langle \phi_{\nu' r'}^{J\pi T} | V_{3N} \left[ 1 - \sum_{i=1}^{A-1} T_{i,A} \right] | \phi_{\nu r}^{J\pi T} \rangle \\
 &= \frac{(A-1)(A-2)}{2} \langle \phi_{\nu' r'}^{J\pi T} | V_{A-2, A-1, A} | \phi_{\nu r}^{J\pi T} \rangle \\
 &\quad - \frac{(A-1)(A-2)}{2} \langle \phi_{\nu' r'}^{J\pi T} | V_{A-2, A-1, A} T_{A-2, A} | \phi_{\nu r}^{J\pi T} \rangle \\
 &\quad - \frac{(A-1)(A-2)}{2} \langle \phi_{\nu' r'}^{J\pi T} | V_{A-1, A-2, A} T_{A-1, A} | \phi_{\nu r}^{J\pi T} \rangle \\
 &\quad - \frac{(A-1)(A-2)(A-3)}{2} \langle \phi_{\nu' r'}^{J\pi T} | V_{A-3, A-2, A} T_{A-1, A} | \phi_{\nu r}^{J\pi T} \rangle
 \end{aligned}$$



“direct” kernel

Handling  
of 3-body density  
challenging

$$= \underbrace{\dots}_{\alpha \langle \Phi'^{(A-1)} | a^\dagger a^\dagger a a | \Phi^{(A-1)} \rangle} - \underbrace{\dots}_{\alpha \langle \Phi'^{(A-1)} | a^\dagger a^\dagger a^\dagger a a a | \Phi^{(A-1)} \rangle}$$

# Handling of Three-Body Density

Hupin, Quaglioni, Navrátil

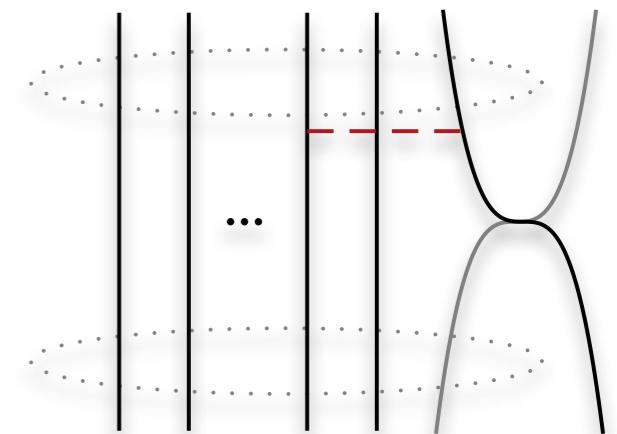
## ① Precomputed coupled densities

$$\sum_{\substack{j_0 j'_0 \\ K J_0}} \sum_{\substack{t_0 t'_0 \\ \tau T_0}} \sum_{\substack{n_a l_a j_a \\ n_b l_b j_b}} \sum_{\substack{n_a' l_a' j_a' \\ n_b' l_b' j_b'}} \frac{1}{12} \hat{\tau} \hat{K} \hat{j}_0 \hat{T}_0 \hat{g}' \hat{t}'_g (-1)^{j'_a + j'_b - j'_0 + j' + K + I_1 + J} (-1)^{3/2 - t'_0 + j' + \tau + T_1 + T}$$

$$\left\{ \begin{matrix} I_1 & K & I'_1 \\ j' & J & j \end{matrix} \right\} \left\{ \begin{matrix} j' & K & j \\ g' & j'_0 & J_0 \end{matrix} \right\} \left\{ \begin{matrix} T_1 & \tau & T'_1 \\ \frac{1}{2} & T & \frac{1}{2} \end{matrix} \right\} \left\{ \begin{matrix} \frac{1}{2} & \tau & \frac{1}{2} \\ t'_g & t'_0 & T_0 \end{matrix} \right\}$$

$$_a \langle ((nlj'_a, nlj'_b)j'_0 t'_0, nlj')J_0 T_0 | V_{3N} | ((nlj_a, nlj_a)j_0 t_0, nlj_b)J_0 T_0 \rangle_a$$

$$\langle \Phi'^{(A-1)} I'_1 T'_1 \left\| \left[ (a_{nlj}^\dagger (a_{nlj'_b}^\dagger a_{nlj'_a}^\dagger)^{j'_0 t'_0})^{g' t'_g} ((\tilde{a}_{nlj_a} \tilde{a}_{nlj_a})^{j_0 t_0} \tilde{a}_{nlj_b})^{J_0 T_0} \right]^{K \tau} \right\| \Phi^{(A-1)} I_1 T_1 \rangle$$



- Make use of  $JT$ -coupled 3N matrix elements
- Three-body density cannot be stored...use a trick

# Handling of Three-Body Density

Hupin, Quaglioni, Navrátil

## ① Precomputed coupled densities

$$\sum_{\substack{j_0 j'_0 \\ J_0 T_0}} \sum_{\substack{t_0 t'_0 \\ n_b l_b j_b}} \sum_{\substack{n_\alpha l_\alpha j_\alpha \\ n'_\alpha l'_\alpha j'_\alpha}} \sum_{\substack{n'_b l'_b j'_b \\ g' t'_g}} \sum_{\Phi'' I_\beta T_\beta}^{\frac{1}{12} \hat{J}_0 \hat{T}_0 \hat{g}' \hat{t}'_g (-1)^{j'_a + j'_b + J_0 + g' + I_\beta - I_1 + j} (-1)^{3/2 + T_0 + t'_g - T_1 + T_\beta}}$$

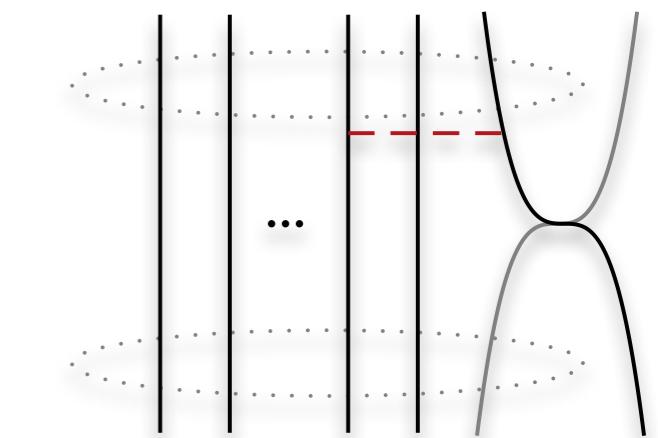
$$\left\{ \begin{matrix} I_\beta & g' & I'_1 \\ J_0 & j'_0 & j' \\ J_1 & j & J \end{matrix} \right\} \left\{ \begin{matrix} T_\beta & t'_g & T'_1 \\ T_0 & t'_0 & \frac{1}{2} \\ T_1 & \frac{1}{2} & T \end{matrix} \right\}$$

$$_a \langle ((nlj'_a, nlj'_b)j'_0 t'_0, nlj')J_0 T_0 | V_{3N} | ((nlj_a, nlj_a)j_0 t_0, nlj_b)J_0 T_0 \rangle_a$$

$$\langle \Phi'^{(A-1)} I'_1 T'_1 \left\| (a_{nlj}^\dagger (a_{nlj'_b}^\dagger a_{nlj'_a}^\dagger)^{j'_0 t'_0})^{g' t'_g} \right\| \Phi''^{(A-4)} I_\beta T_\beta \rangle$$

$$\langle \Phi''^{(A-4)} I_\beta T_\beta \left\| (\tilde{a}_{nlj_\alpha} \tilde{a}_{nlj_\alpha})^{j_0 t_0} \tilde{a}_{nlj_b} \right\| J_0 T_0 \right\| \Phi^{(A-1)} I_1 T_1 \rangle$$

- Make use of  $JT$ -coupled 3N matrix elements
- Three-body density cannot be stored...use a tridiagonal representation
- Use reduced density matrix elements



Applicable to  
 ${}^4\text{He}$  targets

# Handling of Three-Body Density

Langhammer, Roth, Navrátil

## ② Compute uncoupled densities on-the-fly

$$\sum_{jj'} \sum_{M_1 m_j M_{T_1} m_t} \sum_{M'_1 m'_j M'_{T_1} m'_t} \sum_{\frac{1}{12}} (-1)^{I_1 + I'_1 + 2J + j + j'} \begin{Bmatrix} I_1 & \frac{1}{2} & s \\ l & J & j \end{Bmatrix} \begin{Bmatrix} I'_1 & \frac{1}{2} & s' \\ l' & J & j' \end{Bmatrix}$$

$$\left( \begin{array}{cc|c} I_1 & j & J \\ M_1 & m_j & M_J \end{array} \right) \left( \begin{array}{cc|c} T_1 & \frac{1}{2} & T \\ M_{T_1} & m_t & M_T \end{array} \right) \left( \begin{array}{cc|c} I'_1 & j' & J \\ M'_1 & m'_j & M'_J \end{array} \right) \left( \begin{array}{cc|c} T'_1 & \frac{1}{2} & T \\ M'_{T_1} & m' & M' \end{array} \right)$$

$$\sum_{\beta_{A-3} \beta_{A-2} \beta'_{A-3} \beta'_{A-2} \beta'_{A-1}}$$

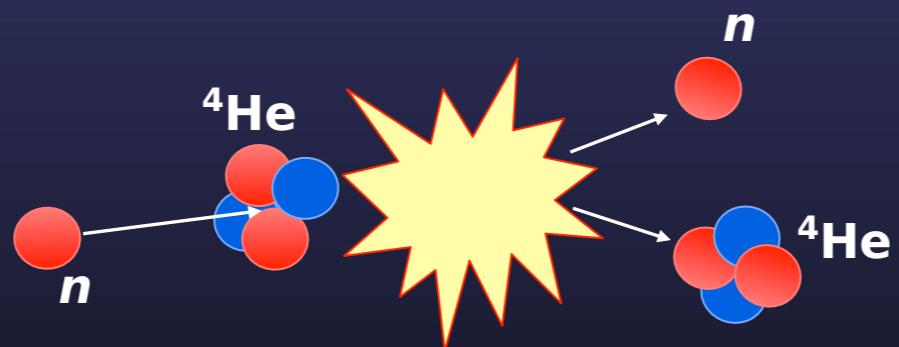
$$a \langle \beta_{A-3} \beta_{A-2} n l j' m'_j \frac{1}{2} m'_t | V_{3N} | \beta'_{A-3} \beta'_{A-2} \beta'_{A-1} \rangle_a$$

$$\langle \Phi'^{(A-1)} I'_1 M'_1 T'_1 M'_{T_1} | a^\dagger_{nljm_j \frac{1}{2} m_t} a^\dagger_{\beta_{A-2}} a^\dagger_{\beta_{A-3}} a_{\beta'_{A-3}} a_{\beta'_{A-2}} a_{\beta'_{A-1}} | \Phi^{(A-1)} I_1 M_1 T_1 M_{T_1} \rangle$$

- Use  $m$ -scheme matrix elements  $\Rightarrow$  efficient
- Perfectly parallel

Treatment of  
 $M_J, M_T$  via  $J_\pm, T_\pm$  ladder  
operators

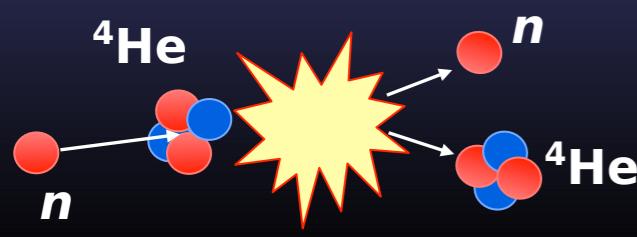
Access to  
heavier targets



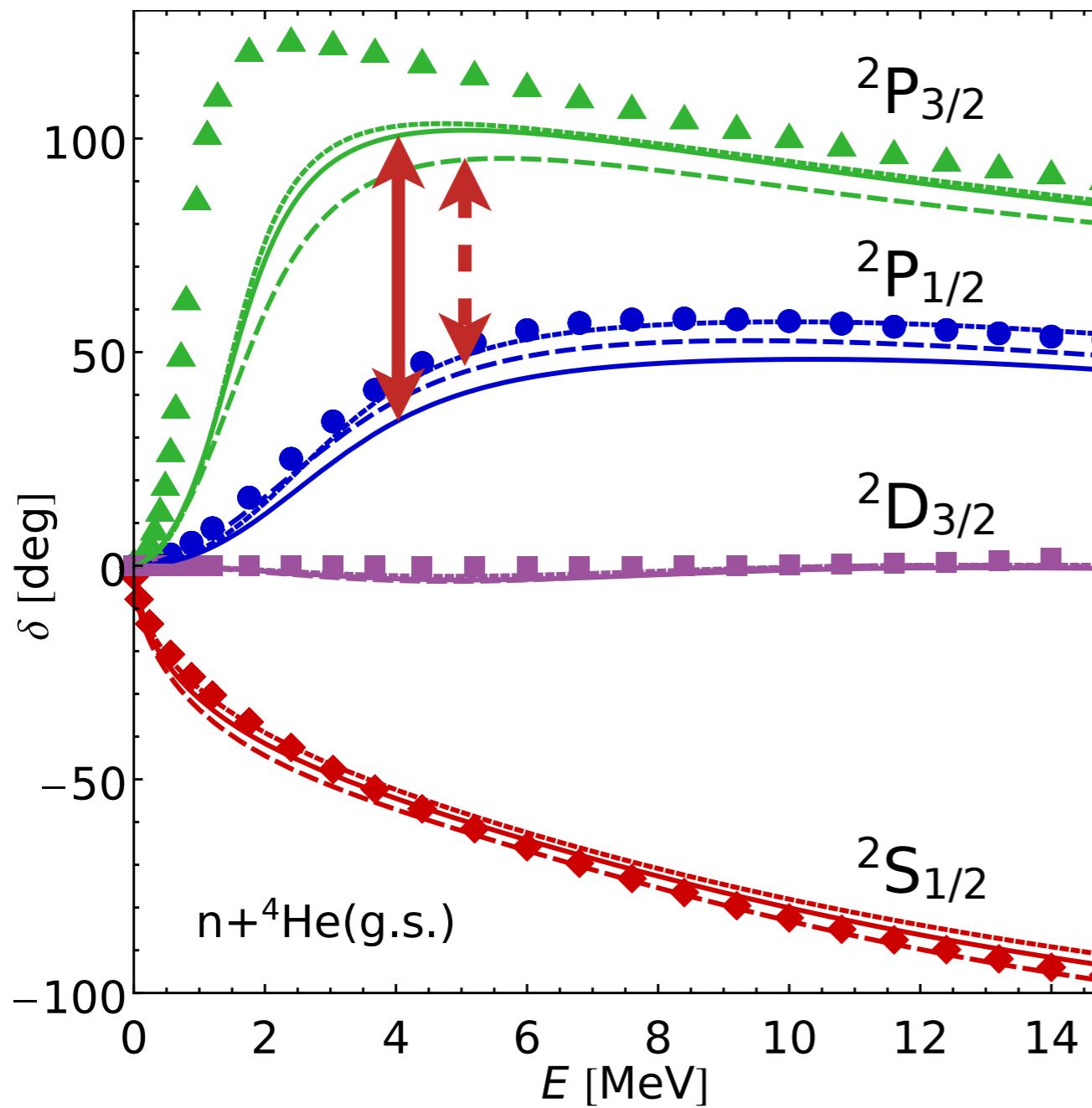
# $n + ^4\text{He}$ Scattering

In collaboration with  
G. Hupin, S. Quaglioni, P. Navrátil & R. Roth

G. Hupin, J. Langhammer et al. ----- in prep.



# 3N Force Effects on Phase Shifts



$\blacklozenge \blacktriangle \blacksquare$  Experiment  
 ..... NN-only  
 - - - 3N-induced  
 — 3N-full

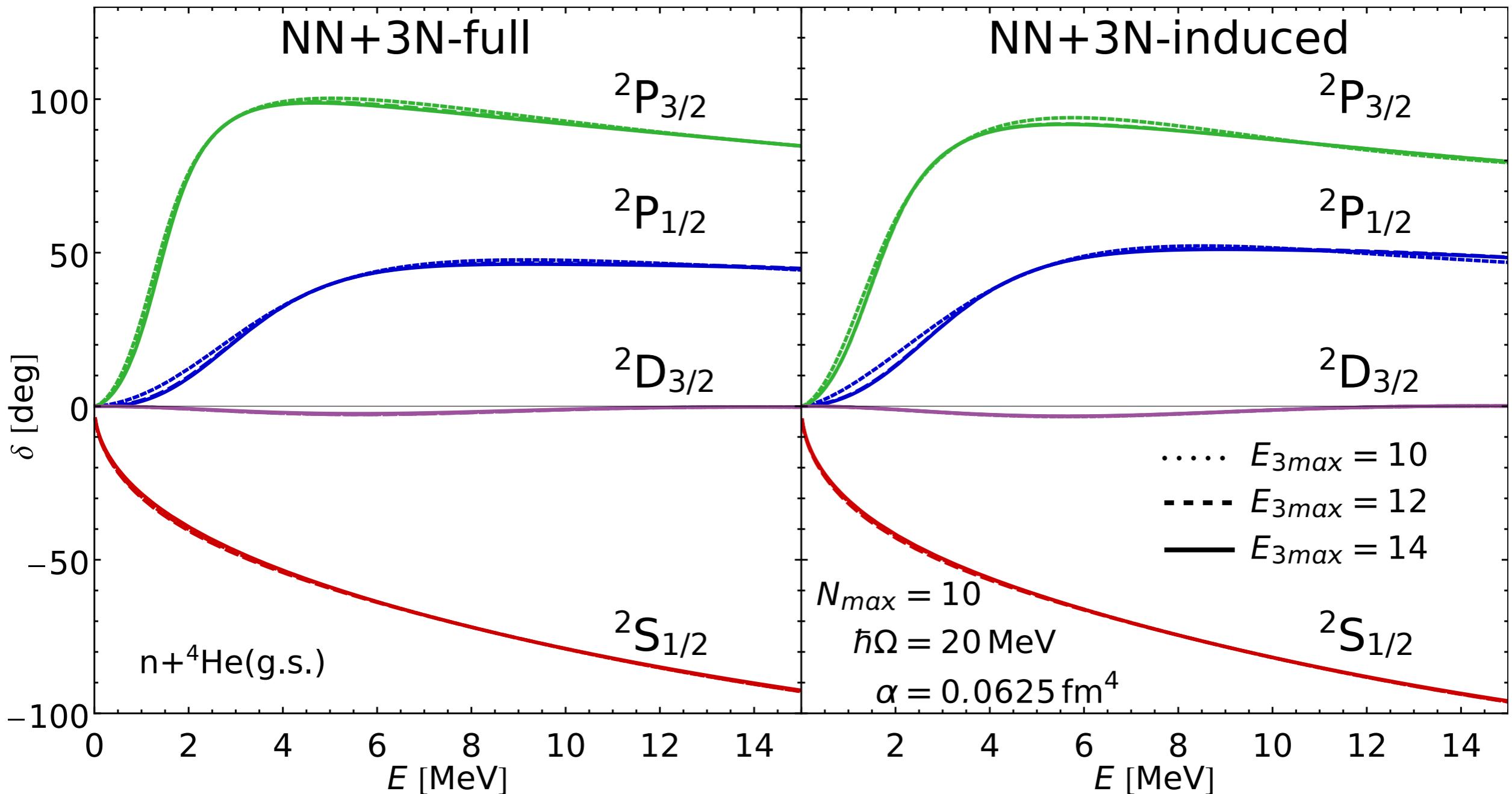
3N increases  
spin-orbit splitting

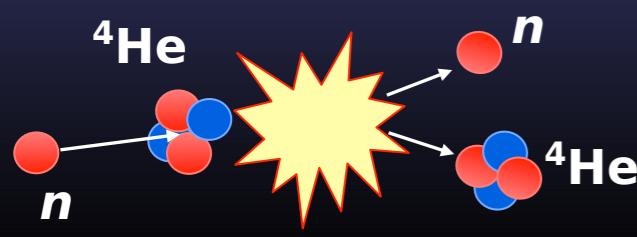
$$\begin{aligned}
 N_{max} &= 12 \\
 E_{3max} &= 14 \\
 \hbar\Omega &= 20 \text{ MeV} \\
 \alpha &= 0.0625 \text{ fm}^4 \\
 \lambda &= 2.0 \text{ fm}^{-1}
 \end{aligned}$$



# 3N Force Effects on Phase Shifts

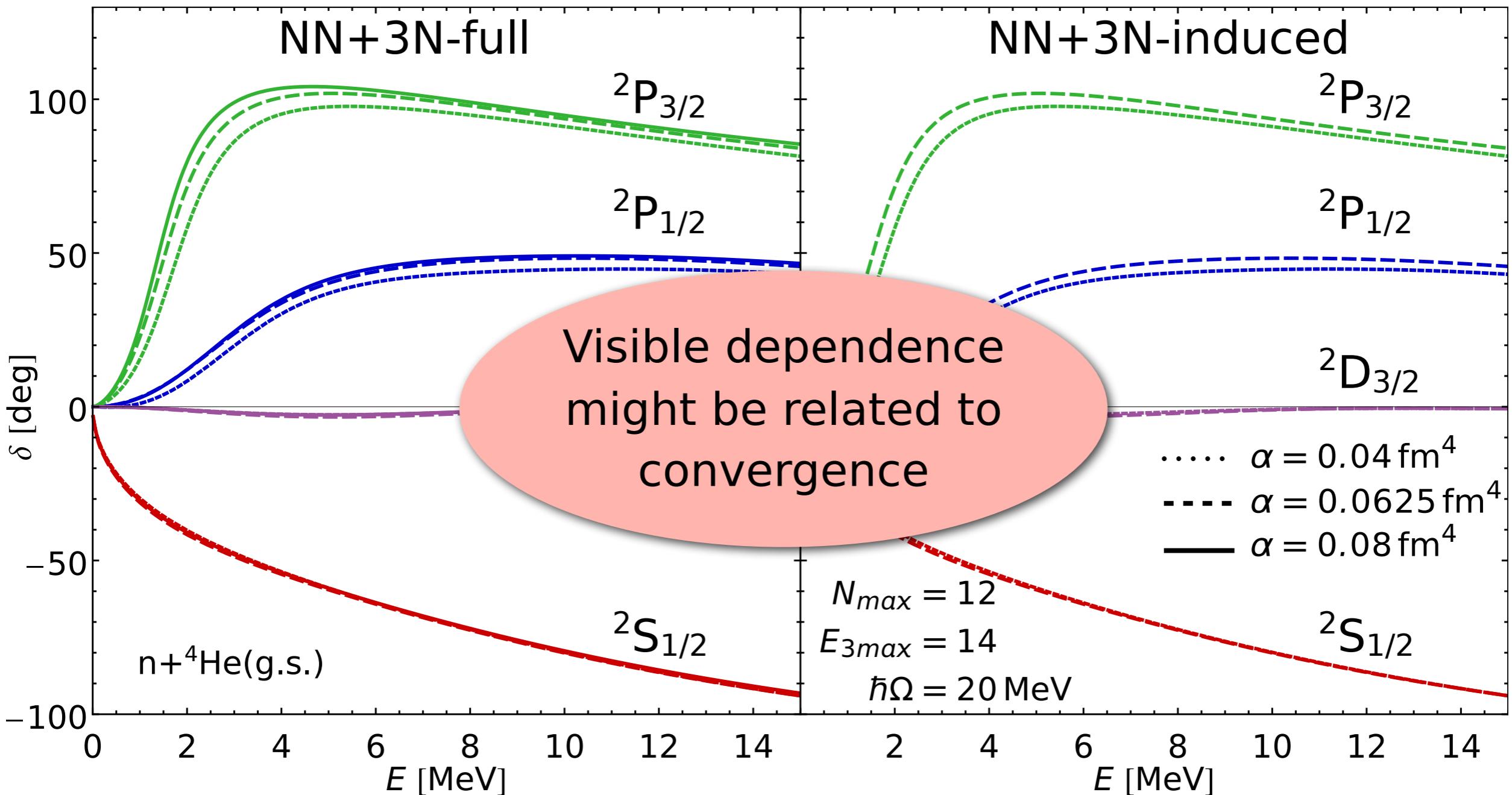
Study dependence  
on  $E_{3max}$  parameter

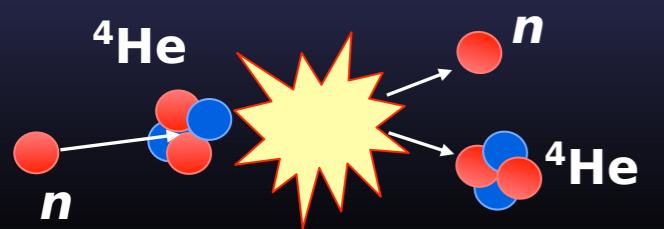




# 3N Force Effects on Phase Shifts

Study dependence  
on SRG parameter

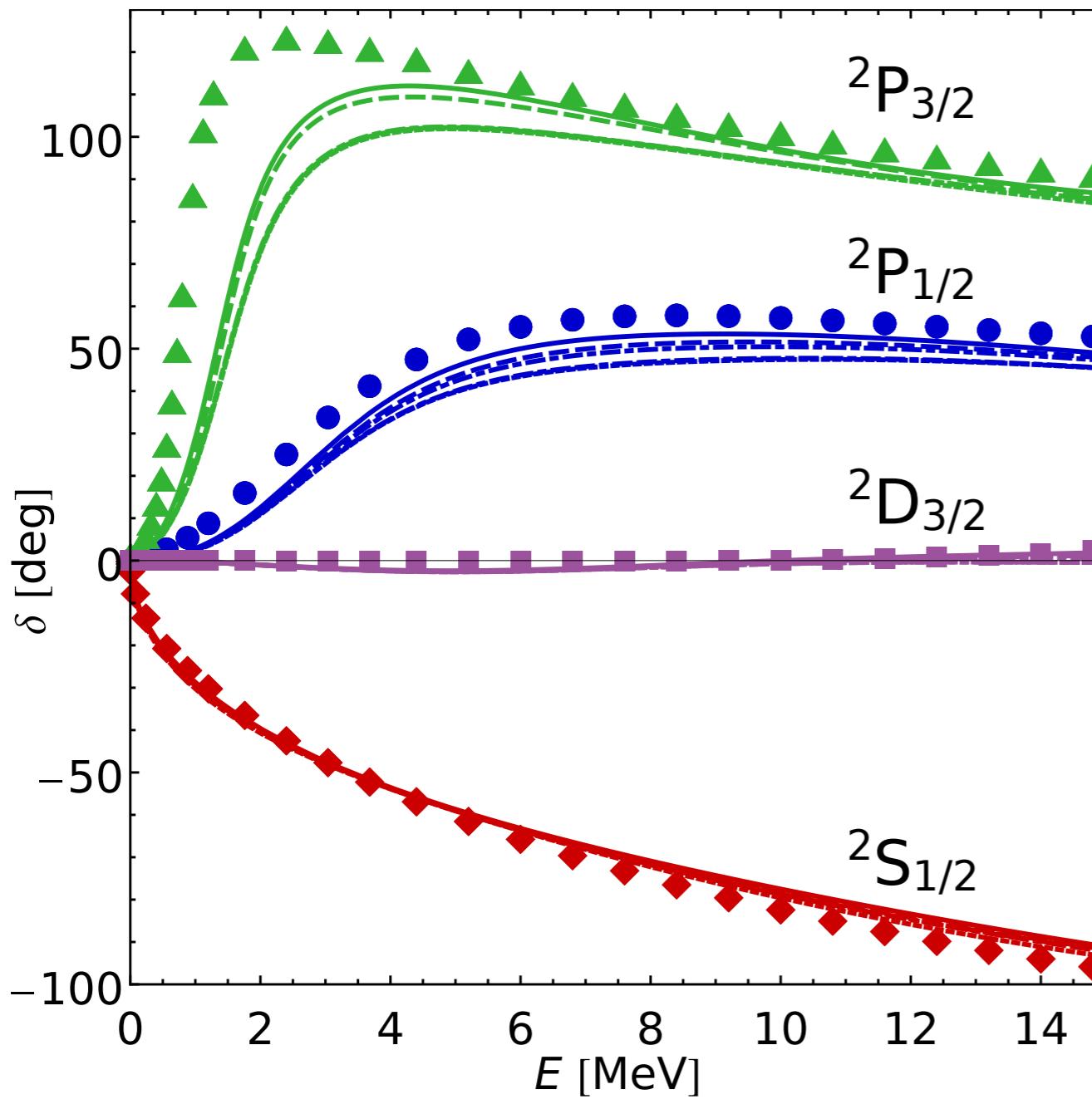




# 3N Force Effects on Phase Shifts

Inclusion of more excited states

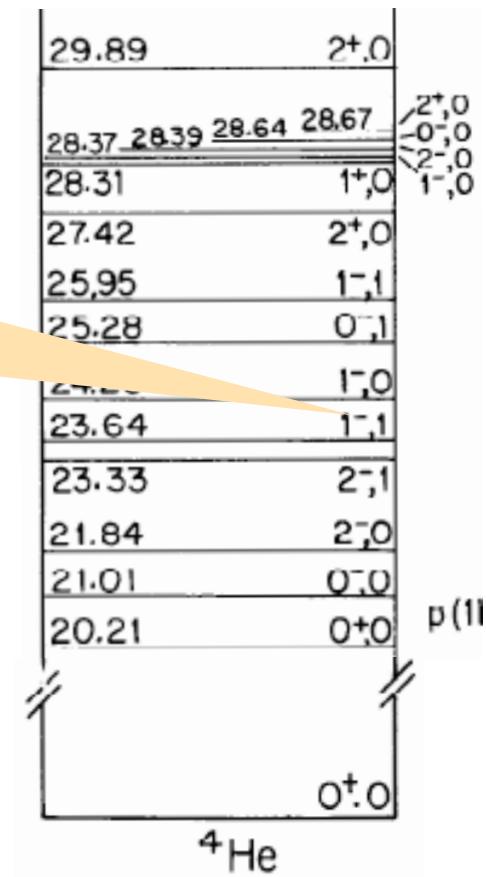
NN+3N-full



- .....  $n + {}^4\text{He(g.s.)}$
- - -  $n + {}^4\text{He(g.s., } 0^+\text{)}$
- - -  $n + {}^4\text{He(g.s., } 0^+, 0^-\text{)}$
- - -  $n + {}^4\text{He(g.s., } 0^+, 0^-, 2^-\text{)}$
- $n + {}^4\text{He(g.s., } 0^+, 0^-, 2^-, 2^-\text{)} T=1$
- ◆▲●■ Experiment

Further inclusion of 1-,1  
& 1-,0 states necessary

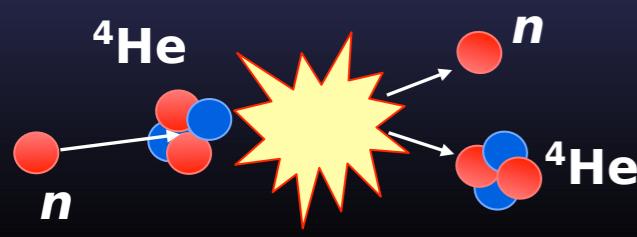
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 \end{aligned}$$



${}^4\text{He}$

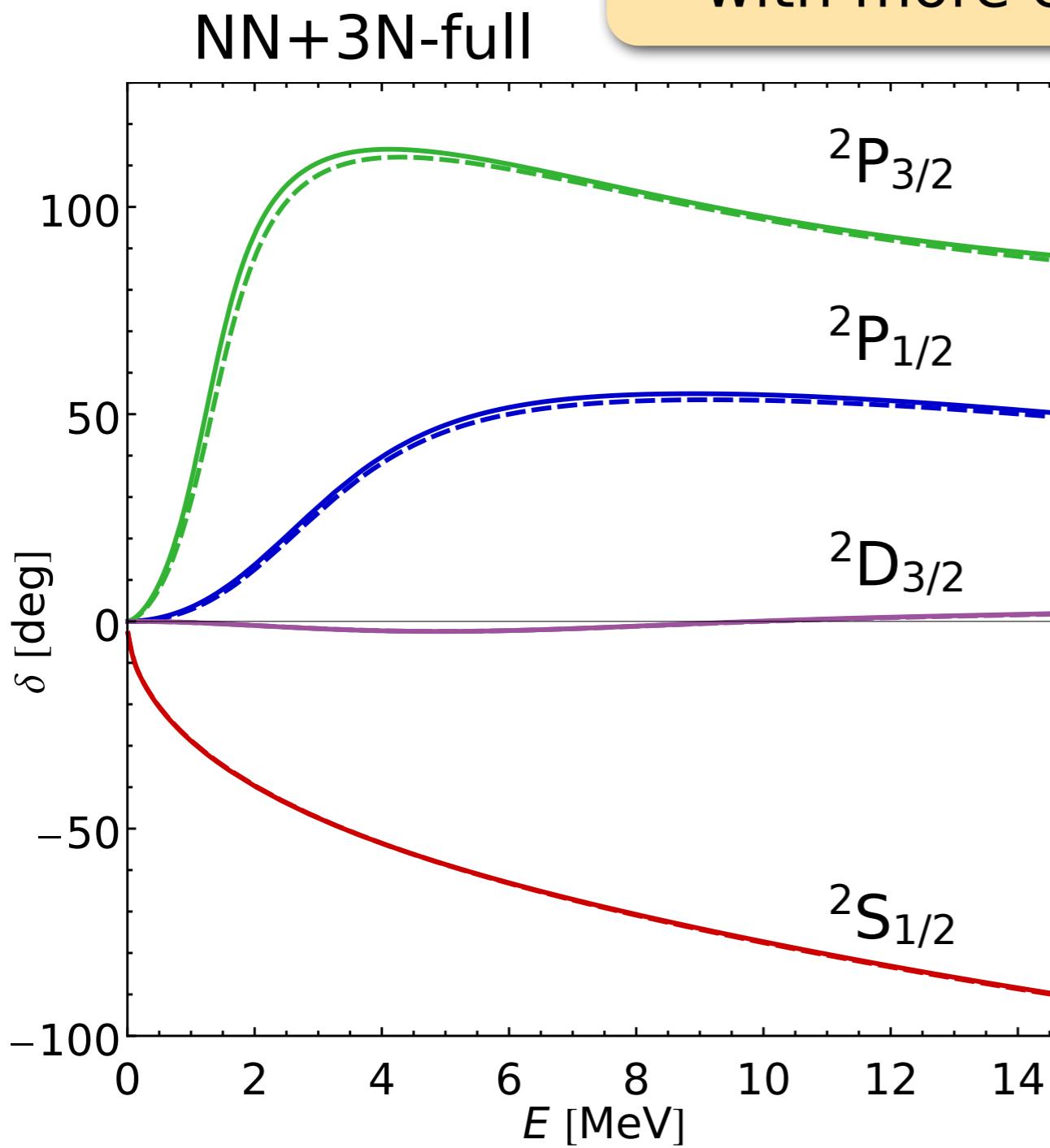
$p(11)$

29.89	$2^+,0$
28.37	$0^-,0$
28.39	$0^-,0$
28.64	$0^-,0$
28.67	$0^-,0$
28.31	$1^+,0$
27.42	$2^+,0$
25.95	$1^-,1$
25.28	$0^-,1$
24.64	$1^-,0$
23.64	$1^-,1$
23.33	$2^-,1$
21.84	$2^-,0$
21.01	$0^-,0$
20.21	$0^+,0$
	$0^-,0$



# 3N Force Effects on Phase Shifts

SRG parameter dependence  
with more excited states



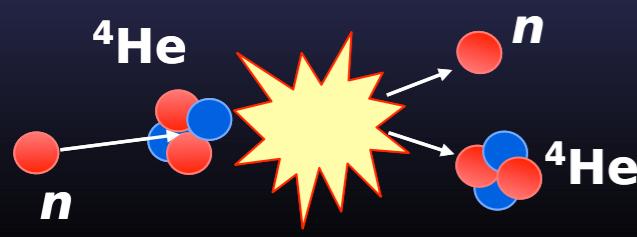
$n + {}^4\text{He}(\text{g.s.}, 0^+, 0^-, 2^-, 2^- \text{ } T=1)$

-----  $\alpha = 0.0625 \text{ fm}^4$

—  $\alpha = 0.08 \text{ fm}^4$

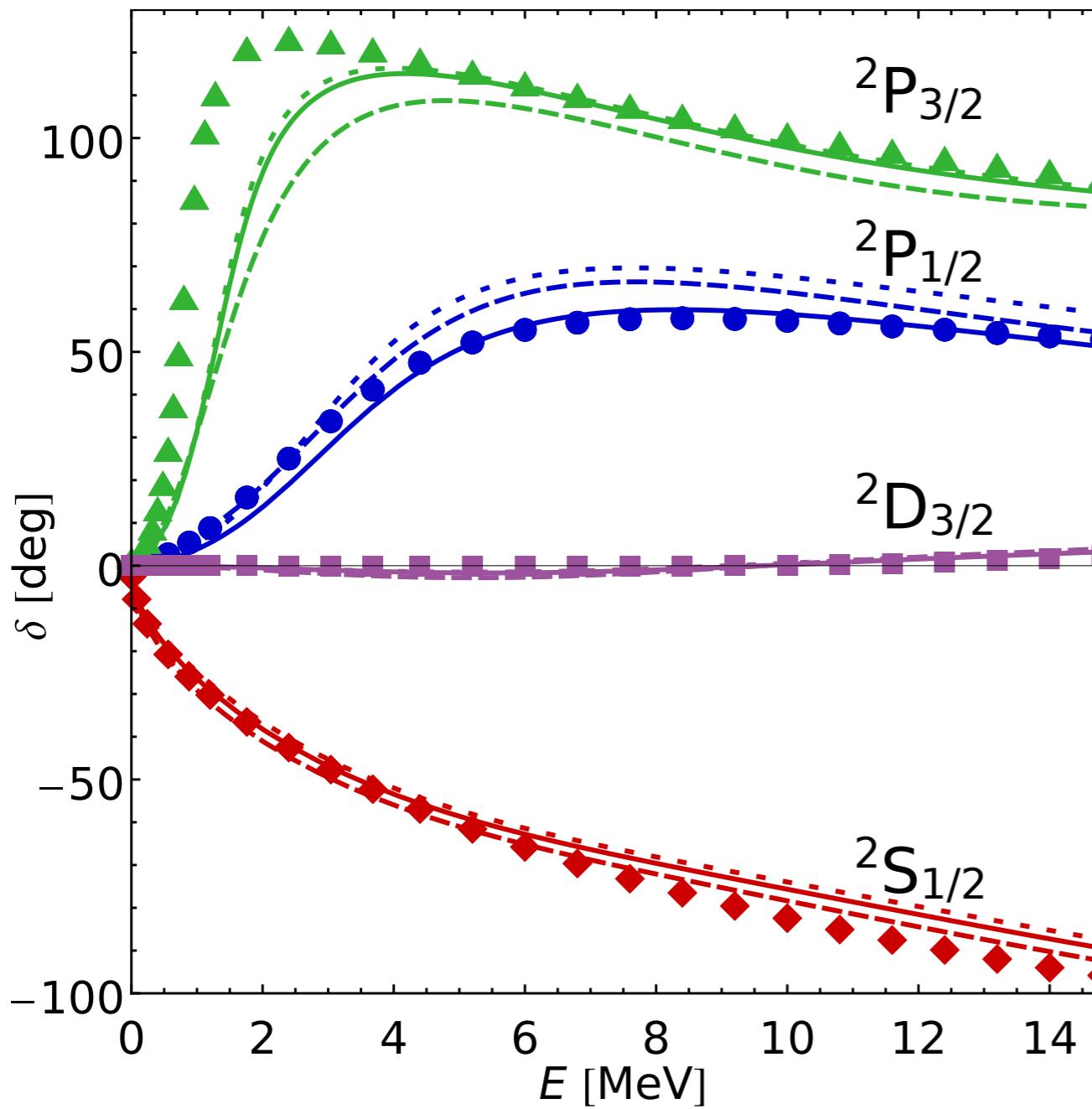
Minimal  
dependence on SRG  
parameter

$N_{max} = 12$   
 $E_{3max} = 14$   
 $\hbar\Omega = 20 \text{ MeV}$   
 $\alpha = 0.0625 \text{ fm}^4$   
 $\lambda = 2.0 \text{ fm}^{-1}$



# 3N Force Effects on Phase Shifts

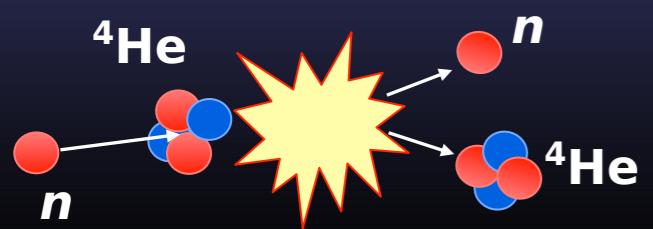
$n + {}^4\text{He}(\text{g.s.}, 0^+, 0^-, 2^-, 2^- \text{T}=1, 1^-, 1^- \text{T}=1)$



- ◆▲●■ Experiment
- ..... NN-only
- 3N-induced
- 3N-full

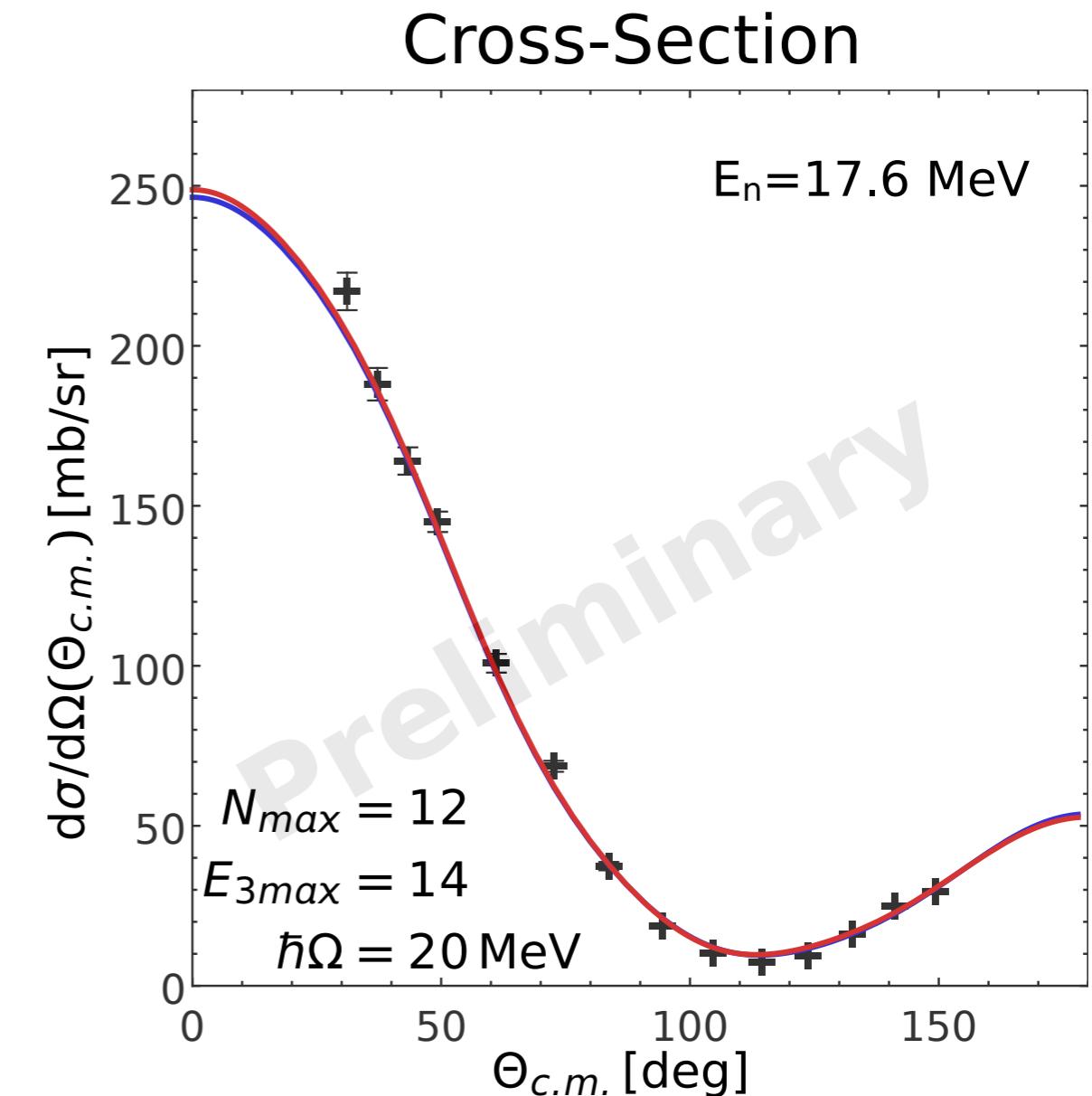
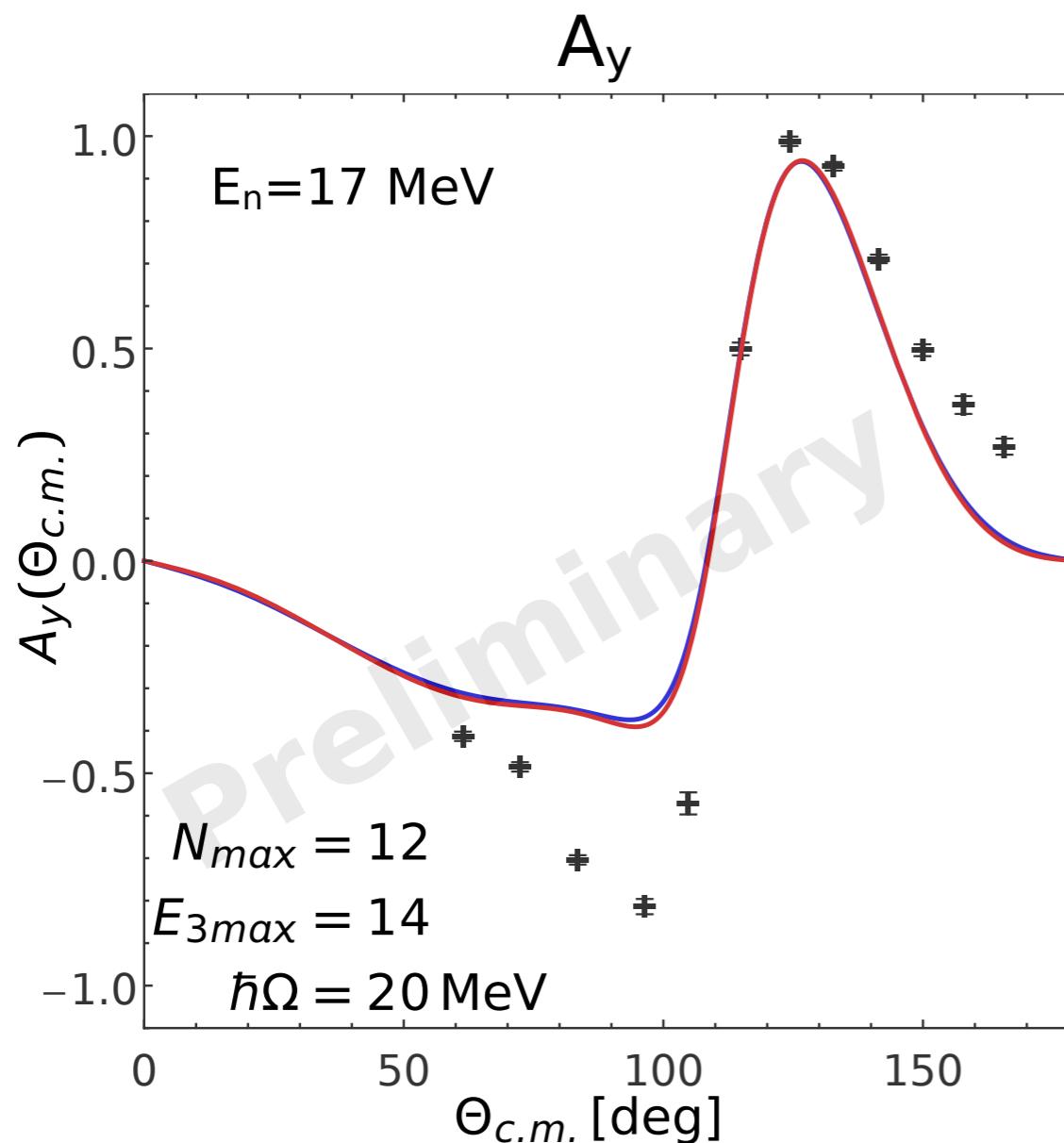
Good agreement  
with data for  ${}^2P_{1/2}$ ,  ${}^2D_{3/2}$   
and  ${}^2S_{1/2}$

$N_{max} = 12$   
 $E_{3max} = 14$   
 $\hbar\Omega = 20 \text{ MeV}$   
 $\alpha = 0.0625 \text{ fm}^4$   
 $\lambda = 2.0 \text{ fm}^{-1}$



# Cross Section & Analyzing Power

$n + {}^4\text{He}(\text{g.s.}, 0^+, 0^-, 2^-, 2^- \text{ T}=1)$



$\alpha = 0.0625 \text{ fm}^4$   
 $\lambda = 2.0 \text{ fm}^{-1}$

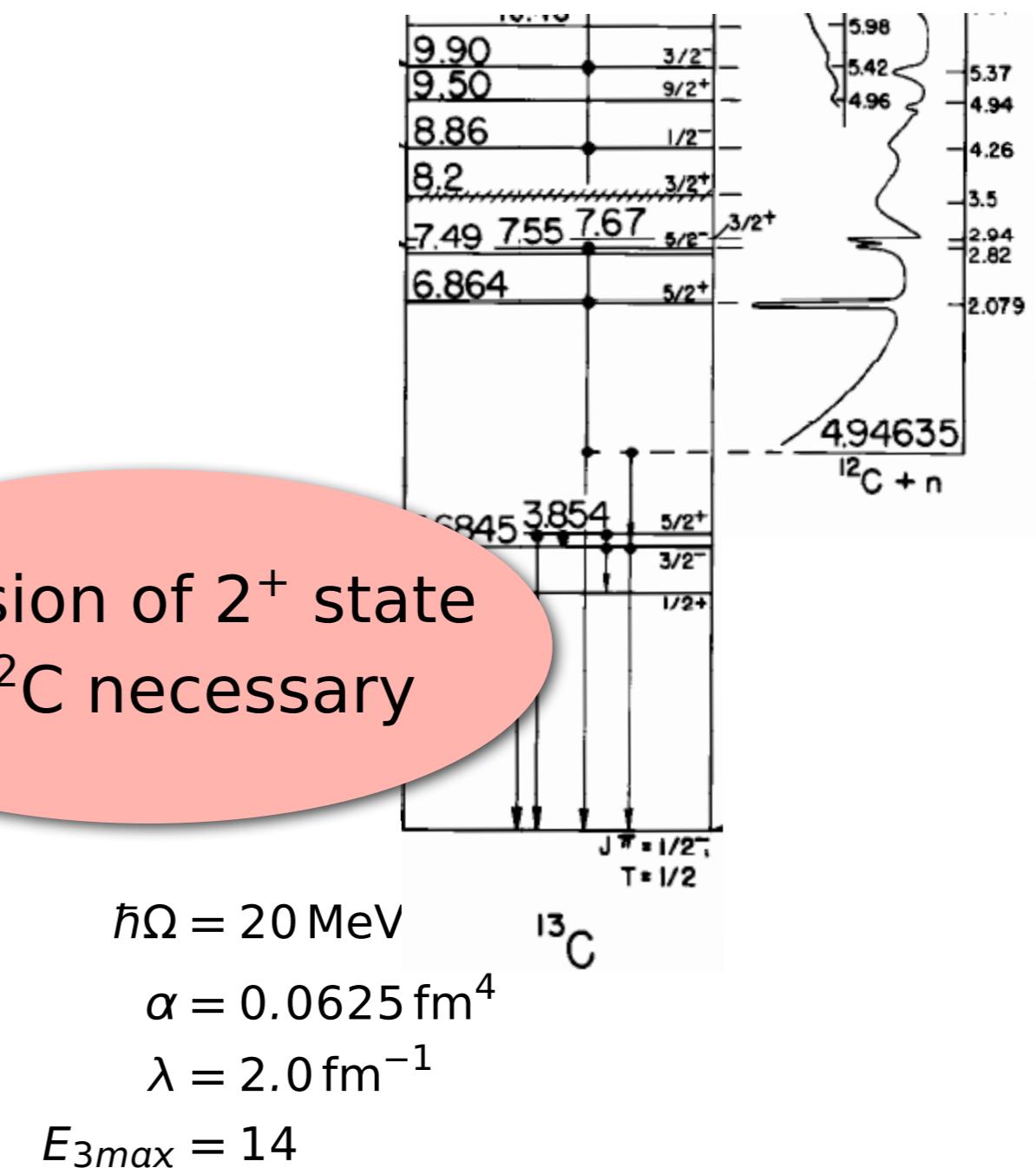
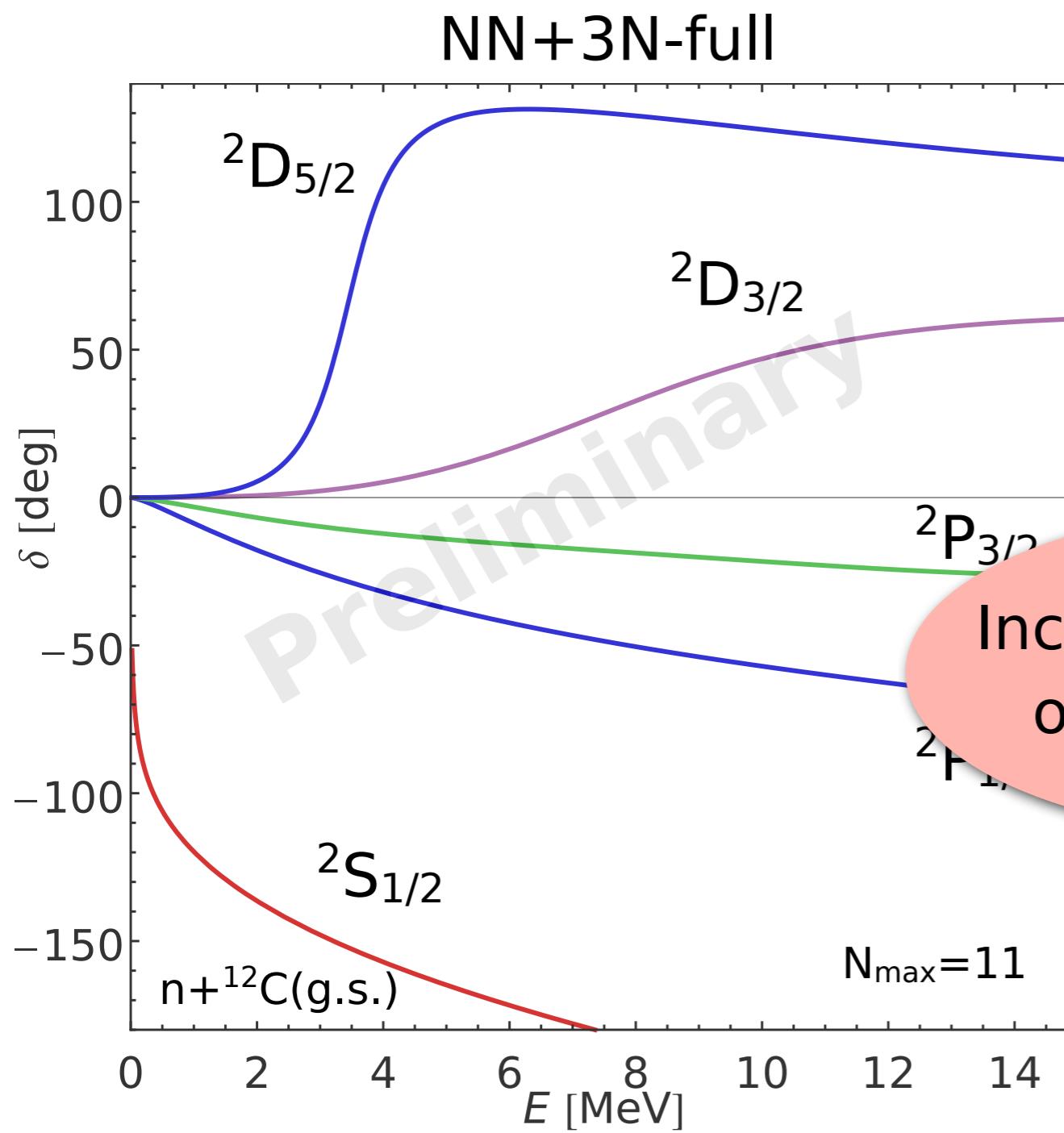
$\alpha = 0.08 \text{ fm}^4$   
 $\lambda = 1.88 \text{ fm}^{-1}$

+ Exp.

[Krupp et al., Phys.Rev.C **30**, 1810]  
[M. Drosg, Los Alamos Scientific Laboratory Report, LA-7269-MS]

# $n + ^{12}\text{C}$ Scattering

- Accessible due to new computational scheme



# Conclusions

# Conclusions

Nuclear spectroscopy with chiral NN+3N forces via IT-NCSM

- ▶ **Propagate uncertainties** to nuclear spectra **and provide feedback**
- ▶ Importance of LEC  $c_3$  for certain states in  $p$ -shell nuclei
- ▶ **Reach** of exact methods **extended, e.g. Oxygen isotopes** in very good agreement with experiments and dripline reproduced

Nuclear reactions with full 3N treatment via NCSM/RGM

- ▶ **Inclusion of 3N forces** challenging but **completed**
- ▶  $n + {}^4He$  scattering phase shifts show **enhanced spin-orbit splitting**
- ▶ New computational scheme  $\implies$  **heavier targets accessible**

Strict tests of chiral interactions in both,  
nuclear structure and reactions possible

# Epilogue

## ■ thanks to my group & collaborators

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■ P. Navrátil

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IDP

Thanks for  
your attention!

Computing Time



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